

Design and Construction of the MicroBooNE Detector

The MicroBooNE Collaboration

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ABSTRACT: This paper describes the design and construction of the MicroBooNE liquid argon time projection chamber and associated systems. MicroBooNE is the first phase of the Short Baseline Neutrino program, located at Fermilab, and will utilize the capabilities of liquid argon detectors to examine a rich assortment of physics topics. In this document details of design specifications, assembly procedures, and acceptance tests are reported.

KEYWORDS: Time projection chambers; Noble liquid detectors; Neutrino detectors

Dedicated to W. J. Willis

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1 Introduction and Physics Motivation

2 The **Micro Booster Neutrino Experiment**, henceforth referred to as MicroBooNE, is the first large
3 (~ 100 tons) Liquid Argon Time Projection Chamber (LArTPC) to operate in the United States.
4 MicroBooNE is the latest among a family of detectors that exploit the potential of liquified noble
5 gases employed as the detection medium for neutrino interactions. These detectors combine the
6 advantages of high spatial resolution and calorimetry with the potential to scale to very large
7 volumes.

8 The potential of liquified noble gases to be employed as the detection media in particle
9 detectors with high spatial resolution was recognized in the 1960's [1]. Large calorimeters for
10 the measurement of particle energy were then realized by using cryogenic noble liquids as active
11 components [2]. The LArTPC concept was then proposed at CERN by C. Rubbia in 1977 [3] and
12 developed within the ICARUS program [4–6] culminating in the realization of the ICARUS T600
13 detector [7]

14 In the US, the first LArTPC to be successfully operated in a neutrino beam was the ArgoNeuT
15 detector [8] with ~ 170 kg mass. ArgoNeuT, which performed a series of detailed studies on the
16 interaction of medium-energy neutrinos [9–12], is the direct predecessor of MicroBooNE, which
17 is, in turn, the progenitor of the SBND [13] and DUNE [14], experiments.

18 MicroBooNE's principal physics goal is to address short baseline neutrino oscillations, specifically
19 the MiniBooNE low energy excess result [15], at Fermi National Accelerator Laboratory
20 (Fermilab). MicroBooNE will be exposed to the 0.5–2 GeV on-axis Booster Neutrino Beam (BNB),
21 the same beam used for the MiniBooNE experiment, at a ~ 500 m baseline, also the same as Mini-
22 BooNE, but by employing the LArTPC technology to efficiently separate signal electrons from
23 background photons.

24 In addition to MicroBooNE's signature oscillation analyses, a suite of precision cross-section
25 measurements will be performed, critical both for future LArTPC oscillation experiments and for
26 what can be learned about neutrino interactions in general. In the neutrino interactions from the
27 BNB, multiple interaction processes (quasi-elastic, resonances, deep inelastic scattering) are pos-
28 sible, and complicated nuclear effects in neutrino interactions on argon result in a variety of final
29 states. These can range from the emission of several nucleons to more complex topologies with
30 multiple pions, all in addition to the leading lepton in charged-current events. The LArTPC technol-
31 ogy employed by MicroBooNE is particularly well suited for complicated topologies because of its
32 excellent particle identification capability and calorimetric energy reconstruction down to very low
33 detection thresholds. MicroBooNE will also be supernova live and capable of detecting a galactic
34 supernova, and prepare the analysis of the search for proton decays. For a proton decay search, the
35 experiment will not have a competitive sensitivity due to too small target mass. However, this will
36 be an important proof of principle measurement towards future searches in larger detectors.

37 MicroBooNE is located 470 m from the BNB production target and 600 m from the NuMI
38 production target. MicroBooNE began operations in late 2015 for an anticipated ~ 3 year data taking
39 run. The next phase of the Short Baseline Neutrino (SBN) program at Fermilab includes continued
40 operation of MicroBooNE along with the SBND and ICARUS experiments, located at 110 m and
41 600 m respectively from the BNB target. SBND and ICARUS are now under construction [16] with
42 operations to begin in ~ 2018 . MicroBooNE will definitively address whether or not the MiniBooNE

1 low energy excess in neutrino mode is due to electrons or photons in its initial run. SBND will look
2 for this low energy excess at the near location and ICARUS, with its larger mass, will enable the
3 three detector program to cover the entire LSND allowed region in neutrino parameter space to 5σ
4 through ν_e appearance.

5 More information on the LArTPC technology can be found in existing reviews (see, e.g., [17]
6 and references therein).

7 **Organization of Document** This document describes the design, construction, and technical
8 details of MicroBooNE. Section 2 gives a brief review of the LArTPC technique and its imple-
9 mentation in MicroBooNE. Section 3 describes the cryogenic and purification systems which are
10 required for maintaining a stable volume of highly purified liquid argon. The LArTPC detector de-
11 scribed in section 4 is the centerpiece of the experiment, providing fine-grained images of neutrino
12 interactions. A light collection system, described in section 5, provides precise timing information
13 by viewing the detector volume and recording signals from the scintillation light that is produced
14 therein. Signals from the light collection system and from the LArTPC are amplified, sampled, and
15 recorded by a custom-designed electronic and readout system, as described in section 6. Section 7
16 describes the auxiliary instrumentation that monitor and control the detector and all of its associated
17 systems, as well as provide an electrically quiet environment for the experiment to operate. Finally,
18 one of the main calibration sources for the experiment is an ultraviolet laser system, described in
19 section 8, that provides the capability to map out the performance of the LArTPC detector. A
20 cosmic ray tagger system, under construction at the time of the writing of this paper, will surround
21 the detector to improve cosmic ray identification and rejection. This system will be described in a
22 subsequent publication.



Figure 1. Aerial diagram showing location of MicroBooNE along the BNB (orange dashed line) at Fermilab.

2 Experiment Overview

The MicroBooNE detector at Fermilab in Batavia, Illinois is sited on axis on the BNB, 470 m downstream from the neutrino production target. The BNB delivers a beam of predominantly muon neutrinos with energies peaking at 700 MeV produced primarily from pion decays[18]. MicroBooNE is also exposed to an off-axis component of the NuMI beam [19] produced from pion and kaon decays with average neutrino energies of about 2 GeV and 0.5 GeV respectively. MicroBooNE is located about 600 m downstream from the NuMI neutrino production target. The characteristics of the BNB beamline are well measured and understood from many years of data taking and analysis on the MiniBooNE experiment [20], which operated directly downstream of the MicroBooNE location. Figure 1 shows the arrangement of MicroBooNE with respect to the BNB beamline at Fermilab. The physics program of MicroBooNE will utilize both BNB and NuMI samples. MicroBooNE will also collect data that is out-of-time with either beam, which will be useful for developing analyses (e.g. proton decay searches) relevant for next-generation detectors.

2.1 The MicroBooNE LArTPC

Charged particles traversing a volume of highly-purified liquid argon leave trails of ionization electrons in their wake and also create prompt vacuum ultraviolet (VUV) scintillation photons. In a LArTPC the liquid argon is highly purified and the ionization trails are transported practically undistorted over distances of the order of meters [21] under the influence of a uniform electric field in the detector volume, until they reach anode planes located along one side of the active volume. The uniform electric field is created by introducing voltage onto a cathode plane and gradually stepping that voltage down in magnitude across a field cage, which is formed from a series of equipotential rings surrounding the drift volume. The anode plane is arranged parallel to the cathode plane, and in MicroBooNE, parallel to the beam direction. There are three planes comprised of sense wires with a characteristic pitch, held at a predetermined bias voltage, that continuously sense the signals induced by the ionization electrons drifting towards them. The electrostatic potentials of the sequence of anode planes allow ionization electrons to pass undisturbed by the first two planes before ultimately ending their trajectory on a wire in the last plane. The drifting ionization thus

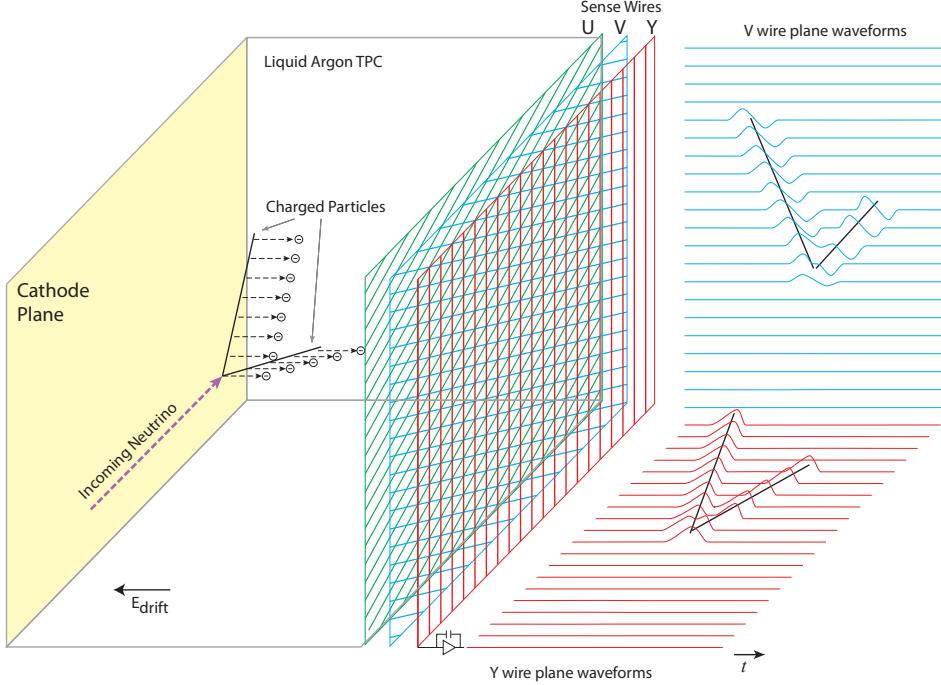


Figure 2. Operational principle of the MicroBooNE LArTPC.

1 induces signals on the first planes (referred to as Induction planes) and directly contributes to the
 2 signals in the final plane (referred to as the Collection plane). Figure 2 depicts the arrangement of
 3 the MicroBooNE LArTPC.

4 The charged particle trajectory is reconstructed using the known positions of the anode plane
 5 wires and the recorded drift time of the ionization. The drift time is the difference between the
 6 arrival times of ionization signals on the wires and the time the interaction took place in the detector
 7 (t_0) which is provided by an accelerator clock synced to the beam (e.g. - BNB or NuMI) or from
 8 a trigger provided by the light collection system. The characteristics of the waveforms observed
 9 by each wire provide a measure of the energy deposition of the traversing particles near that wire,
 10 which, when taken as a whole for each contained particle's trajectory, allow for determination of
 11 momentum and particle identity.

12 The scintillation photons are detected by a light collection system that is immersed in the liquid
 13 argon and faces into the detector volume. This system provides signals that can establish the event
 14 t_0 and supplies trigger information to an electronic readout system. The light collection system
 15 signals are vital in distinguishing detector activity that is in-time with the beam (e.g. originating
 16 from beam interactions) from that which is out-of-time (e.g. likely not originating from beam-
 17 related interactions), benefiting event reconstruction. Spatial information transverse to the drift
 18 direction can also be inferred from the light collection signals, further aiding in the reconstruction

1 of the activity occurring inside the detector.

2 The choice of liquid argon as both the neutrino target and detector medium in LArTPCs is well
3 motivated, though not without introducing unique considerations for experimental design. Liquid
4 argon as a target for neutrinos is attractive due to its density, allowing a more compact detector with
5 a substantial boost in event rate over a comparable detector using less dense media, such as water.
6 A tradeoff to this aspect is the short radiation length (14 cm) and that the complicated structure
7 of the argon atom, relative to simpler targets such as hydrogen or helium, does introduce nuclear
8 effects that must be considered during data analysis. The cryogenic temperatures at which the noble
9 elements are in the liquid phase also introduces the need for additional design considerations to
10 ensure stable and safe operations.

11 Table 1 lists some of the properties of liquid argon that are salient for LArTPC design. The noble
12 liquids are all characterized by their excellent dielectric properties, able to maintain high voltages
13 without suffering electrical breakdown. The long drift lengths over which ionization electrons must
14 travel in large LArTPC detectors requires the presence of a uniform electric field over that distance.
15 This is achieved via the introduction of extremely high voltage ($O(100$ kV magnitude) onto the
16 LArTPC cathode; this voltage is suitably maintained due to the dielectric properties of liquid argon
17 and by ensuring the detector components have no sharp edges. The noble liquids are all very bright
18 scintillators, with wavelengths deep in the UV, and also produce copious amounts of ionization.
19 Both the scintillation and the ionization signals are necessary for a robust accounting of the activity
20 occurring inside the LArTPC. Finally, the desire to build LArTPCs on increasingly larger scales,
21 such as those necessary in neutrino experiments, is bolstered by the abundance (1% of atmosphere)
22 and low cost of argon.

Table 1. Selected properties of liquid argon.

Property	Value	Reference
Atomic number	18	
Atomic weight [g/mol]	39.95	
Boiling point [K] @ 1 atm	87.3	[22]
Density [g/cm ³] @ 1 atm	1.4	[22]
Dielectric constant	1.505	[23]
Radiation length [cm]	14.0	[24]
Molière radius [cm]	10.0	[24]
W-value for ionization [eV/pair]	23.6	[25, 26]
Minimum specific energy loss [MeV/cm]	2.12	[24]
Electron transverse diffusion coef. [cm ² /s]	13	[5, 27, 28]
Electron longitudinal diffusion coef. [cm ² /s]	5	[5, 29]

23 The successful implementation of the LArTPC technique depends critically on several factors.
24 The liquid argon must be purified of any electronegative contaminants, such as water or oxygen, to
25 accommodate the very long drift path of ionization through a MicroBooNE-sized LArTPC without
26 significant charge loss. The signals that the ionization electrons create on the anode wires are very
27 small, requiring low-noise electronics to discern signal pulses. The MicroBooNE collaboration
28 has designed and constructed an experiment that addresses all of these considerations, providing

¹ critical technological development for the next generation of LArTPC experiments to build upon.

² **2.2 MicroBooNE LArTPC Implementation**

³ MicroBooNE’s LArTPC active volume, which is defined as the volume immediately within the
⁴ confines of the LArTPC field cage, is a rectangular liquid argon volume with dimensions 2.3 m
⁵ vertical x 2.6 m horizontal x 10.4 m along the beam direction. This is the maximum volume that
⁶ can be used for physics analyses. The cathode (anode) defines the beam-left (beam-right) side of
⁷ the active volume. The end of the LArTPC that the beam first encounters is referred to as the
⁸ “upstream” end, while the opposite end is referred to as “downstream.” Anode plane-to-plane
⁹ spacing is 3 mm, and each plane has 3 mm wire pitch. The induction plane wires are oriented at
¹⁰ $\pm 60^\circ$ relative to vertical, and the collection plane has vertically oriented wires. Field cage loops
¹¹ are employed to maintain uniformity of the electric field across the entire width of the detector, and
¹² these loops also act to define the top, bottom, upstream, and downstream sides of the active volume.

¹³ MicroBooNE uses a right-handed Cartesian coordinate system, with the origin defined to be
¹⁴ located on the upstream face of the LArTPC, centered halfway up the vertical height of the active
¹⁵ volume and horizontally centered on the anode plane closest to the cathode (the innermost anode
¹⁶ plane). In this system, x ranges from 0.0 m at the innermost anode plane to +2.6 m at the cathode, y
¹⁷ ranges from -1.15 m on the bottom of the active volume to +1.15 m at the top of the active volume,
¹⁸ and z ranges from 0.0 m at the upstream end of the active volume to +10.4 m at the downstream
¹⁹ end.

²⁰ The light collection system, which is an array of photomultiplier tubes (PMTs) and lightguide
²¹ paddles, is located directly behind the anode planes on beam-right, facing the detector volume
²² through the anode planes. The LArTPC and light collection system are immersed in liquid argon
²³ contained within a single-walled cryostat with a 170 ton capacity. Electronics mounted directly
²⁴ on the LArTPC amplify the signals on the wires, which are then passed out of the cryostat for
²⁵ further processing and storage on disk. Table 2 lists the primary detector design parameters of
²⁶ MicroBooNE, and figure 3 shows a schematic of the cross section of the detector. Details of
²⁷ these design parameters and construction of all detector systems will be provided in the subsequent
²⁸ sections.



Figure 3. Schematic of the cross section of the MicroBooNE LArTPC. In this view, the beam would be directed out of the page (in the z direction).

Table 2. Primary detector design parameters for MicroBooNE.

Parameter	Value
LArTPC Dimensions	2.325 m vertically 2.560 m horizontally 10.368 m longitudinally
LArTPC argon mass	90 tons
Total Number of Wires	8256
Drift field	500 V/cm
Light collection	30 200 mm (8 in) diameter PMTs 4 lightguide paddles
Cryostat liquid argon capacity	170 tons
Operating temperature	87 K
Operating pressure	0.14 bar

¹ **3 Cryogenic System**

² The use of large quantities of highly-purified liquid argon as a detector medium in MicroBooNE
³ requires a sophisticated cryogenic infrastructure that can maintain stable operations for years at a
⁴ time. Not only must the purity of the liquid argon be maintained, but the pressure and temperature
⁵ gradients within the LArTPC active volume must be tightly controlled as the drift velocity of
⁶ electrons is dependent on these quantities. A customized cryogenic system that serves these
⁷ purposes has been built, and the requirements for this system are shown in table 3.

Table 3. Primary design requirements for MicroBooNE cryogenic and purification systems.

Parameter	Value	Motivation
Argon purity	<100 ppt O ₂	MIP identification at longest drift
Argon purity	<2 ppm N ₂	Scintillation light output
LAr Temperature gradient	<0.1 K	Drift-velocity uniformity
LAr recirculation rate	1 volume change/day	Maintain purity
Cryostat heat load	<15 W/m ²	Minimize convection currents and bubbles
Cryogenic capacity	10 kW	Capacity to deal with expected heat load
Cryostat maximum pressure	2.1 bar	Determines relief sizing

⁸ The MicroBooNE cryogenic system is represented in figure 4. The central component of
⁹ the system is a cryostat that houses the complete LArTPC and light-collection detector systems.
¹⁰ The cryostat is supported by three major subsystems: the argon purification system, the nitrogen
¹¹ refrigeration system, and the controls and monitoring system. These systems each represent the next
¹² generation of LArTPC cryogenic system after the Liquid Argon Purity Demonstrator (LAPD) [30]
¹³ and make considerable use of the expertise gained during the design and implementation of that
¹⁴ apparatus.

¹⁵ **3.1 Cryostat Design Overview**

¹⁶ Three major components make up the MicroBooNE cryostat: a type 304 stainless steel vessel to
¹⁷ contain the liquid argon and all the active detector elements, front and rear supports to carry the
¹⁸ weight of the fully loaded cryostat, and foam insulation covering the cryostat outer surfaces. The
¹⁹ foam insulation serves to reduce heat input from the ambient environment to a sufficiently low
²⁰ level to prevent large temperature gradients and boiling of the liquid argon. The cryostat and the
²¹ cryogenic systems are designed to achieve the high-purity liquid argon needed to allow ionization
²² electrons to drift to the anode wires with low probability of capture, and the high degree of thermal
²³ homogeneity needed to avoid the introduction of non-constant drift velocities for the ionization
²⁴ electrons. Finally, the outer diameter of the vessel is designed to be the maximum standard size for
²⁵ over-the-road transport.

²⁶ Ionization electrons must not be significantly attenuated, via attachment to electronegative
²⁷ contaminants in the liquid argon, as they drift up to 2.5 m across the active volume. This dictates
²⁸ that the argon be kept free of electronegative contaminants to the level of 100 parts-per-trillion (ppt).
²⁹ The cryostat is designed to minimize outgassing (desorption) and to avoid leakage and diffusion
³⁰ of air into the system. This requirement imposes strict quality assurance demands on all welds

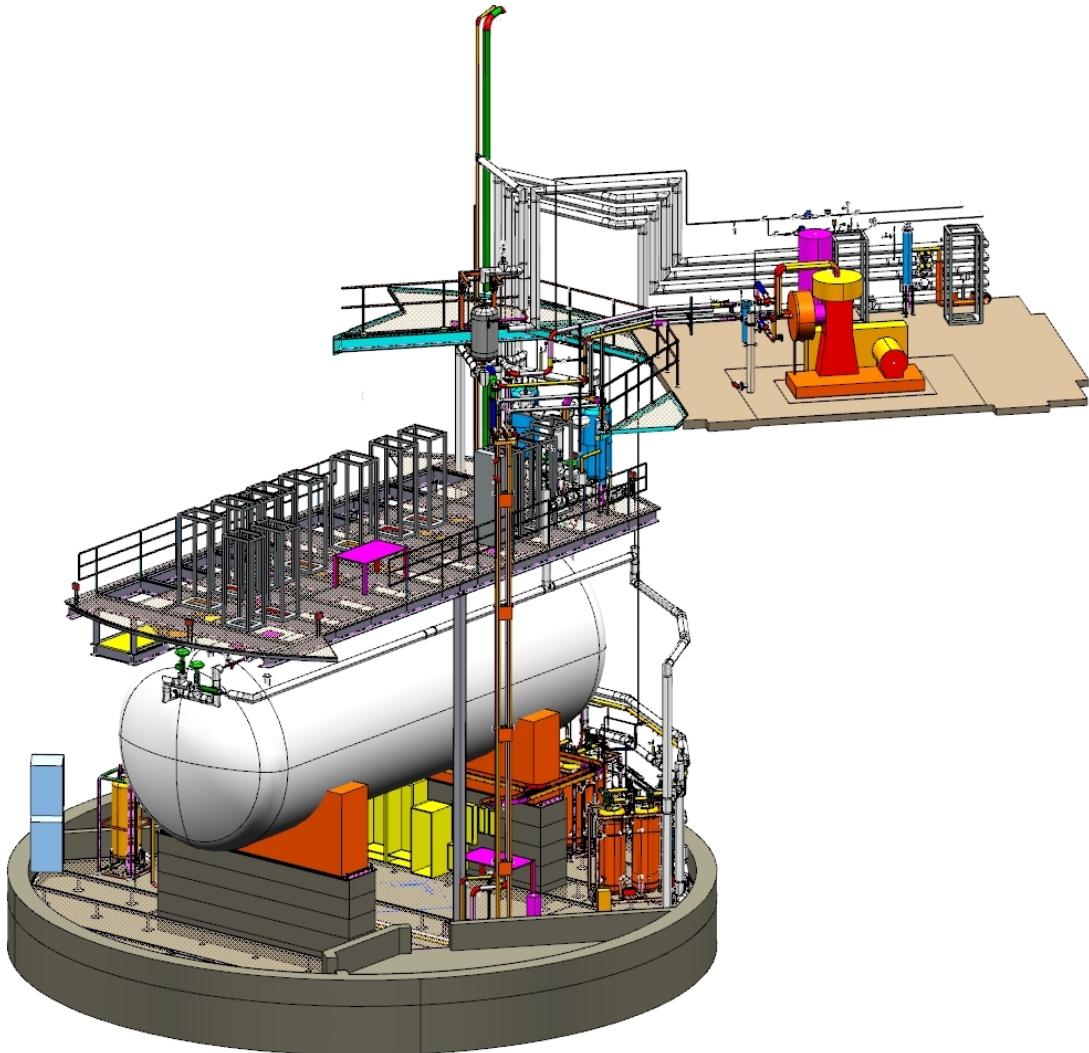


Figure 4. Three-dimensional renderings of the MicroBooNE cryogenic system installed at LArTF.

for penetrations into the cryostat and on cleaning and handling procedures for the finished vessel.
Achieving the required level of purity is accomplished with a purification system, described in section 3.2, that removes electronegative contaminants from the argon during the initial fill and those introduced over time by leaks and outgassing of system components.

The electron drift velocity ($v_d = 1600$ m/s at an electric field of 500 V/cm, with a liquid argon temperature dependence $\Delta v_d/v_d = -0.019\Delta T$) must remain constant in magnitude and direction throughout the active liquid argon volume to avoid distortion of the mapping of drift time into the position along the drift (\hat{x}) direction. This requirement limits the allowable temperature variations of the liquid argon to less than 0.1 K and the laminar and turbulent flow rate of liquid argon to less than 1 m/s. These requirements limit fractional errors in velocity, and therefore in the drift-coordinate determination, to be less than 0.1%. The constraints on constancy of drift velocity affect the design by imposing limits on the acceptable heat flux through the insulation.

