

Vertical shear instability in protoplanetary disks

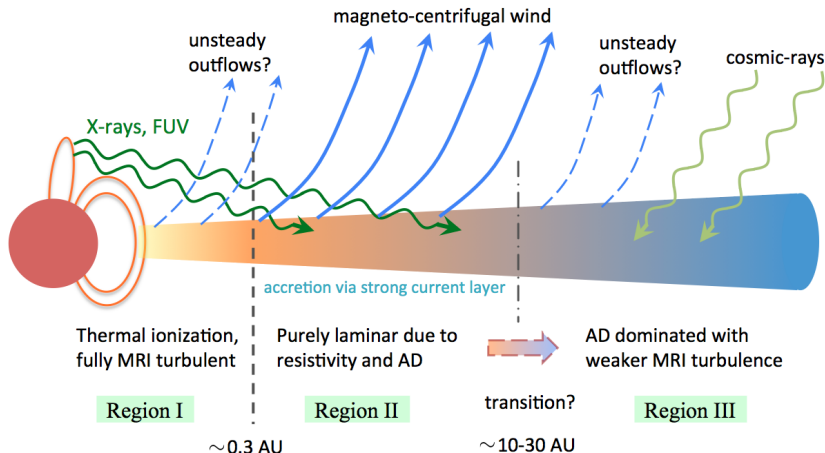
Min-Kai Lin
minkailin@email.arizona.edu

Steward Theory Fellow
University of Arizona

June 18 2015

Transport in protoplanetary disks

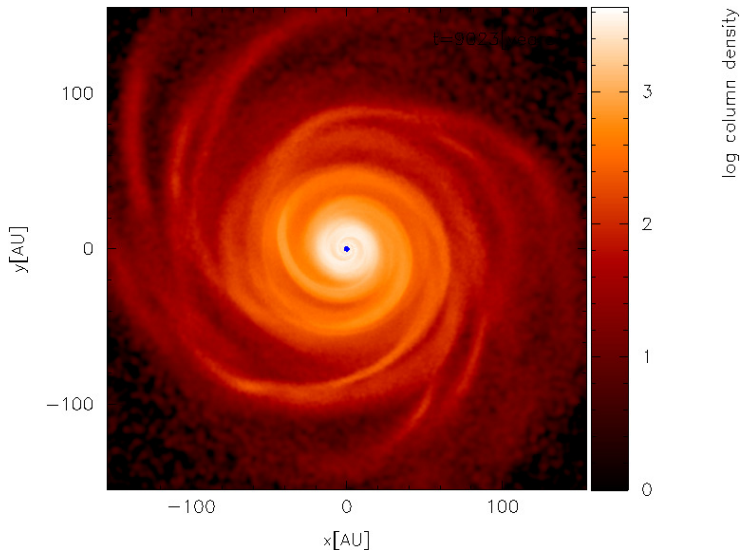
- Magnetic



(Bai, 2013)

Transport in protoplanetary disks

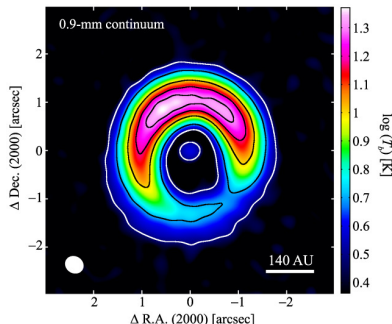
- Self-gravity



(Forgan et al., 2011)

Hydrodynamic instabilities

- Vortices



(Fukagawa et al., 2013)

- 'Rossby wave instability' (Lovelace et al., 1999; Li et al., 2000): Kelvin-Helmholtz in a disk
- 'Convective overstability' (Klahr & Hubbard, 2014; Lyra, 2014): growing epicycles
- 'Baroclinic instability' (Lesur & Papaloizou, 2009): non-linear amplification of vortices

