

# MA579H Scientific Computing

## Numerics of first order ODEs-II

Rafikul Alam  
Department of Mathematics  
IIT Guwahati

# Lecture outline

- Runge-Kutta method for ODEs
- Vector version of Euler's method for systems of first order ODEs.
- Converting higher order ODEs into a system of first order ODEs.

## Second order Runge-Kutta method/Heun's method

Integrating  $y' = f(x, y)$  on  $[x_j, x_{j+1}]$ , we have

$$y(x_{j+1}) - y(x_j) = \int_{x_j}^{x_{j+1}} f(t, y(t)) dt.$$

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the trapezoid quadrature rule gives

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$$\begin{aligned} Y_{j+1} &= y_j + hf(x_j, y_j) \\ y_{j+1} &= y_j + \frac{h}{2} [f(x_j, y_j) + f(x_{j+1}, Y_{j+1})], \quad j = 0 : n - 1. \end{aligned}$$

The second order RK method is also known as Heun's method.

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$$y(x_{j+1}) - y(x_j) \approx hf(x_{j+\frac{1}{2}}, y(x_{j+\frac{1}{2}})).$$

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$$y_{j+1} = y_j + hf(x_{j+1/2}, y_{j+1/2}), \quad j = 0 : n - 1,$$

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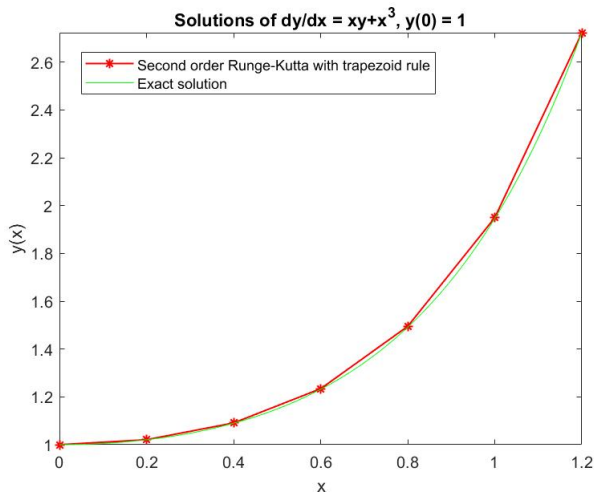
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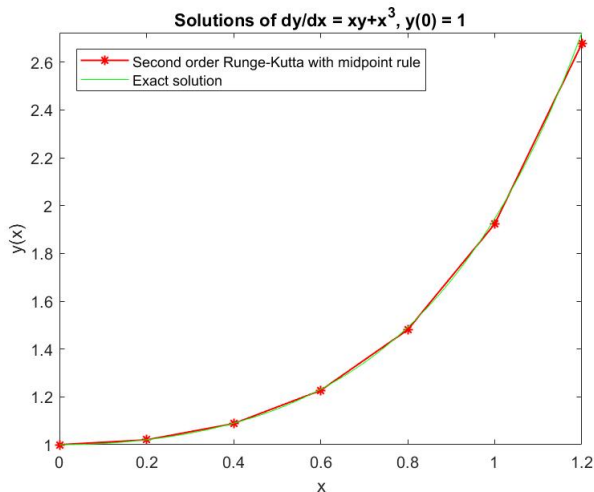
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# Second order Runge-Kutta: Example



**Figure :** Exact solution  $y(x) = 3e^{x^2/2} - x^2 - 2$  of the non-autonomous ODE  $\frac{dy}{dx} = xy + x^3$  satisfying  $y(0) = 1$  along with solution via Second order Runge Kutta method with trapezoid rule and  $h = 0.2$ .

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for  $j = 0 : n - 1$ .



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# Fourth order Runge-Kutta: Example

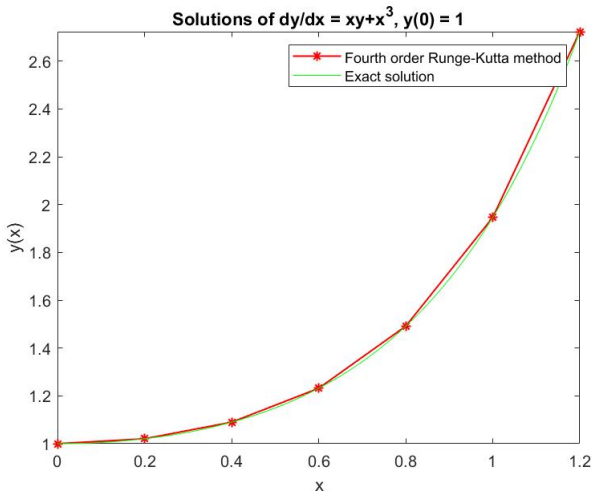


Figure : Exact solution  $y(x) = 3e^{x^2/2} - x^2 - 2$  of the non-autonomous ODE  $\frac{dy}{dx} = xy + x^3$  satisfying  $y(0) = 1$  along with solution via Fourth order Runge Kutta method and  $h = 0.2$ .