The cryostat is constructed to the latest American Society of Mechanical Engineers (ASME)

boiler code requirements [31] and features a single-walled construction, cylindrical shape, and domed caps closing each end, as shown in figure 5. One end was removed for installation of the active detectors, welded back in place upon completion of that task, and then recertified to the ASME code requirements.

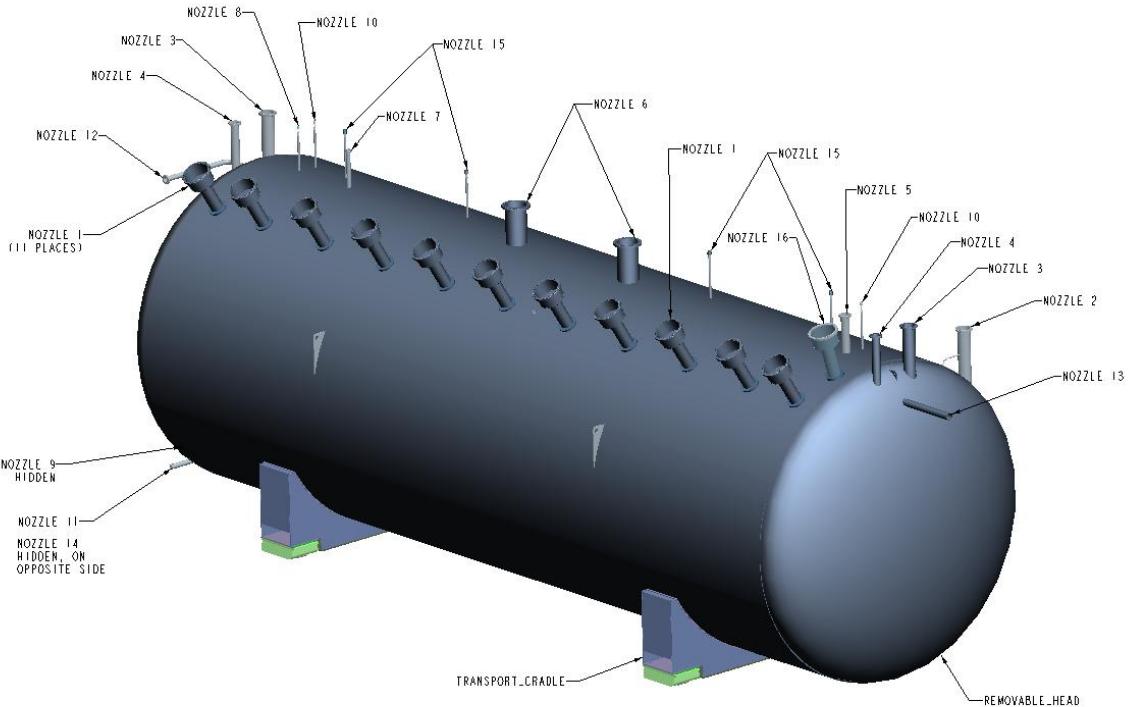


Figure 5. MicroBooNE cryostat with nozzle penetrations labeled.

Upon installation of the sealed cryostat in its final location at the Liquid Argon Test Facility (LArTF), 41 cm of spray-on, closed cell, Polyurethane insulation was applied to the exterior of the cryostat, as shown in figure 6. At LArTF, to avoid ground loops that could interfere with the LArTPC signals, the cryostat vessel is grounded in only one place, allowing it to act as a Faraday cage. This grounding scheme is explained further in section 7.1.

The vessel surface has 34 nozzle penetrations for cryogenic and electrical services, detailed in table 4. All nozzles are sealed with feedthroughs, flanges, or pipes that are suitable for operation at the nominal pressure and temperature of the cryostat.

3.2 Liquid Argon Purification Subsystem

The heart of the cryogenic system is the liquid argon purification subsystem. The primary requirement of this subsystem is to keep the level of electronegative contamination to below 100 ppt of oxygen-equivalent contaminants. This requirement was determined by the physics needs of the experiment, namely the need to keep the attenuation of ionization electrons to less than 20% over the longest drift distance in the LArTPC. In addition to the requirement on the electronegative contamination, the system must maintain the level of nitrogen contamination in the argon at less

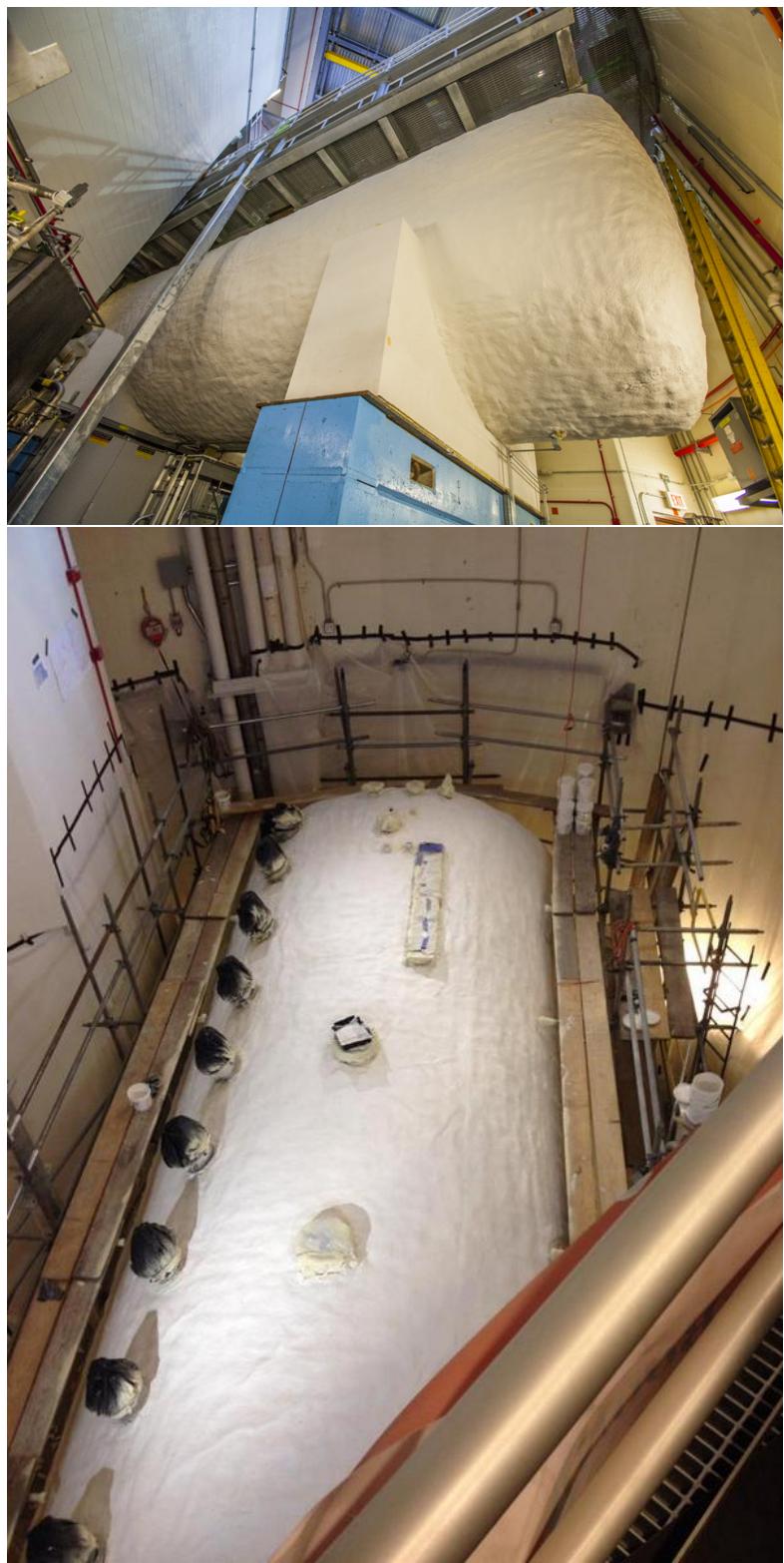


Figure 6. Photographs of the cryostat after application of exterior foam insulation.

Table 4. List of nozzle penetrations in MicroBooNE cryostat, their function, and their flange/pipe type and outer-diameter dimension. CF=ConFlat flanges, RFWN=raised face weld neck flanges, SS = stainless steel pipe.

Nozzle ID	Function	Flange
N1A-N1K	LArTPC Signal Feedthrough	356 mm CF
N2	LArTPC HV Feedthrough	203 mm CF
N3A-N3B	Purity Monitor	203 mm CF
N4	Temperature Signals	152 mm CF
N5	Safety Vent	102 mm RFWN
N6A-N6B	Vacuum Pump-Out	254 mm RFWN
N7	Condensor	76 mm SS
N8	Top Instrument Port	19 mm SS
N9	Bottom Instrument Port	19 mm SS
N10A-N10B	Liquid Level Probe	19 mm SS
N11	Gas Circulation In	51 mm SS
N12	Gas Circulation Out	51 mm SS
N13	From LAr Filters	76 mm SS
N14	To LAr Pumps	51 mm SS
N15A-N15B	Laser Calibration	70 mm CF
N16	PMT Signal Feedthrough	356 mm CF
N17	Spare	152 mm CF
N18	Temperature Signals	152 mm CF
N19	Spare	152 mm CF

than 2 parts per million (ppm) [32] to keep the attenuation of the scintillation photons in the argon to a minimum.

The MicroBooNE argon purification subsystem consists of liquid argon pumps and filters that serve to circulate the argon and remove impurities that degrade the quality of the data collected by the active detectors. There are two pumps in the system arranged in parallel in order to allow for continuous recirculation while one pump is being serviced. Similarly, there are two sets of filters arranged in parallel in the system. Figure 7 schematically depicts the flow of liquid and gaseous argon in the MicroBooNE cryogenic system.

The recirculation pumps are Barber-Nichols [33] BNCP-32B-000 magnetically-driven partial-emission centrifugal pumps. Each pump isolates the liquid argon from the electric motor. The impeller, inducer, and driving section of the magnetic coupling each have their own bearings that are lubricated by the liquid argon at the impeller. The motor is controlled by a variable frequency drive (VFD) that allows adjustment of the pump speed to produce the desired head pressure and flow within the available power range of the motor.

Each filter skid contains two filters, as depicted in figure 8, each having identically-sized filtration beds of 77 liters. The first filter that the argon stream enters contains a 4A molecular sieve supplied by Sigma-Aldrich [34] that primarily removes water contamination but can also remove small amounts of nitrogen and oxygen. The second filter contains BASF CU-0226 S, a highly-

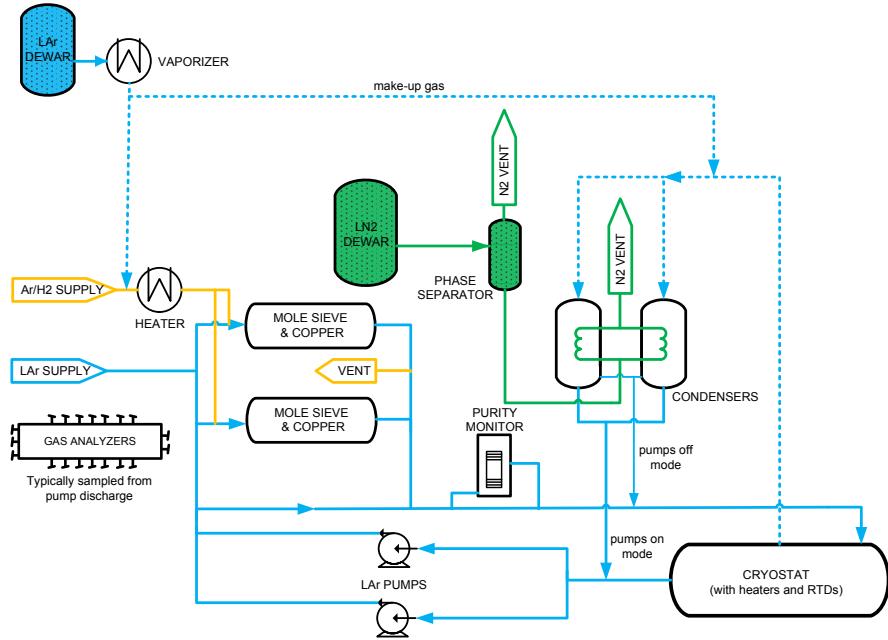


Figure 7. Flow diagram of argon in MicroBooNE, showing direction of liquid and gaseous argon in the cryogenic system. Gaseous argon from the cryostat is condensed and directed through the purification subsystem. Liquid argon drawn from the cryostat volume is directed into the purification subsystem.

1 dispersed copper oxide impregnated on a high-surface-area alumina, which removes oxygen [35]
 2 and, to a lesser extent, water. Thus, the oxygen filter is placed downstream of the molecular sieve
 3 to maximize oxygen filtration. The oxygen-filtering media must be reduced to copper with the
 4 procedure described below before it can remove oxygen from the liquid argon. The filters are
 5 insulated with vacuum jackets and aluminum radiation shields. The metallic radiation shields were
 6 chosen because the filter regeneration temperatures, described below, would damage traditional
 7 aluminized mylar insulation. Pipe supplying the filter regeneration gas is insulated both inside the
 8 filter vacuum-insulation space and outside the filter with Pyrogel XT which is an aerogel-based
 9 insulation [36] that can withstand temperatures up to 920 K.

10 The filters are regenerated *in situ* using heated gas, by a procedure developed for LAPD. The
 11 filters are regenerated using a flow of argon gas that is heated to 473 K, supplied by a commercial
 12 500 liter liquid argon dewar. Once the argon gas reaches 473 K, a small flow of hydrogen is mixed
 13 into the primary argon flow and exothermically combines with oxygen captured by the filter to create
 14 water. Too much hydrogen mixed in with the primary argon flow would induce temperatures that
 15 are sufficiently high to damage the copper-based filter media. The damage is induced by sintering of
 16 the copper, which reduces the available filter surface area. Thus, precautions are taken to maintain
 17 a hydrogen fraction below 2.5% of the heated gas mixture. During the heated gas regeneration, five
 18 filter-bed temperature sensors monitor the filter-material temperature and the water content of the
 19 regeneration exhaust gas is measured. To remove any remaining trace amounts of water, the filters

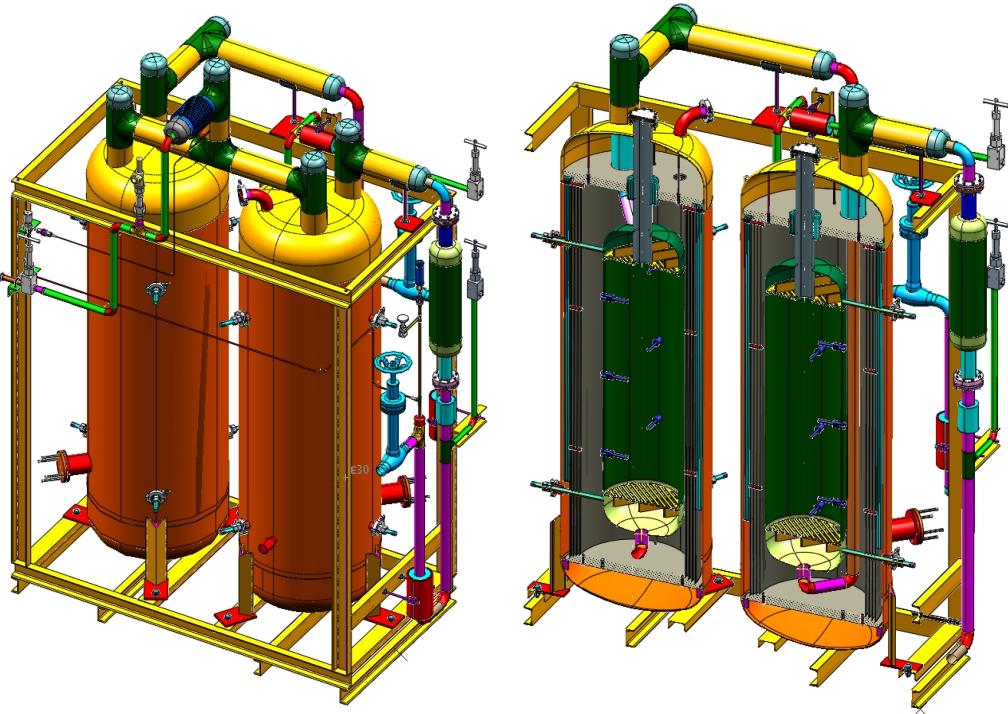


Figure 8. Three-dimensional rendering of a MicroBooNE filter skid. The left drawing shows the full skid, while the right drawing shows a cut-away of the vessels.

1 are then evacuated using turbomolecular vacuum pumps while they cool.

2 A particulate filter with an effective filtration of 10 microns, positioned between the cryostat
 3 and the filter skids, prevents any debris in the piping from being introduced into the cryostat.
 4 The particulate filter consists of a commercial stainless steel sintered-metal cylinder mounted in a
 5 custom cryogenic housing and vacuum jacket. Filtration is accomplished by flowing liquid argon
 6 to the interior, then outward through the walls, of the sintered-metal cylinder. Flanges on the argon
 7 piping, along with flanges and edge-welded bellows on the vacuum jacket, allow removal of the
 8 particulate filter.

9 The argon-purification piping is 2.54 cm diameter stainless steel that was pre-insulated by the
 10 manufacturer with 10.2 cm of polyurethane foam. During the fabrication process, all piping was
 11 washed with distilled water and detergent to remove oil and grease, then cleaned with ethanol. All
 12 valves associated with the argon-purification piping utilize a metal seal with respect to ambient air,
 13 either through a bellows or a diaphragm, to prevent the diffusion of oxygen and water contamination.
 14 The exhaust side of each relief valve is continuously purged with argon gas to prevent diffusion of
 15 oxygen and water from ambient air across the o-ring seal. Where possible, ConFlat flanges with
 16 copper seals are used on both cryogenic and room-temperature argon piping. Pipe flanges in the
 17 system are sealed using spiral-wound graphite gaskets. Smaller connections are made with VCR
 18 fittings with stainless steel gaskets.

3.3 Nitrogen Refrigeration

The cryostat and purification systems that contain the liquid argon are subject to heat load from the environment, as well as from the active detectors that have electrical power enabled. To keep these systems operating at a stable temperature and pressure, a liquid nitrogen refrigeration system is present to provide the necessary cooling power. The liquid nitrogen system contains two condensers that are arranged in parallel. One of these is utilized for normal operations and one serves as a backup on standby. Each condenser contains two liquid nitrogen coils, an inner and an outer, with the gas argon on the shell side, as shown in figure 9. Typically only one coil is actively running and the second can be manually activated during situations where the system heat load is higher than usual. Each condenser is sized to handle a heat load of approximately 9.5 kW. With the vessel full of liquid argon and no pump or liquid argon circulation running, the condenser uses ~2350 liters of liquid nitrogen per day, which equates to a 3.9 kW system heat load. Once liquid argon begins circulating using the pumps, this usage rate will increase by 50-100%.

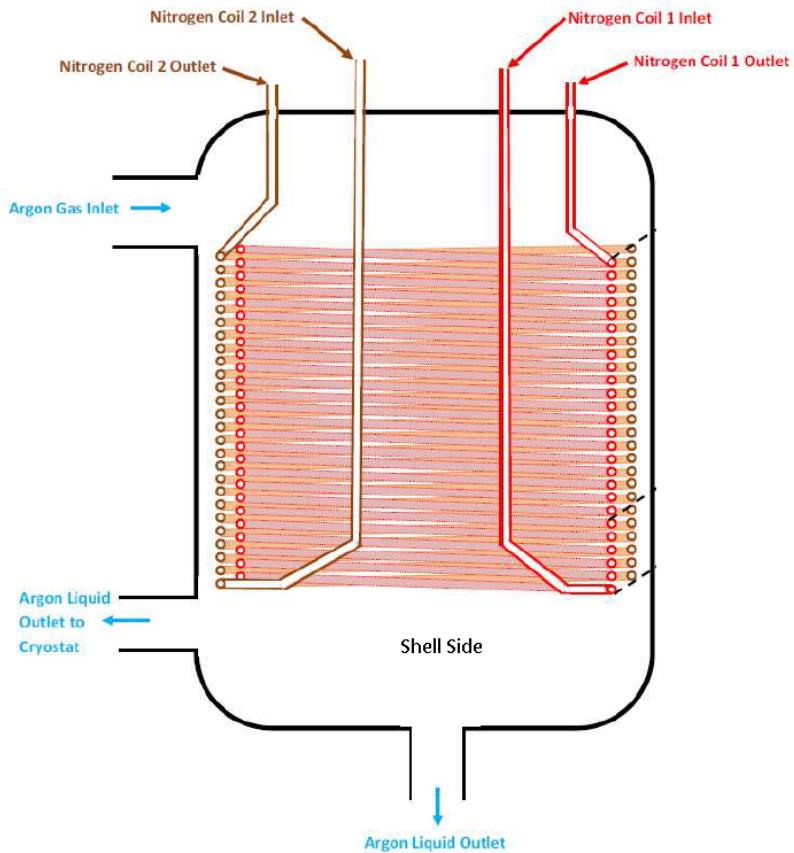


Figure 9. Diagram of a condenser.

3.4 Controls and Purity Monitoring

MicroBooNE makes use of resistive thermal devices (RTDs) to measure temperatures throughout the experimental infrastructure. Twelve RTDs are located along the walls of the cryostat, and another ten RTDs are mounted inside screws attached to the structure of the LArTPC. Each of the filter vessels in the purification system contain nine RTDs. The RTDs within the filter vessels prevent overheating which could potentially occur during filter regeneration with heated argon-hydrogen gas.

Liquid argon contaminations ranging between 300 and 50 ppt oxygen equivalent can be measured using double-gridded ion chambers, henceforth referred to as purity monitors, immersed in liquid argon. The design of the purity monitors is based on the design presented by Carugno et al. [37]. A thorough description of the purity monitors, the data-acquisition hardware and software used in LAPD can be found in [30]. MicroBooNE uses the same type of purity monitors, and the same data-acquisition hardware and software.

A measure of the electronegative impurities is made with the purity monitor by examining the fraction of electrons generated at the purity monitor cathode that subsequently arrive at the purity monitor anode (Q_A/Q_C) after a drift time t . The ratio of (Q_A/Q_C) is related to electron lifetime, τ , such that

$$Q_A/Q_C = e^{-t/\tau}. \quad (3.1)$$

Measurement of liquid argon purity in the MicroBooNE cryogenic system are provided by three purity monitors of varying lengths. One purity monitor with a drift distance of 50 cm sits in a vessel just downstream of the filters and is used to monitor filter effectiveness. Two purity monitors, one with a drift distance of 19 cm and the other with a drift distance of 50 cm, sit within the primary MicroBooNE vessel at each end of the LArTPC.

4 Liquid Argon Time Projection Chamber

The MicroBooNE LArTPC drifts and collects charge to produce fine-grained images of the ionization that is liberated by charged particles traversing a volume of highly-purified liquid argon. This section describes the design and implementation of the LArTPC in the experiment.

The LArTPC is composed of three major structures: the cathode, the field cage, and the anode. A negative voltage is introduced via a feedthrough passing through nozzle N2 on the cryostat and applied at the cathode, which defines an equipotential surface. This cathode voltage is incrementally stepped down in magnitude by means of a voltage divider chain spanning the field cage, creating a region of uniform electric field within the LArTPC. Opposite the cathode, and oriented parallel to it, are the anode wire planes: two induction planes (referred to as the “U” and “V” planes) with wires oriented at $\pm 60^\circ$ from vertical, followed by one collection plane (referred to as the “Y” plane) with vertically-oriented wires. The wires of the anode planes are the sensitive elements that detect the ionization created by charged particles traveling through the LArTPC. Figure 10 depicts the assembled MicroBooNE LArTPC after insertion into the cryostat, showing details of the cathode, field cage, and anode plane. Table 5 lists the main parameters of the MicroBooNE LArTPC, which will be described in detail in this section.

