Renormalization of Symmetry Improved 3PIEA gap equations at 2 loops

Supplement to "Symmetry improvement of 3PI effective actions for O(N) scalar field theory" by Michael J. Brown and Ian B. Whittingham.

Mathematica notebook to compute couter-terms for two loop truncations of the effective action as described in Section IV of the paper.

Hartree-Fock

ClearAll[geom, neom, intrules, msbarrules, mg2soln, cteq, cts, δm , $\delta \lambda$];

Hartree-Fock gap equations with counterterms

Goldstone equation of motion. Quantities in reference to the paper are:

p is the four-momentum flowing through the propagators Δ_{G}^{-1} and $\Delta_{\text{N}}^{-1},$

mg2 is the Goldstone mass squared $m_{\rm G}^2$,

Z and $Z\Delta$ are the wavefunction a propagator renormalization constants,

 m^2 is the (renormalized) Lagrangian mass parameter, δm_1^2 is its counter-term,

 λ is the (renormalized) four point coupling,

 $\delta\lambda_{1a}$, $\delta\lambda_{2a}$, $\delta\lambda_{2b}$ are the independent coupling counter-terms,

v is the scalar field vacuum expectation value,

ħ is the reduced Planck constant,

n is the number of fields in the O(n) symmetry group,

t∞g, t∞n are the divergent tadpole integrals for the Goldstone, Higgs resp.,

tfing, tfinn are the finite parts of the tadpoles for the Goldstone, Higgs resp.

$$\begin{split} \text{geom} &= p^2 - mg2 = \text{Z} \; \text{Z} \Delta \; p^2 - m^2 - \delta m_1^2 - \text{Z} \Delta \; \frac{\lambda + \delta \lambda_{1\,a}}{6} \; \mathbf{v}^2 - \\ & \frac{\hbar}{6} \; \left(\left(\mathbf{n} + \mathbf{1} \right) \; \lambda + \left(\mathbf{n} - \mathbf{1} \right) \; \delta \lambda_{2\,a} + 2 \; \delta \lambda_{2\,b} \right) \; \text{Z} \Delta^2 \; \left(\text{t} \infty \mathbf{g} + \text{tfing} \right) - \frac{\hbar}{6} \; \left(\lambda + \delta \lambda_{2\,a} \right) \; \text{Z} \Delta^2 \; \left(\text{t} \infty \mathbf{n} + \text{tfinn} \right) \\ & - mg2 + p^2 = - m^2 + p^2 \; \text{Z} \; \text{Z} \Delta - \delta m_1^2 - \frac{1}{6} \; \mathbf{v}^2 \; \text{Z} \Delta \; \left(\lambda + \delta \lambda_a \right) - \\ & \frac{1}{6} \; \left(\text{tfinn} + \text{t} \infty \mathbf{n} \right) \; \text{Z} \Delta^2 \; \hbar \; \left(\lambda + \delta \lambda_{2\,a} \right) - \frac{1}{6} \; \left(\text{tfing} + \text{t} \infty \mathbf{g} \right) \; \text{Z} \Delta^2 \; \hbar \; \left((1 + \mathbf{n}) \; \lambda + \left(-1 + \mathbf{n} \right) \; \delta \lambda_{2\,a} + 2 \; \delta \lambda_{2\,b} \right) \end{split}$$

Higgs equation of motion

neom =
$$p^2 - mn2 = \frac{-\lambda v^2}{3} Z\Delta + p^2 - mg2$$

-mn2 + $p^2 = -mg2 + p^2 - \frac{1}{3} v^2 Z\Delta \lambda$

Infinite parts of tadpoles in MSbar

MSbar rules for 4 - 2 ϵ dimensions

$$\begin{split} &\text{msbarrules} = \\ &\left\{ \text{t} \infty \text{g} \rightarrow \frac{\text{-mg2}}{16\,\pi^2} \left(\frac{1}{\epsilon} - \text{EulerGamma} + 1 + \text{Log}[4\,\pi] \right), \ \text{t} \infty \text{n} \rightarrow \frac{\text{-mn2}}{16\,\pi^2} \left(\frac{1}{\epsilon} - \text{EulerGamma} + 1 + \text{Log}[4\,\pi] \right) \right\} \\ &\left\{ \text{t} \infty \text{g} \rightarrow -\frac{\text{mg2}\left(1 - \text{EulerGamma} + \frac{1}{\epsilon} + \text{Log}[4\,\pi] \right)}{16\,\pi^2}, \ \text{t} \infty \text{n} \rightarrow -\frac{\text{mn2}\left(1 - \text{EulerGamma} + \frac{1}{\epsilon} + \text{Log}[4\,\pi] \right)}{16\,\pi^2} \right\} \end{split}$$

Sub in tadpole expressions, eliminate mn2 and solve for mg2

mg2soln = mg2 /. (geom /. msbarrules /. Solve[neom, mn2][[1]] // Solve[#, mg2][[1]] &)
$$\left(-m^2 - p^2 + p^2 \ Z \ Z\Delta - \delta m_1^2 - \frac{1}{6} \ v^2 \ Z\Delta \ (\lambda + \delta \lambda_a) - \frac{1}{6} \ tfinn \ Z\Delta^2 \ \hbar \ (\lambda + \delta \lambda_{2\,a}) + \frac{v^2 \ Z\Delta^3 \ \lambda \ \hbar \ \left(1 - EulerGamma + \frac{1}{\epsilon} + Log[4\ \pi] \right) \ (\lambda + \delta \lambda_{2\,a})}{288\ \pi^2} - \frac{1}{6} \ tfing \ Z\Delta^2 \ \hbar \ \left((1+n) \ \lambda + (-1+n) \ \delta \lambda_{2\,a} + 2 \ \delta \lambda_{2\,b} \right) \right) /$$

$$\left(-1 - \frac{Z\Delta^2 \ \hbar \ \left(1 - EulerGamma + \frac{1}{\epsilon} + Log[4\ \pi] \right) \ (\lambda + \delta \lambda_{2\,a})}{96\ \pi^2} - \frac{1}{96\ \pi^2} \right)$$

$$Z\Delta^2 \ \hbar \ \left(1 - EulerGamma + \frac{1}{\epsilon} + Log[4\ \pi] \right) \ ((1+n) \ \lambda + (-1+n) \ \delta \lambda_{2\,a} + 2 \ \delta \lambda_{2\,b}) \right)$$

