

## Chapter 2:

$\dot{x} = \sin x$  2.1.1.  $\dot{x} = 0 = \sin x$ ;  $x = n\pi$  2.1.2  $(n + \frac{1}{2})\pi$  where  $n$  is even.

2.1.3. a)  $\ddot{x} = \cos x \sin x$  b)  $\frac{1}{2} \sin(2x) = \cos(x) \sin(x)$ ;  $\ddot{x} = \frac{1}{2} \sin(2x)$ ;  $x = (n + \frac{1}{4})\pi$ ;  $n \in \mathbb{Z}$

2.1.4. a)  $x_0 = \pi/4$ ;  $t = \ln |(\csc x_0 + \cot x_0) / (\csc x + \cot x)|$

$$e^t = \frac{\csc x_0 + \cot x_0}{\csc x + \cot x} = \frac{\csc \pi/4 + \cot \pi/4}{\csc x + \cot x} = \frac{\frac{2}{\sqrt{2}} + 1}{\csc x + \cot x} = \frac{\sqrt{2} + 1}{\csc x + \cot x}$$

$$\frac{1}{\sin x + \frac{\cos x}{\sin x}} = \frac{\sin x}{1 + \cos x} = \frac{\sin(2 \cdot \frac{x}{2})}{1 + \cos(2 \cdot \frac{x}{2})} = \frac{2 \cos(\frac{x}{2}) \sin(\frac{x}{2})}{1 + 2 \cos^2(\frac{x}{2}) - 1} = \tan(\frac{x}{2}) = \frac{e^t}{\sqrt{2} + 1}$$

$$x(t) = 2 \tan^{-1} \left( \frac{e^t}{\sqrt{2} + 1} \right); \lim_{t \rightarrow \infty} x(t) = 2 \tan^{-1}(\infty) = 2 \cdot \frac{\pi}{2} = \pi$$

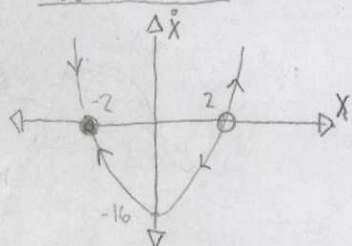
b)  $x(t) = 2 \tan^{-1} \left( \frac{e^t}{\csc x_0 + \cot x_0} \right)$

2.1.5a) A mechanical analog of  $\dot{x} = \sin x$  is the undamped pendulum having an  $x_0$  of the maximal point

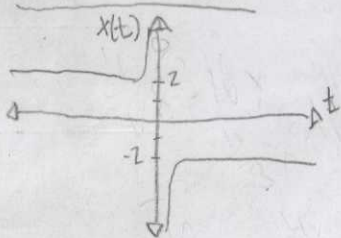
b) Unstable points are described by a positive slope (source) and stable points (sink), a negative slope. The function  $\dot{x} = \sin x$  at  $x^* = 0$  is unstable, while  $x^* = \pi$ , is stable.

$\dot{x} = 4x^2 - 16$  2.2.1

Vector Field:



Plot of  $x(t)$ :



Fixed Points:

$x = 2$

Stability:

source (unstable)

$x = -2$

sink (stable)

Solving for  $x(t)$ :

$$\frac{dx}{(x^2 - 4)} = 4t$$

$$\int \frac{A}{(x-2)} dx + \int \frac{B}{(x+2)} dx = 4t$$

$$A(x+2) + B(x-2) = 1$$

$$A = 1/4 @ x = 2$$

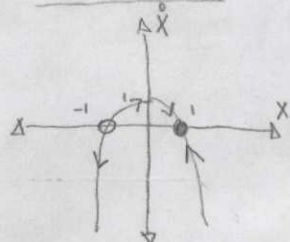
$$B = -1/4 @ x = -2$$

$$\ln|x-2| = 16t + C$$

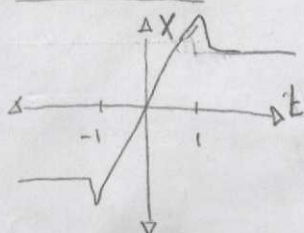
$$x(t) = \frac{2(e^{16t} + 1)}{1 - e^{16t}}$$

$\dot{x} = 1 - x^4$

2.2.2 Vector Field:



Plot of  $x(t)$ :



Fixed Points:

$x = 1$

Stability:

Source (unstable)

$x = -1$

sink (stable)

Solving for  $x(t)$ :

$$t = \int \frac{dx}{1 - x^4}$$

Unsolvable  $e^c + 2$

Analytical Solution of  $x(t)$ :

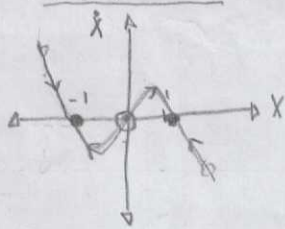
$$x_0 - 2 = (2 + x_0)e^{-t}$$

$$x_0 - 2 = (2 + x_0)e^{-t}$$

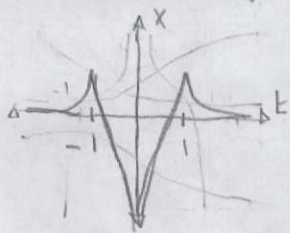
$$2t = \ln \frac{x_0 - 2}{2 + x_0} = C$$



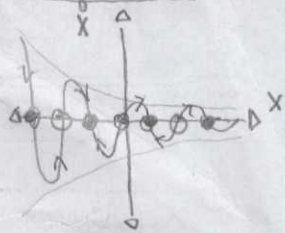
$$\dot{x} = x - x^3 \quad 2.2.3 \quad \text{Vector Field:}$$



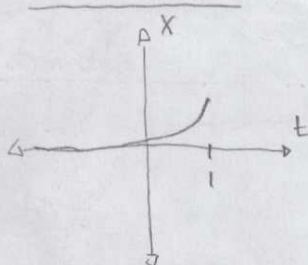
Plot of  $x(t)$ :



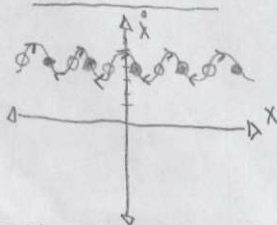
$$\dot{x} = e^{-x} \sin x \quad 2.2.4 \quad \text{Vector Field:}$$



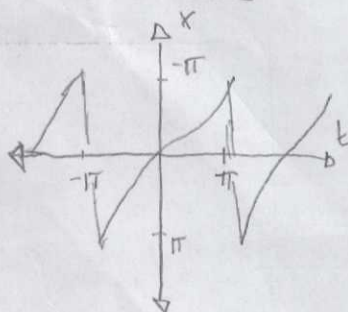
Plot of  $x(t)$ :



$$\dot{x} = 1 + \frac{1}{2} \cos x \quad 2.2.5 \quad \text{Vector Field:}$$



Plot of  $x(t)$ :



Fixed Points: Stability:

$x = -1$  stable (sink)

$x = 0$  unstable (source)

$x = 1$  stable (sink)

Analytical Solution of  $x(t)$ :

$$x(t) = \frac{1}{\sqrt{1 - e^{-2t}}} \quad \text{or} \quad x(t) = \frac{1}{\sqrt{1 - e^{-2t}}}$$

Solving for  $x(t)$ :

$$\begin{aligned} t &= \int \frac{dx}{x - x^3} = \int \frac{dx}{x(1 - x^2)} \\ &= \frac{1}{2} \int \frac{du}{(1-u)u} = \frac{1}{2} \int \frac{A}{1-u} du + \frac{1}{2} \int \frac{B}{u} du \\ &= -\frac{1}{2} \ln |1-u| + \frac{1}{2} \ln |u| = -\frac{1}{2} \ln \left| \frac{1-x^2}{x^2} \right| \\ &= \frac{1}{2} \ln \left| \frac{x^2}{1-x^2} \right| + C \\ (e^{2t} + 1)x^2 &= 1 \quad \Rightarrow \quad x = \frac{1}{\sqrt{1 + e^{2t}}} \end{aligned}$$