Table 5. MicroBooNE LArTPC design parameters and nominal operating conditions.

Parameter	Value
LArTPC (active) dimensions ($h \times w \times l$)	2.325 m \times 2.560 m \times 10.368 m
LArTPC (active) mass	90 tons
# Anode planes	3
Anode planes spacing	3 mm
Wire pitch	3 mm
Wire type	SS, diam. 150 μm
Wire coating	2 μm Cu, 0.1 μm Ag
# wires (total)	8256
# Induction0 plane (U) wires	2400
# Induction1 plane (V) wires	2400
# Collection plane (Y) wires	3456
Wire orientation (w.r.t. vertical)	+60°, -60°, 0° (U,V,Y)
Cathode voltage (nominal)	-128 kV
Bias voltages (U,V,Y)	-200 V, 0 V, +440 V
Drift-field	500 V/cm
Max. Drift Time, Cathode to U (at 500 V/cm)	1.6 ms
# Field-cage steps	64
Ring-to-ring voltage step	2.0 kV

4.1 Cathode

The cathode is assembled from 9 individual stainless steel sheets (Type 304, 2.3 mm thick) that are fastened to a supporting frame by hex button head stainless steel screws. The outer edge of

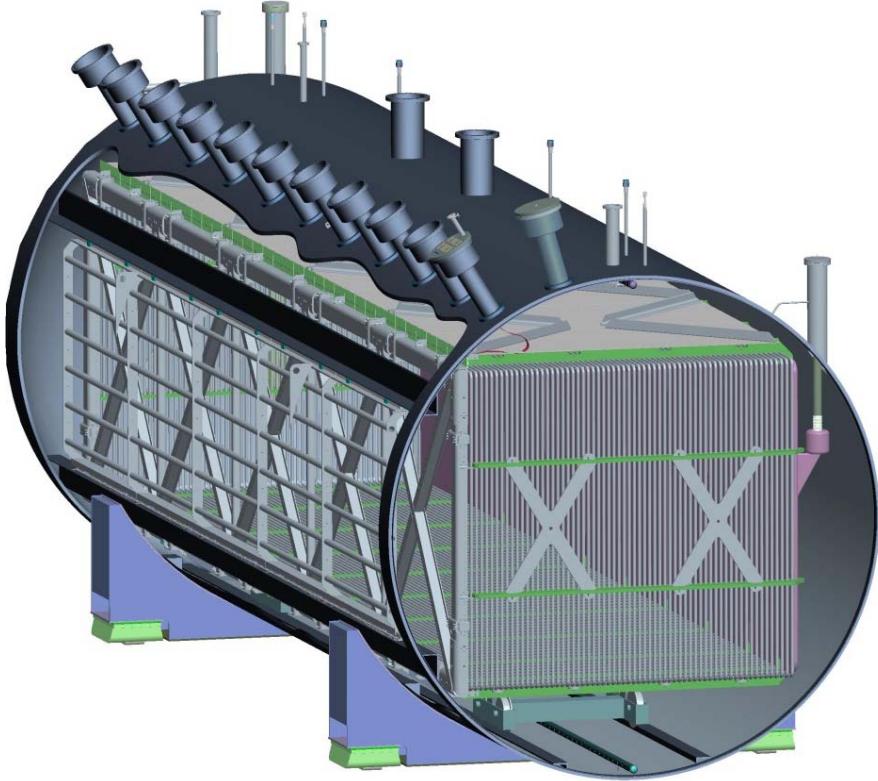


Figure 10. Schematic diagram of the MicroBooNE LArTPC , depicted as it is arranged inside the cryostat.

the cathode frame consists of round stainless steel tubes of 5.08 cm outer diameter and 3.18 mm wall thickness. Within this outer edge, square tubes with 5.08 cm \times 5.08 cm cross-sectional area, and 3.18 mm wall thickness, are fastened together with button head screws, forming a support structure upon which the cathode sheets are attached. The individual components of the support structure are further welded together to eliminate sharp features from this high-potential surface. The exterior frame and support structure of the cathode, and also an interior view, are shown in figure 11. The cathode plane sheets are shimmed according to survey data to make the cathode as flat and as parallel to the anode frame as possible. Flatness of the cathode is evaluated relative to a best fit plane of survey data (more than 10000 survey points recorded with a laser tracker). The largest deviations of the cathode from the best fit plane are +6.6 mm and -6.5 mm. Approximately 55% of the measured survey points fall within +/- 3 mm of the best fit plane, and more than 90% of the points fall within ± 5 mm. Figure 12 shows the results of the survey, with deviations from flat represented as color-coded data extending away from the nominal plane of an ideal cathode.

4.2 Field Cage

The field cage encloses the volume between the cathode plane and the anode wire planes, and in doing so creates a region with a uniform electric field. The volume defined by the interior of the field cage, bounded by the anode and cathode planes, is referred to as the “active” volume. The field cage structure consists of 64 individual thin-walled stainless steel tubes (2.54 cm OD, 0.51 mm



Figure 11. Top: Exterior view showing the cathode frame and structural supports to which cathode sheets are fastened. Bottom: Interior view of cathode plane as viewed from the upstream end of the LArTPC , showing cathode sheets.

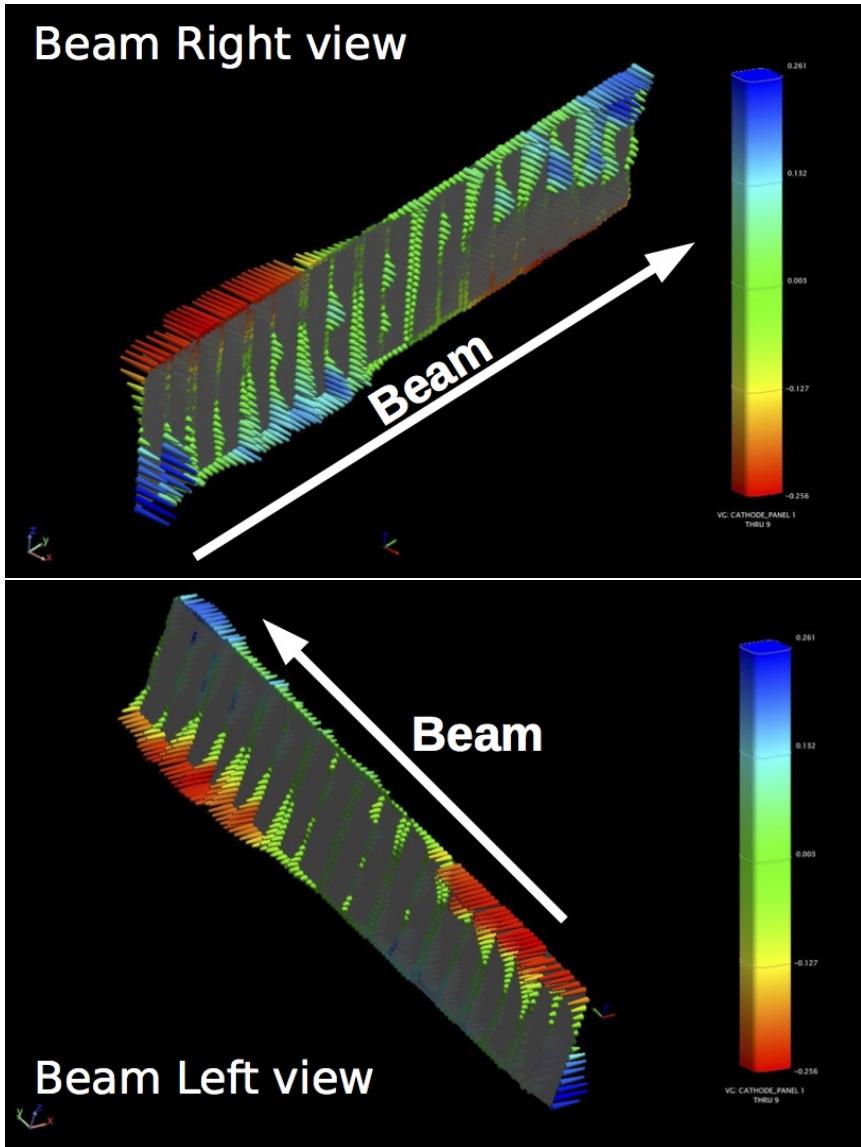


Figure 12. Survey results for interior (top) and exterior (bottom) of the cathode after shimming. Color scale extends from -0.256 cm (red) to +0.261 cm (blue).

1 wall thickness), each shaped into a rectangular loop framing the perimeter of the active volume.
 2 These 64 loops are mounted parallel to the cathode and anode planes, as shown in figure 11, and
 3 are held in place by a G-10 rib support structure. Each field cage loop is electrically connected
 4 to its neighbors via a resistor divider chain (described in following section), causing each loop to
 5 operate at a different electrical potential which in turn maintains a uniform electric field between
 6 the cathode and anode planes. For a nominal -128 kV cathode voltage, the difference in potential
 7 between adjacent field cage loops is 2 kV, ramping down the total potential in equidistant steps from
 8 cathode to anode. The distance from center-to-center of adjacent field cage loops is 4.0 cm.

9 Each field cage loop is assembled from 2.07 m long vertical pipes on the upstream and
 10 downstream ends of the LArTPC , and on the top and bottom from two 5.18 m long horizontal pipes

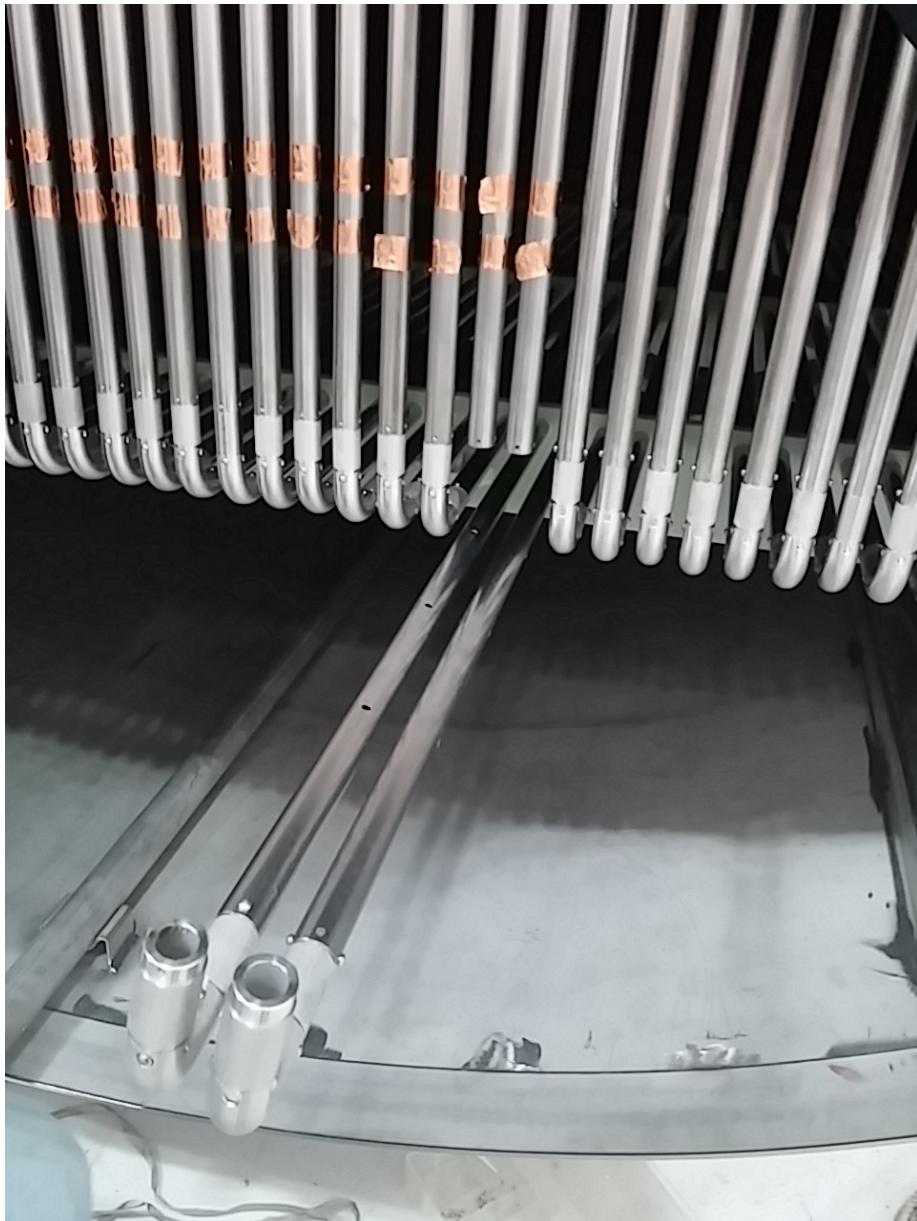


Figure 13. Field cage loops, with two partially disassembled to show elbow and couplings.

1 connected by a stainless steel coupling in the center. Each tube has venting holes approximately
2 every 15 cm so that they fill with liquid argon and allow air and gaseous argon to escape during
3 cryostat cool down and filling.

4 The four corners of each field cage loop are curved with a radius of 5.24 cm. Each corner is
5 formed by three parts: two couplings and an elbow, shown in figure 13. The couplings make the
6 connections between the pipes and the elbow. The thin-walled tubes and elbows slip-fit over the
7 ends of the couplings with a 2.2 cm overlap. Each coupling has two 6-32 NC tapped holes and the
8 connections to the adjoining pieces are made by hex head button screws and lock washers.

9 In order to avoid electrical breakdown between the inner cryostat surface and field cage parts

at high potential on or near the cathode, the electric field strength is minimized at the corners and edges of the field cage. Loops 0, 1, and 2 are each designed differently than the other field cage loops. Loop 0 is a special case in that it is made from larger diameter piping of 5.08 cm OD, and frames the cathode and also acts as its mechanical support. It operates at the same electrical potential as the cathode plane sheets attached to it. Loop 0 has a slightly smaller area than the other field cage loops, as shown in figure 14. Loop 1 is the first of the 64 loops in the field cage with 2.54 cm OD tube, and is the edge of the cathode, operating at the same electrical potential as loop 0. The elbow of loop 1 has a specially designed geometry in order to minimize the electric field potential. The elbow of loop 2 has a softer radius of curvature than the standard elbows, also for the purpose of minimizing the electric field potential. For all three of these loops (loop 0, 1, and 2), connections at corners and joints are made by welding instead of screws to avoid sharp edges that would result in higher electric fields and greater chance of electrical breakdown.

Another precaution to minimize the electrical field between the loops and the cryostat surface is the positioning of the coupling screws: for the first 20 loops, the screws are positioned on the sides facing the screws of the neighboring loops instead of facing inward to the LArTPC active volume and outward toward the grounded cryostat surface. Hex head button screws and lock washers are also used here in order to minimize sharp metal edges.

Figure 15 shows the simulated electric field values inside the MicroBooNE LArTPC , and some of the surrounding volume inside the cryostat, when the cathode is set to an operating voltage of -128 kV.

4.2.1 Resistor Divider Chain

A resistor divider chain installed across the field cage loops steps the voltage down in magnitude from the cathode plane to the anode wire plane in equidistant steps. For a nominal value of -128 kV on the cathode, this results in a potential difference of 2 kV between each pair of loops. The value of the equivalent resistance between loops within the divider chain was chosen to be low enough such that the current flow through the divider circuit is much greater than the signal current flowing through the LArTPC . The signal current in this case is the total energy deposited by cosmic rays, and is estimated to be <50 nA. An equivalent resistance of 250 M Ω between each pair of field cage loops, corresponding to a current flow of 8 μ A, was chosen.

The voltage divider chain is mounted on the inside of the field cage at the upstream end of the detector. The couplings at the top corner of each field cage loop have additional holes facing the inside of the field cage, where the resistors are mounted. On the first 16 field cage loops, pairs of Metallux HVR 969.23 499 M Ω resistors (rated to 23 W, 48 kV in air) are mounted electrically in parallel to establish the beginning of the voltage divider chain. On the remaining loops, thick-film Ohmite Slim-Mox 104E metal-oxide epoxy-coated resistors with a lower power and voltage rating (1.5 W, 10 kV in air) are used. Extensive testing was done on these two types of resistors [38]. For redundancy, the desired 250 M Ω total resistance between each pair of field cage loops is composed of four individual 1 G Ω Slim-Mox 104E resistors placed electrically in parallel and attached to a printed circuit board. Printed circuit boards span across eight field cage gaps and therefore have eight 250 M Ω resistances in series, shown in figure 16. The electrical connection between the boards and each field cage loop is made by metal contact pads on the back side of the boards, held in electrical contact with the field cage tube by a button-head hex screw and lock washer.



Figure 14. Field cage loops closest to (and including) the cathode are modified to reduce sharp edges that would result in higher electric fields.

1 While the designed operating voltage difference across each resistor in the detector is 2 kV
2 with a power flow of 4 mW, there is a slight possibility that these values could temporarily exceed
3 the rating of the resistors in the case of discharge between the cathode plane or field cage loops and
4 the cryostat wall, through the bulk liquid argon. Recent studies [39] have shown that the value of
5 the minimum breakdown electrical field decreases with the increasing argon purity; for purities as
6 high as that required in the MicroBooNE detector, breakdown has been observed at electric fields
7 as low as 40 kV/cm.

8 The field cage behaves like a capacitance network at high frequencies. Based on measurements

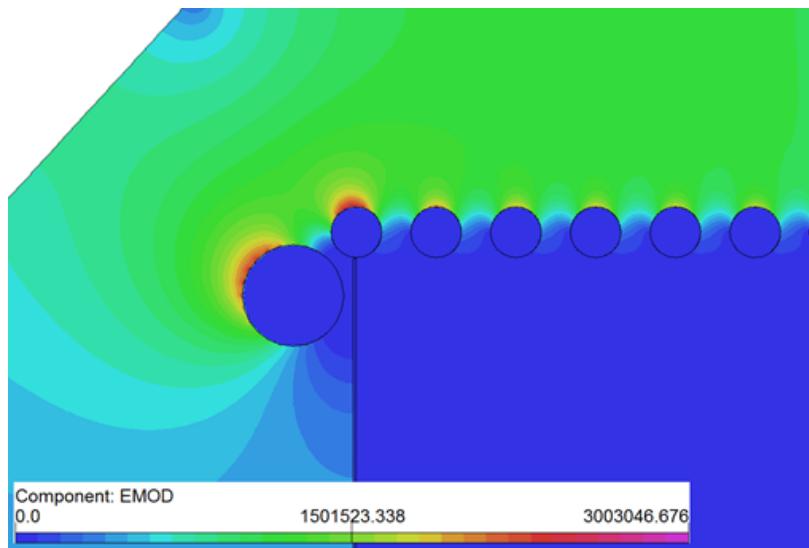


Figure 15. Cross-section view showing electric field simulation inside the field cage when the cathode is set to a voltage of -128 kV. Loop 0 is represented by the larger diameter circle on the left of the image. The legend shows the absolute values of electric field modulus in units of V/m.

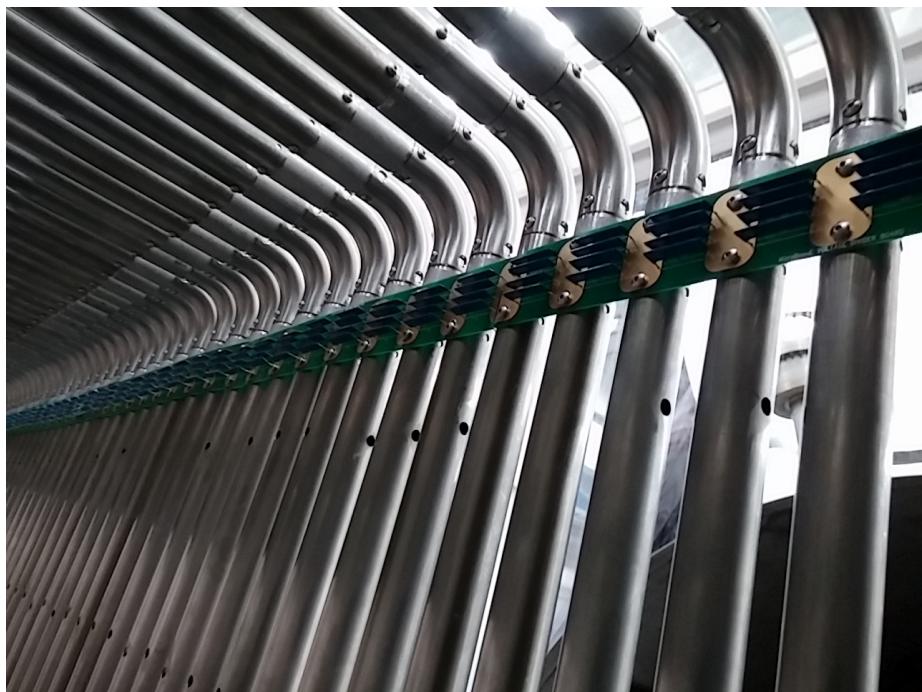


Figure 16. The Ohmite Slim-Mox 104E resistors arranged in parallel sets of four on printed circuit boards that span eight field cage tubes.



Figure 17. The Metallux HVR 969.23 resistors mounted on the 16 field cage loops closest to the cathode.

and simulations, the total energy stored inside the field cage when fully charged is estimated to be approximately 24 J. In the case of a discharge between the cryostat and the cathode or one of the field cage loops close to the cathode, simulations show that voltages of up to 80 kV peak, with a discharge time constant of a few seconds, can develop across the resistors. The observed peak voltages in such discharge scenarios decrease the further the breakdown occurs from the cathode, such that discharges occurring between the cryostat and field cage loops 32 through 63 do not exceed the 10 kV rating of the resistors.

In order to protect the resistors near the cathode from damage which could potentially render the field cage inoperable, two strategies have been implemented. First, due to the knowledge obtained in the breakdown electrical field studies mentioned before, on the first 16 field cage loops near the cathode the Ohmite Slim-Mox 104E resistors are replaced by 499 M Ω Metallux HVR 969.23 resistors that have a higher voltage rating of 48 kV. An extensive study of resistor breakdown in liquid argon of several resistor brands and models [38] has revealed that the breakdown voltage observed in liquid argon exceeds the rating in air substantially for all resistors. No breakdown has been observed up to 32 kV for the Ohmite Slim-Mox 104E resistor. For the Metallux HVR 969.23 resistors, no breakdown has been observed up to the limit of the test apparatus of 130 kV. Since the Metallux HVR 969.23 resistors are significantly larger physically than the loop-to-loop distance, they are mounted diagonally between each pair of field cage loops. They are held by copper brackets, which are attached to studs welded onto the field cage tubes, shown in figure 17.

The second type of protection installed for all resistors on field cage loops 1 through 32 is a surge protection circuit. The chosen surge protection devices are designed to short the circuit in the case of a voltage spike, which protects any other electrical components installed in parallel. Below their



Figure 18. Surge-protecting varistors are installed in parallel with the voltage divider resistors for the first 32 field cage loops. Here, they are shown mounted on small boards in sets of 3, and attached to the field cage by means of 6-32 hex button head screws.

clamping voltage, they exhibit a very high resistance and do not influence the circuit. The behavior of Gas Discharge Tubes (GDTs) and varistors in liquid argon has been studied extensively for application in the MicroBooNE field cage [40]. The surge protection device chosen is a Panasonic ERZ-V14D182 varistor with a clamping voltage of 1700 V. In order to obtain a very high resistance in normal operation and a clamping voltage above the 2 kV in normal operation mode, three of these devices are mounted in series across a block of four Slim-Mox 104E or two Metallux 969.23 resistors. These additional varistor boards make electrical contact with the field cage via brass mounting brackets that are fastened to the field cage with button head screws, as shown in figure 18.

