

(1) Determination of deflection at pt D

- first, determine the reaction forces R<sub>1</sub>, R<sub>2</sub>, R<sub>3</sub>
  - horizontal equilibrium:

$$R_3 = Q$$

• vertical equilibrium:

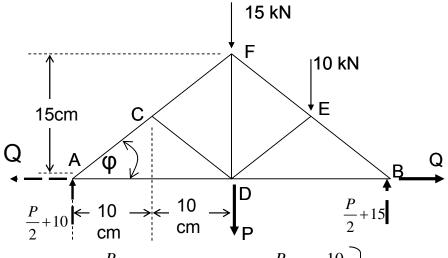
$$R_1 + R_2 = P + 25kN$$

moments about point A:

$$P(20) + 15(20) + 10(30) = R_2(40) \Rightarrow$$

$$R_2 = \frac{P}{2} + 15$$

• then,  $R_1 = \frac{P}{2} + 10$ 



- second, determine the forces in each truss member as a function of P and Q
- note that

$$\tan \varphi = \frac{7.5}{10} \Rightarrow \varphi = 36.87^{\circ}$$

$$F_{AC} \sin \phi + \frac{P}{2} + 10 = 0 \Rightarrow F_{AC} = -\frac{P}{2 \sin \phi} - \frac{10}{\sin \phi}$$

$$\Rightarrow F_{AD} = Q + \left(\frac{P}{2} + 10\right) \cot \phi$$

$$F_{AC} \cos \phi + F_{AD} - Q = 0$$

$$F_{AC} \sin \phi + F_{CD} \sin \phi = F_{CF} \sin \phi$$

$$F_{AC} \cos \phi - F_{CD} \cos \phi = F_{CF} \cos \phi$$

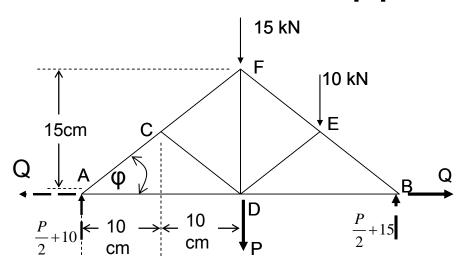
$$\Rightarrow F_{CD} = 0$$

$$F_{CF} = F_{AC} = -\frac{P}{2 \sin \phi} - \frac{10}{\sin \phi}$$

$$F_{BE}\sin\phi + \frac{P}{2} + 15 = 0 \Rightarrow F_{BE} = -\frac{P}{2\sin\phi} - \frac{15}{\sin\phi}$$

$$F_{BE}\cos\phi + F_{BD} - Q = 0$$

$$\Rightarrow F_{BD} = Q + \left(\frac{P}{2} + 15\right)\cot\phi$$



$$F_{AD} = Q + \left(\frac{P}{2} + 10\right) \cot \phi$$

$$F_{CD} = 0$$

$$F_{CF} = F_{AC} = -\frac{P}{2 \sin \phi} - \frac{10}{\sin \phi}$$

$$F_{BD} = Q + \left(\frac{P}{2} + 15\right) \cot \phi$$

$$F_{BE} = -\frac{P}{2 \sin \phi} - \frac{15}{\sin \phi}$$

$$F_{BE}\sin\phi + F_{ED}\sin\phi + 10 = F_{FE}\sin\phi$$

$$F_{BE}\cos\phi - F_{ED}\cos\phi = F_{FE}\cos\phi$$

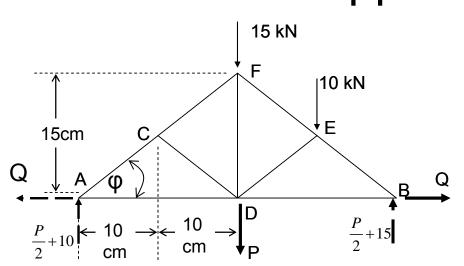
$$2F_{ED}\sin\phi\cos\phi + 10\cos\phi = 0 \Rightarrow F_{ED} = -\frac{10}{2\sin\phi}$$

$$F_{FE} = F_{BE} - F_{ED} = -\frac{P}{2\sin\phi} - \frac{10}{\sin\phi}$$

$$F_{CF} \sin \phi + F_{FD} + F_{FE} \sin \phi + 15 = 0 \Rightarrow F_{FD} = -15 + \frac{P}{2} + 10 + \frac{P}{2} + 10 \Rightarrow F_{FD} = P + 5$$

as a check, horizontal equilibrium at point F:

$$F_{CF}\cos\phi \stackrel{?}{=} F_{FE}\cos\phi \Rightarrow -\frac{P}{2\sin\phi} - \frac{10}{\sin\phi} = -\frac{P}{2\sin\phi} - \frac{10}{\sin\phi}$$



$$F_{AD} = Q + \left(\frac{P}{2} + 10\right) \cot \phi \qquad F_{BE} = -\frac{P}{2 \sin \phi} - \frac{15}{\sin \phi}$$

$$F_{CD} = 0$$

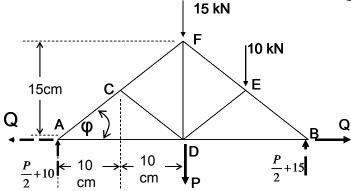
$$F_{CF} = F_{AC} = -\frac{P}{2 \sin \phi} - \frac{10}{\sin \phi} \qquad F_{ED} = -\frac{10}{2 \sin \phi}$$

$$F_{BD} = Q + \left(\frac{P}{2} + 15\right) \cot \phi \qquad F_{FE} = -\frac{P}{2 \sin \phi} - \frac{10}{\sin \phi}$$

$$F_{BE} = -\frac{P}{2 \sin \phi} - \frac{15}{\sin \phi} \qquad F_{FD} = P + 5$$

• third, determine the total energy in the system from eq. (12.23) summed over all truss members  $C_i = \sum_{j=1}^{K_i} \frac{F_j^2 L_j}{2E_j A}$  (12.23)

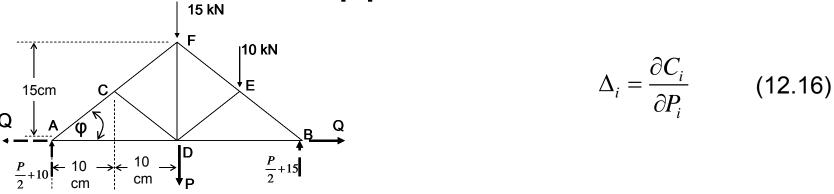
$$C_{i} = \frac{1}{2EA} \begin{bmatrix} \left( -\frac{P}{2\sin\phi} - \frac{10}{\sin\phi} \right)^{2} L_{AC} + \left( Q + \left( \frac{P}{2} + 10 \right) \cot\phi \right)^{2} L_{AD} + \left( \frac{10}{2\sin\phi} \right)^{2} L_{ED} + \left( -\frac{P}{2\sin\phi} - \frac{15}{\sin\phi} \right)^{2} L_{BE} + \left( Q + \left( \frac{P}{2} + 15 \right) \cot\phi \right)^{2} L_{BD} + \left( -\frac{P}{2\sin\phi} - \frac{10}{\sin\phi} \right)^{2} L_{FC} + (P + 5)^{2} L_{FD} + \left( -\frac{P}{2\sin\phi} - \frac{10}{\sin\phi} \right)^{2} L_{FE} \end{bmatrix}$$



