

Hydrodynamic activity in protoplanetary disks

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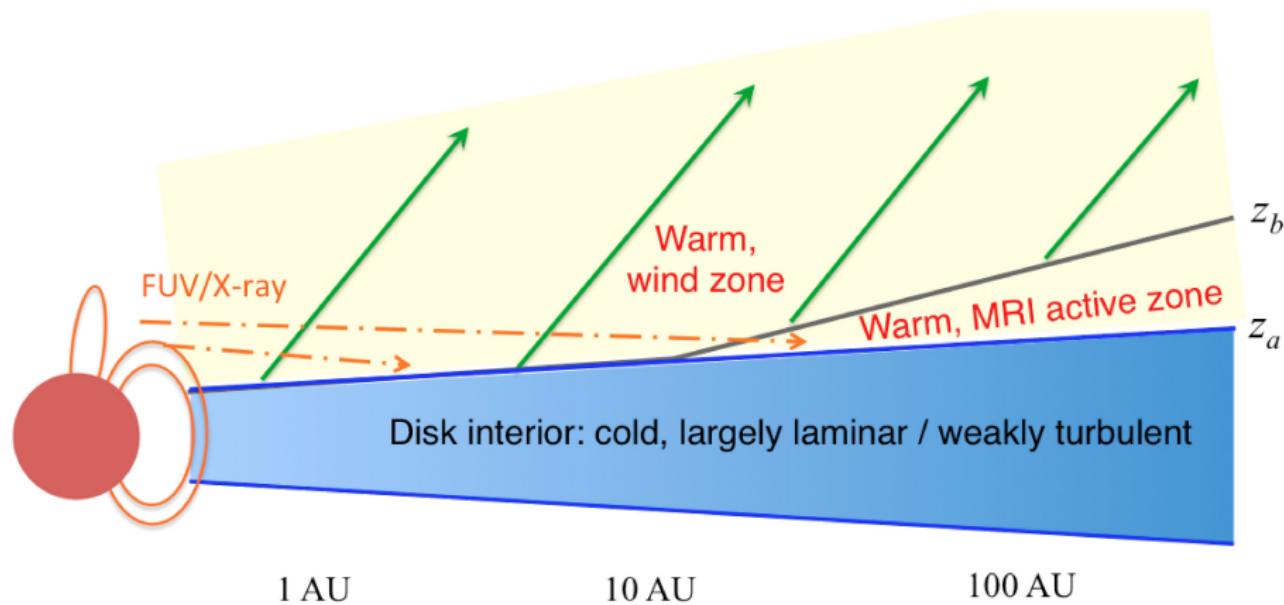
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Interests

- Astrophysical fluid dynamics of accretion/protoplanetary disks
- Disk-planet interactions, orbital migration
- Self-gravitating disks
- Hydrodynamic instabilities
- Magneto-hydrodynamics
- Non-linear numerical simulations
- Analytical/semi-analytical methods

Role of magnetic fields in protoplanetary disks



(Bai, 2016)

- Purely hydrodynamical effects are important in the planet-forming disk bulk

Hydrodynamical processes in protoplanetary disks

- Vortex formation and evolution

- ▶ Vortices as dust traps
- ▶ Explaining asymmetric transition disks
- ▶ Assisting planetesimal formation
- ▶ 3D effects in massive disks (Lin & TBC, in prep.)

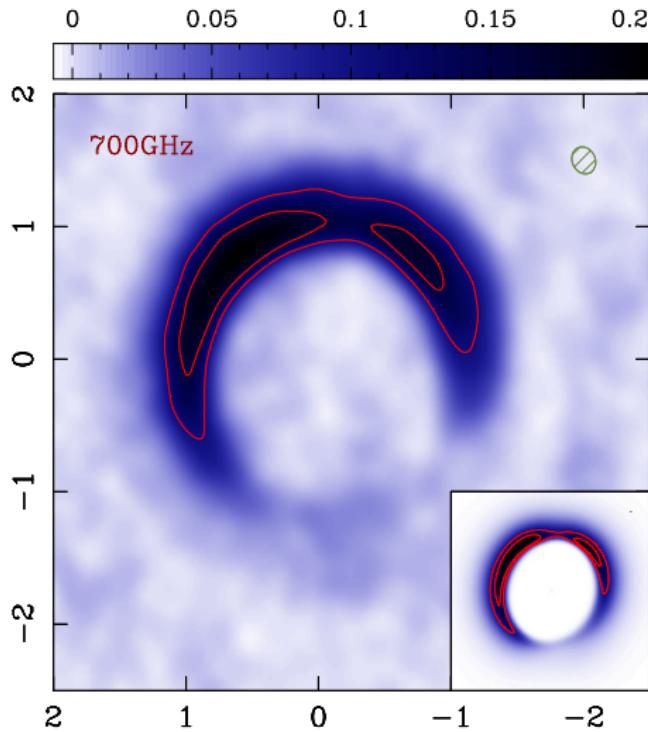
- Vertical shear instability

- ▶ Hydrodynamic turbulence
- ▶ Angular momentum transport
- ▶ Does it occur in realistic PPDs? (Lin & Youdin, 2015)

- Generalized gravitational instabilities

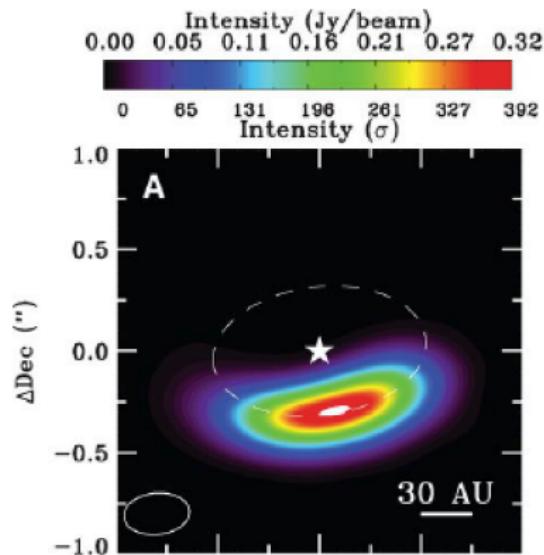
- ▶ 'Gravito-turbulence' and transport
- ▶ Fragmentation to form giant planets/sub-stellar companions
- ▶ Going beyond Toomre and Lin-Shu analyses (Lin & Kratter, re-submitted)

Transition disk asymmetries: vortices?



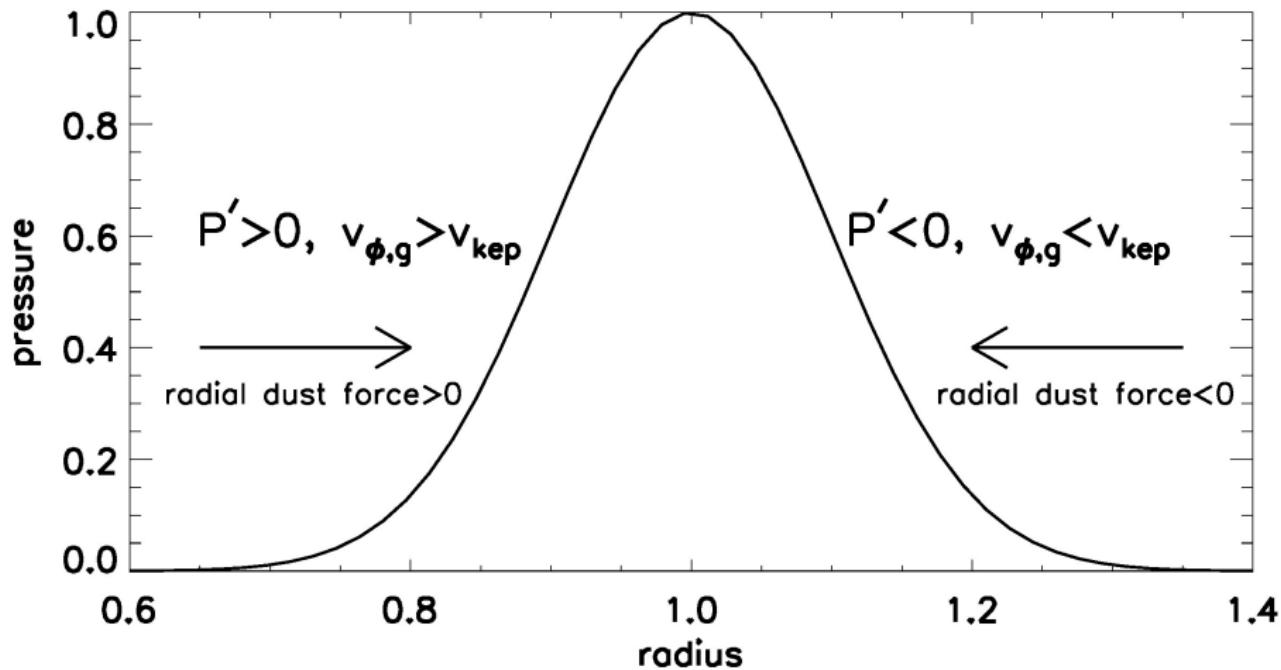
(HD142527, Casassus et al., 2015)

Transition disk asymmetries: vortices?