Baroclinic disks and vertical shear

- Astrophysical disks are generally *baroclinic* ($\nabla P \times \nabla \rho \neq 0$) so that

$$\frac{\partial \Omega}{\partial z} \neq 0.$$

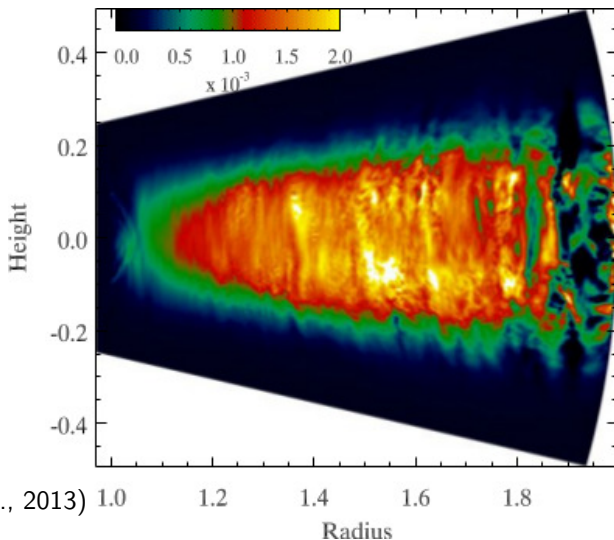
- Example: vertically isothermal thin-disk with radial temperature gradient $T \propto r^q$,

$$r \frac{\partial \Omega}{\partial z} \simeq \frac{1}{2} h q \Omega_{\text{Kep}} \left(\frac{z}{H} \right)$$

$h = H/r$: disk aspect-ratio; Ω_{Kep} : Keplerian rotation

- Shear flow \Rightarrow free energy \Rightarrow instability?

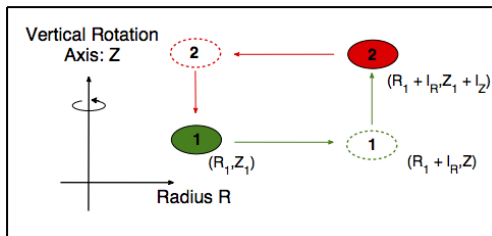
The 'vertical shear instability'



(Nelson et al., 2013)

Hydrodynamic turbulence.

Basic physics of the VSI



(Umurhan et al., 2013)

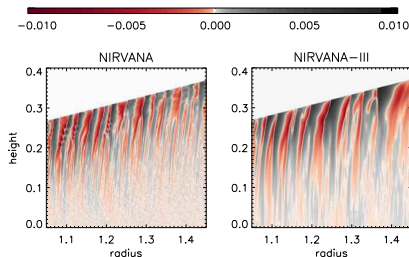
- Change in kinetic energy:

$$\Delta E \sim I_r^2 \left(\Omega^2 + \frac{I_z}{I_r} \cdot r \frac{\partial \Omega^2}{\partial z} \right).$$

- Vertical shear is weak, **BUT**

$$\Delta E < 0 \text{ is possible if } |I_z| \gg |I_r|,$$
$$\Rightarrow \text{ **INSTABILITY!** }$$

Basic physics of the VSI



(Nelson et al., 2013)

- Change in kinetic energy:

$$\Delta E \sim I_r^2 \left(\Omega^2 + \frac{I_z}{I_r} \cdot r \frac{\partial \Omega^2}{\partial z} \right).$$

- Vertical shear is weak, **BUT**

$$\Delta E < 0 \quad \text{is possible if} \quad |I_z| \gg |I_r|, \\ \Rightarrow \quad \mathbf{INSTABILITY!}$$

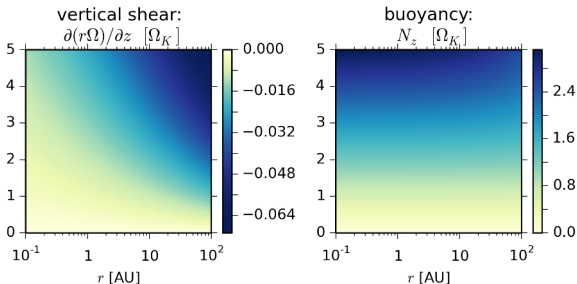
Role of buoyancy

Vertical motion associated with VSI is opposed by buoyancy forces

$\underbrace{r\partial_z\Omega}$
destabilizing vert. shear

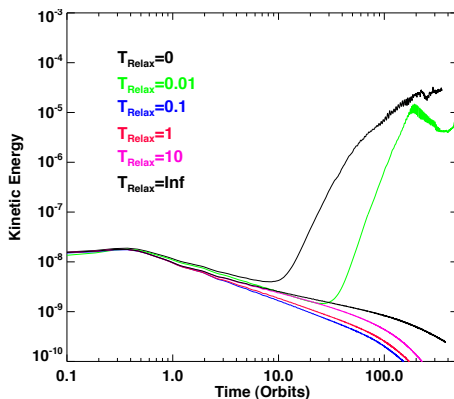
v.s.