Fixed Points: Stability:

$x = 2n\pi$  source (unstable)

$x = (2n+1)\pi$  sink (stable)

Analytical Solution of  $x(t)$ :

$$x(t) = \arcsin^{-1}(C e^{t-1}) \quad \text{where } C = -(1 + \ln(\sin x_0))$$

Solving for  $x(t)$ :

$$\begin{aligned} t &= \int \frac{e^x}{\sin x} dx = \int dx + \int \cot(x) dx \\ &= 1 + \ln(\sin x) + C \\ x(t) &= \arcsin^{-1}(C e^{t-1}) \end{aligned}$$

Fixed Points

Stability:

Solving for  $x(t)$ :

$x = (4n+1)\frac{\pi}{2}$  sink (stable)

$x = (4n-1)\frac{\pi}{2}$  source (unstable)

Analytical solution of  $x(t)$ :

$$x(t) = \frac{2}{\sqrt{3}} \arctan\left(\frac{\tan(\frac{x}{2})}{\sqrt{3}}\right) + C$$

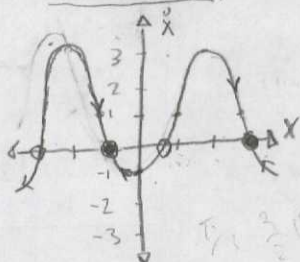
$$\begin{aligned} t &= \int \frac{1}{1 + \frac{1}{2} \cos x} dx \\ &= \int \frac{dx}{\frac{1}{2} + \cos(\frac{x}{2})} = \int \frac{\sec^2(\frac{x}{2}) dx}{\frac{\sec^2(\frac{x}{2})}{2} + 1} \\ &= \int \frac{\sec^2(\frac{x}{2}) dx}{3 + \tan^2(\frac{x}{2})} \end{aligned}$$

$$\begin{aligned} u &= \tan(\frac{x}{2}); \quad \frac{du}{dx} = \sec^2(\frac{x}{2}) \\ &= \int \frac{2\sqrt{3}}{3u^2 + 3} du = \frac{2}{\sqrt{3}} \int \frac{du}{u^2 + 1} \\ &= \frac{2}{\sqrt{3}} \arctan(u) + C \end{aligned}$$

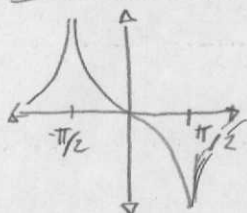
$$x(t) = \frac{2}{\sqrt{3}} \arctan\left(\frac{\tan(\frac{x}{2})}{\sqrt{3}}\right) + C$$

$$\ddot{x} = 1 - 2\cos x \quad 2.2.6.$$

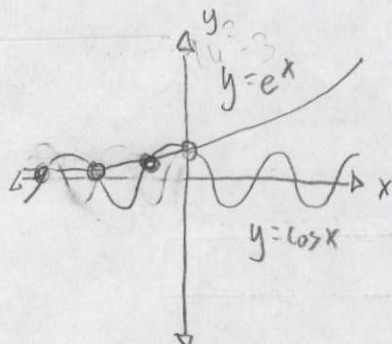
Vector Field



Plot of  $x(t)$



$$\dot{x} = e^x - \cos x \quad 2.2.7.$$



Fixed Points

$$x = (n + \frac{1}{2})\pi; n = \text{even}$$

Stability

sink (stable)

$$x = (n + \frac{1}{2})\pi; n = \text{odd} \quad \text{source (unstable)}$$

Analytical Solution of  $x(t)$

$$x(t) = \frac{\ln}{\sqrt{3}} \left| \frac{3\tan(\frac{x}{2}) - \sqrt{3}}{3\tan(\frac{x}{2}) + \sqrt{3}} \right|$$

Points of stability

$$e^x = \cos x$$

$$x_1 = 0 \quad \text{source (unstable)}$$

$$x_2 = -1.29 \quad \text{sink (stable)}$$

$$x_3 = -4.72 \quad \text{source (unstable)}$$

Solving for  $x(t)$

$$t = \int \frac{dx}{1 - 2\cos x} = \frac{1 - \tan^2(\frac{x}{2})}{1 + \tan^2(\frac{x}{2})} = \cos x$$

$$= \int \frac{dx}{2 \left[ \frac{1 - \tan^2(\frac{x}{2})}{1 + \tan^2(\frac{x}{2})} \right] - 1}$$

$$= \int \frac{dx}{2 \left[ \frac{1 - u^2}{1 + u^2} \right] - 1} \quad ; u = \tan(\frac{x}{2}) \quad \frac{du}{dx} = \sec^2(\frac{x}{2}) = \frac{1}{1 + u^2}$$

$$= \int \frac{du}{2 \sec^2(\frac{x}{2}) \left[ \frac{1 - u^2}{1 + u^2} \right] - 1}$$

$$= \int \frac{2 du}{[u^2 + 1] \left[ 2 \left[ \frac{1 - u^2}{1 + u^2} \right] - 1 \right]}$$

$$= \int \frac{2 du}{2 - 2u^2 - u^2 - 1}$$

$$= \int \frac{2 du}{-3u^2 - 1}$$

$$= -2 \left( \frac{3}{1} \right) \frac{du}{(3u - \sqrt{3})(3u + \sqrt{3})}$$

$$= -6 \left[ \int \frac{A du}{(3u - \sqrt{3})} + \int \frac{B du}{(3u + \sqrt{3})} \right]$$

$$= -6 \left[ \frac{\sqrt{3}}{2} \int \frac{du}{(3u - \sqrt{3})} - \frac{\sqrt{3}}{2} \int \frac{du}{(3u + \sqrt{3})} \right]$$

$$= \frac{\ln(3u + \sqrt{3})}{\sqrt{3}} - \frac{\ln(3u - \sqrt{3})}{\sqrt{3}}$$

$$= \frac{\ln(3\tan(\frac{x}{2}) + \sqrt{3})}{\sqrt{3}} - \frac{\ln(3\tan(\frac{x}{2}) - \sqrt{3})}{\sqrt{3}}$$

$$= \frac{\ln \left| \frac{3\tan(\frac{x}{2}) - \sqrt{3}}{3\tan(\frac{x}{2}) + \sqrt{3}} \right|}{\sqrt{3}} + C$$

$$x_0 = Q = V_0 C (1 - e^{-t/RC})$$

$$t = RC \ln \frac{V_0 C}{V_0 C - Q}$$

$$Q = \frac{V_0}{R} - \frac{Q}{RC} \quad 2.2.11. \quad Q(0) = 0; \quad t = RC \int \frac{dQ}{V_0 C - Q} = -RC \ln V_0 C - Q + C; \quad C = RC \ln V_0 C; \quad t = RC \ln \frac{V_0 C}{V_0 C - Q}$$

$$\dot{Q} = g(v) - \frac{Q}{RC} \quad 2.2.12. \quad V = g(v) - V_{cap} = V_0 - \frac{Q}{C}; \quad -g(v) + RI + \frac{Q}{C} = 0; \quad -g(v) + RI + \frac{Q}{C} = -g(v) + RI + \frac{Q}{C} = 0$$

$$\dot{Q} = g(v) - \frac{Q}{RC};$$

$$\text{Fixed Points: } g(v) = Q/RC$$

Stability: source (unstable)

The nonlinearity of the resistor has a relationship to resistance.

