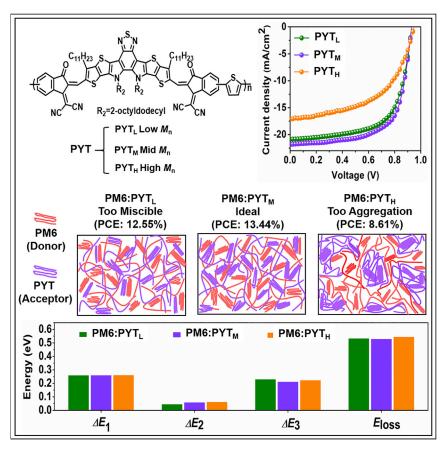
**Joule** 



### **Article**

# Controlling Molecular Mass of Low-Band-Gap Polymer Acceptors for High-Performance All-Polymer Solar Cells



A narrow-band-gap polymer acceptor ( $P_A$ ), namely PYT, is reported, while a series of PYT polymer acceptors with controlled average molecular weight ( $M_n$ ) values were synthesized for fine-tuning the molecular crystallinity and miscibility. When fabricated into all-PSCs with a polymer donor ( $P_D$ ) PM6, we observed a clear  $M_n$  dependence on device performance. The PM6:PYT<sub>M</sub> device gives a record-high PCE of 13.44% for the all-PSCs, resulting from the good  $P_D$ - $P_A$  pair miscibility, suitable blend microstructure, and improved charge transport properties.

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#### **HIGHLIGHTS**

A fused-aromatic-ringconstructed polymer acceptor is reported

Donor-acceptor miscibility can be modified by optimizing molecular mass

A PYT<sub>M</sub>-based all-PSC exhibits an excellent PCE of 13.44%

The molecular weight of polymers plays a vital factor for highperformance all-PSCs

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#### **Article**

## Controlling Molecular Mass of Low-Band-Gap Polymer Acceptors for High-Performance All-Polymer Solar Cells

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#### **SUMMARY**

Recent advances in the development of polymer acceptors and the investigation of molecular mass have boosted the power conversion efficiency (PCE) of all-polymer solar cells (all-PSCs) to approximately 11%. Here, a fused-aromatic-ring-constructed polymer acceptor PYT (Poly[(2,2'-((2Z,2'Z)-((12,13-bis(2-octyldodecyl)-3,9-diundecyl-12,13-dihydro[1,2,5]thiadiazolo[3,4e]thieno[2",3":4',5']thieno[2',3':4,5] pyrrolo[3,2-g]thieno[2',3':4,5]thieno[3,2-b]-indole-2,10-diyl)bis(methanylylidene))bis(3-oxo-2,3-dihydro-1H-indene-2,1-diylidene)) dimalononitrile-alt-2,5-thiophene)]) is reported, while a series of PYT polymers with different molecular masses (designated as PYT<sub>L</sub>, PYT<sub>M</sub>, and PYT<sub>H</sub>) are prepared to fine-tune the molecular crystallinity and miscibility. Benefiting from the advantages of PYT series, which possess broad absorption with a narrow band of 1.40-1.44 eV and high absorption coefficients of over  $1.00 \times 10^5$  cm<sup>-1</sup>, we investigated the blend miscibility and device performance of all-PSCs based on a wide-bandgap polymer donor, PM6. The PYT<sub>M</sub>-based all-PSCs exhibit an excellent PCE of 13.44%, outperforming those with PYT<sub>L</sub> (12.55%) and PYT<sub>H</sub> (8.61%). Our results provide insight into polymer acceptor backbone and molecular mass and suggest guidelines to rationally select polymers for all-PSCs.

#### **INTRODUCTION**

Solution-processed polymer solar cells (PSCs), recognized as a promising technology of converting solar energy to electricity, have attracted much attention due to their advantages of flexibility, low-cost, light-weight, transparent, and easy of large-scale fabrication. 1-6 Intense research efforts over the past years have provided a deeper understanding of the efficient photovoltaic materials, 7-11 interface materials, 12-14 and device engineering, 5,15 now enabling fullerene-free PSCs with power conversion efficiencies (PCEs) beyond 16%. Among the various systems in non-fullerene PSCs, 16,17 all-polymer solar cells (all-PSCs), 18 which are composed of binary blends of conjugated polymer donor (PD) and polymer acceptor (PA), attract particular interest, because of enhanced polymer electronic structure tunability, excellent thermal stability, and mechanical bulk heterojunction (BHJ) film robustness. 19-21 In addition, all-polymer system maintains a good compatibility with large-area manufacture during operation and stress testing protocols.<sup>22-24</sup> Recently, the PCEs of all-PSCs have been dramatically enhanced to exceeding 11%. <sup>25,26</sup> This progress is mainly resulting from optimal polymer design, 27-30 BHJ morphology control, 25,31-34 and device fabrication engineering.3

#### **Context & Scale**

All-polymer solar cells (all-PSCs), composed of blends of polymer donor (P<sub>D</sub>) and polymer acceptor (P<sub>A</sub>), attract much attention. However, only a few all-PSCs have yielded the power conversion efficiency (PCE) values of >10%, mainly due to the lack of highperformance polymer acceptor and poor PD-PA pair miscibility because of the self-aggregation of polymers. Here, we reported a polymer acceptor PYT with a low optical band gap and a high absorption coefficient and established a friendly relationship between the average molecular weight  $(M_n)$  of PYT and  $P_D$ - $P_A$ miscibility to get an ideal morphology. The medium  $M_n$  of PYT (PYT<sub>M</sub>) leads to the suitable PM6:PYT<sub>M</sub> pair miscibility. Consequently, the resultant device shows an impressive PCE of 13.44%, which is among the top values of all-PSCs. By this research, we point out the true potential of the polymer backbone and molecular mass to further improve the photovoltaic performance in all-PSCs.



## Joule Article



In 1995, Friend and Heeger initially made pioneering independent studies on all-PSCs employed P<sub>A</sub> CN-PPV (Poly[2,5-di(hexyloxy) cyanoterephthalylidene]). 36,37 After that, many efforts have been devoted to the development of PA materials, such as perylene diimide (PDI), <sup>29,38</sup> naphthalene diimide (NDI), <sup>39–41</sup> dicyanobenzothiadiazole moieties, <sup>42,43</sup> B ← N-bridged bipyridine, <sup>44–46</sup> and their derivatives <sup>47,48</sup> based copolymers to enhance the photovoltaic performance of all-PSCs. For instance, in 2007, Zhan et al. synthesized a high-mobility electron-transport polymer based on alternating PDI and dithienothiophene units,<sup>27</sup> which showed a PCE of over 1% when a polythiophene derivative as a donor was used. Furthermore, Sirringhaus and Loi independently pioneered all-PSCs based on the NDI-based  $P_A$  N2200 (poly[[N,N'-bis(2-octyldodecyl)-naphthalene-1,4,5,8-bis(dicarboximide)-2,6-diyl]-alt-5,5'-(2,2'-bithiophene)]). 31,49 However, the NDI- and PDI-based polymer derivatives generally show low absorption coefficients in films, which limited their short-circuit current density  $(J_{sc})$  in devices, <sup>38,39</sup> even though the PCEs of 10%-11% of the N2200-based all-PSCs were achieved by carefully selecting polymer donors, processing solvents, and molecular weights 26,50,51 Another design approach to enhance the absorption coefficient and PCE was reported in 2017 by Li et al., who designed a fused-aromatic-ring-constructed PA PZ1 (Poly[(2,2'-((2Z,2'Z)-((4,4,9,9-tetrahexadecyl-4,9-dihydro-s-indaceno[1,2-b:5,6-b']dithiophene-2,7-diyl) bis-(methanylylidene))bis(3-oxo-2,3-dihydro-1H-indene-2,1-diylidene)) dimalononitrile-alt-2,5-thiophene)]).<sup>43</sup> However, PZ1 as well as its derivatives show the medium optical band gaps ( $E_a^{opt}$ ) in the range of 1.55–1.62 eV. 42,43,48 Obviously, further development of all-PSCs is hampered by the  $P_A$  materials.

