

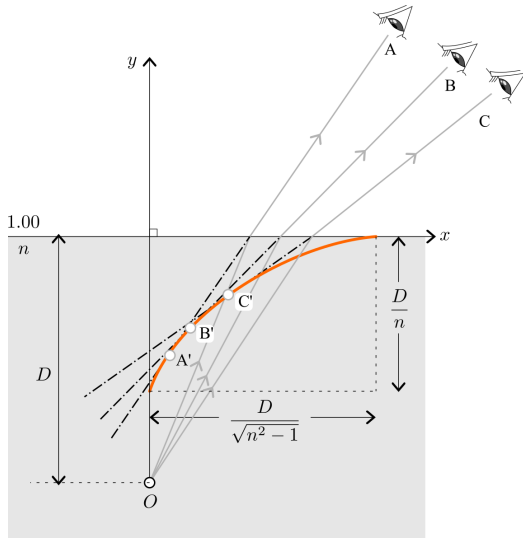
# Rediscovery of astroid from the refraction image by a flat boundary

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January 26, 2025

## 1 Introduction

If you see a point object under water from above the flat surface of water, the position of image depends on the point of view(POV). The image seen within a normal plane traces out a specific curve that is called a *squashed astroid* as the POV moves around.



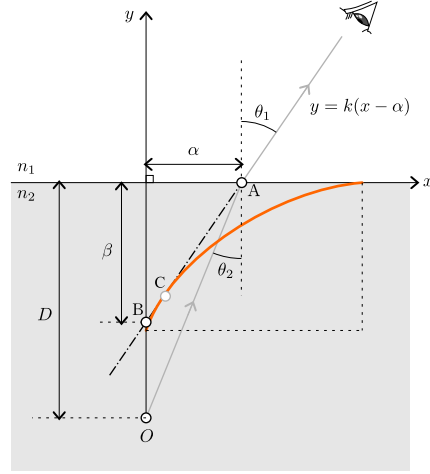
Let the normal plane that contains the object and POV be the  $xy$ -plane, and let the intersection of normal plane and the surface of water be the  $x$ -axis, and the normal through object  $y$ -axis. Then the trace of image is part of the curve

$$\left| \frac{x}{M} \right|^{2/3} + \left| \frac{y}{N} \right|^{2/3} = 1,$$

where  $M = D/\sqrt{n^2 - 1}$  and  $N = D/n$ , and  $D$  is the depth of object and  $n$  is the index of refraction of the water relative to the air.

## 2 Derivation of the formula

Let the indices of refraction of air and of water be  $n_1$  and  $n_2$ , respectively. The point object  $O$  is at the depth  $D$  below the boundary of air and water. A ray starts off the object and enters the boundary of media at  $\alpha$  away from the  $y$ -axis with angle  $\theta_2$  from the normal at that point and then refracts into air with angle  $\theta_1$  from the same normal.



From Snell's law we have

$$\sin \theta_1 = \frac{n_2}{n_1} \sin \theta_2 = n \sin \theta_2.$$

The extension of refracted ray is described by the equation

$$y = k(x - \alpha),$$

where

$$k = \frac{1}{\tan \theta_1} = \frac{\cos \theta_1}{\sin \theta_1},$$

and considering the Snell's law,

$$k = \frac{\sqrt{1 - n^2 \sin^2 \theta_2}}{n \sin \theta_2}.$$

This line meets the  $y$ -axis at  $y = \beta$ , thus

$$\beta = -k\alpha.$$

By the geometry we have

$$\alpha = D \tan \theta_2 = \frac{D \sin \theta_2}{\cos \theta_2},$$

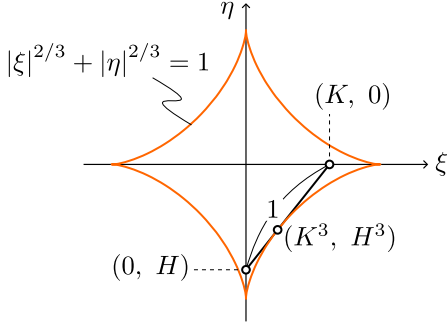
and

$$\begin{aligned} \beta &= -k\alpha \\ &= -\frac{D \sin \theta_2}{\cos \theta_2} \frac{\sqrt{1 - n^2 \sin^2 \theta_2}}{n \sin \theta_2} \\ &= -\frac{D \sqrt{1 - n^2 \sin^2 \theta_2}}{n \cos \theta_2}. \end{aligned}$$

