



Topological Phases in Quantum Systems with Quantum Group Symmetries

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Matrix Product States

- We analyse quantum systems on a 1D lattice.
- Consider states $|\psi\rangle$ of the form

$$|\psi\rangle = \text{tr}(\mathcal{B}^{[1]}\mathcal{B}^{[2]}\dots\mathcal{B}^{[L]})$$

where each $\mathcal{B}^{[i]}$ is a matrix with vector entries (rank-3 tensor) associated to site i .

$$|\psi\rangle = \text{tr} \left(\mathcal{B}^{[1]} \mathcal{B}^{[2]} \mathcal{B}^{[3]} \mathcal{B}^{[4]} \right)$$


Matrix Product States

Example: **spin-1 AKLT ground state** [Affleck, Kennedy et al. '87]

$$\mathcal{B}^{[i]} = \frac{1}{\sqrt{3}} \begin{bmatrix} -|0\rangle & \sqrt{2}|1\rangle \\ -\sqrt{2}|-1\rangle & |0\rangle \end{bmatrix} \quad i = 1, 2, \dots, L.$$

where $|1\rangle, |0\rangle, |-1\rangle \in V$, the one-site state space.

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where $|1\rangle, |0\rangle, |-1\rangle \in V$, the one-site state space.

$$|\psi_{AKLT}\rangle = \text{tr}(\mathcal{B}^{[1]}\mathcal{B}^{[2]}\dots\mathcal{B}^{[L]})$$

$$\stackrel{(L=2)}{=} \frac{2}{3}(\underbrace{|0\rangle \otimes |0\rangle}_{\text{entanglement}} - \underbrace{|1\rangle \otimes |-1\rangle}_{\text{entanglement}} + \underbrace{|-1\rangle \otimes |1\rangle}_{\text{entanglement}}).$$

Generalised AKLT states

- The underlying construction:

[Affleck, Kennedy et al. '87]



- Uses representation theory to assert $\mathfrak{su}(2)$ (rotational) symmetry:

$$g |\psi_{AKLT}\rangle = |\psi_{AKLT}\rangle \quad \forall g \in \mathfrak{su}(2).$$

- Choice of $\dim V \Rightarrow$ different AKLT states of 'spin' $S \in \mathbb{Z}^+$.

Generalised AKLT states

$$S = 1 : \quad \mathcal{B}_{(S=1)} = \frac{1}{\sqrt{3}} \begin{bmatrix} -|0\rangle & \sqrt{2}|1\rangle \\ -\sqrt{2}|-1\rangle & |0\rangle \end{bmatrix}$$

$$S = 2 : \quad \mathcal{B}_{(S=2)} = \frac{1}{\sqrt{10}} \begin{bmatrix} |0\rangle & -\sqrt{3}|1\rangle & \sqrt{6}|2\rangle \\ \sqrt{3}|-1\rangle & -2|0\rangle & \sqrt{3}|1\rangle \\ \sqrt{6}|-2\rangle & -\sqrt{3}|-1\rangle & |0\rangle \end{bmatrix}$$

\vdots

$$\text{General } S : \quad \mathcal{B}_S = (\text{some } (S+1) \times (S+1) \text{ matrix})$$

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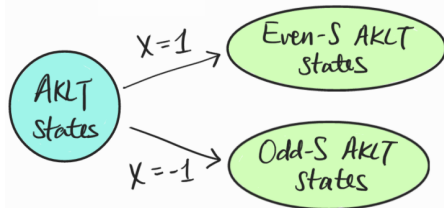
\vdots

$$\text{General } S : \quad \mathcal{B}_S = (\text{some } (S+1) \times (S+1) \text{ matrix})$$

Are all these states 'the same' in some physical sense?

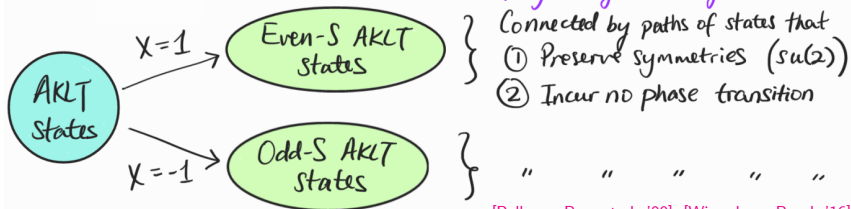
If not, how can we classify them?