Apart from the challenges reported above, achieving suitable BHJ microstructure is also a particularly challenging one for the all-polymer photovoltaic system because the reduced entropic contributions significantly suppress  $P_D$ - $P_A$  pair miscibility. <sup>52</sup> In general, poor  $P_D$ - $P_A$  miscibility results in extreme donor-acceptor phase separation, <sup>27,53–56</sup> limiting charge transport property and thereby promoting charge carrier recombination in the BHJ microstructure. The current studies show that achieving suitable  $P_D$ - $P_A$  miscibility has an enormous impact on the whole process of converting solar light into electricity and consequently determines photovoltaic performance of all-PSCs. <sup>57–59</sup> However, rational design of  $P_A$  has been elusive, because the photovoltaic performance is very sensitive to their number-average molecular weight ( $M_D$ ) values, which has also been highlighted in all-PSCs. <sup>60–63</sup> Overall, finding a  $P_A$  with a low-band-gap and high-absorption coefficient and establishing the friendly relationship between the  $M_D$  of polymer materials and  $P_D$ - $P_A$  miscibility will be very important for the development of highly efficient all-PSCs.

Herein, we report a  $\pi$ -conjugated P<sub>A</sub> PYT (Poly[(2,2'-((2Z,2'Z)-((12,13-bis(2-octyldodecyl)-3,9-diundecyl-12,13-dihydro[1,2,5]thiadiazolo[3,4e]thieno[2'',3'':4',5']thieno [2',3':4,5]pyrrolo[3,2-g]thieno[2',3':4,5]thieno[3,2-b]-indole-2,10-diyl)bis(methanylylidene))bis(3-oxo-2,3-dihydro-1H-indene-2,1-diylidene)) dimalononitrile-alt-2,5-thiophene)]) as shown in Figure 1A, composed of A-DAD-A constructed small molecule acceptor Y5-C20 as the key building block and thiophene as the linking unit. PYT possesses broad absorption with a narrow band gap of approximately 1.40–1.44 eV and high absorption coefficient of over 1.00 ×  $10^5$  cm $^{-1}$ . In addition, we systematically investigate the effect of molecular weight of PYT on the blend morphology and photovoltaic performance of all-PSCs based on a wide-band-gap polymer PM6 (Poly[(2,6-(4,8-bis(5-(2-ethylhexyl-3-fluoro)thiophen-2-yl)-benzo[1,2-b:4,5-b']dithiophene))-alt-(5,5-(1',3'-di-2-thienyl-5',7'-bis(2-ethylhexyl)benzo[1',2'-c:4',5'-c'] dithiophene-4,8-dione)])<sup>64</sup> as donor. The medium  $M_n$  of PYT (PYT<sub>M</sub>) leads to the suitable PM6:PYT<sub>M</sub> pair miscibility. Consequently, the PM6:PYT<sub>M</sub> all-PSCs exhibit more

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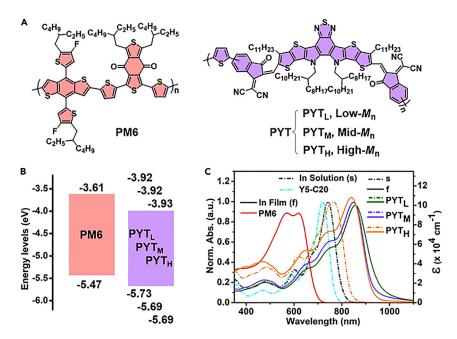


Figure 1. Molecular Structures: Electrochemical and Optical Properties

- (A) Molecular structures of PD PM6 and PA PYT.
- (B) Energy level diagram of PM6 and PYT series.
- (C) Normalized absorption spectra of materials in dilute CF solution and the corresponding absorption coefficients in the thin films.

balanced charge transport property, less carrier recombination loss, and thus enhanced photovoltaic performance as compared to the other two photovoltaic systems based on the PYT with low and high molecular weight (PYT $_{L}$  and PYT $_{H}$ ), respectively. Careful studies are also addressed how the relevant energy loss and blend microstructure determined the device performance. Importantly, the resultant PM6:PYT $_{M}$  device shows an impressive PCE of 13.44%, which is among the top values of all-PSCs.

#### **RESULTS AND DISCUSSION**

#### Synthesis and Characterization of PA PYT

The synthetic routes and chemical structures of the monomers and  $P_A$  PYT are depicted in Scheme S1 in the Supplemental Information. The synthetic details, nuclear magnetic resonance (NMR) spectra and mass spectrometry (MS) spectrum (Figures S1–S7; Supplemental Information) were given in Supplemental Experimental Procedures. The target polymer PYT (Figure 1A) was obtained through a still cross coupling reaction between Y5-C20-Br and 2,5-bis(trimethylstannyl)thiophene. PYT with different molecular weights were synthesized by controlling the reaction time in the Stille polycondensation. The relevant  $M_n$  and polydispersity index (PDI) of three PYT samples were determined to be 7.2 kg mol<sup>-1</sup> and 1.16 (low- $M_n$ , PYT<sub>L</sub>), 12.3 kg mol<sup>-1</sup> and 1.67 (mid- $M_n$ , PYT<sub>M</sub>), and 20.6 kg mol<sup>-1</sup> and 1.86 (high- $M_n$ , PYT<sub>H</sub>), respectively (Figure S8; Table S1). These data were determined by high-performance gel chromatography with polystyrene standards. These PYT series, including PYT<sub>L</sub>, PYT<sub>M</sub>, and PYT<sub>H</sub>, show good solubility in common solvents (e.g., chloroform [CF], and o-dichlorobenzene [o-DCB]) and can be easily cast into uniform films.

The thermogravimetric analysis (TGA) measurements show that PYT series also have excellent thermal stability as evidenced by the high onset thermal





decomposition (5%) temperature of 300°C for PYT<sub>L</sub>, 305°C for PYT<sub>M</sub>, and 339°C for PYT<sub>H</sub>, respectively (Figure S9). In addition, for PYT with different  $M_n$ , melting or glass transition was not observed in the differential scanning calorimetry (DSC) in the temperature range of room temperature (RT) to 300°C (Figure S10). The results indicate its good thermal stability characteristics and the amorphous nature of polymer PYT with different  $M_n$ . Besides, the energy levels of PM6 and PYT series were calculated, as shown in Figure 1B. The lowest unoccupied molecular orbital (LUMO) and highest occupied molecular orbital (HOMO) energy levels of PYT only have a minute difference (Figure S11; Table S2), suggesting the  $M_n$  has little impact on energy levels.