$$C_{i} = \frac{1}{2EA} \left[ \left( -\frac{P}{2\sin\phi} - \frac{10}{\sin\phi} \right)^{2} L_{AC} + \left( Q + \left( \frac{P}{2} + 10 \right) \cot\phi \right)^{2} L_{AD} + \left( \frac{10}{2\sin\phi} \right)^{2} L_{ED} + \left( -\frac{P}{2\sin\phi} - \frac{15}{\sin\phi} \right)^{2} L_{BE} + \left[ Q + \left( \frac{P}{2} + 15 \right) \cot\phi \right)^{2} L_{BD} + \left( -\frac{P}{2\sin\phi} - \frac{10}{\sin\phi} \right)^{2} L_{FC} + (P+5)^{2} L_{FD} + \left( -\frac{P}{2\sin\phi} - \frac{10}{\sin\phi} \right)^{2} L_{FE} \right]$$

fourth, obtain the required deflection using eq. (12.16)

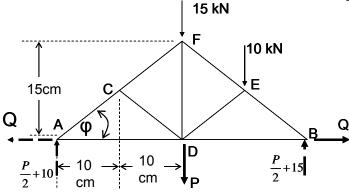
$$\Delta_i = \frac{\partial C_i}{\partial P_i} \tag{12.16}$$



$$C_{i} = \frac{1}{2EA} \left[ \left( -\frac{P}{2\sin\phi} - \frac{10}{\sin\phi} \right)^{2} L_{AC} + \left( Q + \left( \frac{P}{2} + 10 \right) \cot\phi \right)^{2} L_{AD} + \left( \frac{10}{2\sin\phi} \right)^{2} L_{ED} + \left( -\frac{P}{2\sin\phi} - \frac{15}{\sin\phi} \right)^{2} L_{BE} + \left( Q + \left( \frac{P}{2} + 15 \right) \cot\phi \right)^{2} L_{BD} + \left( -\frac{P}{2\sin\phi} - \frac{10}{\sin\phi} \right)^{2} L_{FC} + (P + 5)^{2} L_{FD} + \left( -\frac{P}{2\sin\phi} - \frac{10}{\sin\phi} \right)^{2} L_{FE} \right]$$

• the vertical displacement at point D is then:

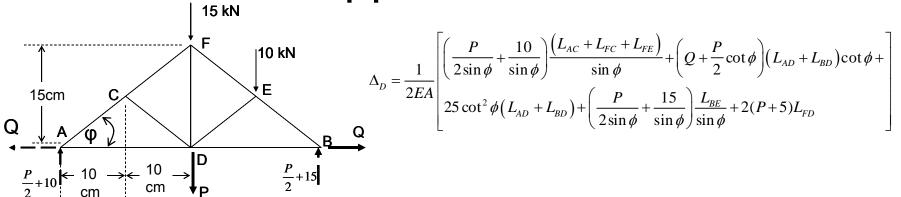
$$\Delta_{D} = \frac{1}{2EA} \begin{bmatrix} 2\left(\frac{P}{2\sin\phi} + \frac{10}{\sin\phi}\right) \frac{L_{AC}}{2\sin\phi} + 2\left(Q + \left(\frac{P}{2} + 10\right)\cot\phi\right) \frac{L_{AD}\cot\phi}{2} + 2\left(\frac{P}{2\sin\phi} + \frac{15}{\sin\phi}\right) \frac{L_{BE}}{2\sin\phi} + \\ 2\left(Q + \left(\frac{P}{2} + 15\right)\cot\phi\right) \frac{L_{BD}\cot\phi}{2} + 2\left(\frac{P}{2\sin\phi} + \frac{10}{\sin\phi}\right) \frac{L_{FC}}{2\sin\phi} + 2(P + 5)L_{FD} + 2\left(\frac{P}{2\sin\phi} + \frac{10}{\sin\phi}\right) \frac{L_{FE}}{2\sin\phi} \end{bmatrix}$$



$$\Delta_{D} = \frac{1}{2EA} \begin{bmatrix} 2\left(\frac{P}{2\sin\phi} + \frac{10}{\sin\phi}\right) \frac{L_{AC}}{2\sin\phi} + 2\left(Q + \left(\frac{P}{2} + 10\right)\cot\phi\right) \frac{L_{AD}\cot\phi}{2} + 2\left(\frac{P}{2\sin\phi} + \frac{15}{\sin\phi}\right) \frac{L_{BE}}{2\sin\phi} + 2\left(Q + \left(\frac{P}{2} + 15\right)\cot\phi\right) \frac{L_{BD}\cot\phi}{2} + 2\left(\frac{P}{2\sin\phi} + \frac{10}{\sin\phi}\right) \frac{L_{FC}}{2\sin\phi} + 2\left(P + 5\right)L_{FD} + 2\left(\frac{P}{2\sin\phi} + \frac{10}{\sin\phi}\right) \frac{L_{FE}}{2\sin\phi} \end{bmatrix}$$

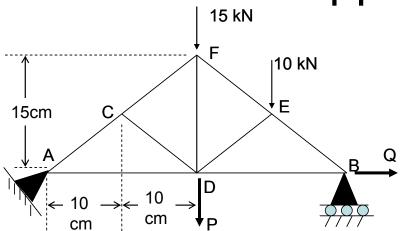
#### which simplifies to:

$$\Delta_{D} = \frac{1}{2EA} \left[ \frac{P}{2\sin\phi} + \frac{10}{\sin\phi} \frac{(L_{AC} + L_{FC} + L_{FE})}{\sin\phi} + (Q + \frac{P}{2}\cot\phi)(L_{AD} + L_{BD})\cot\phi + 25\cot^{2}\phi(L_{AD} + L_{BD}) + (Q + \frac{P}{2}\cot\phi)(L_{AD} + L_{BD})\cot\phi + 25\cot^{2}\phi(L_{AD} + L_{BD}) + (Q + \frac{P}{2}\cot\phi)(L_{AD} + L_{BD})\cot\phi + 25\cot^{2}\phi(L_{AD} + L_{BD}) + (Q + \frac{P}{2}\cot\phi)(L_{AD} + L_{BD})\cot\phi + 25\cot^{2}\phi(L_{AD} + L_{BD}) + (Q + \frac{P}{2}\cot\phi)(L_{AD} + L_{BD})\cot\phi + 25\cot^{2}\phi(L_{AD} + L_{BD}) + (Q + \frac{P}{2}\cot\phi)(L_{AD} + L_{BD})\cot\phi + 25\cot^{2}\phi(L_{AD} + L_{BD}) + (Q + \frac{P}{2}\cot\phi)(L_{AD} + L_{BD})\cot\phi + 25\cot^{2}\phi(L_{AD} + L_{BD}) + (Q + \frac{P}{2}\cot\phi)(L_{AD} + L_{BD})\cot\phi + 25\cot^{2}\phi(L_{AD} + L_{BD}) + (Q + \frac{P}{2}\cot\phi)(L_{AD} + L_{BD})\cot\phi + 25\cot^{2}\phi(L_{AD} + L_{BD}) + (Q + \frac{P}{2}\cot\phi)(L_{AD} + L_{BD})\cot\phi + 25\cot^{2}\phi(L_{AD} + L_{BD}) + (Q + \frac{P}{2}\cot\phi)(L_{AD} + L_{BD})\cot\phi + 25\cot^{2}\phi(L_{AD} + L_{BD}) + (Q + \frac{P}{2}\cot\phi)(L_{AD} + L_{BD})\cot\phi + (Q + \frac{P}{2}\cot\phi)(L_{AD} + L_{D$$