Gather divergences proportional v, tfing and tfinn and set independently to zero

First we subtract the finite equation of motion, then gather coefficients of the remainder into a list and set each to zero (after some trimming and simplifying).

cteq =

$$\left(\left(\text{CoefficientList} \left[\text{mg2soln} + \left(-\text{m}^2 - \frac{\lambda}{6} \, \text{v}^2 - \frac{\hbar}{6} \, \left(\left(\text{n} + 1 \right) \, \lambda \right) \, \left(\text{tfing} \right) - \frac{\hbar}{6} \, \left(\lambda \right) \, \left(\text{tfinn} \right) \right), \, \left\{ \text{p, v, tfing, tfinn} \right\} \right] / / \, \text{Flatten} \right) / /$$

DeleteDuplicates // Simplify // FullSimplify == 0 // Thread

$$\left\{ \left(96 \, \delta m_1^2 \right) \middle/ \left(96 + \frac{1}{\pi^2} Z \Delta^2 \, \hbar \left(1 - \text{EulerGamma} + \frac{1}{\epsilon} + \text{Log}[4] + \text{Log}[\pi] \right) \, \left(\lambda + \delta \lambda_{2 \, a} \right) + \frac{1}{\pi^2} Z \Delta^2 \, \hbar \left(1 - \text{EulerGamma} + \frac{1}{\epsilon} + \text{Log}[4] + \text{Log}[\pi] \right) \, \left(\lambda + n \, \lambda + (-1 + n) \, \delta \lambda_{2 \, a} + 2 \, \delta \lambda_{2 \, b} \right) \right) + \\ m^2 \left(-1 + 96 \middle/ \left(96 + \frac{1}{\pi^2} Z \Delta^2 \, \hbar \left(1 - \text{EulerGamma} + \frac{1}{\epsilon} + \text{Log}[4] + \text{Log}[\pi] \right) \, \left(\lambda + \delta \lambda_{2 \, a} \right) + \frac{1}{\pi^2} Z \Delta^2 \, \hbar \right) \\ \left(1 - \text{EulerGamma} + \frac{1}{\epsilon} + \text{Log}[4] + \text{Log}[\pi] \right) \, \left(\lambda + n \, \lambda + (-1 + n) \, \delta \lambda_{2 \, a} + 2 \, \delta \lambda_{2 \, b} \right) \right) = 0 \, , \, \frac{1}{6} \, \hbar \right) \\ \left(-\lambda + \left(96 \, Z \Delta^2 \, \left(\lambda + \delta \lambda_{2 \, a} \right) \right) \middle/ \left(96 + \frac{1}{\pi^2} Z \Delta^2 \, \hbar \left(1 - \text{EulerGamma} + \frac{1}{\epsilon} + \text{Log}[4] + \text{Log}[\pi] \right) \, \left(\lambda + n \, \lambda + (-1 + n) \, \delta \lambda_{2 \, a} + 2 \, \delta \lambda_{2 \, b} \right) \right) \right) = 0 \, , \\ \frac{1}{\pi^2} Z \Delta^2 \, \hbar \left(1 - \text{EulerGamma} + \frac{1}{\epsilon} + \text{Log}[4] + \text{Log}[\pi] \right) \, \left(\lambda + n \, \lambda + (-1 + n) \, \delta \lambda_{2 \, a} + 2 \, \delta \lambda_{2 \, b} \right) \right) \right) = 0 \, , \\ \left(96 + \frac{1}{\pi^2} Z \Delta^2 \, \hbar \left(1 - \text{EulerGamma} + \frac{1}{\epsilon} + \text{Log}[4] + \text{Log}[\pi] \right) \, \left(\lambda + \delta \lambda_{2 \, a} \right) + \frac{1}{\pi^2} Z \Delta^2 \, \hbar \left(1 - \text{EulerGamma} + \frac{1}{\epsilon} + \text{Log}[4] + \text{Log}[\pi] \right) \, \left(\lambda + \delta \lambda_{2 \, a} \right) + \frac{1}{\pi^2} Z \Delta^2 \, \hbar \left(1 - \text{EulerGamma} + \frac{1}{\epsilon} + \text{Log}[4] + \text{Log}[\pi] \right) \, \left(\lambda + n \, \lambda + (-1 + n) \, \delta \lambda_{2 \, a} + 2 \, \delta \lambda_{2 \, b} \right) \right) \right) = 0 \, , \\ \left(2 \, Z \Delta^3 \, \lambda \, \hbar \left(1 - \text{EulerGamma} + \frac{1}{\epsilon} + \text{Log}[4] + \text{Log}[\pi] \right) \, \left(\lambda + n \, \lambda + (-1 + n) \, \delta \lambda_{2 \, a} + 2 \, \delta \lambda_{2 \, b} \right) \right) - \\ \left(2 \, Z \Delta^3 \, \lambda \, \hbar \left(1 - \text{EulerGamma} + \frac{1}{\epsilon} + \text{Log}[4] + \text{Log}[\pi] \right) \, \left(\lambda + n \, \lambda + (-1 + n) \, \delta \lambda_{2 \, a} + 2 \, \delta \lambda_{2 \, b} \right) \right) - \\ \left(2 \, Z \Delta^3 \, \lambda \, \hbar \left(1 - \text{EulerGamma} + \frac{1}{\epsilon} + \text{Log}[4] + \text{Log}[\pi] \right) \, \left(\lambda + n \, \lambda + (-1 + n) \, \delta \lambda_{2 \, a} + 2 \, \delta \lambda_{2 \, b} \right) \right) - \\ \left(1 - \text{EulerGamma} + \frac{1}{\epsilon} + \text{Log}[4] + \text{Log}[\pi] \right) \, \left(\lambda + n \, \lambda + (-1 + n) \, \delta \lambda_{2 \, a} + 2 \, \delta \lambda_{2 \, b} \right) \right) \right) = 0 \, , \\ - \left(\left(96 \, \pi^2 \, (-1 + 2 \, Z \Delta^2 \, \hbar \, \left(1 - \text{EulerGamma} + \frac{1}{\epsilon} + \text{Log}[4] + \text{Log}[\pi] \right) \, \left(\lambda + n \, \lambda + (-1 + n) \, \delta \lambda_{2 \, a}$$