4.3 Anode

The anode frame holds the induction and collection plane sense wires at tension and provides overall structural support for the beam-right side of the LArTPC. Individual sense wires for all anode planes are held in place by wire carrier boards, which are printed circuit board assemblies that locate the wires as well as provide the electrical connection to the electronic readout system of the experiment.

4.3.1 Mechanical structure

The anode frame is comprised of a stainless steel C-channel hosting adjustable tensioning bars to which the wire carrier boards are attached. The C-channel and tensioning bar assembly is depicted for one corner of the anode frame in figure 19. Wire carrier boards attach to precision alignment pins distributed along the length of the tensioning bars.

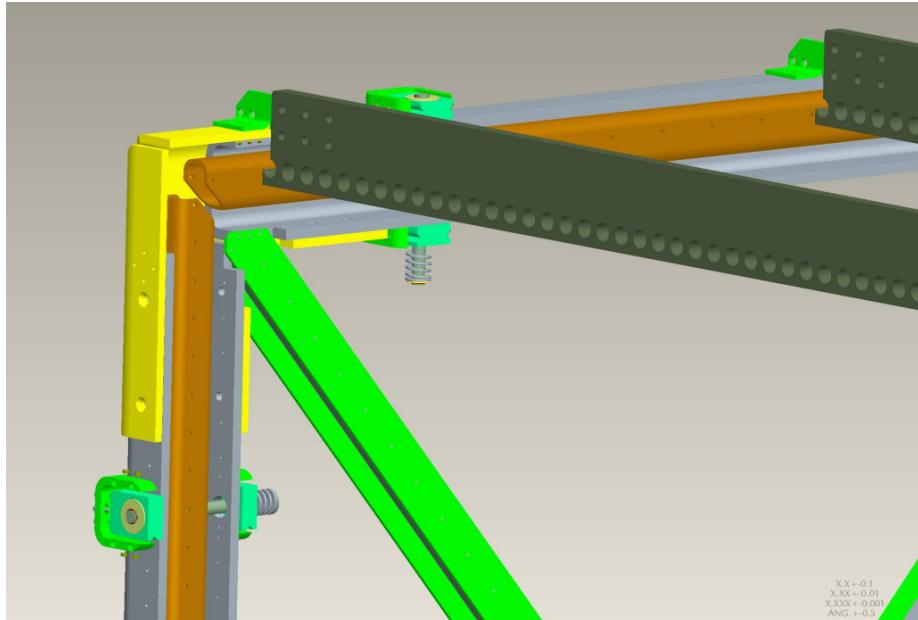


Figure 19. Rendering of the anode frame assembly. The C-channel is depicted in gray, and the adjustable tensioning bar assembly is shown in orange.

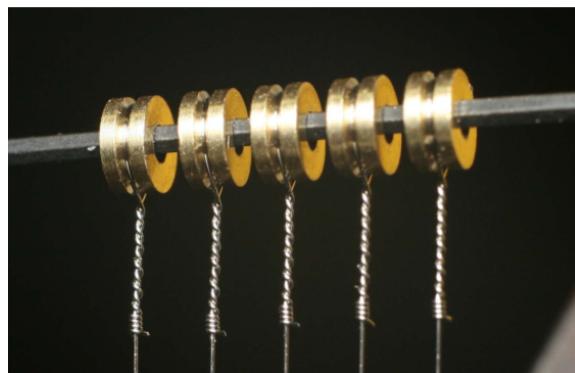


Figure 20. Photograph of the wire termination on the brass ferrules. Each ferrule is 3 mm in diameter, and 1.5 mm thick.

4.3.2 Wire winding and quality assurance

The three anode planes are constructed from wire carrier boards that have had individually-prepared wires attached to them in groups of 16 (for the U- and V- angled planes) or 32 (for the vertical Y-plane). Consistent quality in wire preparation was achieved by a semi-automated winding machine, which terminated the ends of each wire via wrapping around 3 mm diameter brass ferrules as shown in figure 20. Each wire was tested for strength by being placed on a tensioning stand where a load of 2.5 kg (more than 3 times the nominal load of 0.7 kg) was applied for 10 minutes, ensuring that the wire preparation did not leave any weaknesses that could result in a breakage later on. Upon successful completion of the quality assurance testing for each wire, it was placed onto a wire carrier board, shown in figure 21

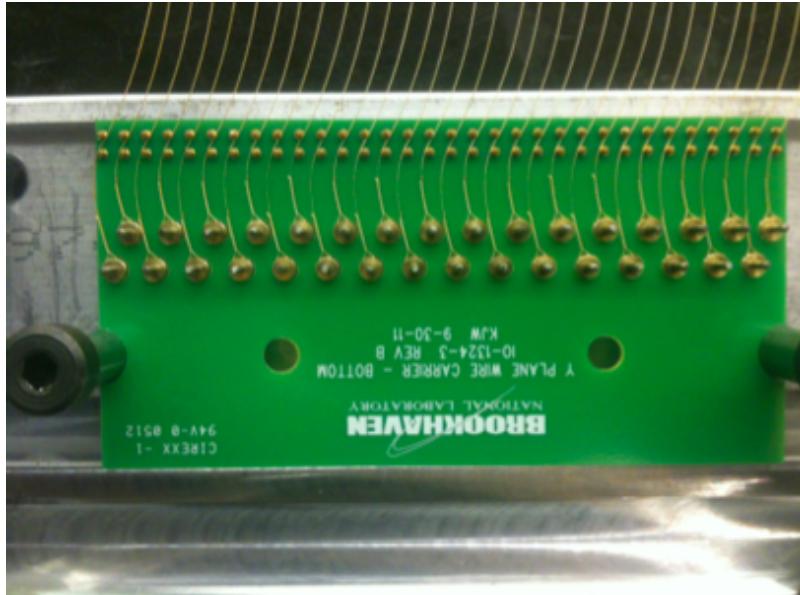


Figure 21. Photograph of a collection plane wire carrier board that has been filled with wires, but has not yet had the cover plate installed.

When installed, the wires make contact with gold pins positioned on the wire carrier board, and these pins are connected to a trace that routes to the cold electronics described in section 6.1. After fully filling a wire carrier board with wires, a cover plate was installed and press-fit rivets were installed to hold the assembly together. The assembled wire carrier board was then placed onto a tension stand, to reapply a 2.5 kg tension/wire to the whole board for 10 minutes. This is to ensure that the wires were not weakened during the board assembly process. The tension stand is depicted in figure 22. A comprehensive description of the MicroBooNE wire preparation and associated quality assurance studies can be found in [41].

4.3.3 Wire installation and tension measurements

During detector construction the completed wire carrier assemblies were manually installed onto the adjustable tensioning bars residing in the C-channel of the supporting anode frame. A team of two people installed each assembly (consisting of wires and supporting carrier boards on either end) onto the anode frame. The collection plane was the first installed, followed by the middle induction plane, and then finally the inner induction plane. Once all three anode planes were completely installed, the tensioning bars were adjusted and a survey was taken of the tension of all anode wires. Tension was measured via measuring the resonant frequency of a laser beam reflected off of a plucked wire and incident on a photodiode connected to a spectrum analyzer program [42]. The tension measuring equipment was developed and produced by the University of Wisconsin Physical Sciences Laboratory. The tensioning bars were adjusted iteratively until the surveyed tension of all wires was within a range, approximately ± 1.0 N of the nominal value of 6.9 N, where no single wire was too taught or loose to create detector performance issues. Figure 23 shows the final surveyed tension of the wires within the LArTPC.



Figure 22. Photograph of a collection plane wire carrier board on the tension stand.

4.4 Parts Preparation

The majority of the parts that make up the LArTPC are either stainless steel or G-10. These two material types, as well as any others used in the LArTPC, were tested in the Fermilab Materials Test Stand (MTS) [43], whose purpose was to investigate the suitability of materials for use in LArTPCs. The MTS confirmed that none of the materials used in the LArTPC assembly would contaminate the liquid argon. Before assembly, all LArTPC parts were cleaned according to the procedures described in the following sections.

4.4.1 Cleaning stainless steel

The delivered stainless steel parts were often greasy due to machining, and those with holes or interior cavities generally had a significant amount of trapped metal shavings due to the machining processes. Many of the pieces also had markings from permanent ink pens, dirt smears, rust spots,

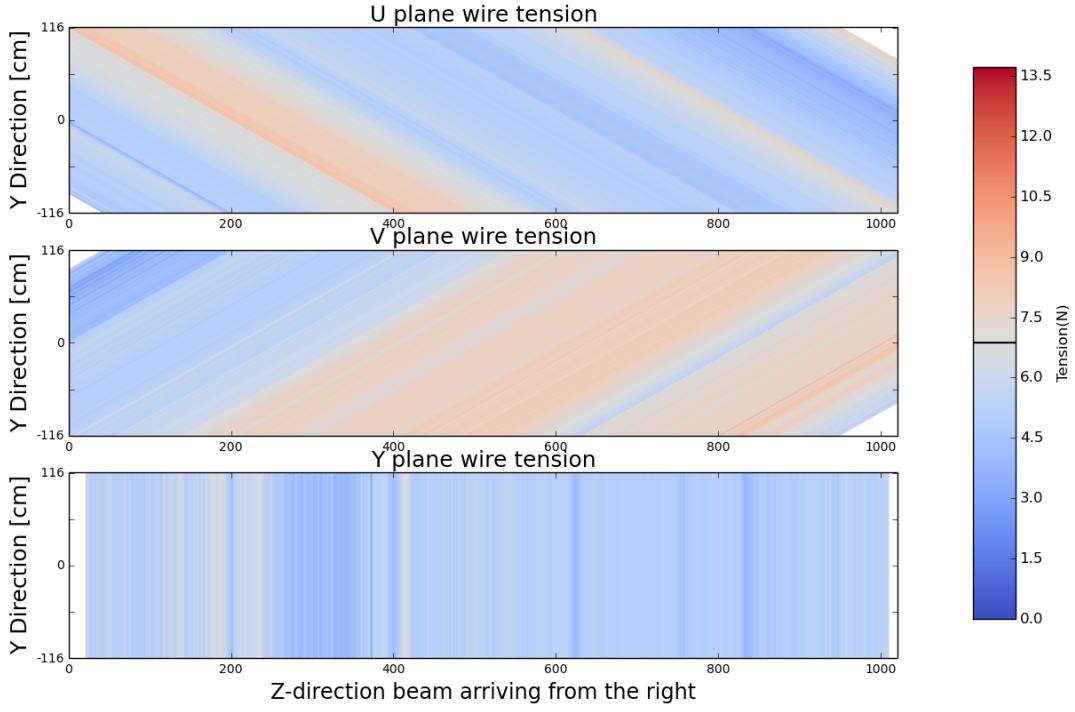


Figure 23. Final survey results for wire tension of the MicroBooNE LArTPC.

and/or dried oil from machining. These pieces were scrubbed with ScotchBrite 7447 general-purpose hand pads before cleaning.

Parts that were small enough to fit in an ultrasonic bath were prepared according to the following prescription. A pre-rinse with tap water was performed to remove particulate matter, followed by deburring of sharp edges. The first ultrasonic wash was 15 minutes in heated distilled water with a 3% solution of Citranox acid detergent [44]. After a first rinse in distilled water, a second ultrasonic wash was performed, again for 15 minutes, but using heated distilled water with a weak solution of Simple Green detergent [45]. A second rinse in distilled water was performed, followed by a final rinse in a fresh bath of distilled water. The parts were then wiped dry with lint-free cloths, air dried completely, and wrapped in plastic film for storage.

Stainless steel parts that were too large to fit in the ultrasonic bath were prepared by a simpler prescription out of necessity. A pre-rinse with tap water was followed by deburring to remove sharp edges. Both the first and second washes were done with tap water and a weak solution of Simple Green detergent, scrubbing with brushes and lint-free sponges. Two tap water rinses were done, and the parts were then wiped dry with lint-free cloths. As a final additional step, each part was wiped with 200-proof ethyl alcohol, and then air dried completely. These parts were also wrapped in plastic film for storage.

4.4.2 Cleaning G-10

G-10 is known to absorb large quantities of water, which would outgas in the argon and could inhibit reaching the required argon purity in the detector. For this reason all G-10 parts were cleaned and then baked to remove moisture. The largest G-10 parts on the detector are beams that span the

1 distance between the cathode and anode. These were washed in 1900-liter ultrasonic baths that are
2 overseen by Fermilab Accelerator Division, typically used for cleaning large sections of accelerator
3 beam pipes. An initial pre-wash was done with tap water to remove as much particulate matter as
4 possible, since the machining process left a large amount of dust on the machined edges. Pieces
5 were then placed in the ultrasonic bath with heated deionized water and a 2% solution of Elma
6 Clean 65 (EC 65) neutral cleanser [46]. Two ultrasonic bath rinses were performed, and the pieces
7 were then sealed in plastic bags with clean dry nitrogen gas. In order to remove the absorbed water,
8 the large G-10 parts were then transported to Fermilab Technical Division where they underwent
9 an outgassing procedure to remove any remaining absorbed moisture. They were baked in a large
10 oven under vacuum until a plateau in the outgassing was reached, as reported by a monitor inside
11 the oven. Upon completion of the outgassing procedure, the parts were resealed in plastic bags for
12 storage.

13 **4.5 Assembly**

14 The full mechanical structure of the LArTPC is shown in figure 24, with the top image depicting
15 the cathode frame on the left, made semi-transparent to show the support structures on which the
16 cathode sheets are attached. The anode frame is on the right of this image, with I-beams configured
17 in a crossed pattern to maintain the shape and rigidity of the outer C-channel structure. Ribs of G-10
18 connect the anode and cathode, electrically isolating them from each other while also providing
19 mounting holes to hold in place each of 64 field cage loops that define the active volume of the
20 LArTPC. The field cage loops are visible in the photograph in the bottom frame of figure 24.

21 Assembly was done inside of a clean tent, shown in figure 25, on a flat surface made up of
22 adjustable-height metal platforms that were installed on the assembly room floor. These platforms
23 were leveled to better than 0.5 mm before beginning assembly. The anode frame was the first part
24 of the detector to be assembled on this surface, shown in figure 26. It was temporarily placed aside,
25 and the cathode frame was assembled on the same set of platforms along with the G-10 ribs, which
26 stood vertically with the help of temporary unistrut support pieces. The combined cathode and G-10
27 frame was then lifted and rotated to the proper orientation, with G-10 ribs extending horizontally
28 from the cathode to the anode, as shown in figure 27. Finally, the anode frame was brought back
29 over and attached to the G-10 ribs, and the stainless steel tubes that make up the field cage loops
30 were fed through the holes in the G-10 ribs to complete the mechanical structure of the LArTPC.

31 **4.6 High Voltage System**

32 To create the drift field, negative voltage (referred to as the “high-voltage” or “HV”) is supplied to
33 the LArTPC cathode. This potential is generated outside of the cryostat by a Glassman LX150N12
34 power supply. Before entering the cryostat, the output of the power supply is passed through a
35 current-limiting resistor series that serves as both a low-pass filter for the power supply ripple, and
36 a partition for the stored energy in the system. The resistor series is a set of eight $10\text{ M}\Omega$ resistors
37 submerged in a transformer oil in an aluminum container. This assembly was successfully tested to
38 -200 kV.

39 The potential is introduced into the cryostat by a custom-designed HV feedthrough. The
40 feedthrough is based on an ICARUS design [7]; a 2.54 cm diameter stainless steel inner conductor

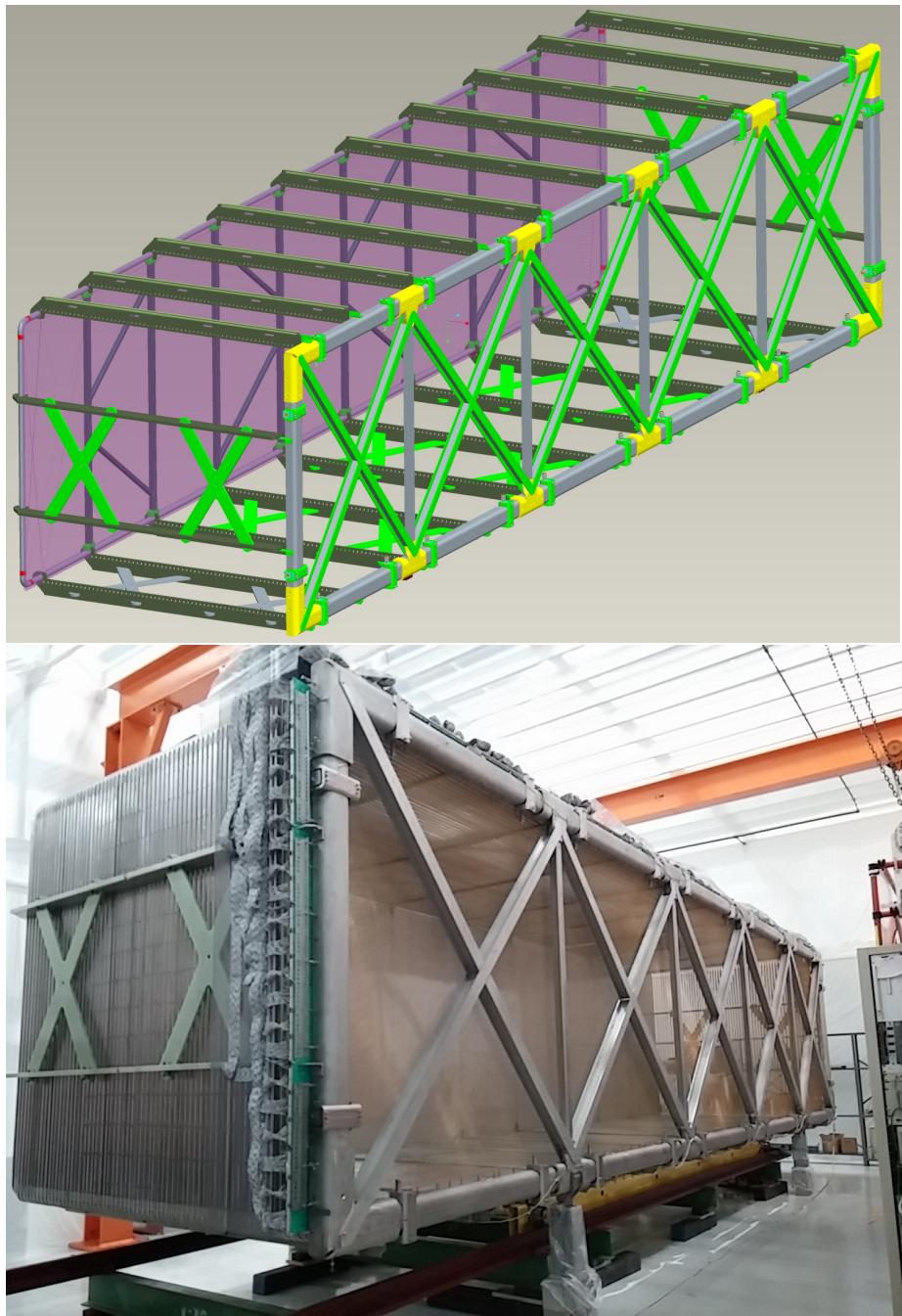


Figure 24. Top: Rendering of the full LArTPC frame assembly. Bottom: Assembled LArTPC after wire and electronics installation.



Figure 25. Clean tent where MicroBooNE LArTPC assembly was conducted.



Figure 26. Anode frame in the process of being moved from the metal assembly platforms.



Figure 27. Cathode frame and G-10 ribs on metal assembly platforms.

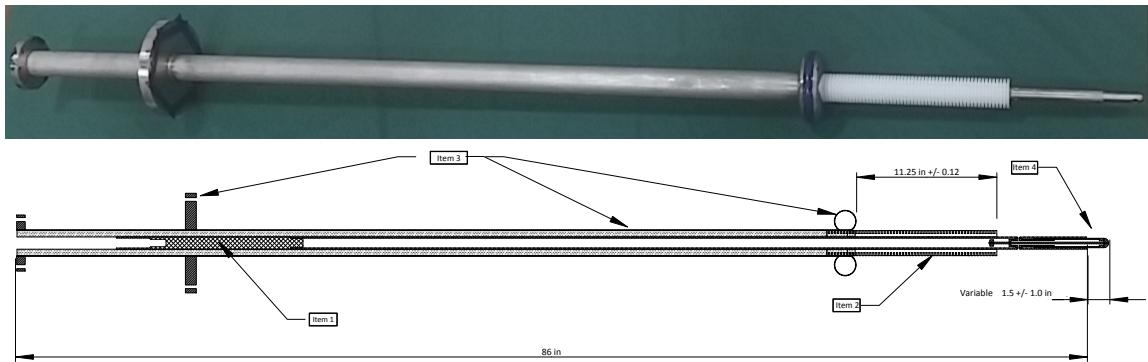


Figure 28. Photograph and drawing of the production HV feedthrough. The spherical probe tip is attached to the end of the inner conductor, on the right side of these figures.

1 is surrounded by a 5.08 cm outer diameter ultra-high molecular weight polyethylene (UHMW
 2 PE) tube that is encased in an outer ground tube. A photograph and drawing of the production
 3 feedthrough are shown in figure 28.

4 Figure 29 shows the HV feedthrough extending into its receptacle cup attached to the LArTPC
 5 cathode. The lower termination of the outer ground tube is a torus chosen to reduce the electric
 6 field between both the feedthrough and the cathode plane and along the feedthrough itself. The
 7 electrical connection to the cathode is made with a hemispherical spring-loaded tip attached to the
 8 inner conductor of the feedthrough. The UHMW PE tube is machined with grooves and extends



Figure 29. The production HV feedthrough inserted into the cathode receptacle cup inside the cryostat.

into the receptacle cup attached to the cathode.

5 Light Collection System

Liquid argon is a very bright scintillator, and collecting the light produced when charged particles travel through the LArTPC is a critical aspect of fully understanding the interactions taking place in the detector. The light collection system is designed to meet MicroBooNE’s physics goals for light collection, which are twofold. First, for accelerator-induced events, the light collection system is designed to trigger on 40 MeV kinetic energy protons produced in neutrino interactions. Second, for non-beam events, the system is designed for efficient observation of 5 to 10 MeV electrons from supernova neutrinos.

The light produced by neutrino interactions in MicroBooNE is an important input for both event selection and reconstruction. One of the critical capabilities the light collection system provides is the ability to form a beam-event trigger when a pulse of light is observed in coincidence with the beam spill. Because a vast majority of beam spills will not produce a neutrino interaction in the detector, such a trigger will substantially reduce the data output rate. For non-beam physics studies, the light system provides triggering and an event t_0 . For accelerator-induced events, the start time of a beam spill is, in principle, sufficient. However, because of the long window over which the ionization electrons of an event drift (about 1.6 ms maximum drift time at 500 V/cm field), there are many accelerator-induced events in which a cosmic ray muon crosses the detector during the drift time. Utilizing the position of hits in the photodetectors, one can better reject cosmic ray muon tracks. The light can also be used to trigger and select specific types of cosmic-ray calibration events (Michel electrons, straight-through muons, etc.) and non-beam events (supernova neutrinos, cosmic background events to proton decay studies, etc.).

The light collection system consists of primary and secondary sub-systems. The primary light collection system is made up of “optical units,” each one consisting of a PMT located behind a wavelength-shifting plate. In total, 32 optical units were installed, yielding 0.85% photocathode coverage. The secondary system consists of four light guide paddles. These paddles are introduced for R&D studies for future LArTPCs, and are placed near the primary optical units to allow a comparison of their performances. A flasher system, used for calibration, consists of optical fibers bringing visible light from an LED to each PMT face.

The light collection detectors are located in the y - z plane behind the anode planes of the LArTPC, as shown in figure 30. The combined transparency of the three anode planes is 86% for light at normal incidence. This transparency value assumes 100% of VUV photons impinging on the wires (150 μm diameter and 3 mm pitch) are absorbed. The detectors were placed so as not to be obscured by the LArTPC structural cross-bars, shown in figure 30. Locating the light detectors behind the anode plane places them in a very weak electric field due to the +440 V bias of the collection plane. To test for an effect from weak electric fields, the response of a PMT placed between a +700 V mesh and ground, separated by 50 cm, was studied. The PMT zenith was 25 cm from the HV. No effects on the signal were observed. This was expected, as the photocathode is held at ground, effectively acting as a Faraday cage.

