

Nuclear Engineering 150 – Discussion Section
Team Exercise Solutions #12

Problem 1

Consider a critical sphere composed of a homogeneous multiplying medium, and surrounded by a reflecting shell of non-multiplying material with a thickness equal to half the radius of the sphere. Outside of the reflector is vacuum.

- Write the diffusion equation and boundary conditions describing this system.
- Calculate the flux in both regions that are not vacuum.
- Determine the reflector savings.

Problem 1 Solution

- In a symmetric spherical system, the one-speed, steady-state diffusion equation is

$$-D\nabla^2\phi(r) + \Sigma_a(r)\phi(r) = \nu\Sigma_f(r)\phi(r) + S(r).$$

For any arbitrary function, $f(r, \theta, \varphi)$, the full Laplacian in spherical coordinates is

$$\nabla^2 f(r, \theta, \varphi) = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial f}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial f}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 f}{\partial \varphi^2}.$$

If the function $f(r, \theta, \varphi)$ is radially symmetric then $f(r, \theta, \varphi) = f(r)$, and there will be no variation in the θ or φ directions ($\frac{\partial f}{\partial \theta} = 0$ and $\frac{\partial f}{\partial \varphi} = 0$). The second and third terms vanish, and

$$\nabla^2 f(r) = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial f}{\partial r} \right).$$

We substitute this back into our one-speed, steady-state diffusion equation, using $\phi(r)$ instead of $f(r)$.

$$-D \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial \phi(r)}{\partial r} \right) + \Sigma_a(r)\phi(r) = \nu\Sigma_f(r)\phi(r) + S(r).$$

Finally, we recognize that our problem has two regions. There is no external source in either, so $S(r) = 0$. Since the regions are homogeneous, the diffusion coefficients and cross sections are constant in each: D_C , $\Sigma_{a,C}$, ν , and Σ_f in the core; D_R and $\Sigma_{a,R}$ in the reflector. We will define the outer radius of the core as a , and the outer radius of the reflector as $\frac{3a}{2}$. Then, our diffusion equations for both regions are:

$$\begin{aligned} -D_C \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial \phi(r)}{\partial r} \right) + \Sigma_{a,C}\phi(r) &= \nu\Sigma_f\phi(r), \quad 0 < r < a \quad (\text{core}) \\ -D_R \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial \phi(r)}{\partial r} \right) + \Sigma_{a,R}\phi(r) &= 0, \quad a < r < \frac{3a}{2} \quad (\text{reflector}). \end{aligned}$$

For our final solution, we will make two more substitutions, $\kappa_C^2 = \frac{\nu\Sigma_f - \Sigma_{a,C}}{D_C}$ and $L_R^2 = \frac{D_R}{\Sigma_{a,R}}$, to get

$$\begin{aligned} \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial \phi(r)}{\partial r} \right) - \kappa_C^2 \phi(r) &= 0, \quad 0 < r < a \quad (\text{core}) \\ \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial \phi(r)}{\partial r} \right) - \frac{1}{L_R^2} \phi(r) &= 0, \quad a < r < \frac{3a}{2} \quad (\text{reflector}). \end{aligned}$$

There are also four boundary conditions that accompany these equations. A vacuum boundary condition (where we will use the extrapolated distance d to indicate where the flux goes to zero), a finiteness condition, and two interface conditions (both continuous current and flux at the boundary).

$$\begin{aligned}
(1) \quad & \phi_R \left(\frac{3a}{2} + d \right) = 0 && \text{(vacuum)} \\
(2) \quad & \phi(r) < \infty && \text{(finiteness)} \\
(3) \quad & J_C(a) = J_R(a) && \text{(continuous current)} \\
(4) \quad & \phi_C(a) = \phi_R(a) && \text{(continuous flux)}
\end{aligned}$$

- b) Since the diffusion equation takes a different form in each region, the solutions describing the flux shapes in those regions are also different. In the core, the flux has the functional form

$$\phi_C(r) = \frac{A_C}{r} \sin(\kappa_C r) + \frac{B_C}{r} \cos(\kappa_C r), \quad 0 < r < a,$$

and in the reflector,

$$\phi_R(r) = \frac{A_R}{r} e^{-r/L_R} + \frac{B_R}{r} e^{r/L_R}, \quad a < r < \frac{3a}{2}.$$

The second term in the flux solution for the core goes to infinity as r goes to zero ($\frac{1}{r} \rightarrow \infty$ and $\cos(\kappa r) \rightarrow 1$). For the finiteness condition to be satisfied, we must set $B_C = 0$.

c)

Problem 2

Consider a bare sphere composed of a homogeneous multiplying medium.

- a) Give the steady-state, continuous energy diffusion equation. Assume that the diffusion coefficient (D) and the average neutrons produced from fission (ν) are constant for all energies.
- b) Derive the multigroup equation corresponding to the case where there are three energy groups. Assume that there is no upscattering, all groups are directly coupled (no inscattering), and fission is only induced by the slowest group while only producing neutrons in the fastest group.
- c) Write the multigroup equation you found as a matrix-equation.

Problem 2 Solution

- a)
- b)
- c)