Solve for counterterms

 $cts = \left\{ \delta m_1^2, \, \delta \lambda_{1a}, \, \delta \lambda_{2a}, \, \delta \lambda_{2b}, \, Z, \, Z\Delta \right\} /. \, \, Solve[cteq, \, \left\{ \delta m_1, \, \delta \lambda_{1a}, \, \delta \lambda_{2a}, \, \delta \lambda_{2b}, \, Z, \, Z_\Delta \right\}] \, // \, \, description + 1 \, \, descri$ FullSimplify // DeleteDuplicates

Solve::svars: Equations may not give solutions for all "solve" variables. >>>

$$\left\{ \left\{ \frac{m^2}{-1 + \frac{96 \, \pi^2 \, \epsilon}{(2+n) \, \lambda \, \hbar \, (1+\epsilon \, (1- \mathrm{EulerGamma} + \mathrm{Log} \, [4] + \mathrm{Log} \, [\pi]))}}, \right. \right. \\ \left. - \left(\left(\lambda \, \left(96 \, \pi^2 \, \left(-1 + \mathrm{Z} \Delta \right) \, \epsilon + \, (4+n) \, \mathrm{Z} \Delta \, \lambda \, \hbar \, \left(-1 + \epsilon \, \left(-1 + \mathrm{EulerGamma} - \mathrm{Log} \, [4 \, \pi] \right) \right) \right) \right) \right) \right. \\ \left. \left. \left(\mathrm{Z} \Delta \, \left(96 \, \pi^2 \, \epsilon + \, (2+n) \, \lambda \, \hbar \, \left(-1 + \epsilon \, \left(-1 + \mathrm{EulerGamma} - \mathrm{Log} \, [4 \, \pi] \right) \right) \right) \right) \right) \right. \\ \left. - \left(\left(\lambda \, \left(96 \, \pi^2 \, \left(-1 + \mathrm{Z} \Delta^2 \right) \, \epsilon + \, (2+n) \, \mathrm{Z} \Delta^2 \, \lambda \, \hbar \, \left(-1 + \epsilon \, \left(-1 + \mathrm{EulerGamma} - \mathrm{Log} \, [4 \, \pi] \right) \right) \right) \right) \right) \right. \\ \left. - \left(\left(\lambda \, \left(96 \, \pi^2 \, \left(-1 + \mathrm{Z} \Delta^2 \right) \, \epsilon + \, (2+n) \, \lambda \, \hbar \, \left(-1 + \epsilon \, \left(-1 + \mathrm{EulerGamma} - \mathrm{Log} \, [4 \, \pi] \right) \right) \right) \right) \right) \right. \\ \left. \left. \left(2\Delta^2 \, \left(96 \, \pi^2 \, \left(-1 + \mathrm{Z} \Delta^2 \right) \, \epsilon + \, (2+n) \, \lambda \, \hbar \, \left(-1 + \epsilon \, \left(-1 + \mathrm{EulerGamma} - \mathrm{Log} \, [4 \, \pi] \right) \right) \right) \right) \right) \right. \\ \left. \left. \left(\mathrm{Z} \Delta^2 \, \left(96 \, \pi^2 \, \epsilon + \, (2+n) \, \lambda \, \hbar \, \left(-1 + \epsilon \, \left(-1 + \mathrm{EulerGamma} - \mathrm{Log} \, [4 \, \pi] \right) \right) \right) \right) \right) \right. \right. \\ \left. \left. \left(\mathrm{Z} \Delta^2 \, \left(96 \, \pi^2 \, \epsilon + \, (2+n) \, \lambda \, \hbar \, \left(-1 + \epsilon \, \left(-1 + \mathrm{EulerGamma} - \mathrm{Log} \, [4 \, \pi] \right) \right) \right) \right) \right) \right. \right. \\ \left. \left. \left(\mathrm{Z} \Delta^2 \, \left(96 \, \pi^2 \, \epsilon + \, (2+n) \, \lambda \, \hbar \, \left(-1 + \epsilon \, \left(-1 + \mathrm{EulerGamma} - \mathrm{Log} \, [4 \, \pi] \right) \right) \right) \right) \right) \right. \right. \\ \left. \left. \left(\mathrm{Z} \Delta^2 \, \left(96 \, \pi^2 \, \epsilon + \, (2+n) \, \lambda \, \hbar \, \left(-1 + \epsilon \, \left(-1 + \mathrm{EulerGamma} - \mathrm{Log} \, [4 \, \pi] \right) \right) \right) \right) \right) \right. \right. \right. \\ \left. \left. \left. \left(\mathrm{Z} \Delta^2 \, \left(96 \, \pi^2 \, \epsilon + \, (2+n) \, \lambda \, \hbar \, \left(-1 + \epsilon \, \left(-1 + \mathrm{EulerGamma} - \mathrm{Log} \, [4 \, \pi] \right) \right) \right) \right) \right) \right. \right. \right. \right. \\ \left. \left. \left(\mathrm{Z} \Delta^2 \, \left(96 \, \pi^2 \, \epsilon + \, (2+n) \, \lambda \, \hbar \, \left(-1 + \epsilon \, \left(-1 + \mathrm{EulerGamma} - \mathrm{Log} \, [4 \, \pi] \right) \right) \right) \right) \right) \right. \right. \\ \left. \left. \left(\mathrm{Z} \Delta^2 \, \left(96 \, \pi^2 \, \epsilon + \, (2+n) \, \lambda \, \hbar \, \left(-1 + \epsilon \, \left(-1 + \mathrm{EulerGamma} - \mathrm{Log} \, [4 \, \pi] \right) \right) \right) \right) \right) \right. \right. \right. \right. \right. \\ \left. \left. \left(\mathrm{Z} \Delta^2 \, \left(96 \, \pi^2 \, \epsilon + \, (2+n) \, \lambda \, \hbar \, \left(-1 + \epsilon \, \left(-1 + \mathrm{EulerGamma} - \mathrm{Log} \, [4 \, \pi] \right) \right) \right) \right) \right) \right. \right. \\ \left. \left. \left(\mathrm{Z} \Delta^2 \, \left(96 \, \pi^2 \, \epsilon + \, \left(2 + n \right) \, \lambda \, \hbar \, \left(-1 + \epsilon \, \left(-1 + \mathrm{EulerGamma} - \mathrm{EulerGamma} \right) \right)$$