$$V_{RC} = Q/C$$



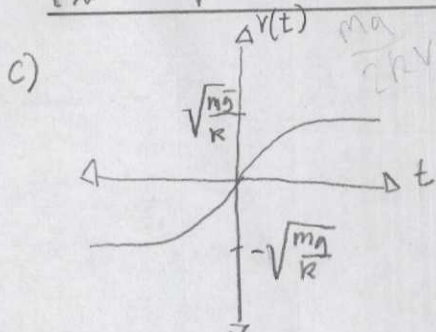
$m\dot{v} = mg - kv^2$  2.3.13 where  $m$  = mass,  $g$  = acceleration,  $k > 0$  = air resistance

a)  $\int g - \frac{k}{m}v^2 = \frac{1}{g} \int \frac{dv}{1 - \frac{k}{mg}v^2} = \frac{1}{g} \left[ \int \frac{dv}{1 - \sqrt{\frac{k}{mg}}v} + \int \frac{dv}{1 + \sqrt{\frac{k}{mg}}v} \right] = \sqrt{\frac{mg}{k}} \frac{1}{2g} \left[ \ln \left| 1 + \sqrt{\frac{k}{mg}}v \right| - \ln \left| 1 - \sqrt{\frac{k}{mg}}v \right| \right]$

$t = \frac{1}{2} \sqrt{\frac{m}{kg}} \ln \left| \frac{1 + \sqrt{\frac{k}{mg}}v}{1 - \sqrt{\frac{k}{mg}}v} \right| \Rightarrow \tanh^{-1} \left( \sqrt{\frac{k}{mg}}v \right)$

b)  $\lim_{t \rightarrow \infty} v(t) = \sqrt{\frac{mg}{k}}$  = "terminal velocity"

$v(t) = \sqrt{\frac{mg}{k}} \tanh \sqrt{\frac{kg}{m}} t$



d.  $V_{avg} = \frac{(31,400 - 2100) \text{ ft}}{116 \text{ sec}} = 252 \frac{\text{ft}}{\text{sec}}$

e.  $s = \frac{ds}{dt} = v = V \tanh \sqrt{\frac{kg}{m}} t$ ;  $s(t) = V \int \tanh \sqrt{\frac{kg}{m}} t dt$

$29,300 = V^2 \ln \cosh \frac{32.2 \text{ ft/sec}^2}{V} 116 \text{ sec} = V \int \frac{\sinh \sqrt{\frac{kg}{m}} t}{\cosh \sqrt{\frac{kg}{m}} t} dt$

$e = \frac{e^{\frac{3735 \text{ ft/sec}}{V}} + e^{-\frac{3735 \text{ ft/sec}}{V}}}{2} = V \int \frac{1}{u} du$

$V = 266 \text{ ft/sec}$

$V \approx V_{avg} = 252 \text{ ft/sec}$

$\frac{gt}{V} = \frac{32.2 \text{ ft/sec}^2 \cdot 116 \text{ sec}}{252 \text{ ft/sec}} = 14.8$

$\frac{V^2}{g} \ln \cosh \frac{gt}{V} \approx \frac{V^2}{g} \left[ \frac{gt}{V} - \ln 2 \right] = 265 \text{ ft/sec}$

$Ce^{-rt} = \frac{1-N/K}{N}$ ;  $N(1 + \frac{Ce^{-rt}}{K}) = 1$

$N = \frac{1}{1 + \frac{Ce^{-rt}}{K}} = \frac{No}{1 + e^{-rt}}$

General Solution:  $x = Ce^{-rt} + \frac{1}{K}$

b.  $x = 1/N$ ;  $\dot{x} = -rx(1 - \frac{1}{Kx}) = \frac{r}{K} - rx$ ;  $\dot{x} + rx - \frac{r}{K} = 0$

$\dot{x} = r(\frac{1}{K} - x)$

$N(t) = \frac{K}{KCe^{-rt} + 1}$

$N = \frac{No}{e^{-rt}}$

$= \frac{1}{K_1 a} \left[ \ln |x| - \ln \left| 1 - \frac{K-1}{K_1 a} x \right| \right] = \frac{1}{K_1 a} \ln \left| \frac{x}{1 - \frac{K-1}{K_1 a} x} \right| + C$

$x(t) = \frac{K-1}{K_1 a} + Ce^{-t/K_1 a}$ ;  $C = \frac{1}{x_0} - \frac{K-1}{K_1 a}$ ; Fixed points of stability:  $x = \frac{K_1 a}{K-1}$  source unstable

$\dot{N} = rN(1-N/K)$  2.3.1 a.  $t = \frac{1}{r} \int \frac{dN}{N(1-N/K)}$

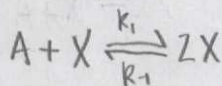
$= \frac{1}{r} \left[ \int \frac{A dN}{N} + \int \frac{B dN}{(1-N/K)} \right]$

$= \frac{1}{r} \left[ \int \frac{1 dN}{N} + \int \frac{1 dN}{K(1-N/K)} \right]$

$= \frac{1}{r} \ln N - \ln |1-N/K| + C$

$= \frac{1}{r} \ln \frac{N}{1-N/K} + C$

$\dot{x} = k_1 x - k_2 x^2$  2.3.2

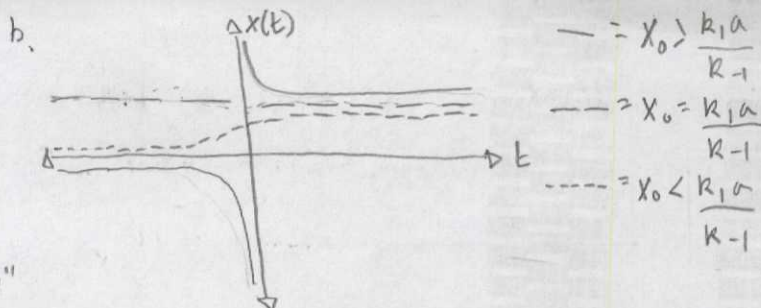


a.  $t = \int \frac{dx}{k_1 x - k_2 x^2} = \frac{1}{k_1 a} \int \frac{dx}{x - \frac{k_2}{k_1 a} x^2} = \frac{1}{k_1 a} \int \frac{dx}{x(1 - \frac{k_2}{k_1 a} x)}$

$= \frac{1}{k_1 a} \left[ \int \frac{A dx}{x} + \int \frac{B dx}{(1 - \frac{k_2}{k_1 a} x)} \right] = \frac{1}{k_1 a} \left[ \int \frac{1 dx}{x} + \frac{k_2}{k_1 a} \int \frac{dx}{1 - \frac{k_2}{k_1 a} x} \right]$

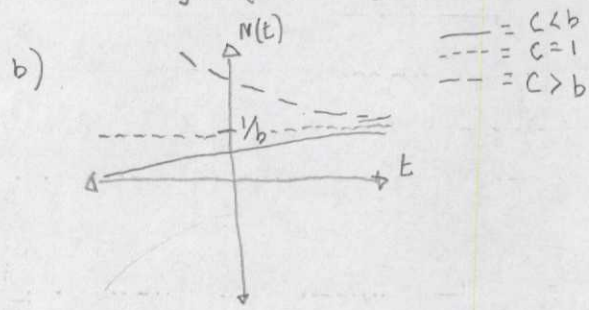
$= \frac{1}{k_1 a} \left[ \ln |x| - \ln \left| 1 - \frac{k_2}{k_1 a} x \right| \right] = \frac{1}{k_1 a} \ln \left| \frac{x}{1 - \frac{k_2}{k_1 a} x} \right| + C$





"Gompertz Law"

2.3.3 a)  $\dot{N} = -a N \ln(bN)$   $t = -\frac{1}{a} \int \frac{dN}{N \ln(bN)} = -\frac{b}{a} \int \frac{du}{u} = -\frac{b}{a} \ln[\ln(bN)]$ ;  $N(t) = C \frac{e^{-\frac{abt}{b}}}{b} = \frac{C}{b} e^{-at}$ ;  $a = \text{rate constant}$   
 $b = \text{Max amount of cells.}$



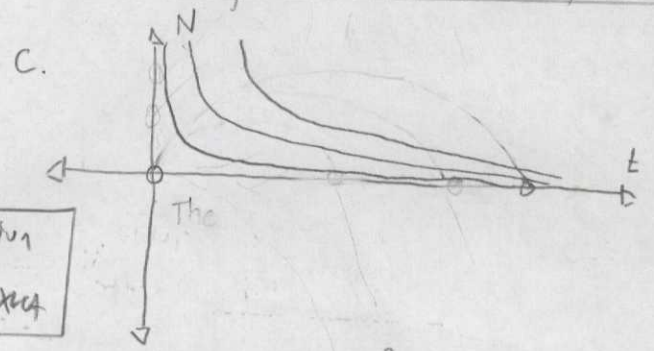
2.3.4. a)  $\frac{\dot{N}}{N} = r - a(N-b)^2$   $\lim_{N \rightarrow 0} \frac{\dot{N}}{N} = \lim_{N \rightarrow 0} r - a(N-b)^2 = r - ab^2 = \infty$ ;  $r = \infty$   
 $\lim_{N \rightarrow \infty} \frac{\dot{N}}{N} = \lim_{N \rightarrow \infty} r - a(N-b)^2 = r - \infty = 0$ ;  $r = \infty$

Each case of competition model at infinite population or extremely small populations that food amount or rate of consumption are insignificant to the competitors.

b) Fixed points of stability:

$N = 0$  source (unstable)  
 $N = \sqrt{\frac{r}{a}} + b$  sink (stable)

d) The solutions of the logistic equation  $y = C e^{-\frac{r}{a}t} + \frac{1}{k}$  are similar, if not exact to the Allele Effect.