 fifth, set to zero any fictitious forces (in this case both P and Q); then, the required deflection at point D is

$$\Delta_D(P=0,Q=0) = \frac{1kN}{2EA} \left[ \frac{10}{\sin^2 \phi} \left( L_{AC} + L_{FC} + L_{FE} \right) + 25\cot^2 \phi \left( L_{AD} + L_{BD} \right) + 15\frac{L_{BE}}{\sin^2 \phi} + 10L_{FD} \right]$$

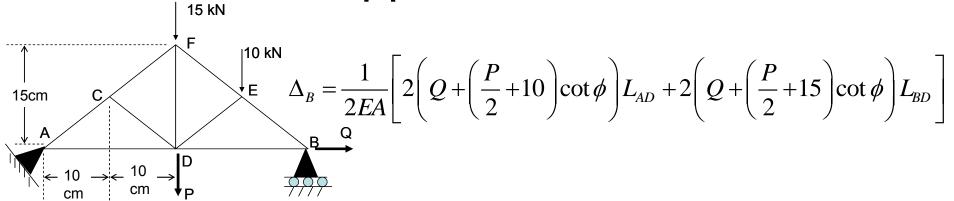


(2) Determination of force Q such that the horizontal deflection at point B is zero

$$\Delta_i = \frac{\partial C_i}{\partial Q}$$

• the first three steps, up to obtaining the total energy expression are the same as before; to determine the horizontal deflection at point B, differentiate the total energy with respect to Q to obtain

$$\Delta_{B} = \frac{1}{2EA} \left[ 2\left(Q + \left(\frac{P}{2} + 10\right)\cot\phi\right) L_{AD} + 2\left(Q + \left(\frac{P}{2} + 15\right)\cot\phi\right) L_{BD} \right]$$

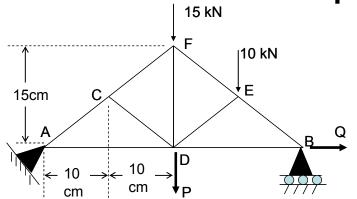


• recalling now that P is a fictitious force (but Q is no longer fictitious because we want to determine its value so that the displacement at B is zero) we can simplify:

$$\Delta_B = \frac{10kN}{2EA} \left[ 2(Q+10\cot\phi)L_{AD} + 2(Q+15\cot\phi)L_{BD} \right]$$

• for the deflection at B to be zero, set  $\Delta_B$ =0 and solve for Q:

$$Q = -5\cot\phi \frac{\left(2L_{AD} + 3L_{BD}\right)}{L_{AD} + L_{BD}}$$



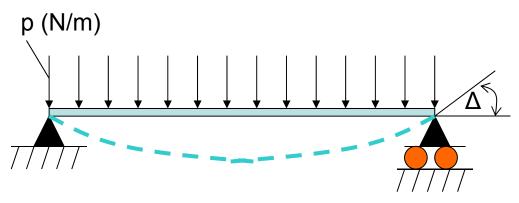
(3) With the Q obtained acting, determine the new vertical deflection at point D

$$\Delta_{D} = \frac{1}{2EA} \left[ \left( \frac{P}{2\sin\phi} + \frac{10}{\sin\phi} \right) \frac{\left( L_{AC} + L_{FC} + L_{FE} \right)}{\sin\phi} + \left( Q + \frac{P}{2}\cot\phi \right) \left( L_{AD} + L_{BD} \right) \cot\phi + \frac{1}{2} \left( \frac{P}{2\sin\phi} + \frac{15}{\sin\phi} \right) \frac{L_{BE}}{\sin\phi} + 2(P+5)L_{FD} \right] \right]$$

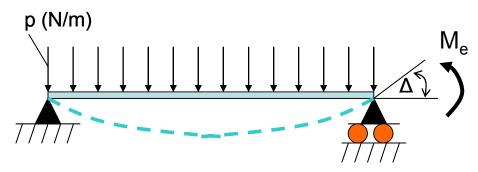
- the derivation is the same as in part (1) up to the expression for  $\Delta_D$  with P and Q included
- now only P is fictitious, Q was determined in part (2); set P=0

$$\Delta_{D} = \frac{1}{2EA} \left[ \frac{10kN}{\sin^{2} \phi} \left( L_{AC} + L_{FC} + L_{FE} \right) + Q \left( L_{AD} + L_{BD} \right) \cot \phi + 25kN \cot^{2} \phi \left( L_{AD} + L_{BD} \right) + \frac{15kN}{\sin^{2} \phi} L_{BE} + 10kNL_{FD} \right]$$

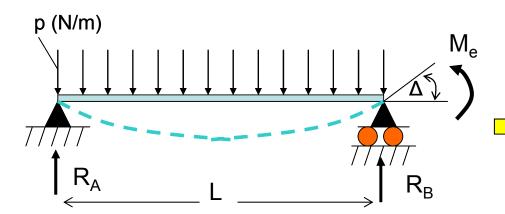
with Q as determined in part (2)



- for a simply supported beam under distributed load p, determine the angle of rotation  $\Delta$  at the right end of the beam
- to do this we need to obtain the energy stored in the beam and differentiate it with respect to the load that causes the rotation  $\boldsymbol{\Delta}$
- but there is no such load applied



- since a rotation is caused by a moment, apply a fictitious end moment  $\rm M_e$  and compute the energy in the beam due to both p and  $\rm M_e$  acting
- first obtain the reaction forces