 $Z\Delta$ is redundant in this truncation, can remove it :

cts /. $Z\Delta \rightarrow 1$ // FullSimplify

$$\left\{ \left\{ \frac{m^2}{-1 + \frac{96\,\pi^2\,\varepsilon}{(2+\mathrm{n})\,\lambda\,\hbar\,\,(1+\varepsilon\,\,(1-\mathrm{EulerGamma+Log}[4]+\mathrm{Log}[\pi]))}}, \frac{(4+\mathrm{n})\,\lambda^2}{-(2+\mathrm{n})\,\lambda + \frac{96\,\pi^2\,\varepsilon}{\hbar\,\,(1+\varepsilon\,\,(1-\mathrm{EulerGamma+Log}[4]+\mathrm{Log}[\pi]))}}, \frac{\lambda^2}{-\lambda + \frac{96\,\pi^2\,\varepsilon}{(2+\mathrm{n})\,\hbar\,\,(1+\varepsilon\,\,(1-\mathrm{EulerGamma+Log}[4]+\mathrm{Log}[\pi]))}}, \frac{\lambda^2}{-\lambda + \frac{96\,\pi^2\,\varepsilon}{(2+\mathrm{n})\,\hbar\,\,(1+\varepsilon\,\,(1-\mathrm{EulerGamma+Log}[4]+\mathrm{Log}[\pi]))}}, \frac{1}{\lambda^2} \right\}$$

Series[cts, $\{\epsilon, 0, 1\}$] /. $\{Z\Delta \rightarrow 1\}$

$$\left\{ \left\{ -m^2 - \frac{96 \left(m^2 \pi^2\right) \epsilon}{(2+n) \lambda \hbar} + O[\epsilon]^2, -\frac{(4+n) \lambda}{2+n} - \frac{96 \left(4 \pi^2 + n \pi^2\right) \epsilon}{(2+n)^2 \hbar} + O[\epsilon]^2, -\lambda - \frac{96 \pi^2 \epsilon}{(2+n) \hbar} + O[\epsilon]^2, 1, 1 \right\} \right\}$$

Sunset

NOTE: this uses some of the same variable names as the Hartree-Fock code! Be careful not to clobber what you need to keep.

ClearAll[geom, neom, intrules, msbarrules, mg2soln, cteq, δ m, $\delta\lambda$, $\delta\lambda$, $\delta\lambda$, $\delta\lambda$);

Equations of motion

Additional variables relative to the Hartree-Fock case:

Ing is the sunset integral $I_{NG}(p)$

Ifingp is the finite sunset integral $I_{NG}^{fin}(p)$,

Ifing 0 is $I_{NG}^{fin}(m_G)$,

Ifingn is $I_{NG}^{fin}(m_N)$,

 $\delta\lambda$ is the sunset graph coupling counter-term, $I\mu$, $t\mu$ and $c\mu$ are the auxiliary integrals I_{μ} , T_{μ} and c_{μ} respectively.

$$\begin{split} &\text{geom} = \mathbf{p}^2 - \text{mg2} + \mathbf{i}\,\hbar \, \left(\frac{(\lambda)\,\mathbf{v}}{3}\right)^2 \, \left(\text{Ifingp-Ifing0}\right) = \\ &\mathbf{z}\,\mathbf{z}\Delta\,\mathbf{p}^2 - \mathbf{m}^2 - \delta\mathbf{m}_1^2 - \mathbf{z}\Delta\,\,\frac{\lambda + \delta\lambda_{1\,a}}{6}\,\,\mathbf{v}^2 - \frac{\hbar}{6}\, \left(\left(\mathbf{n} + \mathbf{1}\right)\,\lambda + \left(\mathbf{n} - \mathbf{1}\right)\,\delta\lambda_{2\,a} + 2\,\delta\lambda_{2\,b}\right)\,\mathbf{z}\Delta^2 \, \left(\text{tg}\right) - \\ &\frac{\hbar}{6}\, \left(\lambda + \delta\lambda_{2\,a}\right)\,\mathbf{z}\Delta^2 \, \left(\text{tn}\right) + \mathbf{i}\,\hbar \, \left(\frac{(\lambda + \delta\lambda)\,\mathbf{v}}{3}\right)^2\,\mathbf{z}\Delta^3 \, \text{Ing} \\ &- \text{mg2} + \mathbf{p}^2 + \frac{1}{9}\,\,\mathbf{i} \, \left(-\text{Ifing0} + \text{Ifingp}\right)\,\,\mathbf{v}^2\,\lambda^2\,\hbar = -\mathbf{m}^2 + \mathbf{p}^2\,\mathbf{z}\,\,\mathbf{z}\Delta + \frac{1}{9}\,\,\mathbf{i}\,\,\text{Ing}\,\,\mathbf{v}^2\,\,\mathbf{z}\Delta^3 \, \left(\delta\lambda + \lambda\right)^2\,\hbar - \delta\mathbf{m}_1^2 - \frac{1}{6}\,\,\mathbf{v}^2\,\,\mathbf{z}\Delta \, \left(\lambda + \delta\lambda_a\right) - \frac{1}{6}\,\,\text{tn}\,\,\mathbf{z}\Delta^2\,\hbar \, \left(\lambda + \delta\lambda_{2\,a}\right) - \frac{1}{6}\,\,\text{tg}\,\,\mathbf{z}\Delta^2\,\hbar \, \left(\left(1 + \mathbf{n}\right)\,\lambda + \left(-1 + \mathbf{n}\right)\,\,\delta\lambda_{2\,a} + 2\,\,\delta\lambda_{2\,b}\right) \\ &\mathbf{neom} = \mathbf{p}^2 - \mathbf{mn2} + \mathbf{i}\,\hbar \, \left(\frac{(\lambda)\,\,\mathbf{v}}{3}\right)^2 \, \left(\mathbf{Ifingp-Ifingn}\right) = \\ &\frac{-\mathbf{z}\Delta\, \left(\lambda + \delta\lambda\right)\,\,\mathbf{v}^2}{3} + \mathbf{p}^2 - \mathbf{mg2} + \mathbf{i}\,\hbar \, \left(\frac{(\lambda)\,\,\mathbf{v}}{3}\right)^2 \, \left(\mathbf{Ifingp-Ifing0}\right) \\ &- \mathbf{mn2} + \mathbf{p}^2 + \frac{1}{9}\,\,\mathbf{i}\, \left(-\mathbf{Ifingn} + \mathbf{Ifingp}\right)\,\,\mathbf{v}^2\,\lambda^2\,\hbar = \\ &- \mathbf{mg2} + \mathbf{p}^2 - \frac{1}{3}\,\,\mathbf{v}^2\,\,\mathbf{z}\Delta\, \left(\delta\lambda + \lambda\right) + \frac{1}{9}\,\,\mathbf{i}\, \left(-\mathbf{Ifing0} + \mathbf{Ifingp}\right)\,\,\mathbf{v}^2\,\lambda^2\,\hbar \end{split}$$