2.3.5 a)  $\dot{x} = ax$   
 $\dot{y} = by$   
 a)  $x(t) = \frac{x(t)}{x(t)+y(t)} = \frac{e^{at}}{e^{at} + e^{bt}}$ ;  $\lim_{t \rightarrow \infty} x(t) \approx 1$

b)  $\dot{x}(t) = \frac{a e^{at} (e^{at} + e^{bt}) + e^{at} (a e^{at} + b e^{bt})}{(e^{at} + e^{bt})^2}$   
 $= \frac{e^{at} [a-b] e^{bt}}{(e^{at} + e^{bt})^2} = \frac{[a-b] e^{at}}{(e^{at} + e^{bt})} \left( \frac{e^{bt} + a e^{at}}{e^{at} + e^{bt}} \right)$   
 $= x[a-b](1-x)$

$t = \int \frac{du}{(b+u)(r-au^2)} = \int \frac{A}{b+u} du + \int \frac{Bu+C}{r-au^2} du$

$A(r-au^2) + (Bu+C)(b+u) = 1$   
 $@ u = \sqrt{\frac{r}{a}}; (B\sqrt{\frac{r}{a}} + C)(b + \sqrt{\frac{r}{a}}) = 1$   
 $B = \sqrt{\frac{a}{r}}; C = \sqrt{\frac{r}{a}}$

@  $u = -b$   $A = \frac{1}{r-(ab)^2}$

$= \frac{1}{r-(ab)^2} \int \frac{du}{b+u} + \frac{1}{\sqrt{r}} \int \frac{ab \cdot u + r}{r-au^2} du$  Partial Fractions (x2)  
 $= \frac{\ln N}{r-(ab)^2} + \frac{b}{4} \sqrt{\frac{a}{r}} \tanh^{-1}[(N+b)^2] + \frac{1}{\sqrt{a}} \tanh^{-1}(\sqrt{\frac{a}{r}}(N+b)) + C$

$\dot{X} = (1-X)P_{yx} - X P_{xy}$  2.3.6 a.  $X=0$   
 $P_{yx} = sX^a$ ;  $P_{xy} = (1-s)(1-X)^a$   $X=1$

$$X = \frac{a-1\sqrt{(1-s)}}{s}$$

$$\frac{1+\sqrt{(1-s)}}{s}$$

b) A plot of  $s(1-x)x^a$  and  $-(1-s)x(1-x)^a$  demonstrate  $-(1-s)x(1-x)^a > s(1-x)x^a$  for  $x=0$  and  $x=1$ , indicating, each fixed point is stable.

c. For  $X = \frac{a-1\sqrt{(1-s)}}{s}$  the plot of  $s(1-x)x^a > (1-s)x(1-x)^a$  suggesting a source.

$\dot{X} = X(1-X)$

2.4.1  $\dot{X} = f(X) = f(X^* + X) = f(X^*) + X f'(X^*) + O(X^2)$   
 $= X f'(X^*) + O(X^2)$   
 $= X(1-2X)$

$X=0$ ;  $f'(X^*) = 1$  : Unstable (source)  
 $X=1$ ;  $f'(X^*) = -1$  : Stable (sink)

$\dot{X} = X(1-X)(2-X)$

2.4.2  $\dot{X} = f(X) = f(X^* + X) = X f'(X^*)$   
 $= X(2X(1-X))$

$X=0$   $f'(X^*) = 0$  Half-stable  
 $X=1$   $f'(X^*) = 0$  Half-stable  
 $X=2$   $f'(X^*) = -4$  sink (stable)

$\dot{X} = \tan X$

2.4.3  $\dot{X} = f(X) = f(X^* + X) = X \sec^2(X)$

$\dot{X} = X^2(6-X)$

2.4.4  $\dot{X} = f(X) = f(X^* + X) = X[12X-3X^2]$

$X=0$   $f'(X) = 0$  Half-stable  
 $X=6$   $f'(X) = -36$  sink (stable)

$\dot{X} = 1 - e^{-X^2}$

2.4.5  $\dot{X} = f(X) = f(X^* + X) = X[2e^{-X^2}]$

$\dot{X} = \ln X$

2.4.6  $\dot{X} = f(X) = f(X^* + X) = \frac{1}{X}$

$\dot{X} = aX - X^3$

2.4.7  $\dot{X} = f(X) = f(X^* + X) = X[a - 3X^2]$

$X=0$   $f'(X^*) = 0$  Half-stable  
 $X=1$   $f'(X^*) = 1$  source (unstable)

	(+)	(-)	(0)
$X \neq 0$	source	sink	Half-stable
$X = \sqrt{a}$	sink	source	Half-stable
$X = -\sqrt{a}$	source	sink	Half-stable

2.4.8  $\dot{N} = -aN \ln(bN)$   $\dot{N} = f(N) = f(N + N^*) = -\frac{aN}{b} [1 + b \ln(bN)]$

$\dot{X} = -X^3$

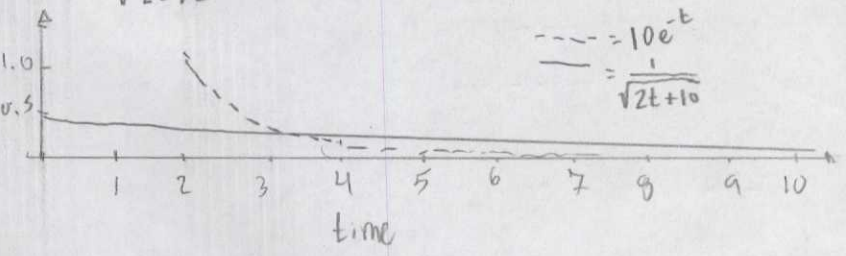
2.4.9 a.  $t = -\int \frac{dx}{x^3} = \frac{1}{2x^2} + C$ ;  $X(t) = \sqrt{\frac{1}{2t+C}}$

$\lim_{x \rightarrow 0} t = \frac{1}{0} + C = \frac{\infty}{x(t)}$

b. if  $X_0 = 10$

$t = -\int \frac{1}{X} = -\ln X$

$X(t) = X_0 e^{-t} = 10e^{-t}$



$\dot{X} = -X^0$

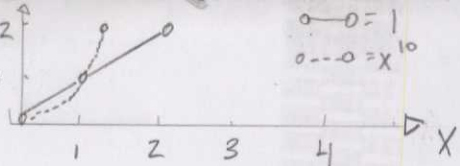
2.5. 1 a.  $c=0$

b.  $dx = -dt$ ;  $x(t) = -t$ ; if  $t=0$  is considered finite time, then yes.