The ultraviolet-visible (UV-vis) absorption spectra of the three PYT series were measured in dilute CF solution and in solid film (Figure 1C), and the corresponding data are summarized in Table S3. In dilute solution, PYT<sub>1</sub> exhibits a similar absorption spectrum to PYT<sub>M</sub>, while the spectra of PYT<sub>H</sub> are redshifted by approximately 23 nm to  $PYT_L$  and  $PYT_M$ , which can be attributed to the more obvious aggregation of polymer PYT caused by a high  $M_{\rm n}$  of 20.6 kg  ${\rm mol}^{-1}$ . It should be noted that the solution absorption spectra of PYT series show a redshift compared to the solution spectra of the small molecular acceptor Y5-C20 (Figure 1C), which results from the extended  $\pi$ -conjugated backbone of the polymers. Due to the  $\pi$ - $\pi$  interactions formed in the solid films, the absorption spectra of the three PYT samples are all redshifted compared to their absorption spectra in solution. However, as shown in Figure 1C, the redshifts of  $PYT_L$  and  $PYT_M$  are 64 and 57 nm, respectively, while PYT<sub>H</sub> are redshifted by only 27 nm, implying that PYT<sub>L</sub> and PYT<sub>M</sub> exhibit stronger  $\pi$ - $\pi$  interactions than the PYT<sub>H</sub>. In addition, as shown in Figure 1C, PM6 and PYT polymers showed complementary absorption and high absorption coefficient in the range of 350 to 900 nm that ensures good light harvesting ability to get a high  $J_{sc}$ .

Figure 1C further displays the UV-vis absorption spectra of the PYT series in films with the absorption coefficients of 1.00  $\times$  10<sup>5</sup> cm<sup>-1</sup> for PYT<sub>L</sub>, 1.03  $\times$  10<sup>5</sup> cm<sup>-1</sup> for  $PYT_M$ , and 1.08 x  $10^5$  cm<sup>-1</sup> for  $PYT_H$ , respectively, indicating that the polymer acceptor PYT with larger  $M_n$  has higher absorption coefficient ( $\varepsilon$ ). <sup>65</sup> Of note is that the PYT series show the broader optical absorption edge corresponding to band gaps ( $E_{\rm g}^{\rm opt}$ ) of 1.40–1.44 eV compared to the Y5-C20 film with a  $E_{\rm g}^{\rm opt}$  of 1.49 eV (Figure S12A; Table S3), even though they exhibit the comparable absorption coefficients. Importantly, it is also found that the  $E_q^{\text{opt}}$  values (1.40–1.44 eV) of PYT series are much lower than those of the highly efficient polymer acceptors PZ1 (1.55 eV)<sup>43</sup> and P-BNBP-fBT (1.86 eV),<sup>46</sup> as shown in Figure S12B. Additionally, the absorption coefficients of PYT series are much higher than those of the efficient polymer acceptors N2200 with a  $\varepsilon$  of 3.48  $\times$  10<sup>4</sup> cm<sup>-1</sup> at 697 nm<sup>20</sup> and DCNBT-IDT (Poly[5,6-dicyano-2,1,3-benzothiadiazole-alt-indacenodithiophene]) with a  $\varepsilon$  of  $6.15 \times 10^5 \text{ cm}^{-1}$  at 759 nm, respectively. 66 This analysis indicates that the PYT series-based devices should have great potential to absorb more energy photons, get higher  $J_{sc}$ , and offer increased power conversion efficiency, which will be discussed below.

We performed grazing incident wide-angle X-ray diffraction (GIWAXS) measurements to further investigate the molecular packing and crystallinity of the pristine PYT<sub>L</sub>, PYT<sub>M</sub>, and PYT<sub>H</sub> films, as shown in Figure 2A. The corresponding in-plane (IP) and out-of-plane (OOP) line cuts of the corresponding neat films are exhibited in Figures 2B and 2C, respectively. The three polymer acceptors adopt a preferential face-on orientation showing more prominent (010) diffraction peaks located at  $q_z = 1.64 \text{ Å}^{-1}$  for PYT<sub>L</sub>,  $q_z = 1.64 \text{ Å}^{-1}$  for PYT<sub>M</sub>, and  $q_z = 1.62 \text{ Å}^{-1}$  for PYT<sub>H</sub> in the





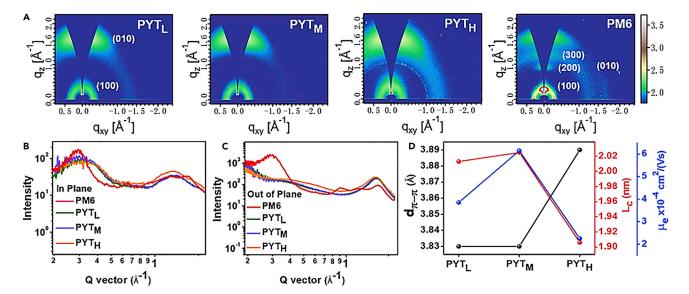


Figure 2. Morphological Analysis of the Pristine Films (A) GIWAXS patterns for the pure films of PYT series and PM6. (B and C) Scattering profiles of IP (B) and OOP (C) for PYT series. (D) The  $d_{\pi^-\pi}$ ,  $L_c$  and  $\mu_e$  of the PYT series as a function of  $M_n$ .

OOP direction (Table S4), which are similar to the reported Y5 acceptor that show preferential face-on ordering.  $^{67}$  A closer look at them reveled that PYT<sub>M</sub> shows a narrower angular distribution of OOP (010) peak, meanwhile PYT<sub>H</sub> has more randomly oriented stacking, and PYT<sub>L</sub> falls in between them, suggesting PYT<sub>M</sub> had the strongest crystalline order compared to the others.  $^{68}$  Besides, for the pristine donor polymer PM6, the (100), (200), and (300) diffraction peaks are present in the OOP direction, whereas the weak (010) peak appears in the IP direction, indicating that face-on and edge-on orientation exist simultaneously with respect to the substrate.

As the tightness of backbone stacking in the OOP direction is an important parameter for photovoltaic materials, we further calculate the OOP (010)  $\pi\text{-}\pi$  stacking distance  $(d_{\pi\text{-}\pi})$  of all three acceptors (Table S4). The plot in Figure 2D exhibits  $d_{\pi\text{-}\pi}$  as a function of  $M_n$ . It was found that a smaller  $d_{\pi\text{-}\pi}$  of 3.83 Å was obtained in the PYT\_L and PYT\_M films than the neat PYT\_H film. This indicates that PYT\_L and PYT\_M show stronger intermolecular aggregations, which is in accordance with the larger redshift in the absorption spectrum of their thin films as mentioned above. In contrast, the  $d_{\pi\text{-}\pi}$  of PYT\_H is significantly increased to 3.89 Å, indicating the unfavorable molecular  $\pi\text{-}\pi$  stacking. Apart from the  $d_{\pi\text{-}\pi}$ , we also use the full-width at half-maximum of the OOP (010) peak to analyze the  $\pi\text{-}\pi$  coherence length ( $L_c$ ). As shown in Figure 2D, the  $L_c$  values of the three acceptors are 2.01 nm for PYT\_L, 2.03 nm for PYT\_M, and 1.91 nm for PYT\_H, respectively, and PYT\_M clearly shows the most ordered features.