Divergent parts subtracted with auxiliary integrals and MSbar

intrules =
$$\left\{\operatorname{Ing} \to \operatorname{I}\mu + \operatorname{Ifingp} + \operatorname{Ifing0}, \right\}$$

 $\operatorname{tg} \to \operatorname{t}\mu - \operatorname{ii} \left(\operatorname{mg2} - \mu^2\right) \operatorname{I}\mu + \operatorname{\hbar} \left(\frac{(\lambda + \delta\lambda) \operatorname{v}}{3}\right)^2 \operatorname{c}\mu + \operatorname{tfing}, \right\}$
 $\operatorname{tn} \to \operatorname{t}\mu - \operatorname{ii} \left(\operatorname{mn2} - \mu^2\right) \operatorname{I}\mu + \operatorname{\hbar} \left(\frac{(\lambda + \delta\lambda) \operatorname{v}}{3}\right)^2 \operatorname{c}\mu + \operatorname{tfinn} \right\}$
 $\left\{\operatorname{Ing} \to \operatorname{Ifing0} + \operatorname{Ifingp} + \operatorname{I}\mu, \operatorname{tg} \to \operatorname{tfing} + \operatorname{t}\mu - \operatorname{ii} \operatorname{I}\mu \left(\operatorname{mg2} - \mu^2\right) + \frac{1}{9} \operatorname{c}\mu \operatorname{v}^2 \left(\delta\lambda + \lambda\right)^2 \operatorname{\hbar}, \right\}$
 $\operatorname{tn} \to \operatorname{tfinn} + \operatorname{t}\mu - \operatorname{ii} \operatorname{I}\mu \left(\operatorname{mn2} - \mu^2\right) + \frac{1}{9} \operatorname{c}\mu \operatorname{v}^2 \left(\delta\lambda + \lambda\right)^2 \operatorname{\hbar} \right\}$
 $\operatorname{msbarrules} = \left\{\operatorname{I}\mu \to \frac{-\operatorname{ii}}{16\pi^2} \left(\frac{1}{\epsilon} - \operatorname{EulerGamma} + \operatorname{1} + \operatorname{Log}\left[4\pi\right]\right), \operatorname{c}\mu \to \left(\frac{\operatorname{a0}}{\epsilon^2} + \frac{\operatorname{a1}}{\epsilon} + \operatorname{a2}\right)\right\}$
 $\left\{\operatorname{I}\mu \to -\frac{\operatorname{ii} \left(-\operatorname{EulerGamma} + \frac{1}{\epsilon} + \operatorname{Log}\left[4\pi\right]\right)}{16\pi^2}, \operatorname{c}\mu \to \operatorname{a2} + \frac{\operatorname{a0}}{\epsilon^2} + \frac{\operatorname{a1}}{\epsilon}\right\}$

Sub everything in, eliminate mn2 and solve for mg2

mg2soln = ((geom /. intrules(*/.msbarrules*) /. Solve[neom, mn2][[1]]) // Solve[#, mg2] &)[[1]]
{mg2 →
$$\left(-m^2 - p^2 + p^2 \ Z \ Z\Delta - \frac{1}{9} \ i \ (-Ifing0 + Ifingp) \ v^2 \ λ^2 \ \hbar + \frac{1}{9} \ i \ (Ifing0 + Ifingp + Iμ) \ v^2 \ Z\Delta^3 \ (δλ + λ)^2 \ \hbar - δm_1^2 - \frac{1}{6} \ v^2 \ Z\Delta \ (λ + δλ_a) - \frac{1}{6} \ tfinn \ Z\Delta^2 \ \hbar \ (λ + δλ_{2a}) - \frac{1}{6} \ t\mu \ Z\Delta^2 \ \hbar \ (λ + δλ_{2a}) + \frac{1}{18} \ i \ Iμ \ v^2 \ Z\Delta^3 \ δλ \ \hbar \ (λ + δλ_{2a}) + \frac{1}{6} \ i \ Iμ \ Z\Delta^2 \ μ^2 \ \hbar \ (λ + δλ_{2a}) - \frac{1}{6} \ i \ Iμ \ Z\Delta^2 \ \hbar \ (Ifing0 \ v^2 \ λ^2 \ \hbar - Ifingn \ v^2 \ λ^2 \ \hbar)$$