$t = -\int \frac{dx}{x^c} = -\frac{x^{1-c}}{1-c}$ ;  $t(x=1) - t(x=0) = -\frac{1^{1-c}}{1-c} + \frac{0^{1-c}}{1-c} = \frac{1}{c+1}$



$$\dot{x} = 1 + x^2; 2.5.2. \quad \lim_{x \rightarrow \infty} \dot{x} = \infty$$



$$\dot{x} = rx + x^3; 2.5.3$$

$$t = \int \frac{dx}{x(r+x^2)} = \int \frac{A}{x} dx + \int \frac{Bx+C}{r+x^2} dx; A(r+x^2) + (Bx+C)x = 1$$

$$x=0; A = \frac{1}{r}; 1 + \frac{x^2}{r} + Bx^2 + Cx = 1$$

$$(B + \frac{1}{r})x^2 + Cx = 0$$

$$(B + \frac{1}{r})x = -C$$

$$x=0; C=0$$

$$x=1; B = -\frac{1}{r}$$

$$\text{If } x_0 \neq 0; \lim_{t \rightarrow \infty} x(t) = \infty$$

$$\dot{x} = x^{1/3}; 2.5.4. \quad x(0)=0; t = \int \frac{dx}{x^{1/3}} = \frac{3}{2} x^{2/3}; x(t) = \sqrt{\frac{2}{3} t - \frac{2}{3} C}$$

$$\dot{x} = |x|^{p/q}; 2.5.5. \quad x(0)=0; a) t = \frac{q}{p+q} (x)^{\frac{p+q}{q}}; x(t) = \left( \frac{p+q}{q} (t+C) \right)^{q/(p+q)}; C = \text{many solutions at zero because of root.}$$

$$x(t) = \left( \frac{p+q}{q} (t+C) \right)^{q/(p+q)}$$

$$b) x(t) = \left( \frac{p+q}{q} (t+C) \right)^{q/(p+q)}; \text{if } p > q; x(0) = \left( \frac{p+q}{q} (0+C) \right)^{q/(p+q)} = 0; C=0$$

h(t) = height; t = time; A = cross-sectional area; v(t) = velocity; a v(t) = A h'(t); 2.5.6 a) Newton's first law that for every force there exists an equal and opposite counter force.

$$b) \frac{1}{2} m v^2 = m g h; v^2 = 2 g h$$

$$c) h'(t) = -\sqrt{\frac{a}{A} 2 g h}; d) h(0)=0; t = -\sqrt{\frac{A h}{2 a}} 2 \sqrt{h}; h(t) = -\sqrt{\frac{a}{2 A}} t$$

$$m \ddot{x} = -kx; 2.6.1$$

The text states, there are no periodic solutions to  $\dot{x} = F(x)$  because undamped systems do not oscillate, and, damped oscillations do not occur for first order systems. Strogatz' statement does not fit the equation of 2.6.1.

$$\dot{x} = F(x); 2.6.2.$$

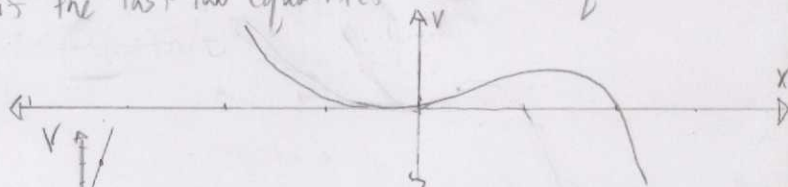
$$x(t) = x(t+T)$$

$$\int_t^{t+T} F(x) \frac{dx}{dt} dt = \int_t^{t+T} F(x) \dot{x}(t) dt = \int_t^{t+T} F(x) \dot{x}(t+T) dt$$

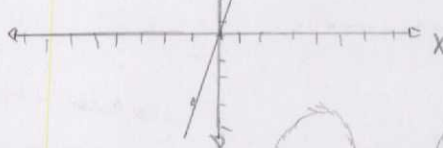
The contradiction of the argument is the last two equalities are unequal.

$$\dot{x} = x(1-x); 2.7.1$$

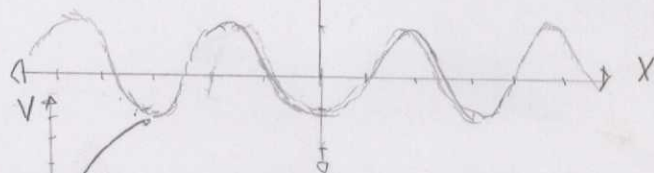
$$\frac{dV}{dx} = \dot{x} = x(1-x); V = (1-\frac{x}{3}) x^2$$



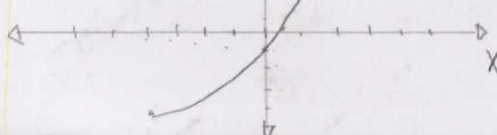
$$\dot{x} = 3; 2.7.2 \quad \frac{dV}{dx} = \dot{x} = 3; V = 3x$$



$$\dot{x} = \sin x; 2.7.3 \quad \frac{dV}{dx} = \dot{x} = \sin x; V = -\cos(x)$$



$$\dot{x} = 2 + \sin x; 2.7.4 \quad \frac{dV}{dx} = \dot{x} = 2 + \sin x; V = 2x - \cos(x)$$

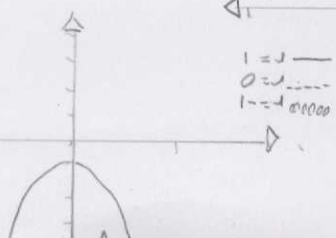
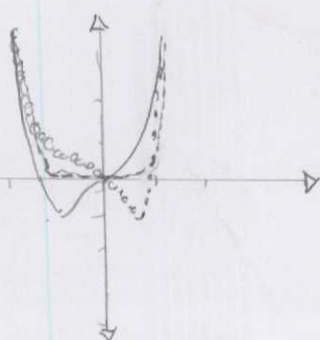


$$2.7.5. \frac{dx}{dt} = -\sinh x; V = -\cosh(x)$$

$$2.7.6. \frac{dx}{dt} = r + x - x^3; V = r + \frac{x^2}{2} - \frac{x^4}{4}$$

$$2.7.7. \frac{dx}{dt} = f(x) = x = f(x); V = \int f(x) dx = \frac{x^2}{2} + C$$

$$f(x) = \frac{dx}{dt} = d(V-C)/dx$$



The solution  $x(t)$  cannot oscillate because of the existence and uniqueness of  $f(x)$ , and the solutions for  $f(x)=0$ , that  $V=C$  or  $C=0$ , withstanding,  $\frac{dx}{dt} = \frac{d(V-C)}{dx} = \frac{dx}{dt}$ , then the solution  $x(t)$  also corresponds to a nonperiodic function.

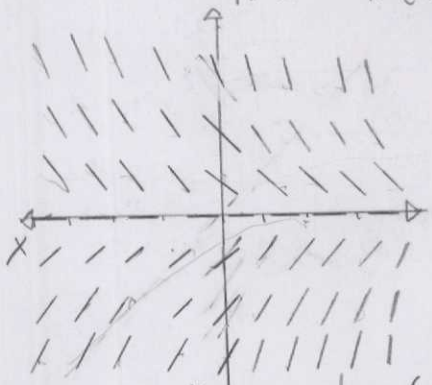
2.8.1 The horizontal lines are to be expected in Figure 2.8.2 because of the slope being zero at  $x=1$ .