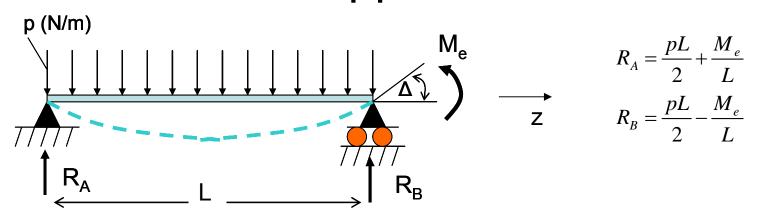
$$R_{A} + R_{B} = pL$$

$$M_{e} \qquad R_{A}L - \int_{0}^{L} pzdz = M_{e} \Rightarrow R_{A}L - p\frac{L^{2}}{2} = M_{e}$$

$$R_{A} = \frac{pL}{2} + \frac{M_{e}}{L}$$

$$R_{B} = \frac{pL}{2} - \frac{M_{e}}{L}$$

13

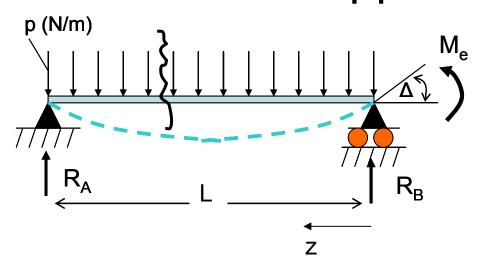


• to use Castigliano's theorem we need eq. (12.24)

$$C_{i} = \int_{0}^{L} \frac{M^{2}}{2EI} dz$$
 (12.24)

which expresses the energy in the beam in terms of the bending moment M along the beam

 therefore, the bending moment M must be determined as a function of z



$$R_{A} = \frac{pL}{2} + \frac{M_{e}}{L}$$

$$R_{B} = \frac{pL}{2} - \frac{M_{e}}{L}$$

$$C_{i} = \int_{0}^{L} \frac{M^{2}}{2EI} dz \qquad (12.24)$$

15

 cutting the beam at some location z and applying moment equilibrium:

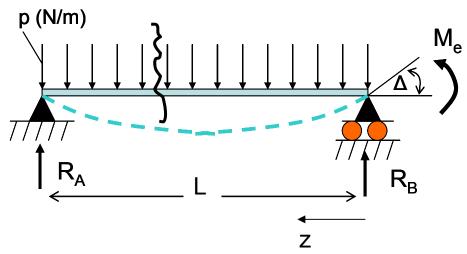
Taking moments about the cut:

R<sub>A</sub>

$$M = -p\frac{z^{2}}{2} + \left(p\frac{L}{2} + \frac{M_{e}}{L}\right)z$$

$$M = -p\frac{z^{2}}{2} + \left(p\frac{L}{2} + \frac{M_{e}}{L}\right)z$$

$$M = -p\frac{z^{2}}{2} + \left(p\frac{L}{2} + \frac{M_{e}}{L}\right)z$$



$$C_i = \int_0^L \frac{M^2}{2EI} dz$$

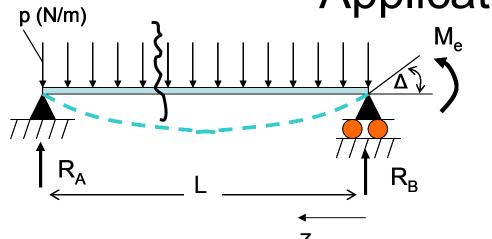
$$M = -p\frac{z^2}{2} + \left(p\frac{L}{2} + \frac{M_e}{L}\right)z$$

• then,

$$M^{2} = p^{2} \frac{z^{4}}{4} + \left(\frac{p^{2}L^{2}}{4} + \frac{M_{e}^{2}}{L^{2}} + pM_{e}\right)z^{2} - pz^{3}\left(p\frac{L}{2} + \frac{M_{e}}{L}\right)$$

and substituting in the energy expression

$$C_{i} = \frac{1}{2EI} \left[ p^{2} \frac{L^{5}}{20} + \left( \frac{p^{2}L^{2}}{4} + \frac{M_{e}^{2}}{L^{2}} + pM_{e} \right) \frac{L^{3}}{3} - p \frac{L^{4}}{4} \left( p \frac{L}{2} + \frac{M_{e}}{L} \right) \right]$$



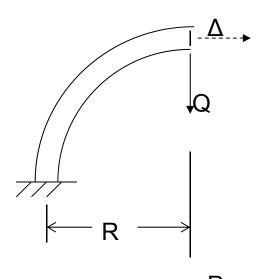
$$C_{i} = \frac{1}{2EI} \begin{bmatrix} p^{2} \frac{L^{5}}{20} + \left(\frac{p^{2}L^{2}}{4} + \frac{M_{e}^{2}}{L^{2}} + pM_{e}\right) \frac{L^{3}}{3} - \frac{L^{4}}{4} \left(p \frac{L}{2} + \frac{M_{e}}{L}\right) \end{bmatrix}$$

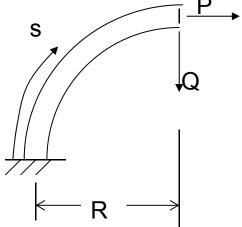
• to determine the rotation Δ at the end, apply Castigliano's theorem:

$$\Delta = \frac{\partial C_i}{\partial M_e}$$
• to obtain,  $\Delta = \frac{1}{2EI} \left[ \frac{2}{3} M_e L + p \frac{L^3}{12} \right]$ 

• finally, since  $M_e$  was a fictitious moment, setting  $M_e$ =0 gives the final answer:  $\Delta = \frac{pL^3}{\Delta}$ 

17





Determine the horizontal displacement  $\Delta$  for a curved beam of radius R and bending stiffness EI under a vertical force Q

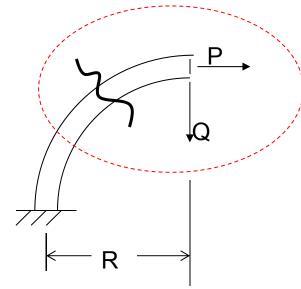
- we introduce a fictitious force P in the direction of the required displacement:
- the energy expression (12.24),

$$C_i = \int_0^L \frac{M^2}{2EI} dz$$

becomes for a curved beam:

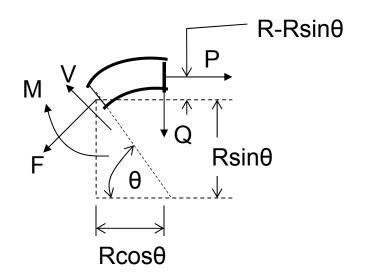
$$C_i = \int_0^{\pi R/2} \frac{M^2}{2EI} ds$$

• and since ds=Rd $\theta$ ,  $C_i = \int_0^{\pi/2} \frac{M^2}{2EI} Rd\theta$ 



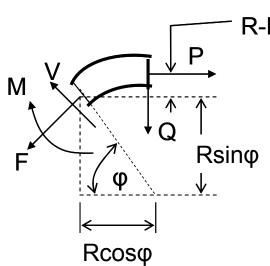
$$C_i = \int_0^{\pi/2} \frac{M^2}{2EI} Rd\theta$$

• to determine the internal moment M, make a cut and place it in moment equilibrium:



$$M + PR(1 - \sin \theta) + QR\cos \theta = 0 \Rightarrow$$
$$M = -PR(1 - \sin \theta) - QR\cos \theta$$