$\underbrace{N_z}$
stabilizing vert. buoyancy



- Vertical shear is *weak*, $r\partial_z \ln \Omega \sim O(h) \ll 1$
- Vertical buoyancy is *strong*, $N_z/\Omega \sim O(1)$

The need for rapid cooling



(Nelson et al., 2013)

- Buoyancy ineffective if cooling times are short \rightarrow VSI can operate
- Stoll & Kley (2014): VSI in radiation-hydrodynamic simulations only when external heating included

Is there a quantitative thermodynamic criteria?

Previous stability analyses and our contribution

- Vertically and radially local analyses: (Urpin & Brandenburg, 1998; Urpin, 2003)
- Vertically global, radially local analyses, no buoyancy: (Nelson et al., 2013; Barker & Latter, 2015)

Previous stability analyses and our contribution

- Vertically and radially local analyses: (Urpin & Brandenburg, 1998; Urpin, 2003)
- Vertically global, radially local analyses, no buoyancy: (Nelson et al., 2013; Barker & Latter, 2015)

Lin & Youdin (2015)

- Vertically global, radially local, including energy equation (i.e. with buoyancy)
- Both constant cooling and realistic cooling functions

Linear theory: simplified model

- Axisymmetric perturbations in a vertically isothermal disk
- Wave-ansatz radial dependence with wavenumber k_x , low-frequency limit
 $|\omega| \ll \Omega_{\text{Kep}}$
- Vertically constant cooling time, $t_c \equiv \beta \Omega_{\text{Kep}}^{-1}$

Linear theory: simplified model

- Axisymmetric perturbations in a vertically isothermal disk
- Wave-ansatz radial dependence with wavenumber k_x , low-frequency limit $|\omega| \ll \Omega_{\text{Kep}}$
- Vertically constant cooling time, $t_c \equiv \beta \Omega_{\text{Kep}}^{-1}$

Reduction to single ODE

$$0 = \delta v_z''(z) - zA\delta v_z'(z) + (B - Cz^2) \delta v_z(z).$$

- Hermite differential equation after transformation
- Dispersion relation $\omega = \omega(k_x; \beta)$ for the frequency

Linear theory: simplified model

- Axisymmetric perturbations in a vertically isothermal disk
- Wave-ansatz radial dependence with wavenumber k_x , low-frequency limit $|\omega| \ll \Omega_{\text{Kep}}$
- Vertically constant cooling time, $t_c \equiv \beta \Omega_{\text{Kep}}^{-1}$

Reduction to single ODE

$$0 = \delta v_z''(z) - zA\delta v_z'(z) + (B - Cz^2) \delta v_z(z).$$

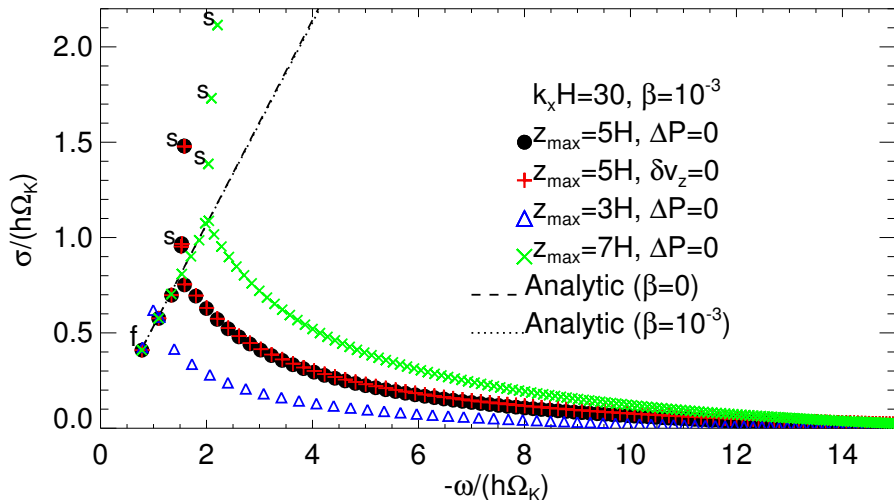
- Hermite differential equation after transformation
- Dispersion relation $\omega = \omega(k_x; \beta)$ for the frequency

