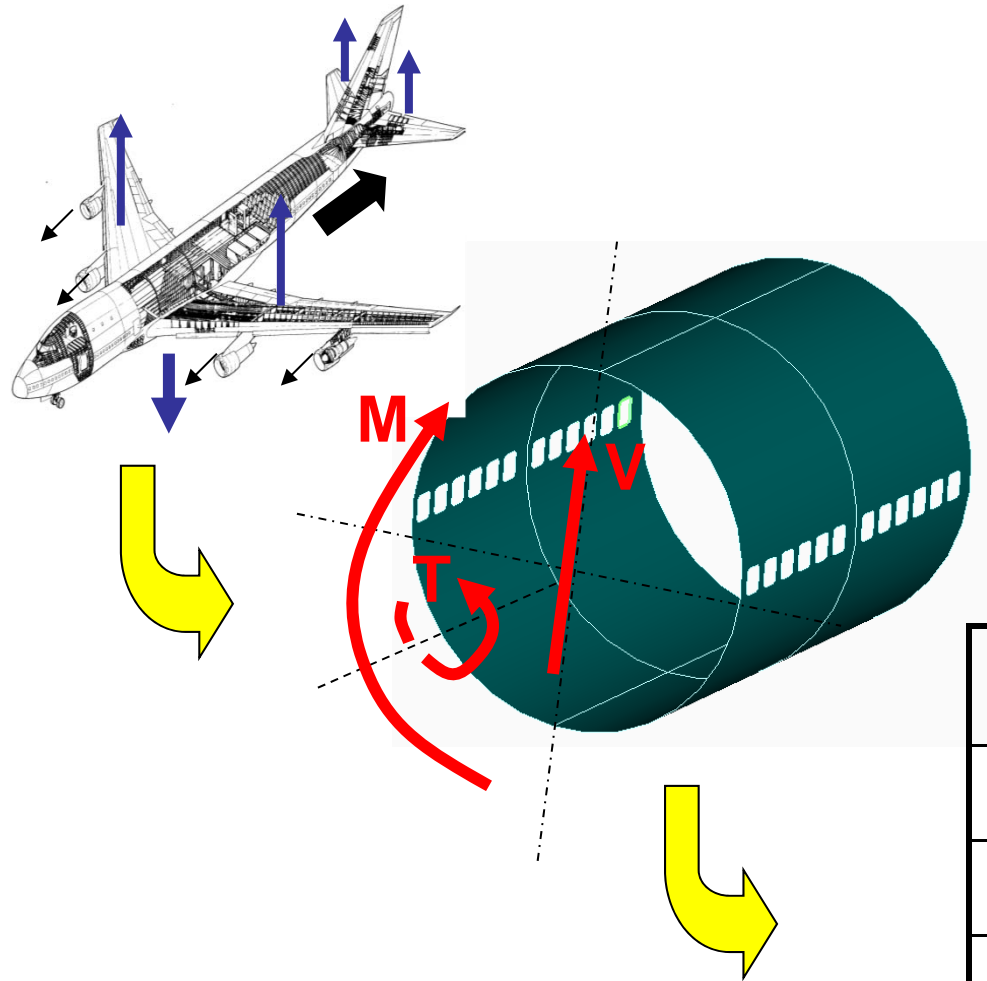
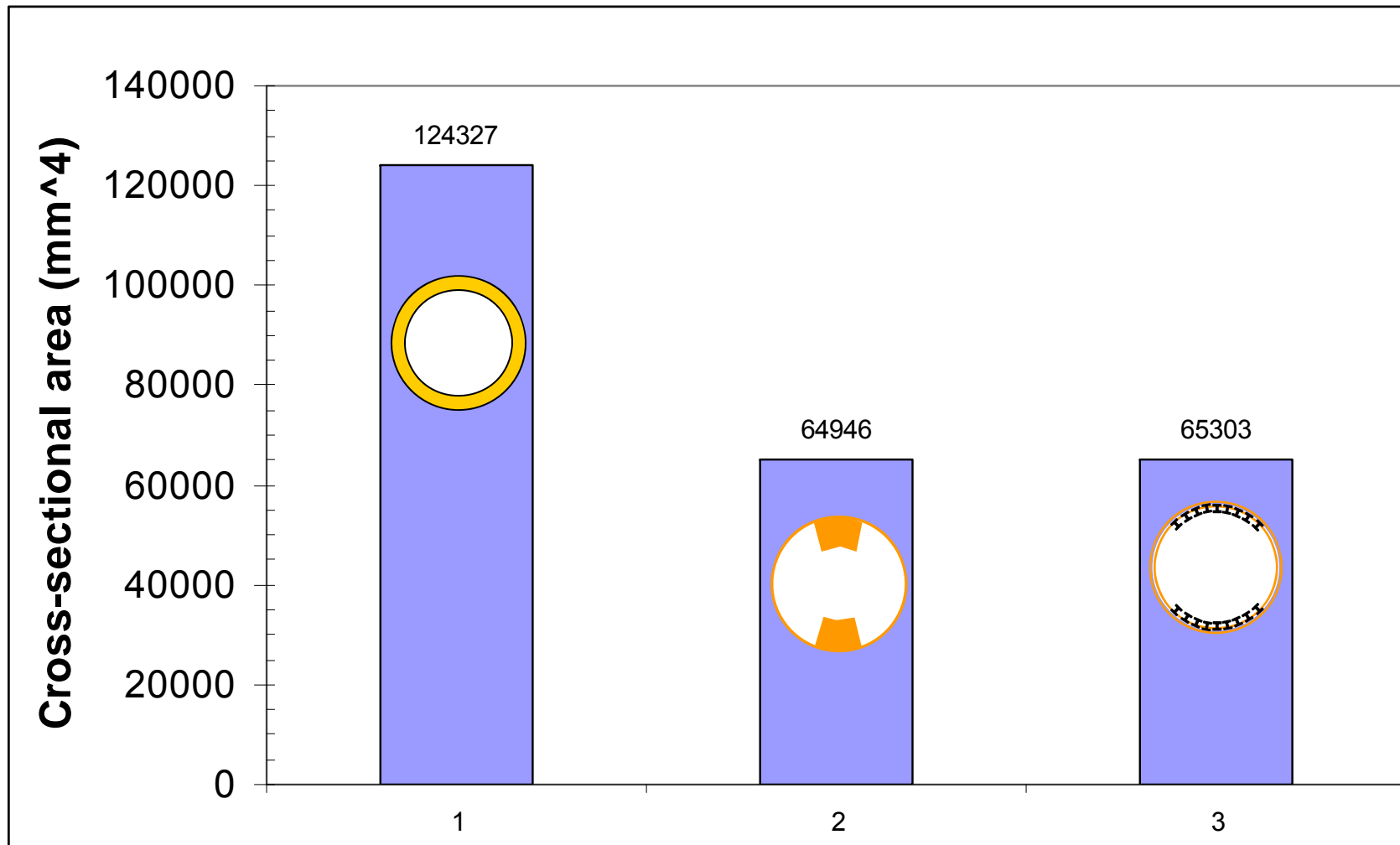


“Running” example – Fuselage cross-section



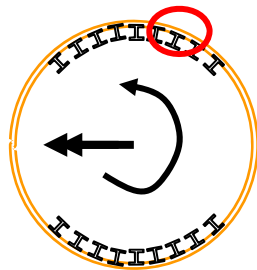
Property	Value
Diameter(m)	4.0
M (MNm)	60
V (kN)	660
T (kNm)	30

So far...