# Application 4 R-Rsin $\varphi$ $C_i = \int_0^{M^2} \frac{M^2}{2EI} R d\theta$ M = 0Castigliano's Second Theorem –



$$C_{i} = \int_{0}^{\pi/2} \frac{M^{2}}{2EI} R d\theta$$

$$M = -PR(1 - \sin \theta) - QR \cos \theta$$

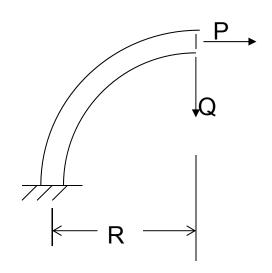
now

$$M^{2} = P^{2}R^{2} (1 - \sin \theta)^{2} + Q^{2}R^{2} \cos^{2} \theta + 2PQR^{2} (1 - \sin \theta) \cos \theta =$$

$$P^{2}R^{2} \left( 1 + \frac{1 - \cos 2\theta}{2} - 2\sin \theta \right) + Q^{2}R^{2} \frac{1 + \cos 2\theta}{2} + 2PQR^{2} \left( \cos \theta - \frac{1}{2}\sin 2\theta \right)$$

 substituting in the energy expression and evaluating the integrals:

$$C_{i} = \frac{R}{2EI} \begin{bmatrix} P^{2}R^{2} \left[ \frac{\pi}{2} + \frac{1}{2} \left( \theta - \frac{1}{2} \sin 2\theta \right)_{0}^{\pi/2} + 2 \left( \cos \theta \right)_{0}^{\pi/2} \right] + \frac{Q^{2}R^{2}}{2} \left[ \frac{\pi}{4} + \frac{1}{2} \left( \sin 2\theta \right)_{0}^{\pi/2} \right] + \frac{1}{2} \left[ \sin 2\theta \right]_{0}^{\pi/2} + \frac{1}{4} \left( \cos 2\theta \right)_{0}^{\pi/2} \right] \\ \bullet \text{ which simplifies to: } C_{i} = \frac{R^{3}}{2EI} P^{2} \left( \frac{3\pi}{4} - 2 \right) + \frac{Q^{2}\pi}{4} + PQ \end{bmatrix}$$



$$C_i = \frac{R^3}{2EI} \left[ P^2 \left( \frac{3\pi}{4} - 2 \right) + \frac{Q^2 \pi}{4} + PQ \right]$$

• to determine the horizontal displacement at the top right:

$$\Delta = \frac{\partial C_i}{\partial P}$$

• which gives: 
$$\Delta = \frac{R^3}{2EI} \left[ 2P \left( \frac{3\pi}{4} - 2 \right) + Q \right]$$

 and since P is a fictitious force, setting P=0 gives the final answer:

$$\Delta = \frac{QR^3}{2EI}$$

#### Castigliano's first theorem

- it is completely analogous to Castigliano's second theorem
- for a body under loads  $Q_1$ ,  $Q_2$ , ... (where  $Q_i$  can be either forces or moments), if the strain energy stored U can be expressed in terms of displacements  $q_1$ ,  $q_2$ , ... (where  $q_i$  can be either deflections or rotations) then:

$$Q_i = \frac{\partial U}{\partial \Delta_i}$$
 (12.29)

that is, to find the force applied to a body in a specific direction, need to differentiate the internal strain energy U with respect to the corresponding deflection

compare with Castigliano's 2<sup>nd</sup> theorem, eq (12.16): 
$$\Delta_i = \frac{\partial C_i}{\partial Q_i} \quad (12.16)$$

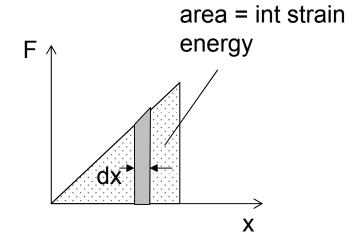
## Castigliano's first theorem – linear spring

• what makes the use of this theorem more difficult is the need to express the internal strain energy in terms of deflections (unlike the 2<sup>nd</sup> theorem where expressing the complementary energy in terms of forces and moments was straightforward)

**Linear Spring** 



linear spring with spring constant k; force F causes displacement x=> F = kx



$$U = \int_{0}^{x} F dx$$

but F=kx substituting

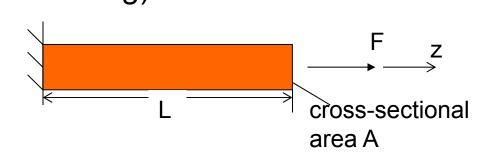
$$U = \int_{0}^{x} kx dx = \frac{1}{2} kx^{2}$$

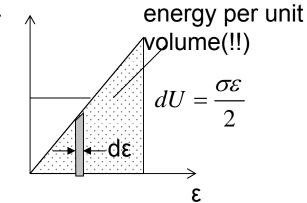
well known results from basic physics and dynamics

 $(12.30)_{23}$ 

### Castigliano's first theorem – truss member

Bar (or beam) in tension or compression (no buckling)



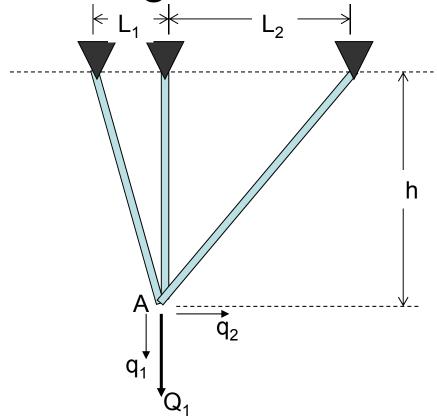


area = internal

$$U = \iiint_{vol} dU = \iiint_{vol} \frac{\sigma\varepsilon}{2} dx dy dz = \int_{0}^{L} \frac{A\sigma\varepsilon}{2} dz$$
but  $\sigma = E\varepsilon$ 

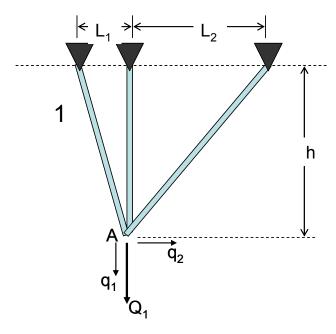
$$\varepsilon = \frac{\Delta}{L}$$

• and since  $\Delta$ , E, A, L are independent of z:  $U = \frac{EA\Delta^2}{2L}$  (12.31)



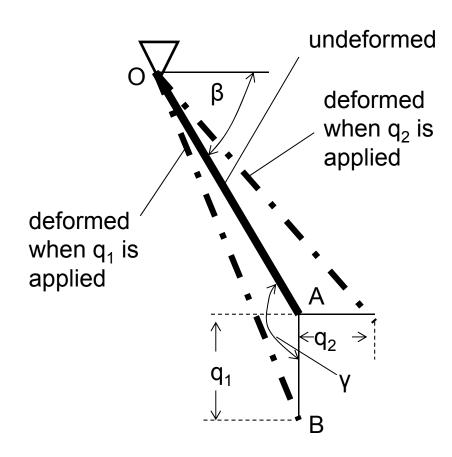
all 3 members have stiffness E and cross-sectional area A

• determine the displacements  $q_1$  and  $q_2$  of point A when a known load  $Q_1$  is applied



all 3 members have stiffness E and cross-sectional area A

• first determine the deflections of the end of each truss member **along the member axis**, if the end moves by  $q_1$  and  $q_2$ 



Consider a beam of length L, stiffness E and cross-sectional area A pinned at one end with applied displacements q<sub>1</sub> and q<sub>2</sub> at the other

