

On a new method towards proof of Riemann's Hypothesis

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Abstract

We consider the analytic continuation of Riemann's Zeta Function derived from **Riemann's Xi function** $\xi(s)$ which is evaluated at $s = \frac{1}{2} + \sigma + i\omega$, given by $\xi(\frac{1}{2} + \sigma + i\omega) = E_{p\omega}(\omega)$, where σ, ω are real and $-\infty \leq \omega \leq \infty$ and compute its inverse Fourier transform given by $E_p(t)$.

We use a new method and show that the Fourier Transform of $E_p(t)$ given by $E_{p\omega}(\omega) = \xi(\frac{1}{2} + \sigma + i\omega)$ **does not have zeros** for finite and real ω when $0 < |\sigma| < \frac{1}{2}$, corresponding to the critical strip **excluding** the critical line and prove Riemann's hypothesis.

More importantly, the new method **does not** contradict the existence of non-trivial zeros on the critical line with real part of $s = \frac{1}{2}$ and **does not** contradict Riemann Hypothesis. It is shown that the new method is **not** applicable to Hurwitz zeta function and related functions and **does not** contradict the existence of their non-trivial zeros away from the critical line.

If the specific solution presented in this paper is incorrect, it is **hoped** that the new method discussed in this paper will lead to a correct solution by other researchers.

Keywords: Riemann, Hypothesis, Zeta, Xi, exponential functions

1. Introduction

It is well known that Riemann's Zeta function given by $\zeta(s) = \sum_{m=1}^{\infty} \frac{1}{m^s}$ converges in the half-plane where the real part of s is greater than 1. Riemann proved that $\zeta(s)$ has an analytic continuation to the whole s-plane apart from a simple pole at $s = 1$ and that $\zeta(s)$ satisfies a symmetric functional equation given by $\xi(s) = \xi(1-s) = \frac{1}{2}s(s-1)\pi^{-\frac{s}{2}}\Gamma(\frac{s}{2})\zeta(s)$ where $\Gamma(s) = \int_0^{\infty} e^{-u}u^{s-1}du$ is the Gamma function.^{[4] [5]} We can see that if Riemann's Xi function has a zero in the critical strip, then Riemann's Zeta function also has a zero at the same location. Riemann made his conjecture in his 1859 paper, that all of the non-trivial zeros of $\zeta(s)$ lie on the critical line with real part of $s = \frac{1}{2}$, which is called the Riemann Hypothesis.^[1]

Hardy and Littlewood later proved that infinitely many of the zeros of $\zeta(s)$ are on the critical line with real part of $s = \frac{1}{2}$.^[2] It is well known that $\zeta(s)$ does not have non-trivial zeros when real part of $s = \frac{1}{2} + \sigma + i\omega$, given by $\frac{1}{2} + \sigma \geq 1$ and $\frac{1}{2} + \sigma \leq 0$. In this paper, **critical strip** $0 < \text{Re}[s] < 1$ corresponds to $0 \leq |\sigma| < \frac{1}{2}$.

In this paper, a **new method** is discussed and a specific solution is presented to prove Riemann's Hypothesis. If the specific solution presented in this paper is incorrect, it is **hoped** that the new method discussed in this paper will lead to a correct solution by other researchers.

In Section 2, we prove Riemann's hypothesis by taking the analytic continuation of Riemann's Zeta Function derived from Riemann's Xi function $\xi(\frac{1}{2} + \sigma + i\omega) = E_{p\omega}(\omega)$ and compute inverse Fourier transform of $E_{p\omega}(\omega)$ given by $E_p(t)$ and show that its Fourier transform $E_{p\omega}(\omega)$ does not have zeros for finite and real ω when $0 < |\sigma| < \frac{1}{2}$, corresponding to the critical strip **excluding** the critical line.

In Appendix A to Appendix D, well known results which are used in this paper are re-derived.

We present an **outline** of the new method below.

1.1. Step 1: Inverse Fourier Transform of $\xi(\frac{1}{2} + i\omega)$

Let us start with Riemann's Xi Function $\xi(s)$ evaluated at $s = \frac{1}{2} + i\omega$ given by $\xi(\frac{1}{2} + i\omega) = \Xi(\omega) = E_{0\omega}(\omega)$, where $-\infty \leq \omega \leq \infty$. Its inverse Fourier Transform is given by $E_0(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} E_{0\omega}(\omega) e^{i\omega t} d\omega$, where ω, t are real, as follows (link).^[3]

$$E_0(t) = \Phi(t) = 2 \sum_{n=1}^{\infty} [2n^4 \pi^2 e^{\frac{9t}{2}} - 3n^2 \pi e^{\frac{5t}{2}}] e^{-\pi n^2 e^{2t}} = 2 \sum_{n=1}^{\infty} [2\pi^2 n^4 e^{4t} - 3\pi n^2 e^{2t}] e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}} \quad (1)$$

We see that $E_0(t) = E_0(-t)$ is a real and **even** function of t , given that $E_{0\omega}(\omega) = E_{0\omega}(-\omega)$ because $\xi(s) = \xi(1-s)$ and hence $\xi(\frac{1}{2} + i\omega) = \xi(\frac{1}{2} - i\omega)$ when evaluated at $s = \frac{1}{2} + i\omega$.

The inverse Fourier Transform of $\xi(\frac{1}{2} + \sigma + i\omega) = E_{p\omega}(\omega)$ is given by the real function $E_p(t)$. We can write $E_p(t)$ as follows for $0 < |\sigma| < \frac{1}{2}$ and this is shown in detail in Appendix A using contour integration.

$$E_p(t) = E_0(t) e^{-\sigma t} = 2 \sum_{n=1}^{\infty} [2\pi^2 n^4 e^{4t} - 3\pi n^2 e^{2t}] e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}} e^{-\sigma t} \quad (2)$$

We can see that $E_p(t)$ is an analytic function in the interval $|t| \leq \infty$, given that the sum and product of exponential functions are analytic in the same interval and hence infinitely differentiable in that interval.

1.2. Step 2: On the zeros of a related function $G(\omega)$

Statement 1: Let us assume that Riemann's Xi function $\xi(\frac{1}{2} + \sigma + i\omega) = E_{p\omega}(\omega)$ has a zero at $\omega = \omega_0$ where ω_0 is real and finite and $0 < |\sigma| < \frac{1}{2}$, corresponding to the critical strip excluding the critical line. We will prove that this assumption leads to a **contradiction**.

Let us consider $0 < \sigma < \frac{1}{2}$ at first. Let us consider a new function $g(t) = f(t) e^{-\sigma t} u(-t) + f(t) e^{\sigma t} u(t)$, where $f(t) = e^{-2\sigma t_0} f_1(t) + e^{2\sigma t_0} f_2(t)$ and $f_1(t) = e^{\sigma t_0} E'_p(t+t_0)$ and $f_2(t) = e^{-\sigma t_0} E'_p(t-t_0)$ and $E'_p(t) = e^{-\sigma t_2} E_p(t-t_2) - e^{\sigma t_2} E_p(t+t_2)$ and t_0, t_2 are real and $g(t)$ is a real function of variable t and $u(t)$ is Heaviside unit step function. We can see that $g(t)h(t) = f(t)$ where $h(t) = [e^{\sigma t} u(-t) + e^{-\sigma t} u(t)]$.

In Section 2.1, we will show that the Fourier transform of the **even function** $g_{\text{even}}(t) = \frac{1}{2}[g(t) + g(-t)]$ given by $G_R(\omega)$ must have **at least one zero** at $\omega = \omega_z(t_2, t_0) \neq 0$, for every value of t_0 , to satisfy Statement 1, where $\omega_z(t_2, t_0)$ is real and finite.

1.3. Step 3: On the zeros of the function $G_R(\omega)$

In Section 2.2, we compute the Fourier transform of the function $g(t)$ and compute its real part given by $G_R(\omega) = G_R(\omega, t_2, t_0)$ and we can write as follows.

$$G_R(\omega, t_2, t_0) = e^{-2\sigma t_0} \int_{-\infty}^0 [E'_0(\tau + t_0) e^{-2\sigma \tau} + E'_{0n}(\tau - t_0)] \cos(\omega \tau) d\tau \\ + e^{2\sigma t_0} \int_{-\infty}^0 [E'_0(\tau - t_0) e^{-2\sigma \tau} + E'_{0n}(\tau + t_0)] \cos(\omega \tau) d\tau \quad (3)$$

We require $G_R(\omega, t_2, t_0) = 0$ for $\omega = \omega_z(t_2, t_0)$ for every value of t_0 , for **each fixed value** of t_2 , to satisfy **Statement 1**. In general $\omega_z(t_2, t_0) \neq \omega_0$. Hence we can see that $P(t_2, t_0) = G_R(\omega_z(t_2, t_0), t_0) = 0$.

1.4. **Step 4: Zero Crossing function** $\omega_z(t_2, t_0)$ **is an even function of variable** t_0

In Section 2.3, we show the result in Eq. 4 and that $\omega_z(t_2, t_0) = \omega_z(t_2, -t_0)$. It is shown that $P(t_2, t_0) = G_R(\omega_z(t_2, t_0), t_0) = P_{\text{odd}}(t_2, t_0) + P_{\text{odd}}(t_2, -t_0)$ is an **odd** function of t_0 , for all t_0 , for a given value of t_2 as follows.

$$\begin{aligned} P_{\text{odd}}(t_2, t_0) &= [\cos(\omega_z(t_2, t_0)t_0) \int_{-\infty}^{t_0} E'_0(\tau) e^{-2\sigma\tau} \cos(\omega_z(t_2, t_0)\tau) d\tau \\ &\quad + \sin(\omega_z(t_2, t_0)t_0) \int_{-\infty}^{t_0} E'_0(\tau) e^{-2\sigma\tau} \sin(\omega_z(t_2, t_0)\tau) d\tau] \\ &+ e^{2\sigma t_0} [\cos(\omega_z(t_2, t_0)t_0) \int_{-\infty}^{t_0} E'_{0n}(\tau) \cos(\omega_z(t_2, t_0)\tau) d\tau + \sin(\omega_z(t_2, t_0)t_0) \int_{-\infty}^{t_0} E'_{0n}(\tau) \sin(\omega_z(t_2, t_0)\tau) d\tau] \end{aligned} \quad (4)$$

1.5. **Step 5: Final Step**

In Section 3, we set $t_0 = t_{0c}$ and $t_2 = t_{2c} = 2t_{0c}$, such that $\omega_z(t_{2c}, t_{0c})t_{0c} = \frac{\pi}{2}$ and substitute in the equation for $P_{\text{odd}}(t_2, t_0)$ in Eq. 4 and show that this leads to the result in Eq. 5. We use $E'_0(t) = E_0(t - t_2) - E_0(t + t_2)$ and $E'_{0n}(t) = E'_0(-t)$.

$$\int_0^{t_{0c}} (E_0(\tau - t_{2c}) - E_0(\tau + t_{2c})) (\cosh(2\sigma t_{0c}) - \cosh(2\sigma\tau)) \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau = 0 \quad (5)$$

We show that the **each** of the terms in the integrand in Eq. 5 are **greater than zero**, in the interval $\tau = [0, t_{0c}]$ where $t_{0c} > 0$. For $\omega_z(t_{2c}, t_{0c})t_{0c} = \frac{\pi}{2}$, we see that $\omega_z(t_{2c}, t_{0c})\tau = \frac{\pi}{2t_{0c}}\tau$ lies in the range $[0, \frac{\pi}{2}]$ and hence $\sin(\omega_z(t_{2c}, t_{0c})\tau) > 0$ in that interval $\tau = [0, t_{0c}]$.

Hence the result in Eq. 5 leads to a **contradiction** for $0 < \sigma < \frac{1}{2}$.

We have shown this result for $0 < \sigma < \frac{1}{2}$ and then use the property $\xi(\frac{1}{2} + \sigma + i\omega) = \xi(\frac{1}{2} - \sigma - i\omega)$ to show the result for $-\frac{1}{2} < \sigma < 0$. Hence we have produced a **contradiction** of **Statement 1** that the Fourier Transform of the function $E_p(t) = E_0(t)e^{-\sigma t}$ has a zero at $\omega = \omega_0$ for $0 < |\sigma| < \frac{1}{2}$.

2. An Approach towards Riemann's Hypothesis

Theorem 1: Riemann's Xi function $\xi(\frac{1}{2} + \sigma + i\omega) = E_{p\omega}(\omega)$ does not have zeros for any real value of $-\infty < \omega < \infty$, for $0 < |\sigma| < \frac{1}{2}$, corresponding to the critical strip excluding the critical line, given that $E_0(t) = E_0(-t)$ is an even function of variable t , where $E_p(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} E_{p\omega}(\omega) e^{i\omega t} d\omega$, $E_p(t) = E_0(t)e^{-\sigma t}$ and $E_0(t) = 2 \sum_{n=1}^{\infty} [2\pi^2 n^4 e^{4t} - 3\pi n^2 e^{2t}] e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}}$.

Proof: We assume that Riemann Hypothesis is false and prove its truth using proof by contradiction.

Statement 1: Let us assume that Riemann's Xi function $\xi(\frac{1}{2} + \sigma + i\omega) = E_{p\omega}(\omega)$ has a zero at $\omega = \omega_0$ where ω_0 is real and finite and $0 < |\sigma| < \frac{1}{2}$, corresponding to the critical strip excluding the critical line. We will prove that this assumption leads to a **contradiction**.

We will prove it for $0 < \sigma < \frac{1}{2}$ first and then use the property $\xi(\frac{1}{2} + \sigma + i\omega) = \xi(\frac{1}{2} - \sigma - i\omega)$ to show the result for $-\frac{1}{2} < \sigma < 0$ and hence show the result for $0 < |\sigma| < \frac{1}{2}$.

We know that $\omega_0 \neq 0$, because $\zeta(s)$ has no zeros on the real axis between 0 and 1, when $s = \frac{1}{2} + \sigma + i\omega$ is real, $\omega = 0$ and $0 < |\sigma| < \frac{1}{2}$. [3] This is shown in detail in first two paragraphs in Appendix C.1.

2.1. New function $g(t)$

Let us consider the function $E'_p(t) = E'_p(t, t_2) = e^{-\sigma t_2} E_p(t - t_2) - e^{\sigma t_2} E_p(t + t_2) = (E_0(t - t_2) - E_0(t + t_2))e^{-\sigma t} = E'_0(t)e^{-\sigma t}$, where t_2 is finite and real, and $E'_0(t) = E'_0(t, t_2) = E_0(t - t_2) - E_0(t + t_2)$. Its Fourier transform is given by $E'_{p\omega}(\omega) = E'_{p\omega}(\omega, t_2) = E_{p\omega}(\omega)(e^{-\sigma t_2} e^{-i\omega t_2} - e^{\sigma t_2} e^{i\omega t_2})$ which has a zero at the **same** $\omega = \omega_0$.

Let us consider the function $f(t) = f(t, t_2, t_0) = e^{-2\sigma t_0} f_1(t) + e^{2\sigma t_0} f_2(t)$ where $f_1(t) = f_1(t, t_2, t_0) = e^{\sigma t_0} E'_p(t + t_0)$ and $f_2(t) = f_2(t, t_2, t_0) = e^{-\sigma t_0} E'_p(t - t_0)$ where t_0 is finite and real and we can see that the Fourier Transform of this function $F(\omega) = F(\omega, t_2, t_0) = E'_{p\omega}(\omega)(e^{-\sigma t_0} e^{i\omega t_0} + e^{\sigma t_0} e^{-i\omega t_0})$ also has a zero at the **same** $\omega = \omega_0$.

Let us consider a new function $g(t) = g(t, t_2, t_0) = g_-(t)u(-t) + g_+(t)u(t)$ where $g(t)$ is a real function of variable t and $u(t)$ is Heaviside unit step function and $g_-(t) = f(t)e^{-\sigma t}$ and $g_+(t) = f(t)e^{\sigma t}$. We can see that $g(t)h(t) = f(t)$ where $h(t) = [e^{\sigma t}u(-t) + e^{-\sigma t}u(t)]$.

We **note** that we use the **shorthand** notation for the functions $f(t), g(t), f_1(t), f_2(t), F(\omega)$ and $G(\omega)$ which are also functions of variables t_2, t_0 . Similarly we use the shorthand notation for the functions $E'_p(t), E'_0(t)$ and $E'_{p\omega}(\omega)$ which are also functions of variable t_2 .

We can show that $E_p(t), E'_p(t), h(t), g(t)$ are real absolutely integrable functions and go to zero as $t \rightarrow \pm\infty$. Hence their respective Fourier transforms given by $E_{p\omega}(\omega), E'_{p\omega}(\omega), H(\omega), G(\omega)$ are finite for $|\omega| \leq \infty$ and go to zero as $|\omega| \rightarrow \infty$, as per Riemann Lebesgue Lemma (link). This is shown in detail in Appendix C.1.