$$(λ + δλ_{2a}) - \frac{1}{6} \ tfing \ Z\Delta^2 \ \hbar \ ((1 + n) \ λ + (-1 + n) \ δλ_{2a} + 2 \ δλ_{2b}) - \frac{1}{6} \ tμ \ Z\Delta^2 \ \hbar$$

$$((1 + n) \ λ + (-1 + n) \ δλ_{2a} + 2 \ δλ_{2b}) - \frac{1}{6} \ i \ Iμ \ Z\Delta^2 \ μ^2 \ \hbar \ ((1 + n) \ λ + (-1 + n) \ δλ_{2a} + 2 \ δλ_{2b}) - \frac{1}{6} \ i \ Iμ \ Z\Delta^2 \ \hbar \ ((1 + n) \ λ + (-1 + n) \ δλ_{2a} + 2 \ δλ_{2b}) \right)$$

$$- \frac{1}{54} \ cμ \ v^2 \ Z\Delta^2 \ (δλ + λ)^2 \ \hbar^2 \ ((1 + n) \ λ + (-1 + n) \ δλ_{2a} + 2 \ δλ_{2b}) \right)$$

$$- \frac{1}{6} \ i \ Iμ \ Z\Delta^2 \ \hbar \ ((1 + n) \ λ + (-1 + n) \ δλ_{2a} + 2 \ δλ_{2b}) \right)$$

Gather kinematically distinct divergences for Goldstone EOM

```
cteq = \left( \left( mg2 - m^2 - \frac{\lambda}{\epsilon} v^2 - \frac{\hbar}{\epsilon} \left( (n+1) \lambda \right) (tfing) - \frac{\hbar}{\epsilon} (\lambda) (tfinn) / . mg2soln \right) / / 
                            CoefficientList[#, {p, v, tfing, tfinn, Ifingp}] & // Flatten //
                     Simplify // DeleteDuplicates = 0 // Thread
 \left\{ \left( -6 \text{ i } \delta \text{m}_{1}^{2} + \text{Z}\Delta^{2} \left( -\text{ i } \text{t}\mu + \text{I}\mu \left( -\text{m}^{2} + \mu^{2} \right) \right) \text{ } \hbar \text{ } \left( \left( 2 + \text{n} \right) \text{ } \lambda + \text{n } \delta \lambda_{2 \text{ a}} + 2 \text{ } \delta \lambda_{2 \text{ b}} \right) \right) \right/ 
           \left(-6 \text{ i} + 2 \text{ I}\mu \text{ Z}\Delta^2 \lambda \hbar + \text{I}\mu \text{ n} \text{ Z}\Delta^2 \lambda \hbar + \text{I}\mu \text{ n} \text{ Z}\Delta^2 \hbar \delta \lambda_{2 \text{ a}} + 2 \text{ I}\mu \text{ Z}\Delta^2 \hbar \delta \lambda_{2 \text{ b}}\right) == 0,
   True, \frac{1}{6}\hbar\left(-\lambda-\left(6\ \text{ii}\ Z\Delta^{2}\ (\lambda+\delta\lambda_{2\,\text{a}})\right)\right)
                      \left(-6 \text{ i} + 2 \text{ I}\mu \text{ Z}\Delta^{2} \lambda \hbar + \text{I}\mu \text{ n} \text{ Z}\Delta^{2} \lambda \hbar + \text{I}\mu \text{ n} \text{ Z}\Delta^{2} \hbar \delta \lambda_{2 \text{ a}} + 2 \text{ I}\mu \text{ Z}\Delta^{2} \hbar \delta \lambda_{2 \text{ b}}\right)\right) == 0,
    \frac{1}{6} \hbar \left( -(1+n) \lambda - \left( 6 i Z\Delta^{2} (\lambda + n \lambda + (-1+n) \delta \lambda_{2a} + 2 \delta \lambda_{2b}) \right) \right)
                      \left(-6 \text{ i} + 2 \text{ I} \mu \text{ Z} \Delta^2 \lambda \hbar + \text{I} \mu \text{ n} \text{ Z} \Delta^2 \lambda \hbar + \text{I} \mu \text{ n} \text{ Z} \Delta^2 \hbar \delta \lambda_{2 \text{ a}} + 2 \text{ I} \mu \text{ Z} \Delta^2 \hbar \delta \lambda_{2 \text{ b}}\right)\right) = 0,
    - ((i (-18 \lambda + 18 \Sigma\Delta \lambda - 12 i Ifing0 \Sigma\Delta^3 \delta\lambda^2 \hbar - 12 i I\mu \Sigma\Delta^3 \delta\lambda^2 \hbar - 24 i Ifing0 \Sigma\Delta^3 \delta\lambda \lambda \hbar -
                                 30 i I\mu Z\Delta^3 \delta\lambda \lambda \hbar – 12 i Ifing0 \lambda^2 \hbar – 6 i I\mu Z\Delta^2 \lambda^2 \hbar – 3 i I\mu n Z\Delta^2 \lambda^2 \hbar –
                                 12 i Ifing0 Z\Delta^3 \lambda^2 \hbar – 18 i I\mu Z\Delta^3 \lambda^2 \hbar + 4 c\mu Z\Delta^2 \delta \lambda^2 \lambda \hbar^2 + 2 c\mu n Z\Delta^2 \delta \lambda^2 \lambda \hbar^2 +
                                 8 c\mu Z^2 \delta\lambda \lambda^2 \hbar^2 + 4 c\mu n Z^2 \delta\lambda \lambda^2 \hbar^2 + 4 c\mu Z^2 \lambda^3 \hbar^2 + 2 Ifing0 I\mu Z^2 \lambda^3 \hbar^2 -
                                 2 Ifingn I\mu Z\Delta^2 \lambda^3 \hbar^2 + 2 c\mu n Z\Delta^2 \lambda^3 \hbar^2 + 18 Z\Delta \delta\lambda_a + Z\Delta^2 \hbar (2 c\mu n (\delta\lambda + \lambda) ^2 \hbar +
                                           I\mu \left(-6 \text{ is } Z\Delta \left(\delta \lambda + \lambda\right) + \lambda \left(-3 \text{ is } n + 2 \left(\text{Ifing 0} - \text{Ifing n}\right) \lambda \hbar\right)\right) \delta \lambda_{2a} -
                                 6 i I\mu Z\Delta^2 \lambda \hbar \delta\lambda_{2b} + 4 c\mu Z\Delta^2 \delta\lambda^2 \hbar^2 \delta\lambda_{2b} + 8 c\mu Z\Delta^2 \delta\lambda \lambda \hbar^2 \delta\lambda_{2b} + 4 c\mu Z\Delta^2 \lambda^2 \hbar^2 \delta\lambda_{2b}) /
                   (18 (-6 i + 2 I \mu Z \Delta^2 \lambda \hbar + I \mu n Z \Delta^2 \lambda \hbar + I \mu n Z \Delta^2 \hbar \delta \lambda_{2a} + 2 I \mu Z \Delta^2 \hbar \delta \lambda_{2b}))) = 
       0, -((2(-\lambda^2 + Z\Delta^3(\delta\lambda + \lambda)^2)\hbar)/
                  \left(3\left(-6\text{ i}+2\text{ I}\mu\text{ Z}\Delta^{2}\lambda\hbar+\text{I}\mu\text{ n}\text{ Z}\Delta^{2}\lambda\hbar+\text{I}\mu\text{ n}\text{ Z}\Delta^{2}\hbar\delta\lambda_{2\text{ a}}+2\text{ I}\mu\text{ Z}\Delta^{2}\hbar\delta\lambda_{2\text{ b}}\right)\right)\right)=0,
     (6 i (-1 + Z Z\Delta)) / (-6 i + 2 I\mu Z\Delta^2 \lambda \hbar + I\mu n Z\Delta^2 \lambda \hbar + I\mu n Z\Delta^2 \hbar \delta\lambda_{2a} + 2 I\mu Z\Delta^2 \hbar \delta\lambda_{2b}) =
```