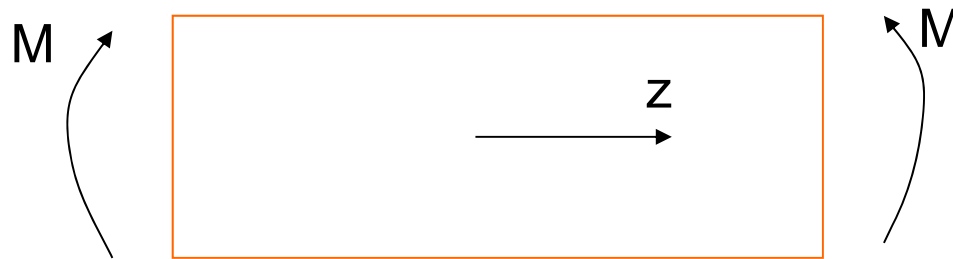


Design now for torsion

- from previous lecture:
 - pure torsion results in constant shear flow q in the wall of closed section beam
 - $T = 2Aq$ (3.44)
- what happens when we have combined bending moment M and torque T ? is q constant?



front view



edge view

- considering M alone, the normal stress σ_z is the same at any $z \Rightarrow$ the load on each stiffener is constant!

Design for torsion (cont'd)

- Now apply M and T at the same time, isolate the region around a single stiffener and place it in equilibrium:

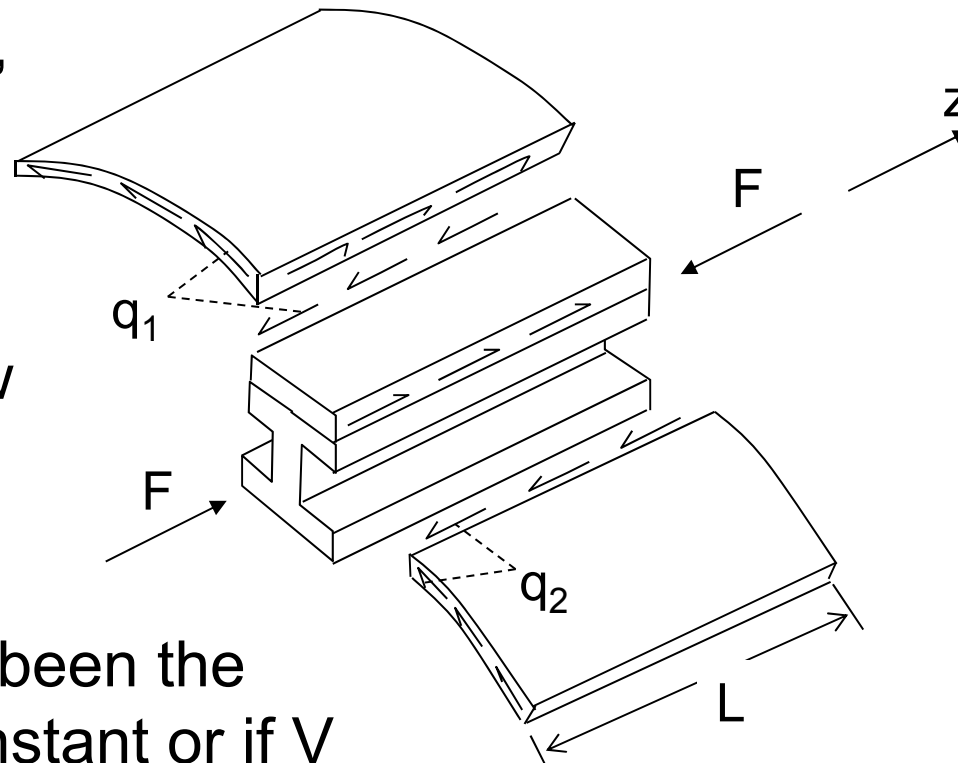
for single stiffener,

$$\sum F_z = 0 \Rightarrow$$

$$-q_1 L - F + q_2 L + F = 0 \Rightarrow$$

$$q_1 = q_2$$

and the shear flow
is still constant!



- This would not have been the case if M were not constant or if V were applied (the latter in a future lecture)

if M were not constant, F
would not be constant!

Design for torsion...

- so our basic relationship, eq. (3.44) $T = 2Aq$ (3.44) is still valid

- we can now relate the shear flow to the shear stress, through eq. (3.43)

$$q = t\tau \quad (3.43)$$

- and combining (3.43) and (3.44) we can solve for t:

$$t = \frac{T}{2A\tau}$$

- at failure, $\tau = \tau_y$ so

$$t = \frac{T}{2A\tau_y}$$

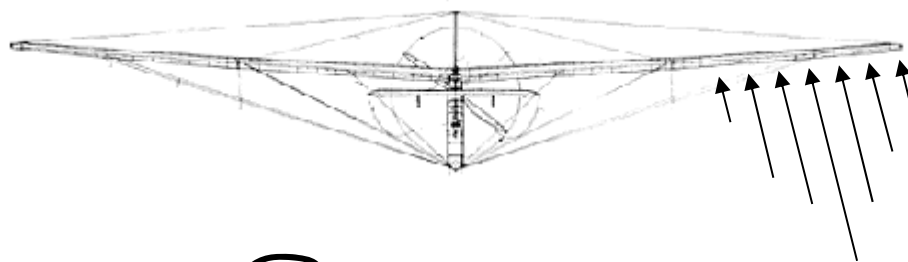
Design for torsion...

$$t = \frac{T}{2A\tau_y} = \frac{T}{2\pi R^2\tau_y}$$

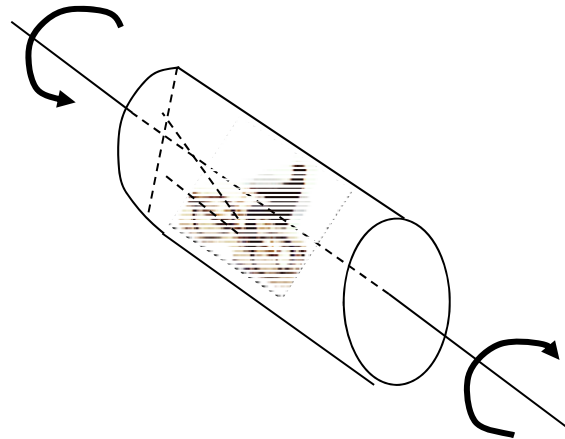
- Substituting values, $t=0.004$ mm !!
- This value is much smaller than the value currently used of 0.5 mm in our example so no change to our design