If we take the Fourier transform of the equation $g(t)h(t) = f(t)$ where $h(t) = [e^{\sigma t}u(-t) + e^{-\sigma t}u(t)]$, we get $\frac{1}{2\pi}[G(\omega) * H(\omega)] = F(\omega) = E'_{p\omega}(\omega)(e^{-\sigma t_0} e^{i\omega t_0} + e^{\sigma t_0} e^{-i\omega t_0}) = F_R(\omega) + iF_I(\omega)$ as per convolution theorem (link), where $*$ denotes convolution operation given by $F(\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} G(\omega')H(\omega - \omega')d\omega'$ and $H(\omega) = H_R(\omega) = [\frac{1}{\sigma - i\omega} + \frac{1}{\sigma + i\omega}] = \frac{2\sigma}{(\sigma^2 + \omega^2)}$ is real and is the Fourier transform of the function $h(t)$ and $G(\omega) = G_R(\omega) + iG_I(\omega)$ is the Fourier transform of the function $g(t)$. This is shown in detail in Appendix B.1.

For **every value** of t_0 , we require the Fourier transform of the function $f(t)$ given by $F(\omega)$ to have a zero at $\omega = \omega_0$. This implies that the Fourier transform of the **even** function $g(t)$ given by $G(\omega) = G_R(\omega)$ must have **at least one real zero** at $\omega = \omega_z(t_0)$ for **every value** of t_0 . Because $H(\omega) = \frac{2\sigma}{(\sigma^2 + \omega^2)}$ does not have real zeros, if $G_R(\omega)$ does not have real zeros, then $F_R(\omega) = G_R(\omega) * H_R(\omega)$ obtained by the convolution of $H_R(\omega)$ and $G_R(\omega)$, cannot possibly have real zeros, which goes against **Statement 1**.

We can write $g(t) = g_{\text{even}}(t) + g_{\text{odd}}(t)$ where $g_{\text{even}}(t)$ is an even function and $g_{\text{odd}}(t)$ is an odd function of variable t . If Statement 1 is true, then the **real** part of the Fourier transform of the **even function** $g_{\text{even}}(t) = \frac{1}{2}[g(t) + g(-t)]$ given by $G_R(\omega)$ must have **at least one zero** at $\omega = \omega_z(t_2, t_0) \neq 0$ where $\omega_z(t_0)$ is real and finite and can be different from ω_0 in general. We call this **Statement 2**.

Because $H(\omega) = \frac{2\sigma}{(\sigma^2 + \omega^2)}$ is real and does not have zeros for any finite value of ω , **if** $G_R(\omega)$ does not have at least one zero for some $\omega = \omega_z(t_2, t_0) \neq 0$, **then** the **real part** of $F(\omega)$ given by $F_R(\omega) = \frac{1}{2\pi}[G_R(\omega) * H(\omega)]$, obtained by the convolution of $H(\omega)$ and $G_R(\omega)$, **cannot** possibly have zeros for any non-zero finite value of ω , which goes against **Statement 1**. This is shown in detail in Lemma 1.

Lemma 1: If Riemann's Xi function $\xi(\frac{1}{2} + \sigma + i\omega) = E_{p\omega}(\omega)$ has a zero at $\omega = \omega_0 \neq 0$ where ω_0 is real and finite, then the **real** part of the Fourier transform of the **even function** $g_{\text{even}}(t) = \frac{1}{2}[g(t) + g(-t)]$ given by $G_R(\omega)$ must have **at least one zero** at $\omega = \omega_z(t_2, t_0) \neq 0$ for **every value** of t_0 , where $\omega_z(t_0)$ is real and finite, where $g(t)h(t) = f(t) = e^{-2\sigma t_0} f_1(t) + e^{2\sigma t_0} f_2(t)$ where $f_1(t) = e^{\sigma t_0} E'_p(t + t_0)$ and $f_2(t) = e^{-\sigma t_0} E'_p(t - t_0)$, $E'_p(t) = e^{-\sigma t_2} E_p(t - t_2) - e^{\sigma t_2} E_p(t + t_2)$, and $h(t) = e^{\sigma t}u(-t) + e^{-\sigma t}u(t)$ and $0 < \sigma < \frac{1}{2}$.

Proof: If $E_{p\omega}(\omega)$ has a zero at finite $\omega = \omega_0 \neq 0$ to satisfy Statement 1, then $F(\omega) = E'_{p\omega}(\omega)(e^{-\sigma t_0} e^{i\omega t_0} + e^{\sigma t_0} e^{-i\omega t_0}) = E_{p\omega}(\omega)(e^{-\sigma t_2} e^{-i\omega t_2} - e^{\sigma t_2} e^{i\omega t_2})(e^{-\sigma t_0} e^{i\omega t_0} + e^{\sigma t_0} e^{-i\omega t_0})$ also has a zero at $\omega = \omega_0$ and its real part given

by $F_R(\omega)$ also has a zero at the same location $\omega = \omega_0 \neq 0$.

Let us consider the case where $G_R(\omega)$ **does not** have at least one zero for finite $\omega = \omega_z(t_0) \neq 0$ and show that $F_R(\omega)$ does not have at least one zero at finite $\omega \neq 0$ for this case, which **contradicts** Statement 1. Given that $H(\omega)$ is real, we can write the convolution theorem only for the real parts as follows.

$$F_R(\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} G_R(\omega') H(\omega - \omega') d\omega' \quad (6)$$

We can show that the above integral converges for all $|\omega| \leq \infty$, given that $G(\omega)$ and $H(\omega)$ have fall-off rate of $\frac{1}{\omega^2}$ as $|\omega| \rightarrow \infty$ because the first derivatives of $g(t)$ and $h(t)$ are discontinuous at $t = 0$. (Appendix C.2)

We substitute $H(\omega) = \frac{2\sigma}{(\sigma^2 + \omega^2)}$ in Eq. 6 and we get

$$F_R(\omega) = \frac{\sigma}{\pi} \int_{-\infty}^{\infty} G_R(\omega') \frac{1}{(\sigma^2 + (\omega - \omega')^2)} d\omega' \quad (7)$$

We can split the integral in Eq. 7 as follows.

$$F_R(\omega) = \frac{\sigma}{\pi} \left[\int_{-\infty}^0 G_R(\omega') \frac{1}{(\sigma^2 + (\omega - \omega')^2)} d\omega' + \int_0^{\infty} G_R(\omega') \frac{1}{(\sigma^2 + (\omega - \omega')^2)} d\omega' \right] \quad (8)$$

We see that $G_R(-\omega) = G_R(\omega)$ because $g(t)$ is a real function (Appendix B.2). We can substitute $\omega' = -\omega''$ in the first integral in Eq. 8 and substituting $\omega'' = \omega'$ in the result, we can write as follows.

$$F_R(\omega) = \frac{\sigma}{\pi} \int_0^{\infty} G_R(\omega') \left[\frac{1}{(\sigma^2 + (\omega - \omega')^2)} + \frac{1}{(\sigma^2 + (\omega + \omega')^2)} \right] d\omega' \quad (9)$$

In Appendix C.1 last paragraph, it is shown that $G(\omega)$ is finite for $|\omega| \leq \infty$ and goes to zero as $|\omega| \rightarrow \infty$. We can see that for $\omega' = 0$ and $\omega' = \infty$, the integrand in Eq. 9 is zero. For finite $\omega > 0$, and $0 < \omega' < \infty$, we can see that the term $\frac{1}{(\sigma^2 + (\omega - \omega')^2)} + \frac{1}{(\sigma^2 + (\omega + \omega')^2)} > 0$.

• **Case 1:** $G_R(\omega') > 0$ for all finite $\omega' > 0$

We see that $F_R(\omega) > 0$ for all finite $\omega > 0$. We see that $F_R(-\omega) = F_R(\omega)$ because $f(t)$ is a real function (Appendix B.2). Hence $F_R(\omega) > 0$ for all finite $\omega < 0$.

This **contradicts** Statement 1 which requires $F_R(\omega)$ to have at least one zero at finite $\omega \neq 0$ because we showed that $\omega_0 \neq 0$ in **Section 2** paragraph 5. Therefore $G_R(\omega')$ must have **at least one zero** at $\omega' = \omega_z(t_2, t_0) \neq 0$, where $\omega_z(t_2, t_0)$ is real and finite.

• **Case 2:** $G_R(\omega') < 0$ for all finite $\omega' > 0$

We see that $F_R(\omega) < 0$ for all finite $\omega > 0$. We see that $F_R(-\omega) = F_R(\omega)$ because $f(t)$ is a real function (Appendix B.2). Hence $F_R(\omega) < 0$ for all finite $\omega < 0$.

This **contradicts** Statement 1 which requires $F_R(\omega)$ to have at least one zero at finite $\omega \neq 0$. Therefore $G_R(\omega')$ must have **at least one zero** at $\omega' = \omega_z(t_2, t_0) \neq 0$, where $\omega_z(t_2, t_0)$ is real and finite.

We have shown that, $G_R(\omega)$ must have **at least one zero** at finite $\omega = \omega_z(t_2, t_0) \neq 0$ to satisfy **Statement 1**. We call this **Statement 2**.

2.2. On the zeros of a related function $G(\omega)$

We can compute the fourier transform of the function $g_{even}(t) = \frac{1}{2}[g(t) + g(-t)]$ given by $G_R(\omega)$. We require $G_R(\omega) = 0$ for $\omega = \omega_z(t_2, t_0)$ for **every value** of t_0 , for a given value of t_2 , to satisfy **Statement 1**. In general, $\omega_z(t_2, t_0) \neq \omega_0$.

First we compute the Fourier transform of the function $g_1(t)$ given by $G_1(\omega) = G_{1R}(\omega) + iG_{1I}(\omega)$. We use $g_1(t) = f_1(t)e^{-\sigma t}u(-t) + f_1(t)e^{\sigma t}u(t) = e^{\sigma t_0}E'_p(t+t_0)e^{-\sigma t}u(-t) + e^{\sigma t_0}E'_p(t+t_0)e^{\sigma t}u(t)$.

We **note** that we use the **shorthand** notation for the functions $f(t), g(t), f_1(t), f_2(t), g_1(t), G(\omega)$ and $G_1(\omega)$ which are also functions of variables t_2, t_0 . Similarly we use the shorthand notation for the functions $E'_p(t), E'_0(t)$ and $E'_{0n}(t)$ which are also functions of variable t_2 .

$$\begin{aligned} G_1(\omega) &= \int_{-\infty}^{\infty} g_1(t)e^{-i\omega t}dt = \int_{-\infty}^0 g_1(t)e^{-i\omega t}dt + \int_0^{\infty} g_1(t)e^{-i\omega t}dt \\ G_1(\omega) &= \int_{-\infty}^0 e^{\sigma t_0}E'_p(t+t_0)e^{-\sigma t}e^{-i\omega t}dt + \int_0^{\infty} e^{\sigma t_0}E'_p(t+t_0)e^{\sigma t}e^{-i\omega t}dt \end{aligned}$$

(10)

We use $E'_p(t) = E'_0(t)e^{-\sigma t}$ where $E'_0(t) = E_0(t-t_2) - E_0(t+t_2)$ and $E'_p(t+t_0) = E'_0(t+t_0)e^{-\sigma t}e^{-\sigma t_0}$. Substituting $t = -t$ in the second integral in Eq. 10, we have

$$\begin{aligned} G_1(\omega) &= \int_{-\infty}^0 E'_0(t+t_0)e^{-2\sigma t}e^{-i\omega t}dt + \int_0^{\infty} E'_0(t+t_0)e^{-i\omega t}dt \\ G_1(\omega) &= \int_{-\infty}^0 E'_0(t+t_0)e^{-2\sigma t}e^{-i\omega t}dt + \int_{-\infty}^0 E'_0(-t+t_0)e^{i\omega t}dt \end{aligned}$$

(11)

We define $E'_{0n}(t) = E'_0(-t)$ and get $E'_0(-t+t_0) = E'_{0n}(t-t_0)$ and write Eq. 11 as follows.

$$G_1(\omega) = \int_{-\infty}^0 E'_0(t+t_0)e^{-2\sigma t}e^{-i\omega t}dt + \int_{-\infty}^0 E'_{0n}(t-t_0)e^{i\omega t}dt = G_R(\omega) + iG_I(\omega)$$

(12)

The above equations can be expanded as follows using the identity $e^{i\omega t} = \cos(\omega t) + i\sin(\omega t)$. Comparing the **real parts** of $G(\omega)$, we have

$$G_{1R}(\omega) = G_{1R}(\omega, t_0) = \int_{-\infty}^0 E'_0(t+t_0)e^{-2\sigma t}\cos(\omega t)dt + \int_{-\infty}^0 E'_{0n}(t-t_0)\cos(\omega t)dt$$

(13)

2.3. Zero crossing function $\omega_z(t_2, t_0)$ is an even function of variable t_0

Now we consider the function $f(t) = e^{-2\sigma t_0}f_1(t) + e^{2\sigma t_0}f_2(t) = e^{-\sigma t_0}E'_p(t+t_0) + e^{\sigma t_0}E'_p(t-t_0)$ where $f_1(t) = e^{\sigma t_0}E'_p(t+t_0)$ and $f_2(t) = f_1(t, -t_0) = e^{-\sigma t_0}E'_p(t-t_0)$ and $g(t)h(t) = f(t)$ where $g(t) = f(t)e^{-\sigma t}u(-t) + f(t)e^{\sigma t}u(t)$ and $h(t) = [e^{\sigma t}u(-t) + e^{-\sigma t}u(t)]$ and compute the Fourier transform of the function $g(t)$ and compute its real part using the procedure in above section, similar to Eq. 13 and we can write as follows. We substitute $t = \tau$.

$$\begin{aligned}
G_R(\omega, t_0) &= e^{-2\sigma t_0} G_{1R}(\omega, t_0) + e^{2\sigma t_0} G_{1R}(\omega, -t_0) \\
G_{1R}(\omega, t_0) &= \int_{-\infty}^0 [E'_0(\tau + t_0)e^{-2\sigma\tau} + E'_{0n}(\tau - t_0)] \cos(\omega\tau) d\tau \\
G_R(\omega, t_0) &= e^{-2\sigma t_0} \int_{-\infty}^0 [E'_0(\tau + t_0)e^{-2\sigma\tau} + E'_{0n}(\tau - t_0)] \cos(\omega\tau) d\tau \\
&\quad + e^{2\sigma t_0} \int_{-\infty}^0 [E'_0(\tau - t_0)e^{-2\sigma\tau} + E'_{0n}(\tau + t_0)] \cos(\omega\tau) d\tau
\end{aligned} \tag{14}$$

We require $G_R(\omega, t_0) = 0$ for $\omega = \omega_z(t_2, t_0)$ for every value of t_0 , for **every given fixed value** of t_2 , to satisfy **Statement 1**. In general $\omega_z(t_2, t_0) \neq \omega_0$. Hence we can see that $P(t_2, t_0) = G_R(\omega_z(t_2, t_0), t_0) = 0$ and we can rearrange the terms as follows.

$$\begin{aligned}
P(t_2, t_0) &= \int_{-\infty}^0 [e^{-2\sigma t_0} E'_0(\tau + t_0)e^{-2\sigma\tau} + e^{2\sigma t_0} E'_{0n}(\tau + t_0)] \cos(\omega_z(t_2, t_0)\tau) d\tau \\
&\quad + \int_{-\infty}^0 [e^{2\sigma t_0} E'_0(\tau - t_0)e^{-2\sigma\tau} + e^{-2\sigma t_0} E'_{0n}(\tau - t_0)] \cos(\omega_z(t_2, t_0)\tau) d\tau = 0
\end{aligned} \tag{15}$$

We can write as follows, where $P_{odd}(t_2, t_0)$ is an **odd** function of variable t_0 .

$$\begin{aligned}
P(t_2, t_0) &= P_{odd}(t_2, t_0) + P_{odd}(t_2, -t_0) = 0 \\
P_{odd}(t_2, t_0) &= \int_{-\infty}^0 [e^{-2\sigma t_0} E'_0(\tau + t_0)e^{-2\sigma\tau} + e^{2\sigma t_0} E'_{0n}(\tau + t_0)] \cos(\omega_z(t_2, t_0)\tau) d\tau
\end{aligned} \tag{16}$$

We see that $f(t, t_0) = e^{-\sigma t_0} E'_p(t + t_0) + e^{\sigma t_0} E'_p(t - t_0) = f(t, -t_0)$ is **unchanged** by the substitution $t_0 = -t_0$ and hence $\omega_z(t_2, t_0)$ is an **even** function of variable t_0 , for **every fixed value** of t_2 .

3. Final Step

We expand $P_{odd}(t_2, t_0)$ in Eq. 16 as follows, using the substitution $\tau + t_0 = \tau'$ and substituting back $\tau' = \tau$. We use $E'_{0n}(\tau) = E'_0(-\tau)$ and $E'_0(\tau) = E_0(\tau - t_2) - E_0(\tau + t_2)$.

We **note** that we use the **shorthand** notation for the functions $E'_0(t)$ and $E'_{0n}(t)$ which are also functions of variable t_2 .

$$\begin{aligned}
P_{odd}(t_2, t_0) &= [\cos(\omega_z(t_2, t_0)t_0) \int_{-\infty}^{t_0} E'_0(\tau)e^{-2\sigma\tau} \cos(\omega_z(t_2, t_0)\tau) d\tau \\
&\quad + \sin(\omega_z(t_2, t_0)t_0) \int_{-\infty}^{t_0} E'_0(\tau)e^{-2\sigma\tau} \sin(\omega_z(t_2, t_0)\tau) d\tau] \\
&\quad + e^{2\sigma t_0} [\cos(\omega_z(t_2, t_0)t_0) \int_{-\infty}^{t_0} E'_{0n}(\tau) \cos(\omega_z(t_2, t_0)\tau) d\tau + \sin(\omega_z(t_2, t_0)t_0) \int_{-\infty}^{t_0} E'_{0n}(\tau) \sin(\omega_z(t_2, t_0)\tau) d\tau]
\end{aligned} \tag{17}$$

In Section 2.1, $\omega_z(t_2, t_0)$ is shown to be **finite** for all $|t_0| < \infty$, for a given value of t_2 . This means there are **no** Dirac delta functions present in $\omega_z(t_2, t_0)$.