# Taylor's Method

Suppose that the solution  $y(x)$  of the IVP  $\frac{dy}{dx} = f(x, y)$ ,  $y(x_0) = y_0$  belongs to  $C^{k+1}$ .

By Taylor's expansion,

$$y(x_i + h) = y(x_i) + hy'(x_i) + \cdots + \frac{h^k}{k!}y^{(k)}(x_i) + \frac{h^{k+1}}{(k+1)!}y^{(k+1)}(\xi)$$

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Now,

$$y'(x) = f(x, y) \Rightarrow y^{(j)}(x) = \frac{d^{(j-1)} f(x, y)}{dx^{j-1}} := f^{(j-1)}(x, y), \quad j = 2 : k + 1.$$

Therefore,

$$\begin{aligned} y(x_i + h) = y(x_i) + hf(x_i, y(x_i)) + \cdots &+ \frac{h^k}{k!} f^{(k-1)}(x(i), y(x_i)) \\ &+ \frac{h^{k+1}}{(k+1)!} f^{(k)}(\xi, y(\xi)) \end{aligned}$$

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**Taylor's Method of order  $k$ :**

$$y_{i+1} = y_i + hf(x_i, y_i) + \cdots + \frac{h^k}{k!} f^{(k-1)}(x_i, y_i), \quad \text{for } i = 0 : n.$$

For  $k = 1$ , this is Forward Euler Method  $y_{i+1} = y_i + hf(x_i, y_i)$ ,  $i = 0 : n$ .

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Since

$$f'(x, y) = \frac{df(x, y)}{dx} = f_x(x, y) + f_y(x, y) \frac{dy}{dx} = f_x(x, y) + f_y(x, y) f(x, y),$$

therefore **Taylor's Method of order 2** is

$$y_{i+1} = y_i + hf(x_i, y_i) + \frac{h^2}{2} (f_x(x_i, y_i) + f_y(x_i, y_i) f(x_i, y_i)), \quad i = 0 : n,$$

where clearly the **error** is  $\mathcal{O}(h^3)$ .



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**Example:** For the IVP  $\frac{dy}{dx} = xy + x^3$ ,  $y(0) = 1$ ,

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Hence we have 2nd order Taylor's method

$$\begin{aligned} y_{i+1} &= y_i + hf(x_i, y_i) + \frac{h^2}{2}f'(x_i, y_i) \\ &= y_i + h(x_i y_i + x_i^3) + \frac{h^2}{2}(y_i + 3x_i^2 + x_i^2 y_i + x_i^4). \end{aligned}$$

for  $i = 0, \dots, n$ .

# Taylor's Method: Example

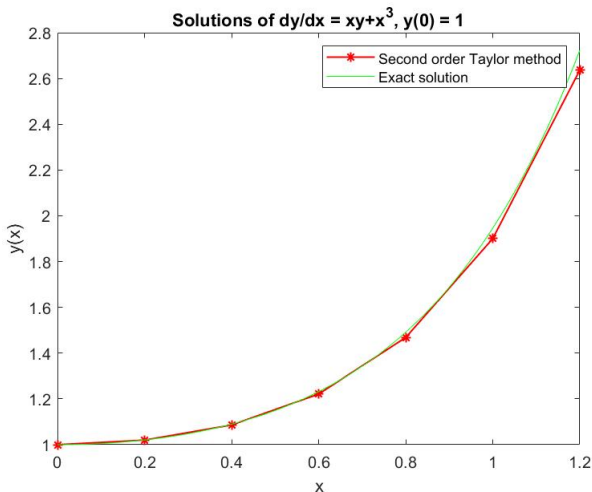


Figure : Exact solution  $y(x) = 3e^{x^2/2} - x^2 - 2$  of the non-autonomous ODE  $\frac{dy}{dx} = xy + x^3$  satisfying  $y(0) = 1$  along with the solution via second order Taylor's method with  $h = 0.2$ .