Topological phases in AKLT states



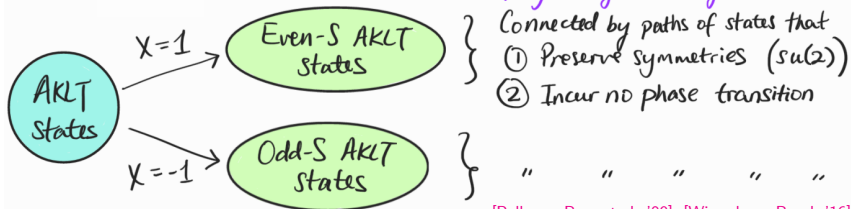
[Pollman, Berg et al. '09], [Wierschem, Beach '16]

Topological phases in AKLT states



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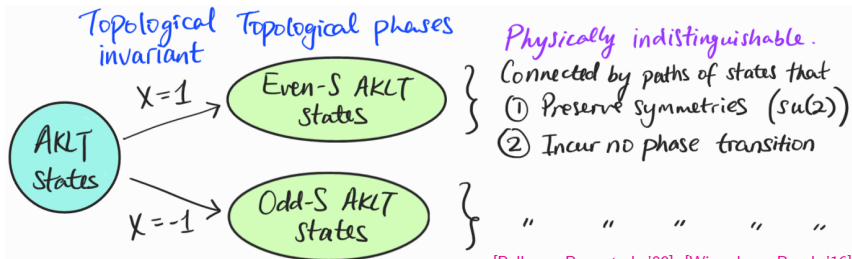
Topological phases in AKLT states



[Pollman, Berg et al. '09], [Wierschem, Beach '16]

- Use paths of tensors $\mathcal{B}_t, 0 \leq t \leq 1$.
- No phase transition \Leftrightarrow dominant eigenvalue of the **transfer matrix** $\sum_{\sigma} \mathcal{B}_t^{\sigma} \otimes \overline{\mathcal{B}_t^{\sigma}}$ is unique $\forall t$.

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Paths between even- S AKLT states

[Pollman, Berg et al. '09]

Connect every even- S AKLT state to a **trivial state**:



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Path from the $S = 2$ AKLT state to a trivial state:

Paths between even- S AKLT states

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Path from the $S = 2$ AKLT state to a trivial state:

$$\mathcal{B}_t = \left[\begin{array}{c|c} (1-t)A & t(1-t)B \\ \hline 0 & tC \end{array} \right] \left[\begin{array}{c|c} (1-t)A & 0 \\ \hline t(1-t)D & tE \end{array} \right]$$

where $A \cdots E$ are derived using rep. theory of $\mathfrak{su}(2)$.

Paths between even- S AKLT states

[Pollman, Berg et al. '09]

$$A = \frac{1}{\sqrt{10}} \begin{bmatrix} |0\rangle & -\sqrt{3}|1\rangle & \sqrt{6}|2\rangle \\ \sqrt{3}|-1\rangle & -2|0\rangle & \sqrt{3}|1\rangle \\ \sqrt{6}|-2\rangle & -\sqrt{3}|-1\rangle & |0\rangle \end{bmatrix}$$

$$B = \begin{bmatrix} -\frac{1}{\sqrt{3}}|-1\rangle & \frac{1}{\sqrt{2}}|0\rangle & -\frac{1}{\sqrt{2}}|1\rangle & \frac{1}{\sqrt{3}}|2\rangle & 0 \\ -\sqrt{\frac{2}{3}}|-2\rangle & \frac{1}{\sqrt{6}}|-1\rangle & 0 & -\frac{1}{\sqrt{6}}|1\rangle & \sqrt{\frac{2}{3}}|2\rangle \\ 0 & -\frac{1}{\sqrt{3}}|-2\rangle & \frac{1}{\sqrt{2}}|-1\rangle & -\frac{1}{\sqrt{2}}|0\rangle & \frac{1}{\sqrt{3}}|1\rangle \end{bmatrix}$$

$$C = \frac{1}{\sqrt{5}} \begin{bmatrix} |-2\rangle \\ -|-1\rangle \\ |0\rangle \\ -|1\rangle \\ |2\rangle \end{bmatrix}^T$$

$$D = \begin{bmatrix} \frac{1}{\sqrt{3}}|1\rangle & -\sqrt{\frac{2}{3}}|2\rangle & 0 \\ \frac{1}{\sqrt{2}}|0\rangle & -\frac{1}{\sqrt{6}}|1\rangle & -\frac{1}{\sqrt{3}}|2\rangle \\ \frac{1}{\sqrt{2}}|-1\rangle & 0 & -\frac{1}{\sqrt{2}}|1\rangle \\ \frac{1}{\sqrt{3}}|-2\rangle & \frac{1}{\sqrt{6}}|-1\rangle & -\frac{1}{\sqrt{2}}|0\rangle \\ 0 & \sqrt{\frac{2}{3}}|-2\rangle & -\frac{1}{\sqrt{3}}|-1\rangle \end{bmatrix}$$

$$E = \begin{bmatrix} |2\rangle \\ |1\rangle \\ |0\rangle \\ |-1\rangle \\ |-2\rangle \end{bmatrix}$$