In Section 5, it is shown that $\omega_z(t_2, t_0)$ is a **continuous** function of variable t_0 for all $|t_0| < \infty$, for **every given fixed value** of t_2 .

In Section 4, it is shown that $E_0(t)$ is **strictly decreasing** for $t > 0$.

Given $\omega_z(t_2, t_0)$ is a continuous function of both t_0 and t_2 , we can find a suitable value of $t_0 = t_{0c}$ and $t_2 = t_{2c} = 2t_{0c}$ such that $\omega_z(t_{2c}, t_{0c})t_{0c} = \frac{\pi}{2}$. Given that $\omega_z(t_2, t_0)$ is a continuous function of t_0 , for every value of t_2 , and t_0 is a continuous function, we see that the **product** of two continuous functions $\omega_z(t_2, t_0)t_0$ is a **continuous** function as well. Given that $0 < \omega_z(t_2, t_0) < \infty$, we see that $\omega_z(t_2, t_0)t_0$ will **certainly pass through** π , as t_0 is increased from zero to ∞ .

We use $P_{odd}(t_2, t_0) + P_{odd}(t_2, -t_0) = 0$ as follows. We set $t_0 = t_{0c}$ and $t_2 = t_{2c} = 2t_{0c}$ such that $\omega_z(t_{2c}, t_{0c})t_{0c} = \frac{\pi}{2}$ in Eq. 17 as follows. We use the fact that $\omega_z(t_{2c}, -t_{0c}) = \omega_z(t_{2c}, t_{0c})$ shown in Section 2.3.

$$\begin{aligned} & \int_{-\infty}^{t_{0c}} E'_0(\tau) e^{-2\sigma\tau} \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau + e^{2\sigma t_{0c}} \int_{-\infty}^{t_{0c}} E'_{0n}(\tau) \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau \\ & - \int_{-\infty}^{-t_{0c}} E'_0(\tau) e^{-2\sigma\tau} \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau - e^{-2\sigma t_{0c}} \int_{-\infty}^{-t_{0c}} E'_{0n}(\tau) \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau = 0 \end{aligned} \quad (18)$$

We split the integral in the left hand side of Eq. 18 and write as follows.

$$\begin{aligned} & \left[\int_{-\infty}^{-t_{0c}} E'_0(\tau) e^{-2\sigma\tau} \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau + \int_{-t_{0c}}^{t_{0c}} E'_0(\tau) e^{-2\sigma\tau} \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau \right] \\ & + e^{2\sigma t_{0c}} \left[\int_{-\infty}^{-t_{0c}} E'_{0n}(\tau) \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau + \int_{-t_{0c}}^{t_{0c}} E'_{0n}(\tau) \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau \right] \\ & - \int_{-\infty}^{-t_{0c}} E'_0(\tau) e^{-2\sigma\tau} \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau - e^{-2\sigma t_{0c}} \int_{-\infty}^{-t_{0c}} E'_{0n}(\tau) \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau = 0 \end{aligned} \quad (19)$$

We combine the terms with common integrals and cancel common terms in Eq. 19 as follows.

$$\begin{aligned} & \int_{-t_{0c}}^{t_{0c}} E'_0(\tau) e^{-2\sigma\tau} \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau + e^{2\sigma t_{0c}} \int_{-t_{0c}}^{t_{0c}} E'_{0n}(\tau) \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau \\ & = -2 \sinh(2\sigma t_{0c}) \int_{-\infty}^{-t_{0c}} E'_{0n}(\tau) \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau \end{aligned} \quad (20)$$

We can rearrange the terms in Eq. 20 as follows.

$$\begin{aligned} & \int_{-t_{0c}}^{t_{0c}} [E'_0(\tau) e^{-2\sigma\tau} + E'_{0n}(\tau) e^{2\sigma t_{0c}}] \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau \\ & = -2 \sinh(2\sigma t_{0c}) \int_{-\infty}^{-t_{0c}} E'_{0n}(\tau) \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau \end{aligned} \quad (21)$$

274 We denote the right hand side of Eq. 21 as RHS . We can split the integral in Eq. 21 using $\int_{-t_{0c}}^{t_{0c}} = \int_{-t_{0c}}^0 + \int_0^{t_{0c}}$ as
 275 follows.

$$\begin{aligned} & \int_{-t_{0c}}^0 [E'_0(\tau)e^{-2\sigma\tau} + E'_{0n}(\tau)e^{2\sigma t_{0c}}] \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau \\ & + \int_0^{t_{0c}} [E'_0(\tau)e^{-2\sigma\tau} + E'_{0n}(\tau)e^{2\sigma t_{0c}}] \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau = RHS \end{aligned} \quad (22)$$

277 We substitute $\tau = -\tau$ in the first integral in Eq. 22 as follows. We use $E'_0(-\tau) = E'_{0n}(\tau)$ and $E'_{0n}(-\tau) = E'_0(\tau)$.

$$\begin{aligned} & \int_{t_{0c}}^0 [E'_{0n}(\tau)e^{2\sigma\tau} + E'_0(\tau)e^{2\sigma t_{0c}}] \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau \\ & + \int_0^{t_{0c}} [E'_0(\tau)e^{-2\sigma\tau} + E'_{0n}(\tau)e^{2\sigma t_{0c}}] \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau = RHS \end{aligned} \quad (23)$$

279 Given that $\int_{t_{0c}}^0 = -\int_0^{t_{0c}}$, we can simplify as follows.

$$\int_0^{t_{0c}} [E'_0(\tau)(e^{-2\sigma\tau} - e^{2\sigma t_{0c}}) + E'_{0n}(\tau)(-e^{2\sigma\tau} + e^{2\sigma t_{0c}})] \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau = RHS \quad (24)$$

281 We substitute $\tau = -\tau$ in the right hand side of Eq. 21 as follows. We use $E'_{0n}(-\tau) = E'_0(\tau)$.

$$RHS = 2 \sinh(2\sigma t_{0c}) \int_{t_{0c}}^{\infty} E'_0(\tau) \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau \quad (25)$$

283 We split the integral on the right hand side in Eq. 25 as follows.

$$RHS = 2 \sinh(2\sigma t_{0c}) \left[\int_0^{\infty} E'_0(\tau) \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau - \int_0^{t_{0c}} E'_0(\tau) \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau \right] \quad (26)$$

285 We consolidate the integrals with the term $\int_0^{t_{0c}} E'_0(\tau)$ in Eq. 24 and Eq. 26 as follows. We use $2 \sinh(2\sigma t_{0c}) =$
 286 $e^{2\sigma t_{0c}} - e^{-2\sigma t_{0c}}$.

$$\begin{aligned} & \int_0^{t_{0c}} [E'_0(\tau)(e^{-2\sigma\tau} - e^{2\sigma t_{0c}} + e^{2\sigma t_{0c}} - e^{-2\sigma t_{0c}}) + E'_{0n}(\tau)(-e^{2\sigma\tau} + e^{2\sigma t_{0c}})] \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau \\ & = 2 \sinh(2\sigma t_{0c}) \int_0^{\infty} E'_0(\tau) \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau \end{aligned} \quad (27)$$

288 We cancel common terms in Eq. 27 as follows.

$$\begin{aligned} & \int_0^{t_{0c}} [E'_0(\tau)(e^{-2\sigma\tau} - e^{-2\sigma t_{0c}}) + E'_{0n}(\tau)(-e^{2\sigma\tau} + e^{2\sigma t_{0c}})] \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau \\ & = 2 \sinh(2\sigma t_{0c}) \int_0^{\infty} E'_0(\tau) \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau \end{aligned}$$

We substitute $E'_0(\tau) = E_0(\tau - t_{2c}) - E_0(\tau + t_{2c})$ and $E'_{0n}(\tau) = E'_0(-\tau) = E_0(-\tau - t_{2c}) - E_0(-\tau + t_{2c})$. We see that $E_0(-\tau - t_{2c}) = E_0(\tau + t_{2c})$ and $E_0(-\tau + t_{2c}) = E_0(\tau - t_{2c})$ given that $E_0(\tau) = E_0(-\tau)$. Hence we see that $E'_{0n}(\tau) = E_0(\tau + t_{2c}) - E_0(\tau - t_{2c}) = -E'_0(\tau)$. We can write Eq. 28 as follows.

$$\begin{aligned} & \int_0^{t_{0c}} (E_0(\tau - t_{2c}) - E_0(\tau + t_{2c}))(e^{-2\sigma\tau} - e^{-2\sigma t_{0c}} + e^{2\sigma\tau} - e^{2\sigma t_{0c}}) \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau \\ &= 2 \sinh(2\sigma t_{0c}) \int_0^\infty (E_0(\tau - t_{2c}) - E_0(\tau + t_{2c})) \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau \end{aligned}$$

We substitute $2 \cosh(2\sigma\tau) = e^{2\sigma\tau} + e^{-2\sigma\tau}$ and $2 \cosh(2\sigma t_{0c}) = e^{2\sigma t_{0c}} + e^{-2\sigma t_{0c}}$ and cancel the common factor of 2 in Eq. 29 as follows.

$$\begin{aligned} & \int_0^{t_{0c}} (E_0(\tau - t_{2c}) - E_0(\tau + t_{2c}))(\cosh(2\sigma\tau) - \cosh(2\sigma t_{0c})) \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau \\ &= \sinh(2\sigma t_{0c}) \int_0^\infty (E_0(\tau - t_{2c}) - E_0(\tau + t_{2c})) \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau \end{aligned}$$

Next Step:

We substitute $\tau + t_{2c} = \tau'$ in the right hand side of Eq. 30 and then substitute $\tau' = \tau$. Similarly we substitute $\tau - t_{2c} = \tau'$ as follows.

$$\begin{aligned} RHS = & \sinh(2\sigma t_{0c}) [\cos(\omega_z(t_{2c}, t_{0c})t_{2c}) \int_{-t_{2c}}^\infty E_0(\tau) \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau \\ & + \sin(\omega_z(t_{2c}, t_{0c})t_{2c}) \int_{-t_{2c}}^\infty E_0(\tau) \cos(\omega_z(t_{2c}, t_{0c})\tau) d\tau \\ & - \cos(\omega_z(t_{2c}, t_{0c})t_{2c}) \int_{t_{2c}}^\infty E_0(\tau) \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau + \sin(\omega_z(t_{2c}, t_{0c})t_{2c}) \int_{t_{2c}}^\infty E_0(\tau) \cos(\omega_z(t_{2c}, t_{0c})\tau) d\tau] \end{aligned}$$

In Eq. 31, given that $\omega_z(t_{2c}, t_{0c})t_{0c} = \frac{\pi}{2}$ and $t_{2c} = 2t_{0c}$ and hence $\omega_z(t_{2c}, t_{0c})t_{2c} = 2\frac{\pi}{2} = \pi$ and $\sin(\omega_z(t_{2c}, t_{0c})t_{2c}) = 0$ and $\cos(\omega_z(t_{2c}, t_{0c})t_{2c}) = -1$. Hence we cancel common terms and write Eq. 31 and Eq. 30 as follows.

$$\begin{aligned} & \int_0^{t_{0c}} (E_0(\tau - t_{2c}) - E_0(\tau + t_{2c}))(\cosh(2\sigma\tau) - \cosh(2\sigma t_{0c})) \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau \\ &= -\sinh(2\sigma t_{0c}) [\int_{-t_{2c}}^\infty E_0(\tau) \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau - \int_{t_{2c}}^\infty E_0(\tau) \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau] \end{aligned}$$

We use $\int_{-t_{2c}}^\infty E_0(\tau) \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau = \int_{-t_{2c}}^{t_{2c}} E_0(\tau) \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau + \int_{t_{2c}}^\infty E_0(\tau) \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau$ and cancel the common term $\int_{t_{2c}}^\infty E_0(\tau) \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau$ in Eq. 32 as follows. Given that $E_0(\tau) \sin(\omega_z(t_{2c}, t_{0c})\tau)$ is an **odd** function of variable τ , we get $\int_{-t_{2c}}^{t_{2c}} E_0(\tau) \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau = 0$.

$$\int_0^{t_{0c}} (E_0(\tau - t_{2c}) - E_0(\tau + t_{2c}))(\cosh(2\sigma\tau) - \cosh(2\sigma t_{0c})) \sin(\omega_z(t_{2c}, t_{0c})\tau) d\tau = 0$$

We can multiply Eq. 33 by a factor of -1 as follows.

$$\int_0^{t_{0c}} [E_0(\tau - t_{2c}) - E_0(\tau + t_{2c})](\cosh 2\sigma t_{0c} - \cosh(2\sigma\tau) \sin(\omega_z(t_{2c}, t_{0c})\tau)) d\tau = 0$$

(34)

In Eq. 34, given that $\omega_z(t_{2c}, t_{0c})t_{0c} = \frac{\pi}{2}$, as τ varies over the interval $[0, t_{0c}]$, $\omega_z(t_{2c}, t_{0c})\tau = \frac{\pi\tau}{2t_{0c}}$ varies from $[0, \frac{\pi}{2}]$ and hence the sinusoidal function varies over a **half cycle** and is > 0 , in the interval $0 < \tau < t_{0c}$, for $t_{0c} > 0$.

In Eq. 34, we see that in the interval $0 < \tau < t_{0c}$, the integral on the left hand side is > 0 for $t_{0c} > 0$, because each of the terms in the integrand are > 0 , in the interval $0 < \tau < t_{0c}$ as follows. Given that $E_0(t)$ is a **strictly decreasing** function for $t \geq \frac{1}{8}$, we see that $E_0(\tau - t_{2c}) - E_0(\tau + t_{2c})$ is > 0 (Section 4.7). The term $(\cosh(2\sigma t_{0c}) - \cosh(2\sigma\tau))$ is > 0 in the interval $0 < \tau < t_{0c}$ and the integrand is zero at $\tau = 0$ and $\tau = t_{0c}$ and hence the integral **cannot** equal zero, as required by the right hand side of Eq. 34. Hence this leads to a **contradiction** for $0 < \sigma < \frac{1}{2}$.

For $\sigma = 0$, both sides of Eq. 34 is zero and **does not** lead to a contradiction.

We have shown this result for $0 < \sigma < \frac{1}{2}$ and then use the property $\xi(\frac{1}{2} + \sigma + i\omega) = \xi(\frac{1}{2} - \sigma - i\omega)$ to show the result for $-\frac{1}{2} < \sigma < 0$. Hence we have produced a **contradiction** of **Statement 1** that the Fourier Transform of the function $E_p(t) = E_0(t)e^{-\sigma t}$ has a zero at $\omega = \omega_0$ for $0 < |\sigma| < \frac{1}{2}$.

Therefore, the assumption in **Statement 1** that Riemann's Xi Function given by $\xi(\frac{1}{2} + \sigma + i\omega) = E_{p\omega}(\omega)$ has a zero at $\omega = \omega_0$, where ω_0 is real and finite, leads to a **contradiction** for the region $0 < |\sigma| < \frac{1}{2}$ which corresponds to the critical strip excluding the critical line. This means $\zeta(s)$ does not have non-trivial zeros in the critical strip excluding the critical line and we have proved Riemann's Hypothesis.

4. Strictly decreasing $E_0(t)$ for $t > 0$

Let us consider $E_0(t) = 2 \sum_{n=1}^{\infty} [2\pi^2 n^4 e^{4t} - 3\pi n^2 e^{2t}] e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}}$ whose Fourier Transform is given by the entire function $E_{0\omega}(\omega) = \xi(\frac{1}{2} + i\omega)$. (link)

$$E_0(t) = \Phi(t) = 2 \sum_{n=1}^{\infty} [2\pi^2 n^4 e^{4t} - 3\pi n^2 e^{2t}] e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}} = \sum_{n=1}^{\infty} [4\pi^2 n^4 e^{4t} - 6\pi n^2 e^{2t}] e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}}$$

$$E_0(t) = \sum_{n=1}^{\infty} 2\pi n^2 e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}} [2\pi n^2 e^{4t} - 3e^{2t}]$$

(35)

We show that $X(t) = \frac{E_0(t)}{2\pi}$ is a **strictly decreasing** function for $t \geq 0$ as follows.

- In Section 4.1, it is shown that the second derivative of $X(t)$, given by $X_2(t) = \frac{d^2 X(t)}{dt^2} < 0$ for $t = 0$.

- In Section 4.2, it is shown that, as t increases from zero, $\frac{dX(t)}{dt}$ starts from zero and reaches a **negative minimum** value at $t = t_{min}$ and then starts increasing towards zero, for $t > t_{min}$. (example plot) Hence $\frac{dX(t)}{dt} < 0$ for $0 < t \leq t_{min}$.

- In Section 4.3, it is shown that $\frac{dX(t)}{dt} < 0$ for $t \geq t'_{min}$, where t'_{min} corresponds to the minima of $X_{11}(t)$ which is the partial term in $\frac{dX(t)}{dt}$ corresponding to $n = 1$.