$$\dot{x} = x$$

$$\frac{dx}{dt} = x$$
  

$$\ln x = t$$
  

$$x = e^t$$

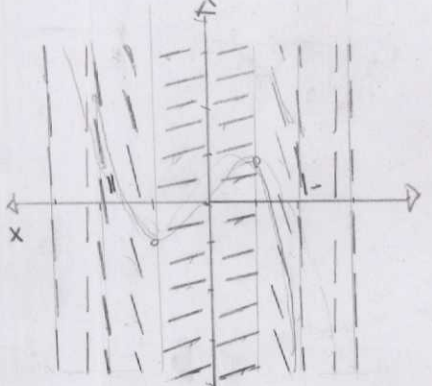


$$\dot{x} = 1 - x^2$$

$$\frac{dx}{dt} = 1 - x^2$$
  

$$\int \frac{dx}{1-x^2} = \int dt$$
  

$$\frac{1}{2} \ln \frac{1+x}{1-x} = t$$



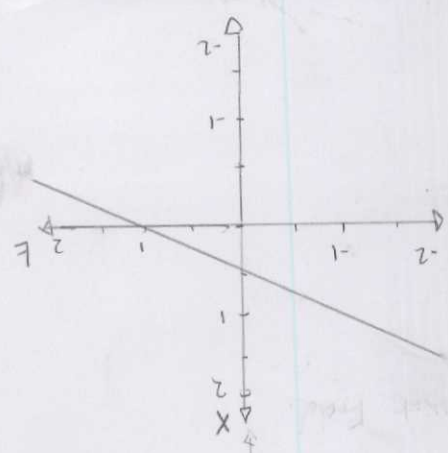
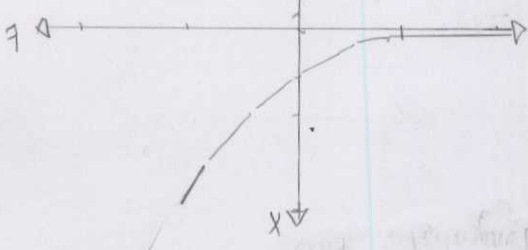
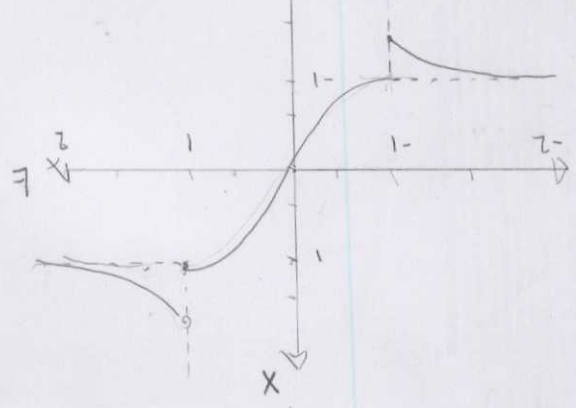
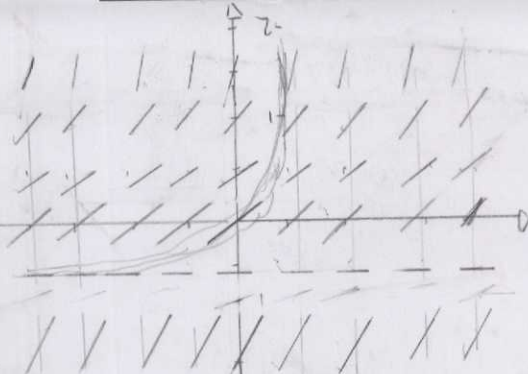
$$\dot{x} = 1 - 4x(1-x)$$

$$\frac{dx}{dt} = 1 - 4x(1-x)$$
  

$$t = \int \frac{dx}{1 - 4x(1-x)}$$
  

$$t = \frac{1}{5} \ln \frac{1-x}{1-4x} = t$$
  

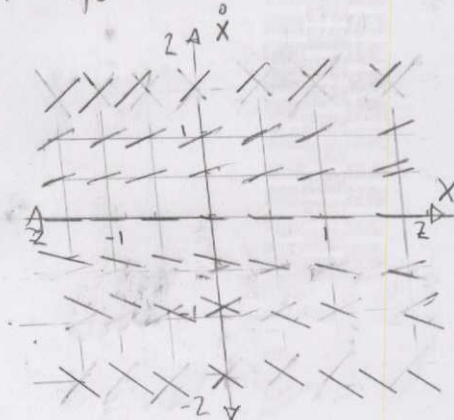
$$x = 0.5 - \frac{1}{4} e^{-5t}$$





$$\dot{x} = \sin x$$

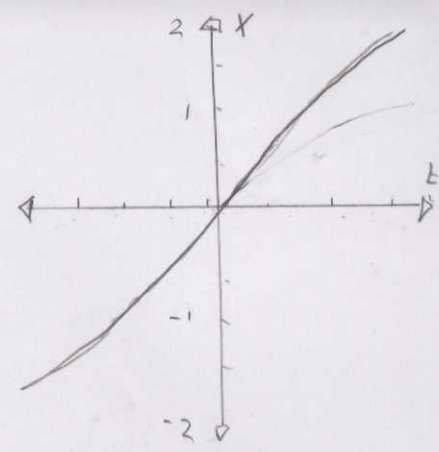
2.32 d) Slope Field



$$\frac{dx}{dt} = \sin x$$

$$\int \cosh x = t$$

$$x = \sinh^{-1}(t)$$



$\dot{x} = -x; x(0) = 1$  2.3.3. a)  $x(t) = Ce^{-t}; C=1; x(t) = e^{-t}$  Int  $-t$

Euler's Method

b)  $\Delta t = 1; x(t_0 + \Delta t) \approx x_0 + f(x_0)\Delta t; x(1 + 1) = 0 + e^{-1} \cdot 1 = 0.3679$

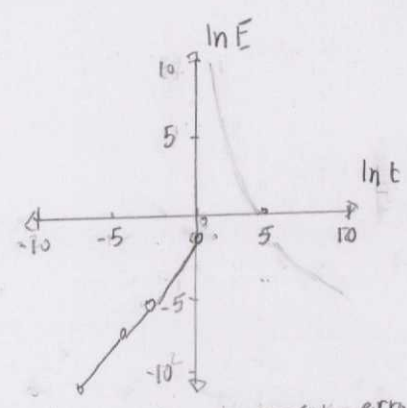
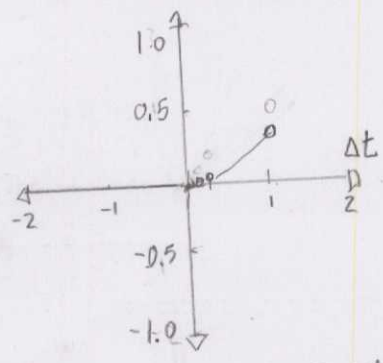
$\Delta t = 10^{-n} \quad n=1 \quad x_1 = x_0 + e^{-x_0} 10^{-1} = 0.36341$

$n=2 \quad x_2 = x_1 + e^{-x_1} 10^{-2} = 0.36697$

$n=3 \quad x_3 = x_2 + e^{-x_2} 10^{-3} = 0.36795$

$n=4 \quad x_4 = x_3 + e^{-x_3} 10^{-4} = 0.36787$

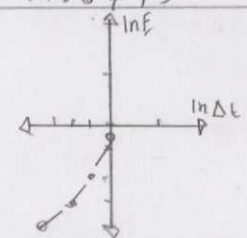
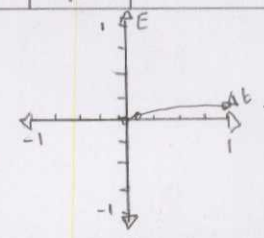
c)  $E = |\hat{x}(1) - x(1)|$



The results of  $E = |\hat{x}(1) - x(1)|$  vs  $\Delta t$  represent error of Euler's method. While the plot of  $\ln E$  vs  $\ln t$  characterizes nothing informative.