Paths between even- S AKLT states

[Pollman, Berg et al. '09]

The path $\{\mathcal{B}_t\}_{t \in [0,1]}$ must satisfy the following:

- 1 $\mathfrak{su}(2)$ symmetry $\forall t$
- 2 No phase transition

Paths between even- S AKLT states

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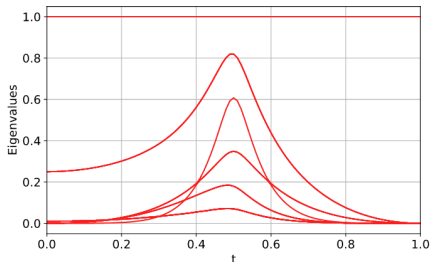
- ① $\mathfrak{su}(2)$ symmetry $\forall t$: by construction.
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Paths between even- S AKLT states

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The path $\{\mathcal{B}_t\}_{t \in [0,1]}$ must satisfy the following:

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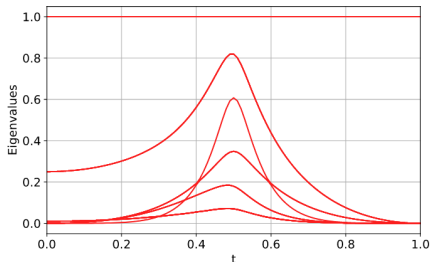
Similarly done for other even values of S .

Paths between even- S AKLT states

[Pollman, Berg et al. '09]

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- ② No phase transition: for $S = 2$,



Can we prove this is true for all even S ?

What it looks like for $S = 2$

```
transfer = ConstantArray[0, {16, 16}];
For[t = 1, t < 25, t++,
  kron = KroneckerProduct[NPSSM[t]], NPSSM[t]]; kron = Simplify[kron]; transfer = transfer + kron;
transfer = Simplify[transfer]; MatrixForm[transfer]
```

Out[17]: MatrixForm

$$\begin{pmatrix} \frac{1}{450} (-1+t)^2 (207-360t+685t^2-650t^3+325t^4) & 0 & 0 & 0 & 0 & \frac{1}{900} (-1+t)^2 (2 \\ 0 & \frac{1}{900} (-1+t)^2 (117-810t+1135t^2-650t^3+325t^4) & 0 & -\frac{(-1+t)^2 t (9-5t+5t^2)}{6\sqrt{30}} & 0 & 0 \\ 0 & 0 & \frac{1}{450} (-1+t)^2 (3-90t+265t^2-350t^3+175t^4) & 0 & 0 & 0 \\ 0 & \frac{(-1+t)^2 t (9-5t+5t^2)}{6\sqrt{30}} & 0 & \frac{t (9+5t+10t^2+5t^3)}{3\sqrt{5}} & 0 & 0 \\ 0 & 0 & 0 & 0 & \frac{1}{900} (-1+t)^2 (117-810t+1135t^2-650t^3+325t^4) & 0 \\ \frac{1}{900} (-1+t)^2 (297+90t+235t^2-650t^3+325t^4) & 0 & 0 & 0 & 0 & \frac{1}{450} (-1+t)^2 (1 \\ 0 & -\frac{1}{125} (-1+t)^2 (-27+135t-85t^2-100t^3+50t^4) & 0 & -\frac{(-1+t)^2 t (9-5t+5t^2)}{6\sqrt{30}} & 0 & 0 \\ -\frac{(-1+t)^2 t (9-5t+5t^2)}{6\sqrt{30}} & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & -\frac{1}{125} (-1+t)^2 (-27+135t-85t^2-100t^3+50t^4) & 0 \\ \frac{1}{900} (3-2t-10t^2+5t^3)^2 & 0 & 0 & 0 & 0 & \frac{1}{900} (-1+t)^2 (2 \\ 0 & 0 & 0 & 0 & -\frac{(-1+t)^2 t (9-5t+5t^2)}{6\sqrt{30}} & 0 \\ 0 & 0 & 0 & 0 & \frac{(-1+t)^2 t (9-5t+5t^2)}{6\sqrt{30}} & 0 \\ -\frac{(-1+t)^2 t (9-5t+5t^2)}{6\sqrt{30}} & 0 & 0 & 0 & 0 & 0 \\ 0 & -\frac{(-1+t)^2 t (9-5t+5t^2)}{6\sqrt{30}} & 0 & \frac{1}{3} \sqrt{5} (-1+t)^2 t^2 & 0 & 0 \\ \frac{1}{3} (-1+t)^2 t^2 & 0 & 0 & 0 & 0 & 0 \end{pmatrix}$$

