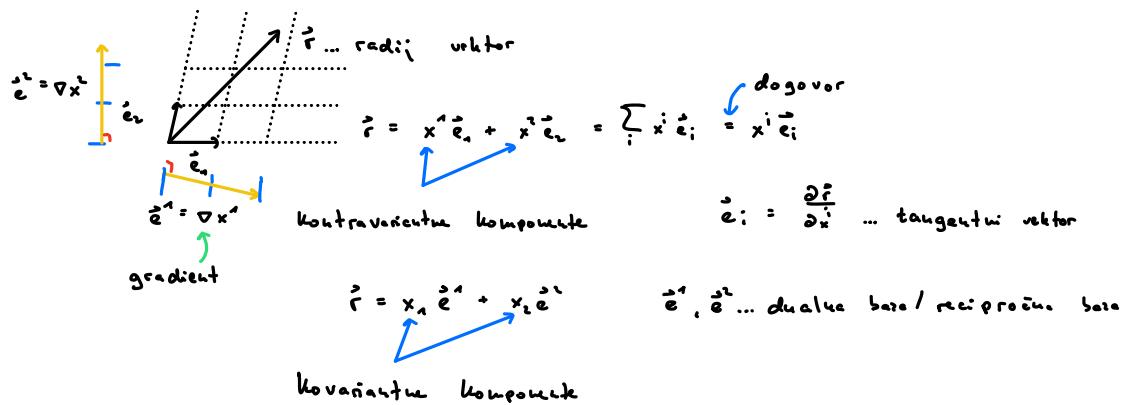


## Kontravariantnost in kovariantnost



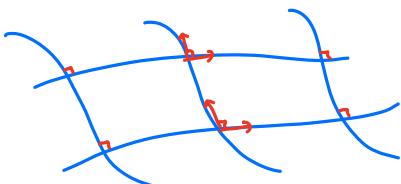
U ortogonalni bazi sva bazi u dualni bazi su isti, zato su tada kontravariantne i kovariantne komponente enake.

## Orthogonalni koordinatni sistemi

- Kartezki koord. sis.

$$\vec{r} = \sum x^i \hat{e}_i = x \hat{e}_x + y \hat{e}_y + z \hat{e}_z = x \hat{i} + y \hat{j} + z \hat{k}$$

- Krivoočrtni koordinatni sistemi



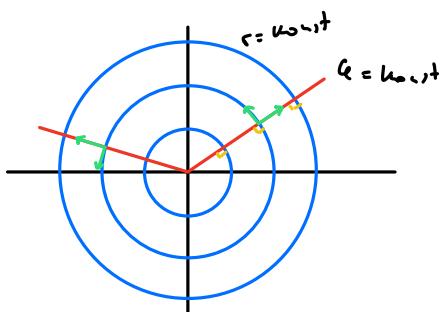
Pričini:

- polarni koord. sis.
- cilindrični koord. sis.
- sferski koord. sis.

} ortogonalni

Baza je lokalna, nije globalna

- Polarni koord. sistem



$$\begin{aligned} \vec{r} &= x \hat{i} + y \hat{j} = r \cos \theta \hat{i} + r \sin \theta \hat{j} \\ x &= r \cos \theta & \hat{e}_r = \frac{\partial \vec{r}}{\partial r} = \cos \theta \hat{i} + \sin \theta \hat{j} \\ y &= r \sin \theta & \hat{e}_\theta = \frac{\partial \vec{r}}{\partial \theta} = -r \sin \theta \hat{i} + r \cos \theta \hat{j} \end{aligned}$$

$$\text{Homotetija } \vec{r} = r \hat{e}_r + \theta \hat{e}_\theta = r \hat{e}_r$$

$$\begin{aligned} |\hat{e}_r|^2 &= \cos^2 \theta |\hat{i}|^2 + \sin^2 \theta |\hat{j}|^2 = 1 \\ |\hat{e}_\theta|^2 &= r^2 \sin^2 \theta |\hat{i}|^2 + r^2 \cos^2 \theta |\hat{j}|^2 = r^2 \end{aligned}$$

1) Newtonov zakon u polarnih koordinatih

$$\vec{F} = m\vec{a} = m\vec{\ddot{v}} = m\vec{\ddot{r}}$$

$$\text{Hitrost } \vec{v} = \frac{d}{dt} \vec{r} = \frac{d\vec{r}}{dt} \frac{dt}{dt} + \frac{d\vec{r}}{dt} \frac{dt}{dt} = \dot{r}\vec{e}_r + \dot{\theta}\vec{e}_\theta$$

$$\text{Pospešek } \vec{a} = \vec{\ddot{v}} = \ddot{r}\vec{e}_r + \dot{r}\vec{e}_r + \dot{\theta}\vec{e}_\theta + \dot{\theta}\vec{e}_\theta$$

$$\begin{aligned}\dot{\vec{e}}_r &= \frac{d}{dt} (\cos \theta \hat{i} + \sin \theta \hat{j}) = \dot{\theta}(-\sin \theta \hat{i} + \cos \theta \hat{j}) = \dot{\theta} \frac{\vec{e}_\theta}{r} \\ \dot{\vec{e}}_\theta &= \frac{d}{dt} (-r \sin \theta \hat{i} + r \cos \theta \hat{j}) = \dot{r}(-\sin \theta \hat{i} + \cos \theta \hat{j}) - r \dot{\theta}(\sin \theta \hat{i} + \cos \theta \hat{j}) = \\ &= \dot{r} \frac{\vec{e}_\theta}{r} - \vec{e}_r \dot{\theta}\end{aligned}$$

$$\vec{a} = \ddot{r}\vec{e}_r (\ddot{r} - \dot{\theta}^2 r) + \dot{r}\vec{e}_\theta (\dot{r} + 2\dot{\theta} \frac{\dot{r}}{r})$$

$$\vec{F} = m (\ddot{r}\vec{e}_r (\ddot{r} - \dot{\theta}^2 r) + \dot{r}\vec{e}_\theta (\dot{r} + 2\dot{\theta} \frac{\dot{r}}{r}))$$

Kinetična energija

$$U_k = T = \frac{1}{2} m v^2 = \frac{1}{2} m (\dot{r} \vec{e}_r + \dot{\theta} \vec{e}_\theta)^2 = \frac{1}{2} m (\dot{r}^2 + r^2 \dot{\theta}^2)$$

2) Sferični koordinatni sistem - izračunaj T

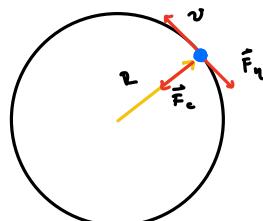
DN

$$x = r \cos \theta \sin \phi$$

$$y = r \sin \theta \sin \phi$$

$$z = r \cos \theta$$

3) Kretanje po obroču z viskoznošću



$$\vec{r}(t) = ?$$

$$m \ddot{\vec{r}} = \vec{F}_c + \vec{F}_d \quad \vec{F}_d = -\eta \vec{v} \quad r = R \quad \dot{r} = \ddot{r} = 0$$

$$-mR\dot{\theta}^2 \vec{e}_r + m\ddot{\theta} \vec{e}_\theta = -F_c \vec{e}_r - \eta (\dot{r} \vec{e}_r + \dot{\theta} \vec{e}_\theta)$$

$$\vec{e}_r: -mR\dot{\theta}^2 = -F_c \Rightarrow F_c = mR\dot{\theta}^2$$

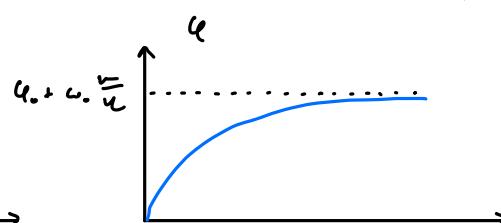
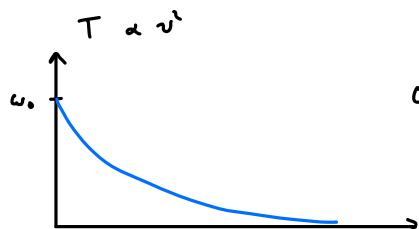
$$\vec{e}_\theta: m\ddot{\theta} = -\eta \dot{\theta}$$

$$x = \dot{r} \quad m\dot{x} = -\eta x \quad \frac{dx}{dt} = -\frac{\eta}{m} dt \quad \ln x = -\frac{\eta}{m} t$$

$$x = \dot{r} = \omega_0 e^{-\frac{\eta}{m} t}$$

$$\int d\theta(t) = -\frac{m\omega_0}{\eta} e^{-\frac{\eta}{m} t} \Big|_0^t$$

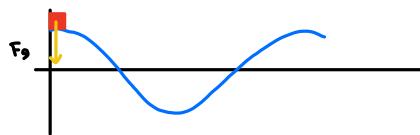
$$\theta(t) = \omega_0 \frac{m}{\eta} \left(1 - e^{-\frac{\eta}{m} t}\right) + \theta_0$$



④ Enak na loga let 3 ke da je zraven še gravitacija.

DP

⑤ Klada na hossiuhi podlagi - ali se klada kdaj odlepi:



$$y = y_0 \cos\left(\frac{x}{x_0}\right)$$

$$F_{xc} = 0$$



$$\tan q = \frac{dy}{dx} = y'$$

$$\hat{e}_x = \cos q \hat{i} + \sin q \hat{j} \dots \text{tangential vector}$$

$$\hat{e}_z = -\sin q \hat{i} + \cos q \hat{j} \dots \text{normalni vektor}$$

$$\begin{aligned} \vec{v} &= v \hat{e}_x \\ \vec{a} &= \dot{\vec{v}} = \dot{v} \hat{e}_x + v \dot{\hat{e}}_x \\ &= \dot{v} \hat{e}_x + v q \hat{e}_z \end{aligned}$$

$$\dot{\hat{e}}_x = \frac{d}{dt} (\cos q \hat{i} + \sin q \hat{j}) = q \hat{e}_z$$

$$\text{Odlepi: se ko } F_t = 0 \Rightarrow m \ddot{a} = \vec{F}_g$$

$$m(\dot{v} \hat{e}_x + v q \dot{\hat{e}}_z) = -mg \hat{j} = -mg (\sin q \hat{e}_x + \cos q \hat{e}_z)$$

$$\begin{aligned} \hat{e}_x &\quad m v q = -m g \cos q \\ \tan q &= y' \end{aligned}$$

$$\frac{d}{dt} \tan q = \frac{1}{\cos^2 q} q'' = \frac{d}{dt} y' = \frac{dy'}{dx} \frac{dx}{dt} = \dot{x} y''$$

$$\hat{j} = \sin q \hat{e}_x + \cos q \hat{e}_z$$

$$\dot{x}^2 + \dot{y}^2 = v^2 = \left(\frac{dx}{dt}\right)^2 + \left(\frac{dy}{dt}\right)^2 = \left(\frac{dx}{dt}\right)^2 \left(1 + \left(\frac{dy}{dx}\right)^2\right) = \dot{x}^2 (1 + y'^2)$$

$$\Rightarrow \dot{x} = \frac{v}{\sqrt{1+y'^2}}$$

$$\Rightarrow \frac{1}{\cos^2 q} q'' = \dot{x} y'' = \frac{v y''}{1+y'^2}$$

$$\cos^2 q = \frac{1}{1+y'^2} = \frac{1}{1+y''^2}$$

$$q'' = \frac{v^2 y''}{(1+y'^2)^2}$$

$$\frac{v^2 y''}{(1+y'^2)^2} = -\frac{g}{\sqrt{1+y''^2}}$$

$$\underline{v^2 y'' + g(1+y'^2) = 0} \quad \text{Kdaj se telo odlepi od podlage}$$

$$\begin{aligned} q &= q_0 \cos \frac{x}{x_0} \\ q' &= -\frac{q_0}{x_0} \sin \frac{x}{x_0} = -\frac{q_0}{x_0} \sqrt{1 - \cos^2 \frac{x}{x_0}} = -\frac{1}{x_0} \sqrt{q_0^2 - y^2} \\ q'' &= -\frac{q_0}{x_0^2} \cos \frac{x}{x_0} = -\frac{q_0}{x_0^2} \end{aligned}$$

$$v^2: \frac{1}{x_0^2} v^2 = mg \Delta h = mg (q_0 - q) \quad v^2 = 2g (q_0 - q)$$

$$-2g(q_0 - q) \frac{q_0}{x_0^2} + g \left(1 + \frac{1}{x_0^2} (q_0^2 - q^2)\right) = 0$$

$$2q(q_0 - q) + x_0^2 + q_0^2 - q^2 = 0$$

$$(q - q_0)^2 + x_0^2 = 0$$

Eenaka nima razide v x\_0 > 0

zato se telo se odlepi od podlage

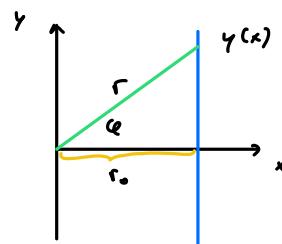
## ⑥ Prost. dreh. u. polaruhr Koordinatent

$$F = ma$$

$$m\ddot{x} = 0 \quad v_x = \text{konst.} \quad x = x_0 + v_x t$$

$$m\ddot{y} = 0 \quad v_y = \text{konst.} \quad y = y_0 + v_y t$$

$$\gamma(x) = y_0 + \frac{v_y}{v_x} (x - x_0)$$



$$r(\theta) = ?$$

$$r_0 = r \cos \varphi$$

$$r = r_0 / \cos \varphi$$

Newtonov zchen u. pol. Koord.

$$\ddot{r} = \dot{\theta} \dot{r} - r \dot{\theta}^2 + \ddot{\theta} r + 2 \dot{r} \dot{\theta}$$

$$\dot{\theta} = r \dot{\theta}_0$$

$$e_1: \quad \ddot{r} - r \dot{\theta}^2 = 0$$

$$e_2: \quad r \ddot{\theta} + 2 \dot{r} \dot{\theta} = 0 \quad | \cdot \downarrow$$

$$\textcircled{i} \quad r \dot{\theta} + 2 \dot{r} \theta = 0$$

$$\frac{\dot{\theta}}{\theta} = -2 \frac{\dot{r}}{r}$$

$$\frac{dt}{dt} \ln \theta = \frac{dt}{dt} \ln r^2$$

$$\frac{\theta}{\theta_0} = \frac{r^2}{r_0^2}$$

$$r^2 \theta = \text{konst.}$$

$$r^2 \dot{\theta} = \text{konst.}$$

$$| \cdot r$$

$$\textcircled{ii} \quad r^2 \dot{\theta} + 2 \dot{r} r \dot{\theta} = 0$$

$$\frac{d}{dt} (r^2 \dot{\theta}) = 0$$

$$r^2 \dot{\theta} = \text{konst.}$$

$$r^2 \dot{\theta} = p_\theta \dots \text{vertikale Winkelgeschw. ex. abhangt von } r \\ \text{aber } r = r \dot{\theta} r$$

$$e_1: \quad \ddot{r} - \frac{r \dot{\theta}^2}{r^2} = 0 \quad | \cdot r \quad r^3 \ddot{r} = p_\theta^2$$

$$\dot{r} \ddot{r} - \frac{\dot{r} p_\theta^2}{r^2} = 0$$

$$\frac{d}{dt} \left( \frac{1}{2} \dot{r}^2 + \frac{p_\theta^2}{2r^2} \right) = 0$$

$$\frac{\dot{r}^2}{2} + \frac{p_\theta^2}{2r^2} = \text{konst.}$$

Kinetische Energie u. pol. Koord.

integrale gibts genau  $\Leftrightarrow$  obigeine holt eine  
symmetrie u. problem

$$\frac{1}{2} \dot{r}^2 + \frac{1}{2} \frac{p_\theta^2}{r^2} = e \text{ konst.} \quad \text{lassen } r(u) \quad r(\theta(t))$$

$$\dot{r} = \frac{dr}{dt} = \frac{dr}{d\theta} \dot{\theta} = r' \frac{p_\theta}{r^2}$$

$$\frac{1}{2} r'^2 \frac{p_\theta^2}{r^2} + \frac{1}{2} \frac{p_\theta^2}{r^2} = e$$

$$\frac{1}{2} \frac{r'^2}{r^2} + \frac{1}{2} \frac{1}{r^2} = \frac{e}{p_\theta^2}$$

$$u = \frac{1}{r} \quad u' = -\frac{1}{r^2} r' \quad \text{d.h.} \quad r' = -\frac{1}{u^2} u'$$

$$\frac{1}{2} u'^2 + \frac{1}{2} \frac{1}{u^2} = \frac{e}{p_\theta^2}$$

$$u'^2 + u^2 = \text{konst.}$$

$$u = u_0 \cos(\theta - \theta_0)$$

~~$$r = r_0 / \cos(\theta - \theta_0)$$~~

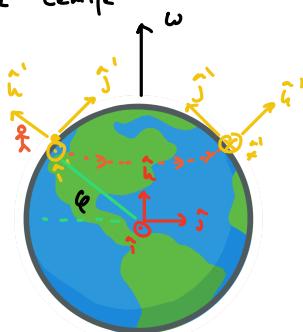
## Vrčaci sistemi

-Inercialni sistemi ... mirujući / euholomno gibanjoci

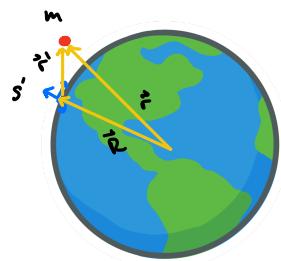
-Neinercialni sistemi ... posjeceni sistemi (npr. vrčaci sistemi)  $S'$

$$\hookrightarrow m \ddot{\vec{a}}_{\text{rel}} = \sum \vec{F} + \vec{F}_{\text{sist}}$$

### ③ Vrčanje Zemlje



$$\begin{aligned}\vec{i}' &= \hat{i} \cos(\omega t) + \hat{j} \sin(\omega t) \\ \vec{j}' &= \hat{k} \cos \varphi + \sin \varphi (\hat{j} \cos(\omega t) - \hat{i} \sin(\omega t)) \\ \vec{k}' &= \hat{k} \sin \varphi + \cos \varphi (-\hat{j} \cos(\omega t) + \hat{i} \sin(\omega t))\end{aligned}$$



$$2. Niz: m \ddot{\vec{r}} = \sum \vec{F} \quad (S)$$

$$\begin{aligned}\ddot{\vec{r}} &= \ddot{\vec{r}}_{\text{rel}} + \ddot{\vec{r}}' \\ \ddot{\vec{r}}' &= \dot{x}' \hat{i}' + \dot{y}' \hat{j}' + \dot{z}' \hat{k}' \\ \ddot{\vec{r}}' &= \ddot{\vec{a}}_{\text{rel}} + \vec{\omega} \times \vec{r}' \\ &\quad \cancel{\ddot{\vec{r}}' = \dot{x}' \hat{i}' + \dot{y}' \hat{j}' + \dot{z}' \hat{k}'}$$

$$\begin{aligned}\hat{i}' \times \hat{j}' &= \hat{k}' \\ \hat{j}' \times \hat{k}' &= \hat{i}' \\ \hat{k}' \times \hat{i}' &= \hat{j}'\end{aligned}$$

$$\begin{aligned}\frac{d}{dt} \hat{i}' &= \frac{d}{dt} (\cos \omega t \hat{i} + \sin \omega t \hat{j}) = \omega (-\sin \omega t \hat{i} + \cos \omega t \hat{j}) \\ &\quad \cancel{- \hat{k}' \times \hat{i}'} \quad \cancel{\hat{k}' \times \hat{i}'} \\ &= \vec{\omega} \hat{k} \times (\underbrace{\sin \omega t \hat{j} + \cos \omega t \hat{i}}_{\hat{r}'}) = \vec{\omega} \times \hat{r}'\end{aligned}$$

$$\begin{aligned}\dot{\hat{i}'} &= \vec{\omega} \times \hat{i}' \\ \dot{\hat{j}'} &= \vec{\omega} \times \hat{j}' \\ \dot{\hat{k}'} &= \vec{\omega} \times \hat{k}'\end{aligned}$$

$$\vec{\omega} = \omega \hat{k}$$

$$\begin{aligned}\ddot{\vec{r}}' &= \frac{d}{dt} (\ddot{\vec{a}}_{\text{rel}} + \vec{\omega} \times \vec{r}') \\ &= \frac{d}{dt} (\dot{x}' \hat{i}' + \dot{y}' \hat{j}' + \dot{z}' \hat{k}' + \vec{\omega} \times \vec{r}') = \vec{\omega} \times \frac{d}{dt} \vec{r}' \\ \ddot{\vec{r}}' &= \ddot{\vec{a}}_{\text{rel}} + \vec{\omega} \times \ddot{\vec{a}}_{\text{rel}} + \vec{\omega} \times (\vec{\omega} \times \vec{r}') \\ &\quad \cancel{\ddot{\vec{r}}' = \dot{x}' \hat{i}' + \dot{y}' \hat{j}' + \dot{z}' \hat{k}'}\end{aligned}$$

$$\vec{r} = R \hat{k}$$

$$\dot{\vec{r}} = R \vec{\omega} \times \hat{k} = \vec{\omega} \times \vec{r}$$

$$\ddot{\vec{r}} = \vec{\omega} \times (\vec{\omega} \times \vec{r})$$

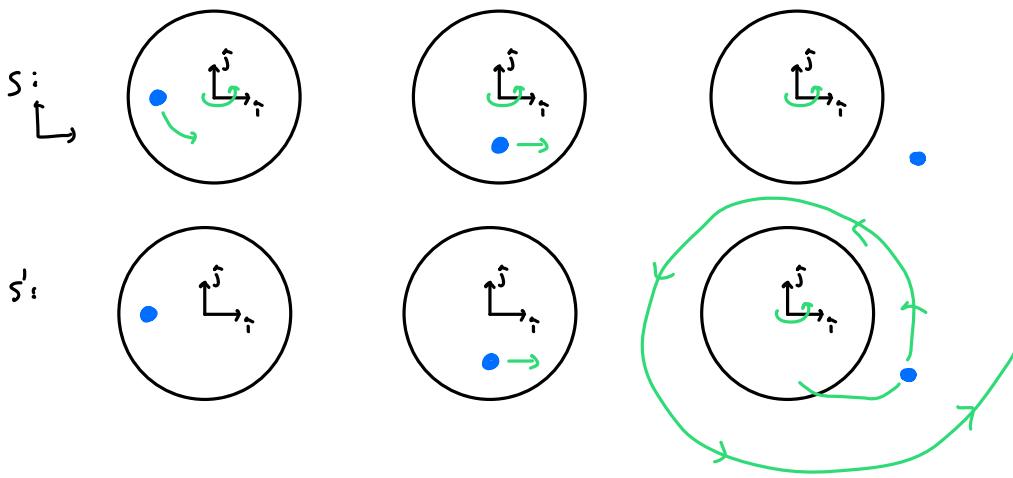
Z.P. z.