(Note: $\tau_y=275.1$ MPa for
7075-T6 Al)

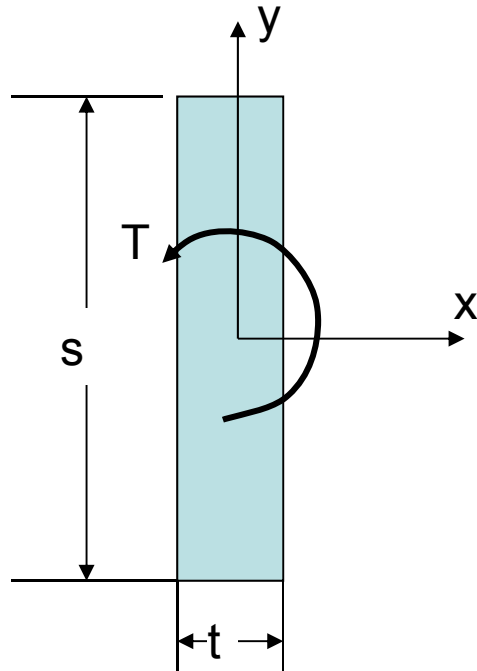
Torsion (cont'd)



e.g. wind gust acting on
one wing tip twists the
fuselage and puts it
under torsion



Torsion of narrow rectangular cross-section



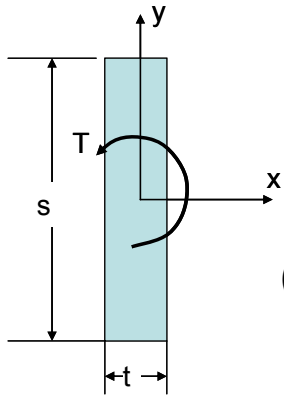
- assume $s \gg t$ (bar infinite in y dir)
- this means that if we fix x and move along y nothing changes
- if nothing changes with y , $\partial/\partial y = 0$
- then, the governing eq (3.18) can be combined with (3.30) to give

$$\begin{aligned}
 (3.18) \quad & \frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} = \nabla^2 \phi = F \\
 (3.30) \quad & -2G \frac{d\theta}{dz} = F
 \end{aligned}
 \left. \vphantom{\begin{aligned} (3.18) \\ (3.30) \end{aligned}} \right\} \frac{\partial^2 \phi}{\partial x^2} = -2G \frac{d\theta}{dz} \quad \text{or} \quad \frac{d^2 \phi}{dx^2} = -2G \frac{d\theta}{dz} \quad (4.1)$$

(since there is no dependence on y)

- integrating twice with respect to x : $\phi = -2G \frac{d\theta}{dz} \frac{x^2}{2} + Bx + C = -G \frac{d\theta}{dz} x^2 + Bx + C$

Torsion of narrow rectangular cross-section



$$\varphi = -G \frac{d\theta}{dz} x^2 + Bx + C$$

- to determine B and C, use the BC (3.22)

(3.22) $\varphi = 0$ on the boundary: $x = \pm t/2$
(how about $y = \pm s/2$??)

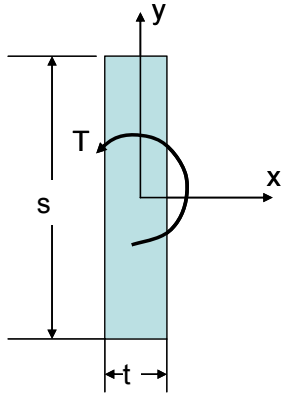
- applying the BC's:

$$\left. \begin{aligned} -G \frac{d\theta}{dz} \left(-\frac{t}{2}\right)^2 + B \left(-\frac{t}{2}\right) + C &= 0 \\ -G \frac{d\theta}{dz} \left(\frac{t}{2}\right)^2 + B \left(\frac{t}{2}\right) + C &= 0 \end{aligned} \right\} \Rightarrow \begin{aligned} B &= 0 \\ C &= G \frac{d\theta}{dz} \left(\frac{t}{2}\right)^2 \end{aligned}$$

- and, therefore,

$$\varphi = -G \frac{d\theta}{dz} \left(x^2 - \left(\frac{t}{2}\right)^2 \right) \quad (4.2)$$

Torsion of narrow rectangular cross-section: stress determination



$$\varphi = -G \frac{d\theta}{dz} \left(x^2 - \left(\frac{t}{2} \right)^2 \right) \quad (4.2)$$

- stresses can now be determined using eqs (3.4) and (3.5)

$$(3.4) \quad \tau_{xz} = \frac{\partial \varphi}{\partial y} \quad \left. \begin{array}{l} \tau_{xz} \approx 0 \quad (\text{except near the ends } y=\pm s/2 !!) \end{array} \right\} \quad (4.3)$$

$$(3.5) \quad \tau_{yz} = -\frac{\partial \varphi}{\partial x} \quad \left. \begin{array}{l} \tau_{yz} = 2Gx \frac{d\theta}{dz} \end{array} \right\} \quad (4.4)$$

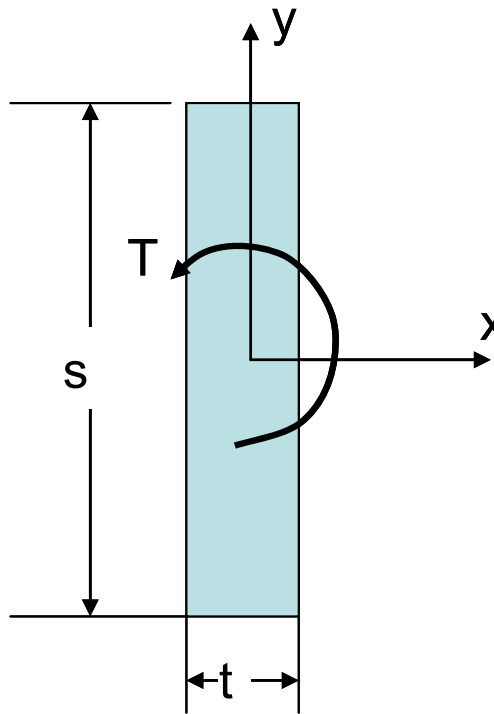
- now from (3.25), and using (4.2):

$$(3.25) \quad T = 2 \iint \varphi dx dy \xrightarrow{(4.2)} T = 2 \int_{-s/2}^{s/2} dy \int_{-t/2}^{t/2} -G \frac{d\theta}{dz} \left(x^2 - \left(\frac{t}{2} \right)^2 \right) dx = -2sG \frac{d\theta}{dz} \left[\frac{x^3}{3} - x \left(\frac{t}{2} \right)^2 \right]_{-t/2}^{t/2} = G \frac{d\theta}{dz} \frac{st^3}{3}$$

- and comparing to (3.25), $T = GJ \frac{d\theta}{dz}$ we get $J = \frac{st^3}{3}$ (4.5)

