

Methods

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1 Single contour method

1.1 Standard

The time-domain integral is

$$I(t) = \int d^2x \delta(\phi(\mathbf{x}) - t) \quad (1)$$

where the Fermat potential is

$$\phi(x_1, x_2) = \frac{1}{2}(x_1^2 + x_2^2) + x_1 \cos(\theta) + \frac{1}{2}y^2 - \psi(x_1, x_2) \quad (2)$$

and

$$t = \tau + t_{\min} \quad (3)$$

We can change to polar coordinates

$$x_1 = x_1^0 + R \cos(\theta) \quad (4)$$

$$x_2 = x_2^0 + R \sin(\theta) \quad (5)$$

where (x_1^0, x_2^0) correspond to the position of the minimum

$$\phi(x_1^0, x_2^0) = t_{\min} \quad (6)$$

With this change of coordinates, the integral can be expressed as

$$I(t) = \int R dR d\theta \delta(\phi(R, \theta) - t) \quad (7)$$

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If $\partial_R \phi \neq 0$, we can invert

$$\phi(R, \theta) = t \quad (8)$$

to obtain $R(\theta, t)$. We can then solve the integral over the δ function, plugging in this solution

$$I(t) = \int_0^{2\pi} \frac{R(\theta, t)}{|\partial_R \phi|} d\theta \quad (9)$$

Finally, the system of differential equations that must be solved to find both the curve $R(\theta, t)$ and $I(t)$ is

$$\frac{dI}{d\theta} = \frac{R}{|\partial_R \phi|} \quad (10)$$

$$\frac{dR}{d\theta} = -\frac{\partial_\theta \phi}{\partial_R \phi} \quad (11)$$

Since we are just interested in $I(t)$, we can integrate this system from $\theta = 0$ to 2π , with initial conditions $I(\theta = 0, t) = 0$ and $R(\theta = 0, t)$ such that $\phi(R(0, t), 0) = t$. A couple of useful intermediate results:

$$\partial_\theta f = -\Delta x_2 \partial_{x_1} f + \Delta x_1 \partial_{x_2} f \quad (12)$$

$$\partial_R f = \frac{\Delta x_1}{R} \partial_{x_1} f + \frac{\Delta x_2}{R} \partial_{x_2} f \quad (13)$$

$$(14)$$

where $\Delta x_i = x_i - x_i^0$.

1.2 Robust (parametric)

In general we cannot solve

$$\phi(R, \theta) = t \quad (15)$$

to obtain a relation $R(\theta, t)$, but instead we can find a parametric curve $R(\sigma, t)$ and $\theta(\sigma, t)$. We can think of it as a Hamiltonian system with a Hamiltonian ϕ , and the problem is to find the equations of motion for an energy t . The equations for these curves are

$$\frac{dR}{d\sigma} = -\partial_\theta \phi \quad (16)$$

$$\frac{d\theta}{d\sigma} = \partial_R \phi \quad (17)$$

Now, making the coordinate transformation $(R, \theta) \rightarrow (\sigma, t')$, we end up with

$$I(t) = \int \frac{R(\sigma, t')}{|J|} d\sigma dt' \delta(t - t') \quad (18)$$

where

$$J = \begin{vmatrix} \partial_\sigma R & \partial_t R \\ \partial_\sigma \theta & \partial_t \theta \end{vmatrix} = 1 \quad (19)$$

so we end up with

$$I(t) = \oint R(\sigma, t) d\sigma \quad (20)$$

The final system of ODEs that we must solve is

$$\frac{dI}{d\sigma} = R \quad (21)$$

$$\frac{dR}{d\sigma} = -\partial_\theta \phi \quad (22)$$

$$\frac{d\theta}{d\sigma} = \partial_R \phi \quad (23)$$

This time we must integrate from $\sigma = 0$ until we close the curve¹, i.e. $\theta(\sigma_f, t) = 2\pi$ and $R(\sigma_f, t) = R(0, t)$. The initial conditions are chosen as before.

2 Single integral method

Let us start again with the general expression

$$I(t) = \int d^2x \delta(\phi(\mathbf{x}) - t) \quad (24)$$

This time we will assume that the lens is axisymmetric, so choosing coordinates centered in the lens position

$$x_1 = r \cos(\varphi) \quad (25)$$

$$x_2 = r \sin(\varphi) \quad (26)$$

we have

$$\psi(x_1, x_2) = \psi(r) \quad (27)$$

Using these coordinates, we can solve the angular integral (with $z \equiv \cos(\varphi)$)

$$I(t) = \int_{-1}^1 \frac{dz}{\sqrt{1-z^2}} \int_0^\infty r dr \delta\left(\frac{1}{2}r^2 - yrz + \frac{1}{2}y^2 - \psi(r) - t\right) \quad (28)$$

$$= \int_0^\infty \frac{2dr}{y\sqrt{1-z_*^2}} \Theta(1-z_*^2), \quad z_* \equiv \frac{1}{2yr} (r^2 - 2\psi(r) - t) \quad (29)$$

We can finally express it as (using the notation in the code)

$$\begin{aligned} I(t) &= \int_0^\infty \alpha(r) dr \\ \alpha &= \begin{cases} 2r/\sqrt{\beta}, & \beta > 0 \\ 0, & \beta < 0 \end{cases} \\ \beta &= -\phi_+ \phi_- \\ \phi_\pm &= \frac{1}{2}r^2 + \frac{1}{2}y^2 - \psi(r) - t \mp ry \end{aligned}$$

The problem now reduces to 1) finding the zeros of ϕ_+ and ϕ_- and 2) performing the integral between them (where $\phi_+ \phi_- < 0$).

3 Area/binning method

Finally, another way to solve the integral

$$I(t) = \int d^2x \delta(\phi(\mathbf{x}) - t) \quad (30)$$

is to compute it directly as a surface integral, representing the delta as

$$\delta_n(x) = \begin{cases} 0, & x < -1/2n \\ n, & -1/2n < x < 1/2n \\ 0, & x > 1/2n \end{cases} \quad (31)$$

when $n \rightarrow \infty$. In this way, we obtain a discrete representation

$$I(t) \simeq I_i, \quad \text{for } t \in [t_i, t_i + \Delta t_i] \quad (32)$$

that converges to the real result as we reduce the size of the boxes Δt_i .

¹Notice that since we assume that (x_1^0, x_2^0) is *inside* the contour we know that the final θ that closes the contour is 2π , otherwise it could also be 0.