$$S: m \ddot{\vec{r}} = \sum \vec{F}$$

$$S': m(\ddot{\vec{r}} + \ddot{\vec{r}}') = \sum \vec{F}$$

$$\begin{aligned}m \ddot{\vec{a}}_{\text{rel}} &= \sum \vec{F} - m \left( \underbrace{2 \vec{\omega} \times \vec{v}_{\text{rel}}}_{\text{coriolisova sila}} + \underbrace{\vec{\omega} \times (\vec{\omega} \times \vec{r}') + \vec{\omega} \times (\vec{\omega} \times \vec{r})}_{\text{centrifugalna sila}} \right) \\ &\quad \underbrace{\text{sistemsku silu}}_{\text{silu}}\end{aligned}$$

8) Prost dalek v různém systému



1. užití:  $\omega \times \vec{r}$

$$\omega \dot{\vec{r}}_{\text{rel}} = 0 - 2\omega \vec{\omega} \times \vec{v}_{\text{rel}} - \omega \vec{\omega} \times (\vec{\omega} \times \vec{r}) \quad | : \omega$$

$$\vec{\omega} \times \vec{v}_{\text{rel}} = \begin{vmatrix} \vec{i} & \vec{j} & \vec{k} \\ 0 & 0 & \omega \\ \dot{x} & \dot{y} & 0 \end{vmatrix} = -\dot{i}\dot{y}\omega + \dot{j}\dot{x}\omega$$

$$\vec{\omega} \times \vec{\omega} \times \vec{r} = \vec{\omega} (\underbrace{\vec{\omega} \cdot \vec{r}}_0) - \vec{\omega}^2 \vec{r} = -\vec{\omega}^2 \vec{r}$$

$$\vec{\omega}_{\text{rel}} = \dot{x}\vec{i} + \dot{y}\vec{j} = -2(-\dot{y}\omega\vec{i} + \dot{x}\omega\vec{j}) + \omega^2(x\vec{i} + y\vec{j})$$

$$\begin{aligned} \ddot{x} &= 2\dot{y}\omega + \omega^2 x \\ \ddot{y} &= -2\dot{x}\omega + \omega^2 y \end{aligned} \quad .i \quad \left. \right\} +$$

$$d = x^i + i y^j \quad \text{Trik}$$

$$\dot{d} = \dot{x}^i + i \dot{y}^j$$

$$id = i\dot{x}^i - \dot{y}^j$$

$$\ddot{d} + 2\omega i \dot{d} - \omega^2 d = 0$$

$$\lambda^2 + 2\omega i \lambda - \omega^2 = 0$$

$$\lambda_{1,2} = \frac{1}{2} (-2\omega i \pm \sqrt{-4\omega^2 + 4\omega^2}) = -\omega i$$

$$d = A e^{\lambda t} + B t e^{\lambda t}$$

$$d(0) = x_0 + i y_0 = A$$

$$\dot{d}(0) = \dot{y}'(0) = 0 \Rightarrow \dot{d} = 0$$

$$d = A \lambda e^{\lambda t} + B e^{\lambda t} + B t \lambda e^{\lambda t}$$

$$\dot{d}(0) = A \lambda + B = 0 \quad B = -A \lambda = \omega i A$$

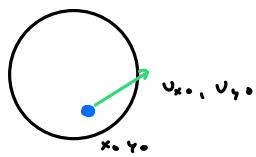
$$d(t) = A e^{-i\omega t} (1 + t; \omega) = e^{-i\omega t} (1 + i\omega t)(x_0 + i y_0)$$

$$= (e^{-i\omega t} - i \sin \omega t)(1 + i\omega t)(x_0 + i y_0) = (e^{-i\omega t} - i \sin \omega t)(x_0 + i y_0 + i \omega t x_0 - \omega t y_0)$$

$$x'(t) = \operatorname{Re}(d) = x_0 (\cos \omega t + \omega t \sin \omega t) + y_0 (\sin \omega t - \omega t \cos \omega t)$$

$$y'(t) = \operatorname{Im}(d) = \cos \omega t (x_0 \omega t + y_0) + \sin \omega t (-x_0 + y_0 \omega t)$$

2. náčin

 $s \xrightarrow{\text{transf.}} s'$ 

$$\vec{v} = \vec{\omega} \times \vec{r}$$

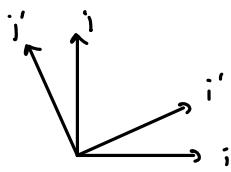
$$(v_{x_0}, v_{y_0}, 0) = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ 0 & 0 & \omega \\ x_0 & y_0 & 0 \end{vmatrix} = (-\omega y_0, \omega x_0, 0)$$

$$\left. \begin{array}{l} x(t) = x_0 - \omega y_0 t \\ y(t) = y_0 + \omega x_0 t \end{array} \right\} \vec{r}$$

$$\left. \begin{array}{l} \hat{i}' = \cos \omega t \hat{i} + \sin \omega t \hat{j} \\ \hat{j}' = -\sin \omega t \hat{i} + \cos \omega t \hat{j} \end{array} \right\} \vec{r}$$

$$\vec{r}' = P \vec{r}$$

Uvodjenje rotacije u kompleksnu



$$x' = x \cos q + y \sin q \quad \vec{x} = R \vec{z}$$

$$y' = -x \sin q + y \cos q$$

$$\vec{z}' = x' + i y' \quad z = x + i y$$

$$\text{Trditev} \quad \vec{z}' = e^{-i q} \vec{z}$$

$$\text{Dokaz} \quad \operatorname{Re} \vec{z}' = x' = \cos q x + \sin q y$$

$$\operatorname{Im} \vec{z}' = y' = -\sin q x + \cos q y$$

⑨ Nitro uihalo na severnem polu

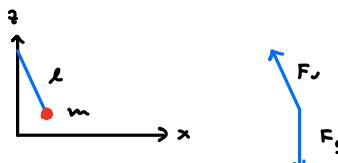
①  $\omega = 0$  žemljic se ne vrati

$$\sin q \approx \frac{x}{L} = q$$

$$z = 1 - \cos q \approx \frac{q^2}{2} \approx 0$$

$$\ddot{z} = 0$$

$$|F_U| \approx m g$$



majhni odniki

$$m \ddot{x} = -mg \frac{x}{L} \Rightarrow \ddot{x} = -\omega_0^2 x \quad \omega_0^2 = \frac{g}{L} \quad \left. \begin{array}{l} \vec{r} = x \hat{i} + y \hat{j} \\ \ddot{\vec{r}} = -\frac{g}{L} \hat{x} \end{array} \right\} \vec{r} = x \hat{i} + y \hat{j}$$

$$\ddot{y} = -\frac{g}{L} y \quad \ddot{\vec{r}} = -\frac{g}{L} \hat{x} \vec{r}$$

②  $\omega > 0$ 

$$\vec{r} = x' \hat{i}' + y' \hat{j}'$$

$$m \ddot{\vec{r}} = -\frac{mg}{L} \vec{r} - 2m \vec{\omega} \times \vec{v}_{rel} - \underbrace{m \vec{\omega} \times \vec{\omega} \times \vec{r}}_{+ m \omega^2 \vec{r}}$$

$$\vec{\omega} \times \vec{v}_{rel} = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ 0 & 0 & \omega \\ x' & y' & 0 \end{vmatrix} = (-\omega y', \omega x', 0)$$

$$\hat{i}' : \quad \ddot{x}' = -\omega_0^2 x' + 2\omega y' + \omega^2 x'$$

$$\hat{j}' : \quad \ddot{y}' = -\omega_0^2 y' - 2\omega x' + \omega^2 y'$$

$$\xi' = x' + iy' \quad \text{je polinomsko z i in sekipenčno}$$

$$\ddot{\xi}' = -\omega_0^2 \xi' - i 2\omega \dot{\xi}' + \omega^2 \xi'$$

$$\text{Naštavlj. } \xi' = \xi_0 e^{-i\omega t}$$

$$-\omega^2 = -\omega_0^2 - 2\omega d + \omega^2$$

$$\omega_d^2 = (\omega - \omega_0)^2$$

$$\omega_d = \omega \pm \omega_0$$

$$\begin{aligned} \xi' &= A e^{-i(\omega + \omega_0)t} + B e^{-i(\omega - \omega_0)t} \\ &= e^{-i\omega t} (A e^{i\omega_0 t} + B e^{i\omega_0 t}) \\ &= e^{-i\omega t} \xi \end{aligned}$$

Vzorci načrtne pogoje:

Os  $t=0$  prilna srednja vrednost v smere  $\hat{n}$

$$\dot{x}' = v_{x0} \quad \dot{y}' = 0$$

$$x' = y' = 0$$

$$\xi' = 0 = A + B \Rightarrow A = -B$$

$$\xi' = A e^{-i\omega t} (e^{-i\omega_0 t} - e^{i\omega_0 t})$$

$$\dot{\xi}' = A e^{i\omega t} (-i(\omega + \omega_0) e^{i\omega_0 t} - (-i\omega + i\omega_0) e^{-i\omega_0 t})$$

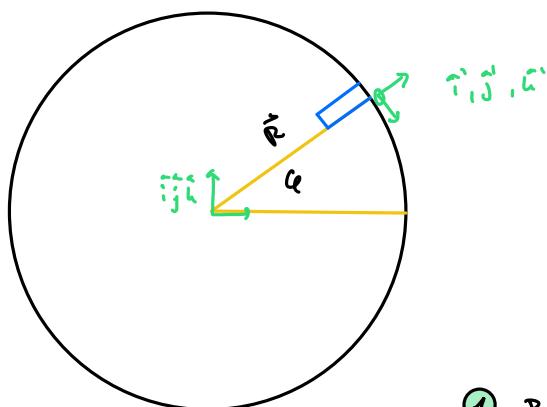
$$\dot{\xi}' = A (-i\omega_0) = v_{x0}$$

$$A = \frac{v_{x0}}{2\omega_0}$$

$$\xi' = \frac{v_{x0}}{2\omega_0} e^{-i\omega t} i (e^{-i\omega_0 t} - e^{i\omega_0 t})$$

$$= \frac{v_{x0}}{\omega_0} e^{-i\omega t} \sin(\omega_0 t)$$

### 10) Padajuči kamnec v jašek



$$S: m \ddot{r} = -mg \frac{\vec{i}}{r}$$

$$S': m \ddot{v}_{rel} = -mg \hat{h}' - 2\vec{\omega} \times \vec{v}_{rel} - \vec{\omega} \times \vec{\omega} \times (\vec{R} + \vec{r}')$$

odvajamo le  
koordinate  
in ne trdi barvih  
vektorjev.

Ker je  $\omega \approx 10^5 \text{ s}^{-1}$

#### ① Perturbativno / iterativno

$$1. \text{ prislužite } m \ddot{v}_{rel}^{(1)} = -mg \hat{h}' \Rightarrow \ddot{v}_{rel}^{(1)} = -gt \hat{h}' \Rightarrow \hat{r}'(t) = -\frac{gt^2}{2} \hat{h}'$$

$$2. \text{ prislužite } \dot{v}_{rel} = -g \hat{h}' - 2\vec{\omega} \times \vec{v}_{rel}^{(1)}$$

$$\vec{\omega} = \omega \hat{h} = \omega (\cos \varphi \hat{j} + \sin \varphi \hat{k})$$

$$\vec{\omega} \times \vec{v}_{\text{rel}} = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ 0 & \omega \cos \varphi & \omega \sin \varphi \\ 0 & 0 & -gt \end{vmatrix} = -gt \cos \varphi \hat{i} \omega$$

Zur Zeit  $t=0$ :

$$\vec{r}(0) = 0$$

$$\vec{v}_{\text{rel}}(0) = 0$$

$$\dot{\vec{v}}_{\text{rel}}^{(0)} = -g \hat{i} + 2gt \omega \cos \varphi \hat{i}$$

$$\vec{v}_{\text{rel}}^{(0)} = -gt \hat{i} + g t^2 \omega \cos \varphi \hat{i}$$

$$\ddot{\vec{r}}^{(0)} = \underbrace{-\frac{gt^2}{2} \hat{i}}_{x'(t)} + \underbrace{\frac{1}{3} g t^3 \omega \cos \varphi \hat{i}}_{x''(t)}$$

$$\dot{\vec{v}}_{\text{rel}}^{(n+1)} = -g \hat{i} t - 2 \vec{\omega} \times \vec{v}_{\text{rel}}^{(n+1)}$$

$$\vec{v}_{\text{rel}}^{(n+1)} = -g \hat{i} t - 2 \vec{\omega} \times \vec{v}_{\text{rel}}^{(n+1)} \quad ] -$$

$$\dot{\vec{v}}_{\text{rel}}^{(n)} - \dot{\vec{v}}_{\text{rel}}^{(n+1)} = -2 \vec{\omega} \times (\vec{v}_{\text{rel}}^{(n+1)} - \vec{v}_{\text{rel}}^{(n)})$$

popravku  $n+1$  korekta

$$\delta \dot{\vec{v}}_{\text{rel}}^{(n)} \quad \delta \vec{v}_{\text{rel}}^{(n+1)}$$

$$\hat{\omega} \times \hat{h}' = (\cos \varphi \hat{j} + \sin \varphi \hat{k}) \times \hat{h}' = \omega \hat{i}'$$

$$\hat{\omega} \times \hat{\omega} \times \hat{h}' = \hat{\omega}(\hat{\omega} \cdot \hat{h}') - \hat{\omega} \hat{h}' = \sin \varphi \hat{\omega} - \hat{h}'$$

$$\hat{\omega} \times \hat{\omega} \times \hat{\omega} \times \hat{h}' = \hat{\omega} \times (\sin \varphi \hat{\omega} - \hat{h}') = -\hat{\omega} \times \hat{h}'$$

$$\delta \vec{v}_{\text{rel}}^{(n)} = -2 \vec{\omega} \times \int_0^t \delta \vec{v}_{\text{rel}}^{(n+1)} dt = -2 \vec{\omega} \times \int_0^t (-2 \vec{\omega} \times \int_0^t \delta \vec{v}_{\text{rel}}^{(n+2)} dt) dt = \dots$$

$$\dots = \underbrace{-2 \vec{\omega} \times (-2 \vec{\omega} \times (-2 \vec{\omega} \times \dots))}_{n-1} \int \dots \int \delta v_{\text{rel}}^{(n)} dt_n dt_{n-1} \dots dt_1$$

$$= (-2 \vec{\omega})^{n-1} (\hat{\omega} \times \hat{\omega} \times \dots \times \hat{\omega}) (-g \frac{t^n}{n!})$$

Komplexe Zapis

Soit: popravku  $\delta \vec{v}_{\text{rel}}^{(n)} = (-2 \vec{\omega})^{n-1} \hat{\omega} \times (\hat{\omega} \times \dots \times \hat{\omega} \times \hat{h}') (-g \frac{t^n}{n!})$

$$= (-2 \vec{\omega})^{n-1} (-i)^n \hat{\omega} \times \hat{h}' (-g \frac{t^n}{(2n)!})$$

lalu: popravku  $\delta \vec{v}_{\text{rel}}^{(2n+1)} = (-i \omega)^{2n} \hat{\omega} \times (\hat{\omega} \times \hat{h}') (-g \frac{t^{2n+1}}{(2n+1)!}) (-i)^n$

$$\vec{v}_{\text{rel}} = \sum_{i=1}^{\infty} \delta \vec{v}_{\text{rel}}^{(i)} = \sum_{n=1}^{\infty} \delta \vec{v}_{\text{rel}}^{(2n+1)} + \sum_{n=0}^{\infty} \delta \vec{v}_{\text{rel}}^{(2n)} =$$

$$= -g \hat{h}' t + \sum_{n=0}^{\infty} \frac{(-2\omega)^{2n+1} (-i)^{2n+1}}{(2n+1)!} \hat{\omega} \times \hat{\omega} \times \hat{h}' (-g \frac{t^{2n+1}}{(2n+1)!}) + \sum_{n=0}^{\infty} \frac{(-2\omega)^{2n}}{(2n)!} (-i)^n \hat{\omega} \times \hat{h}' (-g \frac{t^{2n}}{(2n)!})$$

$$= -g \hat{h}' t + \frac{g}{2\omega} \hat{\omega} \times (\hat{\omega} \times \hat{h}') \underbrace{\sum_{n=0}^{\infty} \frac{(-2\omega t)^{2n+1}}{(2n+1)!} (-i)^n}_{\sin(-2\omega t) - (-2\omega t)} + \frac{g}{2\omega} \hat{\omega} \times \hat{h}' \underbrace{\sum_{n=0}^{\infty} \frac{(-2\omega t)^{2n}}{(2n)!} (-i)^n}_{\cos(-2\omega t) - 1} =$$

$$= -g \hat{h}' t + \frac{g}{2\omega} \hat{\omega} \times (\hat{\omega} \times \hat{h}') (-\sin(2\omega t) - (-2\omega t)) + \frac{g}{2\omega} \hat{\omega} \times \hat{h}' (\cos(2\omega t) - 1)$$

$\sin(-2\omega t) - (-2\omega t)$

$\cos(-2\omega t) - 1$

$$\begin{aligned}
 \ddot{\vec{r}}'(+)&= -g \frac{t^2}{2} \hat{k}' + \frac{g}{2\omega} (\sin \varphi \hat{\omega} - \hat{h}') \left( \frac{\cos(\gamma \omega t)}{2\omega} + \frac{1}{2\omega} + \omega t^2 \right) + \frac{g}{2\omega} (\hat{\omega} \times \hat{h}) \left( \frac{\sin(\gamma \omega t)}{2\omega} - 1 \right) \\
 &= -g \frac{t^2}{2} \hat{k}' + \frac{g}{2\omega} (\sin \varphi (\cos \varphi \hat{j}' + \sin \varphi \hat{h}') - \hat{h}') \left( \frac{\cos(\gamma \omega t) - 1}{2\omega} + \omega t^2 \right) + \frac{g \cos \varphi}{2\omega} \hat{i} \left( \frac{\sin(\gamma \omega t)}{2\omega} - t \right) \\
 &= \hat{i} \frac{g \cos \varphi}{2\omega} \left( \frac{\sin(\gamma \omega t)}{2\omega} - t \right) + \\
 &\quad \hat{j} \frac{g}{2\omega} \sin \varphi \cos \varphi \left( \frac{\cos(\gamma \omega t) - 1}{2\omega} + \omega t^2 \right) + \\
 &\quad \hat{k} (-g \frac{t^2}{2} + \frac{g}{2\omega} (\sin^2 \varphi - 1) \left( \frac{\cos(\gamma \omega t) - 1}{2\omega} + \omega t^2 \right))
 \end{aligned}$$

2. nach:

$$\ddot{\vec{a}}_{rel} + 2\vec{\omega} \times \vec{v}_{rel} = -g \hat{k}$$

$$\vec{\omega} \times \vec{v}_{rel} = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k}' \\ 0 & 0 & \omega \cos \varphi \\ \dot{x} & \dot{y} & \dot{z} \end{vmatrix} = (\omega (\cos \varphi \dot{z}' - \sin \varphi \dot{y}'), \omega \sin \varphi \dot{x}', -\dot{x}' \omega \cos \varphi)$$

$$\dot{x}' = \ddot{x}' + 2\omega (\cos \varphi \dot{z}' - \sin \varphi \dot{y}') = 0$$

$$\dot{y}' = \ddot{y}' + 2\omega \sin \varphi \dot{z}' = 0$$

$$\dot{z}' = \ddot{z}' - 2\omega \cos \varphi \dot{x}' = -g$$

System diff. eqn.