(valid for long rectangular cross-sections)

Torsion of narrow rectangular cross-section: stress determination



$$\tau_{xz} \approx 0 \quad (4.3)$$

$$\tau_{yz} = 2Gx \frac{d\theta}{dz} \quad (4.4)$$

- from (4.4), the max τ_{yz} is $\tau_{yz \max} = tG \frac{d\theta}{dz}$

- but from (3.25) $T = GJ \frac{d\theta}{dz}$

- combining: $\tau_{yz \max} = \frac{Tt}{J}$ $\tau_{yz \max} = \frac{3T}{st^2}$ (4.6)

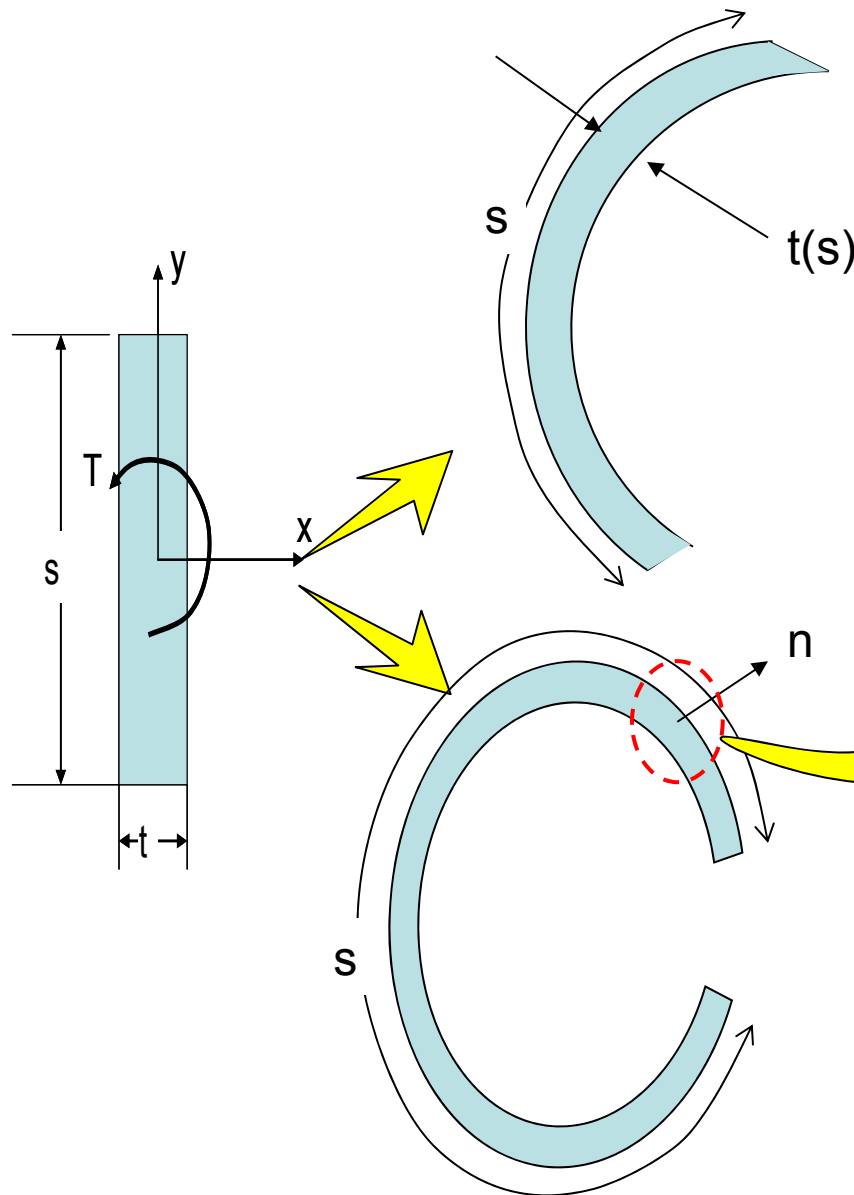
- and using (4.5), $J = \frac{st^3}{3}$ (4.7)

valid for long rectangular cross-sections

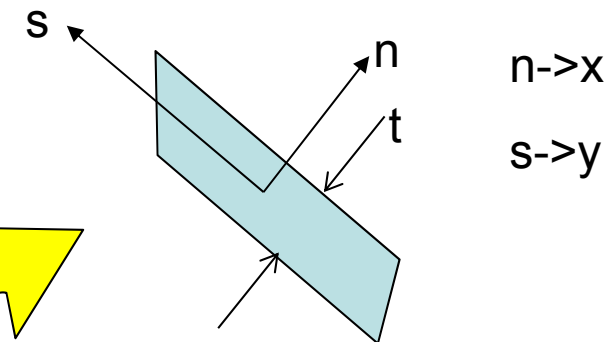
Notes: (1) s in eq (4.5) is always the longer dimension

(2) equations (4.5) and (4.6) are very accurate for $s > 10t$; for smaller s , equation (4.5) is modified but (4.6) is still valid

Torsion of open section beams



- if a long rectangular cross-section is bent to form an open cross-section of a beam (beam axis is perpendicular to the figure) then, locally, the shear stress state is very similar to that of a straight rectangular cross-section:



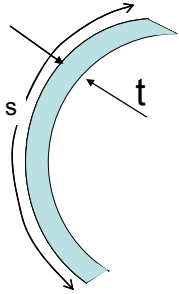
- by analogy to (4.3) and (4.4) we get:

$$\tau_{nz} \approx 0 \quad (4.8)$$

$$\tau_{sz} = 2Gn \frac{d\theta}{dz} \quad (4.9)$$

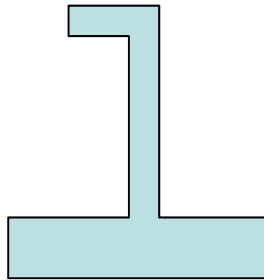
$$\tau_{sz \max} = tG \frac{d\theta}{dz} \quad (4.10) \quad = T/J$$

Torsion of open section beams



- Eq (4.7) for the torsion constant is still valid:

$$J = \frac{st^3}{3} \quad (4.7)$$



- For open sections consisting of multiple segments and/or having variable thickness, eq. (4.7) can be generalized to:

$$J = \frac{1}{3} \sum_{i=1}^N s_i t_i^3 \quad \text{if the thickness } t_i \text{ of each section is constant} \quad (4.11)$$

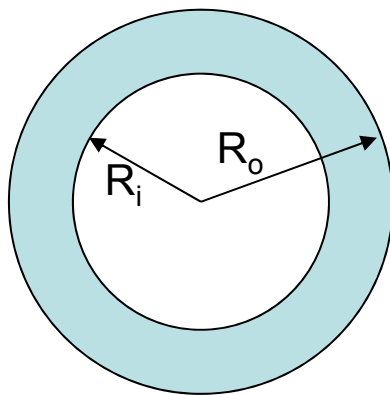
or

$$J = \frac{1}{3} \int_{\text{section}} t^3 ds \quad \text{if the thickness } t_i \text{ of each section is not constant} \quad (4.12)$$