Throughout the design and construction of the light collection system, substantial R&D was performed. The reader should refer to [32, 47–56] for detailed results of these studies. A useful overall review is available in [57].

Figure 31 shows the light observed in two sequential events, consistent with a muon entering

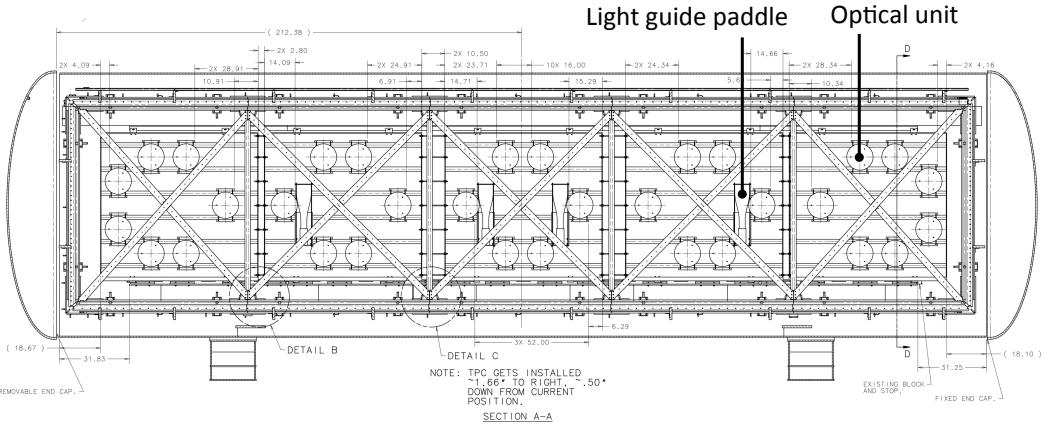


Figure 30. The MicroBooNE light collection system consists of a primary system of 32 optical units and a secondary optical system of four lightguide paddles [47]. These are mounted behind the anode wire planes such that the view is not obscured by structural cross bars of the LArTPC.

the detector followed by a Michel electron from the decay. One can see that the light is relatively well localized. This allows the light to be correlated with specific tracks in the detector. This “flash-track matching” is used to identify and reconstruct the tracks that are in time with the beam spill—an important goal of the light collection system.

5.1 Light Production in Argon

Light produced in liquid argon arises from two processes: scintillation and Cherenkov radiation. Scintillation light is produced by the formation and eventual radiative decay of excited argon dimers (or eximers) and is emitted in an isotropic distribution. Liquid argon is an excellent scintillator: it produces a large amount of light per unit energy deposited (about 24,000 photons per MeV at 500 V/cm drift field) and is transparent to its own scintillation. The scintillation light has a prompt and slow component with decay times of about 6 ns and 1.6 μ s, respectively. The two lifetimes correspond to the two lowest-lying eximer states with the prompt component coming from the decay of a singlet state and the slow from the decay of a triplet state. The prompt to slow ratio is about 1:3 for minimum ionizing particles and varies with ionization density and particle type. Both components consist of photons with a wavelength of 128 nm.

There is a significant uncertainty on the expected triplet lifetime [58] and this may be further modified by quenching (non-radiative dissociation of excimers by impurities) [59]. Other factors that can affect the arrival of the light include Rayleigh scattering, absorption by impurities, and obstructions. For detailed discussion of the physics of scintillation light production and propagation in MicroBooNE, see [57]. Table 6 summarizes information about the scintillation light.

Because scintillation photons have a wavelength of 128 nm they are very difficult to detect using conventional photodetectors. Figure 32 summarizes the challenges involved in detection of the 128 nm scintillation light. In order to detect the scintillation light, with spectrum shown in red, the VUV photons must be shifted into the visible region. MicroBooNE employs tetraphenyl-butadiene (TPB). This organic fluor absorbs in the UV (green line) and emits in the visible with a peak at 425 ± 20 nm (green hatched region), the peak wavelength having a slight dependence

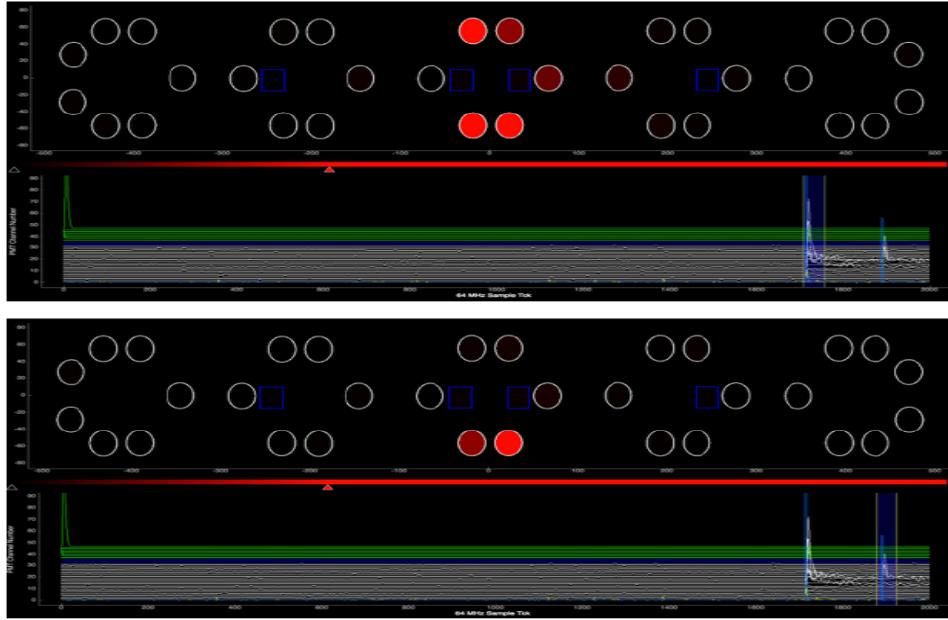


Figure 31. Two sequential event displays for the light collection system. The sequence is consistent with a muon that stops (top) and decays (bottom). The circles correspond to the optical units. Red circles indicate those units with hits. The waveforms versus time are shown below.

Table 6. Important properties of scintillation light in liquid argon that affect detection in MicroBooNE.

Property	Value	Reference
Wavelength	128 nm	[60]
Singlet, Triplet state time constants	6 ns, 1.6 μ s	[58]
Photons/MeV for $E = 500$ V/cm	24000	[7]
Triplet lifetime quench due to N ₂ at 1 ppm	20%	[59]
Attenuation length due to N ₂ at 1 ppm	66 m	[32]
Rayleigh Scattering Length	\sim 90 cm	[61]
Transparency of the wireplanes at normal incidence	86%	-

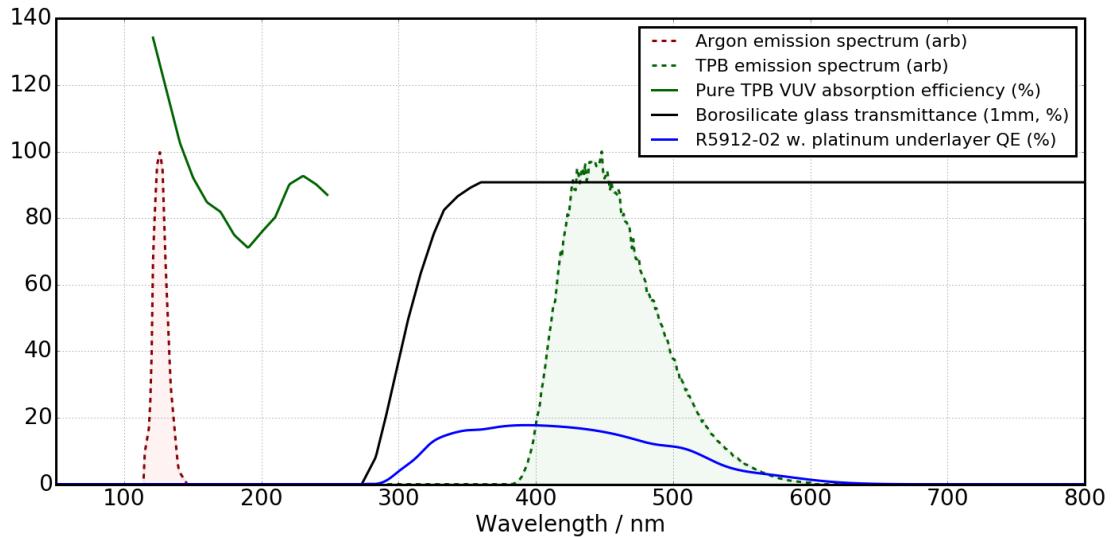


Figure 32. Scintillation light emission spectrum (red) and TPB re-emission spectrum (green), in arbitrary units. Superimposed are important efficiencies (see y axis): Dark green line – absorption of VUV light by TPB; Black line – transmission of borosilicate glass; Blue line – efficiency of a R5912-02mod cryogenic PMT [57]

on the micro-environment of the fluors. This is a favorable wavelength for detection by the PMTs employed by MicroBooNE. The efficiency for transmission through borosilicate PMT glass (black) and the quantum efficiency of the cryogenic tubes used in MicroBooNE (blue) are overlaid on the TPB spectrum.

5.2 The Primary Light Collection System

Each of the 32 optical units of the primary light collection system consist of a cryogenic Hamamatsu 5912-02MOD PMT seated behind an acrylic plate coated with a TPB-rich layer and surrounded by a mu-metal shield. Figure 33 shows a diagram of one unit (left) and a photograph of installed units (right). Past experiments have directly coated PMTs with wavelength shifter [7]. However, the MicroBooNE design separates the PMT from the wavelength-shifting plate for simplicity of quality control and installation. This proved important, as R&D indicated that TPB is particularly vulnerable to environmental degradation (see section 5.2.3). In this section, description is provided for each component of the optical unit, as well as for the overall assembly.

5.2.1 Photomultiplier Tubes and Bases

Reference [52] provides detailed information on the selection and testing of the 200 mm (8.0 in) diameter Hamamatsu R5912-02mod cryogenic PMTs employed in MicroBooNE. In this section a brief summary of the findings from this testing is presented.

The R5912-02mod employs a bi-alkali photocathode. Because the PMT is designed for cryogenic use, the R5912-02mod also features a thin platinum layer between the photocathode and the borosilicate glass envelope to preserve the conductance at low temperatures. While this allows the PMT to function below 150 K, absorption in the platinum reduces the efficiency of the

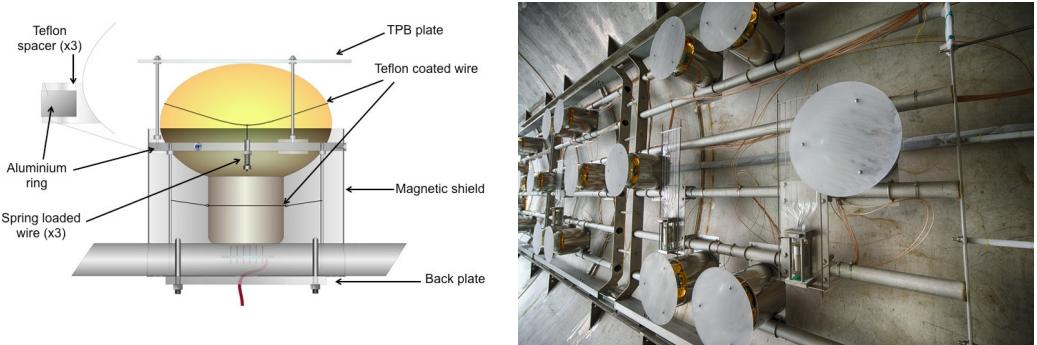


Figure 33. Left: diagram of the optical unit; Right: units mounted in MicroBooNE, immediately prior to LArTPC installation.

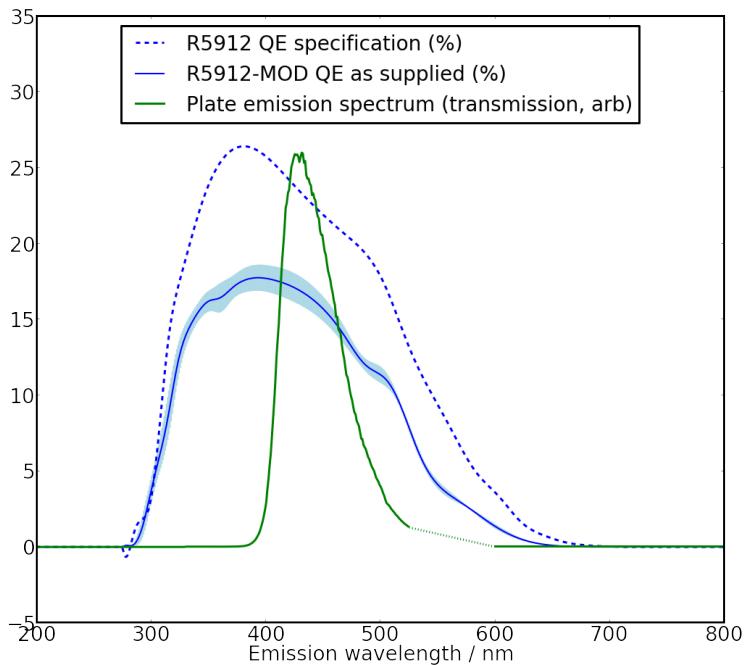


Figure 34. The specification for the non-platinum undercoated PMT from [62]. The blue band shows the mean and standard deviation of the four quantum efficiency curves provided by Hamamatsu for installed MicroBooNE PMTs. Also shown is the measured emission spectra of MicroBooNE wavelength-shifting coatings, discussed in section 5.2.2. [57]

1 PMT by 20%. Figure 34 provides quantum efficiency curves for these PMTs. The manufacturer's
2 specifications do not include the effects of the platinum photocathode coating, but a wavelength
3 dependent quantum efficiency was provided by Hamamatsu for 4 of the 32 installed PMTs. The
4 mean and standard deviation of these curves are shown in figure 34.

5 The R5912-02mod is a 14-stage PMT. The high gain at room temperature (10^9 at ~ 1700 V),
6 compensates for known reduced gain at 87 K in the liquid argon. The high gain also has the
7 additional advantage of allowing operation at lower than nominal voltage which reduces heat-loss

1 in the cables and the potential for high voltage breakdown at the feedthroughs.

2 The PMT base is designed such that the photocathode is grounded and the last dynode in the
3 chain is held at large, positive voltage. Thus the PMT bulb, which is closest to the LArTPC anode
4 plane, is at ground and does not disturb the electric field on the wires. The result is that the high
5 voltage (HV) can be provided and the signal can be extracted from the PMT using a single cable,
6 reducing the cable volume in the vapor region and removing a possible source of out-gassing water
7 impurity.

8 The flat PC-board base was made of Rogers RO4000-series woven glass-reinforced laminate,
9 which is the same material with the LArTPC cryogenic front end boards. Use of Rogers 4000 series
10 as the base material avoids the potential risk of contamination of the liquid argon. The Rogers
11 material has a similar temperature expansion coefficient as the surface mounted components which
12 enhances the reliability of the electronics assembly when operated at cryogenic temperatures. The
13 base is attached \sim 1 cm from the bottom of the PMT and is the closest possible distance of safe
14 approach to the PMT vacuum seal tip. A schematic and photos of the PMT base are provided in
15 reference [52]. The passive components include only metal film resistors and C0G/NP0 capacitors,
16 which have the minimum temperature coefficients, and the performance at cryogenic temperatures
17 were tested. Because the supplied HV and return signal share a single cable, the signal must be
18 split from the HV through an AC-coupling capacitor, as is discussed in section 5.2.5.

19 All installed PMTs were tested in a PMT test stand, both at room temperature and in liquid
20 nitrogen which is at 77 K. The details are described in reference [52]. In brief, the test stand
21 consisted of a light-tight 346 liter, liquid nitrogen filled dewar into which up to four PMTs could be
22 installed. The PMTs were immersed and maintained in the dark environment for up to three days
23 before most measurements of the dark rate and gain were performed. A fiber brought in light from
24 a pulsed blue LED, which was tested for linearity with bias voltage.

25 Among the important results from cryogenic testing were the following [52]:

- 26 • The PMTs could be ramped to voltage quickly in the cold environment, and after 30 minutes,
27 the gains were found to be stable.
- 28 • If the room temperature PMTs were immersed in liquid nitrogen, dramatic changes in the gain
29 were observed after initial turn-on. The PMT gain remained high in the first \sim 5 hours after
30 immersion, and then suddenly dropped by more than a factor of two, afterwards reaching a
31 stable value with a small drift.
- 32 • The PMT response showed good linearity up to 100 photoelectrons (PE), which was the
33 maximum attainable by the PMT test stand LED.
- 34 • The HV for each PMT was selected to produce a gain of 3×10^7 in liquid nitrogen, and was
35 typically chosen to be \sim 1300 V.
- 36 • The dark current plateaus extended up to 1800 V in liquid nitrogen, and the dark current is
37 higher in liquid nitrogen than at room temperature.
- 38 • The PMT performance depended on rate of the pulsed LED.

¹ No PMTs were rejected on the basis of the testing. However, there were three unexpected results to
² note here.

³ The first unexpected behavior was that, at room temperature, most of the PMTs showed
⁴ gains that were 10 to 30% higher than manufacturer’s specifications. As expected at cryogenic
⁵ temperatures, the gain is reduced by ~10% to 50%. To measure the PMT gain, the LED was set to
⁶ produce one to two PEs. The gain was found from the separation of the single PE peak from the
⁷ pedestal, where the single PE response was fit using the procedure described in [63].

⁸ Second, it was found that the PMTs are noisier in the cryogenic environment. It would typically
⁹ be expected that thermal emission is suppressed at cryogenic temperatures and one would expect
¹⁰ a lower dark current for PMTs operating in this regime. However, the dark current measured
¹¹ in the liquid nitrogen is higher than at room temperature. Although the cause is unknown, this
¹² phenomenon has previously been observed[64]. A proposed explanation is provided in [65].

¹³ Third, an LED pulse-rate-dependent gain shift was found during testing, as shown on figure 9
¹⁴ of reference [52]. This behavior is described qualitatively in [66]. With 10 kHz LED pulsing, the
¹⁵ gains of cold and dark-adapted PMTs were shown to steadily increase, requiring nearly 24 hours
¹⁶ from turn-on to stabilize. The effect was not observable at 10 Hz. This is relevant to MicroBooNE
¹⁷ because the cosmic muon rate in the MicroBooNE detector is ~ 5 kHz. Therefore a similar effect
¹⁸ is expected in the MicroBooNE detector. Preliminary results from measurements of the PMTs
¹⁹ installed in the LAr-filled MicroBooNE cryostat shows the expected effect.

²⁰ 5.2.2 Wavelength-Shifting Plates

²¹ In order to be sensitive to 128 nm VUV liquid argon scintillation photons, the optical assemblies use
²² a wavelength-shifting coating to convert this VUV light to visible wavelengths that are detectable
²³ by PMTs. In MicroBooNE, the active ingredient of this coating is TPB, an organic fluor which
²⁴ absorbs efficiently in the vacuum ultraviolet with an emission spectrum peaked around 425 nm [67]
²⁵ as shown in figure 32.

²⁶ The optical unit PMTs observe light transmitted through TPB-coated, 305 mm diameter acrylic
²⁷ plates (see figure 35). The coating consists of a 1:1 TPB-to-polystyrene ratio, with 1 g of each
²⁸ dissolved in 50 ml of toluene. A small amount of ethyl alcohol is added as a surfactant. The coating
²⁹ is applied to the acrylic plate in three layers by brush-coating. The solution dries in air at room
³⁰ temperature. This leads to a final layer which is oversaturated with TPB, and white crystals form
³¹ on the surface as the coating dries. The presence of surface crystallization gives the MicroBooNE
³² plates a white, opaque finish. Details of the process are described in reference [68].

³³ The emission spectrum in figure 34 was measured using a Hitachi F-4500 fluorescence spec-
³⁴ trophotometer with an incident beam of wavelength 270 nm, selected by a diffraction grating from
³⁵ a xenon lamp. A standard rhodamine dye calibration sample was used to correct for drift in the
³⁶ lamp and spectrometer. The measurement was made in transmission mode, and an artifact peak at
³⁷ a harmonic of the twice incident wavelength was observed between 525 and 600 nm. This region
³⁸ is omitted from the reported spectrum of figure 34.

³⁹ As shown in figure 34, although the absolute quantum efficiency for the platinum-undercoated
⁴⁰ cryogenic PMT is lower than the non-cryogenic version, the wavelength dependence is similar,
⁴¹ and overlap between the TPB emission spectrum and the sensitive wavelength range of the PMT
⁴² remains high. Using the measured TPB emission spectrum for plate coatings and the PMT quantum

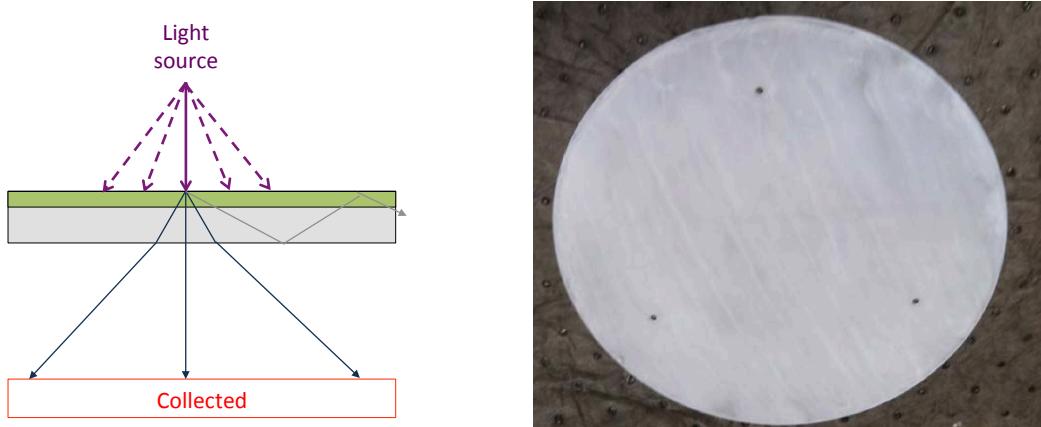


Figure 35. Left : Illustration of transmission mode, used by the optical units. Right: photograph of a coated plate.

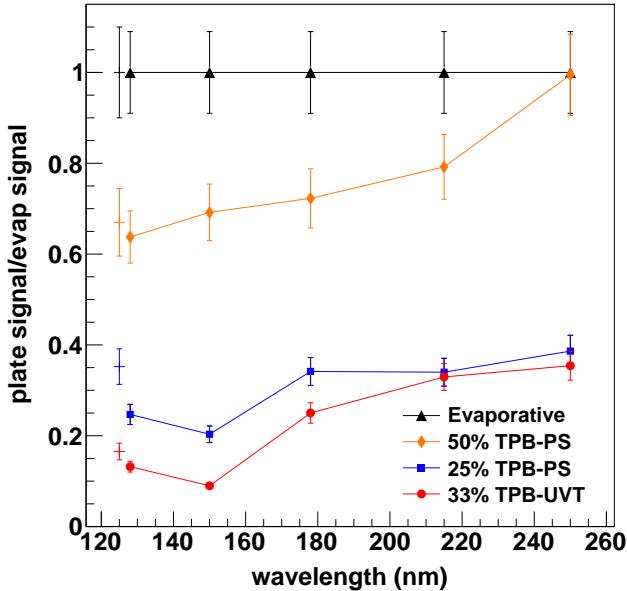


Figure 36. Measured efficiencies of various wavelength-shifting coatings, from [68]. In this plot, 50% TPB-PS is the MicroBooNE plate coating used in the optical units. The 33% TPB-UVT is the light guide coating for the secondary light collection system described below. Connected points were measured in a vacuum monochromator at room temperature, and non-connected points were measured in liquid argon with 128 nm scintillation light. All points are normalized to the performance of an evaporatively coated plate.