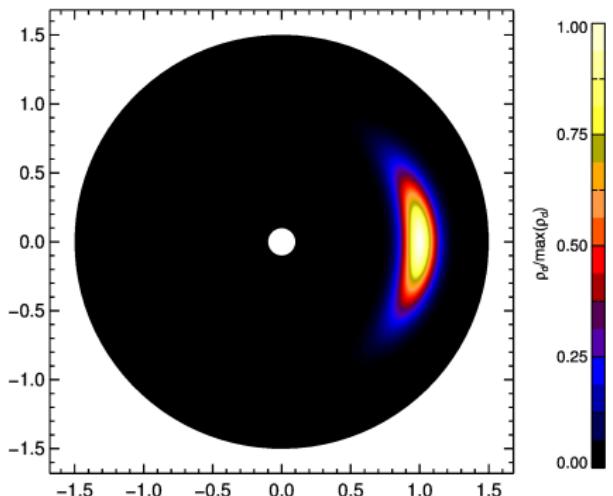
(Oph IRS 48, van der Marel et al., 2013)

Dust trapping at pressure maxima



- Drag forces cause dust to accumulate at pressure bumps

Dust distribution in disk vortices



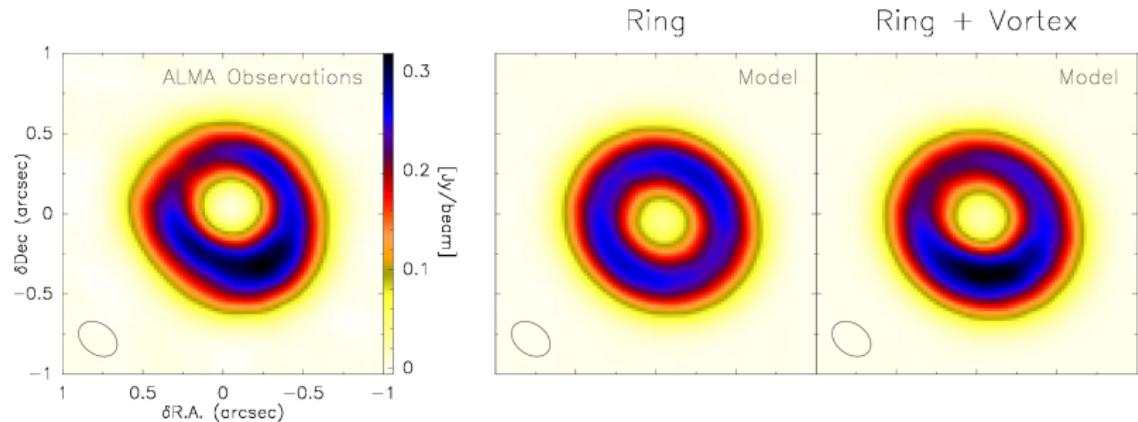
$$\rho_d(a) \propto \exp\left(-\frac{a^2}{2H_v^2}\right),$$

- a : distance from the vortex center
(Lyra & Lin, 2013)

$$H_v(\chi, \delta, \text{St}) = \frac{H_g}{f(\chi)} \sqrt{\frac{\delta}{\delta + \text{St}}}.$$

- χ : vortex aspect-ratio
- δ : turbulence in the vortex
- St: Stokes number (dust-gas friction)
- H_g : gas scale height

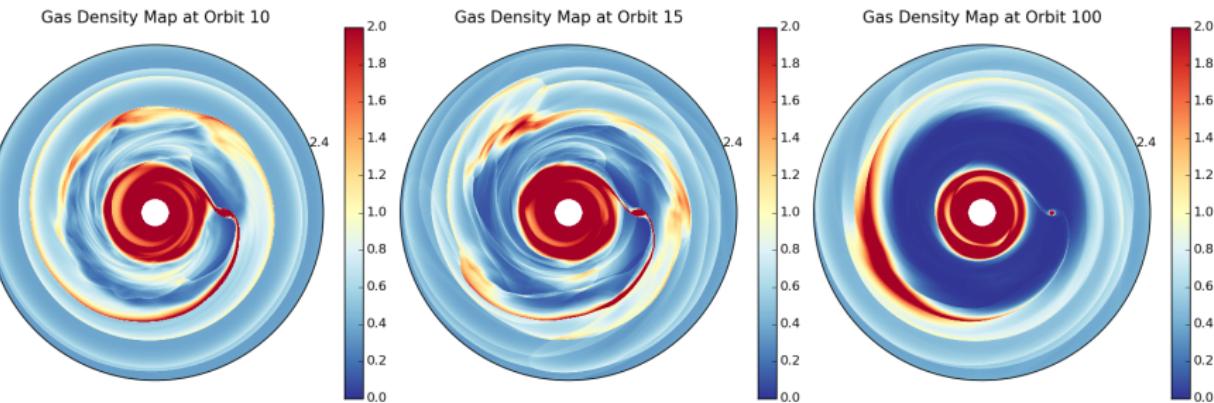
Application to observations



(SAO 206462, Pérez et al., 2014)

$\chi_{\text{obs}} \sim 7$, model+ data $\rightarrow v_{\text{turb}} \sim 0.22 c_s$.

Gap edges as sites for vortex formation



(2D simulations by M. Hammer, 'Vortices and orbital migration')

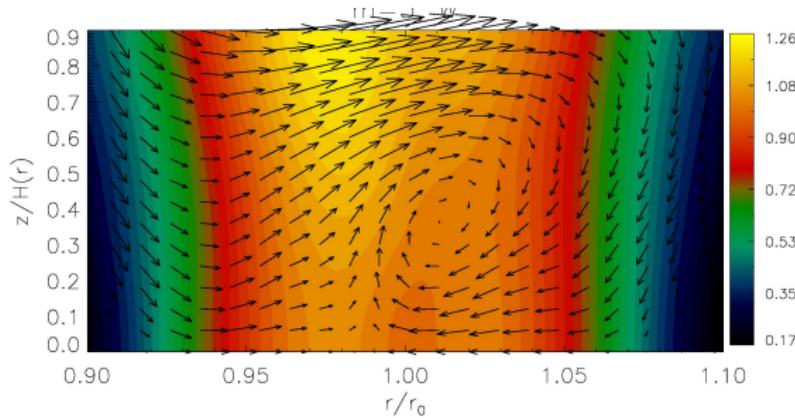
- Disk-planet interaction → gaps
 - Surface density maxima at gap edges, or PV minima due to strong shear
 - Rossby wave instability → edges 'roll up' into vortices
- (Li et al., 2001)