Generally, the  $\pi$ - $\pi$  stacking distance and coherence length results imply the charge transport characteristics of the polymer materials. Thus, the electron mobilities of neat films of these three polymer acceptors as a function of  $M_n$  were measured using the space-charge-limited-current (SCLC) method (Figure S13 for electron-only device). The electron mobility of PYT<sub>M</sub> is 6.15 × 10<sup>-4</sup> cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup>, which is higher than those for PYT<sub>L</sub> (3.85 × 10<sup>-4</sup> cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup>) and PYT<sub>H</sub> (2.25 × 10<sup>-4</sup> cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup>). As shown in Figure 2D, the electron mobility in these acceptors



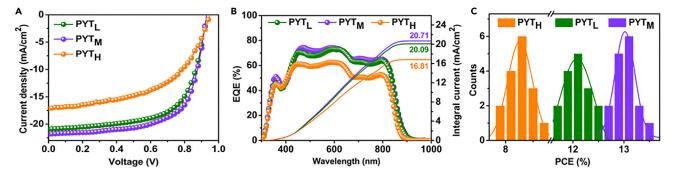


Figure 3. Photovoltaic Characteristics of the PYT-Based All-PSCs

- (A) J-V curves of the PYT-based all-PSCs.
- (B) EQE curves of the PYT-based all-PSCs.
- (C) Histograms of the PCE counts for 16 individual PYT-based devices.

was found to be positively correlated with  $L_{\rm c}$  and also related to  $d_{\pi^-\pi}$ , i.e., the PYT<sub>M</sub> with larger  $L_{\rm c}$  and smaller  $d_{\pi^-\pi}$  has higher electron-only mobility. The results are not only associated with the above-discussed absorption spectra measurements but also imply that PYT<sub>M</sub> is inferred to have much stronger molecular packing properties and electrostatic interactions as compared to the PYT<sub>L</sub> and PYT<sub>H</sub>. It should be noted that although the PYT<sub>L</sub> and PYT<sub>M</sub> samples show the similar  $d_{\pi^-\pi}$  spacing and  $L_{\rm c}$  values, the slightly higher electron mobility of PYT<sub>M</sub> compared to that of PYT<sub>L</sub> is probably attributed to the slightly stronger preferential face-on molecular packing orientation and the different interface morphologies between the pristine films and buffer layers. Thus, it can be predicted that the PYT series as a function of  $M_{\rm n}$  will lead to great changes in optoelectronic property, active layer morphology, charge transport property, and photovoltaic performance.

#### **Photovoltaic Performance of All-PSCs**

To explore the influence of the PYT series with different  $M_n$  on the photovoltaic performance, all-PSCs with a conventional-type structure of indium tin oxide (ITO)/poly(3,4-ethylenedioxythiophene)poly(styrenesulfonate)PEDOT:PSS/PM6:PYT series/PDI functionalized with amino N-oxide (PDINO)/Al were fabricated. To ensure the  $M_n$  of PYT as the unique variable on subsequent tests, the device fabrication parameters, such as D/A weight ratio, material concentration, solvent, and solvent additive, were controlled to the same. Optimized performance was achieved with the same D/A weight ratio of 1:1.2 for all the PYT series and the 0.5 vol % 1-chloronaphthalene (CN) as additive. The detailed procedures for the device fabrication were described in Supplemental Experimental Procedures. The detailed photovoltaic performance of the investigated PM6:PYT all-PSCs based on different spin speeds are shown in Figures S14-S16, and the related photovoltaic parameters are summarized in Tables S5-S7. Additionally, Figure 3A plots the current density-voltage (J-V) curves of the corresponding best-performing solar cells based on the PYT series as acceptors, and Table 1 listed the detailed photovoltaic characteristics of all-PSCs fabricated with the same conditions. All-PSCs based on PM6:PYT<sub>M</sub> exhibit the highest PCE of 13.44% with an open circuit voltage ( $V_{oc}$ ) of 0.93 V, a  $J_{sc}$  of 21.78 mA/cm<sup>2</sup>, and a fill factor (FF) of 66.33%. Notably, the J<sub>sc</sub> and PCE values of the PM6:PYT<sub>M</sub> photovoltaic systems are the highest values reported in literature for all-PSCs based on different polymer acceptors as provided in Figure S17. Of note is that the promising photovoltaic performance of the PM6:PYT<sub>M</sub> system probably results from the low energy loss, which will be discussed below. Slightly lower performance is found for the PM6:PYTL





Table 1. Summary of Photovoltaic Parameters of the Optimized All-PSCs

Devices	V <sub>OC</sub> (V)	$J_{SC}$ (mA/cm <sup>2</sup> )	$J_{SC,EQE}^{a}$ (mA/cm <sup>2</sup> )	FF (%)	PCE (PCE <sup>b</sup> ) (%)
PM6:PYT <sub>L</sub>	0.93	20.92	20.09	64.51	12.55 (12.02 ± 0.53)
PM6:PYT <sub>M</sub>	0.93	21.78	20.71	66.33	13.44 (13.03 ± 0.41)
PM6:PYT <sub>H</sub>	0.95	17.14	16.81	52.91	8.61 (8.29 ± 0.32)

 $<sup>^{\</sup>mathrm{a}}J_{\mathrm{SC,EQE}}$  represents the integrated current density obtained from EQE spectra.

devices with a PCE of 12.55%. However, a much lower PCE of 8.61% was obtained for the PM6:PYT<sub>H</sub> devices, mainly due to lower  $J_{\rm sc}$  and FF values. Furthermore, as compared to PYT<sub>L</sub>- and PYT<sub>M</sub>-based devices with the  $J_{\rm sc}$  values for 20.92 and 21.78 mA cm<sup>-2</sup>, respectively, its lower  $J_{\rm sc}$  value (17.14 mA cm<sup>-2</sup> for the device based on PYT<sub>H</sub>) is confirmed by external quantum efficiency (EQE) measurements for the best devices (Table 1). As compared to the EQE spectra of the PM6:PYT<sub>H</sub> device, the PYT<sub>L</sub>- and PYT<sub>M</sub>-based devices not only possess the increased EQE in the entire absorption range but also show the redshifted EQE spectra.

Figure S18 exhibits the trend in device photocurrent density ( $J_{\rm ph}$ ) measured at  $-5{\rm V}$  with the reverse photocurrent ( $J_{\rm ph}=J_L-J_D$ ) being defined as the difference between the dark ( $J_{\rm D}$ ) and the light current density ( $J_{\rm L}$ ) under one sun. <sup>69</sup> The relative change between  $J_{\rm sc}$  and  $J_{\rm ph}$  ( $-5{\rm V}$ ) is found to be the smaller for the PYT<sub>L</sub>- and PYT<sub>M</sub>-based devices exhibiting negligible voltage dependent recombination in contrast to the PM6:PYT<sub>H</sub> device. Additionally, PYT<sub>L</sub>- and PYT<sub>M</sub>-based systems show a photocurrent of approximately 23 mA cm<sup>-2</sup> at  $-5{\rm V}$ , indicating comparable photogeneration rates. In contrast, the PYT<sub>H</sub>-based blend exhibits a lower photogeneration rate with a  $J_{\rm ph}$  of 19.04 mA cm<sup>-2</sup> at  $-5{\rm V}$ . Since the three blends based on PYT with different  $M_{\rm n}$  possess very similar absorption spectrum and absorption coefficients as shown in Figure S19, the absorption properties were not the primary cause of the different  $J_{\rm sc}$  values (Table 1). Thus, the relative low  $J_{\rm sc}$  and FF values in PM6:PYT<sub>H</sub> devices indicating that this photovoltaic system possess the poor physical mechanisms and unsuitable blend morphology with efficient carrier extraction, which will be discussed below.