- In Section 4.4, it is shown that $\frac{dX(t)}{dt} < 0$ for $t_{min} < t < t'_{min}$. Hence $\frac{dX(t)}{dt} < 0$ for all $t > 0$ and hence $X(t)$ is **strictly decreasing** for all $t > 0$ and $E_0(t) = 2\pi X(t)$ is **strictly decreasing** for all $t > 0$.

4.1. $\frac{d^2 X(t)}{dt^2} < 0$ for $t = 0$

We consider $X(t) = \frac{E_0(t)}{2\pi} = \sum_{n=1}^{\infty} n^2 e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}} [2\pi n^2 e^{4t} - 3e^{2t}]$ and take the first derivative of $X(t)$ as follows. We note that $E_0(t)$ is an analytic function for $|t| \leq \infty$ and is infinitely differentiable in that interval.

$$\frac{dX(t)}{dt} = \sum_{n=1}^{\infty} n^2 e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}} [8\pi n^2 e^{4t} - 6e^{2t} + (2\pi n^2 e^{4t} - 3e^{2t}) (\frac{1}{2} - 2\pi n^2 e^{2t})]$$

$$\frac{dX(t)}{dt} = \sum_{n=1}^{\infty} n^2 e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}} [8\pi n^2 e^{4t} - 6e^{2t} + (\pi n^2 e^{4t} - \frac{3}{2}e^{2t} - 4\pi^2 n^4 e^{6t} + 6\pi n^2 e^{4t})]$$

$$\frac{dX(t)}{dt} = \sum_{n=1}^{\infty} n^2 e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}} [-4\pi^2 n^4 e^{6t} + 15\pi n^2 e^{4t} - \frac{15}{2}e^{2t}]$$

$$\frac{dX(t)}{dt} = \sum_{n=1}^{\infty} n^2 e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}} e^{2t} [-4\pi^2 n^4 e^{4t} + 15\pi n^2 e^{2t} - \frac{15}{2}]$$

(36)

We take the second derivative of $X(t)$ as follows.

$$\frac{d^2 X(t)}{dt^2} = \sum_{n=1}^{\infty} n^2 e^{-\pi n^2 e^{2t}} e^{\frac{5t}{2}} [-16\pi^2 n^4 e^{4t} + 30\pi n^2 e^{2t} + (-4\pi^2 n^4 e^{4t} + 15\pi n^2 e^{2t} - \frac{15}{2})(\frac{5}{2} - 2\pi n^2 e^{2t})]$$

$$\begin{aligned} \frac{d^2 X(t)}{dt^2} &= \sum_{n=1}^{\infty} n^2 e^{-\pi n^2 e^{2t}} e^{\frac{5t}{2}} [-16\pi^2 n^4 e^{4t} + 30\pi n^2 e^{2t} \\ &\quad + (-10\pi^2 n^4 e^{4t} + \frac{75}{2}\pi n^2 e^{2t} - \frac{75}{4} + 8\pi^3 n^6 e^{6t} - 30\pi^2 n^4 e^{4t} + 15\pi n^2 e^{2t})] \\ \frac{d^2 X(t)}{dt^2} &= \sum_{n=1}^{\infty} n^2 e^{-\pi n^2 e^{2t}} e^{\frac{5t}{2}} Z(t), \quad Z(t) = 8\pi^3 n^6 e^{6t} - 56\pi^2 n^4 e^{4t} + \frac{165}{2}\pi n^2 e^{2t} - \frac{75}{4} \end{aligned}$$

(37)

• **Case 1:** $n = 1, t = 0$: The partial term for $n = 1$ and $t = 0$ in $\frac{d^2 X(t)}{dt^2}$ in Eq. 37, given by $A_1 < 0$ shown as follows.

$$A_1 = e^{-\pi}(8\pi^3 - 56\pi^2 + \frac{165}{2}\pi - \frac{75}{4}) < 0$$

(38)

At $n = 1, t = 0$, we see that $Z(t) = 8\pi^3 - 56\pi^2 + \frac{165}{2}\pi - \frac{75}{4} = \pi(8\pi^2 - 56\pi + \frac{165}{2}) - \frac{75}{4} \leq 3*(8*10 - 56*3 + 83) - 18 = -33 < 0$ (**Result K**), because $3 < \pi < 3.1429$ and $\pi^2 < 10$. We note that the term in the brackets is negative and hence we multiply it by the minimum value of π . Hence we see that the partial term for $n = 1$ and $t = 0$ in $\frac{d^2 X(t)}{dt^2}$ given by $A_1 < 0$.

• **Case 2:** $n > 1, t = 0$: the partial terms for $n > 1$ and $t = 0$ in $\frac{d^2 X(t)}{dt^2}$ in Eq. 37, given by $A_2 > 0$ and $A_2 < |A_1|$ shown as follows.

$$A_2 = \sum_{n=2}^{\infty} n^2 e^{-\pi n^2} (8\pi^3 n^6 - 56\pi^2 n^4 + \frac{165}{2}\pi n^2 - \frac{75}{4}) > 0$$

(39)

At $n = 2, t = 0$, we see that $Z(t) = 8\pi^3 n^6 - 56\pi^2 n^4 + \frac{165}{2}\pi n^2 - \frac{75}{4} = \pi n^2(8\pi^2 n^4 - 56\pi n^2 + \frac{165}{2}) - \frac{75}{4} > 3*2^2(8*3^2*2^4 - 56*4*2^2 + 82) - 19 = 4037 > 0$, because $3 < \pi < 4$. At $n > 2, t = 0$, we see that $Z(t) > 0$ due to the dominant term $8\pi^2 n^4$. Hence the partial terms for $n > 1$ and $t = 0$ in $\frac{d^2 X(t)}{dt^2}$ in Eq. 37, given by $A_2 > 0$.

• **We can show** that at $t = 0$, $A_1 \leq -0.4283$ and $A_2 < |A_1|$ and hence $\frac{d^2 X(t)}{dt^2} < 0$ in Eq. 37.

At $n = 1, t = 0$, we see that, given $3 < \pi < \frac{22}{7} = 3.1429$ (link) and $e^{-\pi} \geq e^{-0.1429}(e^{-\frac{1}{2}})^6$ and $e^{-0.1429} \geq 1 - 0.1429 = 0.8571$ and $e^{-\frac{1}{2}} \geq 1 - \frac{1}{2} = \frac{1}{2}$, we can write $e^{-\pi} \geq \frac{0.8571}{2^6}$ and $|A_1| = e^{-\pi}|Z(t)| \geq 33 * \frac{0.8571}{2^6} \geq 32 * \frac{0.8571}{64} = 0.4283$ (using Result K). Hence we see that the partial term for $n = 1, t = 0$ in $\frac{d^2 X(t)}{dt^2}$ in Eq. 37, given by $A_1 \leq -0.4283$ (**Result A**).

At $n > 1, t = 0$, we see that $A_2 = \sum_{n=2}^{\infty} n^2 e^{-\pi n^2} Z(0) \leq \sum_{n=2}^{\infty} n^2 e^{-\pi n^2} * (8\pi^3 n^6 + \frac{165}{2}\pi n^2)$ and we **want to show** $A_2 < |A_1|$ where $A_1 \leq -0.4283$. We use the fact that $(n^2)^4 e^{-\pi n^2} < e^{-(\pi-1)n^2}$ and $(n^2)^2 e^{-\pi n^2} < e^{-(\pi-1)n^2}$.

$$A_2 \leq \sum_{n=2}^{\infty} e^{-(\pi-1)n^2} (8\pi^3 + \frac{165}{2}\pi)$$

(40)

We use $3 < \pi < 3.1429$ and $\sum_{n=2}^{\infty} f(n) \leq \int_{u=2}^{\infty} f(u) du$.

$$A_2 \leq (8\pi^3 + \frac{165}{2}\pi) \int_{u=2}^{\infty} e^{-(\pi-1)u^2} du \leq (8\pi^3 + \frac{165}{2}\pi) \int_{u=2}^{\infty} e^{-2u^2} du$$

(41)

We substitute $\sqrt{2}u = t$ as follows. We use $8\pi^3 + \frac{165}{2}\pi = \pi(8\pi^2 + \frac{165}{2}) \leq 3.1429(8 * 10 + 83) \leq 514$. We use the complementary error function $erfc(z)$ and $\frac{\sqrt{\pi}}{2\sqrt{2}} \leq 1$ as follows. (link)

$$\begin{aligned} erfc(z) &= \frac{2}{\sqrt{\pi}} \int_z^{\infty} e^{-t^2} dt \\ A_2 &\leq \frac{513}{\sqrt{2}} \int_{u=2\sqrt{2}}^{\infty} e^{-t^2} dt \leq \frac{514}{\sqrt{2}} \left(\frac{\sqrt{\pi}}{2} * erfc(2\sqrt{2}) \right) \leq \frac{514}{\sqrt{2}} \frac{\sqrt{\pi}}{2} * erfc(2\sqrt{2}) \\ A_2 &\leq 514 * erfc(2\sqrt{2}) = 0.0326 \end{aligned}$$

(42)

We can see that $A_1 + A_2 < 0$ given that $A_1 \leq -0.4283$ (using Result A).

We use the fact that $erfc(z) \leq \frac{e^{-z^2}}{z\sqrt{\pi}}$ (link) . Hence $A_2 \leq 514 \frac{e^{-8}}{2\sqrt{2}\pi} \leq 514 \frac{e^{-8}}{4} \leq 129e^{-8}$. We can show that $A_1 + A_2 < 0$ as follows, using Eq. 38.

$$\begin{aligned} A_2 &\leq 129e^{-8}, \quad A_1 \leq -33e^{-\pi} \leq -33e^{-3}, \quad A_2 < |A_1| \\ A_1 + A_2 &\leq -33e^{-3} + 129e^{-8} = e^{-3}(-33 + 129e^{-5}) \\ A_1 + A_2 &\leq e^{-3}(-33 + 129 * 2^{-5}) = e^{-3}(-33 + \frac{129}{32}) < 0 \\ A_1 + A_2 &< 0 \\ (\frac{d^2X(t)}{dt^2})_{t=0} &= A_1 + A_2 < 0 \end{aligned}$$

(43)

We have shown that the second derivative of $X(t)$, given by $X_2(t) = \frac{d^2X(t)}{dt^2} < 0$ for $t = 0$. (**Result H**)

4.2. $\frac{dX(t)}{dt} < 0$ **for** $0 < t \leq t_{min}$

• In the sections below, **we will show** that, as t increases from zero, $\frac{dX(t)}{dt}$ starts from zero and reaches a **negative minimum** value at $t = t_{min}$ and then starts increasing towards zero, for $t > t_{min}$. Thus we will show that $X(t)$ is a **strictly decreasing** function of t . We will also show that $\frac{dX(t)}{dt}$ **does not** become positive for any $t > 0$. (example plot)

• We see that $\frac{d^2X(t)}{dt^2} < 0$ at $t = 0$ (using Result H) . Hence, as t increases from zero (point A), $\frac{dX(t)}{dt}$ reaches a **negative minimum** value at $t = t_{min}$ (point B) and then starts increasing towards zero, in the neighborhood of $t > t_{min}$. We see that $\frac{d^2X(t)}{dt^2}$ remains negative in $0 \leq t < t_{min}$ and then reaches a **zero** at $t = t_{min}$ and then becomes **positive** in the neighborhood of $t > t_{min}$. (example plot) Hence $\frac{dX(t)}{dt} < 0$ for $0 < t \leq t_{min}$. (**Result I**)

4.3. $\frac{dX(t)}{dt} < 0$ **for** $t \geq t'_{min}$

We note that $\frac{dX(t)}{dt} = X_{11}(t) + X_{12}(t)$ where $X_{11}(t)$ has partial term corresponding to $n = 1$ and $X_{12}(t)$ has partial terms corresponding to $n > 1$.

$$\begin{aligned}\frac{dX(t)}{dt} &= X_{11}(t) + X_{12}(t) = \sum_{n=1}^{\infty} n^2 e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}} e^{2t} [-4\pi^2 n^4 e^{4t} + 15\pi n^2 e^{2t} - \frac{15}{2}] \\ X_{11}(t) &= e^{-\pi e^{2t}} e^{\frac{t}{2}} e^{2t} [-4\pi^2 e^{4t} + 15\pi e^{2t} - \frac{15}{2}] \\ X_{12}(t) &= \sum_{n=2}^{\infty} n^2 e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}} e^{2t} [-4\pi^2 n^4 e^{4t} + 15\pi n^2 e^{2t} - \frac{15}{2}]\end{aligned}$$

(44)

In the interval $t > t_{min}$, $\frac{d^2 X(t)}{dt^2} = X_{21}(t) + X_{22}(t)$ reaches a **positive maximum** at $t = t_{max2}$ (point C) and then starts decreasing to zero, in the neighborhood $t > t_{max2}$. $X_{21}(t)$ has partial term corresponding to $n = 1$ and $X_{22}(t)$ has partial terms corresponding to $n > 1$. This means that $\frac{dX(t)}{dt}$ **increases** towards zero, in the neighborhood $t > t_{min}$. (example plot)

$$\begin{aligned}\frac{d^2 X(t)}{dt^2} &= X_{21}(t) + X_{22}(t) = \sum_{n=1}^{\infty} n^2 e^{-\pi n^2 e^{2t}} e^{\frac{5t}{2}} Z(t), \quad Z(t) = 8\pi^3 n^6 e^{6t} - 56\pi^2 n^4 e^{4t} + \frac{165}{2} \pi n^2 e^{2t} - \frac{75}{4} \\ X_{21}(t) &= e^{-\pi e^{2t}} e^{\frac{5t}{2}} (8\pi^3 e^{6t} - 56\pi^2 e^{4t} + \frac{165}{2} \pi e^{2t} - \frac{75}{4}) \\ X_{22}(t) &= \sum_{n=2}^{\infty} n^2 e^{-\pi n^2 e^{2t}} e^{\frac{5t}{2}} (8\pi^3 n^6 e^{6t} - 56\pi^2 n^4 e^{4t} + \frac{165}{2} \pi n^2 e^{2t} - \frac{75}{4})\end{aligned}$$

(45)

As t increases from zero (point A), $X_{11}(t)$ in Eq. 44 starts from a positive value and **decreases** to a negative value, passing through **zero** at $t = t_z$. Given that $\lim_{t \rightarrow \infty} X_{11}(t) = 0$, it reaches a **negative minimum** value at $t = t'_{min}$ (point B') and then starts **increasing** towards zero, for $t > t'_{min}$ (**Result L**). (example plot) We note that $X_{11}(t)$ **remains negative** for $t > t_z$ due to the dominant term $-4\pi^2 e^{4t}$ in Eq. 44.

This means that the derivative of $X_{11}(t)$ given by $X_{21}(t)$ in Eq. 45 is **negative** at $0 \leq t < t'_{min}$ and $X_{21}(t)$ reaches a **zero** at $t = t'_{min}$ (**Result N**) and then becomes **positive** for $t > t'_{min}$ (using Result L) and falls to zero as $t \rightarrow \infty$, given that $\lim_{t \rightarrow \infty} X_{21}(t) = 0$. We note that, when $X_{21}(t)$ in Eq. 45 becomes positive after crossing $t = t'_{min}$, it **remains positive** as $t \rightarrow \infty$ due to the dominant term $8\pi^3 e^{6t}$ in Eq. 45 (**Result B**).

Hence $X_{11}(t)$ is a **strictly decreasing** function for $t < t'_{min}$ and then it starts **increasing** towards zero as $t \rightarrow \infty$. We note that $X_{11}(t)$ **cannot** become positive at some $t > t'_{min}$, because if it did become positive again, then it would have to decrease to zero as $t \rightarrow \infty$, which would **require** $X_{21}(t)$ to become negative again, which is **not** the case as shown in Result B.

We will show that $\frac{d^2 X(t)}{dt^2} = X_{21}(t) + X_{22}(t) = 0$ **only for** $t = t_{min}$, and then it becomes positive and then starts decreasing towards zero, as $t \rightarrow \infty$.

In Section 4.5, it is shown that $X_{22}(t)$ is **strictly decreasing** for $t > 0$ (using Result E) and hence $X_{22}(t) > 0$ given that $\lim_{t \rightarrow \infty} X_{22}(t) = 0$. (**Result C**).

We note that $t_{min} < t'_{min}$ given that $X_{21}(t) > 0$ for $t > t'_{min}$ (using Result B) and $X_{22}(t) > 0$ for $t > t'_{min}$ (using Result C) and hence $\frac{d^2 X(t)}{dt^2} = X_{21}(t) + X_{22}(t) > 0$ for $t > t'_{min}$ (**Result D**).

We see that $X_{12}(t) < 0$ for all $t > 0$ in Eq. 44 due to the dominant term $-4\pi^2 n^4 e^{4t}$. Because $\frac{dX(t)}{dt} = X_{11}(t) + X_{12}(t)$ is negative at $t = t'_{min}$, we see that $\frac{dX(t)}{dt}$ starts **increasing** from a negative value towards zero. (using Result D) Hence $\frac{dX(t)}{dt} < 0$ for $t \geq t'_{min}$. (**Result G**)

We note that $\frac{dX(t)}{dt}$ **cannot** become positive at some $t > t'_{min}$, because if it did become positive again, then it would have to decrease to zero as $t \rightarrow \infty$, which would **require** $\frac{d^2X(t)}{dt^2}$ to become negative again, which is **not** the case as shown in Result D.