Gap edges as sites for vortex formation



Rossby vortices are vertically global

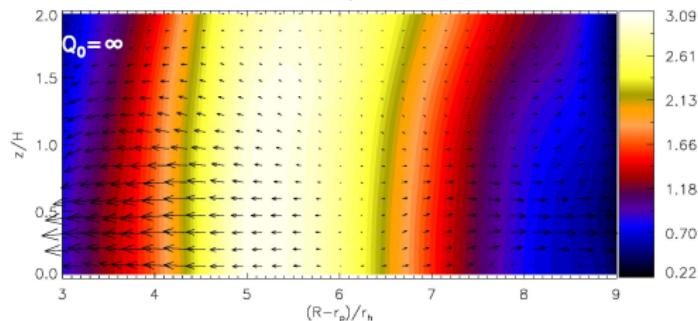
$$\Delta P/\rho$$



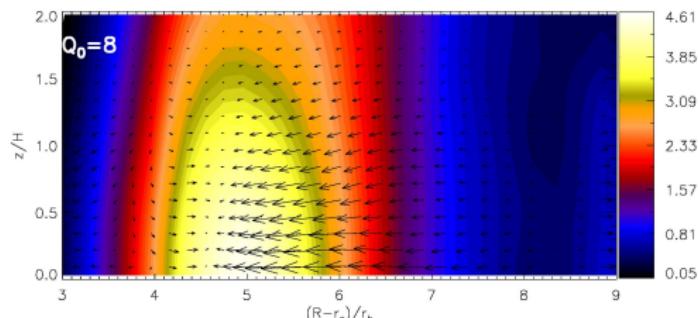
- 3D linear theory
(Lin, 2013b)

Rossby vortices are vertically global

$$\Delta\rho/\rho$$



↑ NO self-gravity

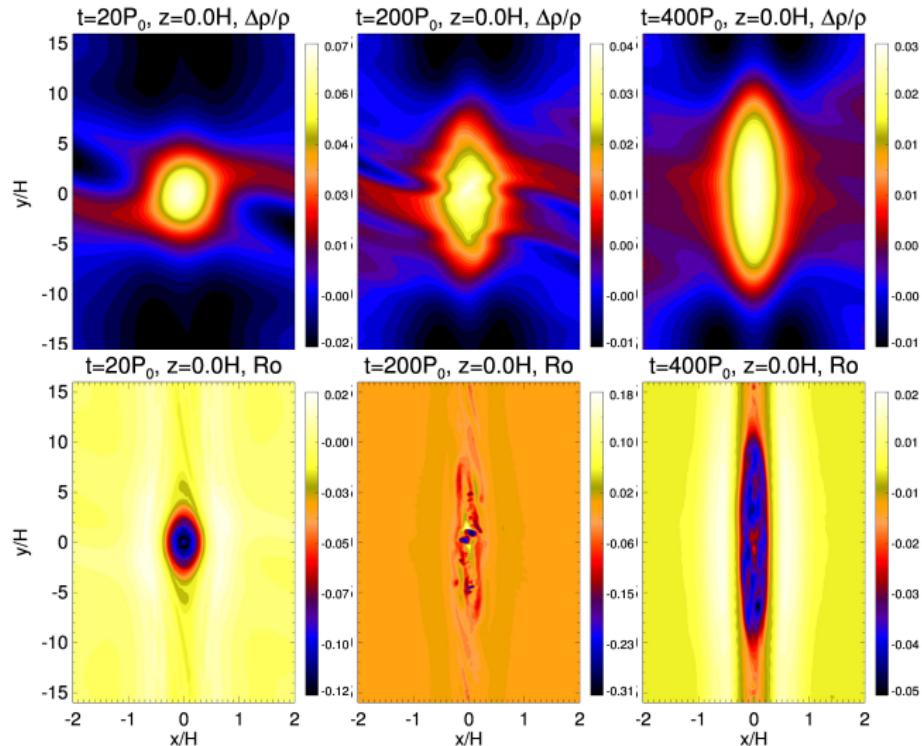


↑ WITH self-gravity

- Global 3D Zeus simulations (Lin, 2012)

Elliptic instability of 3D vortices

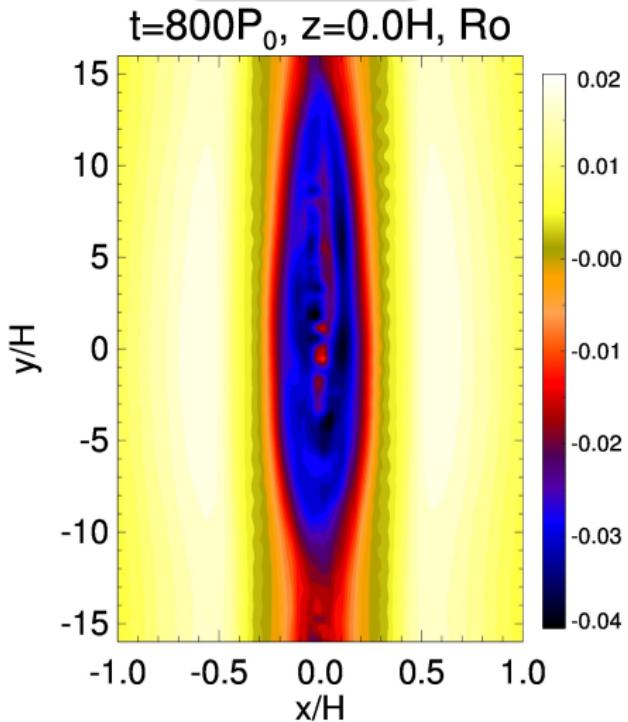
- A 3D instability that weakens/destroys vortices (Lesur & Papaloizou, 2009)



- Athena simulations (Lin & TBC, in prep.)

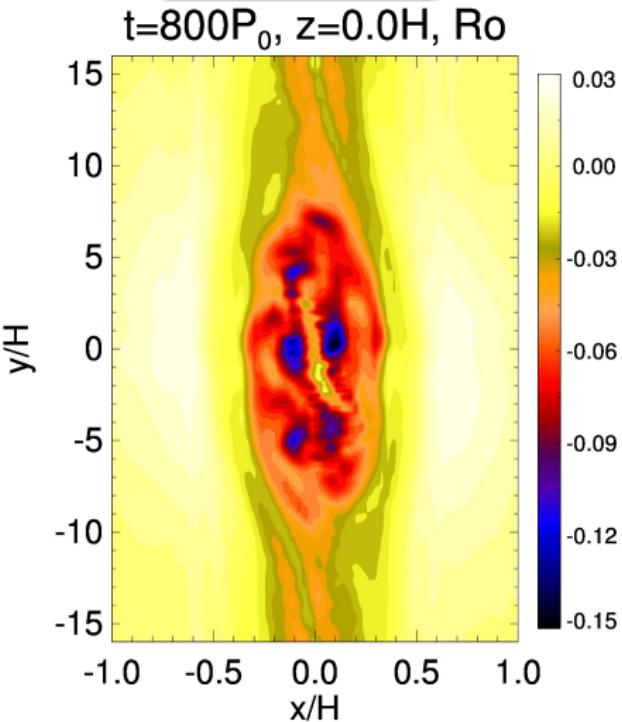
Elliptic instability & self-gravity

LIGHT disk



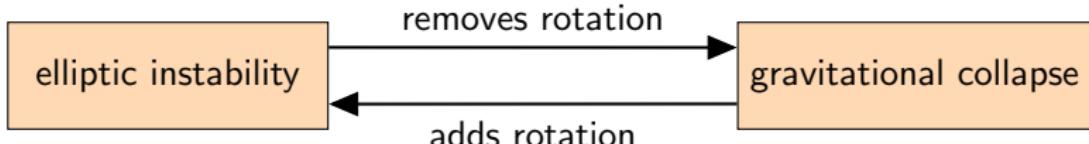
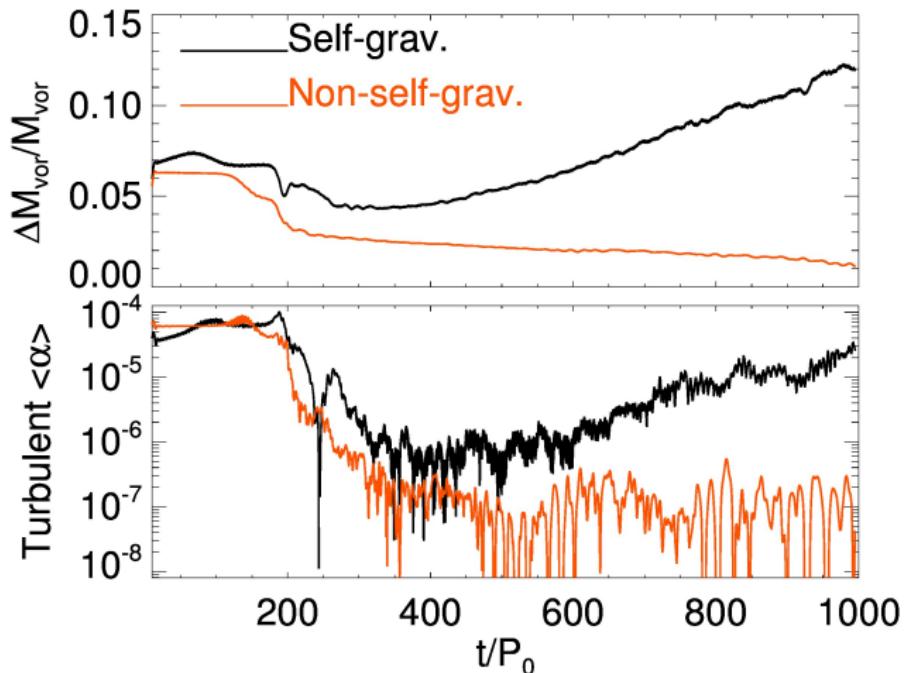
$$(Q = 100)$$