#### **Exciton and Carrier Dynamics**

The selection of molecular weight enables changing the D/A interfacial area and fine-tuning the molecular ordering. This, in turn, is likely to sensitively change the fundamental photo-physical processes. Thus, we first investigated PL spectra under 639 nm photo-excitation to understand the influence of the  $M_n$  values of PYT on the exciton dissociation and the efficiency of exciton separation in the corresponding blends. As provided in Figure S20, PL spectroscopy presented quenching efficiencies of 77.6% for PM6:PYT<sub>L</sub> blend and 80.4% for PM6:PYT<sub>M</sub> blend, respectively. However, the quenching efficiency of PM6:PYT<sub>H</sub> blend was only 52.2%, indicating much exciton was useless for charge generation, which can be the limiting factor for the low photocurrent of the devices (Figure 3A). Notably, these relatively modest PL quenching efficiencies of the PM6:PYT<sub>L</sub> and PM6:PYT<sub>M</sub> blends can be attributed most probably to the large domain size and high crystallinity of PYT<sub>L</sub> and PYT<sub>M</sub>, resulting in significant polymer exciton decay to ground during exciton diffusion to the P<sub>D</sub>-P<sub>A</sub> interface. This exciton decay to ground corresponds to an approximately 20% quantum efficiency loss, and is likely to be the main reason of the subunit maximal EQE data (Figure 3B). In addition, the unsuitable PM6/PYTH blend with larger phase separation discussed below indicates the obvious quantum efficiency loss.<sup>69</sup> Further topographical characterization can shed light on this issue, which

<sup>&</sup>lt;sup>b</sup>The average PCE values with standard deviations were obtained from 16 devices.

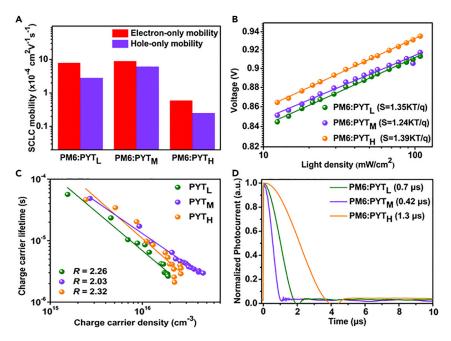


Figure 4. Charge Generation, Transport, Extraction, and Recombination

- (A) The hole and electron mobilities of the PYT-based blend films.
- (B) Measurement of  $V_{oc}$  versus light intensity for the devices.
- (C) Charge carrier lifetime  $\tau,$  obtained from TPV, as a function of charge density n, calculated from
- CE under  $V_{\rm oc}$  conditions (from 0.15 to 2.50 suns). The solid lines represent linear fits of the data. (D) Normalized TPC data for the relevant devices. The illumination pulse intensity was 150 mW cm<sup>-2</sup> (light pulse of 50  $\mu$ s).

we will systematically discuss in the last section. Here, we further calculated the exciton dissociation probabilities  $(P(E,T),P(E,T)=\frac{J_{ph}}{J_{sat}})$ , where  $J_{sat}$  is the saturation current density), which can be used to assessed the exciton dissociation and charge collection efficiency in the corresponding devices. The P(E,T) for the three types of devices were 95.0% (PM6:PYT $_{L}$  blend), 95.9% (PM6:PYT $_{M}$  blend), and 90.2% (PM6:PYT<sub>H</sub> blend), respectively. Here, the exciton dissociation probabilities (charge collection efficiencies) in the relevant devices, in varying degrees, are higher than those of PL quenching efficiencies of the corresponding BHJ blends. On the one hand, this variation in percentage trend may originate from the intrinsic difference between the PL method for measuring exciton dissociation efficiency from a BHJ blend and the P(E,T) model with an actual photovoltaic device. On the other hand, the morphology variation in the film and device may also be a factor, and the various physical mechanisms cannot be totally ruled out. Overall, the PM6:PYT<sub>M</sub> device shows the highest  $J_{ph}/J_{sat}$  value, indicating the more efficient processes of exciton dissociation and charge collection than the other two types of devices.

Then, to gain deeper insight into the impact of  $M_n$  on charge transport in the corresponding blends, SCLC method was also used to map out the charge carrier mobilities investigated by analyzing the J-V characteristics of single carrier devices (Figure S21 for hole-only device and Figure S22 for electron-only device). The hole mobility of PM6:PYT<sub>M</sub> blend is  $6.12 \times 10^{-4}$  cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup>, which is approximately twice higher than that of PM6:PYT<sub>L</sub> blend (2.81  $\times$  10<sup>-4</sup> cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup>) and over one order higher than that of PM6:PYT<sub>H</sub> blend (2.49  $\times$  10<sup>-5</sup> cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup>). Of note is that the neat PM6 film shows a hole mobility of  $6.91 \times 10^{-4}$  cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup>, implying that the





PYT<sub>H</sub> added in the corresponding blend seriously disrupted the hole transport channel. In addition, the average electron mobilities of PM6:PYT<sub>L</sub> and PM6:PYT<sub>M</sub> blends are 7.88 × 10<sup>-4</sup> and 8.91 × 10<sup>-4</sup> cm<sup>2</sup>V<sup>-1</sup>s<sup>-1</sup>, respectively, which is more than one order higher than that of PM6:PYT<sub>H</sub> blend (5.92 × 10<sup>-4</sup> cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup>). In particular, the PM6:PYT<sub>M</sub> blend showed balanced hole and electron mobilities with a  $\mu_h/\mu_e$  value of 0.69, which was closer to 1 than to 0.36 and 0.42 for PM6:PYT<sub>L</sub> and PM6:PYT<sub>H</sub> blends, respectively. Besides, as shown in Figure 4A, the relatively low mobilities of PM6:PYT<sub>H</sub> blend lead to reduced FF values in devices. The sharply reduce of  $\mu_h$  and  $\mu_e$  in PM6:PYT<sub>H</sub> blends should be due to the excessive aggregation of PYT<sub>H</sub> that isolated the acceptor and donor domains, leading to the poor charge carrier transport.