4.4. $\frac{dX(t)}{dt} < 0$ **for** $t_{min} < t < t'_{min}$

We can show that the second derivative $\frac{d^2X(t)}{dt^2} = X_{21}(t) + X_{22}(t)$ becomes zero **only once** at $t = t_{min}$, in the interval $0 < t < t'_{min}$. At $t = t_{min}$, we see that $X_{22}(t_{min}) = X_0 > 0$ and hence $X_{21}(t_{min}) = -X_0 < 0$. In the interval $t_{min} < t < t'_{min}$, we see that $X_{22}(t)$ is **strictly decreasing** and remains positive (using Result E in Section 4.5) and decreases from X_0 further towards zero at a **slower** rate, while $X_{21}(t)$ increases from the negative value $-X_0$ towards zero **faster**, to make $\frac{d^2X(t)}{dt^2} = X_{21}(t) + X_{22}(t) > 0$.

We see that $X_{21}(t)$ is **strictly increasing** for $t_{min} < t < t'_{min}$ (using Result F in Section 4.6) and **reaches zero** at $t = t'_{min}$ at a **faster** rate than $X_{22}(t)$. Hence $\frac{d^2X(t)}{dt^2} = X_{21}(t) + X_{22}(t) > 0$ in the interval $t_{min} < t < t'_{min}$ (**Result M**). Hence $\frac{dX(t)}{dt}$ is **increasing** from a negative value towards zero, in the interval $t_{min} < t < t'_{min}$.

We can **rule out** the Case A that $\frac{dX(t)}{dt} > 0$ somewhere in the interval $t > t_{min}$ and reaches a maximum at $t = t_{max}$, as follows. We see that $\frac{d^2X(t)}{dt^2} = X_{21}(t) + X_{22}(t)$ becomes zero **only once** at $t = t_{min}$, in the interval $0 < t < t'_{min}$ and is positive for $t_{min} < t < t'_{min}$ (using Result M) and **does not** become zero again at $t = t_{max}$, which is required for Case A. Hence $\frac{dX(t)}{dt} < 0$ in the interval $t_{min} < t < t'_{min}$. (**Result J**)

We have shown in earlier sections that $\frac{dX(t)}{dt} < 0$ in the interval $0 < t \leq t_{min}$ (using Result I) and $t \geq t'_{min}$ (using Result G). We see that $\frac{dX(t)}{dt} < 0$ in the interval $t_{min} < t < t'_{min}$ (using Result J) and hence $\frac{dX(t)}{dt} < 0$ for all $t > 0$ and hence $X(t)$ is **strictly decreasing** for all $t > 0$. Hence $E_0(t) = 2\pi X(t)$ is **strictly decreasing** for all $t > 0$.

4.5. **Second derivative given by X_{22} is a strictly decreasing function for $t > 0$**

We consider $X_{22}(t)$ as follows.

$$X_{22}(t) = \sum_{n=2}^{\infty} n^2 e^{-\pi n^2 e^{2t}} e^{\frac{5t}{2}} (8\pi^3 n^6 e^{6t} - 56\pi^2 n^4 e^{4t} + \frac{165}{2}\pi n^2 e^{2t} - \frac{75}{4}) \quad (46)$$

We see that $X_{22}(t)$ in Eq. 46, is **positive** for each $n = 2, 3, \dots$ and hence for **all** $t \geq 0$, as follows. We see that $X_{22}(t) = \sum_{n=2}^{\infty} n^2 e^{-\pi n^2 e^{2t}} e^{\frac{5t}{2}} (\pi^2 n^4 e^{4t} (8\pi n^2 e^{2t} - 56) + \frac{165}{2}\pi n^2 e^{2t} - \frac{75}{4})$. For $n = 2, t = 0$, we see that $X_{22}(t) = 2^2 e^{-\pi 2^2} (8\pi^3 2^6 - 56\pi^2 2^4 + \frac{165}{2}\pi 2^2 - \frac{75}{4}) = 4e^{-4\pi} (\pi^2 2^4 (32\pi - 56) + 165 * 2\pi - \frac{75}{4}) > 0$. For $n \geq 2$ and $t \geq 0$, $X_{22}(t) > 0$ due to the **dominant term** $8\pi^3 n^6 e^{6t}$.

We compute the **derivative** of $X_{22}(t)$ as follows.

$$\begin{aligned} X_{32}(t) &= \frac{dX_{22}(t)}{dt} = \sum_{n=2}^{\infty} n^2 e^{-\pi n^2 e^{2t}} e^{\frac{5t}{2}} [48\pi^3 n^6 e^{6t} - 56 * 4\pi^2 n^4 e^{4t} + 165\pi n^2 e^{2t} \\ &\quad + (8\pi^3 n^6 e^{6t} - 56\pi^2 n^4 e^{4t} + \frac{165}{2}\pi n^2 e^{2t} - \frac{75}{4})(\frac{5}{2} - 2\pi n^2 e^{2t})] \\ X_{32}(t) &= \sum_{n=2}^{\infty} n^2 e^{-\pi n^2 e^{2t}} e^{\frac{5t}{2}} [48\pi^3 n^6 e^{6t} - 56 * 4\pi^2 n^4 e^{4t} + 165\pi n^2 e^{2t} \\ &\quad + (20\pi^3 n^6 e^{6t} - 28 * 5\pi^2 n^4 e^{4t} + \frac{165 * 5}{4}\pi n^2 e^{2t} - \frac{75 * 5}{8} - 16\pi^4 n^8 e^{8t} + 56 * 2\pi^3 n^6 e^{6t} - 165\pi^2 n^4 e^{4t} + \frac{75}{2}\pi n^2 e^{2t})] \\ X_{32}(t) &= \sum_{n=2}^{\infty} n^2 e^{-\pi n^2 e^{2t}} e^{\frac{5t}{2}} Y'(t), \quad Y'(t) = -16\pi^4 n^8 e^{8t} + 180\pi^3 n^6 e^{6t} - 529\pi^2 n^4 e^{4t} + 408.75\pi n^2 e^{2t} - \frac{75 * 5}{8} \end{aligned}$$

We examine the term $Y'(t) = -16\pi^4 n^8 e^{8t} + 180\pi^3 n^6 e^{6t} - 529\pi^2 n^4 e^{4t} + 408.75\pi n^2 e^{2t} - \frac{75*5}{8}$ in Eq. 47. We see that, for $n = 2$ and $t = 0$, $Y'(t) = -\pi^3 2^6 (16\pi * 4 - 180) - \pi^2 (529\pi * 2^2 - 408.75) - \frac{75*5}{8} < 0$. We see that, **for all** $n > 1$ and $t \geq 0$, $Y'(t) < 0$. Hence $X_{32}(t) < 0$ for all $t \geq 0$.

Hence $X_{22}(t)$ is a **strictly decreasing function** for $t > 0$. **(Result E)** (example plot)

4.6. Second derivative given by X_{21} is a strictly increasing function for $t_{min} < t < t'_{min}$

We can show that $X_{21}(t)$ is a **strictly increasing function** for $t_{min} < t < t'_{min}$. We take the derivative of $X_{21}(t) = e^{-\pi e^{2t}} e^{\frac{5t}{2}} (8\pi^3 e^{6t} - 56\pi^2 e^{4t} + \frac{165}{2}\pi e^{2t} - \frac{75}{4})$, given by $X_{31}(t)$ and set $n = 1$ in the summand in Eq. 47 as follows.

$$X_{31}(t) = e^{-\pi e^{2t}} e^{\frac{5t}{2}} Z'(t), \quad Z'(t) = -16\pi^4 e^{8t} + 180\pi^3 e^{6t} - 529\pi^2 e^{4t} + 408.75\pi e^{2t} - \frac{75 * 5}{8}$$

(48)

We see that $X_{21}(t)$ in Eq. 45 is **negative** at $0 \leq t < t'_{min}$ (using Result N). In the interval $t_{min} < t < t'_{min}$, we consider 3 cases for $X_{21}(t)$ as follows.

Case 1: $X_{21}(t)$ starts from the negative value at $t = t_{min}$ and increases to zero at $t = t'_{min}$. This is **possible** because we know that $X_{21}(t)$ in Eq. 45 is **negative** at $0 \leq t < t'_{min}$ and $X_{21}(t) = 0$ at $t = t'_{min}$ (using Result N).

Case 2: $X_{21}(t)$ starts from the negative value at $t = t_{min}$ and decreases to a more negative value and reaches a negative minimum at $t = t_c < t'_{min}$ and then starts increasing towards zero. This requires $X_{31}(t) < 0$ in the interval $t_{min} < t < t_c$ and $X_{31}(t) > 0$ in the interval $t_c < t < t'_{min}$, which is **NOT** possible because of the dominant term $-16\pi^4 e^{8t}$ in Eq. 48.

Case 3: $X_{21}(t)$ starts from the negative value at $t = t_{min}$ and increases to a less negative value and has a point of inflection at $t = t_d < t'_{min}$ and then starts decreasing towards a more negative value and has a second point of inflection at $t = t_e < t'_{min}$ and then starts increasing towards zero. This requires $X_{31}(t) < 0$ in the interval $t_d < t < t_e$ and $X_{31}(t) > 0$ in the interval $t_e < t < t'_{min}$, which is **NOT** possible because of the dominant term $-16\pi^4 e^{8t}$ in Eq. 48.

Similarly, we can extend the argument to the case where $X_{21}(t)$ has many points of inflection in the interval $t_{min} < t < t'_{min}$.

Hence we see that only **Case 1** is possible and hence $X_{21}(t)$ is a **strictly increasing function** for $t_{min} < t < t'_{min}$. **(Result F)** (example plot)

4.7. Result $E_0(t - t_{2c}) - E_0(t + t_{2c}) > 0$

It is shown in Section 4 that $E_0(t)$ is **strictly decreasing** for $t > 0$. In this section, it is shown that $E_0(t - t_{2c}) - E_0(t + t_{2c}) > 0$, for $0 < t < t_{0c}$ and $t_{2c} = 2t_{0c}$ in Eq. 34 .

Given that $E_0(t)$ is a **strictly decreasing** function for $t > 0$ and $E_0(t)$ is an **even** function of variable t , and $t_{2c} = 2t_{0c}$, we see that, in the interval $0 < t < t_{0c}$, $E_0(t + t_{2c}) = E_0(t + 2t_{0c})$ ranges from $E_0(2t_{0c})$ to $E_0(3t_{0c})$, which is **less than** $E_0(t - t_{2c}) = E_0(t - 2t_{0c})$ which ranges from $E_0(-2t_{0c})$ to $E_0(-t_{0c})$. Hence we see that $E_0(t - t_{2c}) > E_0(t + t_{2c})$, in the interval $0 < t < t_{0c}$. At $t = 0$, $E_0(t - t_{2c}) = E_0(t + t_{2c})$. At $t = t_{0c}$, $E_0(t - t_{2c}) > E_0(t + t_{2c})$ because $E_0(-t_{0c}) > E_0(3t_{0c})$.

Hence $E_0(t - t_{2c}) - E_0(t + t_{2c}) > 0$ for $0 < t < t_{0c}$ in Eq. 34 , for $t_{0c} > 0$.

5. $\omega_z(t_2, t_0)$ is a continuous function of t_0

It is shown in this section that $\omega_z(t_2, t_0)$ is a continuous function of t_0 in the neighbourhood $[t_0 - \delta t_0, t_0 + \delta t_0]$ for all $|t_0| < \infty$, for **each** fixed value of t_2 .

• $G_R(\omega) = G_R(\omega, t_2, t_0)$ in Eq. 14 is copied below, which is a **continuous** function of ω which is differentiable **at least** once with respect to ω . (Eq. 50).

$$G_R(\omega) = G_R(\omega, t_2, t_0) = e^{-2\sigma t_0} \int_{-\infty}^0 [E'_0(\tau + t_0, t_2)e^{-2\sigma\tau} + E'_{0n}(\tau - t_0, t_2)] \cos(\omega\tau) d\tau \\ + e^{2\sigma t_0} \int_{-\infty}^0 [E'_0(\tau - t_0, t_2)e^{-2\sigma\tau} + E'_{0n}(\tau + t_0, t_2)] \cos(\omega\tau) d\tau \quad (49)$$

Given that $E_0(\tau) > 0$ for $|\tau| < \infty$ and $\lim_{\tau \rightarrow \pm\infty} E_0(\tau) = 0$ (Appendix C.1), we see that $G_R(\omega) > 0$ at $\omega = 0$. **Set** $t_0 = 0$ and $G_R(\omega)$ passes through its **first zero** at $\omega = \omega_z(t_2, t_0) = \omega_z(t_2, 0)$. In the rest of this section, we consider the **interval** $[-\delta t_0, \delta t_0]$ around $t_0 = 0$, in $\omega_z(t_2, t_0)$. There are 3 possibilities.

Case 1: $G_R(\omega) < 0$ for $\omega = \omega_z(t_2, 0) + dw$, $G_R(\omega) > 0$ for $\omega = \omega_z(t_2, 0) - dw$ for infinitesimal dw (example plot)

In this case, we will show in Section 5.1 that $\omega_z(t_2, t_0)$ is a continuous function of t_0 in the interval $[-\delta t_0, \delta t_0]$, in the neighborhood around the first zero crossing at $\omega = \omega_z(t_2, t_0) = \omega_z(t_2, 0)$.

Case 2: $G_R(\omega) > 0$ for $\omega = \omega_z(t_2, 0) + dw$, $G_R(\omega) > 0$ for $\omega = \omega_z(t_2, 0) - dw$ (example plot)

In this case, $\frac{dG_R(\omega)}{d\omega} = 0$ at the **same** $\omega = \omega_z(t_2, 0)$ because $\frac{dG_R(\omega)}{d\omega} < 0$ at $\omega = \omega_z(t_2, 0) - dw$ and $\frac{dG_R(\omega)}{d\omega} > 0$ at $\omega = \omega_z(t_2, 0) + dw$.

$$\frac{dG_R(\omega, t_2, t_0)}{d\omega} = -[e^{-2\sigma t_0} \int_{-\infty}^0 \tau [E'_0(\tau + t_0, t_2)e^{-2\sigma\tau} + E'_{0n}(\tau - t_0, t_2)] \sin(\omega\tau) d\tau \\ + e^{2\sigma t_0} \int_{-\infty}^0 \tau [E'_0(\tau - t_0, t_2)e^{-2\sigma\tau} + E'_{0n}(\tau + t_0, t_2)] \sin(\omega\tau) d\tau] \quad (50)$$

In this case, we will show in Section 5.2 that $\omega_z(t_2, t_0)$ is a continuous function of t_0 in the interval $[-\delta t_0, \delta t_0]$, in the neighborhood around the first zero crossing at $\omega = \omega_z(t_2, t_0) = \omega_z(t_2, 0)$.

Case 3: $G_R(\omega) = 0$ for $\omega = \omega_z(t_2, 0)$ and $\omega = \omega_z(t_2, 0) + dw$.

In this case, $\frac{dG_R(\omega)}{d\omega} = 0$ at the **same** $\omega = \omega_z(t_2, 0)$ because $\frac{dG_R(\omega)}{d\omega} < 0$ at $\omega = \omega_z(t_2, 0) - dw$ and $\frac{dG_R(\omega)}{d\omega} = 0$ at $\omega = \omega_z(t_2, 0)$. The arguments are similar to that of Case 2 presented in Section 5.2 where it is shown that $\omega_z(t_2, t_0)$ is a continuous function of t_0 in the interval $[-\delta t_0, \delta t_0]$, in the neighborhood around the first zero crossing at $\omega = \omega_z(t_2, t_0) = \omega_z(t_2, 0)$.