MASSIVE disk



$$(Q = 3)$$

Elliptic instability & self-gravity



Astrophysical disks have vertical shear

- Accretion disks are generally baroclinic, $\nabla P \times \nabla \rho \neq 0$

$$\Rightarrow \frac{\partial \Omega}{\partial z} \neq 0$$

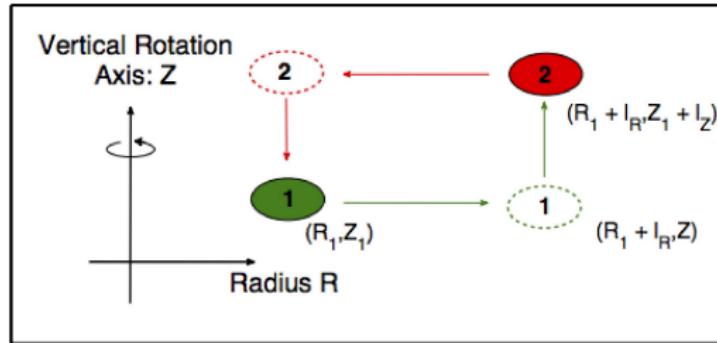
- Vertically isothermal thin-disk with $T \propto r^q$,

$$r \frac{\partial \Omega}{\partial z} \simeq \left(\frac{qz}{2H} \right) \times \frac{H}{r} \Omega_{\text{Kep}}$$

- $O(H/r)$ effect

Vertical shear instability

$\partial_z \Omega \neq 0 \Rightarrow$ free energy \rightarrow instability?



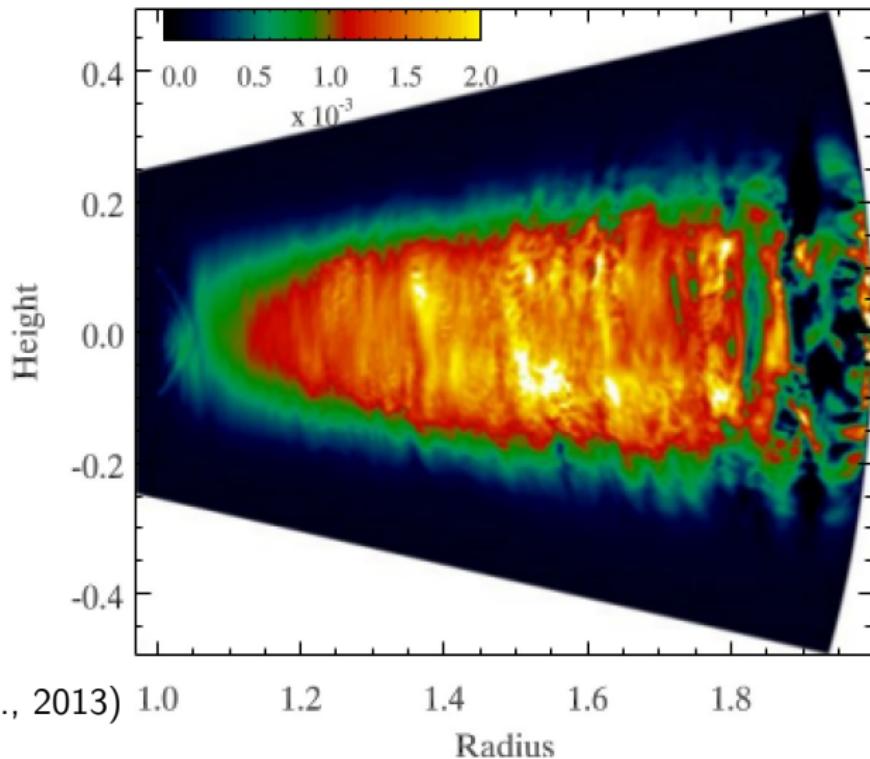
(Umurhan et al., 2013)

- Change in kinetic energy:

$$\Delta E \sim I_r^2 \left(\Omega^2 + \frac{l_z}{I_r} \cdot r \frac{\partial \Omega^2}{\partial z} \right).$$

$\Delta E < 0 \quad \text{if} \quad |l_z| \gg |I_r| \Rightarrow \text{INSTABILITY}$

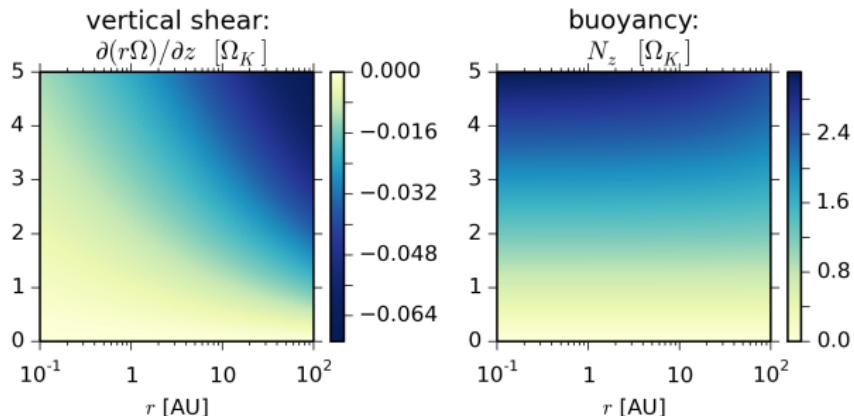
VSI: non-linear outcome



(Nelson et al., 2013) 1.0 1.2 1.4 1.6 1.8
Radius

Hydrodynamic turbulence

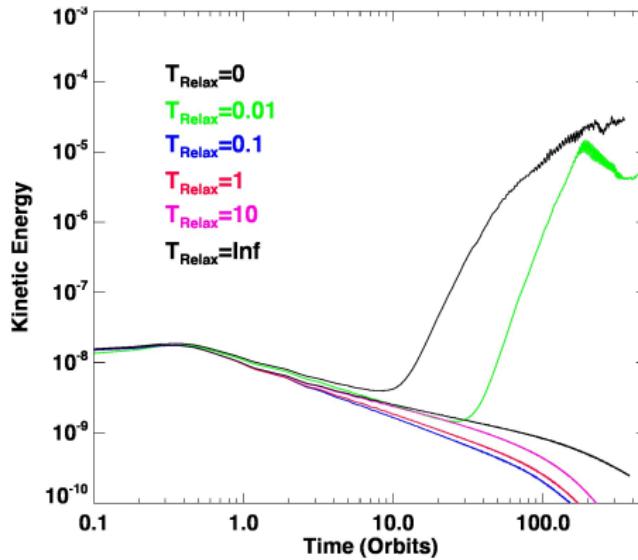
VSI needs to fight buoyancy in real disks



- Vertical shear is weak, $r\partial_z\ln\Omega \sim O(h) \ll 1$ (so need $I_z/I_r \gg 1$)
- Vertical buoyancy is strong, $N_z/\Omega \sim O(1)$

Ultra-fast cooling can overcome buoyancy forces

- Cooling parameterization: $\partial_t T = -(T - T_0)/t_{\text{cool}}$



(Nelson et al., 2013)

Can we quantify this requirement?