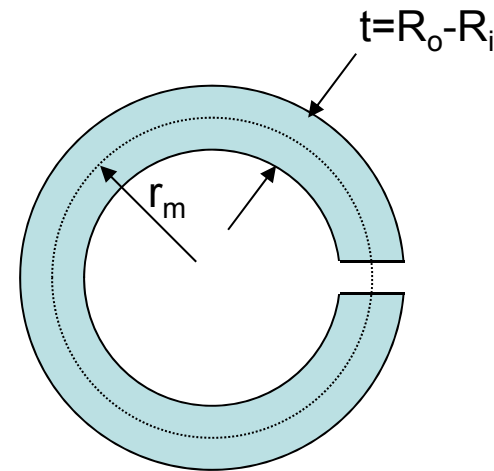
note that if a segment is not “long and narrow” the 1/3 term is corrected as a function of t/s: $J = st^3(1/2 + a(t/s) + b(t/s)^2 + \dots)$

Application: Open versus Closed section beams under torsion

- compare the strength and stiffness of a closed and open cylindrical cross-section beam



versus



$$r_m = (R_o + R_i)/2$$

Torsion of closed and open circular section beams

- to compare the stiffness, compare the amount of twist for a given torque
- to compare the strength, compare the amount of torque needed to create **the same maximum shear stress** in each cross-section

Stiffness comparison

- we know from eq. (3.31)

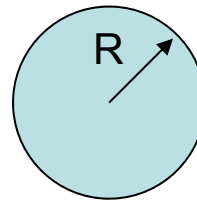
$$T = GJ \frac{d\theta}{dz} \Rightarrow \theta = \frac{T}{GJ} z + \cancel{C}^{\nearrow} \quad (\theta=0 \text{ at } z=0)$$

- so for a given applied torque T , since the shear modulus G is the same, it suffices to compare the J (torsion constant) values
- the configuration with higher J , will rotate less and is, therefore, stiffer

Stiffness comparison

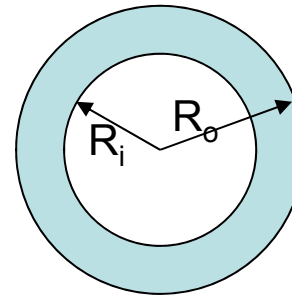
- for a closed circular cross-section, eq (3.35)

$$J = \frac{\pi R^4}{2}$$



- therefore, for a ring of outer radius R_o and inner radius R_i , $J=J_o-J_i$

$$J = \frac{\pi}{2} (R_o^4 - R_i^4)$$



Note: this expression is valid for constant thickness, single material.

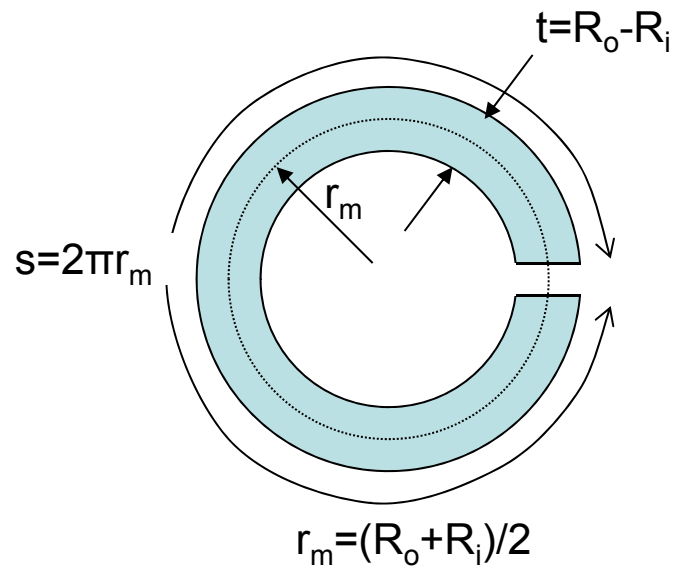
Otherwise, need to use: $\frac{d\theta}{dz} = \frac{1}{2A} \oint \frac{q}{tG} ds$ with $q=\text{const}$

$$T = 2Aq$$

$$\frac{d\theta}{dz} = \frac{T}{GJ}$$

Stiffness comparison

- for the open cross section, eq (4.7) $J = \frac{st^3}{3}$ becomes



$$J = \frac{1}{3} 2\pi r_m t^3$$

Stiffness comparison

- so, comparing the J values

$$\frac{J_{closed}}{J_{open}} = \frac{\frac{\pi}{2}(R_o^4 - R_i^4)}{\frac{1}{3}2\pi r_m t^3}$$

- but,

$$R_o = r_m + \frac{t}{2}$$

$$R_i = r_m - \frac{t}{2}$$

- and, therefore,

$$R_o^4 = r_m^4 + 4r_m^3 \frac{t}{2} + 6r_m^2 \left(\frac{t}{2}\right)^2 + 4r_m \left(\frac{t}{2}\right)^3 + \left(\frac{t}{2}\right)^4$$

$$R_i^4 = r_m^4 - 4r_m^3 \frac{t}{2} + 6r_m^2 \left(\frac{t}{2}\right)^2 - 4r_m \left(\frac{t}{2}\right)^3 + \left(\frac{t}{2}\right)^4$$

Stiffness comparison

- so that

$$R_o^4 - R_i^4 = 4r_m^3 t + 4r_m t^3 \quad \approx 0 \text{ for } t \ll r_m$$

- and finally,

$$\frac{J_{closed}}{J_{open}} = 3 \frac{r_m^2}{t^2}$$

typical values: $r_m = 100 \text{ mm}$, $t = 3 \text{ mm} \Rightarrow J_{closed}/J_{open} = 3333 !!!$

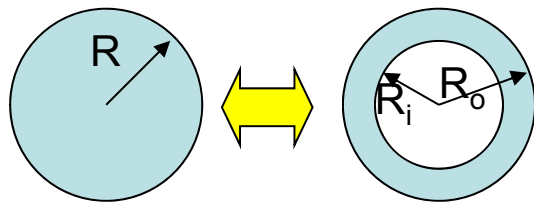
- the closed cross section is thousands of times stiffer and thus rotates thousands of times less under a given torque T (over the same length of beam)

Strength comparison

- the maximum shear stress for a closed circular cross-section is given by eq (3.40)

$$\tau = \frac{TR_o}{J}$$

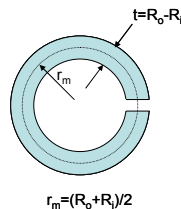
which is the same as for a hollow cylinder



(Note there is no difference in τ_{\max} between solid and hollow circles, i.e., the center of a solid cylinder is NOT buying us anything in carrying shear stresses due to torsion)

- the maximum shear stress for the open cross-section is obtained by combining (4.10) and (3.31)

$$\tau = \frac{Tt}{J}$$



Note: there is a typo in Megson p. 80 right after eq 3.29, a t is missing from the numerator!