¹ efficiency curve provided, the spectrum-averaged PMT quantum efficiency is $15.3\% \pm 0.8\%$ per visible photon incident on the photocathode.

³ The wavelength-shifting performance of the coating was measured at 128 nm relative to evaporative coatings of the type studied in [69]. Coating efficiencies were measured as a function of wavelength between 128 and 250 nm using a vacuum monochromator at room temperature.

1 They were also measured in liquid argon using 128 nm scintillation light, relative to the same
2 evaporatively coated plate. These data are shown in figure 36, and more information about these
3 measurements can be found in [68].

4 The absolute efficiency of the MicroBooNE coatings can be obtained by multiplying the relative
5 efficiencies of figure 36 by the measured absolute efficiencies from [69], and accounting for the
6 temperature dependence in the wavelength-shifting efficiency of pure TPB reported in [70]. The
7 expected number of visible photons per incident 128 nm photon is 0.98 ± 0.17 for the MicroBooNE
8 plate coating.

9 **5.2.3 UV Light Protection for the Wavelength-Shifting Plates**

10 TPB coatings have been shown to degrade under exposure to ultraviolet light [50] through a
11 radical mediated photo-oxidation to the UV-blocker and photo-initiator benzophenone [71]. Several
12 measures were taken to ensure that degradation was minimized during the construction of the
13 experiment. The TPB powder and coated elements were stored in the dark at all times, with coated
14 plates and light guides being kept wrapped in foil and stored in a dark container before installation.
15 The detector construction area was covered with a UV blocking plastic [72], and test plates were
16 placed at various positions in the clean tent to check for degradation from stray light. After several
17 weeks of exposure, one test plate with a clear line of sight to the tent entrance demonstrated a
18 few percent degradation, and all others showed no observable loss of efficiency. The open end
19 of the MicroBooNE cryostat was shielded from light by a black curtain after installation, and the
20 feedthroughs of the cryostat were blocked when not in use to prevent stray light from entering. The
21 coated plates were the final component of the optical system to be installed into the detector to give
22 the minimum possible light exposure during the detector construction process.

23 **5.2.4 Cryogenic Mu Metal Shields**

24 The trajectories of electrons within the PMT can be deflected by the Earth's magnetic field. This
25 effect can be reduced or removed by surrounding the PMT with mu metal, a metal of low magnetic
26 permeability. Commonly used mu-metal fails to provide shielding at cryogenic temperatures. Two
27 types of cryogenic mu metal, Cryoperm 10 and A4K, both products of Amuneal, were identified
28 that did provide shielding at cryogenic temperature.

29 The mu metal shields were tested in the apparatus shown in figure 37. The system allowed for
30 the PMT to be positioned at an angle relative to the vertical axis, with the rotator set to 30 positions
31 from 0° to 348° . The set-up was on a dolly that allowed for rotation about the vertical axis. PMT
32 tests were performed in air and in liquid nitrogen. PMTs were dark adapted for 5 hours before
33 testing. A blue LED provided 1 to 2 single PEs of light through an optical fiber. The fiber was fixed
34 to the PMT mount such that the endpoint was stationary with respect to the tube as the system was
35 rotated.

36 The mean charge from the PMT as a function of angle was recorded. The error was primarily
37 systematic. The 10^6 LED pulses per data point give $< 1\%$ statistical error on the mean. However,
38 24-hour studies of a single point showed $\sim 5\%$ variations in collected charge due to gain drift.
39 Results showed that A4K and Cryoperm function essentially identically, to within the measurement
40 error, and so MicroBooNE chose A4K, based on significant cost savings. A plot of the relative
41 PMT gain variation versus angle from vertical, figure 38, shows that adding the shield significantly

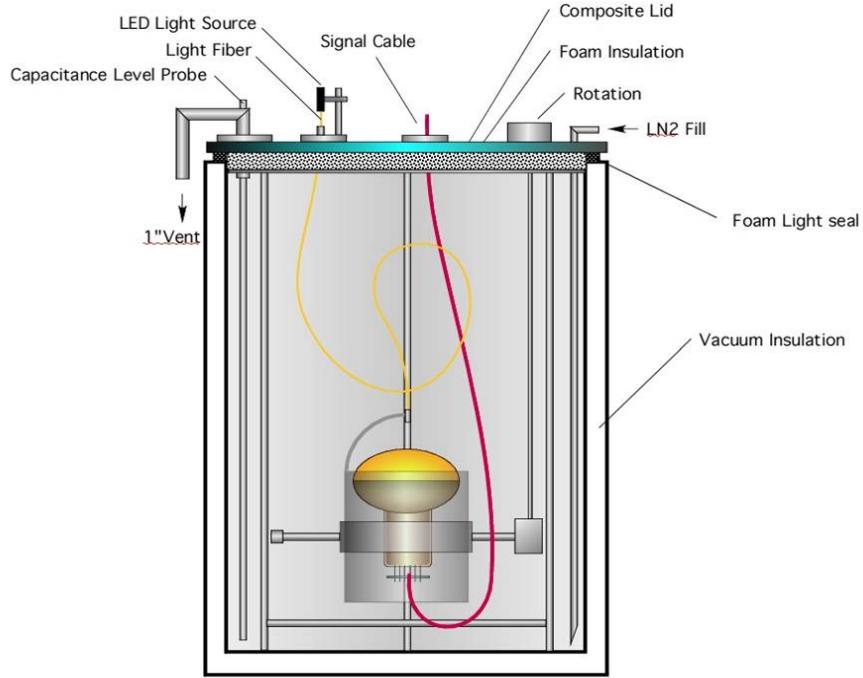


Figure 37. Schematic of the system used to study the mu-metal shields. The design allows rotation along all three axes.

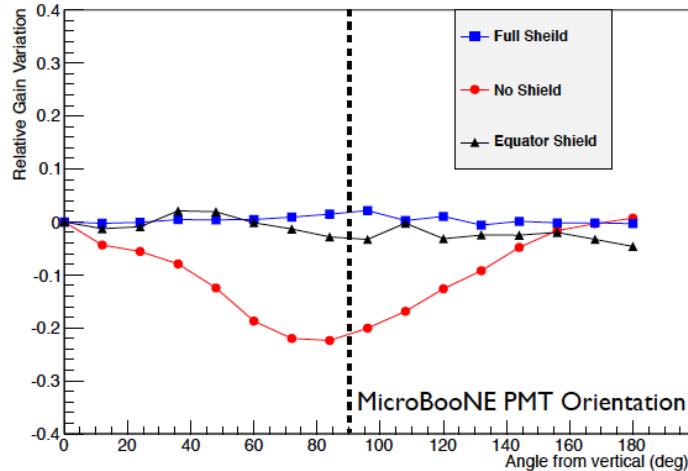


Figure 38. Red: Angular dependence of a PMT response with no shield; Black: for a shield that reaches the tube equator; Blue: for a shield that fully covers the tube to the zenith.

improves PMT performance. This plot shows the effect of the A4K shield with the height aligned with the equator of the PMT (black) and aligned with the zenith of the bulb (blue), compared to no shield (red) as a function of PMT azimuthal angle. Given the 5% systematic error, the two shield positions are indistinguishable.

As a result of these tests, the A4K shields were designed to extend just past the equator of the

- ¹ PMT. The shield has small holes in the backplate that allows the PMT cables to exit the shield.
- ² MicroBooNE is the first LArTPC to use cryogenic mu metal shields in its light collection system.

³ 5.2.5 Implementation of the Primary System

⁴ The light collection system is composed of optical units assembled as shown in the top picture
⁵ of figure 39. The PMT is seated within the mu metal shield on three teflon pads attached to an
⁶ equatorial support ring. The neck of the tube slides inside a loose wire guide-loop that prevents the
⁷ PMT from tipping. The PMT is held within this assembly using teflon-encased wires that extend
⁸ across the bulb and connect to wire hooks attached to the equatorial ring with stainless steel springs.
⁹ Legs extending from the support at the equator are screwed into a backplate for final mounting.
¹⁰ Concern about differences in contraction of the materials led to this design which holds the PMT
¹¹ in place, but with only moderate rigidity. The units were tested to ensure the PMTs would not be
¹² displaced during installation and filling. Three posts extend upward from the equatorial ring to
¹³ hold the plate 3.0 cm above the apex of the bulb. The optical units slide into a cylindrical mu-metal
¹⁴ shield, which screws into the equatorial ring. The unit is then mounted on stainless steel back-plates
¹⁵ affixed to a support rack, as shown in the bottom picture of figure 39.

¹⁶ The support rack consists of five stainless steel components, or modules, for ease of installation.
¹⁷ Each module has vertical height 1.83 m and horizontal length 2.07 m, resulting in a total horizontal
¹⁸ length of 10.36 m. Unlubricated Thomson bearings fitted to the lower edge allow each module
¹⁹ to slide into the cryostat on rails mounted in the vessel. The system was designed to allow the
²⁰ light detection system to slide into the vessel after the LArTPC was installed. However, in the
²¹ end, scheduling permitted installation of the system before the LArTPC installation. This had the
²² advantage of making installation and surveying easier, but the drawback that the system would be
²³ exposed to UV light for a long period. Therefore, the units were installed with dummy clear acrylic
²⁴ plates, and the TPB coated plates were installed only just before the LArTPC was moved in and
²⁵ the detector could be easily protected from light. During optical unit installation, each rack module
²⁶ was supported by a temporary mounting rail. The optical units were then mounted in positions
²⁷ chosen to avoid obstruction by the LArTPC cross-bars, as shown in figure 30. As the units were
²⁸ mounted and slid into the cryostat, the cables were loosely tied to the bars of the rack for support
²⁹ and constraint.

³⁰ The “splitter” circuit, located outside of the cryostat, is shown in figure 40. The splitter
³¹ separates the HV of the PMT from its output signal which is subsequently split into a high-gain
³² (HG) and a low-gain (LG) channel. The HG and LG channels respectively carry 18% and 1.8%
³³ of the output signal. This allows a wide dynamic range for ADC readout of the PMT pulses. The
³⁴ capacitance was chosen to minimize reflections, since the bases are not back-terminated. The HV
³⁵ is supplied to the splitter using BiRa Corporation, Model 4877PS modules.

³⁶ The PMT cable system delivers HV and returns signals between the external splitter and the
³⁷ optical unit in the cryostat. A single cable runs from an external connector, through the feedthrough
³⁸ filled with Stycast 2850 epoxy [73], into the cryostat and to the PMT base. The RG316/U coaxial
³⁹ cable has $50\ \Omega$ impedance. Cables were terminated with Pasternack PE4498 SHV to accommodate
⁴⁰ that the cable carries HV to the tube as well as signals from the tube. The cable carries an AC
⁴¹ voltage rating of 1100 V; however tests showed the DC rating to be at least three times higher and so
⁴² suitable for this use. The cables were routed through feedthroughs consisting of a pipe filled with

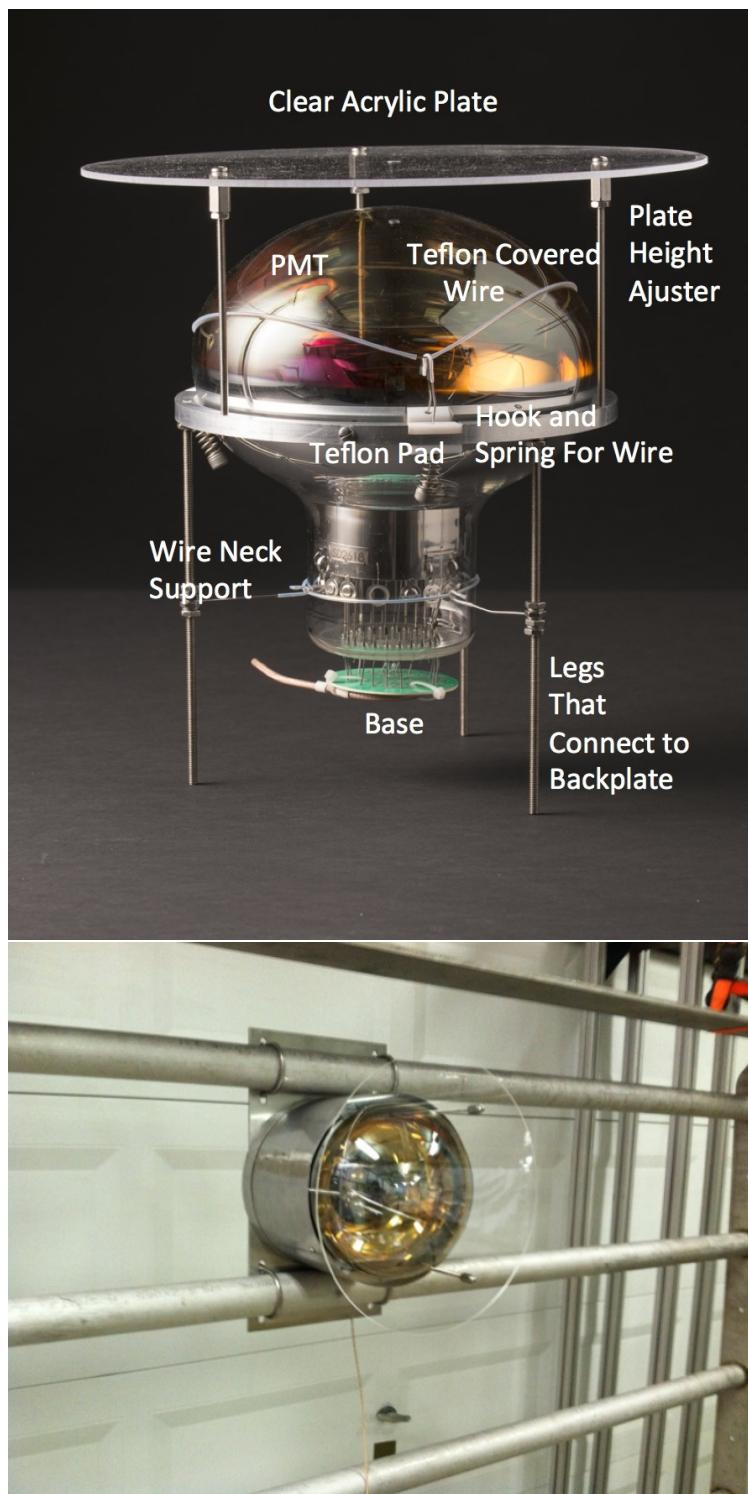


Figure 39. Top The optical unit mount internal to the shield, with components labeled; Bottom: Unit mounted on rails. The clear plates were replaced with TPB-coated plates immediately before LArTPC installation, as discussed in the text [47, 52].

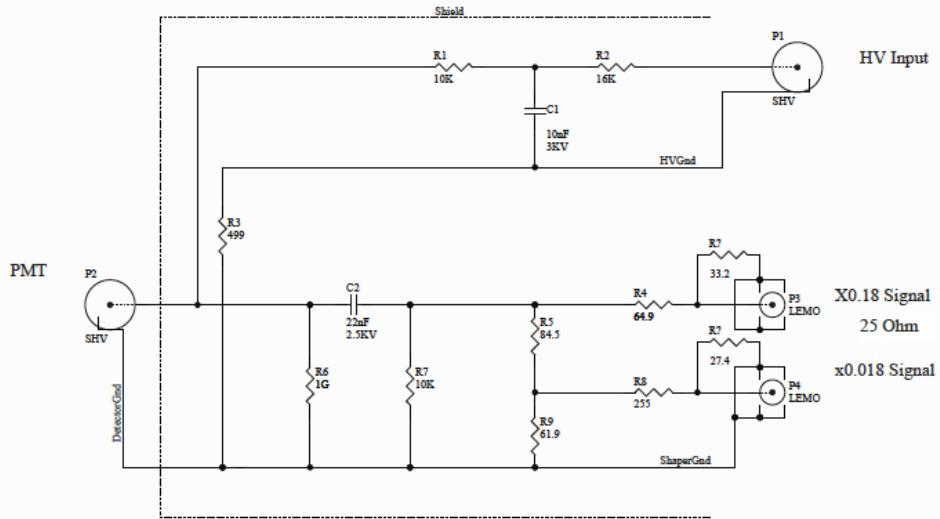


Figure 40. The “splitter” circuit. The circuit connects the HV source to the PMT. It also provides a pathway for signal pulses from the PMT to reach the readout electronics via an AC-coupling capacitor (C_2). The signal is split into two copies, one provided with an attenuation factor of 0.18 and another at 0.018. Both signal sources are recorded by the readout electronics in order to provide two dynamic ranges.

solidified epoxy mounted on a conflat disk. On the warm side of the feedthrough, the cables were terminated at a patch panel with SHV connectors. The SHV connector impedance has a negligible effect on the 20-30 ns PMT signals. SHV cables connect the patch panels to the splitters. The impedance of every channel was tested at the feedthrough patch panel for a stable and correct value for the base resistance, which was $4.04 \pm 0.02 \text{ M}\Omega$.

5.3 PMT Testing and Quality Assurance

The PMTs were tested and characterized before installation in the detector in the “Bo” cryostat at Fermilab. The Bo cryostat is a 250-liter vacuum-insulated vessel with an inner diameter 56 cm and a depth of 102 cm used for R&D studies, and, relevant to this paper, a vertical slice test of the MicroBooNE optical units. The system is described in detail in reference [57]. The cryostat can be filled with purified LAr with parts per billion level contamination of oxygen and water and parts per million level of nitrogen. The light collection system was tested with light from visible (420 nm) and UV (250 nm) LEDs piped in via fiber, as well as scintillation light from ^{210}Po alpha sources and cosmic rays.

A vertical slice test (VST) of the MicroBooNE optical system was performed in the Bo Cryostat. The slice consisted of two PMTs with base electronics, mu-metal shield, TPB plates, cable feed throughs, splitters, the HV power supply and the interlock system. Tests were performed without and with the mu-metal shield. The test made use of the data acquisition components described in section 6, including the shaper, FEM, trigger card, control card, and server. The MicroBooNE trace impurity monitors were also used.

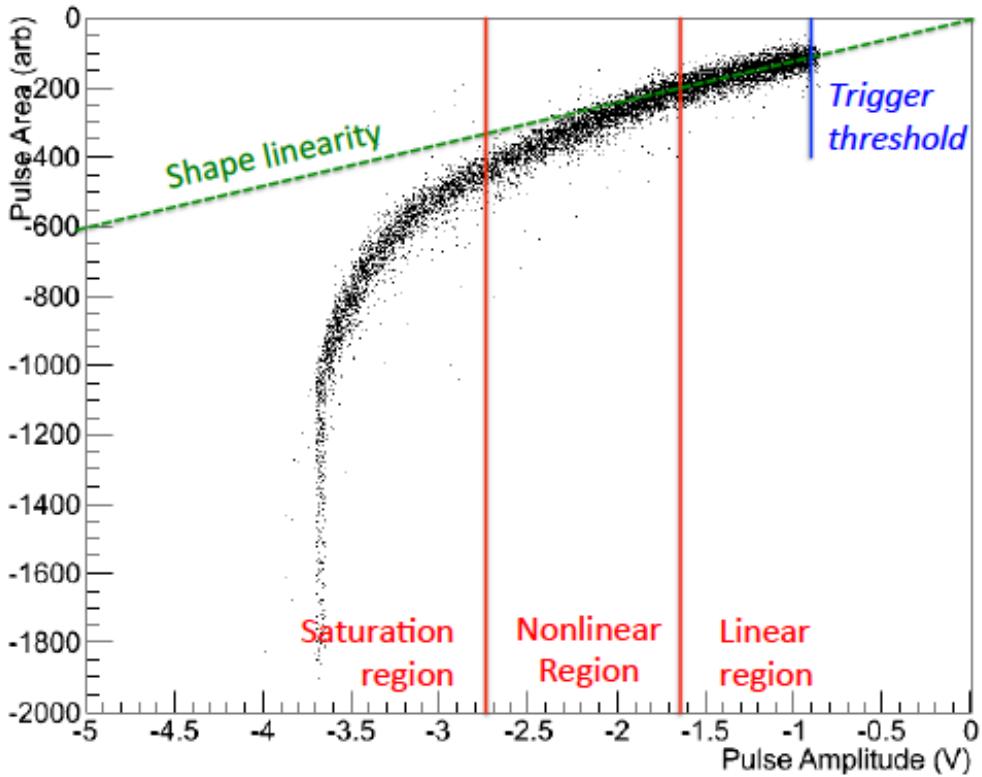


Figure 41. As shown by the VST, linearity of cosmic-ray induced PMT pulses is maintained up to amplitudes of around 1.7 V (300 PE), and amplitude saturation occurs at 3.7 V (670 PE) [57]

1 The VST informed the final design of many components, as well as producing results relevant
 2 to understanding the running conditions and performance expectations. For example, during studies
 3 of the response of the slice to 128 nm scintillation light from the alpha source, valuable information
 4 was gathered on the single photo-electron dark rate and cosmic ray rate that could be scaled to the
 5 MicroBooNE detector expectations. As a second example, these runs allowed characterization of
 6 the pulse shape nonlinearities of the optical units, as seen in figure 41. These were shown to be
 7 significant at \sim 300 PE in pulse amplitude. Full amplitude saturation occurred at \sim 670 PE. Thus,
 8 it is concluded that for pulses of more than 300 PE, pulse shape cannot be described by linear
 9 superposition of single photoelectron pulses.

10 **5.4 Secondary System: Acrylic Light Guides for R&D**

11 A secondary light collection system consisting of four lightguide paddles was also installed. This
 12 concept has several advantages for future large detectors such as DUNE. First, the collection area
 13 per channel is larger than the optical units, providing more coverage for the same number of
 14 electronics channels, cables, and feedthroughs. Second, the detectors have a flat profile so they can
 15 be slid between chambers in a multi-LArTPC detector, minimizing space requirements of the light
 16 collection system. In the case of MicroBooNE, the design gradually guides light in bent acrylic
 17 bars to a PMT. This design was an early alternative to a perfectly flat design that guides the lights

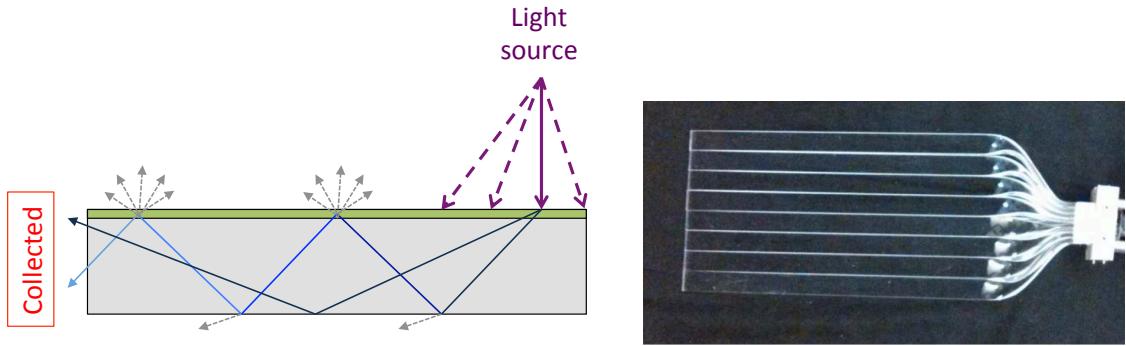


Figure 42. Left : Illustration of guiding mode, used by the paddles. Right: photograph of a coated paddle.

to SiPMs. Running this system will provide long-term information on performance of lightguide based systems. It also enhances the MicroBooNE dynamic range, since the lightguide detectors saturate at a much higher light level than the optical units.

In the case of the lightguides, the 128 nm light is absorbed and shifted by a clear wavelength-shifting coating, and the re-emitted light is guided to a 5.08 cm (2.0 in) Hamamatsu R7725-MOD PMT, as illustrated in figure 42, left. The installed paddles consist of six bars. A photograph of one coated paddle with eight bars is shown in figure 42, right. The active length of each bar is 50.8 cm. This system was added for R&D purposes and made use of 8 spare channels available of HV, cables, feedthroughs, and electronics. The impedance of the lightguide PMTs, tested at the feedthrough, was $4.89 \pm 0.01 \text{ M}\Omega$. As shown in figure 30, each paddle is installed next to an optical unit for direct comparison of performance.

