## The Answer of Assignment 1

## **Problem 1 Solution**

(1) In the position representation, from the eigenvalue equation  $\hat{H}\psi(x) = h_n\psi(x)$  and the Hamiltonian operator  $\hat{H} = -\frac{\hbar^2}{2m}\nabla^2$ , we have:

$$-\frac{\hbar^2}{2m}\nabla^2\psi(x) = h_n\psi(x)$$

The general solution is:

$$\psi(x) = c_1 e^{w_1 x} + c_2 e^{w_2 x}, \quad w_1 = i \sqrt{\frac{2mh_n}{\hbar^2}}, \quad w_2 = -i \sqrt{\frac{2mh_n}{\hbar^2}}$$

Given  $p_n = \frac{2\pi\hbar n}{L}$ ,  $h_n = \frac{2\pi^2\hbar^2n^2}{Lm}$ . Take  $\psi_n(x) = e^{\frac{ip_nx}{\hbar}}$  as an example, the eigenvalue corresponding to  $|\psi_n\rangle$  is  $\frac{2\pi^2\hbar^2n^2}{Lm}$ .

(2)

$$\langle \psi_{n_1} | \psi_{n_2} \rangle = \int_{-\infty}^{+\infty} \psi_{n_1}^*(x) \, \psi_{n_2}(x) \, dx$$

$$= \int_{-\infty}^{+\infty} e^{-\frac{ip_{n_1}x}{\hbar}} e^{\frac{ip_{n_2}x}{\hbar}} dx$$

$$= \int_{-\infty}^{+\infty} e^{\frac{i(p_{n_2} - p_{n_1})x}{\hbar}} dx$$

$$= \int_{-\infty}^{+\infty} e^{\frac{i2\pi(n_2 - n_1)x}{\hbar}} dx$$

$$= \lim_{l \to +\infty} \int_{-l}^{+l} e^{\frac{i2\pi(n_2 - n_1)x}{L}} dx$$

$$= \lim_{l \to +\infty} \frac{L \sin\left[\frac{2\pi(n_2 - n_1)l}{L}\right]}{\pi(n_2 - n_1)}$$

When  $n_2 = n_1$ ,  $\langle \psi_{n_1} | \psi_{n_2} \rangle \to \infty$  and

$$\lim_{l \to +\infty} \int_{-\infty}^{+\infty} \frac{\sin(lx)}{x} dx = \pi$$

So,

$$\langle \psi_{n_1} | \psi_{n_2} \rangle = L \delta_{n_1 n_2}$$

(3) When  $L \to \infty$ ,  $p_n$  becomes continuous and  $\psi_n(x)$  becomes a plane wave.

$$\langle \psi_{p_1} | \psi_{p_2} \rangle = \int_{-\infty}^{+\infty} e^{-\frac{ip_1 x}{\hbar}} e^{\frac{ip_2 x}{\hbar}} dx = \lim_{l \to +\infty} \frac{2\hbar \sin\left[\frac{(p_2 - p_1)l}{\hbar}\right]}{(p_2 - p_1)} = 2\pi \delta(\frac{p_2 - p_1}{\hbar}) = 2\pi \hbar \delta(p_2 - p_1)$$

(4) In the position representation,  $|\psi_{x_0,\epsilon}\rangle = \int_{-\infty}^{+\infty} \psi_{x_0,\epsilon}(x) |x\rangle dx$ 

$$\langle \delta | \delta \rangle = \int_{-\infty}^{+\infty} (x \psi_{x_0, \epsilon}(x) - x_0 \psi_{x_0, \epsilon}(x))^* (x \psi_{x_0, \epsilon}(x) - x_0 \psi_{x_0, \epsilon}(x)) dx$$

Considering the Gaussian wave packet  $\psi_{x_0,\epsilon}(x) = \left(\frac{1}{2\pi\epsilon^2}\right)^{1/4} e^{\frac{-(x-x_0)^2}{4\epsilon^2}}$ 

$$\int_{-\infty}^{+\infty} \psi^* \psi dx = \int_{-\infty}^{+\infty} \frac{1}{\sqrt{2\pi\epsilon^2}} e^{-\frac{(x-x_0)^2}{2\epsilon^2}} dx = \int_{-\infty}^{+\infty} \frac{1}{\sqrt{2\pi}} e^{\frac{-y^2}{2}} dy = 1$$

and

$$\int_{-\infty}^{+\infty} (x - x_0)^2 \frac{1}{\sqrt{2\pi\epsilon^2}} e^{-\frac{(x - x_0)^2}{2\epsilon^2}} dx = -\frac{\partial}{\partial a} \int_{-\infty}^{+\infty} \frac{1}{\sqrt{2\pi\epsilon^2}} e^{-a(x - x_0)^2} dx \bigg|_{a = \frac{1}{2\epsilon^2}}$$

Note that

$$\int_{-\infty}^{+\infty} e^{-a(x-x_0)^2} dx = \sqrt{\frac{\pi}{a}}$$

Therefore,

$$-\frac{\partial}{\partial a}\sqrt{\frac{\pi}{a}} = \frac{1}{2}\sqrt{\pi}a^{-3/2}$$

Substituting  $a = \frac{1}{2\epsilon^2}$ , we get

$$\int_{-\infty}^{+\infty} (x - x_0)^2 \frac{1}{\sqrt{2\pi\epsilon^2}} e^{-\frac{(x - x_0)^2}{2\epsilon^2}} dx = \epsilon^2$$

To ensure  $\langle \delta | \delta \rangle < \epsilon$ , we can set  $\delta = \frac{\sqrt{\epsilon}}{2}$ , so  $\psi_{x_0,\delta}(x) = \frac{1}{\sqrt[4]{2\pi\delta^2}} e^{-\frac{(x-x_0)^2}{4\delta^2}}$ . According to the previous derivation,

$$\langle \delta | \delta \rangle = \int_{-\infty}^{+\infty} (x \psi_{x_0, \delta}(x) - x_0 \psi_{x_0, \delta}(x))^* (x \psi_{x_0, \delta}(x) - x_0 \psi_{x_0, \delta}(x)) dx = \delta^2 = \frac{\epsilon}{4}$$

this satisfies the condition  $\langle \delta | \delta \rangle < \epsilon$ .

## Problem 2 Solution

(1) Consider a qubit with Hamiltonian

$$H = \mu BY$$

and and the initial state at time t = 0 is  $|\uparrow\rangle$ .