Trap-assisted recombination losses in these blends were first detected by  $V_{oc}$ versus light intensity curves in Figure 4B. The light intensity dependence of  $V_{\rm oc}$ data are fitted to the linear law: $V_{oc} \propto \frac{nkT}{q \ln P_{linht}}$ , 70 where k, T, and q are the Boltzmann constant, the temperature in Kelvin and the elementary charge, respectively. The semi-logarithmic plot of  $V_{oc}$  depends linearly on the light intensity with a slope S of kT/q. In general, when S is larger than kT/q, monomolecular recombination and trap-assisted recombination are involved. The slopes were calculated to be 1.35, 1.24, and 1.39 kT/q for PM6:PYT<sub>L</sub>, PM6:PYT<sub>M</sub>, and PM6:PYT<sub>H</sub>, respectively, indicating fewer monomolecular and trap-assisted recombination in the PM6:PYT<sub>M</sub> device than the PYT<sub>I</sub>- and PYT<sub>H</sub>-based devices. In addition, we evaluated the non-geminate recombination order R with the transient photovoltage (TPV) and charge extraction (CE) techniques. The R value ( $R = \lambda + 1$ ) can be calculated via the equation  $\tau = \tau_0 \left(\frac{n_0}{n}\right)^{\lambda}$ , 70 where  $\tau_0$  and  $n_0$  are constants and  $\lambda$  is the so-called recombination exponent. As shown in Figure 4C, we plot the charge carrier lifetimes obtained from TPV measurements (Figure S23) as a function of charge carrier density obtained from CE measurements (Figure S24). A lower recombination order value (R = 2.03) for the PM6:PYT<sub>M</sub> device was achieved, which can be directly correlated to the effective exciton dissociation probability and the spatial uniform charge carrier distributions through active layer. It also strongly implies the lower carrier recombination loss than the other two types of devices with the R values of 2.26 for PM6:PYT<sub>1</sub> and 2.32 for PM6:PYT<sub>H</sub>. Besides, apart from the  $J_{ph}$  and  $J_{sat}$  measurements discussed above, more detailed information about the CE property in the devices can be performed by using transient photocurrent (TPC) measurements. As provided in Figure 4D, the CE time are 0.70  $\mu$ s for the PM6:PYT<sub>L</sub> blend, 0.41  $\mu$ s for the PM6:PYT<sub>M</sub> blend, and 1.3 µs for the PM6:PYT<sub>H</sub> blend, respectively. The shorter extraction lifetime suggests that photo-generated carriers are extracted more efficiently in the PM6:PYT<sub>M</sub> device than the PYT<sub>L</sub>- and PYT<sub>H</sub>-based devices. As a result, the reduced charge recombination and extraction lifetime afforded the increased FF and  $J_{\rm sc}$  values in the PM6:PYT<sub>M</sub> device as compared with the other two types of devices.

#### **Energetic Losses**

In order to investigate the intrinsic reasons that cause the dramatically changed photovoltaic performance of the devices, here, we systematically studied the energy loss ( $E_{\rm loss}$ ) issue of the all-PSCs based on the polymer acceptors PYT with different  $M_{\rm n}$ . We first tested the electroluminescence (EL) of PM6:PYT devices with active layer thicknesses of 100  $\pm$  5 nm that was measured under a constant current injection density of 20 mA cm<sup>-2</sup>. As shown in Figure S25, the higher EL intensity of PM6:PYT<sub>M</sub> device as compared to the PM6:PYT<sub>L</sub> and PM6:PYT<sub>H</sub> devices indicates a relatively higher emissive quantum yield for the charge-transfer (CT) state <sup>71,72</sup>





Table 2. Detailed  $E_{loss}$  and  $V_{loss}$  Parameters of the All-PSCs Based on PM6:PYT Blends

Devices	E <sub>gap</sub> (eV)	V <sub>OC</sub> (V)	E <sub>loss</sub> (eV)	V <sub>oc</sub> <sup>SQa</sup> (V)	△E <sub>1</sub> (eV)	V <sub>oc</sub> (V)	$\Delta E_2^{c}$ (eV)	$\Delta E_3^{\mathbf{d}}$ (eV)
PM6:PYT <sub>L</sub>	1.457	0.93	0.532	1.199	0.258	1.154	0.045	0.229
PM6:PYT <sub>M</sub>	1.466	0.93	0.528	1.207	0.259	1.148	0.058	0.210
PM6:PYT <sub>H</sub>	1.492	0.95	0.544	1.232	0.260	1.170	0.062	0.222

 $<sup>^{</sup>a}V_{oc}^{SQ}$ : Schokley-Queisser limit to  $V_{oc}$ .

for PM6:PYT<sub>M</sub> devices. Furthermore, we quantified the optical band gaps  $(E_{qap})$  of all the blend films based on the cross point between the absorption spectra and EL of blends.<sup>73</sup> As shown in Figure S26, the E<sub>gap</sub> of PM6:PYT<sub>L</sub>, PM6:PYT<sub>M</sub>, and PM6:PYT<sub>H</sub> blends is 1.457, 1.466, and 1.492 eV, respectively. The total  $E_{loss}$  was further determined according to  $E_{loss} = E_{gap} - qV_{oc}$ , where  $E_{gap}$  is determined from the EQE onset,  $^{74-76}$  and the results are summarized in Table 2. We note that the  $M_{\rm n}$  of PYT acceptors have obvious impacts on  $E_{loss}$ , leading to values of 0.532 eV for PYT<sub>L</sub>, 0.528 eV for PYT<sub>M</sub>, 0.544 eV for PYT<sub>H</sub>, respectively. As far as we know, the photon energy loss may facilitate charge separation and improve the exciton dissociation efficiency.  $^{72,75}$  Here, we provide how the EQE relates to  $E_{loss}$  for a range of successful all-polymer photovoltaic systems, as exhibited in Figure S27. Here, we also present the result of the optimized PM6:PYT<sub>M</sub> device with a  $E_{loss}$  of 0.472 eV (Table S9). Although the  $E_{loss}$  of the PM6:PYT<sub>M</sub> devices is the smallest among the published all-polymer systems, this system shows a high EQE of 74%, which is better than most all-polymer systems provided in Figure S27. In addition, there are several mechanisms that can contribute to a low  $E_{\rm loss}$  while maintaining a high EQE. Among these are reducing energetic disorder, eliminating defects that can trap charges, reducing the D/A interface area and so on. 71,72,74,75,77 As expected, further material design and morphology control is imperative to increase the EQE and efficiency of all-PSCs meanwhile keep a low photo energy loss.

In order to probe the impact of  $M_n$  of PYT acceptors on the  $E_{\rm loss}$  of devices, we carried out the Fourier transform photocurrent spectroscopy (FTPS) and EL measurements<sup>77</sup> of the three devices (Figure 5). The detailed components of  $E_{\rm loss}$  can be categorized into three parts based on the SQ limit, as shown in Equation 1.<sup>78</sup>

$$\begin{split} E_{loss} &= E_{gap} - qV_{oc} = \Delta E_1 + \Delta E_2 + \Delta E_3 \\ &= \left(E_{gap} - qV_{oc}^{SQ}\right) + \left(qV_{oc}^{SQ} - qV_{oc}^{rad}\right) + \left(qV_{oc}^{rad} - qV_{oc}\right) \end{split} \tag{Equation 1}$$

Through disentangling the different loss parts to the total  $E_{loss}$  (here,  $E_{loss} = \Delta E_1 + \Delta E_2 + \Delta E_3$ with  $\Delta E_1$  arising from the loss related to the SQ-limit, and  $\Delta E_2$  and  $\Delta E_3$  related to the loss via radiative recombination below band gap and non-radiative recombination, respectively), we found  $\Delta E_2$  and  $\Delta E_3$  in the three devices are obviously affected, as shown in Table 2. Among them, the values of  $\Delta E_2$  gradually increase as the molecular weight of PYT increases. The large  $\Delta E_2$  value of 0.062 eV in the PM6:PYT<sub>H</sub> device mean the radiative recombination loss from the absorption below the band gap is significant. The  $\Delta E_3$  values were 0.229 eV for PM6:PYT<sub>L</sub> device, 0.210 eV for PM6:PYT<sub>M</sub> device, and 0.222 eV for PM6:PYT<sub>H</sub> device, respectively. It is found that the PM6:PYT<sub>M</sub> device showed the reduced  $\Delta E_3$  in comparison with the other two types of devices. Of note, both radiative and non-radiative losses of these devices are relatively low ( $\Delta E_2$  of 0.04–0.07 and  $\Delta E_3$  of 0.21–0.23 eV), probably resulting from the low driving energy ( $\Delta E_{S_1,CT}$ ), which is generally assumed to be correlated to the difference in the HOMO energies of the D/A materilas ( $\Delta E_{HOMO}$  =

 $<sup>{}^{\</sup>rm b}V_{\rm oc}^{\rm rad}$ : radiative limit to  $V_{\rm oc}$ , measured using EQE<sub>EL</sub>.