$\begin{bmatrix} \cdot \sin \varphi & \cdot \\ \cdot \cos \varphi & \cdot \end{bmatrix}$

$$\gamma = \cos \varphi \dot{z}' - \sin \varphi \dot{y}'$$

$$\ddot{x}' + 2\omega \dot{y}' = 0$$

$$\ddot{y}' - 2\omega \dot{x}' = -g \cos \varphi$$

$$\xi = x' + i \gamma$$

$$\ddot{\xi} - 2i\omega \dot{\xi} = -g \cos \varphi$$

$$\xi(t) = A e^{2it} + B t + C \quad \text{Nestural}$$

$$\lambda_1 = 0 \quad \lambda_2 = 2i\omega$$

$$\xi(0) = x'(0) + i\gamma(0)$$

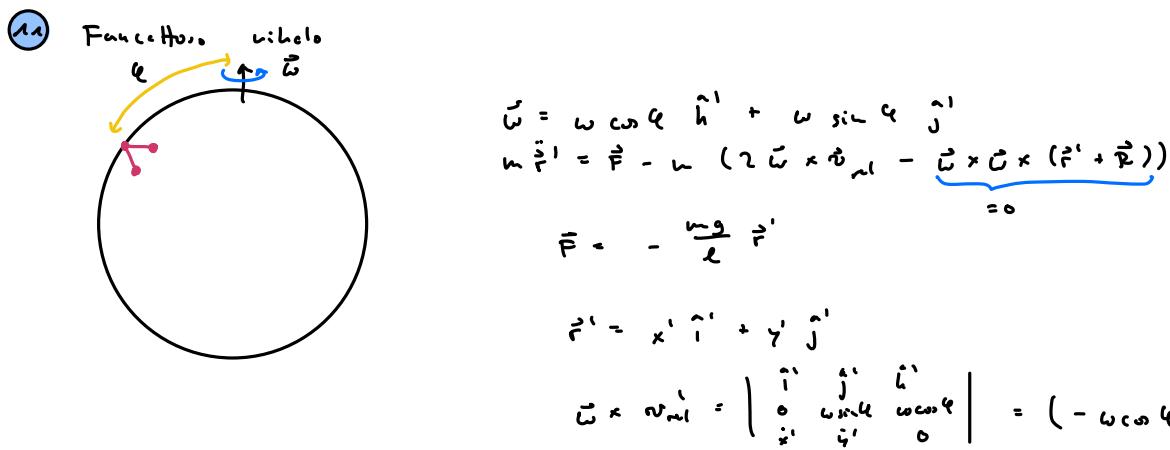
$$\dot{\xi}(0) = \dot{x}'(0) + i\dot{\gamma}(0)$$

$$\Rightarrow A, B, C$$

$$\xi(t) = \frac{g \cos \varphi}{2\omega} \left( \frac{i}{2\omega} (e^{2i\omega t} - 1) + t \right)$$

$$x' = \operatorname{Re} \xi(t)$$

$$\gamma = \operatorname{Im} \xi(t)$$



## D. 2. po koncentr.

$$\ddot{x}' = -\frac{g}{\omega} x' + 2\omega \cos \varphi \dot{\gamma}'$$

$$\ddot{\gamma}' = -\frac{g}{\omega} \dot{\gamma}' - 2\omega \cos \varphi \dot{x}'$$

$$S = x' + i \dot{\gamma}'$$

$$\ddot{S}' = \omega^2 S' - 2\omega \cos \varphi \dot{S}'$$

$$S'(t) = S'_0 e^{-i\omega t}$$

$$-\omega^2 = -\omega_0^2 - 2\omega \cos \varphi \Delta$$

$$\Delta_{\text{res}} = \frac{2\omega \cos \varphi \pm \sqrt{4\omega^2 \cos^2 \varphi + 4\omega_0^2}}{2} = \omega \cos \varphi \pm \omega_0$$

$$S'(t) = e^{-i(\omega \cos \varphi)t} (A e^{i\omega_0 t} + B e^{-i\omega_0 t})$$

## 12 Virialni teorem

Za skupino delcev, ki ima sestoj delitevjo s centralnim potencialom  $V(r) = k/r^d$  velja

$$2\langle T \rangle = d\langle V \rangle$$

$$\langle f \rangle \dots \text{časovna povezava} \quad \langle f(t) \rangle = \lim_{T \rightarrow \infty} \frac{1}{T} \int_0^T f(t) dt$$

$T \dots$  celotna hinc. en.  $V \dots$  celotni pot. en.

- Za en delec

$$m \ddot{\vec{r}} = \vec{F} / r$$

$$m \ddot{\vec{r}} \cdot \vec{r} = \vec{F} \cdot \vec{r}$$

$$m \left( \frac{d}{dt} (\dot{\vec{r}} \cdot \vec{r}) - \dot{\vec{r}} \cdot \dot{\vec{r}} \right) = \vec{F} \cdot \vec{r} \quad \left| \frac{1}{2} \int_0^T dt \right.$$

$$\downarrow \begin{array}{l} \frac{1}{2} m \dot{\vec{r}} \cdot \vec{r} \Big|_0^T - \langle m \dot{\vec{r}} \cdot \dot{\vec{r}} \rangle = \langle \vec{F} \cdot \vec{r} \rangle \\ \rightarrow_0 \text{konst.} \end{array}$$

Če imamo omogočno gibanje  $\vec{r}$  ostane končen

$$\lim_{T \rightarrow \infty}$$

$$0 = 2\langle T \rangle = \langle \vec{F} \cdot \vec{r} \rangle$$

$$\vec{F} = -\text{grad } V(r) \quad V(r) = k/r^d$$

$$= -k d r^{d-1} \text{grad } r$$

$$= -k d r^{d-1} \frac{\vec{r}}{r}$$

$$\text{grad } r = \frac{\vec{r}}{r}$$

$$\vec{F} \cdot \vec{r} = -k d r^d = -dV$$

$$\Rightarrow 2\langle T \rangle = d\langle V \rangle$$

• za u delce

za i-ti delce

$$m \vec{r}_i = \sum_{j \neq i} \vec{F}_{ij} \quad | \cdot \vec{r}_i$$

$$\sum_i \left[ \frac{1}{\pi} \int dt - \underbrace{\left( \frac{d}{dt} (\vec{r}_i \cdot \vec{r}_i) - \dot{\vec{r}}_i \cdot \dot{\vec{r}}_i \right)}_{\rightarrow 0} \right] = \sum_i \sum_{j \neq i} \langle \vec{F}_{ij} \cdot \vec{r}_i \rangle$$

$$- \sum_i \langle m_i \vec{r}_i \cdot \vec{r}_i \rangle = \sum_i \sum_{j \neq i} \langle \vec{F}_{ij} \cdot \vec{r}_i \rangle$$

$$- \sum_j \langle m_i \vec{r}_j \cdot \vec{r}_i \rangle = \sum_j \sum_{i \neq j} \langle \vec{F}_{ji} \cdot \vec{r}_i \rangle$$

$$-k \langle T \rangle = \sum_i \sum_j \langle \vec{F}_{ij} \cdot (\vec{r}_i - \vec{r}_j) \rangle$$

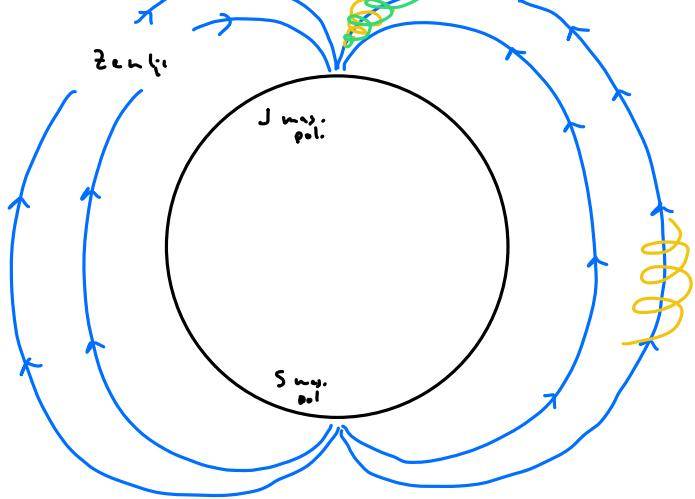
$$F = \text{grad } V = 2 r_{ij}^{d-1} \frac{\vec{F}_i}{r_{ij}} \quad V = k r_{ij}^d \quad \vec{r}_{ij} = \vec{r}_i - \vec{r}_j$$

$$-k \langle T \rangle = - \sum_i \sum_j \langle 2 k r_{ij}^{d-1} \rangle = - \sum_{i \in I} \sum_j 2 \langle k 2 r_{ij}^{d-1} \rangle$$

$$2 \langle T \rangle = 2 \langle V \rangle$$

• Primer: harmoniski potencial  $d=2 \Rightarrow \langle V \rangle = \langle T \rangle$

13) Gibanje nesiteg delca u polju magnetskog monopola (obrazun u severnu ali južnu polu)



$$\vec{B} = \frac{\mu_0 e_m}{4\pi} \frac{\vec{r}}{r^3}$$

$$\vec{F} = e(\vec{E} + \vec{v} \times \vec{B})$$

$$\vec{F} = (j_0 \times \vec{B})e$$

$$m \ddot{\vec{r}} = e \vec{v} \times \vec{B}$$

$$\ddot{\vec{r}} = \frac{e}{m} \vec{v} \times \vec{B}$$

$$\ddot{\vec{r}} = \frac{e \mu_0 e_m}{m 4\pi} \frac{\vec{r}}{r^3} \times \frac{\vec{r}}{r^3}$$

$$\ddot{\vec{r}} = \frac{e^2 g}{m} \frac{\vec{r}}{r^3} \times \frac{\vec{r}}{r^3}$$

(samo konstante slobodne)

$$\textcircled{1} \quad l \cdot \dot{\vec{r}} \quad \ddot{\vec{r}} \cdot \dot{\vec{r}} = 0 = \frac{1}{2} \frac{d}{dt} (\dot{\vec{r}} \cdot \dot{\vec{r}}) \Rightarrow \dot{\vec{r}} \cdot \dot{\vec{r}} = k_0 \text{const} = v_0^2 \Rightarrow |\dot{\vec{r}}| = v_0.$$

bitrost je veća nego

$$\textcircled{2} \quad \vec{l} = \vec{r} \times m \dot{\vec{r}}$$

$$\frac{d}{dt} \vec{l} = \dot{\vec{r}} \times m \vec{r} + \vec{r} \times m \ddot{\vec{r}} = \vec{r} \times m \left( \frac{e}{m} g \vec{r} \times \frac{\vec{r}}{r^3} \right) = \frac{e g}{r^3} \left( \vec{r} r^2 - \vec{r} (\dot{\vec{r}} \cdot \vec{r}) \right)$$

$$\frac{d}{dt} \frac{\vec{r}}{r} = \frac{\dot{\vec{r}}}{r} - \frac{\vec{r} \cdot \ddot{\vec{r}}}{r^2} \quad r = \sqrt{\vec{r} \cdot \vec{r}} \quad \dot{r} = \frac{1}{2} \frac{2 \dot{\vec{r}} \cdot \vec{r}}{\sqrt{\vec{r} \cdot \vec{r}}} = \frac{\dot{\vec{r}} \cdot \vec{r}}{r}$$

$$= \frac{1}{r^2} (\vec{r} r^2 - \vec{r} (\dot{\vec{r}} \cdot \vec{r})) \quad l \cdot \int dt$$

$$\frac{\vec{l}}{g_e} + \frac{\vec{l}_{\text{konst}}}{g_e} = \frac{\vec{r}}{r}$$

$$\textcircled{2} \quad \dot{\vec{L}} = \vec{r} \times (\underbrace{g_e \vec{r} \times \frac{\vec{r}}{r^3}}_{-\frac{g_e \vec{L}}{m r^3}}) = -\vec{r} \times \vec{L} \frac{g_e}{m r^3} \quad \vec{L} \perp \vec{r}, \vec{r}$$

$$\dot{\vec{L}} \cdot \vec{L} = 0 \Rightarrow \frac{d}{dt} \left( \frac{1}{2} \vec{L} \cdot \vec{L} \right) = \vec{L} \cdot \vec{L} = \text{konst.} \Rightarrow L_0 = |\vec{L}| = \text{konst.}$$

zu Zeile 2  $\frac{\vec{r}}{r} - \frac{\vec{L}_{\text{konst}}}{g_e} = \frac{\vec{L}}{g_e} \quad |^2$

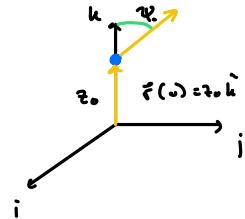
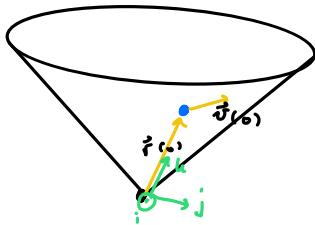
$$\underbrace{\left(\frac{\vec{r}}{r}\right)^2}_{\text{konst.}} + \underbrace{\left(\frac{\vec{L}_{\text{konst}}}{g_e}\right)^2}_{\text{konst.}} - 2 \frac{\vec{r}}{r} \cdot \frac{\vec{L}_{\text{konst}}}{g_e} = \frac{L_0^2}{(g_e)^2}$$

$$\frac{\vec{r}}{r} \cdot \vec{L}_{\text{konst.}} = \text{konst.}$$

$$\hat{e}_r \cdot \vec{L}_{\text{konst.}} = |\vec{L}_{\text{konst.}}| \cos \theta$$



$$\Rightarrow \theta = \text{konst.}$$



$$\vec{v}(t) = v_r(t) \sin \varphi \hat{j} + v_r(t) \cos \varphi \hat{k}$$

$$\begin{aligned} \vec{L}(t) &= m \vec{r}(t) \times \vec{v}(t) \\ &= m z_0 \hat{i} \times v_r(t) \sin \varphi \hat{j} + m z_0 v_r(t) \cos \varphi \hat{k} \\ &= -m z_0 v_r(t) \sin \varphi \hat{i} \end{aligned}$$

$$\begin{aligned} \vec{L}_{\text{konst.}} &= \vec{L}(t) + g_e \frac{\vec{r}}{r}(t) = \\ \vec{L}_{\text{konst.}} &= +m z_0 v_r(t) \sin \varphi \hat{i} + g_e \hat{i} \end{aligned}$$

$$\Theta = \text{konst.}$$

$$\hat{e}_r \cdot \vec{L}_{\text{konst.}} = |\vec{L}_{\text{konst.}}| \cos \theta$$

$$\cos \theta = \frac{\hat{e}_r \cdot \vec{L}_{\text{konst.}}}{|\vec{L}_{\text{konst.}}|} = \frac{g_e}{\sqrt{(g_e)^2 + (m z_0 v_r \sin \varphi)^2}}$$



$$\vec{r}(t) = ? \quad \vec{v}(t) = ?$$

$$\vec{r} \cdot \vec{r} = |\vec{r}| |\vec{r}| \cos \varphi = v_r r \cos \varphi$$

$$\frac{d}{dt} (\vec{r} \cdot \vec{r}) = \underbrace{\vec{r} \cdot \vec{r}}_0 + \vec{r} \cdot \vec{r} = v_r^2 \quad | \int dt$$

$$\vec{r} \cdot \vec{r} = v_r^2 t + c = v_r r \cos \varphi$$

$$\text{p.s. } t=0 \quad c = v_r z_0 \cos \varphi_0$$

$$v_r r \cos \varphi = v_r^2 t + v_r z_0 \cos \varphi_0$$

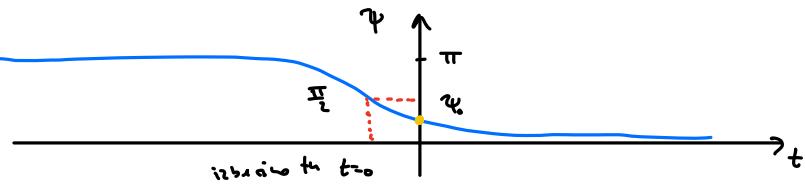
$$\frac{1}{r} |\vec{L}| = \frac{L_0}{r} = |\vec{r} \times m \vec{r}| \frac{1}{r} = |\vec{r} \times \vec{r}| = r v_r \sin \varphi \stackrel{r=t=0}{=} z_0 v_r \sin \varphi_0$$

}  $|^2$  ist wichtig  
definiert

$$r^2 v_r = (z_0 v_r \sin \varphi_0)^2 + (v_r^2 t + v_r z_0 \cos \varphi)^2$$

$$r(t) = \sqrt{(z_0 \sin \varphi_0)^2 + (v_0 t + z_0 \cos \varphi_0)^2}$$

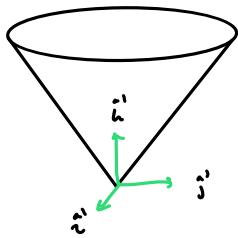
$$\tan \psi = \frac{z_0 v_0 \sin \varphi_0}{v_0 t + z_0 \cos \varphi_0} = \frac{z_0 \sin \varphi_0}{v_0 t + z_0 \cos \varphi_0}$$



zurück:  $\tilde{\omega}_0$   
 $t=0$  da  $\varphi_0 = \frac{\pi}{2}$

$$r(t) = \sqrt{z_0^2 + (v_0 t)^2} \quad \tan \psi = \frac{z_0}{v_0 t} \quad \theta = \text{konst.}$$

Position u. sferischen Koordinaten

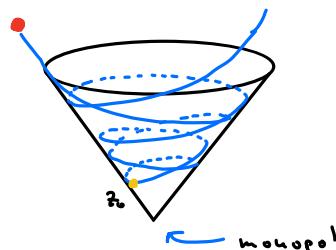


$$\vec{r}(t) = r(t) (\cos \psi \sin \theta, \sin \psi \sin \theta, \cos \theta)$$

$$\ell(t) = \int \frac{z_0 v_0}{r^2 \sin \theta} dt = \frac{z_0 v_0}{\sin \theta} \int \frac{dt}{z_0^2 + (v_0 t)^2},$$

$$\ell(t) = \frac{1}{\sin \theta} \arctan \left( \frac{v_0 t}{z_0} \right)$$

Gitter: diskret



$$N = \frac{\sigma \psi}{2\pi} \quad \text{struktur: außen}$$

$$\Delta \ell = \ell(\infty) - \ell(-\infty) = \frac{\pi}{2 \sin \theta} - \left( - \frac{\pi}{2 \sin \theta} \right) = \frac{\pi}{\sin \theta}$$

$$N = \frac{1}{2 \sin \theta}$$

Lagrangeau formalisieren

①

Newton formulieren

N-diskret

$$\begin{aligned} m_1 \ddot{\vec{r}}_1 &= \vec{F}_1(\vec{r}_1, \dots, \vec{r}_N) \\ &\vdots \\ m_N \ddot{\vec{r}}_N &= \vec{F}_N(\vec{r}_1, \dots, \vec{r}_N) \end{aligned} \quad \left. \right\} 3N \text{ Gleichungen}$$

$$\text{Vor: } f_1(\vec{r}_1, \dots, \vec{r}_N) = 0 \\ \vdots$$

$$f_{N-k}(\vec{r}_1, \dots, \vec{r}_N) = 0 \quad 3N-k \text{ prostostetisch steigend}$$

$$g_1 = g_n(\vec{r}_1, \dots, \vec{r}_N) \quad \text{peripherie Koordinate}$$

$$\vdots$$

$$g_{3N-k} = g_{3N-k}(\vec{r}_1, \dots, \vec{r}_N)$$

## Lagrangeov formulace

$$L = T - V = L(g_1, \dot{g}_1) = L(g_1, \dots, g_{3n-k}, \dot{g}_1, \dots, \dot{g}_{3n-k})$$

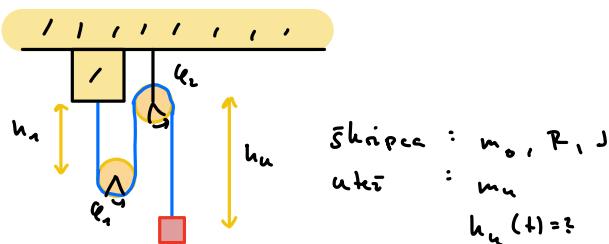
Euler-Lagrangeovu rovnici

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{g}_1} = \frac{\partial L}{\partial g_1} (+ Q_1)$$

físické síly, které vlivují na V.

$$Q_1 = \sum_i \vec{F}_i \cdot \frac{\partial \vec{r}_i}{\partial g_1}$$

(14)



Teorie:

1. zapiši T a V
2. upoštěj užití
3. E-L rovnici

①  $T = \frac{1}{2} m_1 \dot{h}_1^2 + \frac{1}{2} J \dot{\theta}_1^2 + \frac{1}{2} J \dot{\theta}_2^2 + \frac{1}{2} m_2 \dot{h}_2^2$   
 $V = -m_1 g h_1 - m_2 g h_2$