Seek $\text{Im}(\omega) = 0$ for large k_x to find VSI requires

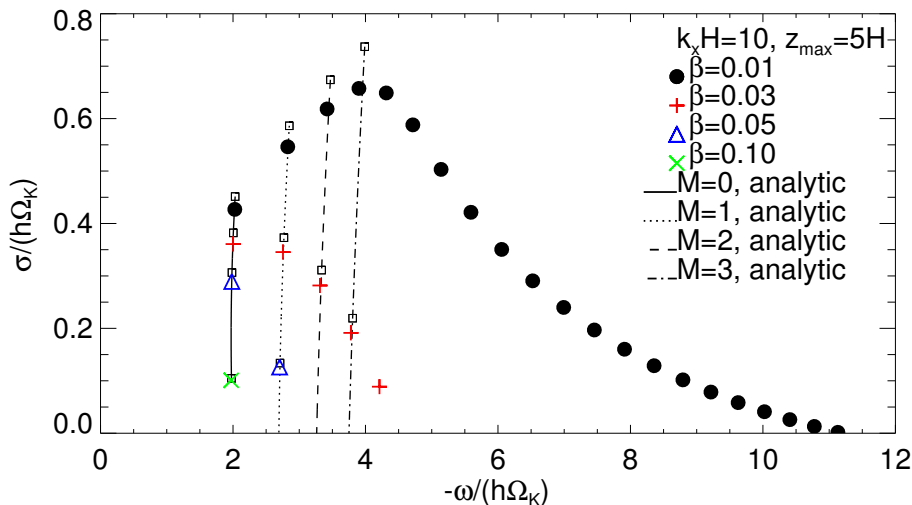
$$t_c \Omega_{\text{Kep}} < \frac{h|q|}{\gamma - 1} \equiv \beta_{\text{crit}}.$$

- $h|q|$: vertical shear
- $\gamma - 1$: vertical buoyancy
- $\beta_{\text{crit}} \ll 1$, i.e. rapid cooling required