 $<sup>^</sup>c \Delta E_2 \ (\Delta V_{oc}^{rad})$ : voltage losses due to non-ideal absorption (it was calculated from EL and FTPS measurements)

 $<sup>^{\</sup>rm d}\Delta E_3$  ( $\Delta V_{oc}^{non-rad}$ ): voltage losses due to non-radiative recombination only.





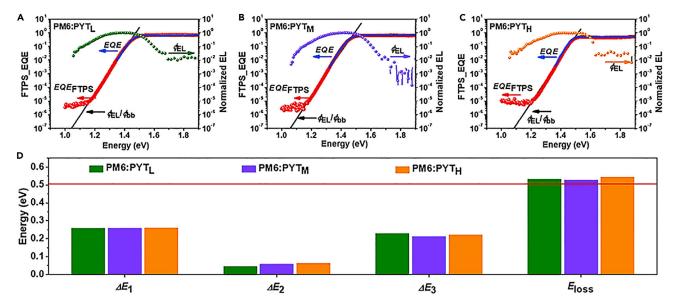


Figure 5. Energy Loss Analysis in All-PSCs

(A–C) Semi-logarithmic plots of normalized EL, measured EQE and EQE calculated by FTPS (EQE<sub>FTPS</sub>) as a function of energy for devices based on PM6:PYT<sub>L</sub> (A), PM6:PYT<sub>M</sub> (B), and PM6:PYT<sub>H</sub> (C). The ratio of  $\varphi_{EL}/\varphi_{bb}$  was used to plot the EQE in the low-energy regime (black line), where  $\varphi_{EL}$  and  $\varphi_{bb}$  represent the emitted photon flux and the room-temperature blackbody photon flux, respectively. Of note is that  $\varphi_{EL}/\varphi_{bb}$  follows experimental EQE<sub>FTPS</sub> faithfully at higher energies as expected from reciprocity.

(D)  $E_{loss}$  and its detailed three parts of  $\Delta E_1$ ,  $\Delta E_2$ , and  $\Delta E_3$  values.

 $E_{HOMO,D} - E_{HOMO,A}$ ). On the basis of the smaller non-radiative recombination loss in the PM6:PYT<sub>M</sub> device, it suggests that with the incorporated P<sub>A</sub> PYT with a medium  $M_n$  having a good compatibility with the polymer donor (P<sub>D</sub>) PM6, not only the PL quenching efficiency and charge transport property are improved, but also the energetic disorder at the D/A interface is decreased, which promotes interfacial charge transfers and dissociation. These merits as mentioned above likely lead to the enhancement of  $J_{sc}$  and FF in the PM6:PYT<sub>M</sub> devices as compared to the other two types of devices, especially for the PM6:PYT<sub>H</sub> devices.

Overall, the results of the exciton dissociation properties and probabilities coupled with the analysis of the charge carrier mobilities and the carrier recombination dynamic as well as the evaluation of energy losses in these devices provide detailed insight into subtle mechanisms being responsible for device parameters as a function of  $M_n$ . Actually, these relationships between physical dynamic and device performance are mainly determined by the morphology characterizations, which will be discussed in detail later.

#### Structural and Morphological Properties of Active Layers

From our spectroscopic measurements thus far, we have found that the PM6:PYT<sub>H</sub> blend suffered from exciton harvesting limitations, whereas exciton harvesting and charge generation can be considered efficient in the PM6:PYT<sub>L</sub> and PM6:PYT<sub>M</sub> blends. We have also demonstrated that the PM6:PYT<sub>H</sub> device endured poorer charge transport and extraction, and more carrier trapping than the PM6:PYT<sub>L</sub> and PM6:PYT<sub>M</sub> devices. To further demonstrate the effect of different  $M_n$  on the above-mentioned photovoltaic performance and physical dynamics in the three devices and correlate them to the corresponding morphological characterizations, tapping-mode atomic force microscopy (AFM) and GIWAXS have been carried out.

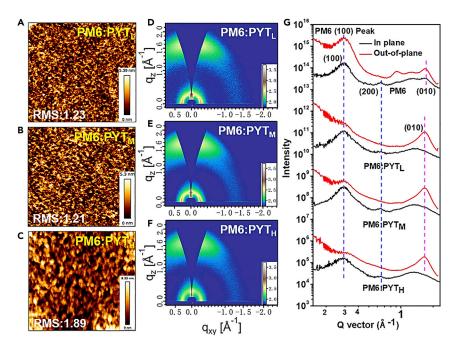


Figure 6. Morphological Properties of Active Layers (A–C) Contact mode AFM surface scans ( $5 \times 5 \, \mu m^2$ ) of PM6:PYT<sub>L</sub> (A), PM6:PYT<sub>M</sub> (B), and PM6:PYT<sub>H</sub> (C). (D–F) GIWAXS patterns for PM6:PYT<sub>L</sub> (D), PM6:PYT<sub>M</sub> (E), and PM6:PYT<sub>H</sub> (F). (G) Scattering profiles of IP and OOP for PM6 and PM6:PYT blends.

Both  $PM6:PYT_L$  and  $PM6:PYT_M$  blends exhibit a smooth surface with a small root mean square (RMS) roughness of 1.23 and 1.21 nm, respectively (Figures 6A and 6B). In addition, it can be found that these two blends form defined phase separation with a bi-continuous interpenetrating network, which is beneficial to suppress the exciton recombination and improve the charge separation and transport properties as mentioned above. However, further increasing  $M_n$  of PYT results in more rough film (a RMS roughness of 1.89 nm) with significant phase separation and large domain size (Figure 6C). Note that the large domain sizes on the length scale of hundred nanometers in the PM6:PYTH blend can seriously result in regions of charge trapping, poor charge separation process, and other losses discussed earlier, which significantly influences the  $J_{\rm sc}$  and FF values of the PM6:PYT<sub>H</sub> devices. In order to further analyze the different origins of the blend morphologies and compare the compatibility of the  $P_D$  and PYT with different  $M_n$ , we delineate the critical role of miscibility expressed in terms of the Flory-Huggins interaction parameter ( $\chi$ ).<sup>80</sup> The  $\chi$  values, between donor and acceptors, are calculated from experimentally measured contact angles (Figure S28, using two solvents, water and ethylene glycol [EG]), and summarized in Table S10. As a consequence, the calculated  $\chi$  values for PYT<sub>L</sub>, PYT<sub>M</sub>, and PYT<sub>H</sub> showed 0.003, 0.066, and 0.200 for blended with PM6. It means that the PM6:PYT<sub>L</sub> and PM6:PYT<sub>M</sub> blends have a more miscible due to relatively low  $\chi$  values compared to that of PM6:PYT<sub>H</sub> blend. In contrast, a high  $\chi$ value of 0.200 implies the severe phase segregation in the PM6:PYT<sub>H</sub> system, which have been demonstrated by the above-discussed AFM image (Figure 6C). These trends were also supported by comparing the PL quenching efficiency of the three different blends (Figure \$20). Thus, it is expected that a proper phase separation of the polymer blends is very important to realize high-performance all-PSCs.