②  $h_1 = -R \dot{\theta}_1$   
 $h_2 = -R \dot{\theta}_2$   
 $l = 2h_1 + h_2 + 2\pi R$  dolejme vrátko

$$\begin{aligned} \dot{h}_1 &= -R \dot{\theta}_1 \\ \dot{h}_2 &= -R \dot{\theta}_2 \\ h_1 &= \frac{1}{2} \dot{h}_2 \end{aligned}$$

$$\begin{aligned} V &= -m_1 g h_1 - \frac{m_2 g}{2} (l - 2\pi R - h_2) = \frac{m_2 g}{2} h_2 - m_1 g h_1 + \text{kost.} \\ T &= \frac{1}{2} m_1 \dot{h}_1^2 + \frac{1}{2} J \frac{\dot{h}_1^2}{R^2} + \frac{1}{2} J \frac{\dot{h}_2^2}{4R^2} + \frac{1}{2} m_2 \frac{1}{4} \dot{h}_2^2 \\ &= \frac{1}{2} \dot{h}_2^2 \left( m_2 + \frac{5}{4} \frac{J}{R^2} + \frac{m_1}{4} \right) \end{aligned}$$

③  $L = T - V = \frac{1}{2} \dot{h}_2^2 \underbrace{\left( m_2 + \frac{5}{4} \frac{J}{R^2} + \frac{m_1}{4} \right)}_{\approx} - h_2 \underbrace{\left( \frac{m_2 g}{2} - m_1 g \right)}_{\approx} + \text{kost.}$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{h}_2} = \frac{\partial L}{\partial h_2}$$

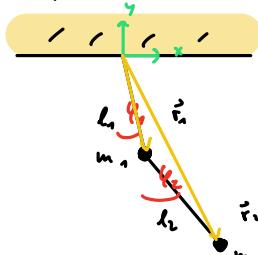
$$\frac{d}{dt} \frac{1}{2} \dot{h}_2^2 = -\frac{1}{2} \ddot{h}_2$$

$$\begin{aligned} \ddot{h}_2 &= -\frac{1}{2} \ddot{h}_2 \\ \ddot{h}_2 &= -\frac{1}{2} \frac{\ddot{h}_2}{\dot{h}_2} = \ddot{\alpha} \end{aligned}$$

$$h_2 = \frac{\ddot{h}_2 t^2}{2} + v_0 t + h_2^0$$

(15)

Dvojina u nicho



-ravníková sítance

-lastní frekvence?

$$(x_1, y_1), (x_2, y_2)$$

$$4 - 2$$

$$2 \text{ verz.}$$

$$\text{Generalizované koordináty: } \vec{r}_1' = \vec{r}_2 - \vec{r}_1$$

$$|\vec{r}_1| = l_1, |\vec{r}_2| = l_2$$

$$T = \frac{1}{2} m_1 \dot{\varphi}_1^2 + \frac{1}{2} m_2 \dot{\varphi}_2^2$$

$$T = \frac{1}{2} m_1 l_1^2 (\cos^2 \varphi_1 + \sin^2 \varphi_1) \dot{\varphi}_1^2 + \frac{1}{2} m_2 \left( \frac{\dot{\varphi}_2^2}{l_1^2} + \frac{\dot{\varphi}_1^2}{l_2^2} + 2 \frac{\dot{\varphi}_1}{l_1} \cdot \frac{\dot{\varphi}_2}{l_2} \right)$$

$$\vec{r}_1 = (l_1 \sin \varphi_1, -l_1 \cos \varphi_1)$$

$$\vec{r}_2 = (l_2 \sin \varphi_2, -l_2 \cos \varphi_2)$$

$$\vec{r}_1 = (l_1 \sin \varphi_1 + l_2 \sin \varphi_2, -l_1 \cos \varphi_1 - l_2 \cos \varphi_2)$$

$$\vec{r}_2 = (l_2 \cos \varphi_2, l_2 \sin \varphi_2)$$

$$T = \frac{1}{2} m_1 l_1^2 \dot{\varphi}_1^2 + \frac{1}{2} m_2 (l_1^2 \dot{\varphi}_1^2 + l_2^2 \dot{\varphi}_2^2 + 2 l_1 l_2 \dot{\varphi}_1 \dot{\varphi}_2 \cos(\varphi_2 - \varphi_1)) \approx 1$$

$$V = m_1 g (-l_1 \cos \varphi_1) + m_2 g (-l_1 \cos \varphi_1 - l_2 \cos \varphi_2)$$

$$\cos x = 1 - \frac{x^2}{2}$$

$$V \approx V_0 + m_1 g l_1 \frac{\dot{\varphi}_1^2}{2} + m_2 g (l_1 \frac{\dot{\varphi}_1^2}{2} + l_2 \frac{\dot{\varphi}_2^2}{2})$$

$$L = T - V$$

$$E - L: \quad \ddot{\varphi}_1: \quad \frac{d}{dt} \frac{\partial L}{\partial \dot{\varphi}_1} = \frac{\partial L}{\partial \varphi_1} \quad m_1 l_1^2 \ddot{\varphi}_1 + m_2 l_1^2 \ddot{\varphi}_1 + m_2 l_1 l_2 \ddot{\varphi}_2 = -m_1 g l_1 \dot{\varphi}_1 - m_2 g l_1 \dot{\varphi}_1$$

$$\ddot{\varphi}_2: \quad \frac{d}{dt} \frac{\partial L}{\partial \dot{\varphi}_2} = \frac{\partial L}{\partial \varphi_2} \quad m_2 l_2^2 \ddot{\varphi}_2 + m_2 l_1 l_2 \ddot{\varphi}_1 + m_2 g l_2 \dot{\varphi}_2 = 0$$

$$\text{Partiell} \quad \dot{\varphi}_1(t) = \dot{\varphi}_{10} e^{-i\omega t} \\ \dot{\varphi}_2(t) = \dot{\varphi}_{20} e^{-i\omega t}$$

$$-m_1 l_1^2 \omega^2 \dot{\varphi}_{10} - m_2 l_1^2 \omega^2 \dot{\varphi}_{10} - m_2 l_1 l_2 \omega^2 \dot{\varphi}_{20} + m_1 g l_1 \dot{\varphi}_{10} + m_2 g l_1 \dot{\varphi}_{10} = 0 \\ -m_2 l_2^2 \omega^2 \dot{\varphi}_{20} - m_2 l_1 l_2 \omega^2 \dot{\varphi}_{10} + m_2 g l_2 \dot{\varphi}_{20} = 0$$

$$A \vec{q}_0 = A \begin{bmatrix} \dot{\varphi}_{10} \\ \dot{\varphi}_{20} \end{bmatrix} = 0 \quad \det A = 0$$

$$\begin{bmatrix} -m_1 l_1^2 \omega^2 - m_2 l_1^2 \omega^2 + m_1 g l_1 + m_2 g l_1 & -m_2 l_1 l_2 \omega^2 \\ -m_2 l_1 l_2 \omega^2 & -m_2 l_2^2 \omega^2 + m_2 g l_2 \end{bmatrix} \vec{q}_0 = 0$$

$$0 = \det A = m_1 m_2 \omega^4 + (\omega_{01}^2 + \omega_{02}^2) m_1 (-m_1 - m_2) \omega^2 + m_2 \omega_{01}^2 \omega_{02}^2 (m_1 + m_2) = 0$$

$$\omega_{01}^2 = \frac{g}{l_1} \quad \omega_{02}^2 = \frac{g}{l_2}$$

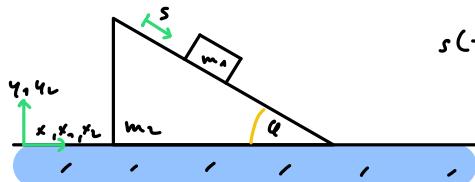
$$\omega_{\pm}^2 = \frac{m_1 + m_2}{2m_1} (\omega_{01}^2 + \omega_{02}^2) \left( 1 \pm \sqrt{1 - \frac{4m_1}{m_1 + m_2} \frac{\omega_{01}^2 \omega_{02}^2}{(\omega_{01}^2 + \omega_{02}^2)^2}} \right)$$

$$\text{zu } m_1 = m_2, \quad l_1 = l_2 \Rightarrow \omega_{01} = \omega_{02}$$

$$\omega_{\pm}^2 = \omega_0^2 (2 \pm \sqrt{2})$$

2 lastige negative Werte

16



$$s(t) = ?.$$

$$x_1 = x + s \cos \theta$$

$$y_1 = -s \sin \theta$$

$$x_2 = x$$

$$y_2 = 0 \quad \text{nach Konst.}$$

$$T = \frac{1}{2} m_1 (\dot{x}_1^2 + \dot{y}_1^2) + \frac{1}{2} m_2 (\dot{x}_2^2 + \dot{y}_2^2)$$

$$T = \frac{1}{2} m_1 ((\dot{x} + \dot{s} \cos \theta)^2 + \dot{s}^2 \sin^2 \theta) + \frac{1}{2} m_2 \dot{x}^2$$

$$T = \frac{1}{2} m_1 (\dot{x}^2 + 2 \dot{x} \dot{s} \cos \theta + \dot{s}^2) + \frac{1}{2} m_2 \dot{x}^2$$

$$V = m_1 g s (-\sin \theta)$$

Generalisieren moment  
x in zibitische Koordinaten

$$x: \frac{d}{dt} \frac{\partial L}{\partial \dot{x}} = \frac{\partial L}{\partial x} \Rightarrow \frac{d}{dt} (m_1 \dot{x} + m_1 \dot{s} \cos \theta + m_2 \dot{x}) = 0 \quad m_1 \ddot{x} + m_1 \ddot{s} \cos \theta + m_2 \ddot{x} = 0$$

$$s: \frac{d}{dt} \frac{\partial L}{\partial \dot{s}} = \frac{\partial L}{\partial s} \Rightarrow \frac{d}{dt} (m_1 \dot{x} \cos \theta + m_1 \dot{s}) = -\sin \theta m_1 g \quad \ddot{x} \cos \theta + \ddot{s} = -\sin \theta g$$

$$\ddot{s} (1 - \frac{m_1}{m_1 + m_2} \cos^2 \theta) = -\sin \theta g \quad \ddot{x} = -\frac{m_1 \cos \theta}{m_1 + m_2} \ddot{s}$$

$$s = \int \int \frac{\sin \theta g}{\frac{m_1}{m_1 + m_2} \cos^2 \theta - 1} dt$$

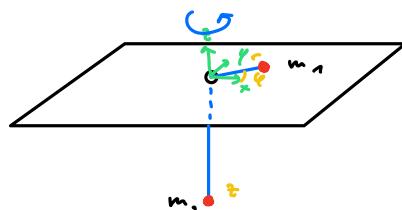
$$s(t) = a \frac{t^2}{2} + At + B$$

Liniär:

$$m_2 \rightarrow \infty \quad a \rightarrow g \sin \theta$$

$$m_2 \rightarrow 0 \quad a \rightarrow \frac{g}{\sin \theta}$$

17) Pložka z dvoj. utěžíma



$$l = r - z \quad 2 \text{ gen. Koordinaten}$$

$$\Rightarrow z = r - l \quad \dot{z} = \dot{r}$$

$$L = \frac{1}{2} m_2 \dot{z}^2 + \frac{1}{2} m_1 (r^2 + r^2 \dot{\theta}^2) + m_2 g z$$

$$= \frac{1}{2} m_2 \dot{r}^2 + \frac{1}{2} m_1 (r^2 + r^2 \dot{\theta}^2) - m_2 g (r - l)$$

$$E-L: \quad r: \frac{d}{dt} (m_2 \dot{r} + m_1 \dot{r}) = m_2 r \dot{\theta}^2 - m_2 g$$

$$(m_1 + m_2) \ddot{r} = m_2 \dot{\theta}^2 r - m_2 g$$

$$\dot{\theta}: \frac{d}{dt} m_1 r^2 \dot{\theta} = 0 \quad p_\theta = m_1 r^2 \dot{\theta}$$

$$m_1 2r \dot{r} \dot{\theta} + m_1 r^2 \ddot{\theta} = 0$$

Stacionární gibaní je v některé oblasti lehká?

$$\begin{aligned} \dot{r} &= 0 \\ \ddot{r} &= 0 \\ r &= r_0 \end{aligned} \Rightarrow m_1 r^2 \dot{\theta} = m_2 g \quad \dot{\theta}^2 = \frac{m_2 g}{m_1 r_0} = \omega_0^2 \quad \theta(t) = \sqrt{\frac{m_2 g}{m_1 r_0}} t$$

Některé zákonům poprvé

$$\begin{aligned} r &= r_0 + x \quad x \ll r_0 \\ (m_1 + m_2) \ddot{x} &- \frac{p_\theta^2}{m_1 r_0^2} + m_2 g = 0 \\ (m_1 + m_2) \ddot{x} &- \frac{p_\theta^2}{m_1 (r_0 + x)^2} + m_2 g = 0 \end{aligned}$$

$$(m_1 + m_2)\ddot{x} - \frac{P_4^2}{m_1 r_0^3 (1 - \frac{x}{r_0})^3} + m_2 g = 0$$

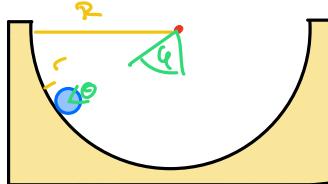
$$(m_1 + m_2)\ddot{x} - \frac{P_4^2}{m_1 r_0^3} (1 - 3\frac{x}{r_0}) + m_2 g = 0$$

$$(m_1 + m_2)\ddot{x} + 3\frac{P_4^2}{m_1 r_0^4} x = 0$$

$$P_4^2 = m_1^2 r_0^4 \dot{\varphi}^2 = m_1^2 r_0^4 \frac{h_1^2 g^2}{h_1 r_0} = m_1 m_2 g r_0^2$$

$$\omega_0^2 = \frac{3 P_4^2}{m_1 r_0^4 (m_1 + m_2)}$$

18) Kroglice v polkrožni ploodi:



$$V_{ext}: \quad x_T = (R-r) \sin \varphi$$

$$y_T = -(R-r) \cos \varphi$$

mrežna skupina

$$r\dot{\theta} = R\dot{\varphi}$$

$$\Rightarrow x, y, \theta \quad 3-2=1 \text{ ges. koord. } \ell$$

$$T = \frac{1}{2} m (\dot{x}_T^2 + \dot{y}_T^2) + \frac{1}{2} J \dot{\theta}^2 \quad J = \frac{2}{3} m r^2$$

$$U = m g y_T = -m g (R-r) \cos \varphi$$

$$L = T - U = \frac{1}{2} m (R-r)^2 \dot{\varphi}^2 + \frac{1}{2} J \left( \frac{R}{r} \dot{\varphi} \right)^2 + m g (R-r) \cos \varphi$$

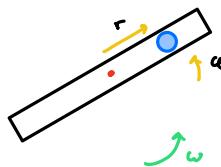
$$E-L: \quad 0 = \frac{d}{dt} \frac{\partial L}{\partial \dot{\varphi}} - \frac{\partial L}{\partial \varphi} = m(R-r)^2 \ddot{\varphi} + 2 \frac{1}{2} \frac{2}{3} m r^2 \frac{R^2}{r^2} \ddot{\varphi} + m g (R-r) \sin \varphi$$

$$0 = \left( m(R-r)^2 + \frac{2}{3} m R^2 \right) \ddot{\varphi} + m g (R-r) \sin \varphi$$

$$\Rightarrow \ddot{\varphi} + \omega_0^2 \sin \varphi = 0 \quad \omega_0^2 = \frac{g (R-r)}{(R-r)^2 + \frac{2}{3} R^2}$$

19) Kroglice v vrteči se cevi:

gledeško drsi:



$$T = \frac{1}{2} m (\dot{r}^2 + r^2 \dot{\varphi}^2) = L$$

$$\text{ver } \dot{\varphi} = \omega \quad l = \frac{1}{2} m (\dot{r}^2 + r^2 \omega^2)$$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{r}} = \frac{d}{dt} m \dot{r} = m \ddot{r} = \frac{\partial L}{\partial r} = m \omega^2 r$$

$$\ddot{r} = \omega^2 r$$

$$r(t) = A e^{\omega t} + B e^{-\omega t} \quad \text{Nastaviš}$$

20) Nihalo na severnem polu

DN



Razni z Lagr. formalizmi in zc opisi: gibanje npravilne krožn.

1

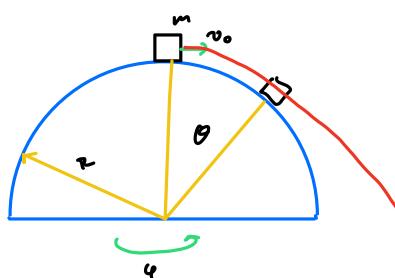
holonomne var:

$$f(g_a) = 0$$

nholonomne var:

$$f(g_a) > 0$$

(2) Utet na ledini polkrogli:



Kdaj se odlep?

- sferične koordinate  $r, \theta, \dot{\theta}$
- $\dot{q} = \omega_{\text{os}}$
- $r = R$  to je delno res  $\Rightarrow r \geq R$

$$\textcircled{1} \quad r = R \quad T = \frac{1}{2} m v^2 = \frac{1}{2} m R^2 \dot{\theta}^2$$

$$V = m g \cos \theta R$$

$$L = T - V = \frac{1}{2} m R^2 \dot{\theta}^2 - m g R \cos \theta$$

$$E-L \quad \theta: \quad m R^2 \ddot{\theta} + m g R (-\sin \theta) = 0$$

$$\ddot{\theta} - \frac{g}{R} \sin \theta = 0 \quad \text{Hiperbolično padanje}$$

2)  $r > R$ 

$$V = m g \cos \theta r$$

$$T = \frac{1}{2} m (\dot{r}^2 + r^2 \dot{\theta}^2)$$

$$L = T - V = \frac{1}{2} m (\dot{r}^2 + r^2 \dot{\theta}^2) - m g r \cos \theta$$

holonomne var, se spremeni s časom, ne velja vek  $r=R$ 

$$\vec{F}_\perp = F_\perp \hat{e}_r$$

$$E-L \quad \frac{d}{dt} \frac{\partial L}{\partial \dot{\theta}} - \frac{\partial L}{\partial \theta} = Q_\theta \quad \text{generalizirani silki}$$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{r}} - \frac{\partial L}{\partial r} = Q_r$$

$$Q_\theta = \sum_i \vec{F}_i \cdot \frac{\partial \vec{r}}{\partial \theta} = F_\perp \hat{e}_r \cdot \hat{e}_\theta = 0$$

$$Q_r = \sum_i \vec{F}_i \cdot \frac{\partial \vec{r}}{\partial r} = F_\perp \hat{e}_r \cdot \hat{e}_r = F_\perp$$

$$\theta: \quad \frac{d}{dt} m r^2 \dot{\theta} - m g r \sin \theta = 0 \\ r^2 \ddot{\theta} + \dot{\theta}^2 r^2 - g r \sin \theta = 0$$

$$r: \quad \frac{d}{dt} m \dot{r} - m r \dot{\theta}^2 + m g \cos \theta = F_\perp$$

nas zadnja količina se odlep:

$$F_\perp = 0$$

$$\dot{r} = \ddot{r} = 0 \quad r = R$$

$$\theta: \quad R^2 \ddot{\theta} - g R \sin \theta = 0$$

$$r: \quad -m R \dot{\theta}^2 + m g \cos \theta = F_\perp = 0$$

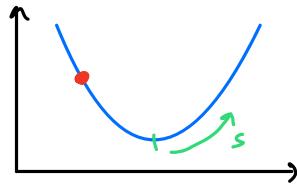
$$\text{Ohranilna energija} \quad E(t) = h_0 \sin t = m g R = m s R \cos \theta + \frac{1}{2} m R^2 \dot{\theta}^2$$

$$\Rightarrow \ddot{\theta} = \frac{2}{\omega^2} \sin \theta (1 - \cos \theta) = \frac{2}{\omega^2} (1 - \cos \theta)$$

$$-\omega^2 \frac{2}{\omega^2} (1 - \cos \theta) + \frac{2}{\omega^2} \omega \theta = 0$$

$$\omega \theta = \frac{2}{\omega}$$

22) Lösung: Vektorlinie  $\gamma(s)$ , ferner, da so nützlich harmonische (wiederholte) oder schwingende



$$T = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2) = \frac{1}{2} m \dot{s}^2$$

$$V = \frac{k s^2}{2} = m g \gamma \quad \gamma = \frac{ds^2}{2} \quad \gamma(s) \quad x(s) \quad ?$$

$$ds^2 = dx^2 + dy^2 \quad / : ds^2$$

$$1 = \left( \frac{dx}{ds} \right)^2 + \left( \frac{dy}{ds} \right)^2$$

$$1 = \left( \frac{dx}{ds} \right)^2 + (ds)^2$$

$$x = \int \sqrt{1 - \dot{x}^2 s^2} \, ds \quad ds = \omega t$$

$$dx = \int \sin t \left( -\frac{\sin t}{\omega} \right) dt$$

$$x = -\frac{1}{\omega} \int_0^t \frac{1 - \cos^2 t}{2} dt = -\frac{1}{\omega} \left( t - \frac{1}{2} \sin 2t \right)$$

$$\gamma = \frac{dx}{ds} = \frac{1}{2\omega} \cos^2 t = \gamma(t) = \frac{1}{4\omega} (1 + \cos 2t)$$