The coating requirements for plate assemblies and light guides are different, and so the composition and coating methods for each were separately optimized. In the case of the light guide coatings, the figure of merit is the light emitted in guided mode. Guided mode light is the light that is detected at one of the ends of a test sample, which is orthogonal to the illuminated face of the sample. In addition to the wavelength-shifting efficiency of the active layer, the detected light yield is affected by the reabsorptive and scattering losses in the coating as visible light propagates along the bar. The light guide assemblies have a TPB coating of 33% TPB to 67% UVT acrylic by mass, also with ethanol surfactant. The coating is applied as a single layer and the TPB remains suspended in the acrylic matrix as the coating dries, leading to a smooth, visibly transparent surface. The performance and attenuation behavior of similar light guides to those installed in MicroBooNE were studied experimentally in [51]. The reported non-exponential attenuation suggests that surface losses dominate over bulk losses as the attenuation mechanism, and that the fractional loss per reflection within the light guide is of order 2-3% [71].

In the light guide coatings, the TPB is suspended in an acrylic matrix which leads to a slight broadening of the emission spectrum compared to the spectrum from the plates. This is an expected effect—TPB fluorescence has been shown to have dependence upon its microenvironment [70, 74–77], and reference [70] demonstrated spectral broadening in the presence of a polystyrene substrate.

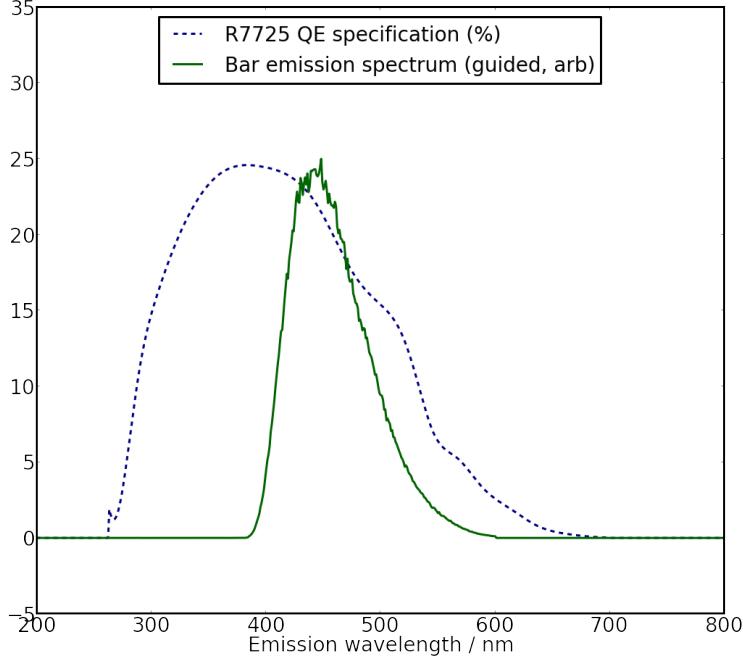


Figure 43. Measured emission spectra of the light guide coating in guided mode, and the R7725 quantum efficiency. Only the quantum efficiency of the non-undercoated PMT model is shown, from [78].

Using the monochromator described for the plate spectrum studies, the light guide coating spectrum was measured in guided mode. A 10 cm section of light guide, with the incident beam perpendicular to the TPB coated surface, was used. Figure 43 shows the results of this measurement. Based on reference [70] it is expected that the emission spectrum for the light guide coating, with TPB embedded in the substrate, will not change significantly as it cools to 87 K. The expected efficiency for the light guide coating is found to be 0.25 ± 0.05 emitted visible photons per incident 128 nm photon.

5.5 Calibration

The flasher system for the optical units and the light guides is described in reference [55]. This system was developed to check the timing of the installed optical units, exercise the optical units during construction and commissioning, and to calibrate them once the detector is operational. The reference provides engineering drawings and details. The system is briefly summarized here.

A control board pulses an array of 400 nm LEDs, each of which is coupled through an optical feedthrough to 10 m optical fibers within the cryostat. The custom feedthrough/patch panel design encases the fibers in Arathane CW 5620 blue with HY 5610 hardener. The internal fibers are Molex FVP polyimide fibers with diameter of 600 nm, cladding of 30 μm , and an additional buffer layer of 25 μm . Each PMT has an individual calibration fiber. Each fiber is routed along the rack and attached to the PMTs by a fiber holder constructed of an aluminum standoff with a nylon-tipped set screw at a distance of about 5 cm from the PMT glass.

Figure 44 shows the results of flasher tests on all 32 optical units and four PMTs on the paddles. Time is along the x axis and PMT channel number is along the y axis. The white region separating

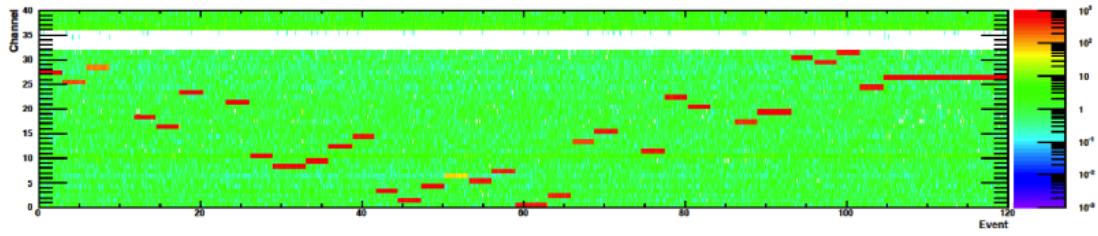


Figure 44. Charge detected in each PMT due to flasher tests, shown as a function of time. Each PMT is flashed for a short period. The white band indicates unused channels.

1 the optical unit PMTs from the paddle PMTs represents unused channels. The colored bars indicate
 2 charge detected in the PMTs during flashing. One can see that all tubes respond properly.

3 **5.6 Coupling of PMT Signals to the Anode Wires**

4 The ICARUS detector has observed an unexpected cross-talk between their PMTs and collection
 5 wire plane [79]. The Argontube test stand at the University of Bern also observed this effect [80].
 6 Therefore, while the LArTPC was under final testing, before being rolled into the cryostat, an
 7 experiment was devised in order to determine the implication of this cross-talk for MicroBooNE.
 8 One of the production 200 mm (8.0 in) Hamamatsu R5912-02mod PMTs was placed 12.7 cm from
 9 the LArTPC collection wire plane. This PMT was encased in a dark box with an optical fiber
 10 delivering light from a LED flasher. The collection wire plane was read out using a test-stand that
 11 reflects the final data acquisition design.

12 A clear signal was observed on the collection plane when the LED fired. The magnitude of the
 13 signal was characterized as a function of photo-electron count (figure 45) and separation between
 14 the PMT and wire-plane. The signal induced on the collection plane was estimated to be no more
 15 than ~ 10 ADC counts for a cosmic ray under normal operating conditions. This was reduced to
 16 ~ 2 ADC counts when the PMT was encased in the μ -metal shields. This was deemed an acceptably
 17 small amount that further shielding was not required.

18 The effect appears to saturate with photo-electron count, and is reduced when shorting the
 19 resistors between the anode and last dynode. A later test with Argontube showed that the signal
 20 was drastically amplified when using capacitors not able to withstand cryogenic temperatures in the
 21 PMT base [80]. This suggests that the effect is electrostatic in nature, and probably due to capacitive
 22 coupling between the PMT and wire-plane. This is also suggested by the similarity between the
 23 signals observed and those produced by anode-coupled readout of a light collection system [56].

24 **5.7 Initial Performance of the MicroBooNE Light Collection System**

25 The system was first powered on after the cryostat had been filled with liquid argon. All the PMTs
 26 in both the primary and secondary system were found to be operational. Figure 46 shows the
 27 waveforms for all the PMT channels around the time that a large pulse, potentially from a cosmic
 28 ray muon, is observed. After initial checks of the system's health, the single photoelectron response
 29 of each PMT was set such that a one photo-electron pulse has an amplitude of 20 ADC counts as
 30 seen by the PMT readout system. The flasher LED system, described previously, is used to set this

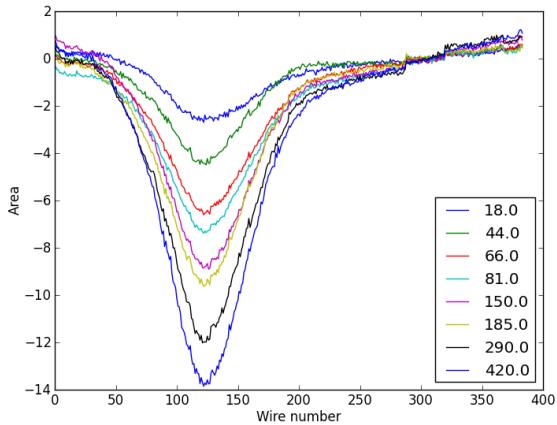


Figure 45. Area of signal pulses recorded on collection plane (in arbitrary units) as a function of wire number (arb. offset). The legend indicates the calibrated photo-electron count. The signal pulses are averaged over 50,000 repetitions. The plots are pinned together at wire 300, as the distribution baseline fluctuated over the course of the experiment due to intermittent noise.

1 response. Figure 47 shows a candidate single photoelectron pulse following arrival of a TTL logic
 2 pulse driving the LED flasher system. The LED light level is set so that the majority (about 80%)
 3 of waveforms see no response in the region where pulses from the LED are expected to occur. This
 4 ensures that for windows with pulses, the pulses are of single photoelectrons. Figure 48 shows an
 5 example of the area vs. maximum amplitude of such pulses seen by a PMT during the flasher runs.

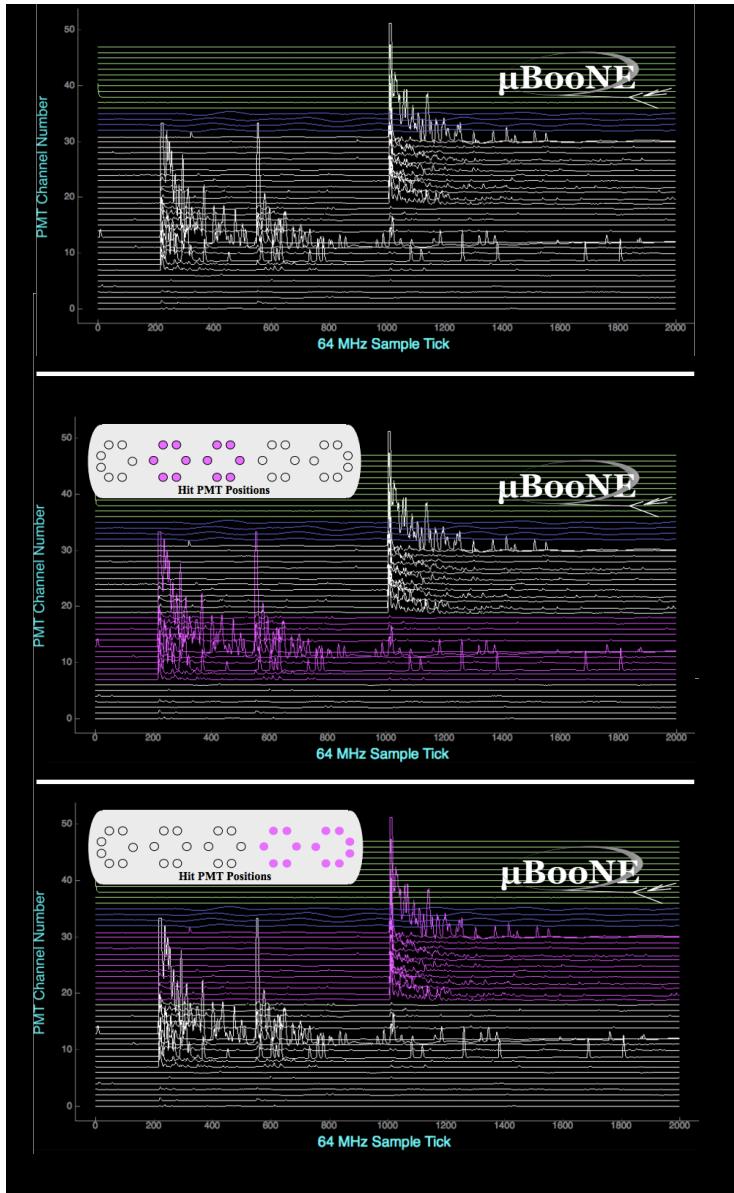


Figure 46. Top: Example waveforms from all 32 PMTs of the primary light collection system over a 31.25 microsecond readout window. The waveforms from the PMTs are in white. The blue waveforms are from the secondary lightguide PMTs, which are off in this picture. Waveforms from channels reserved for logic inputs are shown in green. In this image, the PMTs see two successive flashes of light at different parts of the detector. Middle and Bottom: Magenta highlights the early and late pulse, with the insets showing the PMTs which fired. Each pulse is likely from two cosmic ray muons traveling through the detector.

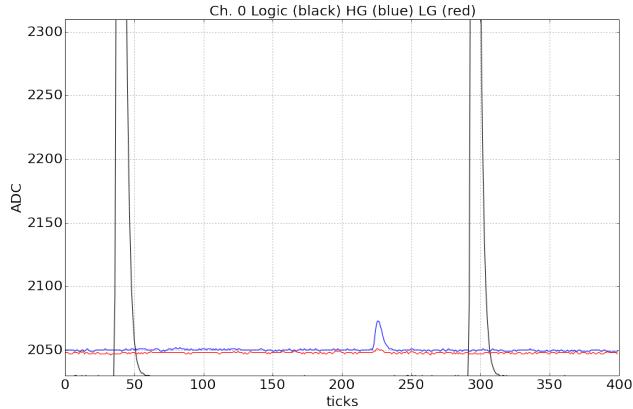


Figure 47. Example waveform captured by the PMT readout electronics during single photoelectron calibrations. The black waveform is of logic pulses that mark time at which an LED in the flasher system is driven. After some delay coming from PMT cable lengths and the flasher system, a candidate single photoelectron pulses is seen.

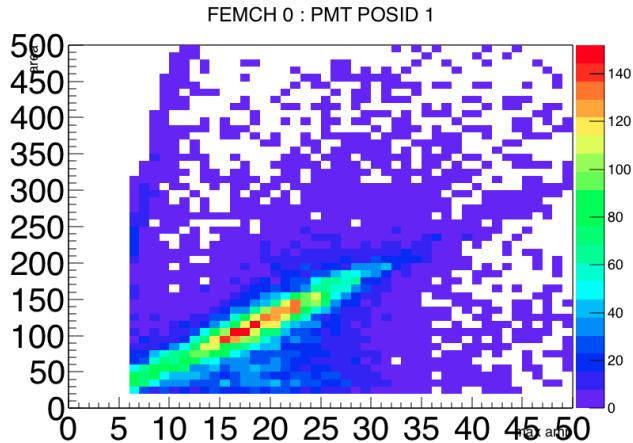


Figure 48. Distribution of the maximum amplitude and charge of pulses collected during an LED flasher calibration run. The central distribution of events is due to single photoelectron pulses.

6 Electronics and Readout Systems

The analog signals that develop on a LArTPC during its operation must be amplified, digitized, and written to disk for use in analysis. Custom low-noise electronics that are capable of operating in the liquid argon environment have been developed for this purpose in MicroBooNE. The data from these LArTPC electronic channels, as well as from the PMTs, is sent to a readout system that digitizes and organizes the information before passing it along to a Data Acquisition (DAQ) system that stores it on disk. The stages of signal processing are illustrated in figure 49. The following subsections describe the LArTPC cold electronics, the LArTPC and PMT readout electronics systems, and the DAQ system in more detail. Details of triggering capabilities available in MicroBooNE are also provided in this section.

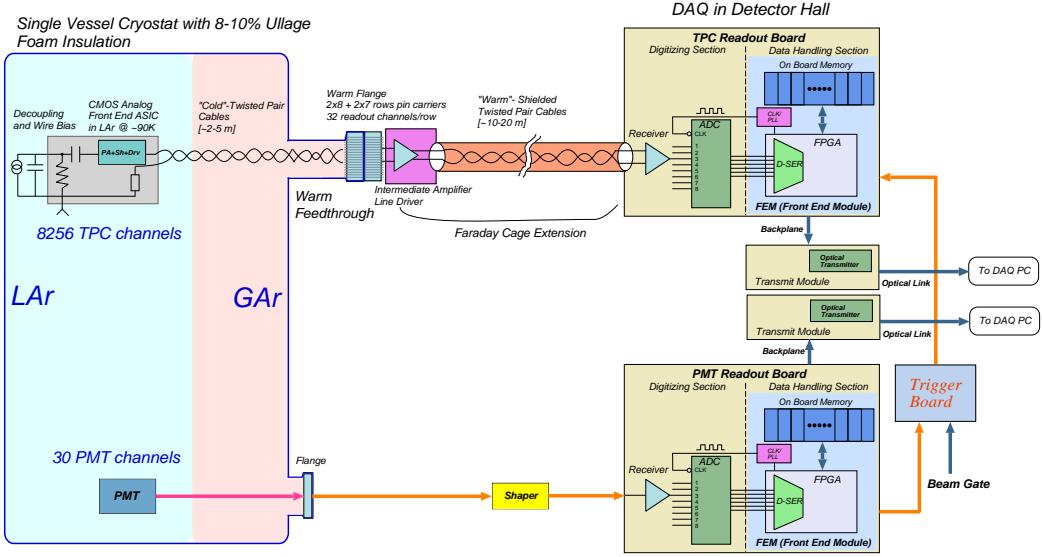


Figure 49. MicroBooNE LArTPC and PMT signal processing and readout stages.

6.1 Cryogenic Low-Noise Electronics

To obtain optimum detector performance, MicroBooNE uses cryogenic low-noise front-end electronics for readout of the LArTPC. To reduce electronic noise, the interconnection length between the LArTPC wires and preamplifier should be as short as possible thus minimizing the total capacitance seen at the preamplifier input. To accomplish this, the analog front-end ASICs, which include a preamplifier, shaper, and signal driver are located inside the cryostat in addition to the wire bias voltage distribution system, decoupling capacitors, and calibration networks. The front-end ASIC and associated circuits are implemented on a cold mother board which is directly attached to wire carrier boards on the LArTPC itself. Cold cables are used to transmit output signals from cold motherboards to warm interface electronics installed on the top of signal feed-through flanges.

6.1.1 CMOS ASIC

The analog front end ASIC is designed in 180 nm CMOS technology, which integrates both the preamplifier and shaper on a single chip. Each chip has 16 channels to read out signals from 16 wires. Each channel also has a charge injection capacitor for precision calibration. In MicroBooNE, the shaper has four programmable gain settings (4.7, 7.8, 14 and 25 mV/fC) and four programmable peaking time settings (0.5, 1.0, 2.0 and 3.0 μ s) that provide increased flexibility to the readout system. The ASIC also has programmable baseline settings (200 or 900 mV) to accommodate different detection wire configurations: either collection or induction plane. It has a selectable AC/DC coupling mode with a 100 μ s time constant for the AC coupling mode which can be used to reduce low frequency noise. The ASIC also has built-in band-gap reference and temperature sensors to facilitate biasing and monitoring.

The CMOS ASICs consume only 6 mW/channel in its default configuration. The front end ASICs of the entire detector generate 50 W of heat load that is easily handled by the cryogenics

1 system. Design guidelines that constrain the electric field and the current density to address the
 2 lifetime of CMOS devices operated at cryogenic temperatures have been applied to every single
 3 transistor (total ~15,000 transistors) in the ASIC design. A picture of the layout of the CMOS ASIC
 4 is shown in figure 50. Test results agree well with simulations and indicate that the analog and the
 5 digital circuits (including the digital interface) operate as expected in the cryogenic environment.

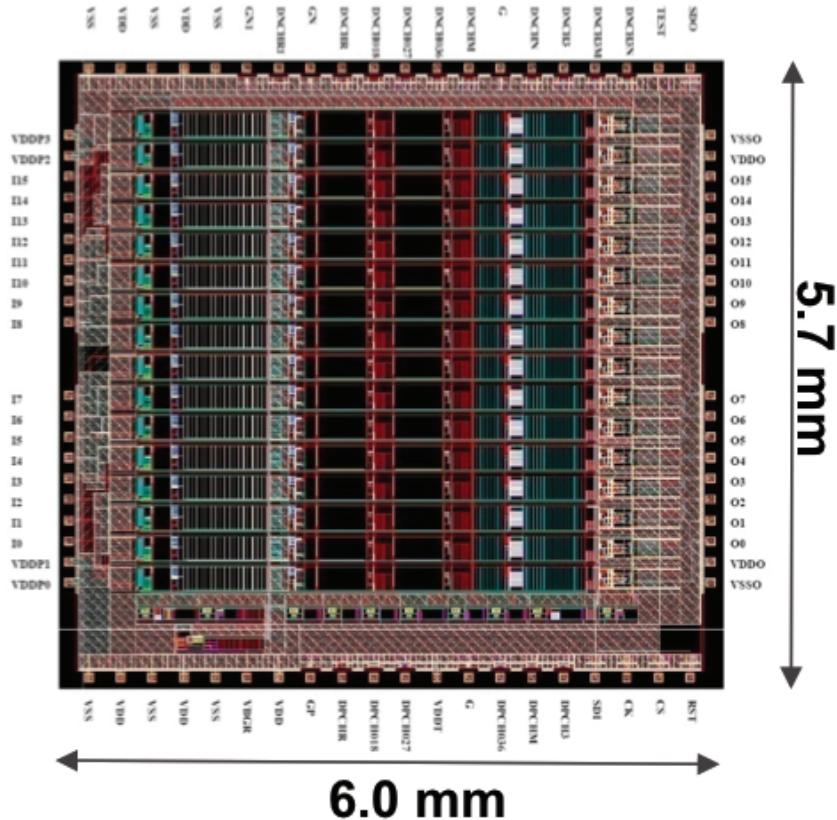


Figure 50. Layout of the CMOS analog front end ASIC

The MicroBooNE LArTPC has 8,256 readout channels and a total of 516 CMOS ASICs are required to fully instrument the detector. The production testing of the CMOS ASICs required two steps: both a warm and cold test. The warm test was performed with a dedicated test board housed in a Faraday box containing a socket to house the ASIC for ease in chip exchange. All programmable parameters (gain, peaking time, baseline, AC/DC coupling etc.) were exercised with the warm test setup for careful screening at room temperature. The yield of the warm testing of the ASICs was 89%. ASICs must have passed the warm test before going through cold testing. The cold test was performed with a dedicated test board containing 6 sockets to facilitate testing of multiple chips in liquid nitrogen at the same time. A total of 201 ASICs went through cold testing with a yield of 97%. Based on this high yield, it was decided not to continue the cold screening test on the rest of the production chips, as they were tested cold after being installed on the motherboard. After enough (~600) ASICs passed the production screening test, they were sent to an assembly house to equip the cold motherboard.

6.1.2 Cold Motherboards

A cold motherboard was designed to house the MicroBooNE CMOS ASICs. In this capacity, the motherboard provides signal interconnections both between the detector wires and preamplifier inputs as well as between the driver outputs and cold cables to the signal feed-through. The cold motherboard design provides sufficient protection of the ASICs against electrostatic discharge during installation. It also provides a calibration network and bias voltage distribution for the wire planes. Specifically, a calibration signal enters the cryostat via a feed-through and reaches the preamplifiers through the motherboard. Each preamplifier channel in the ASIC has a built-in switch to individually cycle the calibration injection. The bias voltage reaches the LArTPC wires via a two-fold redundant path on the motherboard that allows the detector to operate normally even if one bias voltage channel fails. As with the PMT bases, the cold motherboards use Rogers 4000 series as the PC-board material due to the cleanliness and thermal properties previously described in section 5.2.1. The different positions of the wire attachments along the top and sides of the LArTPC requires 2 types of cold motherboard. The top version of the motherboard has 192 readout channels that includes 96 Y channels, 48 U channels, and 48 V channels. The side version of the motherboard has 96 readout channels that are either U or V channels. A picture of the top version of a cold motherboard with 12 mounted ASIC chips is shown in figure 51.