Lin & Youdin (2015): linear theory with finite cooling

From single ODE, reduced model for $T \propto r^q$ disk, find that VSI requires:

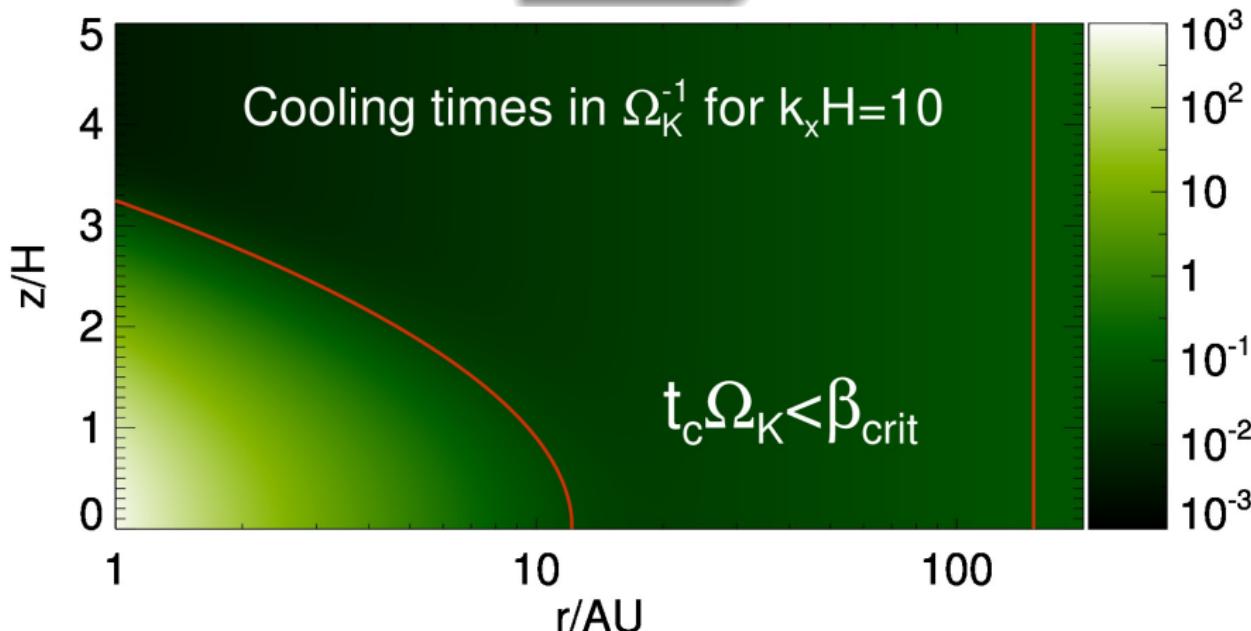
$$t_{\text{cool}}\Omega_K < \frac{h|q|}{\gamma - 1}$$

- $h|q|$: vertical shear ($h \equiv H/r$) — destabilizing
- $\gamma - 1$: vertical buoyancy — stabilizing
- $t_{\text{cool}}\Omega_K \ll 1$ required, i.e. rapid cooling

Vertical shear instability in the Solar Nebula

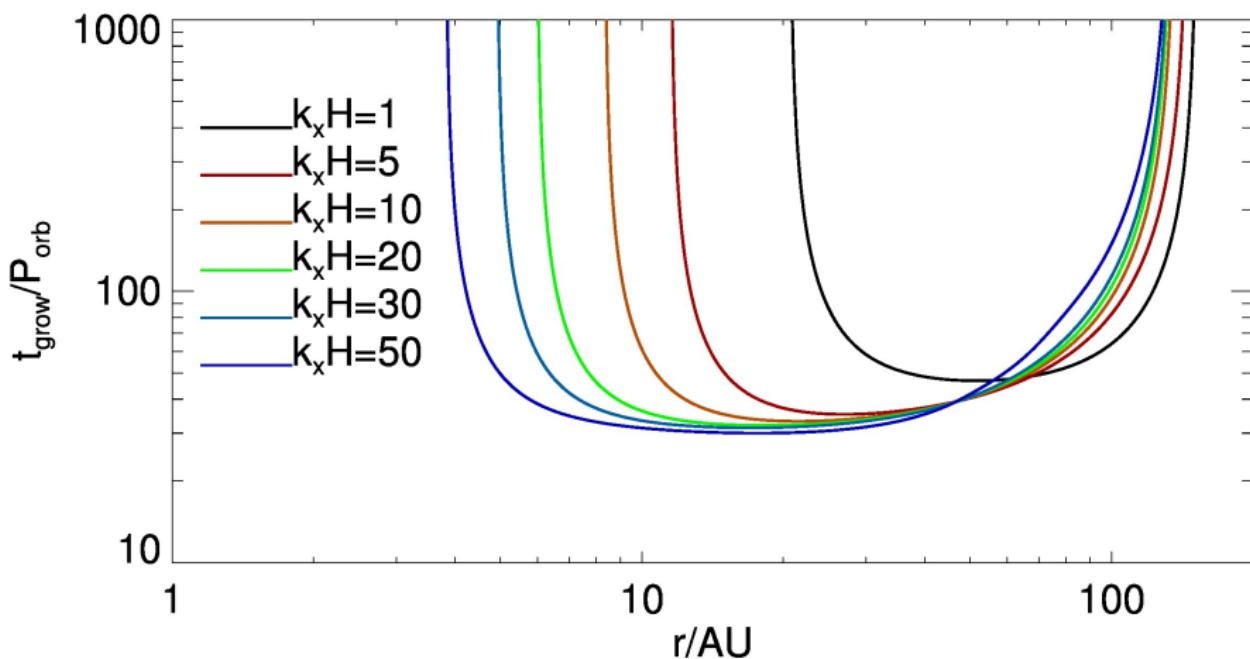
Cooling via dust-opacity ($\propto T^2$) in the Minimum Mass Solar Nebula (Chiang & Youdin, 2010)

$$\beta_{\text{crit}} \equiv \frac{h|q|}{\gamma - 1}$$



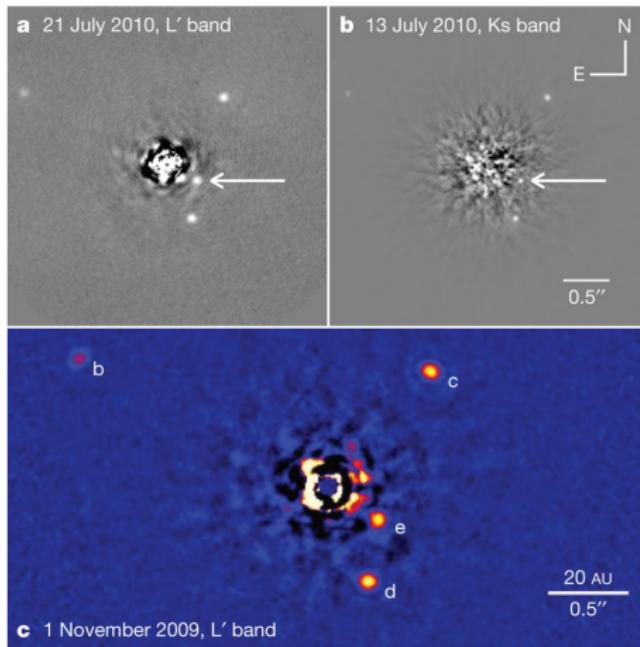
Typical growth times in the Solar Nebula

- Solve the full linearized fluid equations in the radially local approximation, with radiative diffusion/optically-thin cooling



- VSI is most active in the outer disk 10—100AU
- Forced to develop on smaller scales towards inner disk