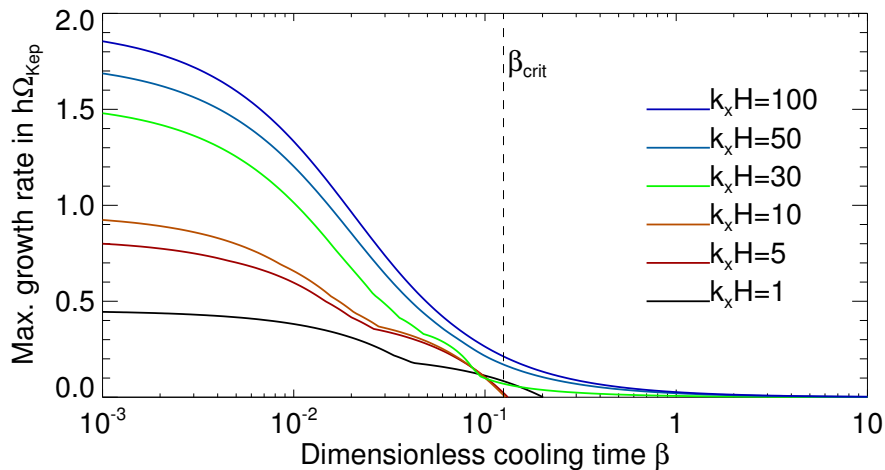
Linear theory: numerical treatment



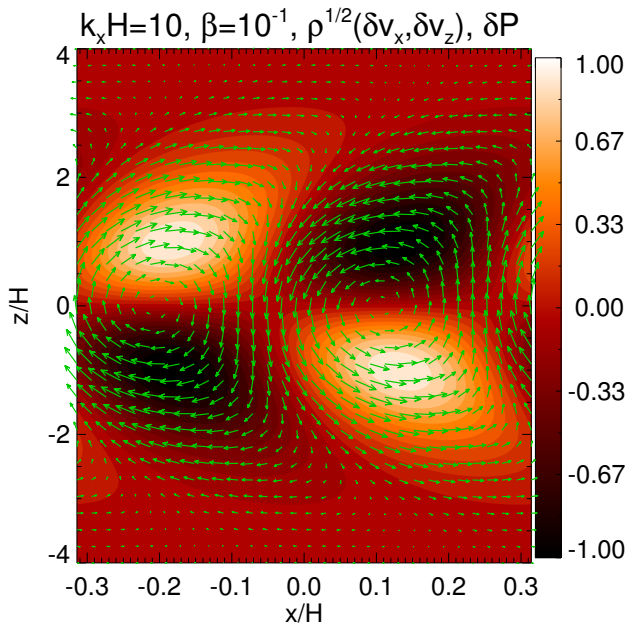
Effect of increasing the cooling time



Testing the critical cooling timescale

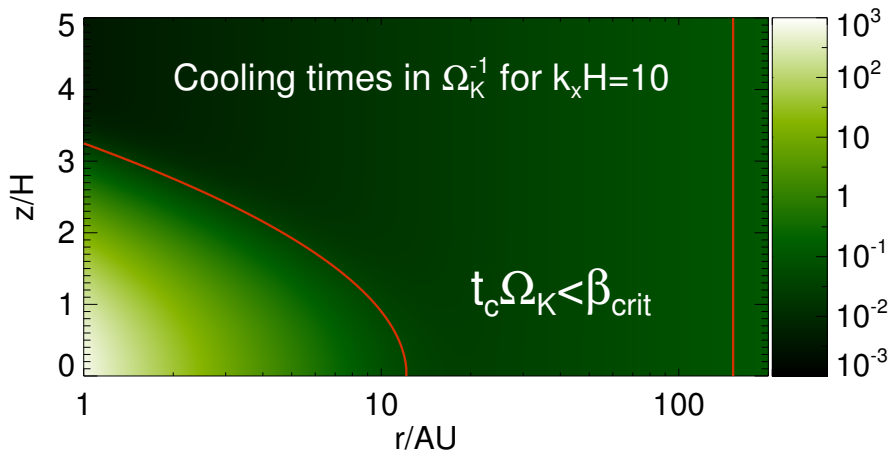


Visualization



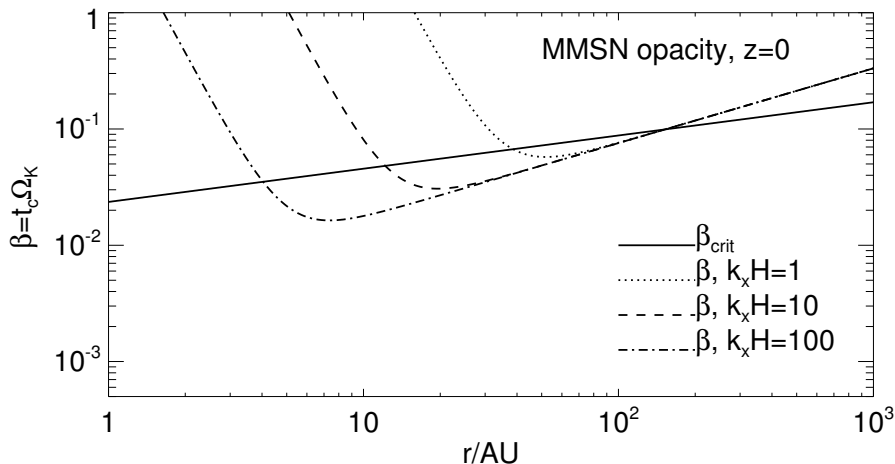
Application to protoplanetary disks

Estimate cooling times in the Minimum Mass Solar Nebula (Chiang & Youdin, 2010) based on dust opacity:



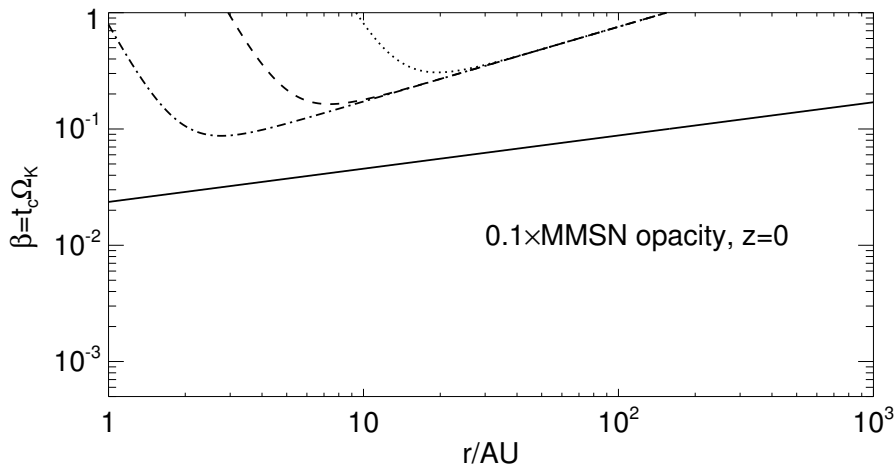
Application to protoplanetary disks

Estimate cooling times in the Minimum Mass Solar Nebula (Chiang & Youdin, 2010) based on dust opacity:

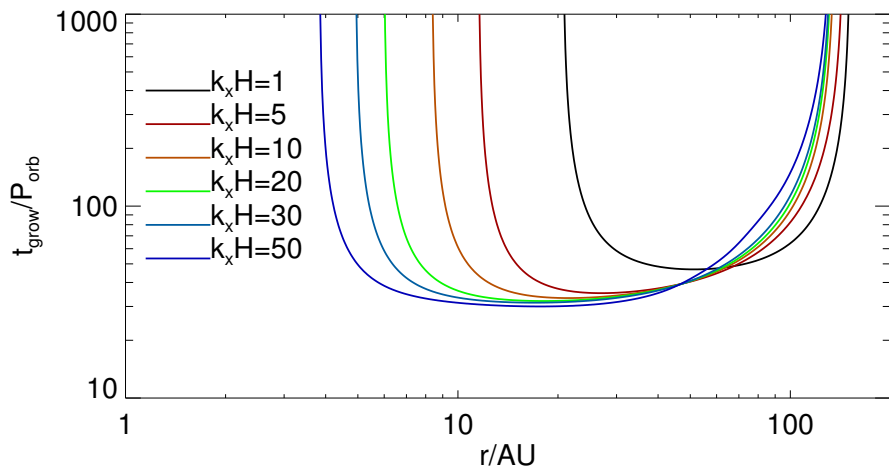


Application to protoplanetary disks

Estimate cooling times in the Minimum Mass Solar Nebula (Chiang & Youdin, 2010) based on dust opacity:



VSI in the Solar Nebula



Conclusions

- Astrophysical disks generally possess vertical shear
- Unstable if buoyancy ineffective: $N_z = 0$ and/or $t_c \Omega_{\text{Kep}} < \beta_{\text{crit}} \ll 1$
- Fast cooling needed because vertical shear is weak but buoyancy is strong
- Stringent thermodynamic requirement satisfied at 10s of AU in typical PPDs

References

- Bai X.-N., 2013, ApJ, 772, 96
- Barker A. J., Latter H. N., 2015, ArXiv e-prints
- Chiang E., Youdin A. N., 2010, Annual Review of Earth and Planetary Sciences, 38, 493
- Forgan D., Rice K., Cossins P., Lodato G., 2011, MNRAS, 410, 994
- Fukagawa M., Tsukagoshi T., Momose M., Saigo K., Ohashi N., Kitamura Y., Inutsuka S.-i., Muto T., Nomura H., Takeuchi T., Kobayashi H., Hanawa T., Akiyama E., Honda M., Fujiwara H., Kataoka A., Takahashi S. Z., Shibai H., 2013, PASJ, 65, L14
- Klahr H., Hubbard A., 2014, ApJ, 788, 21
- Lesur G., Papaloizou J. C. B., 2009, A&A, 498, 1
- Li H., Finn J. M., Lovelace R. V. E., Colgate S. A., 2000, ApJ, 533, 1023
- Lin M.-K., Youdin A., 2015, ArXiv e-prints
- Lovelace R. V. E., Li H., Colgate S. A., Nelson A. F., 1999, ApJ, 513, 805
- Lyra W., 2014, ApJ, 789, 77
- Nelson R. P., Gressel O., Umurhan O. M., 2013, MNRAS, 435, 2610
- Stoll M. H. R., Kley W., 2014, A&A, 572, A77
- Umurhan O. M., Nelson R. P., Gressel O., 2013, in European Physical Journal Web of Conferences Vol. 46 of European Physical Journal Web of Conferences, Breathing Life Into Dead-Zones. p. 3003
- Urpın V., 2003, A&A, 404, 397
- Urpın V., Brandenburg A., 1998, MNRAS, 294, 399