What it looks like for $S = 2$

```
In[40]:= eigs = Simplify[Eigenvalues[transfer]]
```

$$\begin{aligned} & \frac{(-1+t)^2}{\sqrt{5}} - \frac{(-1+t)^2}{\sqrt{5}} - \frac{(-1+t)^2}{\sqrt{5}} + \frac{1}{300}(-1+t)^2(3-90t+265t^2-350t^3+175t^4) + \frac{1}{300}(-1+t)^2(3-90t+265t^2-350t^3+175t^4), \\ & \frac{1}{300}(-1+t)^2(3-90t+265t^2-350t^3+175t^4) + \frac{1}{300}(-1+t)^2(3-90t+265t^2-350t^3+175t^4), \\ & \frac{1}{300}(-1+t)^2(3-90t+265t^2-350t^3+175t^4) + \frac{1}{6}(3-6t+11t^2-20t^3+30t^4-20t^5+5t^6 - \\ & \quad \sqrt{9-36t+66t^2-180t^3+505t^4-920t^5+1330t^6-1700t^7+1810t^8-1400t^9+700t^{10}-200t^{11}+25t^{12}}) + \frac{1}{6}(3-6t+11t^2-20t^3+ \\ & \quad 30t^4-20t^5+5t^6+\sqrt{9-36t+66t^2-180t^3+505t^4-920t^5+1330t^6-1700t^7+1810t^8-1400t^9+700t^{10}-200t^{11}+25t^{12}}), \\ & \frac{1}{360}(45+36(-10+\sqrt{5})t+(880-156\sqrt{5})t^2+10(-91+24\sqrt{5})t^3-60(-7+2\sqrt{5})t^4-100t^5+25t^6 - \\ & \quad \sqrt{5}\sqrt{(-1+t)^2(405-162(35+4\sqrt{5})t-9(-3479+120\sqrt{5})t^2-72(1195+17\sqrt{5})t^3+(131705+11112\sqrt{5})t^4 - \\ & \quad 30(3889+308\sqrt{5})t^5+(58305-2520\sqrt{5})t^6+1600(-10+3\sqrt{5})t^7-75(-61+16\sqrt{5})t^8-750t^9+125t^{10}})), \\ & \frac{1}{360}(45+36(-10+\sqrt{5})t+(880-156\sqrt{5})t^2+10(-91+24\sqrt{5})t^3-60(-7+2\sqrt{5})t^4-100t^5+25t^6 - \\ & \quad \sqrt{5}\sqrt{(-1+t)^2(405-162(35+4\sqrt{5})t-9(-3479+120\sqrt{5})t^2-72(1195+17\sqrt{5})t^3+(131705+11112\sqrt{5})t^4 - \\ & \quad 30(3889+308\sqrt{5})t^5+(58305-2520\sqrt{5})t^6+1600(-10+3\sqrt{5})t^7-75(-61+16\sqrt{5})t^8-750t^9+125t^{10}})), \\ & \frac{1}{360}(45+36(-10+\sqrt{5})t+(880-156\sqrt{5})t^2+10(-91+24\sqrt{5})t^3-60(-7+2\sqrt{5})t^4-100t^5+25t^6 - \\ & \quad \sqrt{5}\sqrt{(-1+t)^2(405-162(35+4\sqrt{5})t-9(-3479+120\sqrt{5})t^2-72(1195+17\sqrt{5})t^3+(131705+11112\sqrt{5})t^4 - \\ & \quad 30(3889+308\sqrt{5})t^5+(58305-2520\sqrt{5})t^6+1600(-10+3\sqrt{5})t^7-75(-61+16\sqrt{5})t^8-750t^9+125t^{10}})), \\ & \frac{1}{360}(45+36(-10+\sqrt{5})t+(880-156\sqrt{5})t^2+10(-91+24\sqrt{5})t^3-60(-7+2\sqrt{5})t^4-100t^5+25t^6 - \\ & \quad \sqrt{5}\sqrt{(-1+t)^2(405-162(35+4\sqrt{5})t-9(-3479+120\sqrt{5})t^2-72(1195+17\sqrt{5})t^3+(131705+11112\sqrt{5})t^4 - \\ & \quad 30(3889+308\sqrt{5})t^5+(58305-2520\sqrt{5})t^6+1600(-10+3\sqrt{5})t^7-75(-61+16\sqrt{5})t^8-750t^9+125t^{10}})), \\ & \frac{1}{360}(45+36(-10+\sqrt{5})t+(880-156\sqrt{5})t^2+10(-91+24\sqrt{5})t^3-60(-7+2\sqrt{5})t^4-100t^5+25t^6 - \\ & \quad \sqrt{5}\sqrt{(-1+t)^2(405-162(35+4\sqrt{5})t-9(-3479+120\sqrt{5})t^2-72(1195+17\sqrt{5})t^3+(131705+11112\sqrt{5})t^4 - \\ & \quad 30(3889+308\sqrt{5})t^5+(58305-2520\sqrt{5})t^6+1600(-10+3\sqrt{5})t^7-75(-61+16\sqrt{5})t^8-750t^9+125t^{10}})), \\ & \frac{1}{360}(45+36(-10+\sqrt{5})t+(880-156\sqrt{5})t^2+10(-91+24\sqrt{5})t^3-60(-7+2\sqrt{5})t^4-100t^5+25t^6 - \\ & \quad \sqrt{5}\sqrt{(-1+t)^2(405-162(35+4\sqrt{5})t-9(-3479+120\sqrt{5})t^2-72(1195+17\sqrt{5})t^3+(131705+11112\sqrt{5})t^4 - \\ & \quad 30(3889+308\sqrt{5})t^5+(58305-2520\sqrt{5})t^6+1600(-10+3\sqrt{5})t^7-75(-61+16\sqrt{5})t^8-750t^9+125t^{10}})) \end{aligned}$$