5.1. Case 1: $G_R(\omega) < 0$ **for** $\omega = \omega_z(t_2, 0) + dw$, $G_R(\omega) > 0$ **for** $\omega = \omega_z(t_2, 0) - dw$

• Consider the **segment S** in $G_R(\omega, t_2, t_0)$ in the neighborhood around the first zero crossing where $\frac{dG_R(\omega, t_2, t_0)}{d\omega} < 0$. (Segment S is the portion between the green lines in example plot)

• In the **segment S**, $G_R(\omega, t_2, t_0)$ in Eq. 49 is a **continuous** function of ω , for **each** value of t_0 and t_2 . Hence $G_R(\omega, t_2, t_0 - \delta t_0)$ and $G_R(\omega, t_2, t_0 + \delta t_0)$ are **continuous** functions of ω , which are differentiable **at least** once, and $G_R(\omega, t_2, t_0 \pm \delta t_0)$ tends to $G_R(\omega, t_2, t_0)$, as infinitesimal $\delta t_0 \rightarrow 0$.

$$\begin{aligned}
G_R(\omega, t_2, t_0) &= e^{-2\sigma t_0} \int_{-\infty}^0 [E'_0(\tau + t_0, t_2)e^{-2\sigma\tau} + E'_{0n}(\tau - t_0, t_2)] \cos(\omega\tau) d\tau \\
&\quad + e^{2\sigma t_0} \int_{-\infty}^0 [E'_0(\tau - t_0, t_2)e^{-2\sigma\tau} + E'_{0n}(\tau + t_0, t_2)] \cos(\omega\tau) d\tau \\
G_R(\omega, t_2, t_0 + \delta t_0) &= e^{-2\sigma(t_0 + \delta t_0)} \int_{-\infty}^0 [E'_0(\tau + t_0 + \delta t_0, t_2)e^{-2\sigma\tau} + E'_{0n}(\tau - t_0 - \delta t_0, t_2)] \cos(\omega\tau) d\tau \\
&\quad + e^{2\sigma(t_0 + \delta t_0)} \int_{-\infty}^0 [E'_0(\tau - t_0 - \delta t_0, t_2)e^{-2\sigma\tau} + E'_{0n}(\tau + t_0 + \delta t_0, t_2)] \cos(\omega\tau) d\tau \\
G_R(\omega, t_2, t_0 - \delta t_0) &= e^{-2\sigma(t_0 - \delta t_0)} \int_{-\infty}^0 [E'_0(\tau + t_0 - \delta t_0, t_2)e^{-2\sigma\tau} + E'_{0n}(\tau - t_0 + \delta t_0, t_2)] \cos(\omega\tau) d\tau \\
&\quad + e^{2\sigma(t_0 - \delta t_0)} \int_{-\infty}^0 [E'_0(\tau - t_0 + \delta t_0, t_2)e^{-2\sigma\tau} + E'_{0n}(\tau + t_0 - \delta t_0, t_2)] \cos(\omega\tau) d\tau \\
\lim_{\delta t_0 \rightarrow 0} G_R(\omega, t_2, t_0 + \delta t_0) &= G_R(\omega, t_2, t_0) \\
\lim_{\delta t_0 \rightarrow 0} G_R(\omega, t_2, t_0 - \delta t_0) &= G_R(\omega, t_2, t_0)
\end{aligned}$$

(51)

• In the **segment S**, $G_R(\omega, t_2, t_0)$ in Eq. 51 is a **continuous** function of ω , for **each** value of t_0 and t_2 and $\frac{dG_R(\omega, t_2, t_0)}{d\omega} < 0$ in the neighborhood around the **first zero crossing**. If we **fix** the X-coordinate ω and t_2 , $G_R(\omega, t_2, t_0)$ is a **continuous** function of t_0 , for **each** fixed value of ω . Hence, for **each** fixed value of ω , as we change t_0 by an infinitesimal δt_0 , $G_R(\omega, t_2, t_0 - \delta t_0)$ and $G_R(\omega, t_2, t_0 + \delta t_0)$ in Eq. 51, move towards $G_R(\omega, t_2, t_0)$ in a **continuous** manner, as $\delta t_0 \rightarrow 0$. Every point in the segment S, moves continuously, as we change t_0 by an infinitesimal δt_0 .

This also applies to the first **zero crossing** in $G_R(\omega, t_2, t_0)$ in the segment S, which corresponds to $\omega_z(t_2, t_0) = \omega_z(t_2, 0)$ at $t_0 = 0$ where $G_R(\omega, t_2, t_0) = 0$ in Eq. 51. The **zero crossing** moves **continuously**, as we change t_0 by an infinitesimal δt_0 . This is explained below.

• **Explanation:** This is shown by an **example** plot. **Red** plot corresponds to $G_R(\omega, t_2, t_0)$ with zero crossing at point P_0 , **Green** plot corresponds to $G_R(\omega, t_2, t_0 + \delta t_0)$ with zero crossing at point P_{11} and **Blue** plot corresponds to $G_R(\omega, t_2, t_0 - \delta t_0)$ with zero crossing at point P_{21} .

We **define** the **point** P_{12} in $G_R(\omega, t_2, t_0 + \delta t_0)$ as the point which has the **fixed X-coordinate** $\omega = \omega_z(t_2, 0)$. We **define** the **point** P_{22} in $G_R(\omega, t_2, t_0 - \delta t_0)$ as the point which has the **fixed X-coordinate** $\omega = \omega_z(t_2, 0)$.

We **define** the **point** P_{11} in $G_R(\omega, t_2, t_0 + \delta t_0)$ as the **zero crossing point** which has the **fixed Y-coordinate** which equals zero. We **define** the **point** P_{21} in $G_R(\omega, t_2, t_0 - \delta t_0)$ as the **zero crossing point** which has the **fixed Y-coordinate** which equals zero.

As we change t_0 by an infinitesimal δt_0 , $G_R(\omega, t_2, t_0 + \delta t_0)$ in Eq. 51 moves towards $G_R(\omega, t_2, t_0)$ in a **continuous** manner, for **each fixed** value of ω and t_2 , including the zero crossing point, as follows. The **point** P_{12} in $G_R(\omega, t_2, t_0 + \delta t_0)$ which corresponds to the **fixed X-coordinate** $\omega = \omega_z(t_2, 0)$, moves towards corresponding point P_0 in $G_R(\omega, t_2, t_0)$, for the **same** $\omega = \omega_z(t_2, 0)$ in a **continuous** manner, as $\delta t_0 \rightarrow 0$. Given that P_0 is a **zero crossing point** in $G_R(\omega, t_2, t_0)$, this is equivalent to the **Zero crossing point** P_{11} in $G_R(\omega, t_2, t_0 + \delta t_0)$ moving towards corresponding **zero crossing point** P_0 in $G_R(\omega, t_2, t_0)$ in a **continuous** manner, as $\delta t_0 \rightarrow 0$.

Similarly, as we change t_0 by an infinitesimal δt_0 , $G_R(\omega, t_2, t_0 - \delta t_0)$ in Eq. 51 moves towards $G_R(\omega, t_2, t_0)$ in a **continuous** manner as follows. The **point** P_{22} in $G_R(\omega, t_2, t_0 - \delta t_0)$ which corresponds to the **fixed X-coordinate** $\omega = \omega_z(t_2, 0)$, moves towards corresponding point P_0 in $G_R(\omega, t_2, t_0)$, for the **same** $\omega = \omega_z(t_2, 0)$ in a **continuous** manner, as $\delta t_0 \rightarrow 0$. Given that P_0 is a **zero crossing point** in $G_R(\omega, t_2, t_0)$, this is equivalent to the **Zero crossing point** P_{21} in $G_R(\omega, t_2, t_0 - \delta t_0)$ moving towards corresponding **zero crossing point** P_0 in $G_R(\omega, t_2, t_0)$ in a **contin-**

596 **uous** manner, as $\delta t_0 \rightarrow 0$.

597

598 • Hence in the **segment** S, $\omega_z(t_2, t_0)$ is a **continuous** function of t_0 in the neighborhood $[-\delta t_0, \delta t_0]$ around the
599 first zero crossing at $\omega = \omega_z(t_2, t_0) = \omega_z(t_2, 0)$ at $t_0 = 0$.

$$\begin{aligned}
 G_R(\omega_z(t_2, t_0), t_2, t_0) &= e^{-2\sigma t_0} \int_{-\infty}^0 [E'_0(\tau + t_0, t_2)e^{-2\sigma\tau} + E'_{0n}(\tau - t_0, t_2)] \cos(\omega_z(t_2, t_0)\tau) d\tau \\
 &\quad + e^{2\sigma t_0} \int_{-\infty}^0 [E'_0(\tau - t_0, t_2)e^{-2\sigma\tau} + E'_{0n}(\tau + t_0, t_2)] \cos(\omega_z(t_2, t_0)\tau) d\tau = 0 \\
 G_R(\omega_z(t_2, t_0 + \delta t_0), t_2, t_0 + \delta t_0) &= \\
 e^{-2\sigma(t_0 + \delta t_0)} \int_{-\infty}^0 [E'_0(\tau + t_0 + \delta t_0, t_2)e^{-2\sigma\tau} + E'_{0n}(\tau - t_0 - \delta t_0, t_2)] \cos(\omega_z(t_2, t_0 + \delta t_0)\tau) d\tau \\
 &\quad + e^{2\sigma(t_0 + \delta t_0)} \int_{-\infty}^0 [E'_0(\tau - t_0 - \delta t_0, t_2)e^{-2\sigma\tau} + E'_{0n}(\tau + t_0 + \delta t_0, t_2)] \cos(\omega_z(t_2, t_0 + \delta t_0)\tau) d\tau = 0
 \end{aligned}$$

600 (52)

601 5.2. **Case 2:** $G_R(\omega) > 0$ **for** $\omega = \omega_z(t_2, 0) + dw$, $G_R(\omega) > 0$ **for** $\omega = \omega_z(t_2, 0) - dw$

602

603 • In this case, $\frac{dG_R(\omega)}{d\omega} = 0$ at the **same** $\omega = \omega_z(t_2, t_0)$ because $\frac{dG_R(\omega)}{d\omega} < 0$ at $\omega = \omega_z(t_2, t_0) - dw$ and $\frac{dG_R(\omega)}{d\omega} > 0$
604 at $\omega = \omega_z(t_2, t_0) + dw$.

605

606 • Consider the **segment** S' in $\frac{dG_R(\omega, t_2, t_0)}{d\omega}$ in the neighborhood around the first zero crossing where $\frac{d^2G_R(\omega, t_2, t_0)}{d\omega^2} > 0$.
607 (Segment S' is the portion between the green lines in example plot) In this segment S', $\frac{dG_R(\omega, t_2, t_0)}{d\omega}$ is a **continuous**
608 function of ω which is differentiable **at least** once.(Section ??)

609

610 • In the **segment** S', $\frac{dG_R(\omega, t_2, t_0)}{d\omega} = 0$ at the **same** $\omega = \omega_z(t_2, t_0)$. The arguments in Section 5.1 can be applied
611 here, with $G_R(\omega, t_2, t_0)$ replaced by $\frac{dG_R(\omega, t_2, t_0)}{d\omega}$.

612

613 Hence $\omega_z(t_2, t_0)$ is a **continuous** function of t_0 in the neighborhood $[-\delta t_0, \delta t_0]$ around the first zero crossing at
614 $\omega = \omega_z(t_2, t_0) = \omega_z(t_2, 0)$ at $t_0 = 0$ in the **segment** S'.

615

616 We can use similar arguments and see that $\omega_z(t_2, t_0)$ is a **continuous** function of t_0 in the neighbourhood $[t_0 -$
617 $\delta t_0, t_0 + \delta t_0]$ for all $|t_0| < \infty$, for **each** fixed value of t_2 .

618 5.3. Further Points

619

620 • Using arguments in previous subsections, we see that $\omega_z(t_2, t_0)$ is a **continuous** function of t_2 in the neighbour-
621 hood $[t_2 - \delta t_2, t_2 + \delta t_2]$ for all $|t_2| < \infty$, for **each** fixed value of t_0 .

622

623 • We **set** $t_2 = Kt_0$ for even positive integer K . Using arguments in previous subsections, we see that $\omega_z(Kt_0, t_0)$
624 is a **continuous** function of t_0 in the neighbourhood $[t_0 - \delta t_0, t_0 + \delta t_0]$ for all $|t_0| < \infty$.

625 6. Hurwitz Zeta Function and related functions

626

627 We can show that the new method is **not** applicable to Hurwitz zeta function and related zeta functions and
628 **does not** contradict the existence of their non-trivial zeros away from the critical line with real part of $s = \frac{1}{2}$.
629 The new method requires the **symmetry** relation $\xi(s) = \xi(1-s)$ and hence $\xi(\frac{1}{2} + i\omega) = \xi(\frac{1}{2} - i\omega)$ when eval-
630 uated at the critical line $s = \frac{1}{2} + i\omega$. This means $\xi(\frac{1}{2} + i\omega) = E_{0\omega}(\omega) = E_{0\omega}(-\omega)$ and $E_0(t) = E_0(-t)$ where

$E_0(t) = 2 \sum_{n=1}^{\infty} [2\pi^2 n^4 e^{4t} - 3\pi n^2 e^{2t}] e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}}$ and this condition is satisfied for Riemann's Zeta function.

It is **not** known that Hurwitz Zeta Function given by $\zeta(s, a) = \sum_{m=0}^{\infty} \frac{1}{(m+a)^s}$ satisfies a symmetry relation similar to $\xi(s) = \xi(1-s)$ where $\xi(s)$ is an entire function, for $a \neq 1$ and hence the condition $E_0(t) = E_0(-t)$ is **not** known to be satisfied^[6]. Hence the new method is **not** applicable to Hurwitz zeta function and **does not** contradict the existence of their non-trivial zeros away from the critical line.

Dirichlet L-functions satisfy a symmetry relation $\xi(s, \chi) = \epsilon(\chi) \xi(1-s, \bar{\chi})$ ^[7] which does **not** translate to $E_0(t) = E_0(-t)$ required by the new method and hence this proof is **not** applicable to them.

We know that $\zeta(s) = \sum_{m=1}^{\infty} \frac{1}{m^s}$ diverges for real part of $s \leq 1$. Hence we derive a convergent and entire function $\xi(s)$ using the well known theorem $F(x) = 1 + 2 \sum_{n=1}^{\infty} e^{-\pi n^2 x} = \frac{1}{\sqrt{x}} (1 + 2 \sum_{n=1}^{\infty} e^{-\pi \frac{n^2}{x}})$, where $x > 0$ is real and then derive $E_0(t) = 2 \sum_{n=1}^{\infty} [2\pi^2 n^4 e^{4t} - 3\pi n^2 e^{2t}] e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}}$ (Appendix D). In the case of **Hurwitz zeta** function and **other zeta functions** with non-trivial zeros away from the critical line, it is **not** known if a corresponding relation similar to $F(x)$ exists, which enables derivation of a convergent and entire function $\xi(s)$ and results in $E_0(t)$ as a Fourier transformable, real, even and analytic function. Hence the new method presented in this paper is **not** applicable to Hurwitz zeta function and related zeta functions.

The proof of Riemann Hypothesis presented in this paper is **only** for the specific case of Riemann's Zeta function and **only** for the **critical strip** $0 \leq |\sigma| < \frac{1}{2}$. This proof requires both $E_p(t)$ and $E_{p\omega}(\omega)$ to be Fourier transformable where $E_p(t) = E_0(t) e^{-\sigma t}$ is a real analytic function and uses the fact that $E_0(t)$ is an **even** function which is **strictly decreasing** function for $t \geq \frac{1}{8}$. These conditions may **not** be satisfied for many other functions including those which have non-trivial zeros away from the critical line and hence the new method may **not** be applicable to such functions.

If the proof presented in this paper is internally consistent and does not have mistakes and gaps, then it should be considered correct, **regardless** of whether it contradicts any previously known external theorems, because it is possible that those previously known external theorems may be incorrect.

References

- [1] Bernhard Riemann, On the Number of Prime Numbers less than a Given Quantity.(Ueber die Anzahl der Primzahlen untereiner gegebenen Grosse.) Monatsberichte der Berliner Akademie, November 1859. (Link to Riemann's 1859 paper)
- [2] Hardy, G.H., Littlewood, J.E. The zeros of Riemann's zeta-function on the critical line. Mathematische Zeitschrift volume 10, pp.283 to 317 (1921).
- [3] E. C. Titchmarsh, The Theory of the Riemann Zeta Function. (1986) pp.254 to 255
- [4] Fern Ellison and William J. Ellison, Prime Numbers (1985).pp147 to 152
- [5] J. Brian Conrey, The Riemann Hypothesis (2003). (Link to Brian Conrey's 2003 article)
- [6] Mathworld article on Hurwitz Zeta functions. (Link)
- [7] Wikipedia article on Dirichlet L-functions. (Link)

Appendix A. Derivation of $E_p(t)$

Let us start with Riemann's Xi Function $\xi(s)$ evaluated at $s = \frac{1}{2} + i\omega$ given by $\xi(\frac{1}{2} + i\omega) = E_{0\omega}(\omega)$. Its inverse Fourier Transform is given by $E_0(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} E_{0\omega}(\omega) e^{i\omega t} d\omega = 2 \sum_{n=1}^{\infty} [2\pi^2 n^4 e^{4t} - 3\pi n^2 e^{2t}] e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}}$ (link). This is

re-derived in Appendix D.

We will show in this section that the inverse Fourier Transform of the function $\xi(\frac{1}{2} + \sigma + i\omega) = E_{p\omega}(\omega)$, is given by $E_p(t) = E_0(t)e^{-\sigma t}$ where $0 \leq |\sigma| < \frac{1}{2}$ is real.

$$\begin{aligned}\xi(\frac{1}{2} + \sigma + i\omega) &= \xi(\frac{1}{2} + i(\omega - i\sigma)) = E_{p\omega}(\omega) = E_{0\omega}(\omega - i\sigma) \\ E_p(t) &= \frac{1}{2\pi} \int_{-\infty}^{\infty} E_{p\omega}(\omega) e^{i\omega t} d\omega = \frac{1}{2\pi} \int_{-\infty}^{\infty} E_{0\omega}(\omega - i\sigma) e^{i\omega t} d\omega\end{aligned}$$

(A.1)

We substitute $\omega' = \omega - i\sigma$ in Eq. A.1 as follows.

$$E_p(t) = e^{-\sigma t} \frac{1}{2\pi} \int_{-\infty - i\sigma}^{\infty - i\sigma} E_{0\omega}(\omega') e^{i\omega' t} d\omega'$$

(A.2)

We can evaluate the above integral in the complex plane using contour integration, substituting $\omega' = z = x + iy$ and we use a rectangular contour comprised of C_1 along the line $x = [-\infty, \infty]$, C_2 along the line $y = [\infty, \infty - i\sigma]$, C_3 along the line $x = [\infty - i\sigma, -\infty - i\sigma]$ and then C_4 along the line $y = [-\infty - i\sigma, -\infty]$. We can see that $E_{0\omega}(z) = \xi(\frac{1}{2} + iz)$ has no singularities in the region bounded by the contour because $\xi(\frac{1}{2} + iz)$ is an entire function in the Z -plane.