- analyze the effect of each deflection separately and superpose the results
- first, consider q<sub>1</sub> only
- from the law of cosines:

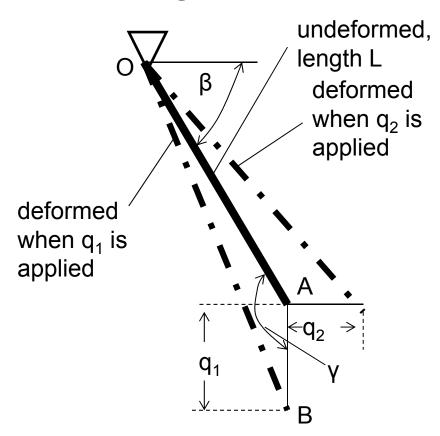
$$OB^2 = OA^2 + AB^2 - 2(OA)(AB)\cos\gamma$$

• But: OA=L, AB= $q_1$  and  $\gamma$ =90+ $\beta$ ; So:

$$OB^2 = L^2 + q_1^2 + 2Lq_1 \sin \beta$$

recall that:

 $\cos(90 + \beta) = \cos 90 \cos \beta - \sin 90 \sin \beta$ 



$$OB^2 = L^2 + q_1^2 + 2Lq_1 \sin \beta$$

the change in beam length is then,

$$\Delta_{q_1} = OB - OA = \sqrt{L^2 + q_1^2 + 2Lq_1 \sin \beta} - L$$

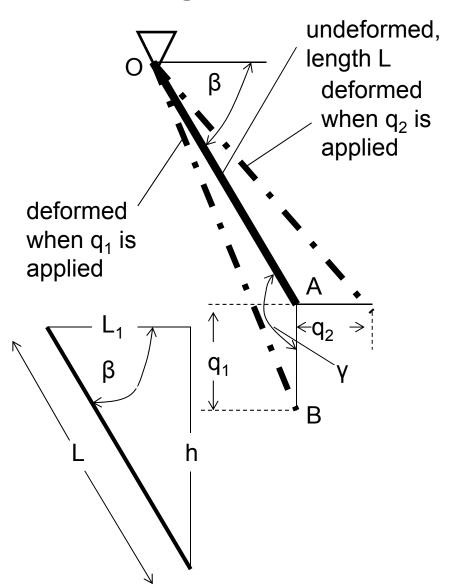
• for small deflections,  $q_1$  is small and  $q_1^2$  can be neglected; then

$$\Delta_{q_1} = OB - OA \doteq \sqrt{L^2 + 2Lq_1 \sin \beta} - L$$

• now,  $2Lq_1sin\beta$  is small compared to  $L^2$  so the square root can be expanded in a Taylor series for small  $q_1$ :

$$\sqrt{L^{2} + 2Lq_{1}\sin\beta} \doteq \sqrt{L^{2} + 2Lq_{1}\sin\beta} \Big]_{q_{1}=0} + \frac{d}{dq_{1}} \Big[ \sqrt{L^{2} + 2Lq_{1}\sin\beta} \Big]_{q_{1}=0} q_{1} + \dots = L + q_{1}\sin\beta$$

$$L + \left[ \frac{1}{2} \frac{2L\sin\beta}{\sqrt{L^{2} + 2Lq_{1}\sin\beta}} \right]_{q_{1}=0} q_{1} + \dots = L + q_{1}\sin\beta$$
28



$$\Delta_{q_1} = OB - OA \doteq \sqrt{L^2 + 2Lq_1 \sin \beta} - L$$
$$\sqrt{L^2 + 2Lq_1 \sin \beta} \doteq L + q_1 \sin \beta$$

• substituting:

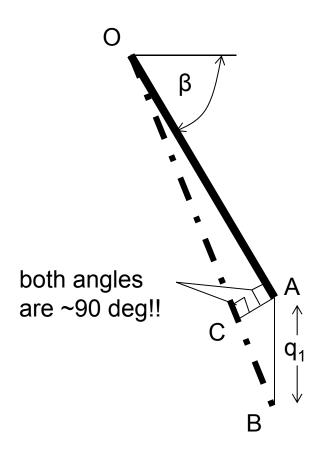
$$\Delta_{q_1} = q_1 \sin \beta$$

and using the fact that

$$\sin \beta = \frac{h}{L}$$

we get finally:

$$\Delta_{q_1} = q_1 \frac{h}{L}$$

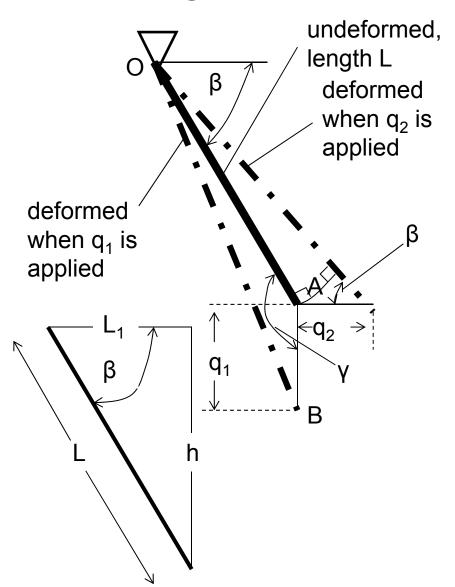


- An alternate approach to get the same result.
   It is simpler but requires some assumptions associated with small deflection theory that can be confusing:
- With radius OA draw a circle segment that intersects OB at C. If the deflection q<sub>1</sub> is small, then the arc AC can be approximated by a straight line AC. Then:

$$OC = OA = L$$
  
 $\Delta_{q_1} = OB - OC = BC$ 

 Now angle CAB=β because it has mutually perpendicular sides with β. Then,

$$BC = q_1 \sin \angle CAB = q_1 \sin \beta$$

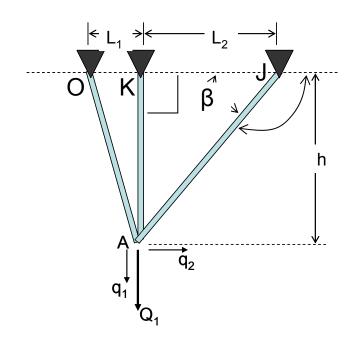


• in a completely analogous fashion, we can find the extension along the beam when it displaces horizontally by q<sub>2</sub>:

$$\Delta_{q_2} = q_2 \cos \beta = q_2 \frac{L_1}{L}$$

• both  $\Delta q_1$  and  $\Delta q_2$  are along the length of the beam so the total deflection when the beam displaces by  $q_1$  and  $q_2$  is the sum of the two:

$$\Delta = \Delta_{q_1} + \Delta_{q_2} = q_1 \sin \beta + q_2 \cos \beta = q_1 \frac{h}{L} + q_2 \frac{L_1}{L}$$



$$\Delta = \Delta_{q_1} + \Delta_{q_2} = q_1 \sin \beta + q_2 \cos \beta = q_1 \frac{h}{L} + q_2 \frac{L_1}{L}$$