$\dot{x} = -x; x(0) = 1$  2.3.4.  $x(t) = e^{-t}$

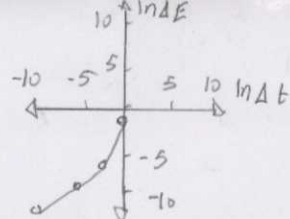
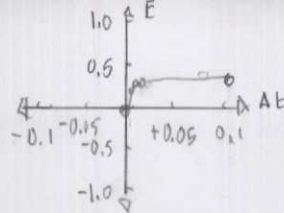
n	$\Delta t$	$f(x_n)$	$x_n = x_{n-1} + f(x_{n-1})\Delta t$	$E =  \hat{x}(1) - x(1) $	$\ln E$
0	$10^0$	$-x_0$	0.36788	0.00	-1.00
1	$10^{-1}$	$\exp(x_{n-1})$	0.33527	0.03261	-3.42
2	$10^{-2}$		0.36577	0.0211	-6.16
3	$10^{-3}$		0.36773	0.0002	-9.93
4	$10^{-4}$		0.36773	0.0000	-10.79



Improved Euler's Method

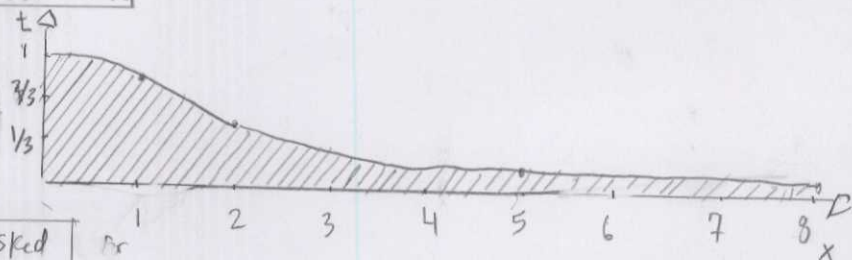
$\dot{x} = -x; x(0) = 1$  2.3.5.  $x(t) = e^{-t}$

n	$\Delta t$	$f(x_n)$	$x_n = x_{n-1} + \frac{1}{2}(k_1 + 2k_2 + k_3)$	$E =  \hat{x}(1) - x(1) $	$\ln E$
0	$10^0$				
1	$10^{-1}$	$\exp(x_{n-1})$			
2	$10^{-2}$				
3	$10^{-3}$				
4	$10^{-4}$				



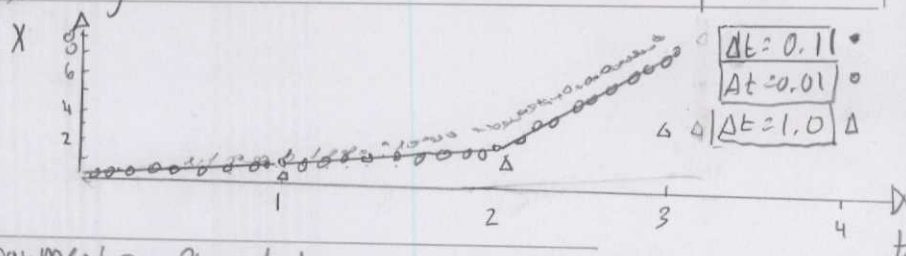
The Euler method aided the analysis of numerical methods; including, Precision, Euler's Improved Method approached the solution of  $f(x) = e^{-x}$  with less round-off error, Runge-Kutta's Routine provided the least round-off errors with  $10^{-20}$  across the spreadsheet, and necessitated high-precision.

$\dot{x} = x + e^{-x}$  2.8.6. a)  $t = \int \frac{1}{x + e^{-x}} dx = \int \frac{e^x}{x e^x + 1} dx$



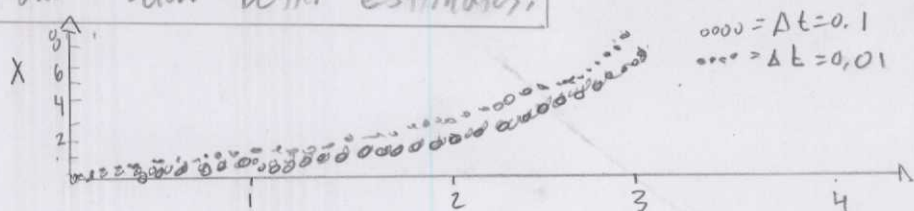
I noticed the book asked for  $x(t)$  (and not  $t(x)$ ).

This led me to investigate a Numerical method of integration; Withstanding, Runge-Kutta Routine aided with the plot of  $x(t)$ .



b. At  $t=0$ , analytical arguments provided an  $x=1.011$ .

c. Stepsizes of  $\Delta t = 0.1$  and  $0.01$  had different results, including, inaccuracies above and below better estimates;



d) See part a.

$x_1 = x_0 + f(x_0) \Delta t$  2.8.7 a)  $x(t_1) = x(t_0 + \Delta t)$

Taylor Series:

$$x(t + \Delta t) = \sum_{n=0}^{\infty} \frac{x^{(n)}(t)}{n!} (\Delta t)^n = x(t) + x'(t) \cdot \Delta t + O(\Delta t^2) + O(\Delta t^3)$$

$$f(t + \Delta t) = f(t) + f'(t) \cdot \Delta t + O(\Delta t^2) = x_0 + f'(t) \cdot \Delta t$$

b)  $|x(t_1) - x| = |x(t_1) - x(t_1) - x'(t) \cdot \Delta t - O(\Delta t^2)| = |O(\Delta t^2)| = \frac{x''(t) \Delta t^2}{2!} = C(\Delta t^2)$

$$C = \frac{x''(t)}{2!}$$



Taylor Series: 2.8.8.  $x = x + e^{-x}$  ;  $|X(t_0) - t_0| = |X(t_0) - X(t_0) - X'(t_0)\Delta t - \frac{X''(t_0)\Delta t^2}{2!}| = \frac{X''(t_0)\Delta t^2}{2} = \frac{X''(t_0)\Delta t^2}{2} = O(\Delta t^2)$

$\dot{x} = x + e^{-x}$

2.8.9 Runge-Kutta :  $X_{n+1} = X_n + \frac{1}{6}(k_1 + 2k_2 + 2k_3 + k_4)$  Where  $k_1 = f(X_n)\Delta t$

$X(t + \Delta t) = X(t_0) + X'(t_0)\Delta t + \frac{X''(t_0)\Delta t^2}{2!} + O(\Delta t^3)$

$k_2 = f(X_n + \frac{1}{2}k_1)\Delta t$

$k_3 = f(X_n + \frac{1}{2}k_2)\Delta t$

$k_4 = f(X_n + k_3)\Delta t$

$k_1 = f(X_n)\Delta t = X'(t_0)\Delta t$

$k_2 = f(X_n + \frac{1}{2}k_1)\Delta t = f(X_n) + f'(X_n)\frac{1}{2}k_1 + O[(\frac{1}{2}k_1)^2]$

$k_3 = f(X_n + \frac{1}{2}k_2)\Delta t = f(X_n) + f'(X_n)\frac{1}{2}k_2 + O[(\frac{1}{2}k_2)^2]$

$= f(X_n) + f'(X_n)\frac{1}{2}[f(X_n) + f'(X_n)\frac{1}{2}k_1 + O[(\frac{1}{2}k_1)^2]] + O[(\frac{1}{2}k_2)^2]$

$k_4 = f(X_n + k_3)\Delta t = f(X_n) + f'(X_n) \cdot k_3 + O[k_3^2]$

$= f(X_n) + f'(X_n)[f(X_n) + f'(X_n)\frac{1}{2}[f(X_n) + f'(X_n)\frac{1}{2}k_1 + O[(\frac{1}{2}k_1)^2]] + O[(\frac{1}{2}k_2)^2]] + O[k_3^2]$

$X_{n+1} = X_n + \frac{1}{6}(k_1 + 2k_2 + 2k_3 + k_4) = X_n + \frac{1}{6}(X'(t_0)\Delta t + 2X'(t_0) + X''(t_0)k_1$

$+ 2X'(t_0) + X''(t_0)[X'(t_0) + X''(t_0)X'(t_0)\Delta t]$

$+ 2X'(t_0) + X''(t_0)[X(t_0) + \frac{X'(t_0)}{2}[X'(t_0) + X''(t_0)X'(t_0)\Delta t]])$