In **Appendix C.1**, we show that $\int_{-\infty}^{\infty} |E_p(t)| dt$ is finite and $E_p(t) = E_0(t)e^{-\sigma t}$ is an absolutely integrable function, for $0 \leq |\sigma| < \frac{1}{2}$.

We use the fact that $E_{0\omega}(z) = \xi(\frac{1}{2} + iz) = \xi(\frac{1}{2} - y + ix) = \int_{-\infty}^{\infty} E_0(t) e^{-izt} dt = \int_{-\infty}^{\infty} E_0(t) e^{yt} e^{-ixt} dt$, **goes to zero** as $x \rightarrow \pm\infty$ when $-\sigma \leq y \leq 0$, as per Riemann-Lebesgue Lemma (link), because $E_0(t)e^{yt}$ is a absolutely integrable function in the interval $-\infty \leq t \leq \infty$. Hence the integral in Eq. A.2 **vanishes** along the contours C_2 and C_4 . Using Cauchy's Integral theroem, we can write Eq. A.2 as follows.

$$\begin{aligned}E_p(t) &= e^{-\sigma t} \frac{1}{2\pi} \int_{-\infty}^{\infty} E_{0\omega}(\omega') e^{i\omega' t} d\omega' \\ E_p(t) &= E_0(t) e^{-\sigma t} = 2 \sum_{n=1}^{\infty} [2\pi^2 n^4 e^{4t} - 3\pi n^2 e^{2t}] e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}} e^{-\sigma t}\end{aligned}$$

(A.3)

Thus we have arrived at the desired result $E_p(t) = E_0(t)e^{-\sigma t}$.

Appendix B. Properties of Fourier Transforms Part 1

In this section, some well-known properties of Fourier transforms are re-derived.

Appendix B.1. Convolution Theorem: Multiplication of $g(t)$ and $h(t)$ corresponds to convolution in Fourier transform domain

We start with the Fourier transform equation $F(\omega) = \int_{-\infty}^{\infty} f(t) e^{-i\omega t} dt$ where $f(t) = g(t)h(t)$ and show that $F(\omega) = \frac{1}{2\pi} [G(\omega) * H(\omega)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} G(\omega') H(\omega - \omega') d\omega'$ obtained by the **convolution** of the functions $G(\omega)$ and $H(\omega)$ which correspond to the Fourier transforms of $g(t)$ and $h(t)$ respectively.

$$F(\omega) = \int_{-\infty}^{\infty} f(t)e^{-i\omega t} dt = \int_{-\infty}^{\infty} g(t)h(t)e^{-i\omega t} dt \quad (\text{B.1})$$

We use the inverse Fourier transform equation $g(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} G(\omega')e^{i\omega' t} d\omega'$ and we interchange the order of integration in equations below using Fubini's theorem (link).

$$\begin{aligned} F(\omega) &= \frac{1}{2\pi} \int_{-\infty}^{\infty} \left[\int_{-\infty}^{\infty} G(\omega')e^{i\omega' t} d\omega' \right] h(t)e^{-i\omega t} dt \\ F(\omega) &= \frac{1}{2\pi} \int_{-\infty}^{\infty} G(\omega') \left[\int_{-\infty}^{\infty} e^{i\omega' t} h(t)e^{-i\omega t} dt \right] d\omega' \\ F(\omega) &= \frac{1}{2\pi} \int_{-\infty}^{\infty} G(\omega') \left[\int_{-\infty}^{\infty} h(t)e^{-i(\omega-\omega')t} dt \right] d\omega' \end{aligned} \quad (\text{B.2})$$

We substitute $\int_{-\infty}^{\infty} h(t)e^{-i(\omega-\omega')t} dt = H(\omega - \omega')$ in Eq. B.2 and arrive at the convolution theorem.

$$F(\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} G(\omega')H(\omega - \omega')d\omega' = \int_{-\infty}^{\infty} g(t)h(t)e^{-i\omega t} dt \quad (\text{B.3})$$

Appendix B.2. *Fourier transform of Real $g(t)$*

In this section, we show that the Fourier transform of a real function $g(t)$, given by $G(\omega) = G_R(\omega) + iG_I(\omega)$ has the properties given by $G_R(-\omega) = G_R(\omega)$ and $G_I(-\omega) = -G_I(\omega)$.

$$\begin{aligned} G(\omega) &= \int_{-\infty}^{\infty} g(t)e^{-i\omega t} dt = G_R(\omega) + iG_I(\omega) \\ G_R(\omega) &= \int_{-\infty}^{\infty} g(t) \cos(\omega t) dt = G_R(-\omega) \\ G_I(\omega) &= - \int_{-\infty}^{\infty} g(t) \sin(\omega t) dt = -G_I(-\omega) \end{aligned} \quad (\text{B.4})$$

Appendix B.3. *Even part of $g(t)$ corresponds to real part of Fourier transform $G(\omega)$*

In this section, we show that the **even part** of real function $g(t)$, given by $g_{\text{even}}(t) = \frac{1}{2}[g(t) + g(-t)]$, corresponds to **real part** of its Fourier transform $G(\omega)$. We use the fact that $G_R(-\omega) = G_R(\omega)$ and $G_I(-\omega) = -G_I(\omega)$ for a real function $g(t)$.

$$\begin{aligned} G(\omega) &= \int_{-\infty}^{\infty} g(t)e^{-i\omega t} dt = G_R(\omega) + iG_I(\omega) \\ \int_{-\infty}^{\infty} g_{\text{even}}(t)e^{-i\omega t} dt &= \int_{-\infty}^{\infty} \frac{1}{2}[g(t) + g(-t)]e^{-i\omega t} dt = \frac{1}{2}[G(\omega) + G(-\omega)] = G_R(\omega) \end{aligned} \quad (\text{B.5})$$

Appendix B.4. Odd part of $g(t)$ corresponds to imaginary part of Fourier transform $G(\omega)$

In this section, we show that the **odd part** of real function $g(t)$, given by $g_{\text{odd}}(t) = \frac{1}{2}[g(t) - g(-t)]$, corresponds to **imaginary part** of its Fourier transform $G(\omega)$. We use the fact that $G_R(-\omega) = G_R(\omega)$ and $G_I(-\omega) = -G_I(\omega)$ for a real function $g(t)$.

$$\begin{aligned} G(\omega) &= \int_{-\infty}^{\infty} g(t)e^{-i\omega t}dt = G_R(\omega) + iG_I(\omega) \\ \int_{-\infty}^{\infty} g_{\text{odd}}(t)e^{-i\omega t}dt &= \int_{-\infty}^{\infty} \frac{1}{2}[g(t) - g(-t)]e^{-i\omega t}dt = \frac{1}{2}[G(\omega) - G(-\omega)] = iG_I(\omega) \end{aligned} \tag{B.6}$$

Appendix C. Properties of Fourier Transforms Part 2

Appendix C.1. $E_p(t), h(t), g(t)$ are absolutely integrable functions and their Fourier Transforms are finite.

The inverse Fourier Transform of the function $E_{p\omega}(\omega) = \xi(\frac{1}{2} + \sigma + i\omega)$ is given by $E_p(t) = E_0(t)e^{-\sigma t} = \frac{1}{2\pi} \int_{-\infty}^{\infty} E_{p\omega}(\omega)e^{i\omega t}d\omega$. We see that $E_0(t) = 2 \sum_{n=1}^{\infty} [2\pi^2 n^4 e^{4t} - 3\pi n^2 e^{2t}]e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}} > 0$ for all $0 \leq t < \infty$. Given that $E_0(t) = E_0(-t)$, we see that $E_0(t) > 0$ and $E_p(t) = E_0(t)e^{-\sigma t} > 0$ for all $-\infty < t < \infty$.

As $t \rightarrow \infty$, $E_p(t)$ goes to zero, due to the term $e^{-\pi n^2 e^{2t}}$. As $t \rightarrow -\infty$, $E_p(t)$ goes to zero, because for every value of n , the term $e^{-\pi n^2 e^{2t}} e^{\frac{5t}{2}} e^{-\sigma t}$ goes to zero, for $0 \leq |\sigma| < \frac{1}{2}$. Hence $E_p(t) = E_0(t)e^{-\sigma t} = 0$ at $t \rightarrow \pm\infty$ and we showed that $E_p(t) > 0$ for all $-\infty < t < \infty$. Hence $E_{p\omega}(\omega) = \int_{-\infty}^{\infty} E_p(t)e^{-i\omega t}dt$, evaluated at $\omega = 0$ **cannot** be zero. Hence $E_{p\omega}(\omega)$ **does not have a zero** at $\omega = 0$ and hence $\omega_0 \neq 0$.

Given that $\xi(\frac{1}{2} + \sigma + i\omega) = E_{p\omega}(\omega)$ is an entire function in the whole of s-plane, it is finite for $|\omega| \leq \infty$ and also for $\omega = 0$. Hence $\int_{-\infty}^{\infty} E_p(t)dt$ is finite. We see that $E_p(t) \geq 0$ for all $|t| \leq \infty$. Hence we can write $\int_{-\infty}^{\infty} |E_p(t)|dt$ is finite and $E_p(t)$ is an absolutely **integrable function** and its Fourier transform $E_{p\omega}(\omega)$ goes to zero as $\omega \rightarrow \pm\infty$, as per Riemann Lebesgue Lemma (link).

Let us consider a new function $g(t) = f(t)e^{-\sigma t}u(-t) + f(t)e^{\sigma t}u(t)$ where $g(t)$ is a real function of variable t and $u(t)$ is Heaviside unit step function and $0 < \sigma < \frac{1}{2}$. We can see that $g(t)h(t) = f(t)$ where $h(t) = e^{\sigma t}u(-t) + e^{-\sigma t}u(t)$.

We can see that $h(t) = e^{\sigma t}u(-t) + e^{-\sigma t}u(t)$ is an absolutely **integrable function** because $\int_{-\infty}^{\infty} |h(t)|dt = \int_{-\infty}^{\infty} h(t)dt = [\int_{-\infty}^{\infty} h(t)e^{-i\omega t}dt]_{\omega=0} = [\frac{1}{\sigma-i\omega} + \frac{1}{\sigma+i\omega}]_{\omega=0} = \frac{2}{\sigma}$, is finite for $0 < \sigma < \frac{1}{2}$ and its Fourier transform $H(\omega)$ goes to zero as $\omega \rightarrow \pm\infty$, as per Riemann Lebesgue Lemma (link).

It is shown in Appendix C.4 that $E_0(t)$ and $E_0(t)e^{-2\sigma t}$ have fall-off rates **at least** $\frac{1}{t^2}$ as $|t| \rightarrow \infty$ and hence are absolutely **integrable** functions and the integrals $\int_{-\infty}^{\infty} |E_0(t)|dt < \infty$ and $\int_{-\infty}^{\infty} |E_0(t)e^{-2\sigma t}|dt < \infty$. Hence $g(t) = E_0(t)e^{-2\sigma t}u(-t) + E_0(t)u(t)$ is an absolutely **integrable function** and $\int_{-\infty}^{\infty} |g(t)|dt = \int_{-\infty}^{\infty} g(t)dt$ is finite and its Fourier transform $G(\omega)$ goes to zero as $\omega \rightarrow \pm\infty$, as per Riemann Lebesgue Lemma (link).

Appendix C.2. Convolution integral convergence

Let us consider $h(t) = e^{\sigma t}u(-t) + e^{-\sigma t}u(t)$ whose **first derivative is discontinuous** at $t = 0$. The second derivative of $h(t)$ given by $h_2(t)$ has a Dirac delta function $A_0\delta(t)$ where $A_0 = -2\sigma$ and its Fourier transform $H_2(\omega)$ has a constant term A_0 , corresponding to the Dirac delta function.

This means $h(t)$ is obtained by integrating $h_2(t)$ twice and its Fourier transform $H(\omega)$ has a term $-\frac{A_0}{\omega^2}$ (link) and has a **fall off rate** of $\frac{1}{\omega^2}$ as $|\omega| \rightarrow \infty$ and $\int_{-\infty}^{\infty} H(\omega) d\omega$ converges.

We see that $E_p(t) = E_0(t)e^{-\sigma t}$ where $[\frac{1}{2}e^{\frac{t}{2}} - \sum_{n=1}^{\infty} e^{-\pi n^2 e^{-2t}} e^{\frac{-t}{2}}]u(-t) + [\frac{1}{2}e^{\frac{-t}{2}} - \sum_{n=1}^{\infty} e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}}]u(t)$.

Let us consider a new function $g(t) = E_p(t)e^{-\sigma t}u(-t) + E_p(t)e^{\sigma t}u(t)$ where $g(t)$ is a real function of variable t and $u(t)$ is Heaviside unit step function and $0 < \sigma < \frac{1}{2}$. We can see that $g(t)h(t) = E_p(t)$ where $h(t) = e^{\sigma t}u(-t) + e^{-\sigma t}u(t)$.

We can see that $G(\omega), H(\omega)$ have **fall-off rate** of $\frac{1}{\omega^2}$ as $|\omega| \rightarrow \infty$ because the **first derivatives** of $g(t), h(t)$ are **discontinuous** at $t = 0$. Also, $h(t), g(t)$ are absolutely integrable functions and their Fourier Transforms are finite. Hence the convolution integral below converges to a finite value for $|\omega| \leq \infty$.

$$F(\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} G(\omega') H(\omega - \omega') d\omega' = \frac{1}{2\pi} [G(\omega) * H(\omega)] \quad (C.1)$$

Appendix C.3. **Fall off rate of Fourier Transform of functions**

Let us consider a real Fourier transformable function $P(t) = P_+(t)u(t) + P_-(t)u(-t)$ whose $(N-1)^{th}$ **derivative** is **discontinuous** at $t = 0$. The $(N)^{th}$ derivative of $P(t)$ given by $P_N(t)$ has a Dirac delta function $A_0\delta(t)$ where $A_0 = [\frac{d^{N-1}P_+(t)}{dt^{N-1}} - \frac{d^{N-1}P_-(t)}{dt^{N-1}}]_{t=0}$ and its Fourier transform $P_N(\omega)$ has a constant term A_0 , corresponding to the Dirac delta function.

This means $P(t)$ is obtained by integrating $P_N(t)$, N times and its Fourier transform $P(\omega)$ has a term $\frac{A_0}{(i\omega)^N}$ (link) and has a **fall off rate** of $\frac{1}{\omega^N}$ as $|\omega| \rightarrow \infty$.

We have shown that if the $(N-1)^{th}$ **derivative** of the function $P(t)$ is **discontinuous** at $t = 0$ then its Fourier transform $P(\omega)$ has a **fall-off rate** of $\frac{1}{\omega^N}$ as $|\omega| \rightarrow \infty$.

Appendix C.4. **Payley-Weiner theorem and Exponential Fall off rate of analytic functions.**

We know that Payley-Weiner theorem relates analytic functions and exponential decay rate of their Fourier transforms (link). Using similar arguments, we will show that the functions $E_0(t), E_p(t)$ and $x(t) = E_0(t)e^{-2\sigma t}$ and $\frac{d^{2r}x(t)}{dt^{2r}}$ have fall-off rates **at least** $\frac{1}{t^2}$ as $|t| \rightarrow \infty$ for $0 < \sigma < \frac{1}{2}$.

We know that the order of Riemann's Xi function $\xi(\frac{1}{2} + i\omega) = E_{0\omega}(\omega) = \Xi(\omega)$ is given by $O(\omega^A e^{-\frac{|\omega|\pi}{4}})$ where A is a constant^[3] (link). Hence both $E_{0\omega}(\omega)$ and $E_{p\omega}(\omega) = \xi(\frac{1}{2} + \sigma + i\omega) = E_{0\omega}(\omega - i\sigma)$ have **exponential fall-off** rate $O(\omega^A e^{-\frac{|\omega|\pi}{4}})$ as $|\omega| \rightarrow \infty$ and they are absolutely integrable and Fourier transformable, given that they are derived from an entire function $\xi(s)$.

Given that $\xi(s)$ is an entire function in the s -plane, we see that $E_{0\omega}(\omega)$ and $E_{p\omega}(\omega)$ are **analytic** functions which are infinitely differentiable which produce no discontinuities for all $|\omega| \leq \infty$ and $0 < \sigma < \frac{1}{2}$. Hence their respective **inverse Fourier transforms** $E_0(t), E_p(t)$ have fall-off rates faster than $\frac{1}{t^M}$ as $M \rightarrow \infty$, as $|t| \rightarrow \infty$ (Appendix C.3) and hence it should have **exponential fall-off** rates as $|t| \rightarrow \infty$.