Strength comparison

- for either case (closed or open beam), failure would occur when the shear stress τ equals the shear strength τ_{ult} of the material

$$\tau_{ult} = \frac{T_{failc} R_o}{J_{closed}} \Rightarrow T_{failc} = \frac{\tau_{ult} J_{closed}}{(r_m + t/2)}$$

$$\tau_{ult} = \frac{T_{failo} t}{J_{open}} \Rightarrow T_{failo} = \frac{\tau_{ult} J_{open}}{t}$$

- therefore, the ratio of failure torques is

$$\frac{T_{failc}}{T_{failo}} = \frac{J_{closed}}{J_{open}} \frac{t}{r_m + \frac{t}{2}}$$

Strength comparison

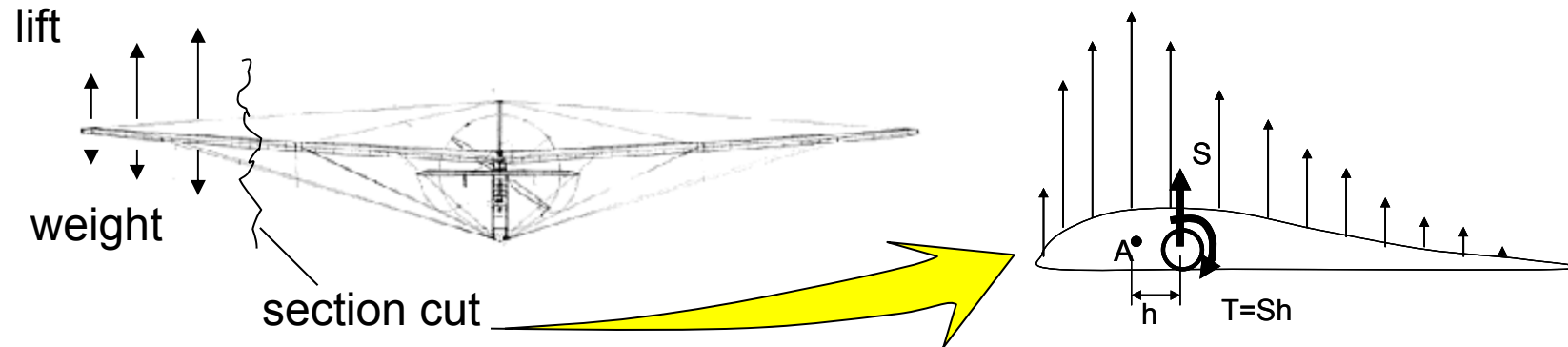
- but $J_{\text{closed}}/J_{\text{open}}$ was calculated earlier; substituting,

$$\frac{T_{\text{failc}}}{T_{\text{failo}}} = 3 \frac{r_m^2}{t^2} \frac{t}{r_m + \frac{t}{2}}$$

for the same typical values as before ($r_m=100$ mm, $t=3$ mm), $T_{\text{failc}}/T_{\text{failo}}=98.5$!!!

- the closed cross-section can take 100 times higher torque before it fails!

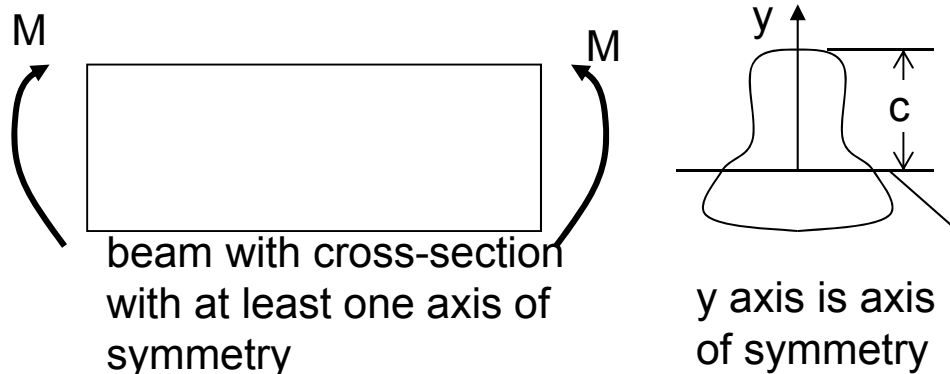
Bending revisited



- structures under bending loads, such as a wing or fuselage, must be designed so that
 - the deflections are minimized
 - the resulting stresses do not cause failure
- the most efficient structural element to minimize deflections and stresses is a beam (with high moment of inertia I)