Now, let  $K = \alpha/M$  and  $H = \beta/N$ , then

$$\begin{aligned} K^2 + H^2 &= \frac{\alpha^2}{M^2} + \frac{\beta^2}{N^2} \\ &= \frac{(n^2 - 1) \sin^2 \theta_2 + 1 - n^2 \sin^2 \theta_2}{\cos^2 \theta_2} \\ &= \frac{1 - \sin^2 \theta_2}{\cos^2 \theta_2} \\ &= 1 \end{aligned}$$

Let  $\xi = x/M$  and  $\eta = y/N$ , then as the POV moves around in the  $xy$ -plane, the points  $A(\alpha, 0)$ , and  $B(0, \beta)$  move accordingly, and the points  $(K, 0)$  and  $(0, H)$  in the  $\xi\eta$ -plane follow suite, while keeping the distance between them a constant; namely 1.



The envelope of such a segment is well-known as an *astroid*<sup>1</sup>, which is described by the equation

$$|\xi|^{2/3} + |\eta|^{2/3} = 1.$$

The image is at the point of tangency C of the segment  $\overline{AB}$  and the envelope of the moving segment, for it is the instant point of divergence of the neighboring rays. Its corresponding point in the  $\xi\eta$ -plane is  $(K^3, H^3)$ .

Thus we can obtain the coordinates of image  $(x_C, y_C)$  from the relation

$$\begin{cases} \xi_C = \frac{x_C}{M} = K^3 = \frac{\alpha^3}{M^3}, \\ \eta_C = \frac{y_C}{N} = H^3 = \frac{\beta^3}{N^3}. \end{cases}$$

That is

$$\begin{cases} x_C = \frac{\alpha^3}{M^2}, \\ y_C = \frac{\beta^3}{N^2} = -\frac{k^3 \alpha^3}{N^2}. \end{cases}$$

Using

$$\sin \theta_2 = \frac{\alpha}{\sqrt{D^2 + \alpha^2}},$$

we have

$$k = \frac{\sqrt{D^2 - (n^2 - 1)\alpha^2}}{n\alpha},$$

and we can derive the position of the image as para-

metric functions w.r.t.  $\alpha$ :

$$\begin{cases} x_C = (n^2 - 1) \frac{\alpha^3}{D^2}, \\ y_C = -\frac{n^2}{D^2} \frac{\alpha^3}{n^3 \alpha^3} \{D^2 - (n^2 - 1)\alpha^2\}^{3/2} \\ \quad = -\frac{D}{n} \left\{1 - (n^2 - 1) \frac{\alpha^2}{D^2}\right\}^{3/2}. \end{cases}$$

### 3 POV under water

If the object is in the air, at height of  $D$  above the boundary, and the POV is under water, the relative index of refraction is  $1/n < 1$ , and the similar reasoning leads to the equation

$$-|\xi|^{2/3} + |\eta|^{2/3} = 1,$$

with  $\xi = \frac{x}{W}$  and  $\eta = \frac{y}{Z}$ , where  $W = \frac{nD}{\sqrt{1 - n^2}}$  and  $Z = nD$ .

I could not find the name for this shape of the curve. The astroid is a member of the family of curves named *superellipse*, which is defined by

$$|\xi|^n + |\eta|^n = 1.$$

Asteroid is the case where  $n = 2/3$ . But as far as I know there's no name for the family of curves in the form

$$|\xi|^n - |\eta|^n = \pm 1,$$

which could be named *super-hyperbola*, although the repeating cognates may bother you<sup>2</sup>.

<sup>1</sup>Not to be confused with *asteroid*.

<sup>2</sup>I might suggest *superbola*.