We can use similar arguments to show that $x(t) = E_0(t)e^{-2\sigma t}$ and $\frac{d^{2r}x(t)}{dt^{2r}}$ have fall-off rates **at least** $\frac{1}{t^2}$ as $|t| \rightarrow \infty$, because their Fourier transforms are **analytic** functions for all $|\omega| \leq \infty$ with **exponential fall-off** rate $O(\omega^A e^{-\frac{|\omega|\pi}{4}})$ as $|\omega| \rightarrow \infty$.

Appendix D. **Derivation of entire function $\xi(s)$**

In this section, we will re-derive Riemann's Xi function $\xi(s)$ and the inverse Fourier Transform of $\xi(\frac{1}{2} + i\omega) = E_{0\omega}(\omega)$ and show the result $E_0(t) = 2 \sum_{n=1}^{\infty} [2\pi^2 n^4 e^{4t} - 3\pi n^2 e^{2t}] e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}}$.

We will use the steps in Ellison's book "Prime Numbers" pages 151-152 and re-derive the steps below^[4] (link). We start with the gamma function $\Gamma(s) = \int_0^\infty y^{s-1} e^{-y} dy$ and substitute $y = \pi n^2 x$ and derive as follows.

$$\begin{aligned}\Gamma\left(\frac{s}{2}\right) &= \int_0^\infty y^{\frac{s}{2}-1} e^{-y} dy \\ \Gamma\left(\frac{s}{2}\right)(\pi n^2)^{-\frac{s}{2}} &= \int_0^\infty x^{\frac{s}{2}-1} e^{-\pi n^2 x} dx\end{aligned}\tag{D.1}$$

For real part of s greater than 1, we can do a summation of both sides of above equation for all positive integers n and obtain as follows. We note that $\zeta(s) = \sum_{n=1}^\infty \frac{1}{n^s}$.

$$\Gamma\left(\frac{s}{2}\right)\pi^{-\frac{s}{2}}\zeta(s) = \sum_{n=1}^\infty \int_0^\infty x^{\frac{s}{2}-1} e^{-\pi n^2 x} dx\tag{D.2}$$

For real part of s (σ') greater than 1, we can use theorem of dominated convergence and interchange the order of summation and integration as follows. We use the fact that $w(x) = \sum_{n=1}^\infty e^{-\pi n^2 x}$ and

$$\sum_{n=1}^\infty \int_0^\infty |x^{\frac{s}{2}-1} e^{-\pi n^2 x}| dx = \Gamma\left(\frac{\sigma'}{2}\right)\pi^{-\frac{\sigma'}{2}}\zeta(\sigma').$$

$$\Gamma\left(\frac{s}{2}\right)\pi^{-\frac{s}{2}}\zeta(s) = \int_0^\infty x^{\frac{s}{2}-1} w(x) dx\tag{D.3}$$

For real part of s less than or equal to 1, $\zeta(s)$ **diverges**. Hence we do the following. In Eq. D.3, first we consider real part of s greater than 1 and we divide the range of integration into two parts: $(0, 1]$ and $[1, \infty)$ and make the substitution $x \rightarrow \frac{1}{x}$ in the first interval $(0, 1]$. We use **the well known theorem** $1 + 2w(x) = \frac{1}{\sqrt{x}}(1 + 2w(\frac{1}{x}))$, where $x > 0$ is real.^[4]

$$\Gamma\left(\frac{s}{2}\right)\pi^{-\frac{s}{2}}\zeta(s) = \int_1^\infty x^{\frac{s}{2}-1} w(x) dx + \int_1^\infty \frac{x^{-(\frac{s}{2}-1)}}{x^2} \frac{(1 + 2w(x))\sqrt{x} - 1}{2} dx\tag{D.4}$$

Hence we can simplify Eq. D.4 as follows.

$$\Gamma\left(\frac{s}{2}\right)\pi^{-\frac{s}{2}}\zeta(s) = \frac{1}{s(s-1)} + \int_1^\infty x^{\frac{s}{2}-1} w(x) dx + \int_1^\infty x^{\frac{-(s+1)}{2}} w(x) dx\tag{D.5}$$

We multiply above equation by $\frac{1}{2}s(s-1)$ and get

$$\xi(s) = \frac{1}{2}s(s-1)\Gamma\left(\frac{s}{2}\right)\pi^{-\frac{s}{2}}\zeta(s) = \frac{1}{2}[1 + s(s-1) \int_1^\infty (x^{\frac{s}{2}} + x^{\frac{1-s}{2}}) w(x) \frac{dx}{x}]\tag{D.6}$$

We see that $\xi(s)$ is an entire function, for all values of $Re[s]$ in the complex plane and hence we get an analytic continuation of $\xi(s)$ over the entire complex plane. We see that $\xi(s) = \xi(1-s)$ ^[4].

Appendix D.1. Derivation of $E_p(t)$ and $E_0(t)$

Given that $w(x) = \sum_{n=1}^{\infty} e^{-\pi n^2 x}$, we substitute $x = e^{2t}$, $\frac{dx}{x} = 2dt$ in Eq. D.6 and evaluate at $s = \frac{1}{2} + \sigma + i\omega$ as follows.

$$\xi\left(\frac{1}{2} + \sigma + i\omega\right) = \frac{1}{2} \left[1 + 2\left(\frac{1}{2} + \sigma + i\omega\right) \left(-\frac{1}{2} + \sigma + i\omega\right) \int_0^{\infty} \sum_{n=1}^{\infty} e^{-\pi n^2 e^{2t}} \left(e^{\frac{t}{2}} e^{\sigma t} e^{i\omega t} + e^{\frac{t}{2}} e^{-\sigma t} e^{-i\omega t} \right) dt \right] \quad (\text{D.7})$$

We can substitute $t = -t$ in the first term in above integral and simplify above equation as follows.

$$\begin{aligned} \xi\left(\frac{1}{2} + \sigma + i\omega\right) &= \frac{1}{2} + \left(-\frac{1}{4} + \sigma^2 - \omega^2 + i\omega(2\sigma)\right) \left[\int_{-\infty}^0 \sum_{n=1}^{\infty} e^{-\pi n^2 e^{-2t}} e^{\frac{-t}{2}} e^{-\sigma t} e^{-i\omega t} dt \right. \\ &\quad \left. + \int_0^{\infty} \sum_{n=1}^{\infty} e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}} e^{-\sigma t} e^{-i\omega t} dt \right] \end{aligned} \quad (\text{D.8})$$

We can write this as follows.

$$\xi\left(\frac{1}{2} + \sigma + i\omega\right) = \frac{1}{2} + \left(-\frac{1}{4} + \sigma^2 - \omega^2 + i\omega(2\sigma)\right) \int_{-\infty}^{\infty} \left[\sum_{n=1}^{\infty} e^{-\pi n^2 e^{-2t}} e^{\frac{-t}{2}} u(-t) + \sum_{n=1}^{\infty} e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}} u(t) \right] e^{-\sigma t} dt \quad (\text{D.9})$$

We define $A(t) = \left[\sum_{n=1}^{\infty} e^{-\pi n^2 e^{-2t}} e^{\frac{-t}{2}} u(-t) + \sum_{n=1}^{\infty} e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}} u(t) \right] e^{-\sigma t}$ and get the **inverse Fourier transform** of $\xi\left(\frac{1}{2} + \sigma + i\omega\right)$ in above equation given by $E_p(t)$ as follows. We use dirac delta function $\delta(t)$.

$$\begin{aligned} E_p(t) &= \frac{1}{2} \delta(t) + \left(-\frac{1}{4} + \sigma^2\right) A(t) + 2\sigma \frac{dA(t)}{dt} + \frac{d^2 A(t)}{dt^2} \\ A(t) &= \left[\sum_{n=1}^{\infty} e^{-\pi n^2 e^{-2t}} e^{\frac{-t}{2}} u(-t) + \sum_{n=1}^{\infty} e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}} u(t) \right] e^{-\sigma t} \\ \frac{dA(t)}{dt} &= \sum_{n=1}^{\infty} e^{-\pi n^2 e^{-2t}} e^{\frac{-t}{2}} e^{-\sigma t} \left[-\frac{1}{2} - \sigma + 2\pi n^2 e^{-2t} \right] u(-t) + \sum_{n=1}^{\infty} e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}} e^{-\sigma t} \left[\frac{1}{2} - \sigma - 2\pi n^2 e^{2t} \right] u(t) \\ \frac{d^2 A(t)}{dt^2} &= \sum_{n=1}^{\infty} e^{-\pi n^2 e^{-2t}} e^{\frac{-t}{2}} e^{-\sigma t} \left[-4\pi n^2 e^{-2t} + \left(-\frac{1}{2} - \sigma + 2\pi n^2 e^{-2t}\right)^2 \right] u(-t) \\ &\quad + \sum_{n=1}^{\infty} e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}} e^{-\sigma t} \left[-4\pi n^2 e^{2t} + \left(\frac{1}{2} - \sigma - 2\pi n^2 e^{2t}\right)^2 \right] u(t) + \delta(t) \left[\sum_{n=1}^{\infty} e^{-\pi n^2} (1 - 4\pi n^2) \right] \end{aligned} \quad (\text{D.10})$$

We can simplify above equation as follows.

$$\begin{aligned} \frac{d^2 A(t)}{dt^2} &= \sum_{n=1}^{\infty} e^{-\pi n^2 e^{-2t}} e^{\frac{-t}{2}} e^{-\sigma t} \left[\frac{1}{4} + \sigma^2 + \sigma + 4\pi^2 n^4 e^{-4t} - 6\pi n^2 e^{-2t} - 4\sigma \pi n^2 e^{-2t} \right] u(-t) \\ &\quad + \sum_{n=1}^{\infty} e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}} e^{-\sigma t} \left[\frac{1}{4} + \sigma^2 - \sigma + 4\pi^2 n^4 e^{4t} - 6\pi n^2 e^{2t} + 4\sigma \pi n^2 e^{2t} \right] u(t) + \delta(t) \left[\sum_{n=1}^{\infty} e^{-\pi n^2} (1 - 4\pi n^2) \right] \end{aligned} \quad (\text{D.11})$$

842 We use the fact that $F(x) = 1 + 2w(x) = \frac{1}{\sqrt{x}}(1 + 2w(\frac{1}{x}))$, where $w(x) = \sum_{n=1}^{\infty} e^{-\pi n^2 x}$ and $x > 0$ is real^[4], and we
843 take the first derivative of $F(x)$ and evaluate it at $x = 1$. We see that $\sum_{n=1}^{\infty} e^{-\pi n^2} (1 - 4\pi n^2) = -\frac{1}{2}$ (Appendix D.2)
844 and hence **dirac delta terms cancel each other** in equation below.

$$\begin{aligned}
E_p(t) &= \frac{1}{2}\delta(t) + (-\frac{1}{4} + \sigma^2)A(t) + 2\sigma\frac{dA(t)}{dt} + \frac{d^2A(t)}{dt^2} \\
E_p(t) &= \sum_{n=1}^{\infty} e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}} e^{-\sigma t} [-\frac{1}{4} + \sigma^2 + 2\sigma(\frac{1}{2} - \sigma - 2\pi n^2 e^{2t}) + \frac{1}{4} + \sigma^2 - \sigma + 4\pi^2 n^4 e^{4t} - 6\pi n^2 e^{2t} + 4\sigma\pi n^2 e^{2t}]u(t) \\
&\quad + \sum_{n=1}^{\infty} e^{-\pi n^2 e^{-2t}} e^{\frac{-t}{2}} e^{-\sigma t} [-\frac{1}{4} + \sigma^2 + 2\sigma(-\frac{1}{2} - \sigma + 2\pi n^2 e^{-2t}) \\
&\quad + \frac{1}{4} + \sigma^2 + \sigma + 4\pi^2 n^4 e^{-4t} - 6\pi n^2 e^{-2t} - 4\sigma\pi n^2 e^{-2t}]u(-t)
\end{aligned}
\tag{D.12}$$

846 We can simplify above equation as follows.

$$\begin{aligned}
E_p(t) &= [E_0(-t)u(-t) + E_0(t)u(t)]e^{-\sigma t} \\
E_0(t) &= 2 \sum_{n=1}^{\infty} [2\pi^2 n^4 e^{4t} - 3\pi n^2 e^{2t}] e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}}
\end{aligned}
\tag{D.13}$$

848 We use the fact that $E_0(t) = E_0(-t)$ because $\xi(s) = \xi(1-s)$ and hence $\xi(\frac{1}{2} + i\omega) = \xi(\frac{1}{2} - i\omega)$ when evaluated at
849 the critical line $s = \frac{1}{2} + i\omega$. This means $\xi(\frac{1}{2} + i\omega) = E_{0\omega}(\omega) = E_{0\omega}(-\omega)$ and $E_0(t) = E_0(-t)$ and we arrive at the
850 desired result for $E_p(t)$ as follows.

$$\begin{aligned}
E_0(t) &= 2 \sum_{n=1}^{\infty} [2\pi^2 n^4 e^{4t} - 3\pi n^2 e^{2t}] e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}} \\
E_p(t) &= E_0(t)e^{-\sigma t} = 2 \sum_{n=1}^{\infty} [2\pi^2 n^4 e^{4t} - 3\pi n^2 e^{2t}] e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}} e^{-\sigma t}
\end{aligned}
\tag{D.14}$$

852 *Appendix D.2. Derivation of* $\sum_{n=1}^{\infty} e^{-\pi n^2} (1 - 4\pi n^2) = -\frac{1}{2}$
853

854 In this section, we derive $\sum_{n=1}^{\infty} e^{-\pi n^2} (1 - 4\pi n^2) = -\frac{1}{2}$. We use the fact that $F(x) = 1 + 2w(x) = \frac{1}{\sqrt{x}}(1 + 2w(\frac{1}{x}))$,
855 where $w(x) = \sum_{n=1}^{\infty} e^{-\pi n^2 x}$ and $x > 0$ is real^[4], and we take the first derivative of $F(x)$ and evaluate it at $x = 1$.

$$\begin{aligned}
F(x) &= 1 + 2w(x) = \frac{1}{\sqrt{x}}(1 + 2w(\frac{1}{x})) \\
F(x) &= 1 + 2 \sum_{n=1}^{\infty} e^{-\pi n^2 x} = \frac{1}{\sqrt{x}}(1 + 2 \sum_{n=1}^{\infty} e^{-\pi n^2 \frac{1}{x}}) \\
\frac{dF(x)}{dx} &= 2 \sum_{n=1}^{\infty} (-\pi n^2) e^{-\pi n^2 x} = \frac{1}{\sqrt{x}} \sum_{n=1}^{\infty} (2\pi n^2) e^{-\pi n^2 \frac{1}{x}} (\frac{1}{x^2}) + (1 + 2 \sum_{n=1}^{\infty} e^{-\pi n^2 \frac{1}{x}}) (\frac{-1}{2}) \frac{1}{x^{\frac{3}{2}}}
\end{aligned}
\tag{D.15}$$

We evaluate the above equation at $x = 1$ and we simplify as follows.

$$\begin{aligned} \left[\frac{dF(x)}{dx}\right]_{x=1} &= 2 \sum_{n=1}^{\infty} (-\pi n^2) e^{-\pi n^2} = \sum_{n=1}^{\infty} (2\pi n^2) e^{-\pi n^2} + (1 + 2 \sum_{n=1}^{\infty} e^{-\pi n^2}) \left(\frac{-1}{2}\right) \\ &\quad \sum_{n=1}^{\infty} e^{-\pi n^2} (1 - 4\pi n^2) = -\frac{1}{2} \end{aligned}$$

(D.16)

Appendix D.3. Derivation of Result 1

In this section, we derive the result $\frac{1}{2}e^{\frac{-t}{2}} - \sum_{n=1}^{\infty} e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}} = \frac{1}{2}e^{\frac{t}{2}} - \sum_{n=1}^{\infty} e^{-\pi n^2 e^{-2t}} e^{\frac{-t}{2}}$ for $|t| < \infty$. We use the well known theorem $F(x) = 1 + 2w(x) = \frac{1}{\sqrt{x}}(1 + 2w(\frac{1}{x}))$, where $w(x) = \sum_{n=1}^{\infty} e^{-\pi n^2 x}$ and $x > 0$ is real^[4]. We substitute $x = e^{2t}$ as follows.

$$\begin{aligned} F(x) &= 1 + 2w(x) = \frac{1}{\sqrt{x}}(1 + 2w(\frac{1}{x})) \\ F(e^{2t}) &= 1 + 2 \sum_{n=1}^{\infty} e^{-\pi n^2 e^{2t}} = \frac{1}{e^t}(1 + 2 \sum_{n=1}^{\infty} e^{-\pi n^2 e^{-2t}}) \end{aligned}$$

(D.17)

We multiply above equation by $\frac{1}{2}e^{\frac{t}{2}}$ and derive the result as follows.

$$\begin{aligned} \frac{1}{2}e^{\frac{t}{2}} + \sum_{n=1}^{\infty} e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}} &= \frac{1}{2}e^{\frac{-t}{2}} + \sum_{n=1}^{\infty} e^{-\pi n^2 e^{-2t}} e^{\frac{-t}{2}} \\ \frac{1}{2}e^{\frac{-t}{2}} - \sum_{n=1}^{\infty} e^{-\pi n^2 e^{2t}} e^{\frac{t}{2}} &= \frac{1}{2}e^{\frac{t}{2}} - \sum_{n=1}^{\infty} e^{-\pi n^2 e^{-2t}} e^{\frac{-t}{2}} \end{aligned}$$

(D.18)