To further investigate the crystallization at nanoscale, GIWAXS measurements of the PM6:PYT<sub>L</sub> (Figure 6D), PM6:PYT<sub>M</sub> (Figure 6E), and PM6:PYT<sub>H</sub> (Figure 6F) blends





were carried out. As mentioned above, PM6 adopts bimodal distribution, evidenced by the simultaneous appearance of (100) lamellar diffraction and (010)  $\pi$ - $\pi$  stacking peak along both IP and OOP directions, while the three polymers showed a preferential face-on molecular packing orientation (Figures 2B and 2C). Interestingly, for the blends, the (010) and (100) peaks are located at both, the OOP and IP directions (Figure 6G), respectively. When compared with the pristine films, the blends of PYT<sub>L</sub> and PYT<sub>M</sub> showed a stronger intensity of (010) peaks, while the PYT<sub>H</sub> exhibited weaker face-on stacking. The GIWAXS data of (100)-IP peaks and (010)-OOP peaks using the Gaussian function are fitted first, as shown in Figures S29 and S30, respectively. And the relevant fitting data are summarized in Tables S11 and S12, respectively. The detailed results of the (100)-IP peak's fitting patterns are in good agreement with the GIWAXS measurements of neat films. This indicates that PM6 and PYT within the blend have kept the same orientation as the pure film, giving their preferential orientations simultaneously. In addition, with Gaussian fitting, there is only one peak found for (010)-OOP peaks in relevant blends. The (010)-OOP peaks of PYT series have a better fit, but not exactly fit, as shown in Figure S30. In Table S12, the data are also well match with the GIWAXS data of neat film.

In Figure 6G, we can also find the lamellar packing peaks were still clear in the PM6:PYT<sub>L</sub> and PM6:PYT<sub>M</sub> blends, demonstrating (010)-OOP peaks of PYT<sub>L</sub> and PYT<sub>M</sub>, and (100) peak from PM6, respectively, at the same position as in the pristine materials, thereby implying that when blended with PM6, PYT<sub>L</sub> and, PYT<sub>M</sub> molecules favored to form more ordered arrangement. In contrast, the PM6:PYT<sub>H</sub> exhibits a different behavior, indicating a relatively strong interaction between PM6 and PYT<sub>H</sub> because of its relatively longer rigid backbones.

It should be noted that the PM6:PYT<sub>M</sub> have the smallest (100) lamellar stacking distance of 21.37 Å and (010)  $\pi$ - $\pi$  stacking distance of 3.78 Å (Table S13) as compared to the PM6:PYT<sub>1</sub> and PM6:PYT<sub>H</sub> blends, indicating that it adopted closer and more ordered molecular orientation. In addition, the  $PM6:PYT_M$  blend remains in the preferred (200) lamellar diffraction of PM6 along the IP direction, responsible for the good charge transport of the all-PSCs. Besides, the coherence length  $(L_c)$ values, determined from the full-width-at-half-maximum of X-ray reflection (the Scherrer equation), are 19.5 Å for the PM6:PYT<sub>L</sub> blend, 19.5 Å for the PM6:PYT<sub>M</sub> blend, and 17.5 Å for the PM6:PYT<sub>H</sub> blend, respectively (Table S13). Since the  $L_c$ is related to the crystalline size and lattice order,  $^{81}$  the larger  $L_c$  values of the corresponding PM6:PYT<sub>L</sub> and PM6:PYT<sub>M</sub> blends have less grain boundaries and therefore render better charge transport properties than the PM6:PYT<sub>H</sub> blend. Since the PM6:PYT<sub>H</sub> blend has the smallest coherence length of 17.5 Å, which are identical to the trend of SCLC analysis (Figure 4A). It is noteworthy that the smaller  $L_c$  value of the PM6:PYT<sub>H</sub> blend can be the critical fact of the poorer  $J_{sc}$  and FF values as compared to those of the PYT<sub>L</sub>- and PYT<sub>M</sub>-based blends (Figure 3A; Table 1). Thus, it can be seen that the morphological characterizations of these blends are consistent with the physical dynamics and device performance of all-PSCs observed above.

#### Conclusion

In this work, applying a fused-ring electron acceptor Y5-C20 as the key building block and thiophene as the  $\pi$ -bridges, we designed and synthesized a  $\pi$ -conjugated P<sub>A</sub> PYT. In particular, the absorption coefficients of PYT gradually increase as its molecular weight increases, but the absorption edge of PYT blueshifts. These drive us to synthesize a series of PYT polymer acceptors with controlled





 $M_{\rm D}$  values for fine-tuning the molecular crystallinity and  $P_{\rm D}$ - $P_{\rm A}$  miscibility. When fabricated into all-PSCs with PM6, we observed a clear molecular weight dependence on device performance. Optimized devices based on PYTL, PYT<sub>M</sub>, and PYT<sub>H</sub> exhibits PCEs of 12.55%, 13.44%, and 8.61%, respectively. As compared to PM6:PYT<sub>L</sub> and PM6:PYT<sub>H</sub> systems, the higher device efficiency of PYT<sub>M</sub> originates from the more efficient exciton dissociation and charge generation, higher carrier mobility, faster CE in combination with suppressed carrier recombination, and reduced energy loss and is further assigned to a better blend morphology with suitable phase separation. Whereas PM6:PYT<sub>H</sub> system shows a very poor device efficiency, mainly resulting from the poor morphology with large domains observed above. The presented results not only point out a promising design strategy of polymer acceptor tuned through molecular engineering on the fused-ring electron acceptor units of the small molecule acceptors but also highlight the importance of the molecular weight of polymers as a vital factor for high-performance all-PSCs. Therefore, the molecular backbone of polymer as well as its molecular mass are necessary to carefully understand their true performance potential and promote further advances in the field of PSCs.

#### **EXPERIMENTAL PROCEDURES**

Full details of experimental procedures can be found in the Supplemental Information.

#### **SUPPLEMENTAL INFORMATION**

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#### **AUTHOR CONTRIBUTIONS**

W.W. and Q.W. conceived the ideas and coordinated the work. Q.W. fabricated all the solar cell samples and conducted the measurements. W.W. and Y.W. contributed to the acceptor polymer-based all-polymer solar cell design and data analysis. R.S. performed the transient measurement unit (TMU) and AFM measurements, and J.G. performed FTPS and PL measurements. M.S., W.Y., and H.L. performed devices morphology measurements. W.W., Q.W., and J.M. contributed to manuscript preparation, and W.W. and Q.W., supervised by J.M., conceived and directed the project. All authors commented on the manuscript.

#### **DECLARATION OF INTERESTS**

The authors declare no competing interests.

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