Solve for counter-terms from Goldstone EOM

Note there are two solutions differing by a sign for $\delta\lambda$.

```
cts =
```

Solve[cteq, $\{\delta m_1, \delta \lambda_{1a}, \delta \lambda_{2a}, \delta \lambda_{2b}, \delta \lambda, Z, Z\Delta\}$] // FullSimplify // DeleteDuplicates;

Solve::svars: Equations may not give solutions for all "solve" variables. >>

$$\left\{ \left\{ -\frac{(2+n) \ \lambda \left(i \ t \mu + I \mu \ (m-\mu) \ (m+\mu) \right) \ \hbar}{6 \ i + I \mu \ (2+n) \ \lambda \ \hbar} \right. \right. \\ \left. \left\{ \left\{ -\frac{(2+n) \ \lambda \left(i \ t \mu + I \mu \ (m-\mu) \ (m+\mu) \right) \ \hbar}{6 \ i + I \mu \ (2+n) \ \lambda \ \hbar} \right. \\ \left. \left. \frac{1}{3 \ Z \Delta^4 \left(6 \ i + I \mu \ (2+n) \ \lambda \ \hbar} \right) \ \lambda \left(6 \ I \mu \ Z \Delta^{5/2} \ \lambda \ \hbar - 2 \ i \ c \mu \ (2+n) \ \lambda^2 \ \hbar^2 - 3 \ Z \Delta^4 \left(6 \ i + I \mu \ (2+n) \ \lambda \ \hbar \right) + 2 \ Z \Delta^3 \left(9 \ i + \lambda \ \hbar \left(- 6 \ (2 \ I fing 0 + I \mu) + i \ I \mu \ (I fing n + I \mu \ (2+n) + I fing 0 \ (3+2 \ n)) \ \lambda \ \hbar \right) \right) \right) , \\ \left(-1 + \frac{6 \ i}{Z \Delta^2 \left(6 \ i + I \mu \ (2+n) \ \lambda \ \hbar \right)} \right) , \left\{ -1 + \frac{6 \ i}{Z \Delta^2 \left(6 \ i + I \mu \ (2+n) \ \lambda \ \hbar \right)} \right) , \\ \left(-1 - \frac{1}{Z \Delta^{3/2}} \right) \lambda, \ \frac{1}{Z \Delta}, \ Z \Delta \right\}, \left\{ -\frac{(2+n) \ \lambda \left(i \ t \mu + I \mu \ (m-\mu) \ (m+\mu) \right) \ \hbar}{6 \ i + I \mu \ (2+n) \ \lambda \ \hbar} \right. \\ \left. \frac{1}{3 \ Z \Delta^4 \left(6 \ i + I \mu \ (2+n) \ \lambda \ \hbar \right)} \right. \lambda \left(-6 \ I \mu \ Z \Delta^{5/2} \ \lambda \ \hbar - 2 \ i \ c \mu \ (2+n) \ \lambda^2 \ \hbar^2 - 3 \ Z \Delta^4 \left(6 \ i + I \mu \ (2+n) \ \lambda \ \hbar \right) + 2 \ Z \Delta^3 \left(9 \ i + \lambda \ \hbar \left(- 6 \ (2 \ I fing 0 + I \mu) + i \ I \mu \ (I fing n + I \mu \ (2+n) + I fing 0 \ (3+2 \ n)) \ \lambda \ \hbar \right) \right) \right) , \\ \left. \lambda \left(-1 + \frac{6 \ i}{Z \Delta^2 \left(6 \ i + I \mu \ (2+n) \ \lambda \ \hbar \right)} \right) , \lambda \left(-1 + \frac{6 \ i}{Z \Delta^2 \left(6 \ i + I \mu \ (2+n) \ \lambda \ \hbar \right)} \right) \right) , \\ \left. \left(-1 + \frac{1}{Z \Delta^{3/2}} \right) \lambda, \ \frac{1}{Z \Delta}, \ Z \Delta \right\} \right\}$$