# Systems of ODE

A first-order system has the form: For  $x \in [a, b]$  solve

$$y_1' = f_1(x, y_1, y_2, \dots, y_n)$$

$$y_2' = f_2(x, y_1, y_2, \dots, y_n)$$

$$\vdots$$

$$y_n' = f_n(x, y_1, y_2, \dots, y_n)$$

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Setting  $\mathbf{y} := [y_1, \dots, y_n]^\top$  and defining  $\mathbf{f} : [a, b] \times \mathbb{R}^n \longrightarrow \mathbb{R}^n$  by

$$\mathbf{f}(x, \mathbf{y}) := [f_1(x, \mathbf{y}), \dots, f_n(x, \mathbf{y})]^\top,$$

we have the IVP

$$\mathbf{y}' = \mathbf{f}(x, \mathbf{y}), \quad x \in [a, b] \text{ and } \mathbf{y}(x_0) = \mathbf{y}_0.$$

We can use vector version of Euler method to solve the IVP.

# Vector version of Euler method

Consider the system of first order ODE

$$\mathbf{y}' = \mathbf{f}(x, \mathbf{y}), \mathbf{y}(x_0) = \mathbf{y}_0.$$

Forward Euler's method for the system

$$\mathbf{y}_{j+1} = \mathbf{y}_j + h \mathbf{f}(x_j, \mathbf{y}_j), \quad j = 0 : n - 1.$$

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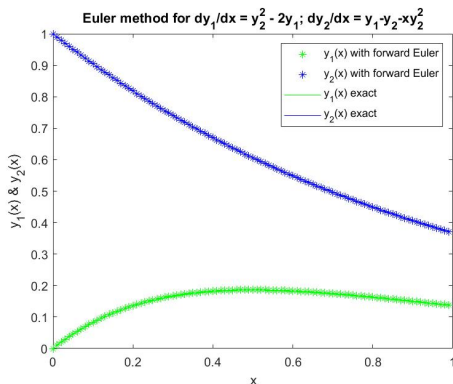
$$\mathbf{y}_{j+1} = \mathbf{y}_j + h \mathbf{f}(x_{j+1}, \mathbf{y}_{j+1}), \quad j = 0 : n - 1.$$

At each step of backward Euler, we have to solve the nonlinear system  $\mathbf{y}_{j+1} - \mathbf{y}_j - h \mathbf{f}(x_{j+1}, \mathbf{y}_{j+1}) = 0$  for  $\mathbf{y}_{j+1}$ .



# Vector version of Euler method

Example: 
$$\begin{bmatrix} y_1' \\ y_2' \end{bmatrix} = \begin{bmatrix} y_2^2 - 2y_1 \\ y_1 - y_2 - xy_2^2 \end{bmatrix}, \quad \begin{bmatrix} y_1(0) \\ y_2(0) \end{bmatrix} = \begin{bmatrix} 0 \\ 1 \end{bmatrix} \Rightarrow \begin{bmatrix} y_1 \\ y_2 \end{bmatrix} = \begin{bmatrix} xe^{-2x} \\ e^{-x} \end{bmatrix}.$$



**Figure :** Exact solutions  $y_1(x) = xe^{2x}$  and  $y_2(x) = e^{-x}$  of the first order system of ODEs in Example 1 and their respective approximations via forward Euler method with  $h = 0.01$ .

# Higher order equation

Higher order ODE can often be converted to a system of first order ODE.  
Consider the  $n$ -th order ODE

$$y^{(n)} = f(x, y, y', \dots, y^{(n-1)}), \quad y^{(j)}(x_0) = \alpha_j, \quad j = 0 : n - 1.$$

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Introducing new variable  $y_1 := y, y_2 := y', \dots, y_n := y^{(n-1)}$ , we have

$$\mathbf{y}' := \begin{bmatrix} y_1' \\ \vdots \\ y_{n-1}' \\ y_n' \end{bmatrix} = \begin{bmatrix} y_2 \\ \vdots \\ y_n \\ f(x, y_1, \dots, y_n) \end{bmatrix} =: \mathbf{f}(x, \mathbf{y}), \quad \mathbf{y}'(x_0) = \mathbf{y}_0,$$

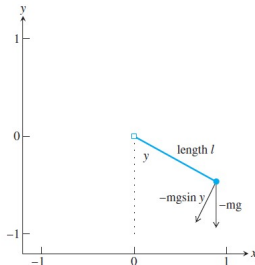
where  $\mathbf{y} := [y_1, \dots, y_n]^\top$  and  $\mathbf{y}_0 := [\alpha_0, \dots, \alpha_{n-1}]^\top$ .