$|X(t_1) - X_1| = |X(t_0 + \Delta t) - X_{n+1}| = O(\Delta t^5)$

### Chapter 3

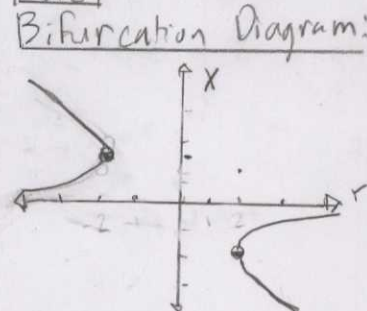
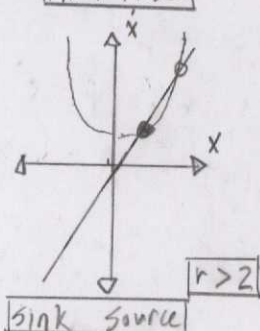
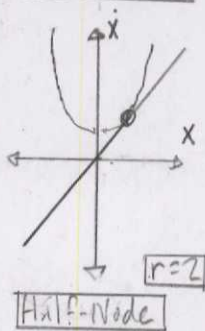
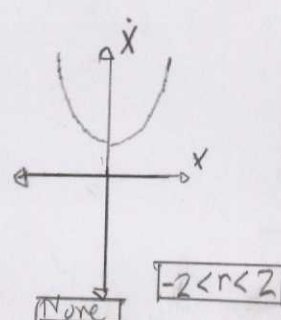
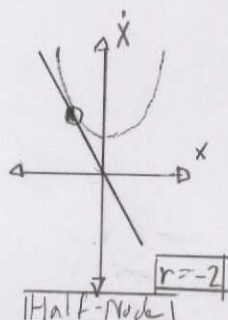
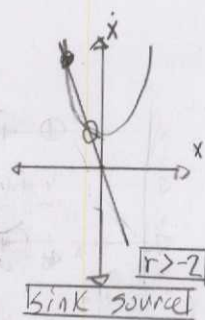
$\dot{x} = 1 + rx + x^2$  3.1.1.

Vector Field:

$x = \frac{r \pm \sqrt{r^2 - 4}}{2}$

$r = \frac{r \pm \sqrt{(r-2)(r+2)}}{2}$

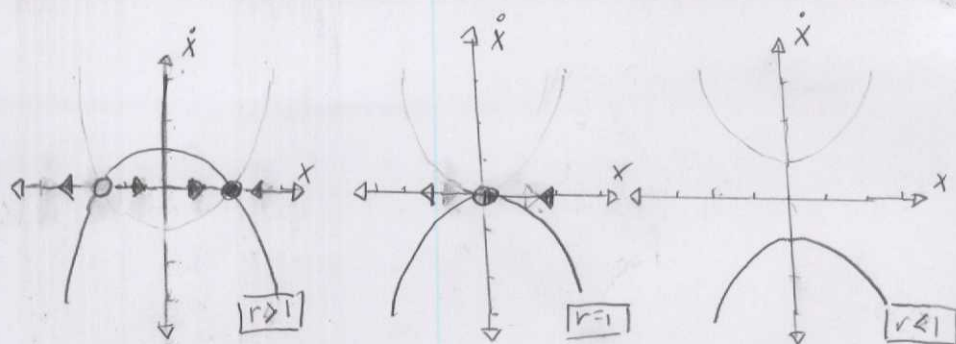
r	Bifurcations
> -2	Two
-2	One
0	zero
2	one
> 2	Two.



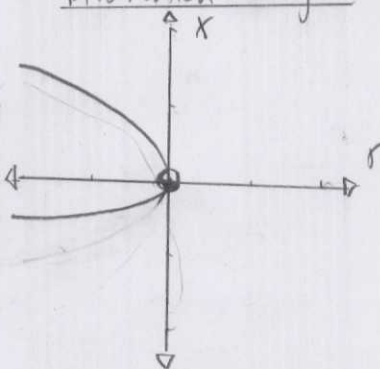
$\dot{x} = r - \cosh x$  3.1.2. Vector Field

$r = \cosh(x)$

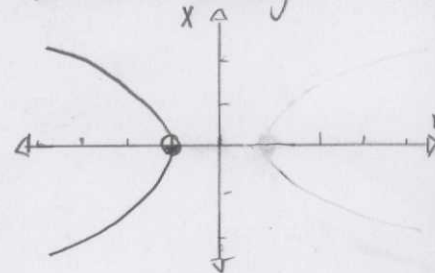
$r$	Bifurcations
$< 1$	Zero
$= 1$	One
$> 1$	Two



Bifurcation Diagram:



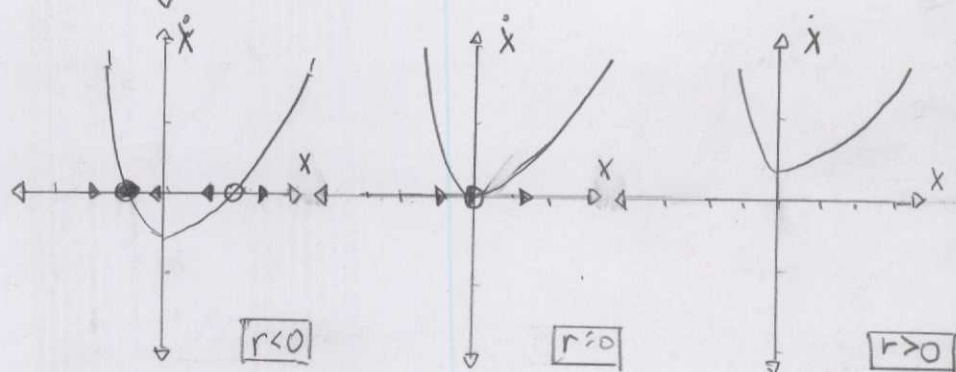
Bifurcation Diagram:



$\dot{x} = r + x - \ln(1+x)$  3.1.3

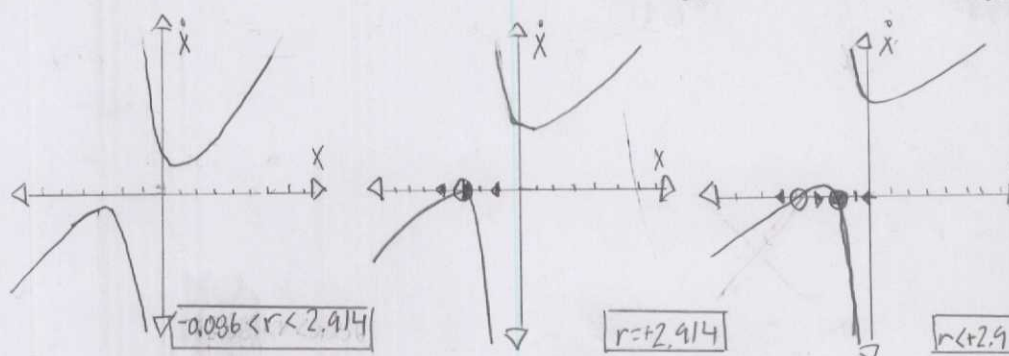
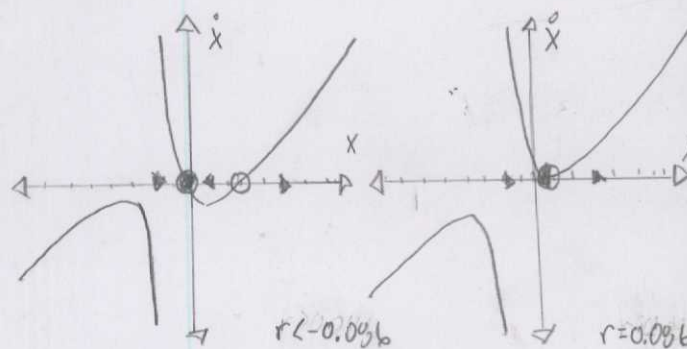
Vector Field

$r$	Bifurcation
$> 0$	Zero
$= 0$	One
$< 0$	Two



$\dot{x} = r + \frac{1}{2}x - x/(1+x)$  3.1.4. Vector Field:

$r$	Bifurcations
$r < -0.096$	Two
$r = -0.096$	One
$-0.096 < r < 2.914$	Zero
$r = 2.914$	One
$r > 2.914$	Two



Bifurcation Diagram:

