



Enhancing the Thermal Stability of Solution-Processed Small-Molecule Semiconductor Thin Films Using a Flexible **Linker Approach**

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Solution-processable small-molecule organic semiconductors have recently emerged as promising materials for application in low-cost thin-film transistors,[1] light-emitting diodes,[2] and photovoltaic cells.[3] In contrast to the more-common conjugated polymer semiconductors, molecular semiconductors offer the advantages of straightforward purification, structural precision, and batch-to-batch consistency. [4] However, the strong tendency of conjugated molecular organic semiconductors to self-assemble into crystalline domains results in drawbacks like dewetting, poor thin-film formation due to low solution viscosity, unpredictable crystalline domain sizes, and grain boundaries and that confound the morphological control and charge transport in devices fabricated from these materials. In particular, the need to control the thin-film morphology is especially important for application in bulk heterojunction (BHJ) photovoltaics where the domain size between the electron donating and electron accepting phases must be precisely controlled to afford high power conversion efficiency (PCE).^[5] For BHJ devices prepared with molecular semiconductor donors and the common small-molecule fullerene derivative, phenyl-C₆₁-butyric acid methyl ester (PC₆₁BM), as the acceptor the crystallization of the donor phase has been recently shown to be the driving force for blend separation.^[6] Thus, developing strategies to afford control over donor crystallization is an important goal. While volatile additives like diiodooctane^[7–9] and more recently non-volatile nucleation promoters[10,11] or insulating polymers^[12] have been employed to offer control over the thin-film formation and crystalline domain size, these approaches do not nominally act on the donor phase nor do they address the intrinsic stability of the thin-film morphology. Indeed, enhancing the long-term morphological stability of BHJ photovoltaics is one of the major remaining challenges in the field. [13,14] Specifically, the inherent thermodynamic immiscibility of the donor and acceptor phases drives phase separation and eventually decreases device performance, even at normal operating temperatures.^[15]

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DOI: 10.1002/adma.201501826



Given the drawbacks of molecular organic semiconductors, a "flexible linker" approach, [16] wherein well-defined conjugated segments are covalently linked with flexible aliphatic chains into polymeric materials, could reasonably provide a route to offer control over molecular donor crystallization and improve morphological stability in thin-film devices, while also maintaining the consistency of the conjugated moiety. Herein, we demonstrate this flexible-linker strategy for molecular organic semiconductors for the first time using a common small-molecule donor, 3,6-bis(5-(benzofuran-2-yl)thiophen-2-yl)-2,5-bis(2ethylhexyl)pyrrolo[3,4-c]pyrrole-1,4(2H,5H)-dione, coded as DPP(TBFu)2. [17] When the flexibly linked material is employed as an additive in thin films, a unique level of control is observed over the morphology and the thermal stability of pure donor films in transistor devices and in blends with PC61BM for BHJ photovoltaics.

The synthesis of the flexibly linked DPP(TBFu)2 derivative [coded FL-DPP(TBFu)₂] is shown schematically in Figure 1a. The Stille-coupling polycondensation between the brominated diketopyrrolopyrrole core and the stannylated benzofuran dimer followed by standard purification and finally preparatory size exclusion chromatography afforded the pure polymer with a number average molecular weight, M_n , of 10 kDa and a PDI of 1.4 ($n \approx 11-15$) in the hundred-milligram scale. Full synthetic procedures and detailed characterization are given in the Experimental Section and the Supporting Information. To gain a first insight into the crystallization behavior of the FL-DPP(TBFu)₂ polymer, we performed differential scanning calorimetry (DSC). The neat material did not exhibit a detectable glass transition nor enthalpic transitions associated with melting/crystallization in the temperature range from 50 to 240 °C. However, when combined with the parent DPP(TBFu)₂ small molecule, a significant change in the crystallization occurred. Cooling curves obtained after melting blends of the FL additive (at a weight fraction, $f_{\rm FL}$) in the small molecule are shown in Figure 1b. A strong dependency on the DPP(TBFu)2 crystallization temperature, Tc, upon addition of the FL additive is clearly observed. With $f_{\rm FL}$ at only 1 wt% a pronounced shift of the exothermic crystallization to higher temperatures is noted compared to the neat small molecule ($f_{\rm FL}=0$). This observation is consistent with the behavior of a nucleation promoter,[10,11,14] suggesting the FL additive acts to reduce the driving force needed to nucleate stable nuclei of DPP(TBFu)2. The maximum shift of the crystallization onset was found to be $\Delta T_{\rm c} = 10$ °C with $f_{\rm FL} = 5$ wt%, which is consistent in magnitude with recent reports of small molecule, non-conjugated www.advmat.de

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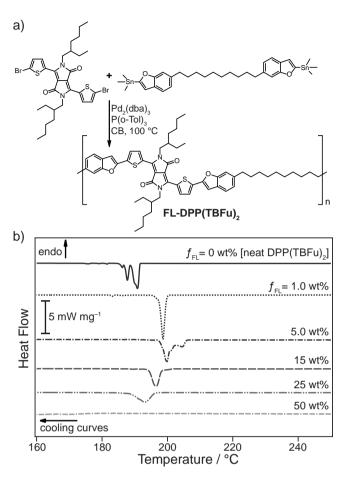


Figure 1. a) The synthetic scheme for FL-DPP(TBFu)₂. b) The differential scanning calorimetry cooling curves for different amounts of FL-DPP(TBFu)₂ in DPP(TBFu)₂.

nucleation promoters.^[14] For higher $f_{\rm FL}$ the onset of crystallization returns to lower temperatures and the enthalpy of the crystallization decreases suggesting the complete disruption of the ability of the blend to crystallize. We note that for $f_{\rm FL}=1$ and 5 wt%, comparable enthalpies of crystallization are observed ($\Delta H_{\rm c}=34.7$ and 30.1 J g⁻¹, respectively) with respect to the neat material (39.9 J g⁻¹) suggesting that the overall degree of crystallinity is not strongly reduced despite the addition of the amorphous FL-DPP(TBFu)₂. In addition, we note that despite the drastic change in the crystallization behavior, only a small change was observed on the melting behavior (see Figure S1 in the Supporting Information) consistent with other reports of nucleation promoters.^[18]

The effect of our FL strategy on the charge transport properties of solution-cast thin films was next investigated in bottom-contact/bottom-gate field-effect transistors (FETs). Devices prepared with the neat FL material showed no field-effect mobility ($\mu_{\rm FET}$) confirming that the absence of crystalline ordering (as indicated by the DSC results) prevents long-range charge carrier transport. However, as-cast films of the FL material blended with DPP(TBFu)₂ gave measurable $\mu_{\rm FET}$ for a wide range of $f_{\rm FL}$ and for $f_{\rm FL}$ up to 5 wt% the values were comparable with the neat DPP(TBFu)₂ film with $\mu_{\rm FET} \approx 1$ –2 × 10⁻⁵ cm² V⁻¹ s⁻¹ (see Figure S2 in the Supporting Information). We note that the

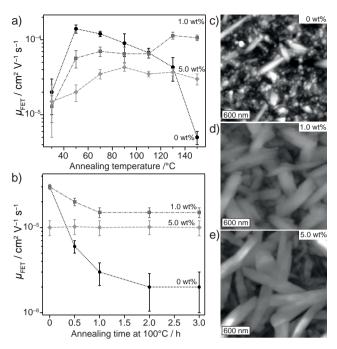


Figure 2. Thin-film transistor performance and morphology. a) The average extracted field-effect mobility as a function of thin-film annealing temperature. b) The mobility as a function of annealing time at 100 °C. c–e) Atomic force images show the topology of the thin-film transistor active layer after 3 h at 100 °C.

values of the as-cast μ_{FET} are similar to previous reports of the as-synthesized small molecule (which contains a mixture of stereoisomers).^[19] As annealing is typically performed to increase the crystalline ordering and the charge carrier transport,[20] we next investigated the effect of the annealing temperature on transistors prepared with neat DPP(TBFu)₂ ($f_{\rm FL} = 0$) and for $f_{\rm FL}$ = 1 and 5 wt%. The results are shown in Figure 2a. Annealing at 50 °C for 30 min improved μ_{FET} in all cases, however, subsequent annealing at higher temperatures causes a decrease in μ_{FET} for neat DPP(TBFu)₂ devices. In contrast, the observed temperature dependence is quite different in the presence of the FL additive. The measured μ_{FET} increases slightly as a function of temperature until 150 °C for the 1 wt% devices while the 5 wt% devices remain fairly constant from 70 to 150 °C, suggesting a more robust charge carrier transport network is created when adding the FL-DPP(TBFu)2.

To better establish the thermal stability of the active layer, a set of devices, prepared by annealing at 100 °C to have similar initial performance, were further subjected to a long-term thermal stress, periodically quenching the films to 30 °C for testing. The measured $\mu_{\rm FET}$ is reported against the annealing time at 100 °C in Figure 2b. While $\mu_{\rm FET}$ drops an order of magnitude for $f_{\rm FL}=0$, a considerably smaller decrease is observed when $f_{\rm FL}=1$ wt%, and notably at 5 wt% the performance remains constant. Visualizing the active layer morphology with atomic force microscopy (AFM) gives further insight into the origin of the improved thermal performance of the transistors with the added FL-DPP(TBFu)₂. As-cast films with 0–5 wt% of the FL additive all exhibit a similar fibril-type morphology (Figure S3, Supporting Information). After the extended

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thermal stress test at 100 °C for a total of 3 h, a drastic difference is observed between the topology of the films as shown in Figure 2c-e. The neat (0 wt%) device exhibits small circular domains (≈100-200 nm) and a few longer shards. The disappearance here of the initial fibril-type morphology may be due to a dewetting from the substrate or reasonably caused by the formation of defects during the rapid quenching process at the semiconductor dielectric interface^[21] which propagate and alter the domain morphology. Indeed, the thermal annealing of solution-processed molecular semiconductor films has recently been shown to induce the formation of trapping states, likely at grain boundaries, that reduce charge carrier transport. [22] In stark contrast to the neat DPP(TBFu)₂ devices, when $f_{\rm FL} = 1$ or 5 wt% films exhibit large banded features 200-500 nm in width and more than micrometers in length after the extended annealing test. In addition, the roughness of the film decreases considerably with $f_{FL} = 5$ wt% (see height profiles in Figure S4

in the Supporting Information). The transistor device thermal stability data together with the evolution of the morphology demonstrate that the FL-DPP(TBFu)₂ actively participates in the stabilization of the thinfilm charge transport network. In contrast to the typical behavior of a nucleation promoter, which increases the number of nucleation events thus leading to a smaller crystalline domain size, [11] the FL additive also appears at this point to promote larger crystalline domains. Given the polymeric structure of our FL additive one possible explanation for this behavior would be that it is acting as a tie molecule to bridge adjacent crystal domains, effectively locking-in the active layer morphology. Given the relatively short length of the FL polymer chain (≈40 nm) compared to the observed domain size, and the amount of FL additive employed, this explanation would be possible if the FL additive is positioned at the crystal grain boundaries and the polymer chain bridged two adjacent domains. This is reasonable as the FL additive is likely positioned at the grain boundary due to its action as a nucleation promoter. However, we do not suspect that the FL additive can act as a traditional tie molecule, which typically would link many crystallite domains. Another aspect that may affect the formation of a more stable charge transport network is the nucleation of a different crystal form of the DPP(TBFu)₂ via the FL additive. Indeed, polymorph formation has been observed with stereoisomers of DPP(TBFu)₂ and are known to have an effect on the charge transport. [18] Analysis of thick films (deposited on SiO₂ and annealed at a high temperature, 190 °C, below the melting point to maximize crystallinity) by grazing incidence wide angle X-ray scattering (GIWAXS) (see Figure S5 in the Supporting Information) supports the notion that a change in crystal formation could play a role. While at $f_{FL} = 0$ the diffraction peaks match previously reported films annealed at low temperature (80 °C),[10] distinctive Bragg rods are observed when $f_{\rm FL}$ = 1 wt% suggesting a thin-film polymorphism.^[23] Analysis of the films annealed at 190 °C by AFM supports this by indicating a clear change from fiber-like features (500 nm in width) observed in the control film (Figure S5a, Supporting Information) to wide platelets observed with $f_{\rm FL}$ = 1 wt% (Figure S5c, Supporting Information). While the complete understanding of the complex role of the FL-DPP(TBFu)₂ on the thin-film crystallinity is undoubtedly beyond the scope of this initial demonstration of the flexible

linker concept, the strong effect of the FL additive on the morphology and performance of donor-only films is clear.

Given the common use of DPP(TBFu)₂ as an electrondonor small molecule blended in BHJ photovoltaic cells with fullerene-based electron acceptors, we next probed the effect of the inclusion of the FL additive in DPP(TBFu)2:PC61BM solar cells. Standard indium tin oxide (ITO)/poly(3,4ethylenedioxythiophene):polystyrene sulfonate (PEDOT:PSS)/ BHJ/Al devices were fabricated by typical procedures as described in the Experimental Section. We compared performance of devices with $f_{\rm FL} = 0-5$ wt% (based on the total mass of the active layer in DPP(TBFu)₂:PC₆₁BM blends at a ratio of 3:2) and found that, for devices tested without annealing, the photovoltaic figures of merit, J_{sc} , V_{oc} , and FF (the short-circuit current density, open-circuit voltage, and fill factor, respectively) were not affected by the presence of the FL additive (see Figure S6 in the Supporting Information) and were consistent with previous reports with "as-cast" active layers.[17] The invariance of $V_{\rm oc}$ supports the notion that the presence of the additive does not change the HOMO level of the acceptor phase, which is reasonable considering the nature of the conjugated segment is unchanged. However, the unchanging I_{sc} and FF values of the as-cast devices with $f_{\rm FL} = 0-5$ wt% suggest that the FL additive does not have a significant effect on the BHJ morphology during the initial film formation (spin-casting from chlorobenzene) despite its established activity as a nucleation promoter for the donor phase.

Annealing devices with $f_{\rm FL}$ = 0 at 100 °C for 10 min gave an average device PCE of 3.4% \pm 0.3%. We note that this value is slightly lower than state-of-the-art values obtained when using $PC_{71}BM$ due to the expected lower J_{sc} . Devices with $f_{\rm FL}$ = 0.5 and 1.5 wt% exhibited equivalent performance when the annealing time was optimized. Interestingly, we found it necessary to extend the annealing to 1 h to afford PCE equal to the $f_{FL} = 0$ case. Optimized "best-case" J-V curves for $f_{FL} = 0-5$ wt% are shown in Figure 3a and the average PCE as a function of the annealing time at 100 °C is shown for $\ensuremath{f_{FL}}=0$ and 0.5 wt% in Figure 3b. When $f_{FL} = 3$ or 5 wt% the best-case J-V performance is significantly decreased, exhibiting lower FF and I_{sc} , which indicates an increased recombination and poorer charge carrier transport. Remarkably, the thermal stability (Figure 3b) of the device performance was found to be significantly different between the $f_{\rm FL}$ = 0 and 0.5 wt%. After only 3 h of thermal stress at 100 °C a decrease in the PCE from 3.4% to 2.3% is observed when $f_{\rm FL} = 0$. This observed trend, common for many systems employing PC₆₁BM, is primarily attributed to a decrease in FF and I_{sc} due to the phase segregation of the PC₆₁BM from the donor phase. In contrast, device performance remains stable at 3.2% when $f_{\rm FL} = 0.5$ wt% even after 22 h of annealing at 100 °C. This suggests a different evolution of the microstructure of the BHJ during the thermal stress test.

Insight into the origin of the improved thermal stability of the donor:acceptor blends with the FL additive was next sought by examining the morphology of the active layer with AFM and nanomechanical mapping^[24,25] to identify donor-rich and acceptor-rich domains. Indeed, as the Young's modulus (YM) of neat PC₆₁BM (\approx 17 GPa) and neat DPP(TBFu)₂ (\approx 3 GPa) are significantly different (see Figure S7 in the Supporting Information), YM mapping can give information on the domain

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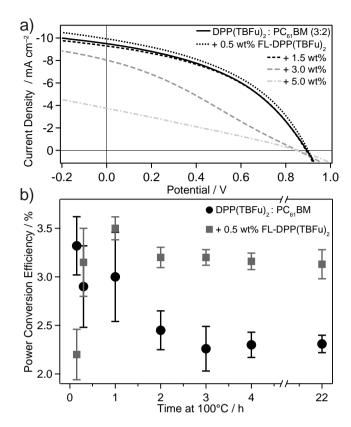


Figure 3. Bulk heterojunction device performance. a) The best case *J–V* curves for devices prepared with a DPP(TBFu)₂:PC₆₁BM blend at a ratio of 3:2 and added FL-DPP(TBFu)₂. b) The average power conversion efficiency as a function of the annealing time at 100 °C.

composition. Combined topographical and YM map results are shown in Figure 4. The active layer topography after annealing at 100 °C for 10 min (unconfined by the Al top electrode) is shown in Figure 4a (left panel), and exhibits rod-shaped domains 100–200 nm in size when $f_{\rm FL} = 0$ (similar to previous reports.[17] We note that the topography of as-cast BHJ films when $f_{\rm FL}$ = 0–5 wt% are similar, except with slightly larger rodshaped domains for $f_{\rm FL}$ = 0.5 and 1.5 wt% (see Figure S8 in the Supporting Information). The YM map corresponding to the topography image of Figure 4a is displayed on the right of the figure. Here, domains of low YM, corresponding roughly to the rod-shaped domains, are distinguishable from areas of higher modulus, confirming that these rod domains are donorrich and further suggesting that the donor and acceptor are mixed at a length scale of ≈100 nm. For films with no FL additive ($f_{\rm FL} = 0$) that have been annealed for 22 h at 100 °C, the topology (Figure 4b) shows the absence of rod-shaped domains and the presence of circular domains similar to the donor-only film shown in Figure 2c. The corresponding YM map gives values <6 GPa suggesting that the surface of the film is composed primarily of DPP(TBFu)₂ (i.e., no PC₆₁BM-rich domains are present). Indeed, micrometer-sized PC₆₁BM-rich domains are observed by AFM, YM mapping, and optical microscopy (see Figure S9 in the Supporting Information) suggesting that the PC₆₁BM has segregated from the DPP(TBFu)₂ on a length scale too large for BHJ operation. We note that the performance of the corresponding photovoltaic device ($f_{FL} = 0$) after annealing for 22 h likely stabilizes at a nonzero value due to the effect of the morphological confinement brought by the Al top contact. [26] With the FL additive present in the BHJ, a considerably different morphology is observed after thermal

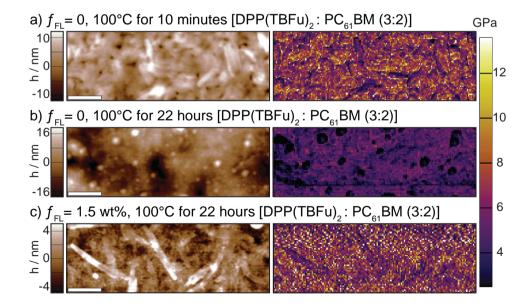


Figure 4. a-c) Bulk heterojunction topography (left panels) and Young's modulus mapping (right panels) of the corresponding area. The scale bars are 500 nm.

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stress for 22 h (Figure 4c). In this case, the topography shows a comparatively smoother film with long rod-shaped domains similar to those observed in the neat films (Figure 2d,e). The YM map indicates that these rod-shaped regions exhibit a lower modulus consistent with the DPP(TBFu)₂. Moreover, we note that Young's moduli higher than 10 GPa are also measured in high proximity to the rod-shaped domains, indicating that PC₆₁BM-rich domains remain present in the film with the FL additive present after thermal treatment for 22 h. This considerable difference in behavior of the morphology is indeed consistent with the observed difference in device performance and suggests that the FL additive strategy is effective at extending the morphological stability of BHJ molecular solar cells.

In addition, we note that the presence of the comparatively longer rod-shaped crystalline domains observed in the morphological study suggests an apparent increase in donor domain size when the FL additive is present in blends with PC₆₁BM and optimized thermal annealing is performed. However, we note that out-of-plane XRD measurements indicate that the donor has similar crystallinity in the blend film without the FL additive after 10 min annealing compared to the $f_{\rm FL} = 0.5$ wt% blend film when annealing for 1 h (see Figure S10 in the Supporting Information); therefore, no increased overall crystallinity of the donor phase is observed using the FL additive consistent with the DSC results which suggest a slight decrease in donor crystallinity. However, supposing that the longer rod-shaped donor domains observed in stabilized BHJ blends with the FL additive could afford improved charge carrier transport over longer distance, we hypothesized that a greater difference in device performance (with and without the FL additive) would be observed if thicker active layers were employed. To provide preliminary evidence to support this notion, we increased the thickness of the active layer by 50% (i.e., from 90 to 140 nm). After optimized annealing at 100 °C, devices with $f_{\rm FL} = 0$ gave a 40% lower PCE compared to devices with $f_{\rm FL}$ = 0.5 wt%. As expected, the difference was due to a drop in FF and $J_{\rm sc}$ which are found to be, respectively, 0.34 and 7.5 mA cm⁻² (for $f_{\rm FL}$ = 0), compared to 0.45 and 9.9 mA cm⁻² (for $f_{\rm FL}$ = 0.5 wt%, see Figure S11 in the Supporting Information). Even though significant PCE improvement was not observed over the thinner cells, this result suggests that in thicker BHJ films—where increased crystallinity is required to afford efficient charge extraction—the FL additive can support a more intimate mixing of the donor and acceptor, while still providing suitable crystallinity for charge transport. Indeed, we note that for BHJ films annealed for 1 h with or without the FL additive, films without the additive exhibited increased donor crystallinity (consistent with the donordriven crystallization in this system, see Figure S10 in the Supporting Information) which drives the phase segregation and the lower J_{SC} . The improved BHJ thermal stability and lower comparative donor crystallinity can be rationalized again by the ability of the FL additive to promote the nucleation of the donor phase. The presence of the amorphous FL additive at the grain boundaries of the donor crystallites could reasonably maintain a stable donor transport network while also providing amorphous regions for PCBM, thereby limiting its exclusion. This combination of beneficial characteristics is unique among additives to bulk heterojunctions. We cannot, however, exclude the

possibility of donor polymorph formation playing a role in the BHJ stability. Indeed, while out-of-plane XRD measurements do not indicate any considerable difference in the observed diffraction peak, a difference is not expected even with polymorphs as reported DPP(TBFu)₂ polymorphs exhibit similar scattering vectors for the main peak. [19] Further analysis of the BHJ crystallinity is a point of interest given these initial results.

Importantly, we note that the FL additive strategy presented here is significantly different than recently reported approaches to reactively crosslink the donor and acceptor after BHJ formation. While these crosslinking approaches have been shown to improve the thermal stability of BHJs when polymer donor phases are used, as a drawback they rely on the uncontrolled reaction of functional groups that likely add electronic defects to the film. In addition, we emphasize that the observed behavior of the FL additive is also distinct from simply blending in insulating (inert) polymer additives, in molecular bulk heterojunctions. While this approach offers some control over the thin-film formation, since the inert polymer segregates from the BHJ, it likely does not improve thermal stability.

Overall the device results and morphological investigations presented here demonstrate that the strategy of linking molecular semiconductors with flexible aliphatic chains into polymers that possess defined conjugated segments, but extended covalent connectivity is a viable way to control the crystalline behavior of the solution-processed molecular semiconductor thin films. We have demonstrated that our new class of polymeric additive can promote donor nucleation in neat films and BHJ blends to improve thermal stability and long-range charge transport. We have also shown, for the first time, that smallmolecule BHJ solar-cell thermal stability can be drastically improved by including our flexibly linked additive. Since the flexible linker approach is straightforward and can be applied to a wide variety of different conjugated small molecules, we believe that this strategy can have broad application in controlling and stabilizing the active layer morphology in molecular BHJ PVs, OLEDs, and TFTs.

Experimental Section

Material Synthesis and Basic Characterization: The DPP(TBFu)₂ and PC₆₁BM small molecules were prepared using standard procedures.^[17] The FL-DPP(TBFu)₂ polymer was prepared using a procedure adapted from previous work.^[16] See the Supporting Information for full synthetic details as well as basic material characterization.

Transistor Fabrication and Testing: Bottom-contact FETs were fabricated on prepatterned test substrates (Fraunhofer Institute for Photonic Microsystems) whose source and drain contacts were composed of a 30 nm thick gold layer on top of a 10 nm thick titanium adhesion layer. A 230 nm thick silicon oxide was used as gate dielectric and n-doped silicon wafer as the substrates and gate electrode. The channel length (L) was 10 or 20 µm and channel width (W) was 10 mm. The transistor substrates were cleaned by sonication in acetone and isopropyl alcohol at RT for 15 min in each solvent. After drying under nitrogen, the substrates were subsequently exposed to a nitrogen plasma for 30 min. Films of ≈60 nm thickness were spin-coated from 10 mg mL⁻¹ DPP(TBFu)₂ + FL-DP(TBFU)₂ solutions in CHCl₃ at 500 rpm (for 1 min). The solutions were prepared by dissolution at 55 °C for 1 h under continuous stirring. All solutions and films were prepared in an argon atmosphere. Unannealed films were left overnight under vacuum

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(0.01 mbar) to remove solvent before testing. Electronic testing of the FETs was carried out in a nitrogen atmosphere using a custom-built probe station and a Keithley 2612A dual-channel source measure unit. Samples were annealed in a nitrogen atmosphere and quenched to room temperature before electronic testing. The field-effect mobility was extracted from the saturation region as previously described. [16] Six devices were measured at each condition and the average values are reported.