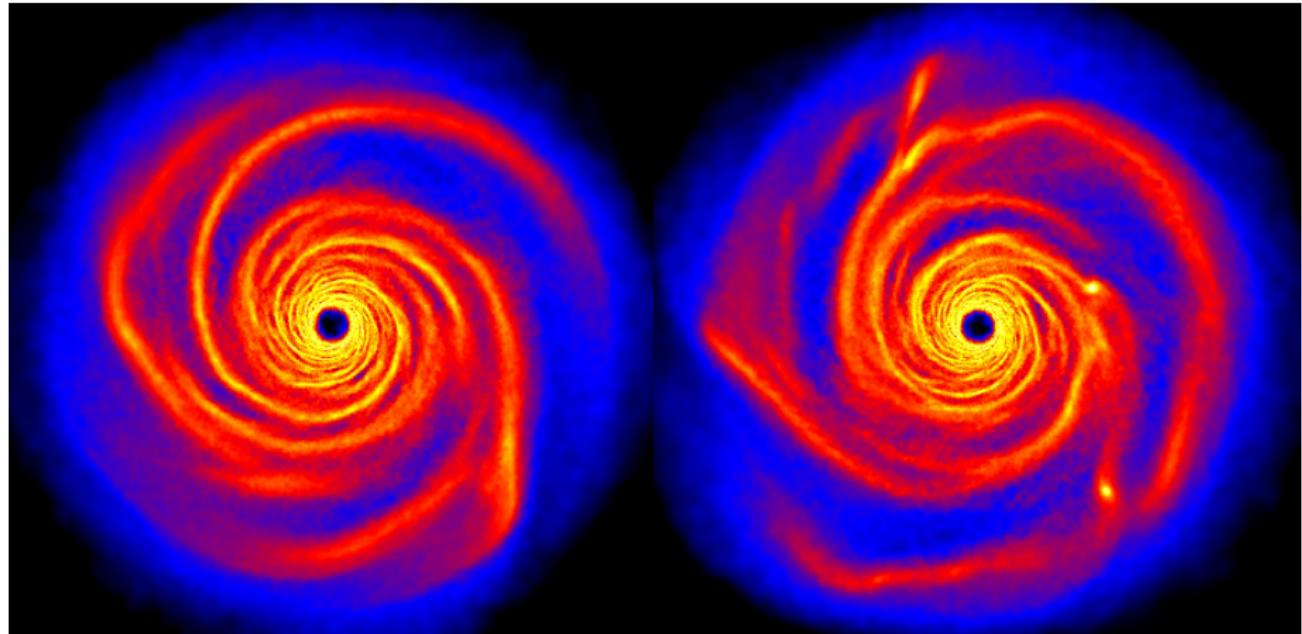
Wide orbit planets



(Marois et al., 2010)

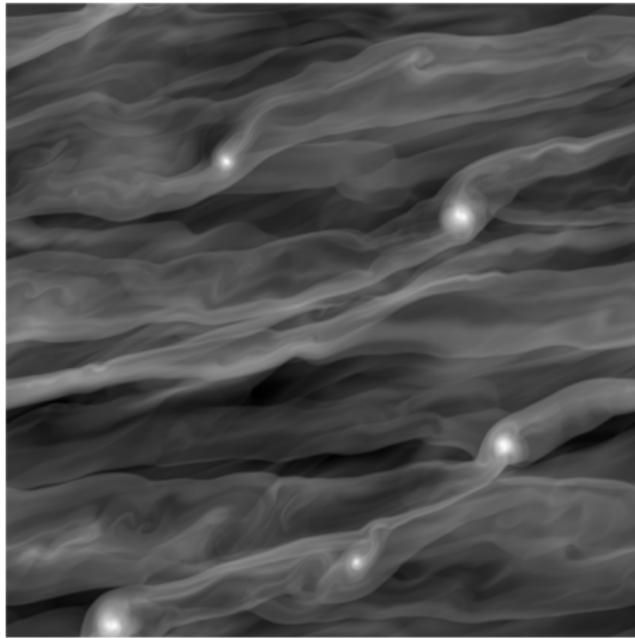
Disk instability theory

- Young, massive protoplanetary disks fragment under its own gravity



(Rice et al., 2005)

When do (simulated) protostellar disks fragment?



(Rice et al., 2011)

- Massive disk

$$Q \equiv \frac{c_s \Omega}{\pi G \Sigma} \lesssim 2$$

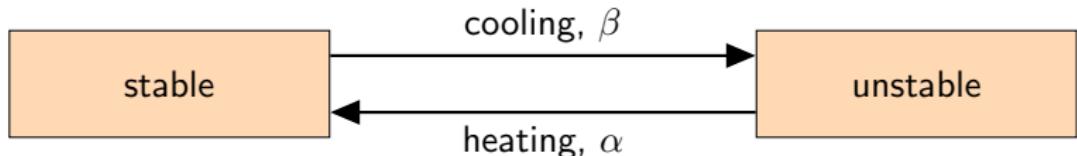
- Fast cooling

$$t_{\text{cool}} \Omega \lesssim 3$$

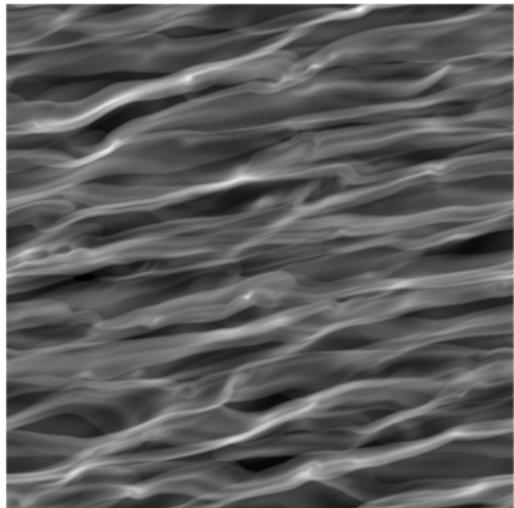
The cooling criterion is empirical.

And when it does not fragment

Gravito-turbulence



- Cooling: $\beta \equiv t_{\text{cool}} \Omega$
- Heating: α viscosity
- Equilibrium $\Rightarrow \alpha = \alpha(\beta)$
- Need $\beta > \beta_c$ or $\alpha < \alpha_c$
But why?



(Rice et al., 2011)

When do realistic protostellar disks fragment?

Work out $\Sigma(R)$, $T(R)$..etc., then ask

- ➊ Is Toomre $Q \sim 1$?
- ➋ Is $t_{\text{cool}}\Omega \sim 1$ or $\alpha \sim 0.1$?

Possible issues:

- Need to choose critical values
- Numerical uncertainties in critical values (e.g. resolution)

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Motivation 1:

Assess disk fragmentation without input from hydrodynamic simulations

Beyond classical gravitational instability

Modern simulations

- Cooling physics, e.g.

$$\frac{\partial E}{\partial t} = -\frac{E}{t_{\text{cool}}}$$

- Turbulent/viscous, e.g.

$$\nu = \alpha \frac{c_s^2}{\Omega}$$

Analytic toolbox

Lin-Shu dispersion relation, Toomre Q

$$\omega^2 = \kappa^2 - 2\pi G \Sigma |k| + c_s^2 k^2$$

$$Q \equiv \frac{c_s \kappa}{\pi G \Sigma}$$

- Isothermal/adiabatic (no cooling)
- Laminar (inviscid)

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Motivation 2:

Generalize analytic treatment of GI to include cooling, irradiation and viscosity

$$\omega = \omega(k; Q, t_{\text{cool}}, \alpha)$$

Quantifying cooling

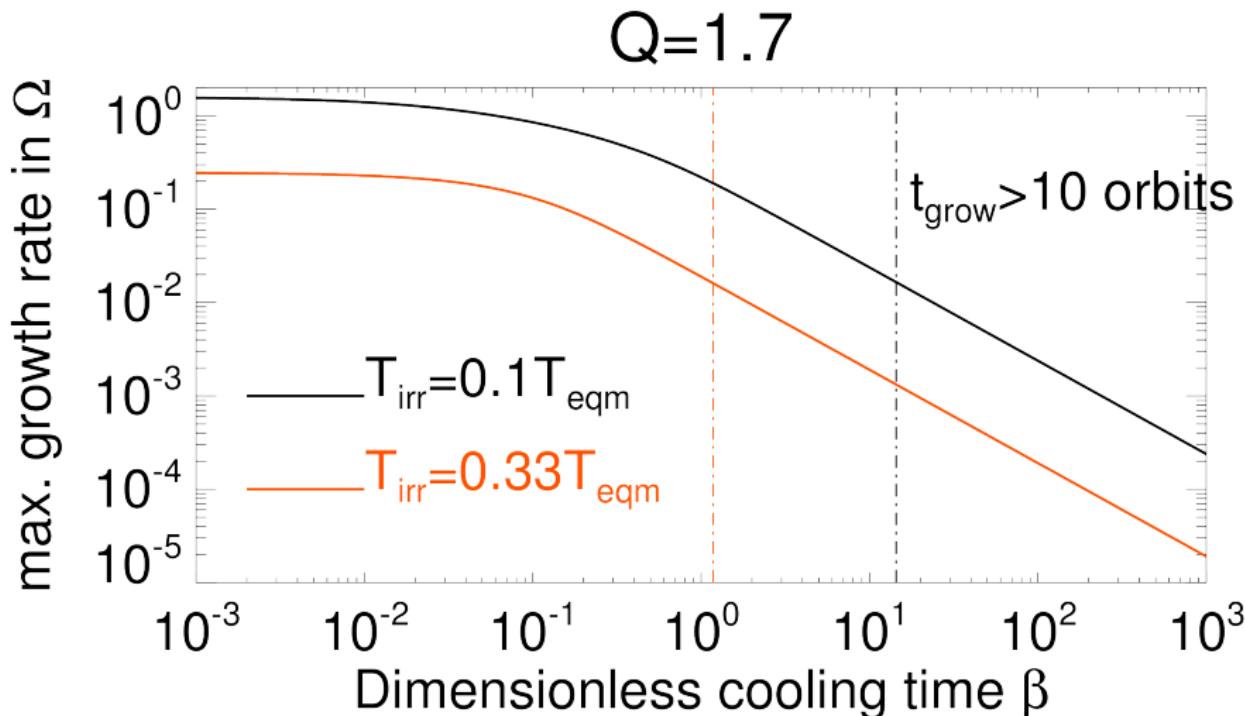
Dispersion relation with cooling

$$\underbrace{s^2}_{\text{growth}} = \underbrace{2\pi G \Sigma |k|}_{\text{+gravity}} - \underbrace{\Omega^2}_{\text{-rotation}} - \underbrace{\left(\frac{T_{\text{irr}}/T + \gamma t_{\text{cool}} s}{1 + t_{\text{cool}} s} \right) c_s^2 k^2}_{\text{-modified pressure}}$$

- T_{irr} : irradiation or floor temperature
- Can be formally unstable for any $t_{\text{cool}} < \infty$
- $T_{\text{irr}} = 0 \sim$ pressureless disk

Just a fancy way to compare compressional heating v.s. thermal losses

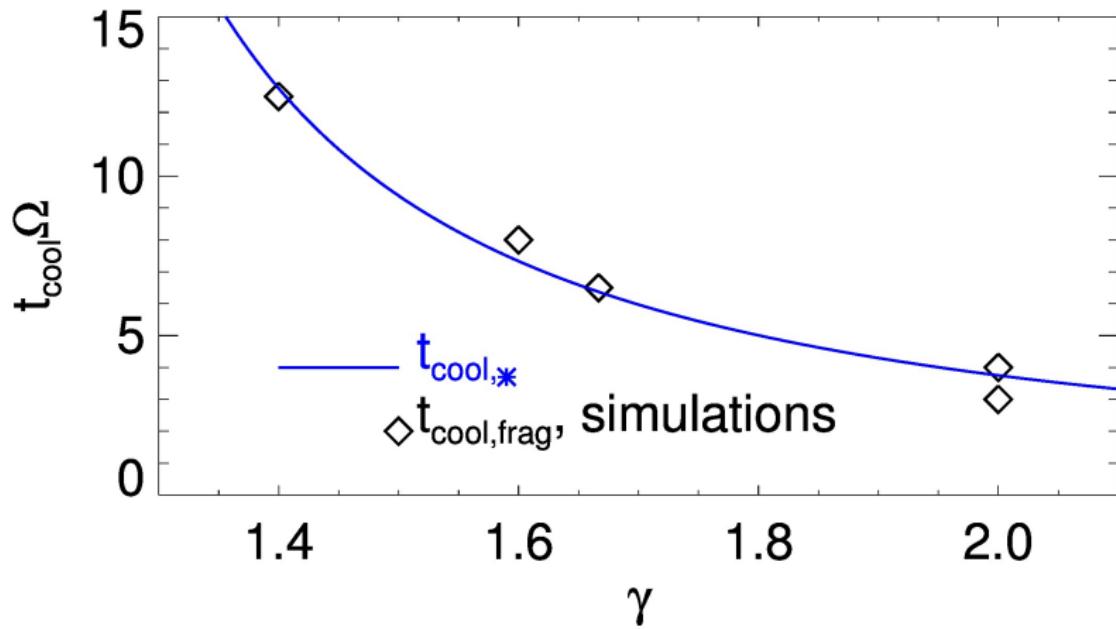
Cooling-driven gravitational instability



Understanding (some) simulations

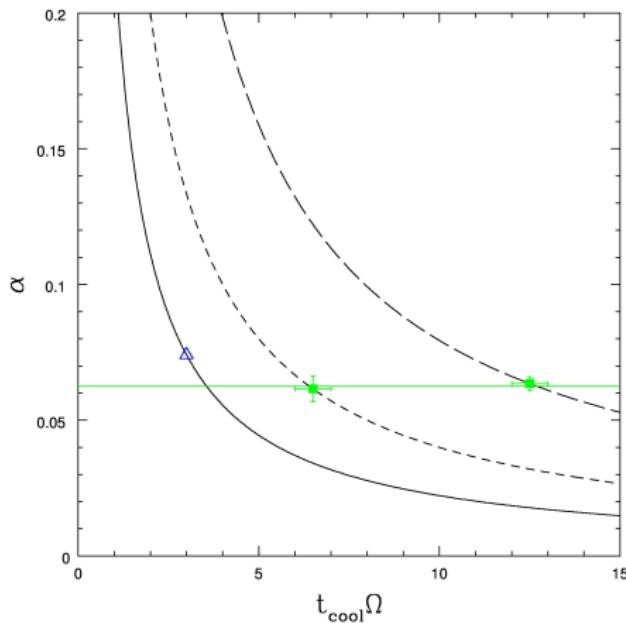
Cooling timescale to remove pressure over a lengthscale $\sim H$

$$t_{\text{cool},*} = (\sqrt{\gamma} - 1)^{-3/2} \Omega^{-1}$$



Simulations: Gammie (2001); Rice et al. (2005, 2011); Paardekooper (2012)