Closed versus open cross-sections in bending

- symmetrical bending about one axis (from



beam with cross-section with at least one axis of symmetry

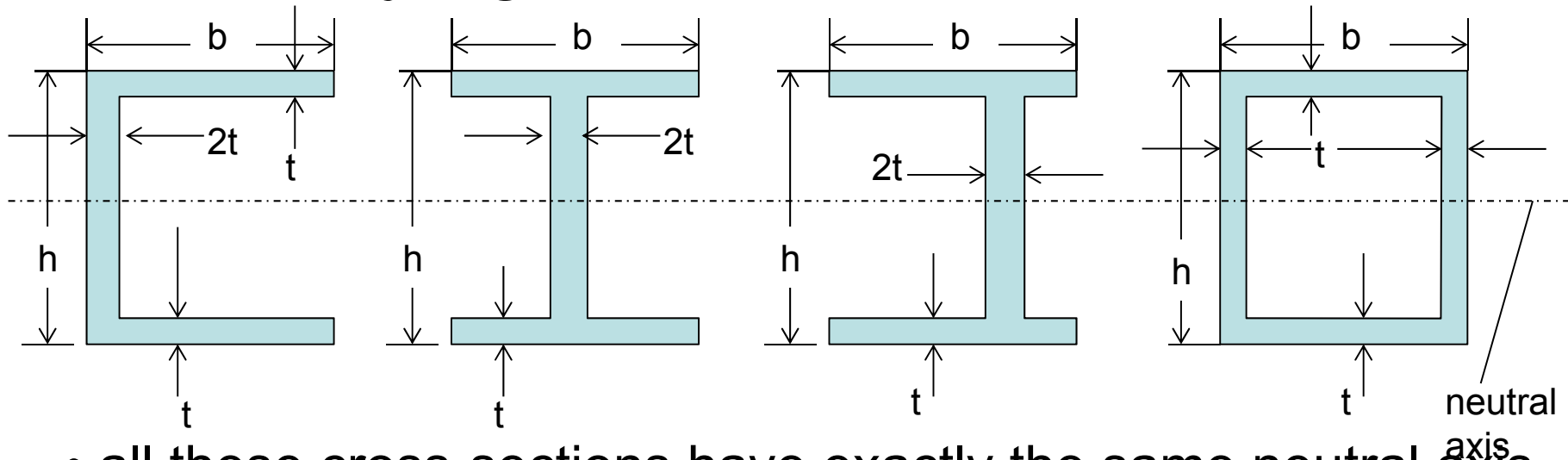
y axis is axis of symmetry

neutral axis

$$\sigma_z = \frac{My}{I} \quad \sigma_{z\max} = \frac{Mc}{I} \quad (4.13)$$

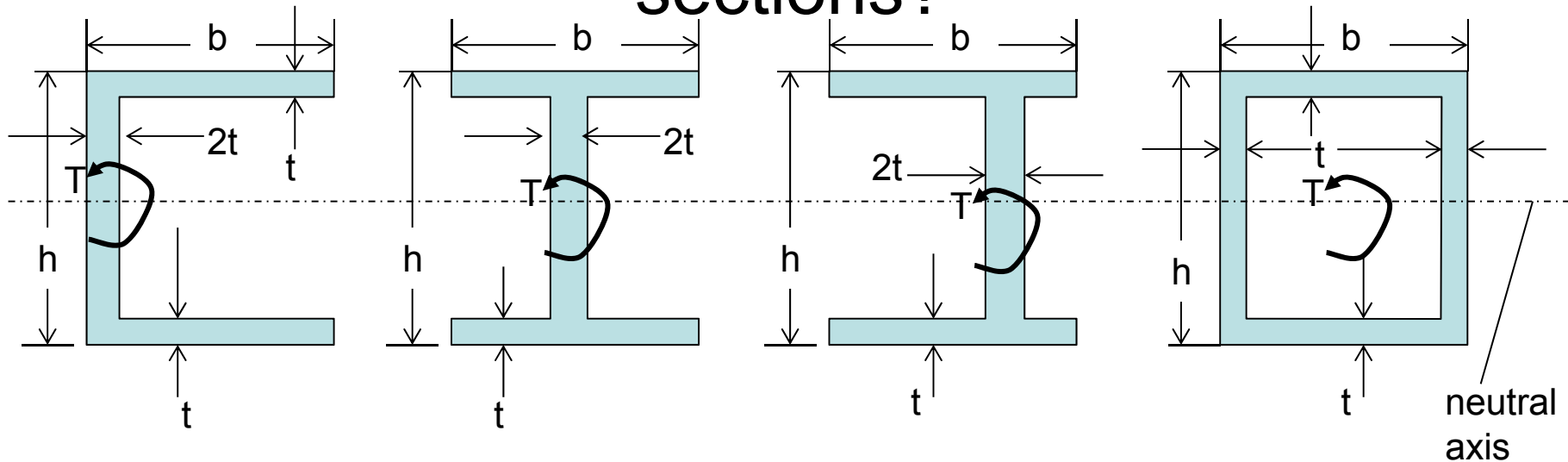
- the normal stress σ_z depends on applied moment and cross-sectional properties (I, c)
- these properties do not change if the beam is closed or open; therefore, **there is no difference in bending stresses between closed and open cross sections**
- this is also valid for unsymmetric cross-sections

Carrying this one step further...



- all these cross-sections have exactly the same neutral axis location, at the mid-height of each cross-section (due to symmetry)
- all these cross-sections have exactly the same moment of inertia
- therefore, all these cross-sections have exactly the same maximum stress and, from a strength perspective **are equivalent**

What about pure torsion of the same cross-sections?

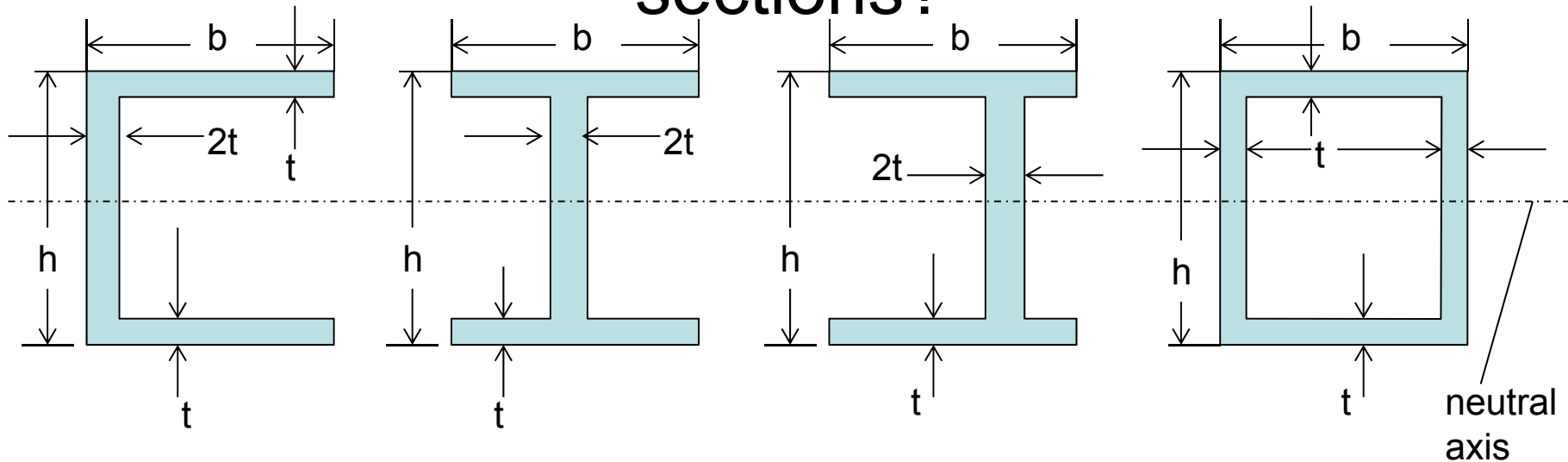


- from eq (4.11) the first three cross sections have exactly the same torsion constant J (as long as $b, h \gg t$):

$$J = \frac{1}{3} \sum_{i=1}^N s_i t_i^3$$

- but the fourth one, being closed, has J that is thousands of times bigger than the other three!

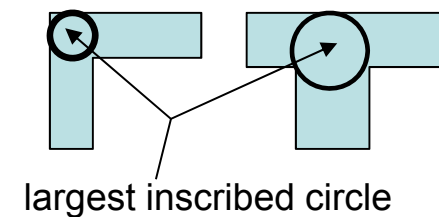
What about pure torsion of the same cross-sections?