Thermal Characterization: Melting temperatures, T_m , and enthalpies of fusion, $\Delta H_{\rm f}$, were measured with a Perkin-Elmer DSC8000 differential scanning calorimeter calibrated with indium and zinc, using a scanning rate of 10 °C min⁻¹. Samples were prepared by drop-casting and slow evaporation of the solvent at 40 °C under argon atmosphere before the measurement. $\Delta H_{\rm f}$ was calculated by integration of the melting endotherm.

Solar-Cell Fabrication and Testing: Solar cells were fabricated with a 35 nm layer of PEDOT:PSS (Ossilla M 121 Al 4083) deposited and annealed (100 °C) on a glass substrate patterned with 300 nm of ITO. The BHJ active layer was spin-cast at 3000 rpm from a solution of DPP(TBFu)₂ and PC₆₁BM in chloroform at a total solid concentration of 20-30 mg mL^{-1} . The active layers were determined to be ≈ 90 nm thick using a Bruker Dektak XT profilometer. An 80 nm thick aluminum cathode was deposited (area 16 mm²) by thermal evaporation (Kurt J. Lesker Mini-SPECTROS). Electronic characterization was performed under simulated AM1.5G irradiation from a 300 W Xe arc lamp set to 100 mW cm² with a calibrated Si photodiode (ThorLabs). Currentvoltage curves were obtained with a Keithley 2400 source measure unit. Device fabrication was performed under an argon atmosphere and testing was performed under in a nitrogen filled glovebox.

Force Peak Mapping: Peak force measurements were recorded in ambient conditions using a Cypher S AFM from Asylum Research. A silicon tip (AC160TS) with a theoretical spring constant of 42 N m⁻¹ was used. The experimental spring constant was obtained using the GetReal automated probe calibration feature of Asylum Research software. Over all the scans, a peak force of 100 nN was kept constant for all the samples. The reduced Young's modulus was obtained by fitting the retraction curve with the Derjaguin, Muller, Toropov (DMT) model, which accounts for the adhesion between the tip and the sample surface.

Diffraction Characterization: GIWAXS measurements were carried out at ANKA synchrotron in KIT on beamline PDIFF. The incident energy was 11 keV. The films were illuminated at a constant incidence angle of 0.15° using a 2D area detector. The distance between the sample and the 2D camera was set at 180 mm. Out-of-plane, specular X-ray scattering was performed with a standard Cu K_{α} laboratory source (0.15418 nm) on a Bruker D8 Advanced goniometer using grazing incidence conditions with an incident angle of 0.3°, a scan speed of 0.05 deg min⁻¹, and a Lynx Eye linear detector.

Supporting Information

Supporting Information is available from the Wiley Online Library or from the author.

Acknowledgements

The authors thank the Swiss National Science Foundation (Project No. 200021_144311) and the European Research Commission (Project No. 336506, CEMOS) for financial support. Dr. Stephen Doyle, at the ANKA Beamline is gratefully acknowledged for assistance with the 2D GIWAXS measurements.

> Received: April 16, 2015 Revised: May 26, 2015 Published online: August 25, 2015