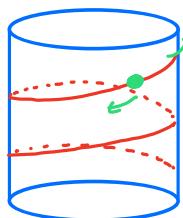
$$\tilde{t} = 2t \quad x(\tilde{t}) = -\frac{1}{4\omega} (\tilde{t} - \sin \tilde{t})$$

$$\gamma(\tilde{t}) = \frac{1}{4\omega} (1 + \cos \tilde{t})$$

cikloide



23) Vektorlinie spiralförmig



Vektoren:

$$r = R$$

$$z = z_0 + p (\varphi - \phi)$$

Koordinatent:

$$-kugelsticke \quad z, \theta, r$$

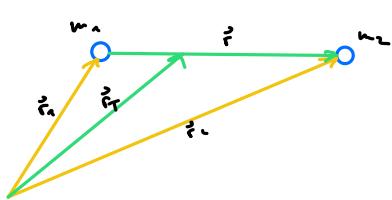
$$-vektoren \quad \phi \dots \text{zusätzl. Vektoren}$$

$T$  ... Kreisbogen in  $r$   $\Rightarrow L$

$V$  ... Kreisbogen

$$\varphi(t) = ? \quad \phi(t) = ?$$

24) Gitarre durch Teilchen perzentral stabil



$$\vec{r}_1, \vec{r}_2 \rightarrow \vec{r}_T, \vec{r} \\ \vec{r}_T = \frac{1}{m_1+m_2} (m_1 \vec{r}_1 + m_2 \vec{r}_2) \\ \vec{r} = \vec{r}_{12} - \vec{r}_T$$

$$L = \frac{1}{2} m_1 \dot{\vec{r}}_1^2 + \frac{1}{2} m_2 \dot{\vec{r}}_2^2 - U(\vec{r}, \vec{r}) \quad \text{zentrale potentiell}$$

$$L = \frac{1}{2} \mu \dot{\vec{r}}^2 + \frac{1}{2} (m_1 + m_2) \dot{\vec{r}}_T^2 - U(\vec{r}, \vec{r})$$

$$\text{reduzierbare } \frac{1}{\mu} = \frac{1}{m_1} + \frac{1}{m_2}$$

$\vec{r}_T = (x_T, y_T, z_T)$  ... cirkular Koordinaten  $\Rightarrow$  paralleler impulse

$$\vec{p}_T = \frac{\partial L}{\partial \dot{\vec{r}}_T} = \text{konst.}$$

$$\tilde{L} = \frac{1}{2} \mu \dot{\vec{r}}^2 - U(\vec{r}, \vec{r})$$

Zuletzt je gesuchte Bewegungslaw?

$$\vec{F} = \vec{r} \times \vec{\vec{F}} = -\nabla U(\vec{r}) = -\frac{\partial U}{\partial \vec{r}} \vec{r} \\ \vec{p} = \vec{r} \times \vec{F} = 0 \quad \vec{r} \parallel \vec{F} \Rightarrow \vec{p} = \vec{p}_T = \text{konst.} \Rightarrow \text{ruhende Gitarre} \\ \vec{n} = \vec{r} \times \vec{p}$$

$$L = \frac{1}{2} \mu \dot{\vec{r}}^2 - U(\vec{r}, \vec{r}) \\ = \frac{1}{2} \mu (r^2 + r^2 \dot{\theta}^2 + r^2 \sin^2 \theta \dot{\varphi}^2) - U(r)$$

$$E-L \quad \theta: \quad p_\theta = \frac{\partial L}{\partial \dot{\theta}} = \mu r^2 \sin^2 \theta \dot{\varphi} \quad \rightarrow \dot{\varphi} = \frac{p_\theta}{\mu r^2 \sin^2 \theta}$$

$$\theta: \quad \frac{d}{dt} \frac{\partial L}{\partial \theta} = \frac{\partial L}{\partial \dot{\theta}} \Rightarrow \frac{d}{dt} (\mu r^2 \dot{\theta}) = \mu r^2 \sin \theta \cos \theta \dot{\varphi}^2$$

$$\frac{d}{dt} = \frac{d\theta}{dt} \frac{d}{d\theta} = \dot{\theta} \frac{d}{d\theta}$$

$$\dot{\theta} \frac{d}{d\theta} (\mu r^2 \dot{\theta} \frac{d\theta}{d\theta}) = \mu r^2 \sin \theta \cos \theta \dot{\varphi}^2$$

$$\frac{d}{d\theta} (\mu r^2 \frac{p_\theta}{\mu r^2 \sin^2 \theta} \frac{d\theta}{d\theta}) = \mu r^2 \sin \theta \cos \theta \frac{p_\theta}{\mu r^2 \sin^2 \theta}$$

$$\frac{d}{d\theta} \left( \frac{1}{\sin^2 \theta} \frac{d\theta}{d\theta} \right) = \cot \theta$$

$$-\frac{d^2}{d\theta^2} (\cot \theta) = \csc \theta$$

$$\cot \theta = A \cos(\alpha - \alpha_0)$$

$$\vec{r} = r \sin \theta \cos \varphi \hat{i} + r \sin \theta \sin \varphi \hat{j} + r \cos \theta \hat{k} \\ \vec{n} = \sin \theta \cos \alpha_0 \hat{i} + \sin \theta \sin \alpha_0 \hat{j} + \cos \theta \hat{k}$$

$$\vec{r} \cdot \vec{n} = 0$$

$$r \sin \theta \cos \varphi \sin \theta \cos \alpha_0 + r \sin \theta \sin \varphi \sin \theta \sin \alpha_0 + r \cos \theta \cos \theta \cos \alpha_0 = 0 \quad | \frac{1}{r \sin \theta \cos \theta}$$

$$0 = \cos\theta \cos\phi_0 + \sin\theta \sin\phi_0 + \cos\theta \cos\theta_0$$

$$0 = \cos(\phi - \phi_0) + \cos\theta \cos\theta_0$$

$$\cos(\phi - \phi_0) = -\cos\theta \cos\theta_0$$

$$A = -\frac{1}{\cos\theta_0}$$

Orbital energy

$$H = T + V = \underbrace{\frac{1}{2} \mu r^2}_{T} + \underbrace{\frac{P_e^2}{2\mu r^2}}_{V} + U(r)$$

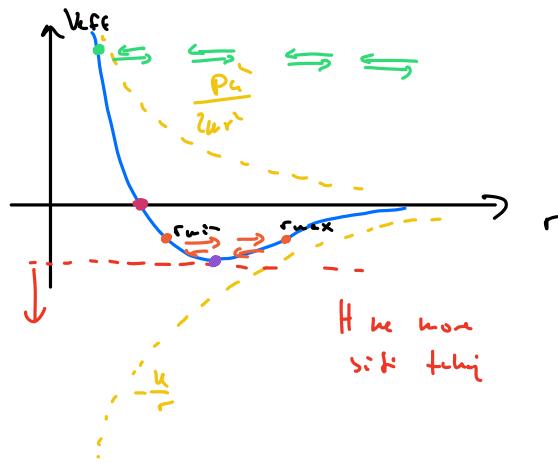
$$H = \frac{1}{2} \mu r^2 + V_{\text{eff}}(r)$$

$$V_{\text{eff}}(r) = U(r) + \frac{P_e^2}{2\mu r^2}$$

Keplar's problem - heliocentric orbit

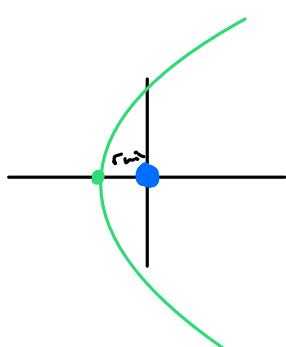
$$U(r) = -\frac{k}{r}, \quad k > 0$$

$$H = \frac{1}{2} \mu r^2 + V_{\text{eff}}(r) \Rightarrow \frac{1}{2} \mu r^2 = H - V_{\text{eff}}(r) \geq 0$$

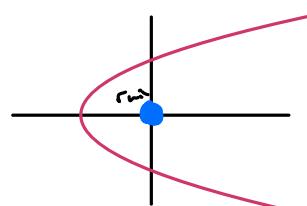


$$H \geq V_{\text{eff}}$$

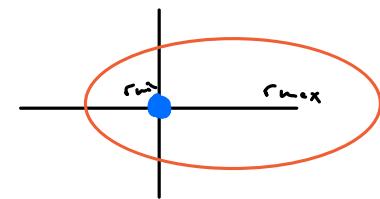
$$V_{\text{eff}}(r) = -\frac{k}{r} + \frac{P_e^2}{2\mu r^2}$$



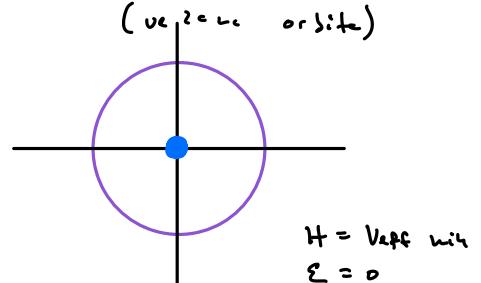
$H > 0$        $\varepsilon > 1$   
hyperbola  
(unbound orbits)



$H = 0$        $\varepsilon = 1$   
parabola  
(boundless orbits)



$H < 0$        $\varepsilon < 1$   
ellipse  
(closed orbits)



$H = V_{\text{eff}}$  with  
 $\varepsilon = 0$

# Orbit 1. Kugelges. problem

$$r(u) = ?$$

$$H = \frac{1}{2} \mu r^2 + \frac{p_u^2}{2\mu r^2} - \frac{k}{r} \quad u = \frac{1}{r} \quad \frac{d}{dt} = \frac{du}{dt} \frac{d}{du} = \dot{u} \frac{d}{du} = \frac{p_u}{\mu r^2} \frac{d}{du}$$

$$\dot{r} = \frac{dr}{dt} = \frac{p_u}{\mu r^2} \frac{du}{dt} = \frac{p_u}{\mu} u^2 \frac{d}{du} \frac{1}{u} = \frac{p_u}{\mu} u^2 \left(-\frac{1}{u^2}\right) \frac{du}{dt} = -\frac{p_u}{\mu} u^3$$

$$H = \frac{1}{2} \mu \left(\frac{p_u}{\mu} u^2\right)^2 + \frac{p_u^2}{2\mu} u^2 - k u = \frac{p_u^2}{2\mu} u^4 + \frac{p_u^2}{2\mu} u^2 - k u \quad 1. \frac{2k}{p_u^2}$$

$$u'' + u^2 - \frac{2ku}{p_u^2} u = \frac{2p_u^2 H}{p_u^2} \quad | \frac{d}{du}$$

$$2u' u'' + 2u u' - \frac{2ku}{p_u^2} u^2 = 0$$

$$2u' (u'' + u - \frac{ku}{p_u^2}) = 0$$

$$\textcircled{1} \quad u' = 0$$

$$\textcircled{2} \quad u'' + u = \frac{ku}{p_u^2}$$

$$\Rightarrow u = \text{konst}$$

$$u = B \cos(\varphi - \varphi_0) + \frac{ku}{p_u^2}$$

$$\Rightarrow r = \frac{1}{u} = \text{konst}$$

$$\Rightarrow \text{Kreisbahnen}$$

$$\varphi = \varphi_0 \quad u(\varphi_0) = \frac{ku}{p_u^2}$$

$$u'(\varphi_0) = -D \sin(\varphi - \varphi_0) \Big|_{\varphi=\varphi_0} = 0$$

$$\begin{aligned} & H = \frac{1}{2} \mu r^2 + \frac{p_u^2}{2\mu} u^2 - k u = \\ & \dot{r} \cdot u' = 0 \quad \xrightarrow{\text{Hilfsgleichung}} \\ & H = \frac{p_u^2}{2\mu} \left( D \cos(\varphi - \varphi_0) + \frac{ku}{p_u^2} \right)^2 - k \left( D \cos(\varphi - \varphi_0) + \frac{ku}{p_u^2} \right) \end{aligned}$$

$$\frac{2p_u^2 H}{p_u^2} = D^2 + \frac{2ku}{p_u^2} D + \left(\frac{ku}{p_u^2}\right)^2 - \frac{2ku}{p_u^2} D - 2\left(\frac{ku}{p_u^2}\right)^2$$

$$D^2 = \frac{2p_u^2 H}{p_u^2} + \left(\frac{ku}{p_u^2}\right)^2$$

$$r(u) = \frac{1}{u} = \frac{1}{\frac{ku}{p_u^2} - \sqrt{\frac{2p_u^2 H}{p_u^2} + \left(\frac{ku}{p_u^2}\right)^2} \cos(\varphi - \varphi_0)} = \\ \text{mit vereinfachter Schreibweise für die Form}$$

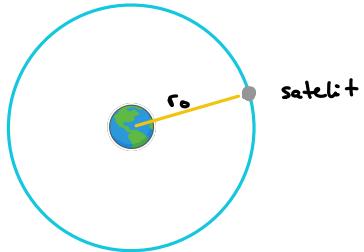
$$p = \frac{p_u^2}{\mu k}$$

$$\varepsilon = \sqrt{1 + \frac{2p_u^2 H}{k^2 \mu}}$$

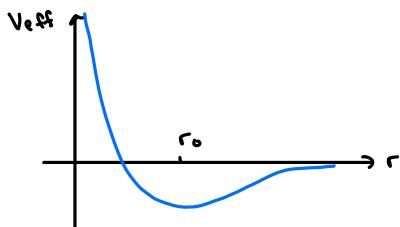
exzentrität

$$= \frac{p}{1 - \varepsilon \cos(\varphi - \varphi_0)}$$

(25) Satellit



$V_{\text{satellit}}$  se satelli asteroide tolle, da j.  $\Delta T = -0,02T$   
 $r_{\min} = ?$  ( $p_0$  trk)



① Punkt trk -  $\frac{p_0^2}{2\mu r^2} - \frac{\alpha}{r}$  ( $\mu = \mu$ )

$$\begin{aligned}\frac{\partial V_{\text{eff}}}{\partial r} &= 0 \\ -\frac{p_0^2}{\mu r^3} + \frac{\alpha}{r^2} &= 0 \quad |_{r=r_0} \\ \Rightarrow r_0 &= \frac{p_0^2}{\alpha \mu}\end{aligned}$$

$$V_{\text{eff}}(r) = \frac{\alpha}{2} \frac{r_0^2}{r^2} - \frac{\alpha}{r}$$

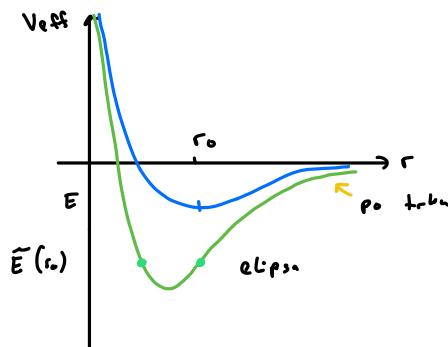
$$E = \frac{p_0^2}{2\mu r} + \frac{1}{2} \mu r^2 - \frac{\alpha}{r} = \underbrace{\frac{\alpha}{2r_0}}_T - \underbrace{\frac{\alpha}{r_0}}_U = \frac{\alpha}{2r_0} \quad (-2T+U)$$

②  $P_0$  trk

$$p_0^2 \rightarrow p_0^2 (1-\eta)$$

$$E \rightarrow \tilde{E} = \frac{\alpha}{2r_0} (1-\eta) - \frac{\alpha}{r_0}$$

$$V_{\text{eff}} \rightarrow \tilde{V}_{\text{eff}}(r) = \frac{\alpha}{2} \frac{r_0^2}{r^2} (1-\eta) - \frac{\alpha}{r}$$



$V$  ekstremum tolle

$$\tilde{E} = V_{\text{eff}}(r)$$

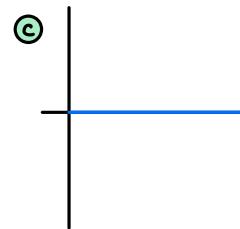
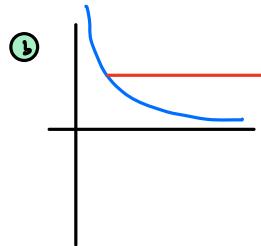
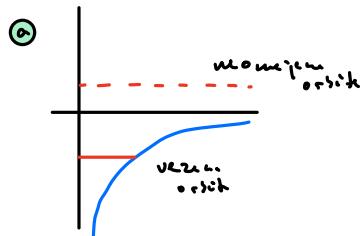
$$\frac{\alpha}{2r_0} (1-\eta) - \frac{\alpha}{r_0} = \frac{\alpha}{2} \frac{r_0 (1-\eta)}{r^2} - \frac{\alpha}{r} \quad \frac{r^2}{r_0 \alpha}$$

$$\frac{r^2}{r_0^2} \left( \frac{1-\eta}{2} \right) = \frac{1-\eta}{2} - \frac{r}{r_0}$$

$$\frac{r}{r_0} = \frac{1 \pm \sqrt{1}}{1+\eta} \quad \Rightarrow \quad r_1 = r_0 \quad r_2 = r_{\min} = r_0 \frac{1-\eta}{1+\eta}$$

(26)

$$\begin{aligned}V &= -\frac{\alpha}{r^2} \\ E &= \frac{1}{2} \mu r \dot{r}^2 + \left( \frac{p_0^2}{2\mu r^2} - \frac{\alpha}{r} \right) \stackrel{!}{=} \underline{V_{\text{eff}}}\end{aligned}$$



$$\frac{dr}{dt} = \frac{dr}{du} \cdot u = r^2 \cdot \frac{p_u}{u^2} \quad p_u = u r^2 \dot{u}$$

$$E = \frac{1}{2} \frac{p_u^2}{r^2} - \frac{r'^2}{r^4} + \left( \frac{p_u^2}{2u} - \alpha \right) \frac{1}{r^2} \quad u = \frac{1}{r} \quad r' = -\frac{1}{u^2} u' \quad r'^2 = \frac{1}{u^4} u''$$

$$E = \frac{p_u^2}{2u} u^2 + \left( \frac{p_u^2}{2u} - \alpha \right) u^2 \quad | \frac{2u}{p_u^2}$$

$$\frac{2u-E}{p_u^2} = u^2 + \left( 1 - \frac{2u}{p_u^2} \right) u^2$$

$$e = u^2 + (1 - \tilde{\alpha}) u^2 \quad | \frac{d}{du}$$

$$0 = 2u' u'' + (1 - \tilde{\alpha}) 2u u' \quad | u' \neq 0$$

$$u'' + (1 - \tilde{\alpha}) u = 0$$

(a)  $\tilde{\alpha} = (1 - \tilde{\alpha}) < 0 \quad u'' - \beta^2 u = 0 \quad u = A e^{-\beta u} + B e^{\beta u}$

(i)  $u = \frac{1}{r_0} \sin \beta u$

(ii)  $u = \frac{1}{r_0} \cosh \beta u$

$$r = \frac{r_0}{\cosh \beta u}$$



(b)  $1 - \tilde{\alpha} = 0$

$$e = u'^2 \quad (i) \quad u' = \sqrt{e}$$

$$u = \sqrt{e} (u - u_0) \quad r = \frac{1}{\sqrt{e}(u - u_0)} \text{ etc}$$

$$(ii) \quad u' = \text{konst.} \Rightarrow \text{lineare}$$

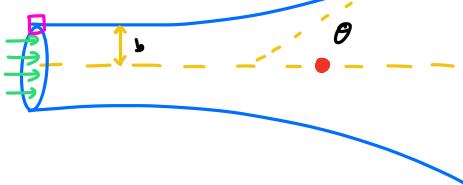


(c)  $\tilde{\alpha} = 1 - \tilde{\alpha} > 0 \quad u'' + \beta^2 u = 0 \quad u = \frac{1}{r_0} \cos (u - u_0) \beta \quad r = \frac{r_0}{\omega(u - u_0) \beta}$



(27) Sipalni je sijalni presel

pri kateri  $d\Omega$  je  $\theta$  nov sijalni presel po vzdoljini delca



$$\frac{dN}{dt} (v d\Omega) = b \, db \, d\Omega \quad \text{gostota flake vzdoljih delcev}$$

$$\frac{dN}{dt} \frac{(v d\Omega)}{j \, d\Omega} = \sigma (\Omega) = \frac{b \, db \, d\Omega}{d\Omega} = \dots \quad \text{diferencialni sijalni presel}$$

$$\sigma (\Omega) = \frac{d\sigma_{TOT}}{d\Omega}$$

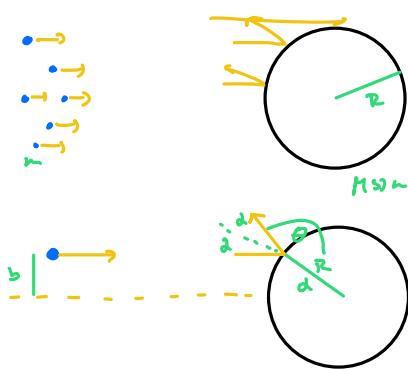
celotni sijalni presel

s tem  
definira da  
deleni sijalni ... =  $\frac{b \, db \, d\Omega}{d(\cos \theta) d\Omega} =$

medenim od

$$j = d\Omega \quad \sigma (\Omega) = \left| \frac{b \, db}{\sin \theta \, d\theta} \right|$$