- for OA:  $\Delta_{OA} = q_1 \sin \beta + q_2 \cos \beta = q_1 \frac{h}{L_{OA}} + q_2 \frac{L_1}{L_{OA}}$  for KA:  $\beta = 90^\circ = \lambda_{KA} = q_1 \sin 90 = q_1$ 

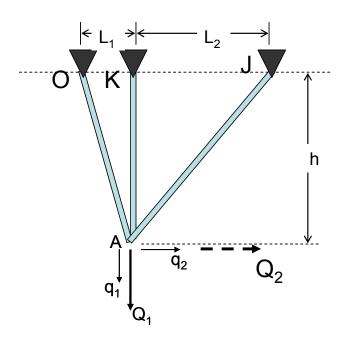
  - for JA:  $\beta \rightarrow 180 \beta = > \Delta_{JA} = q_1 \sin(180 \beta) + q_2 \cos(180 \beta) =$  $q_1 \sin \beta - q_2 \cos \beta = q_1 \frac{h}{L} - q_2 \frac{L_2}{L}$
  - use now eq (12.31) to determine the total strain energy,

$$U = \frac{EA\Delta^2}{2L} \quad (12.31)$$

substituting:

$$U = \frac{EA}{2} \left[ \frac{\Delta_{OA}^{2}}{L_{OA}} + \frac{\Delta_{KA}^{2}}{L_{KA}} + \frac{\Delta_{JA}^{2}}{L_{JA}} \right] = \frac{EA}{2} \left[ q_{1}^{2} \left( \frac{1}{L_{KA}} + \frac{h^{2}}{L_{OA}^{3}} + \frac{h^{2}}{L_{JA}^{3}} \right) + q_{2}^{2} \left( \frac{L_{1}^{2}}{L_{OA}^{3}} + \frac{L_{2}^{2}}{L_{JA}^{3}} \right) + 2q_{1}q_{2} \left( \frac{hL_{1}}{L_{OA}^{3}} - \frac{hL_{2}}{L_{JA}^{3}} \right) \right]$$
32

$$U = \frac{EA}{2} \left[ \frac{\Delta_{OA}^{2}}{L_{OA}} + \frac{\Delta_{KA}^{2}}{L_{KA}} + \frac{\Delta_{JA}^{2}}{L_{JA}} \right] = \frac{EA}{2} \left[ q_{1}^{2} \left( \frac{1}{L_{KA}} + \frac{h^{2}}{L_{OA}^{3}} + \frac{h^{2}}{L_{JA}^{3}} \right) + q_{2}^{2} \left( \frac{L_{1}^{2}}{L_{OA}^{3}} + \frac{L_{2}^{2}}{L_{JA}^{3}} \right) + 2q_{1}q_{2} \left( \frac{hL_{1}}{L_{OA}^{3}} - \frac{hL_{2}}{L_{JA}^{3}} \right) \right]$$



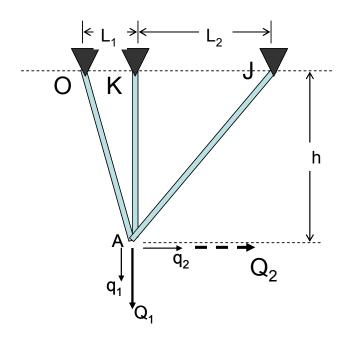
$$Q_i = \frac{\partial U}{\partial \Delta_i} \tag{12.29}$$

• use now eq (12.29) to obtain two equations for the loads  $Q_1$  and  $Q_2$ :

$$Q_{1} = EA \left[ q_{1} \left( \frac{1}{L_{KA}} + \frac{h^{2}}{L_{OA}^{3}} + \frac{h^{2}}{L_{JA}^{3}} \right) + q_{2} \left( \frac{hL_{1}}{L_{OA}^{3}} - \frac{hL_{2}}{L_{JA}^{3}} \right) \right]$$

$$Q_{2} = EA \left[ q_{1} \left( \frac{hL_{1}}{L_{OA}^{3}} - \frac{hL_{2}}{L_{JA}^{3}} \right) + q_{2} \left( \frac{L_{1}^{2}}{L_{OA}^{3}} + \frac{L_{2}^{2}}{L_{JA}^{3}} \right) \right]$$

• in the present problem Q<sub>2</sub>=0 and Q<sub>1</sub> is a given applied force



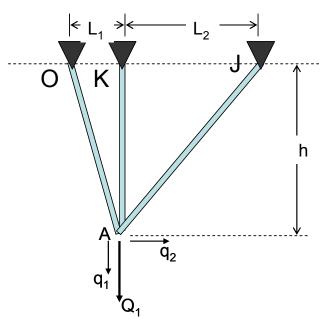
$$Q_{1} = EA \left[ q_{1} \left( \frac{1}{L_{KA}} + \frac{h^{2}}{L_{OA}^{3}} + \frac{h^{2}}{L_{JA}^{3}} \right) + q_{2} \left( \frac{hL_{1}}{L_{OA}^{3}} - \frac{hL_{2}}{L_{JA}^{3}} \right) \right]$$

$$0 = EA \left[ q_{1} \left( \frac{hL_{1}}{L_{OA}^{3}} - \frac{hL_{2}}{L_{JA}^{3}} \right) + q_{2} \left( \frac{L_{1}^{2}}{L_{OA}^{3}} + \frac{L_{2}^{2}}{L_{JA}^{3}} \right) \right]$$

- system of two equations in the two unknowns q<sub>1</sub> and q<sub>2</sub>
- numerical example: EA=10<sup>7</sup> (appropriate units)

 $L_1$ =50a,  $L_2$ =90a, h=120a,  $Q_1$ =10000 (appr. units)

then: L<sub>OA</sub>=130a, L<sub>KA</sub>=120a, L<sub>JA</sub>=150a



$$Q_{1} = EA \left[ q_{1} \left( \frac{1}{L_{KA}} + \frac{h^{2}}{L_{OA}^{3}} + \frac{h^{2}}{L_{JA}^{3}} \right) + q_{2} \left( \frac{hL_{1}}{L_{OA}^{3}} - \frac{hL_{2}}{L_{JA}^{3}} \right) \right]$$

$$0 = EA \left[ q_{1} \left( \frac{hL_{1}}{L_{OA}^{3}} - \frac{hL_{2}}{L_{JA}^{3}} \right) + q_{2} \left( \frac{L_{1}^{2}}{L_{OA}^{3}} + \frac{L_{2}^{2}}{L_{JA}^{3}} \right) \right]$$