- [1] S. Allard, M. Forster, B. Souharce, H. Thiem, U. Scherf, Angew. Chem., Int. Ed. 2008, 47, 4070.
- [2] K. S. Yook, J. Y. Lee, Adv. Mater. 2014, 26, 4218.
- [3] Y. Lin, Y. Li, X. Zhan, Chem. Soc. Rev. 2012, 41, 4245.
- [4] B. Walker, A. Tamayo, D. T. Duong, X. D. Dang, C. Kim, J. Granstrom, T. Q. Nguyen, Adv. Energy Mater. 2011, 1, 221.
- [5] M. T. Dang, L. Hirsch, G. Wantz, J. D. Wuest, Chem. Rev. 2013, 113, 3734
- [6] A. Sharenko, M. Kuik, M. F. Toney, T. Q. Nguyen, Adv. Funct. Mater. 2014, 24, 3543.
- [7] A. Viterisi, F. Gispert-Guirado, J. W. Ryan, E. Palomares, J. Mater. Chem. 2012, 22, 15175.
- [8] Y. M. Sun, G. C. Welch, W. L. Leong, C. J. Takacs, G. C. Bazan, A. J. Heeger, Nat. Mater. 2012, 11, 44.
- [9] L. A. Perez, J. T. Rogers, M. A. Brady, Y. M. Sun, G. C. Welch, K. Schmidt, M. F. Toney, H. Jinnai, A. J. Heeger, M. L. Chabinyc, G. C. Bazan, E. J. Kramer, Chem. Mater. 2014, 26, 6531.
- [10] A. Sharenko, N. D. Treat, J. A. Love, M. F. Toney, N. Stingelin, T. Q. Nguyen, J. Mater. Chem. A 2014, 2, 15717.
- [11] N. D. Treat, J. A. N. Malik, O. Reid, L. Yu, C. G. Shuttle, G. Rumbles, C. J. Hawker, M. L. Chabinyc, P. Smith, N. Stingelin, Nat. Mater. **2013**. 12. 628.
- [12] Y. Huang, W. Wen, S. Mukherjee, H. Ade, E. J. Kramer, G. C. Bazan, Adv. Mater. 2014, 26, 4168.
- [13] B. C. Schroeder, Z. Li, M. A. Brady, G. C. Faria, R. S. Ashraf, C. J. Takacs, J. S. Cowart, D. T. Duong, K. H. Chiu, C. H. Tan, J. T. Cabral, A. Salleo, M. L. Chabinyc, J. R. Durrant, I. McCulloch, Angew. Chem., Int. Ed. 2014, 53, 12870.
- [14] C. Lindqvist, J. Bergqvist, C. C. Feng, S. Gustafsson, O. Backe, N. D. Treat, C. Bounioux, P. Henriksson, R. Kroon, E. G. Wang, A. Sanz-Velasco, P. M. Kristiansen, N. Stingelin, E. Olsson, O. Inganas, M. R. Andersson, C. Muller, Adv. Energy Mater. 2014, 4, 1301437.
- [15] I. Cardinaletti, J. Kesters, S. Bertho, B. Conings, F. Piersimoni, J. D'Haen, L. Lutsen, M. Nesladek, B. Van Mele, G. Van Assche, K. Vandewal, A. Salleo, D. Vanderzande, W. Maes, J. V. Manca, J. Photonics Energy **2014**, 4, 040997.
- [16] A. Gasperini, S. Bivaud, K. Sivula, Chem. Sci. 2014, 5, 4922.
- [17] B. Walker, A. B. Tomayo, X. D. Dang, P. Zalar, J. H. Seo, A. Garcia, M. Tantiwiwat, T. Q. Nguyen, Adv. Funct. Mater. 2009, 19,
- [18] W. J. Liu, H. L. Yang, Z. Wang, L. S. Dong, J. J. Liu, J. Appl. Polym. Sci. 2002, 86, 2145.
- [19] J. Liu, Y. Zhang, H. Phan, A. Sharenko, P. Moonsin, B. Walker, V. Promarak, T.-Q. Nguyen, Adv. Mater. 2013, 25, 3645.
- [20] Z. Li, X. Zhang, Y. Zhang, C. F. Woellner, M. Kuik, J. H. Liu, T. Q. Nguyen, G. Lu, J. Phys. Chem. C 2013, 117, 6730.
- [21] E. Gann, X. K. Gao, C. A. Di, C. R. McNeill, Adv. Funct. Mater. 2014, 24, 7211.
- [22] J. D. A. Lin, O. V. Mikhnenko, T. S. van der Poll, G. C. Bazan, T.-Q. Nguyen, Adv. Mater. 2015.
- [23] S. C. B. Mannsfeld, M. L. Tang, Z. Bao, Adv. Mater. 2011, 23, 127.
- [24] D. Wang, T. P. Russell, T. Nishi, K. Nakajima, ACS Macro Lett. 2013, 2, 757.
- [25] D. Wang, F. Liu, N. Yagihashi, M. Nakaya, S. Ferdous, X. Liang, A. Muramatsu, K. Nakajima, T. P. Russell, Nano Lett. 2014, 14,
- [26] K. Sivula, Z. T. Ball, N. Watanabe, J. M. J. Frechet, Adv. Mater. 2006, 18, 206.
- [27] L. Derue, O. Dautel, A. Tournebize, M. Drees, H. Pan, S. Berthumeyrie, B. Pavageau, E. Cloutet, S. Chambon, L. Hirsch, A. Rivaton, P. Hudhomme, A. Facchetti, G. Wantz, Adv. Mater. **2014**, 26, 5831.
- [28] C.-P. Chen, C.-Y. Huang, S.-C. Chuang, Adv. Funct. Mater. 2015, 25,