- the maximum shear stress is a bit more complicated; it is given by a variation of eq. (4.10)

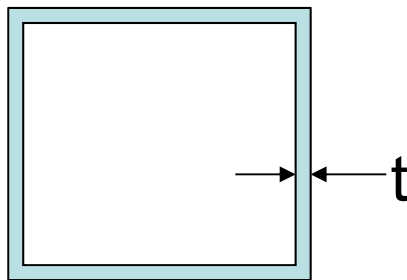
$$\tau_{\max} = \frac{Td}{J}$$

where $d=2t$ and its location is at one of the points where the largest inscribed circle touches the boundary

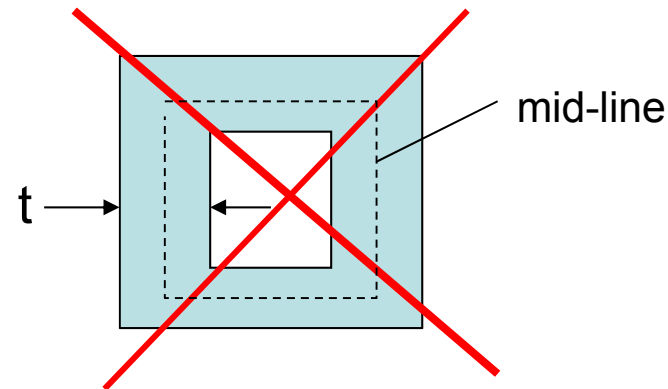


Approximations for thin-walled structure

- thin-walled structure $\Rightarrow t \ll$ other dimensions of cross-section



thin walled

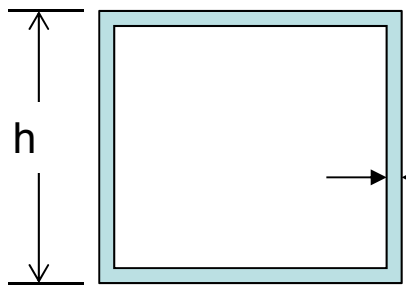


not thin walled

- implications of thin walled assumption:
 - stresses are constant through the thickness
 - cross-section can be represented by its mid-line
 - t^2 , t^3 or higher order terms can be neglected in the calculations

Approximations for thin-walled structure

- example: calculate moment of inertia for following cross-section



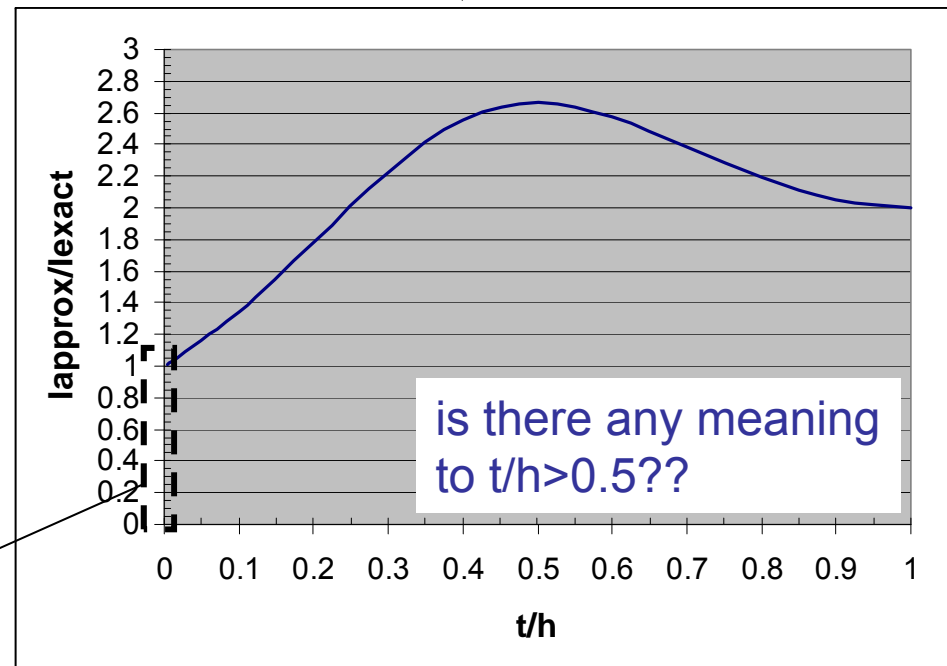
$$I = \frac{hh^3}{12} - \frac{(h-2t)(h-2t)^3}{12} = \frac{h^4}{12} - \frac{1}{12} \left[h^4 - 4h^3(2t) + 6h^2(2t)^2 - 4h^2(2t)^3 + (2t)^4 \right]$$

neglected terms

$$\approx \frac{2}{3}h^3t$$

typical aerospace structure has $t/h < 0.02$ and the error is less than 6% (usually less than 1.5%); at the other extreme, for $t/h = 0.5$ the error is 167%

typical thin-walled structure in this box



Approximations for thin-walled structure

- for torsion, the same assumption of thin walled structure led to the expression for J in the form of eq. (4.7)

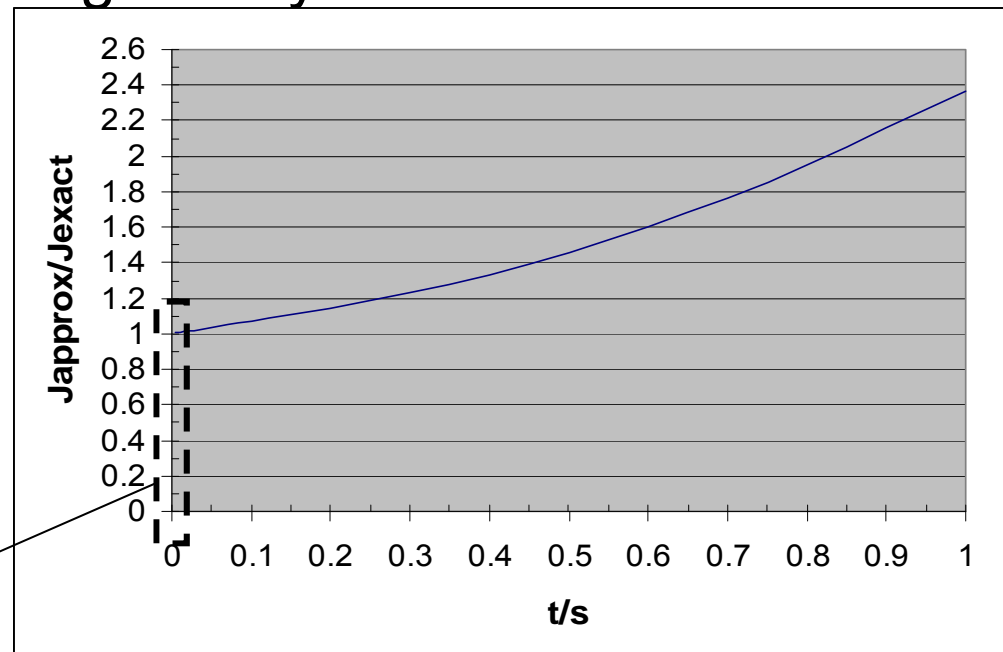
$$J = \frac{st^3}{3}$$

- for (open) sections where $t \ll s$ is not valid, a more accurate expression for J is given by

$$J = \frac{st^3}{3} \left[1 - 0.63 \frac{t}{s} \left(1 - \frac{t^4}{12s^4} \right) \right]$$

typical aerospace structure has $t/s < 0.02$ and the error is less than 1%; at the other extreme, for $t/s = 1$ the error is 137%

typical thin-walled structure in this box



Application Session 1