28 Sípanje u tosi minimoz toci

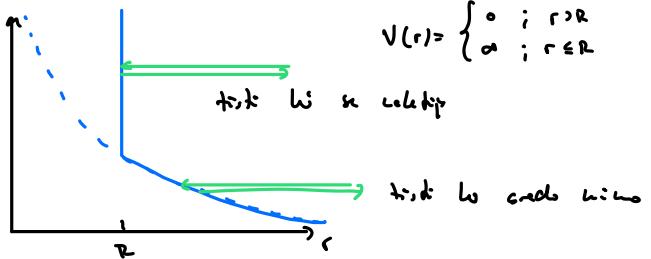


$$H = \frac{1}{2} m v^2 + V_{\text{eff}}(r)$$

$$V_{\text{eff}} = \frac{p_\theta^2}{2mr} + V(r)$$

$V_{\text{eff}}$

$$V(r) = \begin{cases} 0 & r > R \\ \infty & r \leq R \end{cases}$$



Odbijanje zahol

$$2\lambda + \Theta = \pi \Rightarrow b(\Theta) = R \sin \lambda = R \sin\left(\frac{\pi}{2} - \frac{\Theta}{2}\right) = R \cos \frac{\Theta}{2}$$

$$\frac{b}{R} = \sin \lambda$$

$$\sigma(\Theta) = \frac{b}{\sin \Theta} \left| \frac{db}{d\Theta} \right| = \frac{R \cos \frac{\Theta}{2}}{\frac{R}{2} \sin \Theta} \left| \frac{R}{2} \sin \frac{\Theta}{2} \right| = \frac{R^2}{4} = \text{konst.}$$

19. u. 2. uve

? međi

$\sigma_{\text{tot}}$

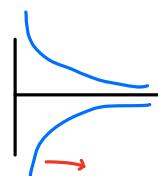
izotropno sípanje

29 Presel u rotirajušem okružju

$$V(r) = -\frac{k}{r^n} \quad n \geq 2, V \geq 0$$

$$\textcircled{1} \quad n=2 \quad H = \frac{1}{2} m v^2 + V_{\text{eff}}$$

$$V_{\text{eff}} = \frac{1}{r^2} \left( \frac{p_\theta^2}{2m} - k \right) \leq 0$$



to posess rotirajušem okružju

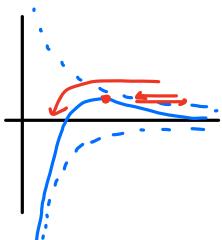
$$\textcircled{2} \quad n=3 \quad \text{minim primitiv} \quad k = \frac{p_\theta^2}{2m} = \frac{(mv_\infty)^2}{2m} \Rightarrow b^2 = \frac{2k}{mv_\infty^2}$$

$$p_\theta = mv_\infty$$



$$\sigma_{\text{tot}} = \pi b^2$$

\textcircled{3} \quad n=2 \quad \text{vremenski} \quad n=3



$$V_{\text{eff}} = \frac{p_\theta^2}{2mr^2} - \frac{k}{r^n}$$

$$0 = \left. \frac{dV_{\text{eff}}}{dr} \right|_{r=r_0} = -\frac{p_\theta^2}{mr^3} + \frac{nk}{r_0^n} \Rightarrow r_0 = \sqrt[n]{\frac{nk}{p_\theta^2}}$$

$$V_{\text{eff}}(r_0) = \frac{1}{r_0^2} \left( \frac{p_\theta^2}{2m} - \frac{nk}{3kR} \right) = \frac{p_\theta^2}{2m r_0^2} \frac{1}{6}$$

$$H = \frac{1}{2} mv_\infty^2 = V_{\text{eff}}(r_0)$$

$$\text{minim primitiv} \quad \frac{1}{2} mv_\infty^2 = \frac{p_\theta^2}{6mr_0^2} = \frac{\left( \frac{m^2 b_{\text{min}}^2 v_\infty^2}{2} \right)^2}{6m \left( \frac{3k}{p_\theta^2} \right)^2} =$$

$$= \frac{m^2 b_{\text{min}}^2 v_\infty^6}{2 \cdot 27 k^2} \Rightarrow b_{\text{min}}^2 = \left( \frac{k}{mv_\infty^2} \right)^{2/3} \Rightarrow \sigma_{\text{tot}} = \pi b_{\text{min}}^2 = \pi \left( \frac{k}{mv_\infty^2} \right)^{2/3}$$

30 Rutherfordova sponja

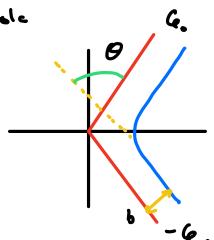
$$V(r) = \frac{q}{r} \quad d = \frac{e_1 e_2}{4\pi \epsilon_0}$$

$$d < 0 \quad r(q) = \frac{p}{1 + \epsilon \cos q} \quad p = \frac{pc^2}{md} \quad \epsilon = \sqrt{1 + \frac{2mc^2 E}{n \alpha^2}}$$

$$d > 0 \quad r(q) = \frac{p}{-1 + \epsilon \cos q} \quad \text{hiperbole}$$

$$\pm q_0 \rightarrow r \rightarrow \infty$$

$$\frac{d\sigma}{d\Omega} \sim \left| \frac{b \sin \theta}{\sin \theta \cot \theta} \right| =$$



$$\theta = \pi - 2q_0$$

$$\sin \frac{\theta}{2} = \cos q_0 = \frac{1}{\epsilon}$$

$$\epsilon \cos q_0 = 1$$

$$q_0 = \arccos \frac{1}{\epsilon}$$

$$\begin{aligned} \cot^2 \frac{\theta}{2} &= \frac{\sin^2 \theta / 2}{\sin \theta / 2} = \frac{1}{\sin^2 \theta / 2} - 1 = c^2 - 1 = \\ &= \frac{2pc^2 E}{md^2} = \frac{2mc^2 E^3}{\alpha^2} = \frac{4E^3 \alpha^2}{d^4} \end{aligned}$$

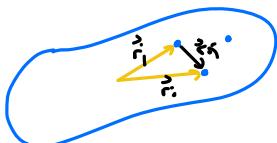
$$\cot \frac{\theta}{2} = \frac{2E^3}{d^2} \quad b = \frac{d}{2E} \cot \frac{\theta}{2}$$

$$\frac{db}{d\theta} = \frac{1}{2E} \left( -\frac{1}{\sin^2 \theta / 2} \right) + \frac{1}{2}$$

$$\sigma(\Omega) = \left( \frac{d}{2E} \right)^2 \frac{\cot \theta / 2}{\sin \theta / 2} \frac{1}{2} \frac{1}{\sin^2 \theta / 2} \frac{1}{2 \sin \theta / 2 \cos \theta / 2} = \left( \frac{d}{4E} \right)^2 \frac{1}{\sin^4 \theta / 2}$$

$$\begin{aligned} \sigma_{tot} : \int \sigma(\Omega) d\Omega \sin \theta d\theta &= 2\pi \left( \frac{d}{4E} \right)^2 \int_0^\pi \frac{1}{\sin^4 \theta / 2} \sin \theta d\theta = 2C \int_0^\pi \frac{\cot \theta / 2}{\sin^3 \theta / 2} d\theta / 2 = \\ &= 4C \left( -\frac{1}{\sin^2 \theta / 2} \right) \Big|_0^\pi = \infty \end{aligned}$$

31 Dinamika toga go telesa



$$|\vec{r}_{ij}| = c_{ij} \dots \text{vsi}$$

prostorialni stupnji - rotacija

$$q, \theta, \omega$$

$$\dot{r} = \vec{v}$$

$$\vec{L} = J \vec{\omega}$$

$\vec{L}$  tensor vibracionski moment

$$(J)_{ij} = \int dm (r^2 \delta_{ij} - r_i r_j)$$

$$\vec{L} = \sum_i m_i \vec{r}_i \times \vec{v}_i \quad \vec{v}_i = \vec{v}_{rel} + \vec{\omega} \times \vec{r}_i$$

$$= \sum_i m_i \vec{r}_i \times \vec{\omega} \times \vec{r}_i =$$

$$= \sum_i m_i (\vec{r}_i \vec{\omega} - (\vec{r}_i \cdot \vec{\omega}) \vec{r}_i) = J \vec{\omega}$$

J simetričan  $\Rightarrow$  lasten vrednost je lasten osi ( $i^*, j^*, k^*$ )

$$\vec{\omega} = \omega_i i^* + \omega_j j^* + \omega_k k^* = \sum_i \omega_i \hat{e}_i$$

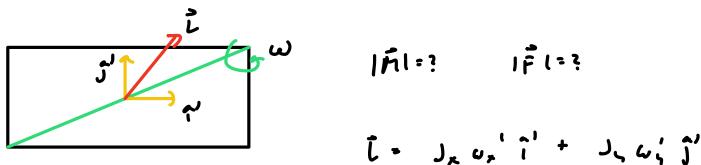
$$\begin{aligned}\vec{\omega} &= \sum_i \omega_i \hat{\epsilon}_i \\ \dot{\vec{\omega}} &= \sum_i \dot{\omega}_i \hat{\epsilon}_i + \sum_i \omega_i \hat{\epsilon}_i' \\ (\vec{\omega} \times \vec{\omega})_k &= \epsilon_{ijk} \omega_i \dot{\omega}_j \\ \dot{\vec{\omega}} &= \sum_i (\dot{\omega}_i \hat{\epsilon}_i + \sum_{j,k} \omega_i \omega_j' \epsilon_{ijk} \hat{\epsilon}_k')\end{aligned}$$

$$\boxed{\dot{\vec{\omega}} = \sum_i (\dot{\omega}_i \hat{\epsilon}_i + \sum_{j,k} \omega_i \omega_j' \epsilon_{ijk} \hat{\epsilon}_k')}$$

$$\begin{aligned}x: M_x &= \dot{\omega}_x = J_x \dot{\omega}_x + \omega_y' \omega_z' (J_z - J_x) \\ y: M_y &= \dot{\omega}_y = J_y \dot{\omega}_y + \omega_z' \omega_x' (J_x - J_y) \\ z: M_z &= \dot{\omega}_z = J_z \dot{\omega}_z + \omega_x' \omega_y' (J_y - J_z)\end{aligned}$$

Berechnung  
einfach

(32) Dauer- und Torsionsschwingungen



$$|F| = ?$$

$$\dot{\vec{\omega}} = J_x \omega_x' \hat{i} + J_y \omega_y' \hat{j} + J_z \omega_z' \hat{k}$$

$$\begin{aligned}J_{11} &= \int d\omega (r^2 \delta_{11} - r_1 r_1) = \int dr g(r) (r^2 \delta_{11} - r_1 r_1) = \int ds g_s (r^2 \delta_{11} - r_1 r_1) \\ &= \frac{m}{ab} \int_{-a/2}^{a/2} dx \int_{-b/2}^{b/2} dy ((x^2 + y^2) \delta_{11} - r_1 r_1)\end{aligned}$$

$$J_{11} = \frac{m}{ab} \iint (x^2 + y^2) dx dy = 0$$

$$J_{22} = \frac{m}{ab} \iint (x^2 + y^2) - x^2 = \frac{m}{ab} a \cdot \frac{1}{2} \left( \left(\frac{a}{2}\right)^2 + \left(\frac{b}{2}\right)^2 \right) = \frac{m b^2}{12}$$

$$J_{33} = \frac{m a^2}{12}$$

$$J_{12} = \frac{m a b}{12}$$

$$J = \frac{m}{12} \begin{pmatrix} b^2 & a^2 & a^2 + b^2 \\ a^2 & b^2 & a^2 + b^2 \\ a^2 + b^2 & a^2 + b^2 & a^2 + b^2 \end{pmatrix}$$

$$M_x = J_x \dot{\omega}_x + (J_y - J_z) \omega_y' \omega_z' = 0$$

$$M_y = J_y \dot{\omega}_y + (J_z - J_x) \omega_z' \omega_x' = 0$$

$$M_z = J_z \dot{\omega}_z + (J_x - J_y) \omega_x' \omega_y' = 0$$

$$\omega = 0$$

$$\omega = \left( \frac{\omega_x}{\sqrt{a^2 + b^2}}, \frac{\omega_y}{\sqrt{a^2 + b^2}}, 0 \right)$$

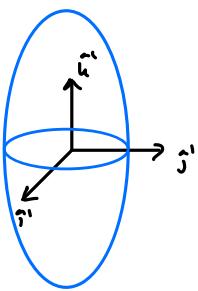
$$M_z = (J_x - J_y) \omega_x' \omega_y'$$

$$M_z = \frac{m}{ab} (a^2 - b^2) \frac{\omega_x \omega_y}{a^2 + b^2}$$

$$\vec{M} \times \vec{\omega} \times \vec{F} = r F = \sqrt{a^2 + b^2} F$$

$$\Rightarrow F = \frac{m}{ab} (a^2 - b^2) \frac{\omega_x \omega_y}{(a^2 + b^2)^{3/2}}$$

33) Proste simmetrische Urteller



$$J = \begin{pmatrix} J & & \\ & J & \\ & & J \end{pmatrix} \quad \bar{\mu} = 0 \quad J_x = J_y$$

Eulerianen Geschle

$$M_x = J_x \dot{\omega}_x + \omega_y \omega_z (J_y - J_z) = 0$$

$$M_y = J_y \dot{\omega}_y + \omega_x \omega_z (J_x - J_z) = 0$$

$$M_z = J_z \dot{\omega}_z + \omega_x \omega_y (J_y - J_x) \Rightarrow \omega_x \dot{\omega}_x = 0 \quad \omega_x = \text{konst}$$

$$\dot{\omega}_x = -\frac{J_z - J_x}{J} \omega_x \omega_y$$

$$\dot{\omega}_y = \frac{J_z - J_x}{J} \omega_x \omega_z$$

$$\Omega_{pr} = \frac{J_z - J_x}{J} \omega_x$$

$$\dot{\omega}_x = -\Omega_{pr} \omega_y$$

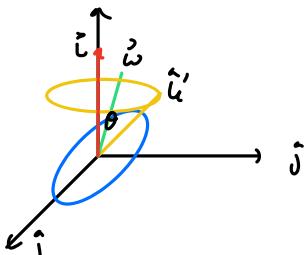
$$\dot{\omega}_y = +\Omega_{pr} \omega_x \quad |.i. \quad \dot{\omega} = \omega_x + i \omega_y$$

$$\dot{\omega} = i \Omega_{pr} \omega \quad \Rightarrow \quad \omega = \omega_0 e^{i \Omega_{pr} t}$$

$$\omega_0 = \omega_{x0} + i \omega_{y0}$$

Woraus ist diese Gleichung in zugehörigen System?

$$\dot{\mu} = 0 \Rightarrow \ddot{\mu} = \text{konst} \quad \ddot{\mu} = L \hat{h}$$



$$\textcircled{1} \quad \Theta = ? \quad \ddot{\mu} = L \cos \Theta = L \frac{1}{\sqrt{1 + \sin^2 \Theta}} = L \frac{1}{\sqrt{1 + \sin^2 \Theta}} = \text{konst} \quad \Rightarrow \quad \Theta = \text{konst.}$$

$$\textcircled{2} \quad \ddot{\mu}, \ddot{\omega}, \ddot{h} \text{ konst. } \& \text{ sind miteinander?} \quad (\ddot{\mu} \times \ddot{\omega}) \cdot \ddot{h} = 0$$

$$(\ddot{\mu} + \ddot{\omega}) \cdot \ddot{h} = 0 = \sum_{ij} \varepsilon_{ijk} L_i \omega_j = L_x \omega_y - L_y \omega_x = \text{konst.} \quad \text{da } J_x = J_y$$

$$\textcircled{3} \quad \ddot{h} = \ddot{\mu} + \ddot{\omega} = (\omega_x \hat{i} + \omega_y \hat{j} + \omega_z \hat{k}) \times \hat{h} = -\omega_x \hat{j} + \omega_y \hat{i}$$

$$|\ddot{h}| = \sqrt{\omega_x^2 + \omega_y^2} = |\ddot{\omega}| = \text{konst.} \quad \left. \begin{array}{l} \text{gibende endokonkrete} \\ \theta = \text{konst.} \end{array} \right\} \text{po stützen}$$

$$\Omega_e = \Omega_e \hat{h}$$

$$\ddot{h} = \Omega_e \hat{h} + \ddot{h}' = \Omega_e \sin \theta \hat{h}$$

$$|\ddot{h}| = \Omega_e |\hat{h} + \ddot{h}'| = \Omega_e \sin \theta = |\ddot{\omega}|$$

$$\Omega_e \sin \theta = \frac{L \sin \theta}{J} \quad \Omega_e = \frac{L}{J}$$

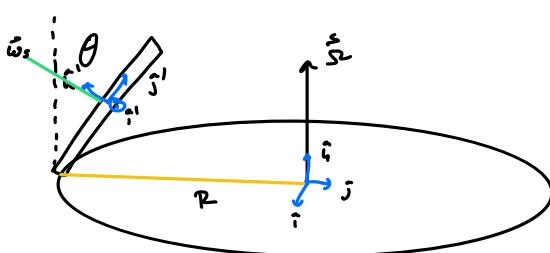
$$\omega_x = \frac{L_x}{J \sin \theta} = \frac{L \cos \theta}{J \sin \theta} \Rightarrow L = \frac{\omega_x J}{\cos \theta}$$

$$\Omega_e = \frac{J}{J \sin \theta} \omega_x$$

$$\textcircled{4} \quad \text{korinik / frische}$$

$$J' = \frac{1}{2} I R^2 \quad \Rightarrow \quad \Omega_e = 2 \omega_x \quad \textcircled{2} \quad \text{Zurück - Chardberg - passend}$$

$$J = \frac{1}{4} I R^2$$



$$\theta = ?$$

$$\begin{aligned}\vec{l} &= \underline{\underline{J}} \vec{r} \\ \underline{\underline{J}} &= \sum \underline{\underline{I}}\end{aligned}$$

$$\underline{\underline{J}} = \begin{pmatrix} J & 0 & 0 \\ 0 & J & 0 \\ 0 & 0 & J' \end{pmatrix} \quad J = \frac{1}{4} \mu r^2 \quad J' = \frac{1}{2} \mu r^2$$

$$\vec{\omega} = \vec{\omega}_s + \vec{\omega}_c$$

$$\vec{\omega}_c = \underline{\underline{J}} \vec{r} = \mu ( \cos \theta \vec{j}' + \sin \theta \vec{h}' )$$

$$\vec{\omega}_s = \omega_s \vec{h}'$$

$$\vec{\omega} = \omega_s \vec{h}' + \mu ( \cos \theta \vec{j}' + \sin \theta \vec{h}' )$$

Vektory (bez svedcovačky)

↳ hmotnost je podleci vektor 0 ( $\vec{r}_p = 0$ )

$$0 = \vec{\omega}_s = \vec{\omega} \times \vec{r} = \vec{\omega}_s \times \vec{r} = \mu R \vec{i}' + \omega_s r \vec{i}' = (\mu R + \omega_s r) \vec{i}' \Rightarrow \omega_s = -\frac{R}{r} \mu$$

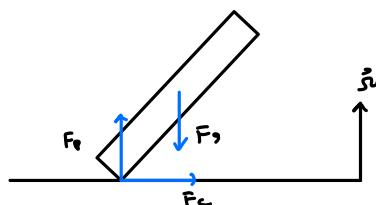
$$\vec{\omega} = \mu ( \cos \theta \vec{j}' + \sin \theta \vec{h}' ) - \mu \frac{R}{r} \vec{h}' = \omega_x \vec{i}' + \omega_y \vec{j}' + \omega_z \vec{h}'$$

$$\dot{\vec{l}} = \frac{d}{dt} ( J \omega_y \vec{j}' + J' \omega_z \vec{h}' ) = J \omega_y \vec{j}' + J' \omega_z \vec{h}'$$

