# Details of each benchmark line

#### Sho Iwamoto

Details (fail safe note) for each of my analyses.

$$\tilde{l}_{L} = (\tilde{e}_{L}, \tilde{\mu}_{L}, \tilde{\tau}_{L}), \qquad \tilde{\nu} = (\tilde{\nu}_{e}, \tilde{\nu}_{\mu}, \tilde{\nu}_{\tau}), \qquad \text{slepton} = (\tilde{l}_{L}, \tilde{\nu}, (\tilde{l}_{R})), \\
l = (e, \mu, \tau), \qquad \ell = (e, \mu), \dots$$

### 1. Decay chain and accepted signal categorization

We label decay chains as

$$chain-3L: \tilde{\chi}^0 \tilde{\chi}^{\pm} \to \left[ \left( \tilde{l}_L^{(*)} l^{(*)}, \tilde{\nu}^{(*)} \tilde{\nu}^{(*)}; 1/12 \text{ each} \right) \left( \tilde{\nu} l, \tilde{l}_L \nu; 1/6 \text{ each} \right) \right] \otimes (\text{slep} \to \text{lep} \tilde{\chi}_{LSP}^0; 100\%)$$
 (1.1)

chain-LLT: 
$$\tilde{\chi}^0 \tilde{\chi}^{\pm} \rightarrow \left[ \left( \tilde{l}_{\rm R}^{(*)} l^{(*)}; 1/6 \text{ each} \right) \left( \tilde{\tau}_{\rm R} \nu; 100\% \right) \right] \otimes (\text{slep} \rightarrow \text{lep} \tilde{\chi}_{\rm LSP}^0; 100\%)$$
 (1.2)

chain-WZ: 
$$\tilde{\chi}^0 \tilde{\chi}^{\pm} \to (Z \tilde{\chi}_{\rm LSP}^0)(W^{\pm} \tilde{\chi}_{\rm LSP}^0)$$
 (possibly virtual), (1.3)

chain-WH: 
$$\tilde{\chi}^0 \tilde{\chi}^{\pm} \to (H \tilde{\chi}^0_{\rm LSP})(W^{\pm} \tilde{\chi}^0_{\rm LSP})$$
 (possibly virtual), (1.4)

and selection criteria as\*1

$$sig-3L: (3\ell)_{SFOS} \tag{1.5}$$

$$sig-LLT: (2\ell)_{SFOS} + \tau_{had}$$
 (1.6)

$$sig-SS2L: (2\ell)_{SS} + \tau_{had}$$
 (1.7)

ATLAS1803 analyzes chain-3L (x = 0.5) by sig-3L ("a statistical combination of the five SR3-slep regions."); note that the lepton  $\ell$  may as well come from leptonic tau decays.

CMS1709 has several analyses; relevant ones are

- chain-3L (x = 0.5) by sig-3L,
- chain-3L (x = 0.05, 0.95) by sig-SS2L×3L (×: statistical combination),
- chain-LLT (x = 0.05, 0.5, 0.95) by sig-3L×LLT,
- chain-WZ by sig-3L,
- chain-WH by combination of various SRs.

#### 2. tab1-0.50

This line is characterized by

$$M_2 = \mu = 2M_1, \quad x = \frac{m_{\tilde{l}_L} - m_{\tilde{\chi}_1^0}}{m_{\tilde{\chi}_2^0} - m_{\tilde{\chi}_1^0}} = 0.5, \quad \tan \beta = 40, \quad \tilde{l}_R, \tilde{q}, \text{ heavy-Higgs: decoupled.}$$
 (2.1)

The mass spectrum is shown below, and we use  $m_{\tilde{\chi}_1^{\pm}}$  (in GeV) as the label of each model point.

P<300 is not considered in our analysis, as neither by ATLAS. P150 has  $\tilde{\nu}$ -LSP. In P $\geq$ 300,  $\tilde{\chi}_{2}^{0}\tilde{\chi}_{1}^{\pm}$  mainly produces chain-3L. Also,  $\tilde{\chi}_{3,4}^{0}$  and  $\tilde{\chi}_{2}^{\pm}$  decays similarly. We safely ignore  $\tilde{\chi}_{3}^{0}$ because of non-degeneracy and, as it has less  $\tilde{W}$ -component, smaller production rate. A degenerate pair

<sup>\*1 &</sup>quot;SF", "OS", "SS" are respectively for same flavor, opposite sign, and same sign.  $(3\ell)_{SFOS}$  means that two of them form a SFOS pair and the other is arbitrary. Particles are "hard" (not soft) unless noted as such. Events with extra leptons are vetoed in some analyses, but I do not care those vetoes as we are anyway not interested in vetoed signatures.

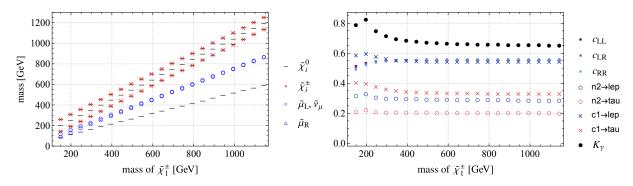


Figure 1: Mass spectrum and c-factors of tab1-0.50 benchmark line. The models are generated with  $M_2 = 200, 250, \dots, 1200 \,\text{GeV}$ , while  $m_{\tilde{\chi}_1^{\pm}}$  is used to label.

 $\tilde{\chi}_4^0 \tilde{\chi}_2^{\pm}$  may serve as the chain-3L target. At this first stage I will ignore it, but it may be interesting to include the effect (but how?).