Some semblance of success

Want exact forms of e-values of all transfer matrices

$$E_t := \sum_{\sigma} \mathcal{B}_t^{\sigma} \otimes \overline{\mathcal{B}_t^{\sigma}}, \quad t \in [0, 1].$$

- ① Recognise that TM can be block-diagonalised by permuting rows/columns
- ② Appeal to rep. theory of $\mathfrak{su}(2)$ and linear algebra to tame the problem (Schur's Lemma, C-G coefficients, Perron-Frobenius theorem...)

$\{\text{eigenvalues of TM}\} = \{\text{eigenvalues of largest block}\}!$

q -deformed AKLT states

- By modifying one of the relations defining $\mathfrak{su}(2)$,

$$S^+S^- - S^-S^+ = [2S^z]_q = \frac{q^{2S^z} - q^{-2S^z}}{q - q^{-1}}$$

we obtain the **quantum group** $\mathcal{U}_q[sl(2)]$.

- Construct the AKLT states exactly the same way.
- Do these q -deformed AKLT states also admit a binary classification?
- Numerical evidence: yes, at least for $q \in \mathbb{R}^+$.

Summary

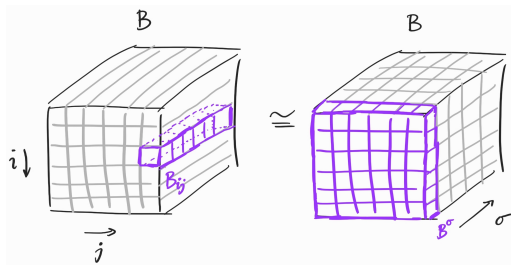
- Despite being constructed identically, the AKLT states exhibit different physical properties for different S .
- Existing numerical evidence supports a classification of AKLT states into 2 distinct phases of matter, depending on the parity of S .
- Our analytic results give way to a formal proof.
- In related news, our numerical evidence supports the same classification of q AKLT states.

More on the transfer matrix

The **transfer matrix** is defined by $E = \sum_{\sigma} \mathcal{B}^{\sigma} \otimes \overline{\mathcal{B}}^{\sigma}$.

- View the object \mathcal{B} as a tuple of (scalar) matrices

$$\mathcal{B} := (\mathcal{B}^{-S}, \mathcal{B}^{-S+1}, \dots, \mathcal{B}^S).$$



More on the transfer matrix

- The symbol \otimes denotes the Kronecker product

$$A \otimes B = \begin{bmatrix} a_{11}B & a_{12}B & \cdots & a_{1n}B \\ a_{21}B & a_{22}B & \cdots & a_{2n}B \\ \vdots & \vdots & \ddots & \vdots \\ a_{n1}B & a_{n2}B & \cdots & a_{nn}B \end{bmatrix}.$$

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