$\theta = \text{konst.}$

$$\begin{aligned}\vec{j}' &= \vec{\omega} \times \vec{j}' = \mu ( \cos \theta \vec{j}' + \sin \theta \vec{h}' ) \times \vec{j}' = -\mu \sin \theta \vec{i}' \\ \vec{h}' &= \vec{\omega} \times \vec{h}' = \mu ( \cos \theta \vec{j}' + \sin \theta \vec{h}' ) \times \vec{h}' = \mu \cos \theta \vec{i}'\end{aligned}$$

$$\Rightarrow \dot{\vec{l}} = \frac{1}{4} \mu r^2 \mu \cos \theta (-\mu \sin \theta) \vec{i}' + \frac{1}{2} \mu r^2 (\mu \sin \theta - \mu \frac{R}{r}) \mu \cos \theta \vec{i}'$$



$$\vec{F}_g = mg (-\vec{h})$$

$$\vec{F}_n = m \vec{a}_n = m \mu^2 (R - r \sin \theta) \vec{j} = m \mu^2 (R - r \sin \theta) (\sin \theta \vec{j}' - \cos \theta \vec{i}')$$

$$\begin{aligned}\vec{F} &= \vec{r} \times (\vec{F}_n + \vec{F}_g) = r(\vec{j}') \times (m g (\cos \theta \vec{j}' + \sin \theta \vec{h}') + m \mu^2 (R - r \sin \theta) (\sin \theta \vec{j}' - \cos \theta \vec{i}')) \\ &= \vec{r}' r (-m g \sin \theta + m \mu^2 (R - r \sin \theta) \cos \theta)\end{aligned}$$

$$\dot{\vec{l}} = \vec{F}$$

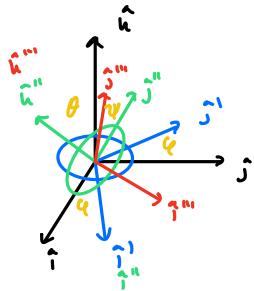
$$\vec{i}' : -\frac{1}{4} \mu r^2 \mu^2 \cos \theta \sin \theta + \frac{1}{2} \mu r^2 \mu^2 \sin \theta \cos \theta - \frac{1}{2} \mu r^2 \frac{R}{r} \cos \theta = -r \mu g \sin \theta + \mu \mu^2 r (R - r \sin \theta) \cos \theta$$

$$\frac{5}{4} \mu r^2 \mu^2 \sin \theta \cos \theta - \frac{3}{2} \mu r^2 \mu^2 \cos \theta = -r \mu g \sin \theta$$

$$\frac{5}{4} r \mu^2 \cos \theta \sin \theta - \frac{3}{2} R \mu^2 \cos \theta = -g r \mu \sin \theta$$

$$\begin{aligned}r \cos \theta &= \frac{3}{2} R \mu^2 \cos \theta = g \sin \theta \\ \tan \theta &= \frac{3}{2} \frac{R \mu^2}{g}\end{aligned}$$

### 35 Eulerjevi koti



$\hat{i}, \hat{j}, \hat{k}$  ... mrežni sistem  
 $\hat{i}', \hat{j}', \hat{k}'$  ... sistem prijet na rotaciju

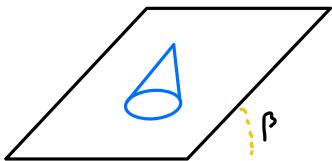
$$\vec{\omega} = \dot{\varphi} \hat{i} + \dot{\theta} \hat{j} + \dot{\psi} \hat{k}$$

$$\begin{aligned}\vec{\omega}' &= \hat{i}' (\dot{\varphi} \sin \theta \sin \psi + \dot{\theta} \cos \psi) + \\ &\quad \hat{j}' (\dot{\varphi} \sin \theta \cos \psi - \dot{\theta} \sin \psi) + \\ &\quad \hat{k}' (\dot{\varphi} + \dot{\theta} \cos \theta)\end{aligned}$$

$$T = \frac{1}{2} \vec{\omega} \cdot \vec{\omega}$$

$$\begin{aligned}T &= \frac{1}{2} J (\dot{\varphi}^2 \sin^2 \theta + \dot{\theta}^2) + \\ &\quad \frac{1}{2} J' (\dot{\varphi} \cos \theta + \dot{\theta})^2\end{aligned}$$

### 36 Stroice u neugivenim podloščem



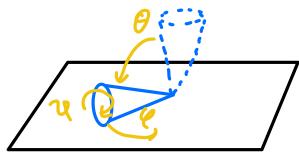
$$V = \frac{\pi r_0^2 h_0}{3} \quad \tan \alpha = \frac{r_0}{h}$$

$$h' = h_0 = \frac{3}{4} h_0$$

$$\begin{aligned}J' &= J_z = \frac{2}{10} m r_0^2 \\ J &= J_x = J_y = \frac{3}{10} m r_0^2 + \frac{2}{5} m h_0^2\end{aligned}$$



$$R = \sqrt{r_0^2 + h^2}$$



$$\text{Pognimo za } \theta \quad \theta = \frac{\pi}{2} - \alpha = 45^\circ$$

$$\text{in zanesave} \quad \dot{\varphi} R = - \dot{\varphi} r_0$$

$$T = \frac{1}{2} J (\dot{\varphi}^2 \sin^2 \theta + \dot{\theta}^2) + \frac{1}{2} J' (\dot{\varphi} \cos \theta + \dot{\theta})^2$$

$$= \frac{1}{2} J (\dot{\varphi}^2 \cot^2 \alpha) + \frac{1}{2} J' (\dot{\varphi} \sin \alpha - \frac{2}{r_0} \dot{\varphi})^2$$

$$T = \frac{1}{2} J \dot{\varphi}^2 \quad \bar{J} = \frac{h^2}{4^2 + h^2} \left( \frac{7}{20} m r_0^2 + \frac{9}{10} m h^2 \right)$$

$$\vec{g} = (\cos \beta (-\hat{i}) + \sin \beta \hat{j}) g$$

$$V = -m \vec{g} \cdot \vec{r}_T$$

$$\vec{r}_T = \frac{2}{3} h (\hat{i} \sin \alpha + \cos \alpha (\hat{i} \cos \alpha + \hat{j} \sin \alpha))$$

$$V = \frac{3}{4} m g h \sin \alpha - \frac{3}{4} m g h \cos \alpha \cos \alpha \sin \beta$$

$$L = \frac{1}{2} \tilde{J} \dot{\varphi}^2 + \frac{3}{4} mgh \cos \alpha \sin \beta \cos \varphi + C$$

$$L = \frac{1}{2} \tilde{J} \dot{\varphi}^2 - \frac{3}{4} mgh \cos \alpha \sin \beta \frac{\omega^2}{2} + C'$$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{\varphi}} = \tilde{J} \ddot{\varphi} = -\frac{3}{4} mgh \cos \alpha \sin \beta \omega^2 \varphi$$

$$\omega^2 = \frac{7mgh \cos \alpha \sin \beta}{4\tilde{J}}$$

### Mala nihauje

Pospolisne koordinate	$g_1, \dots, g_n$	$\underline{g} = \begin{pmatrix} g_1 \\ \vdots \\ g_n \end{pmatrix}$	odnosić
Potencjal	$V(g_1, \dots, g_n)$	$\underline{g} = \underline{g}_0 + \underline{y}$	prawdziwej linię

$$V(\underline{y}) = V(\underline{g}_0) + \frac{1}{2} \sum_{i,j} \frac{\partial^2 V}{\partial g_i \partial g_j} \Big|_{\underline{g}=\underline{g}_0} \underline{y} \cdot \underline{y}$$

Matrycne osłabie

$$V = V_0 + \frac{1}{2} \underline{y}^T \underline{\underline{V}} \underline{y} \quad (\underline{\underline{V}})_{ij} = \frac{\partial^2 V}{\partial g_i \partial g_j} \Big|_{\underline{g}=\underline{g}_0}$$

$$\text{kineticka energia} \quad T = \frac{1}{2} \underline{\dot{y}}^T \underline{\underline{I}} \underline{\dot{y}}$$

$\underline{\underline{V}}, \underline{\underline{I}}$  symetryczne, reelle, pozitivne definitne

$$L = \frac{1}{2} \underline{\dot{y}}^T \underline{\underline{I}} \underline{\dot{y}} - \frac{1}{2} \underline{y}^T \underline{\underline{V}} \underline{y}$$

$$\text{E-L równa} \quad \underline{\underline{I}} \underline{\ddot{y}} + \underline{\underline{V}} \underline{y} = 0$$

$$\underline{y} = \underline{a} e^{-i\omega t} \quad \text{Następstwo}$$

↑   
 relatywne amplitudowe

$$(-\omega^2 \underline{\underline{I}} + \underline{\underline{V}}) \underline{a} = 0 \quad \text{pospolisne lastni problem}$$

$$\det(-\omega^2 \underline{\underline{I}} + \underline{\underline{V}}) = 0 \quad \Rightarrow \quad \omega_n^2 \dots \text{lastni frequencja}$$

↳  $\underline{a}_n \dots$  lastni wektorji / normalni wektori nihauje

Splaszczenie wskutek

$$\underline{y}(t) = \sum_n a_n \begin{cases} d_n e^{-i\omega_n t} & \omega_n^2 > 0 \\ b_n + c_n t & \omega_n^2 = 0 \end{cases} \quad \text{translacja}$$

Ortogonalizacjne zwiazek

$$a_n^T \underline{\underline{I}} a_m = \underline{\underline{I}}^T \delta_{nm} = \delta_{nm}$$

zak. co normalne

$$(\text{metryczne osłabie}) \quad \underline{\underline{A}} = (a_1, a_2, \dots, a_n)$$

$$\underline{\underline{A}}^T \underline{\underline{I}} \underline{\underline{A}} = \underline{\underline{I}}' = \underline{\underline{I}}$$

normalne  
diagonalia

$$\underline{A}^T \underline{V} = \underline{A} \Rightarrow \underline{V}' = \underline{A}^{-1} \text{ diagonal}$$

$\underline{\omega}_n = \begin{pmatrix} \omega_1 & \\ & \ddots & \omega_n \end{pmatrix}$

$$\underline{y}(t) = \begin{pmatrix} \underline{q}_1 \\ \vdots \\ \underline{q}_n \end{pmatrix} = \sum_n q_n \underline{e}_n \quad \underline{e}_n = \begin{pmatrix} 1 \\ \vdots \\ 0 \end{pmatrix}$$

$$\underline{y}(t) = \sum_n p_n(q_n) \underline{e}_n e^{-i\omega_n t} = \sum_n p_n(q_n(t)) \underline{e}_n$$

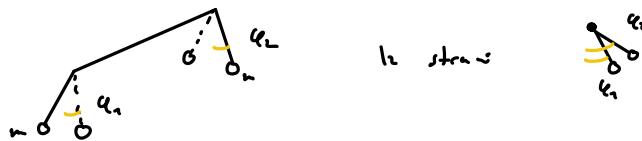
normalisierte Koordinaten

$$L = \sum_n \left( \frac{1}{2} \dot{q}_n^2 - \frac{1}{2} \underline{q}_n^T \underline{d}_n^2 \right) = \sum_n \left( \frac{1}{2} \dot{q}_n^2 - \frac{1}{2} \omega_n^2 q_n^2 \right)$$

$$H = \sum_n \left( \frac{1}{2} \dot{q}_n^2 + \frac{1}{2} \omega_n^2 q_n^2 \right)$$

usoft k  
modellisch  
harmonisch  
oscillatory

③ 2 fachig starke Schwingung mit derselben Amplitude



$$V = mgd((1 - \cos q_1) + (1 - \cos q_2)) + \frac{d}{2}(q_1 - q_2)^2$$

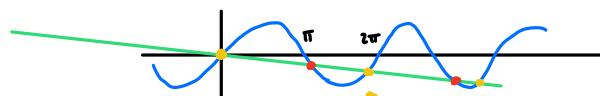
1. Schritt: rechnerische Lsg.

$$\begin{aligned} \frac{dV}{dq_1} &= mgd \sin q_1 + d(q_1 - q_2) = 0 \\ \frac{dV}{dq_2} &= mgd \sin q_2 - d(q_1 - q_2) = 0 \end{aligned} \quad \left. \begin{array}{l} V \text{ rechnerische Lsg.} \\ \hline \end{array} \right.$$

$$\begin{aligned} mgd(\sin q_1 + \sin q_2) &= 0 \quad \Rightarrow \quad q_1 = -q_2 + 2k\pi \quad q_1 > q_2 + \pi + 2k\pi \\ mgd(\sin q_1 - \sin q_2) &= -2d(q_1 - q_2) \end{aligned}$$

④ Rechnerische Lsg., hier  $q_1^0 = -q_2^0$

$$2 \sin q_1 = - \frac{4d}{mgd} q_1$$



$$\begin{array}{ll} q_1^0 & \text{für } r_{\text{sd}} > 2 \\ q_1^0 & \text{für } r_{\text{lih}} \end{array} \quad \begin{array}{ll} q_1 \in [\frac{3\pi}{2}, 2\pi] + 2k\pi & \text{statisch} \\ q_1 \in [\pi, \frac{3\pi}{2}] + 2k\pi & \text{lochig} \end{array}$$

$$\frac{\partial V}{\partial q_1 q_2} = \begin{pmatrix} mgd \cos q_1 + d & -d \\ -d & mgd \cos q_2 + d \end{pmatrix} \stackrel{q_1 = -q_2}{=} \begin{pmatrix} mgd \cos q_1 + d & -d \\ -d & mgd \cos q_1 + d \end{pmatrix} = \underline{U}$$

$$\text{Diagonalmatrix} \quad \underline{a}_1 = (1, 1) \quad \underline{a}_2 = (1, -1)$$

$$\text{Last. rot.} \quad mgd \cos q_1 \quad mgd \cos q_2 + 2d$$

1. Schritt: Rechnerische Lsg.  $q_1^0 = q_2^0 = 0$

$$T = \frac{1}{2} m (\dot{q}_1)^2 + \frac{1}{2} L (\dot{q}_2)^2 = \frac{1}{2} m^2 (\dot{q}_1, \dot{q}_1) \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} \dot{q}_1 \\ \dot{q}_2 \end{pmatrix}$$

$$\underline{I} = m \underline{e}^2 \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$

$$\omega^2 \underline{I} \underline{\alpha} = \underline{V} \underline{\alpha}$$

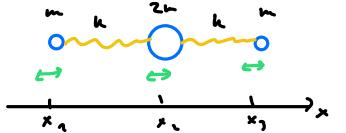
$$\textcircled{1} \quad \underline{\alpha}_1 = \frac{1}{\mu_1} \begin{pmatrix} 1 \\ -1 \end{pmatrix} \quad \omega^2 m \underline{e}^2 = m \underline{e} \cos \underline{\alpha}_1^\circ = m \underline{e} \quad \omega_1^2 = \frac{g}{L}$$

$$\textcircled{2} \quad \underline{\alpha}_2 = \frac{1}{\mu_2} \begin{pmatrix} 1 \\ 1 \end{pmatrix} \quad \omega^2 m \underline{e}^2 = m \underline{e} + 2 \underline{e} \quad \omega_2^2 = \frac{g}{L} + \frac{2 \underline{e}^2}{m \underline{e}^2}$$

$$\textcircled{2} \quad \underline{\alpha}_1^\circ = \underline{\alpha}_2^\circ + \pi (+ 2k\pi)$$



(38)



Obratne vektore u 1D

V ravnini je radijus L

$$\omega_{n1} \underline{\alpha}_1 = ? \quad H_1 = ?$$

$$x_1 = 0 + \underline{\alpha}_1 \quad \dot{x}_1 = \dot{\underline{\alpha}}_1$$

$$x_2 = L + \underline{\alpha}_2 \quad \dot{x}_2 = \dot{\underline{\alpha}}_2$$

$$x_3 = 2L + \underline{\alpha}_3 \quad \dot{x}_3 = \dot{\underline{\alpha}}_3$$

$$\text{Klik.} \quad T = \frac{1}{2} m \dot{x}_1^2 + \frac{1}{2} k L \dot{x}_2^2 + \frac{1}{2} m \dot{x}_3^2 = \frac{k}{2} (\dot{\underline{\alpha}}_1^2 + 2 \dot{\underline{\alpha}}_2^2 + \dot{\underline{\alpha}}_3^2) = \frac{k}{2} \underline{\dot{\alpha}}^T T \underline{\dot{\alpha}} \quad T = k \begin{bmatrix} 1 & 1 & 1 \\ 1 & 2 & 1 \\ 1 & 1 & 1 \end{bmatrix} \underline{\alpha}$$

$$\text{Pot.} \quad V = \frac{1}{2} k ((x_2 - x_1) - L)^2 + \frac{1}{2} k ((x_3 - x_2) - L)^2 = \frac{k}{2} ((\underline{\alpha}_2 - \underline{\alpha}_1)^2 + (\underline{\alpha}_3 - \underline{\alpha}_2)^2) =$$

$$= \frac{k}{2} (\underline{\alpha}_2^2 + \underline{\alpha}_3^2 - 2 \underline{\alpha}_1 \underline{\alpha}_2 + \underline{\alpha}_1^2 + \underline{\alpha}_2^2 - 2 \underline{\alpha}_2 \underline{\alpha}_3)$$

$$= \frac{k}{2} \underline{\dot{\alpha}}^T \underline{V} \underline{\dot{\alpha}} \quad \underline{V} = k \begin{bmatrix} 1 & -1 & 0 \\ -1 & 2 & -1 \\ 0 & -1 & 1 \end{bmatrix}$$

↑  
sistem je  
symmetric

$$\text{Poznato je: prosim} \quad \underline{\alpha}(t) = 2 \underline{\alpha} e^{-i \omega t}$$

$$(-\omega^2 \underline{I} + \underline{V}) \underline{\alpha} = 0$$

$$\frac{1}{\omega_0^2} = \lambda$$

$$0 = \det (\underline{V} - \omega^2 \underline{I}) = \det (k(\tilde{V} - \omega^2 \frac{1}{m} \tilde{T}))$$

$$0 = \det (\tilde{V} - \lambda \tilde{T})$$

$$\det \left( \begin{bmatrix} 1 & -1 & 0 \\ -1 & 2 & -1 \\ 0 & -1 & 1 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 & 0 \\ 0 & 2 & 0 \\ 0 & 0 & 1 \end{bmatrix} \right) = \det \begin{bmatrix} 1-\lambda & -1 & 0 \\ -1 & 2-2\lambda & -1 \\ 0 & -1 & 1-\lambda \end{bmatrix} =$$

$$= (1-\lambda) \begin{vmatrix} 2-2\lambda & -1 \\ -1 & 1-\lambda \end{vmatrix} - (-1) \begin{vmatrix} -1 & -1 \\ 0 & 1-\lambda \end{vmatrix} =$$

$$= (1-\lambda) ((2-2\lambda)(1-\lambda) - 1) - (1-\lambda) = (1-\lambda)(2-\lambda) 2\lambda$$

$$\lambda_1 = 1 \quad \lambda_2 = 2 \quad \lambda_3 = 0$$

$$\omega_1^2 = \omega_0^2 \quad \omega_2^2 = 2\omega_0^2 \quad \omega_3^2 = 0$$

$$\lambda_1 = 1 \quad (\tilde{Y} - \lambda \tilde{I}) \underline{z} = 0$$

$$\begin{pmatrix} 1-1 & -1 & 0 \\ -1 & 2-2 & -1 \\ 0 & -1 & 1-1 \end{pmatrix} \begin{pmatrix} a_{11} \\ a_{12} \\ a_{13} \end{pmatrix} = 0$$

$$\begin{pmatrix} 0 & -1 & 0 \\ -1 & 0 & -1 \\ 0 & -1 & 0 \end{pmatrix} \begin{pmatrix} a_{11} \\ a_{12} \\ a_{13} \end{pmatrix} = 0 \quad a_{12} = 0 \quad a_{13} = -a_{11}$$

$$\underline{a}_1 = \begin{pmatrix} 1 \\ 0 \\ -1 \end{pmatrix} \quad \omega_1^2 = \omega_0^2$$

$$\lambda_2 = 2 \quad \begin{pmatrix} 1-1 & -1 & 0 \\ -1 & 2-2 & -1 \\ 0 & -1 & 1-1 \end{pmatrix} \begin{pmatrix} a_{11} \\ a_{12} \\ a_{13} \end{pmatrix} = 0 \quad a_{21} = -a_{12} \quad a_{23} = -a_{13}$$

$$\underline{a}_2 = \begin{pmatrix} 1 \\ -1 \\ 1 \end{pmatrix} \quad \omega_2^2 = 2\omega_0^2$$

$$\omega_2^2 = 2\omega_0^2$$

$$\lambda_3 = 0 \quad \begin{pmatrix} 1-1 & -1 & 0 \\ -1 & 2-2 & -1 \\ 0 & -1 & 1-1 \end{pmatrix} \begin{pmatrix} a_{11} \\ a_{12} \\ a_{13} \end{pmatrix} = 0 \quad \underline{a}_3 = \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix} \quad \text{Transl.} \quad \text{Diagram}$$