The c-factors are similar and thus we use the factor

$$K_{\sigma} = \text{mean}(c_{\text{LL}}, c_{\text{LR}}, c_{\text{RR}}) \tag{2.2}$$

to compensate the difference in production cross section from pure-wino (ATLAS 1803) case. We thus estimate the production cross section as

$$\sigma(pp \to \tilde{\chi}_2^0 \tilde{\chi}_1^{\pm}) \approx K_{\text{sigma}} \cdot \sigma(pp \to \tilde{W}^{\pm} \tilde{W}^3),$$
 (2.3)

where the pure-wino production cross section  $\sigma(pp \to \tilde{W}^{\pm}\tilde{W}^{3})$  is taken from LHCSUSYXSWG<sup>\*2</sup>. The decays of  $\tilde{\chi}_{2}^{0}\tilde{\chi}_{1}^{\pm}$  does not equally branch because of Yukawas and slepton mass differences;

$$\tilde{\chi}_{2}^{0}: \quad \tau > e = \mu \gtrsim \nu_{e} = \nu_{\mu} = \nu_{\tau} \quad \text{(also to } Z, H),$$
 (2.4)

$$\tilde{\chi}_1^{\pm}: \quad \tau > e = \mu \gtrsim \nu_{\tau} \gtrsim \nu_e = \nu_{\mu} \quad \text{(also to } W^{\pm}),$$

while in ATLAS analysis they are universal since they consider wino-like  $\tilde{\chi}_2^0 \tilde{\chi}_1^{\pm}$  with degenerate sleptons. As seen from Eq. (A.1), this non-universality is compensated by

$$K_{\Gamma} = \frac{1}{0.269} \left[ p(\tilde{\ell}_{L}, \tilde{\nu}_{\ell} | \tilde{\chi}_{1}^{\pm}) + \frac{1}{2} p(\ell | \tau) p(\tilde{\tau}_{L}, \tilde{\nu}_{\tau} | \tilde{\chi}_{1}^{\pm}) \right] \left[ p(\tilde{\ell}_{L}^{(*)} | \tilde{\chi}_{2}^{0}) + \frac{1}{2} p(\ell | \tau)^{2} p(\tilde{\tau}_{L}^{(*)} | \tilde{\chi}_{2}^{0}) \right]$$
(2.6)

(the notation p(daughters|mothers) should be understood). Here one should note that we ignore the kinematics difference due to the origin of lepton (whether  $\tau$ ,  $\tilde{\chi}$ , or sleptons) as well as slepton non-degeneracy. We then compare our cross section with  $\sigma_{\text{UL}}$ ; the points with

$$\frac{K_{\sigma}K_{\Gamma} \cdot \sigma(pp \to \tilde{W}^{\pm}\tilde{W}^{3})}{\sigma_{\text{UL}}} \ge 1 \tag{2.7}$$

are excluded.\*3

The results against ATLAS1803 is shown in Fig. 2, where the black dots correspond to  $K_{\sigma}K_{\Gamma}\sigma(\text{Wino})$ . It shows that the ATLAS1803 analysis nearly excludes below  $\sim 860\,\text{GeV}$ . The wiggles in  $\sigma_{\text{UL}}$  is due to interpolation of the  $\sigma_{\text{UL}}$ -grid ATLAS provides, for which logarithmic interpolation (i.e., linear interpolation on the function  $\log \sigma_{\text{UL}}(m_{\tilde{\chi}_1^\pm}, m_{\tilde{\chi}_1^0})$ ) is used.

# A. N2C1 acceptance

The acceptance of N2C1-to-slep events in relevant selection criteria is summarized here.

 $<sup>^{*2} \</sup>mathtt{https://twiki.cern.ch/twiki/bin/view/LHCPhysics/SUSYCrossSections13TeVn2x1wino}$ 

 $<sup>^{*3}</sup>$  ATLAS 1803: https://doi.org/10.17182/hepdata.81996.v1/t80

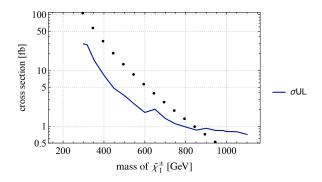


Figure 2: ATLAS 1803 analysis (3l-slep) on tab1-0.50 benchmark line.

Noting that sig-3L requires  $3\ell$  with SFOS pair, the probability that chain-3L with x=0.5 produces signatures falling in sig-3L category is estimated by

$$A(\text{sig-3L}|\text{chain-3L}_{x=0.5}) \approx \left[ p(\tilde{\ell}_{L}, \tilde{\nu}_{\ell} | \tilde{\chi}_{1}^{\pm}) + p(\ell | \tau) p(\tilde{\tau}_{L}, \tilde{\nu}_{\tau} | \tilde{\chi}_{1}^{\pm}) \right] \left[ p(\tilde{\ell}_{L}^{(*)} | \tilde{\chi}_{2}^{0}) + \frac{1}{2} p(\ell | \tau)^{2} p(\tilde{\tau}_{L}^{(*)} | \tilde{\chi}_{2}^{0}) \right]$$

$$= \left( \frac{4}{6} + \frac{2}{6} p(\ell | \tau) \right) \left( \frac{4}{12} + \frac{1}{2} p(\ell | \tau)^{2} \frac{2}{12} \right) = 0.269,$$
(A.1)

where  $p(\ell|\tau) \simeq 0.352$  is the leptonic decay rate of  $\tau$ .

## References