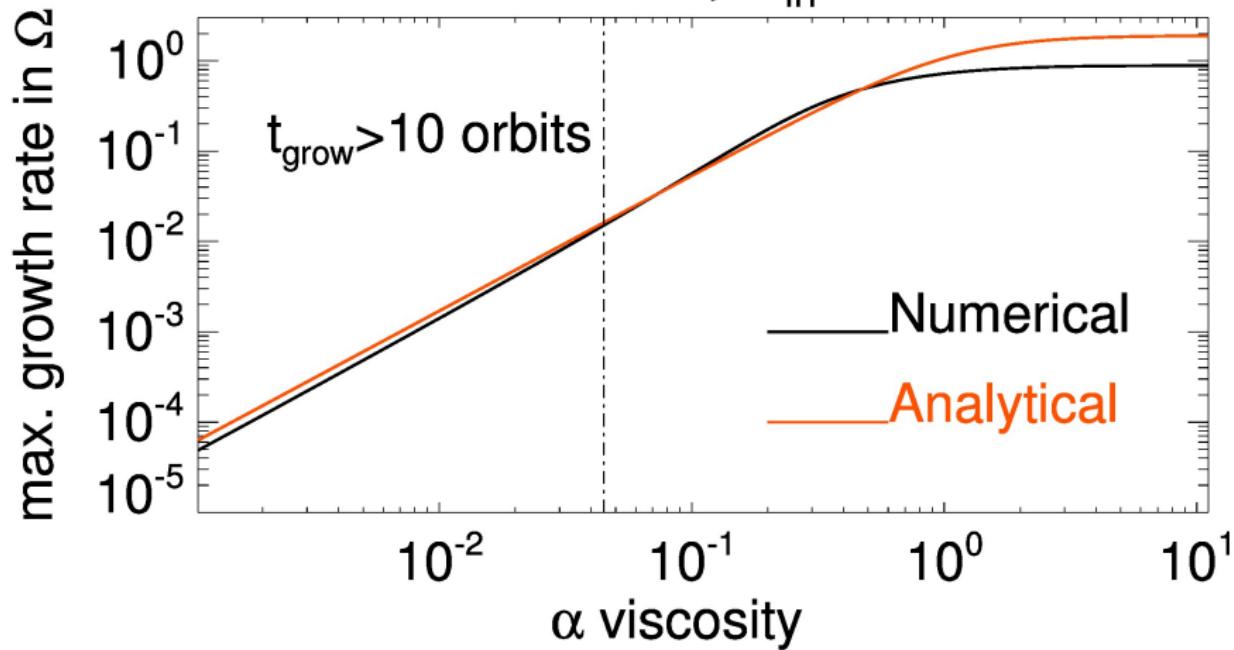
Maximum stress/viscosity in simulations



- Rice et al. (2005) report a α_{max} before fragmentation (also supported by Clarke et al. (2007))

Viscous gravitational instability

$$Q \propto \alpha^{-1/2}, T_{\text{irr}} = 0$$

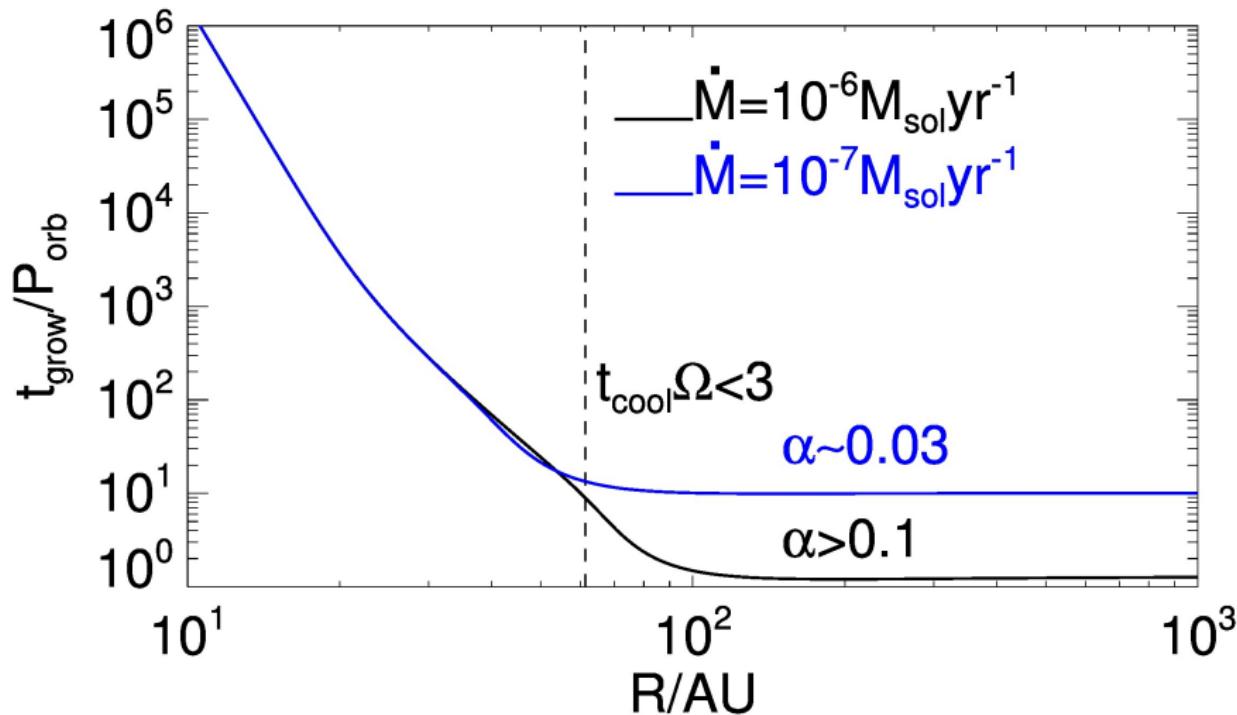


- Viscosity/friction removes rotational stabilization (cf. inwards migration of particles due to gas-dust drag)

(Lynden-Bell & Pringle, 1974; Willerding, 1992; Gammie, 1996)

Application to protoplanetary disks

- Input physical disk model into stability calculation — get growth timescales



- High \dot{M} disk fragments $\gtrsim 60 \text{ AU}$, growth times \sim one orbit

Summary

Vortices in massive 3D disks

- Self-gravity helps vortex survival, but with a turbulent core?

Vertical shear instability

- Feeds off free energy in $\partial_z \Omega \neq 0$
- Require large vertical motions to tap into weak vertical shear
- Enabled by ultra-fast cooling
- VSI possible in the outer disk between 10—100AU

Generalized gravitational instability

- Cooling: reduces thermal support
- Viscosity: reduces rotational support
- Fragmentation: inability to maintain steady gravito-turbulence because cooling and/or turbulent stresses (viscosity)

References

- Bai X.-N., 2016, ArXiv e-prints
- Casassus S., Wright C. M., Marino S., Maddison S. T., Wootten A., Roman P., Pérez S., Pinilla P., Wyatt M., Moral V., Ménard F., Christiaens V., Cieza L., van der Plas G., 2015, *ApJ*, 812, 126
- Chiang E., Youdin A. N., 2010, *Annual Review of Earth and Planetary Sciences*, 38, 493
- Clarke C. J., Harper-Clark E., Lodato G., 2007, *MNRAS*, 381, 1543
- Gammie C. F., 1996, *ApJ*, 457, 355
- Gammie C. F., 2001, *ApJ*, 553, 174
- Lesur G., Papaloizou J. C. B., 2009, *A&A*, 498, 1
- Li H., Colgate S. A., Wendroff B., Liska R., 2001, *ApJ*, 551, 874
- Lin M.-K., 2012, *MNRAS*, 426, 3211
- Lin M.-K., 2013a, *MNRAS*, 428, 190
- Lin M.-K., 2013b, *ApJ*, 765, 84
- Lin M.-K., 2014, *MNRAS*, 437, 575
- Lin M.-K., Youdin A. N., 2015, *ApJ*, 811, 17
- Lynden-Bell D., Pringle J. E., 1974, *MNRAS*, 168, 603
- Lyra W., Lin M.-K., 2013, *ApJ*, 775, 17
- Marois C., Zuckerman B., Konopacky Q. M., Macintosh B., Barman T., 2010, *Nature*, 468, 1080
- Nelson R. P., Gressel O., Umurhan O. M., 2013, *MNRAS*, 435, 2610
- Paardekooper S.-J., 2012, *MNRAS*, 421, 3286
- Pérez L. M., Isella A., Carpenter J. M., Chandler C. J., 2014, *ApJL*, 783, L13
- Rice W. K. M., Armitage P. J., Mamatsashvili G. R., Lodato G., Clarke C. J., 2011, *MNRAS*, 418, 1356
- Rice W. K. M., Lodato G., Armitage P. J., 2005, *MNRAS*, 364, L56
- Umurhan O. M., Nelson R. P., Gressel O., 2013, in *European Physical Journal Web of Conferences* Vol. 46 of *European Physical Journal Web of Conferences*, *Breathing Life Into Dead-Zones*. p. 3003
- van der Marel N., van Dishoeck E. F., Bruderer S., Birnstiel T., Pinilla P., Dullemond C. P., van Kempen T. A., Schmalzl M., Brown J. M., Herczeg G. J., Matthews G. S., Geers V., 2013, *Science*, 340, 1199
- Willerding E., 1992, *Earth Moon and Planets*, 56, 173