

# Renormalization of SI-2PIEA gap equations in the Hartree-Fock approximation

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Supplement to thesis Chapter 3

*Mathematica* notebook to compute counter-terms for two loop truncations of the two particle irreducible effective action

```
ClearAll[veom, geom, neom, regularisedtadpoles, mg2soln, cteq, cts,  $\delta m$ ,  $\delta \lambda$ ];
```

## Hartree-Fock gap equations with counterterms

Goldstone equation of motion. Quantities in reference to the thesis are:

$p$  is the four-momentum flowing through the propagators  $\Delta_G^{-1}$  and  $\Delta_N^{-1}$ ,

$mg^2$  is the Goldstone mass squared  $m_G^2$ ,

$mn^2$  is the Higgs mass squared  $m_H^2$ ,

$Z$  and  $Z\Delta$  are the wavefunction a propagator renormalization constants,

$m^2$  is the (renormalized) Lagrangian mass parameter,  $\delta m_0^2$ ,  $\delta m_1^2$  are its counter-terms,

$\lambda$  is the (renormalized) four point coupling,

$\delta\lambda_0$ ,  $\delta\lambda_{1a}$ ,  $\delta\lambda_{1b}$ ,  $\delta\lambda_{2a}$ ,  $\delta\lambda_{2b}$  are the independent coupling counter-terms,

$v$  is the scalar field vacuum expectation value,

$\hbar$  is the reduced Planck constant,

$n$  is the number of fields in the  $O(n)$  symmetry group,

$t_{\infty g}$ ,  $t_{\infty n}$  are the divergent tadpole integrals for the Goldstone, Higgs resp.,

$t_{fin g}$ ,  $t_{fin n}$  are the finite parts of the tadpoles for the Goldstone, Higgs resp.

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## Equations of motion

Vev equation of motion

```
(*veom=
Z $\Delta$ -1 (m2+ $\delta m_0^2$ ) v +  $\frac{\lambda+\delta\lambda_0}{6}$  v3 +  $\frac{\hbar}{6}$  Z $\Delta$  (n-1) ( $\lambda+\delta\lambda_{1a}$ ) v (t $\infty g$ +t $\infty n$ ) +  $\frac{\hbar}{6}$  Z $\Delta$  (3 $\lambda+\delta\lambda_{1a}+2\delta\lambda_{1b}$ ) v (t $\infty n$ +t $\infty n$ )
finveom=m2v +  $\frac{\lambda}{6}$  v3 +  $\frac{\hbar}{6}$  (n-1)  $\lambda$  v t $\infty n$  +  $\frac{\hbar}{2}$   $\lambda$  v t $\infty n$ *)
veom = v mg2
```

Goldstone equation of motion

$$\text{geom} = p^2 - \text{mg2} == Z \Delta p^2 - m^2 - \delta m_1^2 - Z \Delta \frac{\lambda + \delta \lambda_{1a}}{6} v^2 -$$

$$\frac{\hbar}{6} \left( (n+1) \lambda + (n-1) \delta \lambda_{2a} + 2 \delta \lambda_{2b} \right) Z \Delta^2 (t\omega g + t\text{fing}) - \frac{\hbar}{6} (\lambda + \delta \lambda_{2a}) Z \Delta^2 (t\omega n + t\text{finn})$$

$$\text{finmg2} = \text{mg2} /. \text{Solve}\left[p^2 - \text{mg2} == p^2 - m^2 - \frac{\lambda}{6} v^2 - \frac{\hbar}{6} (n+1) \lambda t\text{fing} - \frac{\hbar}{6} \lambda t\text{finn}, \text{mg2}\right][[1]]$$

Higgs equation of motion

$$\text{neom} = p^2 - \text{mn2} == Z \Delta p^2 - m^2 - \delta m_1^2 - Z \Delta v^2 \frac{(3 \lambda + \delta \lambda_{1a} + 2 \delta \lambda_{1b})}{6} -$$

$$\frac{\hbar}{6} (\lambda + \delta \lambda_{2a}) (n-1) Z \Delta^2 (t\omega g + t\text{fing}) - \frac{\hbar}{6} (3 \lambda + \delta \lambda_{2a} + 2 \delta \lambda_{2b}) Z \Delta^2 (t\omega n + t\text{finn})$$

$$\text{finmn2} = \text{mn2} /. \text{Solve}\left[p^2 - \text{mn2} == p^2 - m^2 - v^2 \frac{\lambda}{2} - \frac{\hbar}{6} \lambda (n-1) t\text{fing} - \frac{\hbar}{2} \lambda t\text{finn}, \text{mn2}\right][[1]]$$