Letzte Lösungsmethode

$$y_1 = \underline{a}_1 \left( \tilde{A}_1 e^{i\omega_0 t} + \tilde{B}_1 e^{-i\omega_0 t} \right) = \underline{a}_1 (A_1 \cos(\omega_0 t) + B_1 \sin(\omega_0 t))$$

$$y_2 = \underline{a}_2 \left( \tilde{A}_2 e^{i\sqrt{2}\omega_0 t} + \tilde{B}_2 e^{-i\sqrt{2}\omega_0 t} \right) = \underline{a}_2 (A_2 \cos(\sqrt{2}\omega_0 t) + B_2 \sin(\sqrt{2}\omega_0 t))$$

$$y_3 = \underline{a}_3 (A_3 + B_3 t)$$

$$\text{Splittung resultiert} \quad y(t) = \sum_{i=1}^3 y_i(t)$$

Zeitliche Position

$$\underline{y}(t) = \begin{pmatrix} \vdots \\ \vdots \\ \vdots \end{pmatrix} \quad \dot{\underline{y}}(t) = \begin{pmatrix} \ddot{v}_0 \\ \vdots \\ \vdots \end{pmatrix}$$

$$\underline{y}(t) = \begin{pmatrix} \underline{a}_1 & \underline{a}_2 & \underline{a}_3 \end{pmatrix} \begin{pmatrix} A_1 \cos(\omega_0 t) + \dots \\ A_2 \cos(\sqrt{2}\omega_0 t) + \dots \\ A_3 + B_3 t \end{pmatrix} \stackrel{t \rightarrow 0}{=} (\underline{a}_1 \ \underline{a}_2 \ \underline{a}_3) \begin{pmatrix} A_1 \\ A_2 \\ A_3 \end{pmatrix} =$$

$$= \begin{pmatrix} -1 & -1 & 1 \\ 0 & -1 & 1 \\ 1 & 1 & 1 \end{pmatrix} \begin{pmatrix} A_1 \\ A_2 \\ A_3 \end{pmatrix} = \begin{pmatrix} \vdots \\ \vdots \\ \vdots \end{pmatrix} \quad \Rightarrow \quad A_1 = A_2 = A_3 = 0$$

$$\dot{\underline{y}}(t) = \begin{pmatrix} \ddot{v}_0 \\ \vdots \\ \vdots \end{pmatrix} = (\underline{a}_1, \underline{a}_2, \underline{a}_3) \begin{pmatrix} B_1 \omega_0 \\ \vdots \\ \vdots \end{pmatrix} = \begin{pmatrix} -B_1 \omega_0 & +B_1 \sqrt{2} \omega_0 & +B_3 \omega_0 \\ 0 & -B_2 \sqrt{2} \omega_0 & +B_3 \omega_0 \\ B_1 \omega_0 & +B_2 \sqrt{2} \omega_0 & +B_3 \omega_0 \end{pmatrix}$$

$$\Rightarrow \quad B_1 = -\frac{1}{2} \frac{v_0}{\omega_0} \quad B_2 = \frac{1}{4\sqrt{2}} \frac{v_0}{\omega_0} \quad B_3 = \frac{1}{4} v_0$$

Splittung resultiert

$$y(t) = \underbrace{\frac{1}{2} \frac{v_0}{\omega_0} \sin(\omega_0 t)}_{d_1(t)} \begin{pmatrix} \vdots \\ \vdots \\ \vdots \end{pmatrix} + \underbrace{\frac{1}{4\sqrt{2}} \frac{v_0}{\omega_0} \sin(\sqrt{2}\omega_0 t)}_{d_2(t)} \begin{pmatrix} \vdots \\ \vdots \\ \vdots \end{pmatrix} + \underbrace{\frac{v_0}{4} t}_{d_3(t)} \begin{pmatrix} \vdots \\ \vdots \\ \vdots \end{pmatrix}$$

höchste Frequenz ist  $\frac{\omega_0}{4}$

Energiesatz: jüL präspezifische potentielle Energie:  $\underline{u}$

$$H_u = \frac{1}{2} d_u^T T_u^T + \frac{1}{2} d_u^T V_u^T \quad T^T = \underline{A}^T \underline{T} \underline{A} \quad \underline{V}^T = \underline{A}^T \underline{V} \underline{A} \quad \underline{A} = (\underline{a}_1 \ \underline{a}_2 \ \underline{a}_3)$$

$$\tau^1 = \begin{pmatrix} 2 & 0 & 0 \\ 0 & 2 & 0 \\ 0 & 0 & 4 \end{pmatrix} m \quad v^1 = \begin{pmatrix} 2 & 0 & 0 \\ 0 & 2 & 0 \\ 0 & 0 & 0 \end{pmatrix} u$$

$$H_1 = \frac{1}{2} \lambda^2 2m + \frac{1}{2} \lambda^2 2k = \left( \frac{1}{2} m \cos \omega_0 t \right)^2 m + \left( \frac{1}{2} m \sin \omega_0 t \right)^2 k = \frac{1}{4} m \omega_0^2$$

$$H_2 = \frac{1}{2} \lambda^2 4m + \frac{1}{2} \lambda^2 8k = \dots = \frac{1}{8} m \omega_0^2$$

$$H_3 = \frac{1}{2} \lambda^2 4m + 0 = \dots = \frac{1}{8} m \omega_0^2$$

$$H = H_1 + H_2 + H_3 = \frac{1}{4} m \omega_0^2$$



$$\underline{\underline{\underline{V}}}_{\underline{\underline{\underline{a}}}} = \omega \underline{\underline{\underline{T}}} \underline{\underline{\underline{a}}}$$

$$u \begin{pmatrix} 1 & -1 & 0 \\ -1 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \underline{\underline{\underline{a}}} = \omega^2 m \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \underline{\underline{\underline{a}}}$$

Let's see which of these is possible simultaneously

$$\underline{\underline{\underline{a}}}_i^T \underline{\underline{\underline{T}}} \underline{\underline{\underline{a}}}_j = \delta_{ij}, \quad \underline{\underline{\underline{a}}}_i^T \underline{\underline{\underline{T}}} \underline{\underline{\underline{a}}}_i$$

1. translation:  $\underline{\underline{\underline{a}}}_1 = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} \quad \underline{\underline{\underline{V}}}_{\underline{\underline{\underline{a}}}_1} = 0 \Rightarrow \omega^2 = 0$

2. circular motion:  $\underline{\underline{\underline{a}}}_2 = \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} \quad \underline{\underline{\underline{a}}}_2^T \underline{\underline{\underline{T}}} \underline{\underline{\underline{a}}}_2 = 0 \quad \underline{\underline{\underline{V}}}_{\underline{\underline{\underline{a}}}_2} = k \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} = \omega^2 \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} \Rightarrow \omega^2 = \frac{k}{m}$

3. octagonal motion:  $\underline{\underline{\underline{a}}}_3 = \begin{pmatrix} 1 \\ 1 \\ 0 \end{pmatrix} \quad \underline{\underline{\underline{a}}}_3^T \underline{\underline{\underline{T}}} \underline{\underline{\underline{a}}}_3 = \begin{pmatrix} 1 \\ 1 \\ 0 \end{pmatrix}^T \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} = a_m - k c \Rightarrow a_m = c$

$$\underline{\underline{\underline{a}}}_4 = \begin{pmatrix} 1 \\ 0 \\ 1 \end{pmatrix} \quad \underline{\underline{\underline{a}}}_4^T \underline{\underline{\underline{T}}} \underline{\underline{\underline{a}}}_4 = \begin{pmatrix} 1 \\ 0 \\ 1 \end{pmatrix}^T \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} = 2m + b m = 0 \quad b = -2 \frac{k}{m}$$

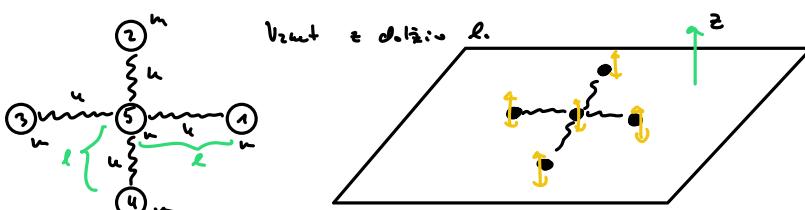
$$\Rightarrow \underline{\underline{\underline{a}}}_4 = \begin{pmatrix} 1 \\ 0 \\ -2 \frac{k}{m} \end{pmatrix}$$

$$\underline{\underline{\underline{V}}}_{\underline{\underline{\underline{a}}}_3} = k \begin{pmatrix} 1+2 \frac{k}{m} \\ -2-4 \frac{k}{m} \\ 1+\frac{k}{m} \end{pmatrix} = \omega^2 \underline{\underline{\underline{T}}} \begin{pmatrix} 1+2 \frac{k}{m} \\ -2-4 \frac{k}{m} \\ 1+\frac{k}{m} \end{pmatrix} = \begin{pmatrix} 1 \\ 1 \\ 0 \end{pmatrix} \omega^2 = \begin{pmatrix} 1 \\ 1 \\ 0 \end{pmatrix} \omega^2 m$$

$$k \left( 1 + \frac{2k}{m} \right) \begin{pmatrix} 1 \\ 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 1 \\ 1 \\ 0 \end{pmatrix} \omega^2 m$$

$$\omega^2 = \frac{k}{m} \left( 1 + \frac{2k}{m} \right)$$

40



Viert = doppelt d.

$$\underline{\underline{\underline{z}}} = \begin{pmatrix} z_1 \\ z_2 \\ \vdots \\ z_N \end{pmatrix}$$

$$T = \frac{1}{2} m \sum_{i=1}^N \dot{z}_i^2 = \frac{1}{2} \dot{\underline{\underline{\underline{z}}}}^T \underline{\underline{\underline{T}}} \dot{\underline{\underline{\underline{z}}}} = \frac{1}{2} \dot{\underline{\underline{\underline{z}}}}^T m \underline{\underline{\underline{T}}} \dot{\underline{\underline{\underline{z}}}}$$

$$\dot{z}_1 = (x_1, 0, z_1)$$

$$\dot{z}_2 = (0, x_2, z_2)$$

$$\dot{z}_3 = (-x_3, 0, z_3)$$

$$\dot{z}_4 = (0, -x_4, z_4)$$

$$\dot{z}_5 = (0, 0, z_5)$$

Loc. d.: to power do so  
vertical projection

$$V = \frac{1}{2} k \sum_{i=1}^4 (r_i - r_s - l_c)^2$$

$$|r_i - r_s| = \sqrt{l^2 + (z_i - z_s)^2} = \sqrt{1 + \left(\frac{z_i - z_s}{l}\right)^2} \approx 1 \left(1 + \frac{1}{2} \left(\frac{z_i - z_s}{l}\right)^2\right) \quad \text{mögliche Näherung}$$

$$V = \frac{1}{2} k \sum_{i=1}^4 ((l - l_0) + \frac{1}{2} (z_i - z_s))^2 = \frac{1}{2} k \left( (l - l_0)^2 + 2(l - l_0) \frac{1}{2} (z_i - z_s) + \frac{1}{4} (z_i - z_s)^2 \right)$$

$$\approx \frac{1}{2} k \sum_{i=1}^4 \underbrace{\frac{l - l_0}{l}}_{>0} (z_i - z_s)^2 + \text{konst.} = \frac{1}{2} k \sum_{i=1}^4 (z_i - z_s)^2 + \text{konst.}$$

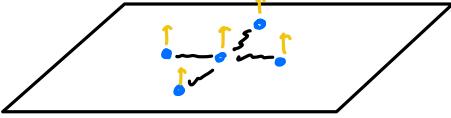
$$= \frac{1}{2} k (z_1^2 + z_2^2 + z_3^2 + z_4^2 - 2z_1 z_2 - 2z_1 z_3 - 2z_1 z_4 - 2z_2 z_3 - 2z_2 z_4 - 2z_3 z_4) =$$

$$= \frac{1}{2} \underline{z}^T \underline{V} \underline{z}$$

$$\underline{V} = \tilde{k} \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 4 \end{bmatrix}$$

$$\text{E-L Gleichung: } (\underline{V} - \omega^2 \underline{I}) \underline{s} = 0 \quad \underline{I} = n \underline{J}$$

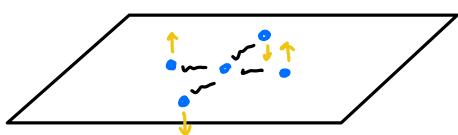
$$(\underline{V} - \omega^2 n \underline{I}) \underline{s} = 0$$

Rechteckig: 

$$\underline{a}_1 = \begin{pmatrix} 1 \\ 1 \\ 1 \\ 1 \end{pmatrix}$$

$$\underline{V} \underline{s}_1 = \tilde{k} \begin{pmatrix} 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \end{pmatrix} \begin{pmatrix} 1 \\ 1 \\ 1 \\ 1 \end{pmatrix} = \begin{pmatrix} 4 \\ 4 \\ 4 \\ 4 \end{pmatrix} = n \omega^2 \underline{a}_1$$

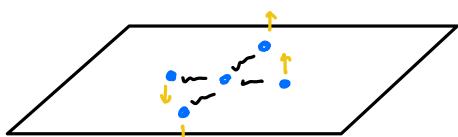
$$\Rightarrow \omega_1^2 = \frac{\tilde{k}}{n} = \omega_0^2$$



$$\underline{a}_2 = \begin{pmatrix} 1 \\ -1 \\ 1 \\ -1 \end{pmatrix}$$

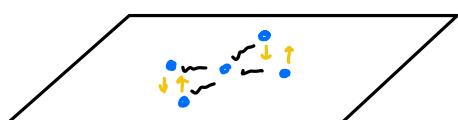
$$\underline{V} \underline{s}_2 = \tilde{k} \begin{pmatrix} 1 & 1 & 0 & 0 \\ 1 & 1 & 0 & 0 \\ 0 & 0 & 1 & 1 \\ 0 & 0 & 1 & 1 \end{pmatrix} \begin{pmatrix} 1 \\ -1 \\ 1 \\ -1 \end{pmatrix} = \tilde{k} \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \end{pmatrix} = n \omega^2 \underline{a}_2$$

$$\Rightarrow \omega_2^2 = \frac{\tilde{k}}{n} = \omega_0^2$$



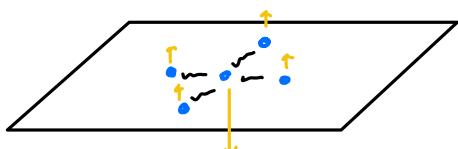
$$\underline{a}_3 = \begin{pmatrix} 1 \\ -1 \\ -1 \\ 1 \end{pmatrix}$$

$$\underline{V} \underline{s}_3 = \dots = \omega_1^2 n \underline{a}_1 \Rightarrow \omega_3^2 = \omega_0^2$$



$$\underline{a}_4 = \begin{pmatrix} 1 \\ -1 \\ 1 \\ 1 \end{pmatrix}$$

$$\dots \Rightarrow \omega_4^2 = \omega_0^2$$



$$\underline{a}_5 = \begin{pmatrix} d_1 \\ d_2 \\ d_3 \\ d_4 \\ d_5 \end{pmatrix}$$

$$\underline{a}_1 \cdot \underline{a}_5 = 0$$

$$\vdots$$

$$\underline{a}_4 \cdot \underline{a}_5 = 0$$

$$\} \Rightarrow d_1 = d_2 = d_3 = d_4 = 1 \quad \underline{a}_5 = \begin{pmatrix} 1 \\ 1 \\ 1 \\ 1 \\ d_5 \end{pmatrix}$$

$$\underline{a}_1 \cdot \underline{a}_5 = 0 = 1 + 1 + 1 + 1 + d_5 = 0 \quad d_5 = -4$$

$$\underline{a}_5 = \begin{pmatrix} 1 \\ 1 \\ 1 \\ 1 \\ -4 \end{pmatrix}$$

$$\underline{\underline{V}}_{SS} = \dots = \omega_3^2 \sim SS \quad \Rightarrow \quad \omega_3^2 = 5\omega_0^2$$

### (T) Hamiltonsches formulieren

$$L = L(g_\alpha, \dot{g}_\alpha) \quad p_\alpha = \frac{\partial L}{\partial \dot{g}_\alpha}$$

$$H = \sum_\alpha p_\alpha \dot{g}_\alpha - L(g_\alpha, \dot{g}_\alpha) = H(g_\alpha, p_\alpha)$$

Präzise: zylindrische Koordinaten

$$L = \frac{1}{2} m \dot{r}^2 + \frac{1}{2} m r^2 \dot{\theta}^2 - V(r)$$

$$p_r = m \dot{r} \quad p_\theta = m r^2 \dot{\theta}$$

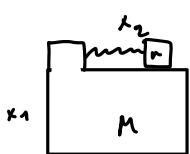
$$H = m \dot{r}^2 + m r^2 \dot{\theta}^2 - \frac{1}{2} m \dot{r}^2 - \frac{1}{2} m r^2 \dot{\theta}^2 + V(r) = \frac{1}{2} m \dot{r}^2 + \frac{1}{2} m r^2 \dot{\theta}^2 + V(r) = T + V$$

$$\text{da } V \neq V(\dot{g}) \rightarrow H = T + V$$

$$H = \frac{p_r^2}{2m} + \frac{p_\theta^2}{2m} + V(r)$$

$$\boxed{\dot{g}_\alpha = \frac{\partial H}{\partial p_\alpha} \quad \dot{p}_\alpha = -\frac{\partial H}{\partial g_\alpha}} \quad \text{Hamilt. eqn.}$$

(41)



$$L = \frac{1}{2} M \dot{x}_1^2 + \frac{1}{2} m (\dot{x}_1 + \dot{x}_2)^2 - \frac{k x_2^2}{2}$$

$$p_1 = \frac{\partial L}{\partial \dot{x}_1} = M \dot{x}_1 + m (\dot{x}_1 + \dot{x}_2)$$

$$p_2 = \frac{\partial L}{\partial \dot{x}_2} = m (\dot{x}_1 + \dot{x}_2)$$

$$H = T + V = \frac{1}{2} M \dot{x}_1^2 + \frac{1}{2} m (\dot{x}_1 + \dot{x}_2)^2 + \frac{k x_2^2}{2} = \frac{1}{2} \frac{(p_1 - p_2)^2}{M} + \frac{p_2^2}{2m} + \frac{k x_2^2}{2}$$

$$\dot{x}_1 = \frac{\partial H}{\partial p_1} = \frac{p_1 - p_2}{M}$$

$$\dot{x}_2 = \frac{\partial H}{\partial p_2} = -\frac{(p_1 - p_2)}{m} + \frac{p_2}{m}$$

$$\dot{p}_1 = -\frac{\partial H}{\partial x_1} = 0 \quad \dot{p}_2 = -\frac{\partial H}{\partial x_2} = -k x_2$$

$$\ddot{x}_2 = -\frac{\dot{p}_2}{m} + \dot{p}_2 \left( \frac{1}{m} + \frac{1}{m} \right) = -k x_2 \left( \frac{1}{m} + \frac{1}{m} \right) \quad \text{nach}$$

(42)

Dreh v. homogenem magnetum Polen

$$L = \frac{1}{2} m \dot{r}^2 - e \dot{\varphi} + e \vec{v} \cdot \vec{A}$$

$$p_\varphi = \frac{\partial L}{\partial \dot{\varphi}} = m \dot{\varphi} + e A_\varphi$$

$$\frac{dp_\varphi}{dt} = m \ddot{\varphi} + e \dot{A}_\varphi = \frac{\partial L}{\partial r} = -e \frac{\partial \vec{A}}{\partial r}$$

$$m\ddot{v} = e \left( -\frac{\partial \Phi}{\partial x} - \frac{\partial A_x}{\partial t} \right)$$

$$m\dot{\tilde{v}} = e \left( -\nabla \Phi - \frac{\partial \tilde{A}}{\partial t} \right) = e \tilde{E}$$

$$H = \sum_i p_i \dot{x}_i - L = \dot{p} \cdot \vec{v} - L = \frac{1}{2} m v^2 + e \vec{A} \cdot \vec{v} - \frac{1}{2} m v^2 + e \Phi - e \vec{A} \cdot \vec{v} = \frac{1}{2} m v^2 + e \Phi \neq T + U$$

$$Q=0 \quad \vec{B} = B \hat{k} \quad H = \frac{(\vec{p} - e\vec{A})^2}{2m} \quad \vec{A} = B(-y\hat{i})$$

$$\vec{A} = \frac{1}{2} B (-y\hat{i} + x\hat{j}) \quad \vec{B} = \nabla \times \vec{A} = (0, 0, B)$$

$$H = \frac{1}{2m} \left( p_x + \frac{eBx}{2} \right)^2 + \frac{1}{2L} \left( p_y - \frac{eBy}{2} \right)^2$$

$$\dot{x} = \frac{\partial H}{\partial p_x} = \frac{1}{m} \left( p_x + \frac{eBx}{2} \right) \quad \dot{p}_x = -\frac{\partial H}{\partial x} = \frac{1}{m} \left( p_x - \frac{eBy}{2} \right) \left( -\frac{eB}{2} \right) = \frac{eB}{2} \dot{y}$$

$$\dot{y} = \frac{\partial H}{\partial p_y} = \frac{1}{m} \left( p_y - \frac{eBy}{2} \right) \quad \dot{p}_y = -\frac{\partial H}{\partial y} = \frac{1}{m} \left( p_y + \frac{eBx}{2} \right) \frac{eB}{2} = \frac{eB}{2} \dot{x}$$

$$\ddot{x} = \frac{\dot{p}_x + \frac{eB}{2} \dot{y}}{m} = \frac{eB}{m} \dot{y} = \omega_c \dot{y} \quad \ddot{y} = -\frac{eB}{m} \dot{x} = -\omega_c \dot{x}$$

$$\Rightarrow v_x = A \cos(\omega_c t)$$

$$v_y = A \sin(\omega_c t)$$