•EA= $10^7$  (appropriate units) L<sub>1</sub>=50a, L<sub>2</sub>=90a, h=120a, Q<sub>1</sub>=10000 (app units)

substituting values, the system to be solved becomes:

$$191544q_1 - 4690q_2 = 10000a$$
$$-4690q_1 + 35378q_2 = 0$$

• solving:  $q_1=0.0524a$ ,  $q_2=0.0069a$ 

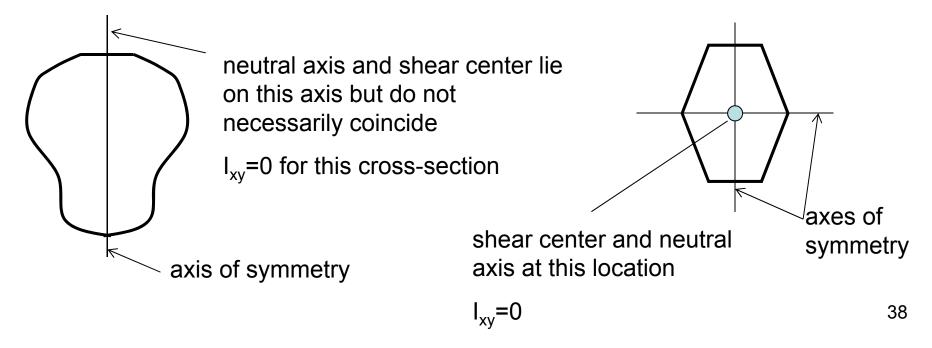
#### Wrap-up

- bending of beams
  - symmetric and unsymmetric cross-sections
  - direct stresses in beams due to bending
- shear of beams
  - direct stresses due to shear
  - shear stresses due to shear
  - closed versus open sections
  - idealization: constant shear flow when skins take no bending loads
  - determination of flange/boom areas

### Wrap-up

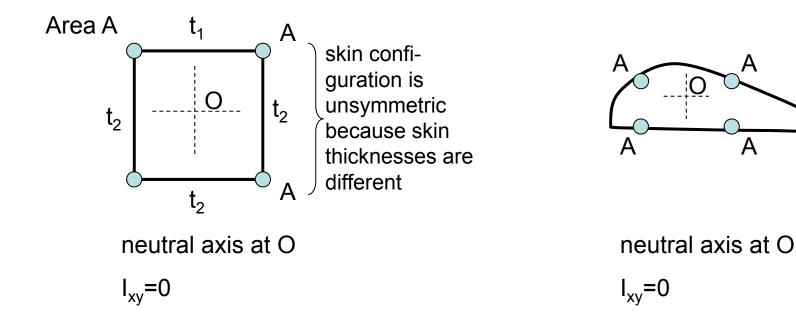
- Torsion of beams
  - stress functions
  - rate of twist
  - torsional rigidity
  - closed versus open sections
  - multi-cell beams
- Effect of taper
- Effect of cut-outs
- Castigliano's theorems (determination of deflections or forces as derivatives of strain energy)
- Buckling of beams
  - perfect
  - with initial imperfections or eccentricities
  - using energy methods

- if there is an axis of symmetry, the neutral axis is on it
- unsymmetric cross-sections have I<sub>xy</sub>≠0; there must be at least one axis of symmetry (for the entire cross-section or for the booms if idealized) for I<sub>xv</sub> to be zero
- if there is an axis of symmetry, the shear center is on it



• if the booms or flanges have one axis of symmetry, in a problem where skin carries only shear, then the neutral axis is on that axis of symmetry and  $I_{xy}$ =0 even if the skin configuration is unsymmetric

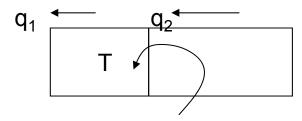
39



- a shear force acting through the shear center causes only bending and no twist
- in bending problems where the skin is idealized to carry only shear, the shear flow between booms/flanges is constant and equal to the average of the shear flow that would be obtained if the skin carried direct stresses also
- for basic beam bending

$$\sigma_z = \frac{My}{I_{xx}}$$

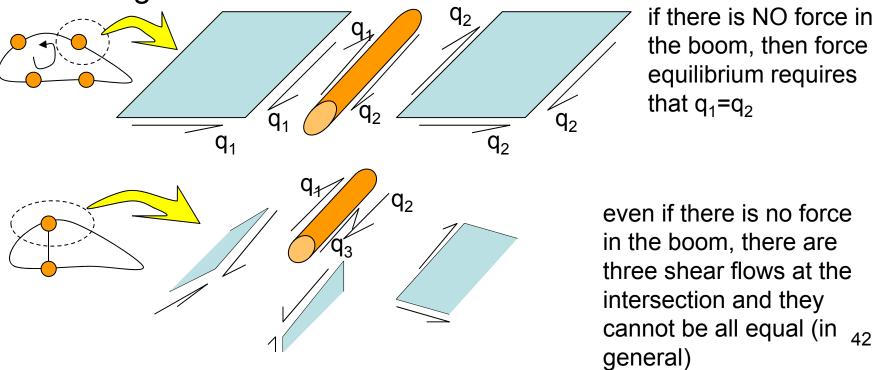
- for a constant shear flow through a skin: T = 2Aq
- there is an important difference between moment equilibrium and torque equivalence



the shear flows are **caused** by T; then, the moment caused by T about a point is equal to the moment caused by the shear flows about the same point (torque equivalence) this is not the same as adding all the moments about a point and setting them equal to 0

$$T=2A_1q_1+2A_2q_2$$
 and NOT:  $T+2A_1q_1+2A_2q_2=0$ 

- if a cross-section is under pure torsion, then for a single cell section, the booms do not change the shear flow; the shear flow is constant across booms
- for multi-cell beams under pure torsion, the shear flow changes across booms

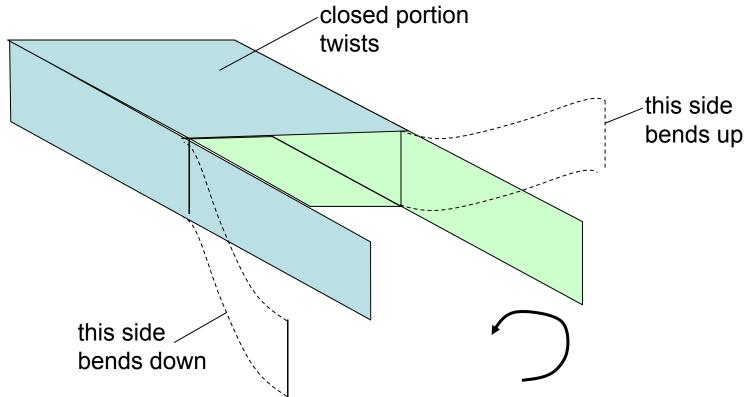


- for multi-cell beams, the rate of twist of each cell equals the rate of twist of the entire cross-section
- open sections are terrible in torsion; the shape of the open section has no effect on its torsional rigidity; it is always given by the thin rectangle approximation



same length same thickness for all => same J

 beams with cutouts take torsional loads at the cutout region through differential bending



- taper introduces additional (vertical) forces; the axial forces are still given by the standard equations
- cutouts are a pain to deal with
- Castigliano  $\Delta_i = \frac{\partial C_i}{\partial P_i} \qquad Q_i = \frac{\partial U}{\partial \Delta_i}$
- the buckling load of a beam fixed (clamped) at both ends is four times the buckling load of a beam simply supported (pinned) at both ends

$$P_{crit} = rac{\pi^2 EI}{L^2}$$
  $P_{cr} = rac{4\pi^2 EI}{L^2}$  simply-supported or pinned fixed or clamped