## Infinite parts of tadpoles

$c_0$ ,  $c_1$ ,  $\Lambda$  and  $\mu$  are regularisation/renormalisation scheme dependent quantities

$$\text{regularisedtadpoles} = \{t\omega g \rightarrow c_0 \Lambda^2 + c_1 \text{mg2} \text{Log}[\Lambda^2 / \mu^2], t\omega n \rightarrow c_0 \Lambda^2 + c_1 \text{mn2} \text{Log}[\Lambda^2 / \mu^2]\}$$

## Sub in tadpole expressions, eliminate mn2 and solve for mg2

```
mn2fromneom = Solve[neom /. regularisedtadpoles, mn2][[1]]
mg2soln = mg2 /. (geom /. regularisedtadpoles /. mn2fromneom // Solve[#, mg2][[1]] &)
mn2soln = mn2 /. mn2fromneom /. mg2 -> mg2soln // Simplify
```

## Gather divergences proportional v, tfing and tfinn and set independently to zero

First we subtract the finite equation of motion, then gather coefficients of the remainder into a list and set each to zero (after some trimming and simplifying).

```
cteq = ((CoefficientList[mg2soln - finmg2, {p, v, tfing, tfinn}] // Flatten) //
DeleteDuplicates // Simplify // FullSimplify) == 0 // Thread
cteq2 = ((CoefficientList[mn2soln - finmn2, {p, v, tfing, tfinn}] // Flatten) //
DeleteDuplicates // Simplify // FullSimplify) == 0 // Thread
```

# Solve for counterterms

## Find counter-terms from the gap equations

```
cteqs = {cteq, cteq2} // Flatten // FullSimplify // DeleteDuplicates
cts = {δm12, δλ1a, δλ2a, δλ1b, δλ2b, Z, ZΔ} /. Solve[cteqs,
  {δm1, δλ1a, δλ2a, δλ1b, δλ2b, Z, ZΔ}] // FullSimplify // DeleteDuplicates
```

ZΔ is redundant in this truncation, can remove it :

```
cts /. ZΔ → 1 // FullSimplify
```

## Verify that the finite gap equations come out right

```
finmg2 ==
  (mg2soln /. Solve[cteqs, {δm1, δλ1a, δλ2a, δλ1b, δλ2b, Z, ZΔ}] /. ZΔ → 1 // FullSimplify //
    DeleteDuplicates)[[2]] // Simplify

finmn2 == mn2 /.
  ((neom /. regularisedtadpoles /. mg2 → mg2soln /. Solve[cteqs, {δm1, δλ1a,
    δλ2a, δλ1b, δλ2b, Z, ZΔ}] /. ZΔ → 1 // FullSimplify //
    DeleteDuplicates) // Solve[#, mn2] &) // FullSimplify
```

## Verify counter-term expressions in text

```
{δm12 ==  $\frac{-\hbar\lambda(n+2)}{6} \left( c_0\Lambda^2 + c_1 m^2 \text{Log}\left[\frac{\Lambda^2}{\mu^2}\right] \right) \frac{\delta\lambda_{1a} + \lambda}{\delta\lambda_{1b} + \lambda}} /.
  Solve[cteqs, {δm1, δλ1a, δλ2a, δλ1b, δλ2b, Z, ZΔ}] /. ZΔ → 1 //
  FullSimplify // Flatten // DeleteDuplicates

{δλ1a == δλ2a, δλ1b == δλ2b} /. Solve[cteqs, {δm1, δλ1a, δλ2a, δλ1b, δλ2b, Z, ZΔ}] /.
  ZΔ → 1 // FullSimplify // Flatten // DeleteDuplicates

{δλ1a / δλ1b} /. Solve[cteqs, {δm1, δλ1a, δλ2a, δλ1b, δλ2b, Z, ZΔ}] /. ZΔ → 1 //
  FullSimplify // Flatten // DeleteDuplicates

δλ1b /. Solve[cteqs, {δm1, δλ1a, δλ2a, δλ1b, δλ2b, Z, ZΔ}] /. ZΔ → 1 // FullSimplify //
  DeleteDuplicates$ 
```

## Total number of independent counter-term equations

```
Length[{cteqs} // Flatten // FullSimplify // DeleteDuplicates] -
  1 (* -1 because one of the "equations" is identically "True" *)
```