# Motion of a pendulum

The motion of a pendulum of length  $\ell$  is governed by the equation

$$y'' = -\frac{g}{\ell} \sin(y), \quad y(0) = a, \quad y'(0) = b,$$

where  $y$  is angle measured in radian and  $g$  is gravity.



**Figure :** The component of force along the tangential direction is  $F = -mg \sin y$  where  $y$  is the angle made by the pendulum bob with the vertical axis.

# Motion of a pendulum

Setting  $y_1 = y$  and  $y_2 = y'$ , we have the first order system

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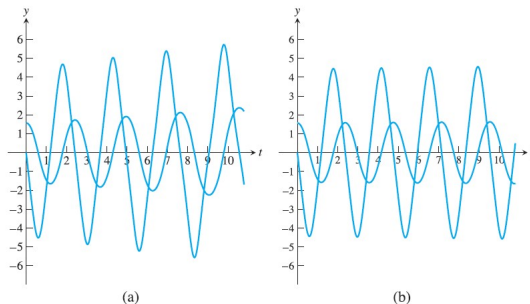
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If the pendulum is started from a position straight out to the right then the initial conditions are  $y_1(0) = \pi/2$  and  $y_2(0) = 0$ . Now, considering  $\ell = 1$  and  $g = 9.81m/sec^2$ , we can test the suitability of Eulers Method as a solver for this system.

# Motion of a pendulum



**Figure :** Euler method for the motion of the pendulum. Smaller oscillations plot angle  $y_1$  and larger oscillations plot angular velocity  $y_2$ . (a)  $h = 0.01$  is too large showing growing amplitude of oscillations (b)  $h = 0.001$  is more accurate.

# Orbit of a satellite

Let  $(x, y)$  denote the position of a satellite. Then Newton's law of motion yields the second order system of ODE

$$\begin{aligned}m_1 x'' &= -\frac{gm_1 m_2 x}{(x^2 + y^2)^{3/2}} \\m_1 y'' &= -\frac{gm_1 m_2 y}{(x^2 + y^2)^{3/2}}\end{aligned}$$

To transform it to a system of first order ODEs, let  $v_x = x'$  and  $v_y = y'$ .



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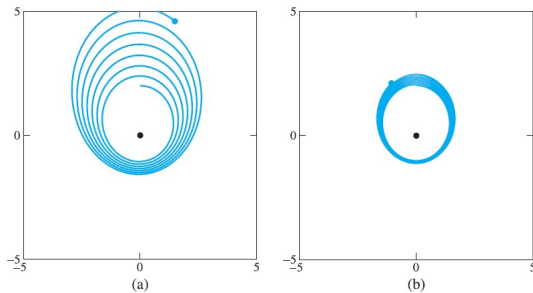
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$$\begin{aligned}m_1 x'' &= -\frac{gm_1 m_2 x}{(x^2 + y^2)^{3/2}} \\m_1 y'' &= -\frac{gm_1 m_2 y}{(x^2 + y^2)^{3/2}}\end{aligned}$$

To transform it to a system of first order ODEs, let  $v_x = x'$  and  $v_y = y'$ . Then the system is rewritten as

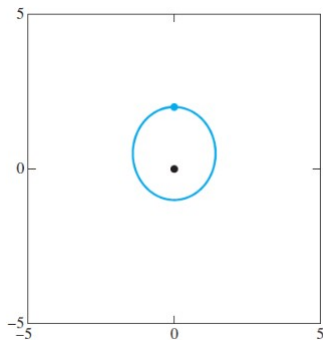
$$\begin{aligned}x' &= v_x \\v_x' &= -\frac{gm_2 x}{(x^2 + y^2)^{\frac{3}{2}}} \\y' &= v_y \\v_y' &= -\frac{gm_2 y}{(x^2 + y^2)^{\frac{3}{2}}}\end{aligned}$$

# Orbit of a satellite



**Figure :** The one body problem approximated with forward Euler method (a)  $h = 0.01$  (b)  $h = 0.001$ .

# Orbit of a satellite



**Figure :** The one body problem approximated with RK2 using trapezoid rule with  $h = 0.01$ .