Gather kinematically distinct divergences for Higgs EOM

Solve for counter-terms from Higgs EOM

cts2 = Solve[cteq2[[2]], {Z
$$\Delta$$
}]
$$\left\{ \left\{ Z\Delta \rightarrow -\frac{9}{\left(3 \text{ i} + \text{Ifing0 } \lambda \text{ } \hbar - \text{Ifingn } \lambda \text{ } \hbar \right)^2} \right\} \right\}$$

Both equations should have the same solution:

$$(Z\Delta /. Solve[cteq2[[3]], {Z\Delta}][[1]]) - (Z\Delta /. cts2[[1]]) == 0$$

True

Final Counterterms

$$\begin{split} &\left(\left\{\delta\mathbf{m}_{2}^{2},\delta\lambda_{1a},\delta\lambda_{2a},\delta\lambda_{2b},\delta\lambda,\mathbf{z},\mathbf{z}\Delta\right\}/.\ \mathbf{cts}/.\ \mathbf{cts2}\,//\,\,\mathbf{Simplify}\right)[[1]]\,//\,\\ &\mathbf{DeleteDuplicates};\\ &\mathbf{counterterms} = \mathbf{Thread}\left[\left\{\delta\mathbf{m}_{1}^{2},\delta\lambda_{1a},\delta\lambda_{2a},\delta\lambda_{2b},\delta\lambda,\mathbf{z},\mathbf{z}\Delta\right\}\rightarrow *[[1]]\right]\\ &\left\{\delta\mathbf{m}_{1}^{2}\rightarrow-\frac{(2+n)\,\lambda\left(\mathrm{i}\,\,\mathrm{t}\mu+\mathrm{I}\mu\left(\mathrm{m}-\mu\right)\,\left(\mathrm{m}+\mu\right)\right)\,\hbar}{6\,\,\mathrm{i}\,+\mathrm{I}\mu\left(2+n\right)\,\lambda\,\hbar}\right.,\\ &\delta\lambda_{a}\rightarrow\frac{1}{1\,\,9\,683\,\left(6\,\,\mathrm{i}\,+\mathrm{I}\mu\left(2+n\right)\,\lambda\,\hbar\right)}\,\lambda\left(3\,\,\mathrm{i}\,+\mathrm{Ifing0}\,\lambda\,\hbar-\mathrm{Ifingn}\,\lambda\,\hbar\right)^{8}}\\ &\left(-2\,\,\mathrm{i}\,\,\mathrm{c}\mu\left(2+n\right)\,\lambda^{2}\,\hbar^{2}\,+1458\,\,\mathrm{I}\mu\,\lambda\,\hbar\right)\left(-\frac{1}{\left(3\,\,\mathrm{i}\,+\mathrm{Ifing0}\,\lambda\,\hbar-\mathrm{Ifingn}\,\lambda\,\hbar\right)^{2}}\right)^{5/2}\,-\\ &\frac{19\,683\,\left(6\,\,\mathrm{i}\,+\mathrm{I}\mu\left(2+n\right)\,\lambda\,\hbar\right)}{\left(3\,\,\mathrm{i}\,+\mathrm{Ifing0}\,\lambda\,\hbar-\mathrm{Ifingn}\,\lambda\,\hbar\right)^{8}}\,-\left(1458\,\left(9\,\,\mathrm{i}\,+\lambda\,\hbar\right)\left(-6\,\left(2\,\,\mathrm{Ifing0}\,+\mathrm{I}\mu\right)\,+\,\mathrm{i}\,\,\mathrm{I}\mu\,\left(\mathrm{Ifingn}\,+\,\mathrm{I}\mu\right)\right)\right)\left(3\,\,\mathrm{i}\,+\,\mathrm{Ifing0}\,\lambda\,\hbar-\mathrm{Ifingn}\,\lambda\,\hbar\right)^{6}\right),\\ &\delta\lambda_{2\,a}\rightarrow\lambda\left(-1+\frac{2\,\,\mathrm{i}\,\left(3\,\,\mathrm{i}\,+\,\mathrm{Ifing0}\,\lambda\,\hbar-\mathrm{Ifingn}\,\lambda\,\hbar\right)^{4}}{27\,\left(6\,\,\mathrm{i}\,+\,\mathrm{I}\mu\left(2+n\right)\,\lambda\,\hbar\right)}\right),\\ &\delta\lambda_{2\,a}\rightarrow\lambda\left(-1+\frac{2\,\,\mathrm{i}\,\left(3\,\,\mathrm{i}\,+\,\mathrm{Ifing0}\,\lambda\,\hbar-\mathrm{Ifingn}\,\lambda\,\hbar\right)^{4}}{27\,\left(6\,\,\mathrm{i}\,+\,\mathrm{I}\mu\left(2+n\right)\,\lambda\,\hbar\right)}\right),\\ &\delta\lambda\rightarrow\lambda\left(-1-\frac{1}{27\,\left(-\frac{1}{\left(3\,\,\mathrm{i}\,+\,\mathrm{Ifing0}\,\lambda\,\hbar-\mathrm{Ifingn}\,\lambda\,\hbar\right)^{2}}\right)^{3/2}}\right),\\ &Z\rightarrow\\ &-\frac{1}{9}\,\left(3\,\,\mathrm{i}\,+\,\mathrm{Ifing0}\,\lambda\,\hbar-\mathrm{Ifingn}\,\lambda\,\hbar\right)^{2},\\ &Z\Delta\rightarrow-\frac{9}{\left(3\,\,\mathrm{i}\,+\,\mathrm{Ifing0}\,\lambda\,\hbar-\mathrm{Ifingn}\,\lambda\,\hbar\right)^{2}}\right\} \end{split}$$

The should be momentum independent:

True