Figure 51. Top version of cold motherboard with 12 ASIC chips, including 6 chips mounted on top layer and 6 chips on bottom layer.

The MicroBooNE LArTPC required a total of 36 top version motherboards and 14 side version motherboards to instrument the full detector. A test stand was built for testing of the front end electronics. This test stand included a full readout chain from the cold motherboard, cold cable, signal feed-through, warm interface electronics, warm cable and Receiver ADC board to a DAQ board based on a Xilinx ML605 FPGA evaluation board which sends acquired data to a PC over a Gigabit Ethernet. The production test of each motherboard also involved both a warm and cold test. Both tests used the same test stand, except the motherboard was placed in a Faraday box for the warm test and submerged in a liquid nitrogen dewar for the cold test. Noise, gain, peaking time, and linearity parameters were measured in both warm and cold to screen the motherboard. Motherboards had to pass both tests before being installed on the detector.

1 6.1.3 Cold Cables

2 Cold cables transmit the detector signals from the cold motherboard to an intermediate amplifier
3 on top of the signal feed-through and distribute power to the CMOS ASICs. The cold cable is a
4 custom-built 32-pair twisted pair flat ribbon cable with Teflon FEP insulation and $100 \Omega (\pm 10\%)$
5 impedance, using AWG 26 stranded wire with silver-plated copper. Custom designed shells with
6 jack screws used in the cable assembly ensured proper alignment of the insertion on the signal feed-
7 through pin carriers. Cold cables of two different formats were assembled by an assembly house:
8 signal cables and service cables. Signal cables are used to transmit amplified detector signals while
9 the service cable is used to transmit calibration pulses and slow control/monitoring signals. Signal
10 cables were produced in three different lengths: 203 cm, 254 cm, and 457 cm to accommodate
11 the different lengths between the cold motherboards and the signal feedthroughs, while the service
12 cables were produced in two different lengths: 254 cm and 457 cm. All of the cold cable assemblies
13 were tested with a cable tester in the assembly house before they were shipped out. In addition, 10%
14 of the cold cables were tested in the test stand at BNL to confirm the quality of the cable assembly.

15 6.1.4 Electronic Calibration

16 The MicroBooNE cold electronics include a precision charge calibration system. Through the cold
17 cable and calibration network on the motherboard, a calibration signal enters the cryostat via a
18 feed-through and reaches the preamplifiers. A built-in switch in the ASIC makes it possible to
19 power cycle the calibration injection for every channel individually. The electronics calibration is
20 based on charge injection through known capacitances (180 fF) in the ASIC. This system enables
21 gain (charge sensitivity) calibration, verification of sense wire integrity and noise measurements.
22 The built-in electronics calibration capability is an important tool in testing and characterizing the
23 overall performance of the detector readout system. It was extensively used in the cold electronics
24 production testing and the electronics checkout during installation, commissioning, and data taking.

25 6.1.5 Performance Tests

26 The development of the analog front end ASICs was initiated using 180 nm CMOS technology
27 and 300 K models, though the performance parameters are extracted at 77 K. CMOS was found
28 to function at cryogenic temperatures with increased gain and lower noise. The noise, gain, and
29 pulse shaping were found to be as expected in evaluation tests of the ASICs. Extensive testing of
30 the ASICs mounted on the motherboards was performed; these tests were done in liquid nitrogen
31 rather than at room temperature, since noise levels and characteristics of the ASIC performance in
32 liquid nitrogen are similar to the performance in liquid argon. Thus, cold tests were performed on
33 all production cold motherboards fully populated with 12 chips. A total of $\sim 2,200$ chip-immersions
34 were accumulated in liquid nitrogen without any failures due to thermal contraction or expansion.

35 The test results show the noise of the front end readout electronics system decreasing uniformly
36 for all 768 channels from $\sim 1,200 e^-$ at 293 K to less than $600 e^-$ at 77 K with 150 pF detector
37 (sense wire) capacitance. A plot of noise versus temperature of 12 ASICs for a total of 192 channels
38 is shown in figure 52. The response of the front end electronics exhibits excellent uniformity at
39 cryogenic temperatures. As shown in figure 53, the gain variation of a cold motherboard with 12
40 ASICs is only 7% peak-to-peak across 192 channels. The spread of the gain variation is only 1%
41 of the gain setting.

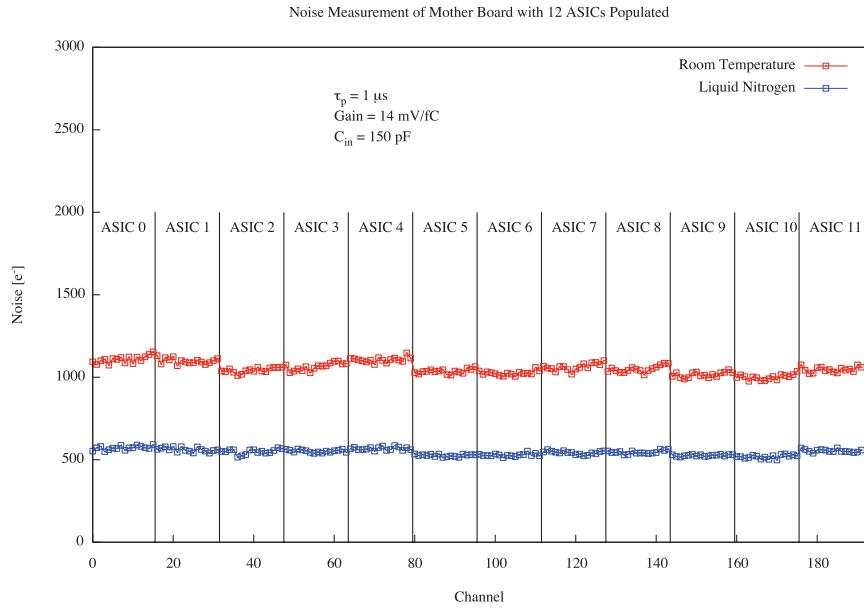


Figure 52. Plot of noise vs. channel number for 192 channels (12 ASICs) and at two different temperatures. Noise is $\sim 1,200 e^-$ at 293 K, and $\sim 550 e^-$ at 77 K with $150 \text{ pF } C_d$

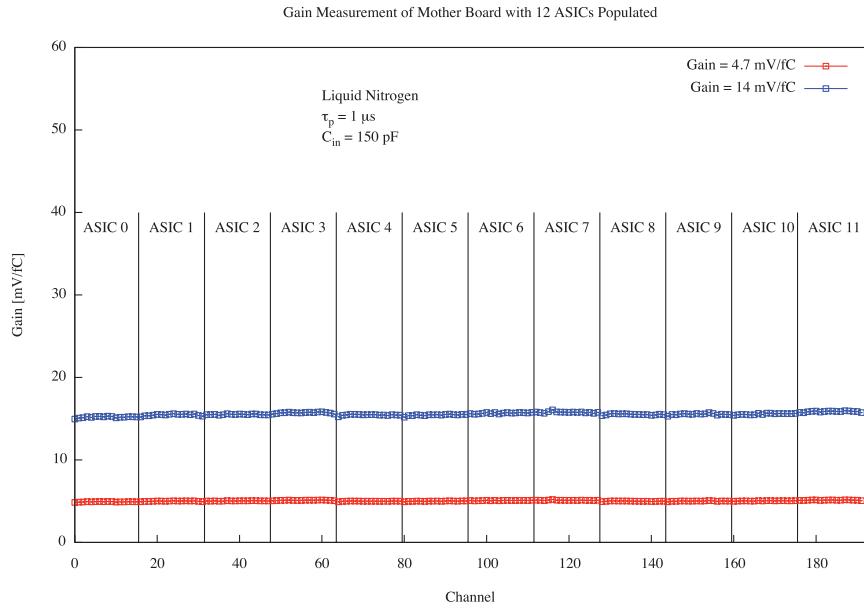


Figure 53. Plot of gain uniformity of 12 ASICs, total 192 channels, at 77 K with two different gain settings

6.2 Warm Electronic Amplification

Signals from the cold electronics are carried over the cold cables to dedicated feedthroughs mounted on the cryostat. The cold cables are connected to pin-carriers located on 356 mm outer-diameter CF signal feedthrough flanges that are mounted on nozzles N1A-N1K of the cryostat (see figure 5). The signal feedthrough design must accommodate 100% hermeticity and high signal density. A design based on the ATLAS pin carrier style was developed for this purpose. Two 8-row pin carriers and two 7-row pin carriers are welded onto the CF flange, as shown in figure 54, and create a vacuum-tight seal. Nine of the 11 signal feedthroughs receive signals from the three LArTPC anode planes (384 Y-plane, 192 U-plane, 192 V-plane), while the remaining two on the extreme ends of the cryostat only receive signals from one of the angled induction planes (672 U-plane on one feedthrough, 672 V-plane on the other).

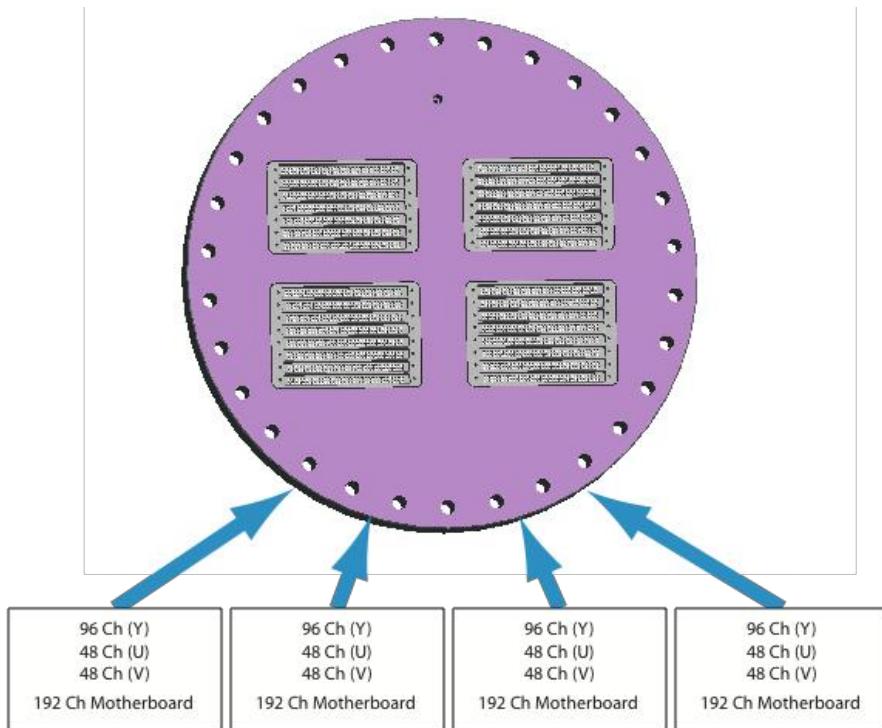


Figure 54. Signal feedthrough flange consisting of a 356 mm CF flange with two 8-row and two 7-row pin carriers welded in place.

A Faraday cage is mounted on the external, warm, side of the signal feedthroughs to provide shielding for the intermediate amplifiers located inside. The bias voltage feedthrough, which supplies anode plane bias voltages into the cryostat, is built onto a small 70 mm CF flange welded onto the signal feedthrough flange. A filter board mounted on the bias voltage flange filters noise and ensures a good ground connection. Figure 55 shows details of the signal feedthrough assembly with electronics boards, bias voltage feedthrough, and Faraday cage.

The intermediate amplifiers provide ~ 12 dB gain to the LArTPC signals to make them suitable for transmission over a 20 m long cable to the readout electronics (see section 6.3). Each intermediate

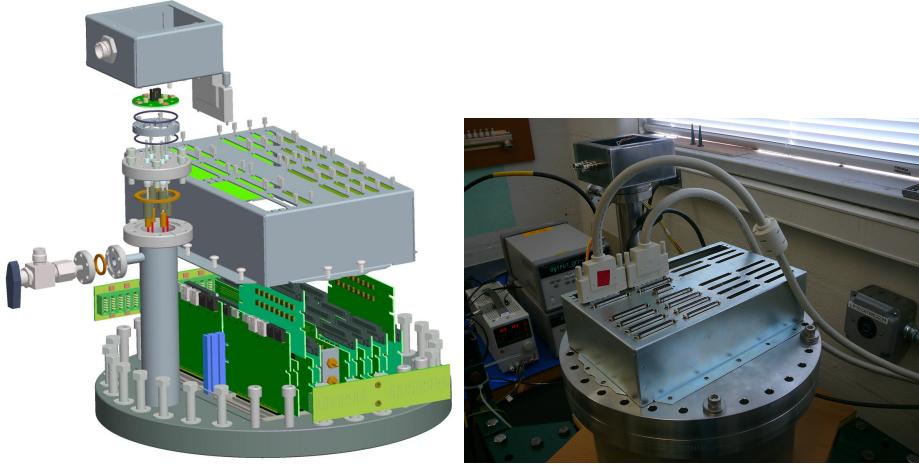


Figure 55. Left: Diagram of the signal feedthrough assembly, which includes intermediate amplifiers, Faraday cage, bias-voltage feedthrough and filtering circuit. Right: Photograph of one of the feedthrough assembly, partially constructed and being tested.

1 amplifier has 32 channels installed on the signal feedthrough flange and housed inside the Faraday
 2 cage to provide noise isolation. Figure 56 shows a picture of a prototype intermediate amplifier
 3 plugged on the signal feedthrough pin carrier. The intermediate amplifier uses a 68-pin SCSI-3
 4 connector to drive the 32 channels of signal differentially for better noise immunity. The layout and
 5 connector position have been carefully designed to ensure the card can be plugged on the pin carrier
 6 in either direction. This efficiently utilizes the limited available space on the top of the feedthrough,
 7 which also makes the design of the Faraday cage easier.

8 In addition to the intermediate amplifiers, there are two service boards mounted on the top of
 9 each signal feedthrough. The service board provides regulated low voltage, control and monitoring
 10 signals to the analog front end ASICs. It also provides pulse injection to the preamplifiers for
 11 precision calibration. The control, monitoring, and calibration signals are provided to the front end
 12 electronics with two-fold redundancy. Should one set of signals become defective the detector can
 13 still operate normally with the redundant set. Each service board plugs onto a 64-pin carrier row.
 14 Figure 56 shows a picture of a prototype service board.

15 **6.3 Readout Electronics**

16 The MicroBooNE readout electronics system consists of two subsystems: the LArTPC and PMT
 17 readout electronics. The LArTPC readout electronics are responsible for the readout, digitization,
 18 and processing of the induction and collection wire signals after amplification. The PMT readout
 19 electronics are responsible for the amplification, shaping, digitization, and handling of PMT signals,
 20 as can be used to provide a trigger signal for the readout and data acquisition systems. While the
 21 LArTPC and PMT readout systems share the same back-end design that organizes and packages
 22 the data for delivery to the DAQ system, they employ different analog front-end and digitization
 23 designs, which are described in this and the following subsection.

24 The LArTPC readout electronics are responsible for processing the signals from the 8,256



Figure 56. Left: Photograph of one of the intermediate amplifier boards plugged into a pin carrier during testing. Right: Photograph of a prototype service board.

wires in MicroBooNE after pre-amplification and shaping in the cold electronics, as described in section 6.1. The pre-amplified and shaped analog signals from the cold electronics are transmitted to the warm electronics outside the cryostat, as described in section 6.2, and then passed to custom-designed LArTPC readout modules distributed evenly over nine readout crates. The readout modules digitize the analog signals and then process and prepare them for shipping to designated DAQ machines (one per readout crate), described in section 6.6.

The LArTPC readout crates communicate with the DAQ machines via three duplex 3.125 GB/s optical links that connect to a crate controller module and data transmitter (XMIT) module on the crate end, and to three PCI Express boards on the DAQ machine end. The controller is responsible for configuration, trigger and run control command distribution as well as the slow monitoring of each readout crate. The controller occupies one of the optical links while the XMIT is responsible for sending two separate streams of readout data to the DAQ machines via the two other optical links. The first XMIT stream contains losslessly compressed LArTPC data associated with event triggers received by the LArTPC readout crates, such as the BNB trigger, and is referred to as the “NU” data stream. The second stream is a continuous LArTPC data stream which is compressed with some data loss. The continuous data stream is used for beam-unrelated physics analyses, such as the study of potential supernova neutrino events, and is referred to as the “SN” data stream. The compression schemes used in the NU and SN streams are described in section 6.3.3.

All readout crates are synchronized to a common 16 MHz clock. The clock sync is provided by a clock fanout board which shares the same ground as all readout crates, and is sent via coaxial cable to a distribution board which is mounted on each crate backplane. The readout frame size is set to 1.6 ms, which is equivalent to the time it takes for charge produced on the far end of the LArTPC to drift to the wire planes at the design cathode voltage of -128 kV.

6.3.1 Data Digitization

The amplified and shaped analog LArTPC signals are differentially received and digitized in the first section of the readout modules. Each ADC module holds 8 AD9222 octal-channel 12-bit ADCs. Each ADC module handles signals from 64 wires. The wire signals are grouped in two sets of 32 consecutive wire channels: either 32 induction wires plus 32 collection wires or two sets of 32

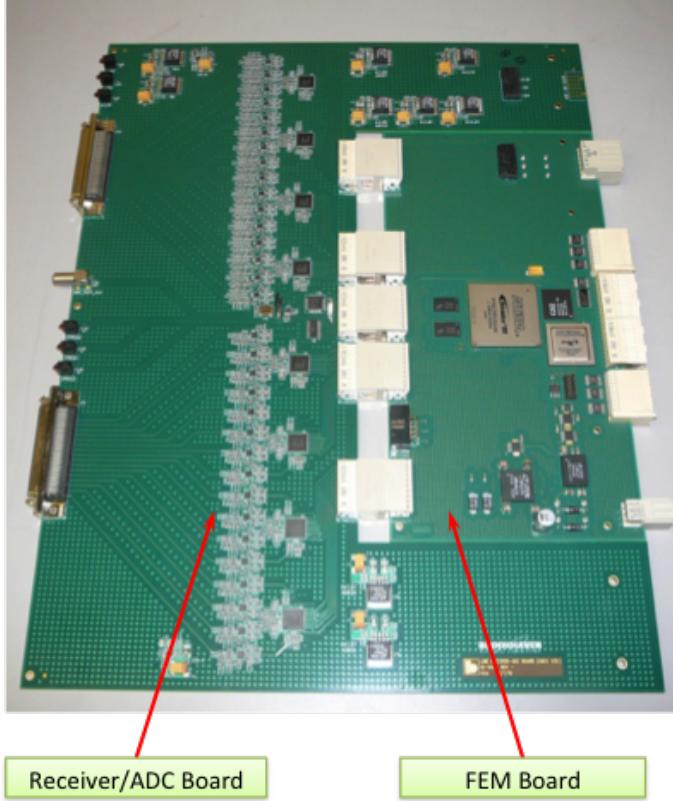


Figure 57. LArTPC ADC+FEM board.

1 induction wires. The induction channel sequence alternates wires between the two induction planes.
 2 The ADC module digitizes the signals continuously at 16 MHz. Each channel has a configurable
 3 baseline, which is either set low (450 ADC counts) for collection channels or at the middle of the
 4 dynamic range (2055 ADC counts) for induction channels, thus ensuring that both the collection
 5 plane unipolar differential signals and the induction plane bipolar differential signals can make use
 6 of the full ADC analog input range. The requirement to observe a MIP produced at the far end
 7 of the LArTPC in the induction plane determines the lower end of the dynamic range, while the
 8 requirement to observe a highly-ionizing stopping proton at the close end of the LArTPC without
 9 saturation sets the upper end. The digitized outputs from the ADC board are passed directly to a
 10 Front End Module (FEM) in the second section of the LArTPC readout module. The FEM houses
 11 an FPGA for data processing, data reduction, and preparation for readout by the DAQ system as
 12 described in the following section.

13 **6.3.2 Data Handling**

14 The FEM board consists of a 14-layer printed circuit board which is mechanically integrated with
 15 the ADC board as illustrated in figure 57. The choice of a smaller board allows for short trace
 16 lengths which is beneficial for high speed signals. The full assembly comes together as a standard
 17 VME 9U card in height, with a 280 mm depth. Differential outputs from the ADCs connect to the
 18 FPGA through HM-Zd connectors that have individual ground shielding on each differential pair.

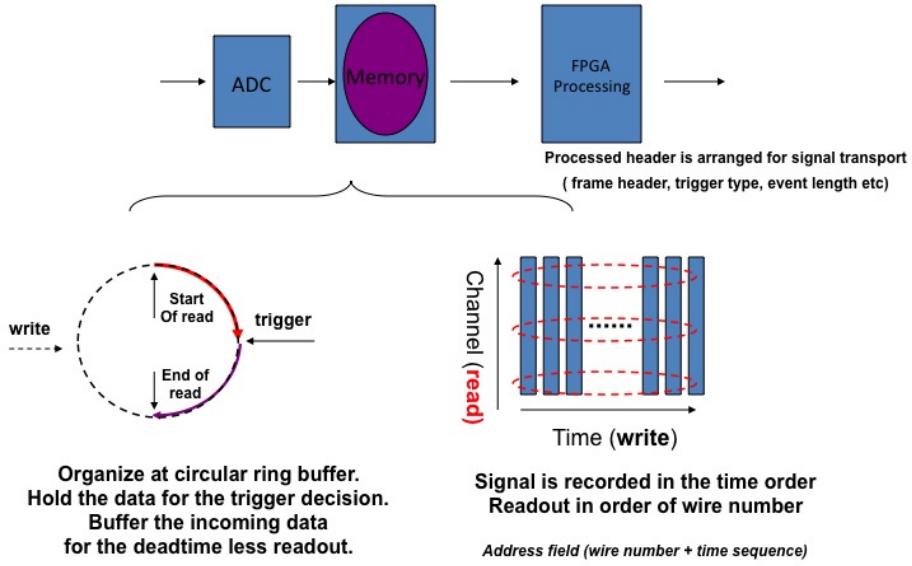


Figure 58. LArTPC readout sequence in the TPC FEM.

1 The digitized data stream moves from the ADCs to a Stratix III Altera FPGA, which reduces the
 2 sampling rate of the ADC from 16 MHz to 2 MHz. The 2 MHz sampling rate is optimized by taking
 3 into consideration the expected pulse shape provided by the convolution of the cold electronics,
 4 the expected LArTPC field responses, and the $O(1\mu s)$ diffusion effects which govern charge drift
 5 within the liquid argon. The FPGA stores the data from all 64 wires per board sequentially in time
 6 in a $1M \times 36$ bit 128 MHz SRAM, grouping two ADC words together in each 36 bit memory
 7 word. This requires a data storage rate of $(64/2) \times 2$ MHz = 64 MHz. The SRAM chip size and
 8 memory access speed allow for continuous readout of the LArTPC data. Since data reduction and
 9 compaction algorithms rely on the sequential time information of a given wire, the data readout out
 10 from this SRAM takes place in wire order in alternate clock cycles, again at the rate of 64 MHz.
 11 This read in/out sequence is illustrated in figure 58.

12 Separate DRAM multi-event buffers on the FEM store the NU and SN data streams. The data
 13 divert into the NU readout stream when a trigger is issued and received (as shown in figure 59) which
 14 signals, for example, an accelerator neutrino-induced event. When an event trigger is received, 4.8
 15 ms worth of data, relevant to that event, are packeted per channel and sent to the DAQ through the
 16 NU data stream. The 4.8 ms readout size is governed by the maximum drift time and spans three
 17 or four frames. In order to reduce the amount of data being transmitted, the FPGA trims the three
 18 or four frames to span the exact 4.8 ms required, 1.6 ms before the trigger plus 3.2 ms after the
 19 trigger. In parallel, the data is continually sent out through the SN data stream, frame by frame.
 20 The compression and data reduction algorithms applied to each of the two streams are described in
 21 the following section.

22 After processing by the FPGA, the data passes to the crate backplane dataway on connectors

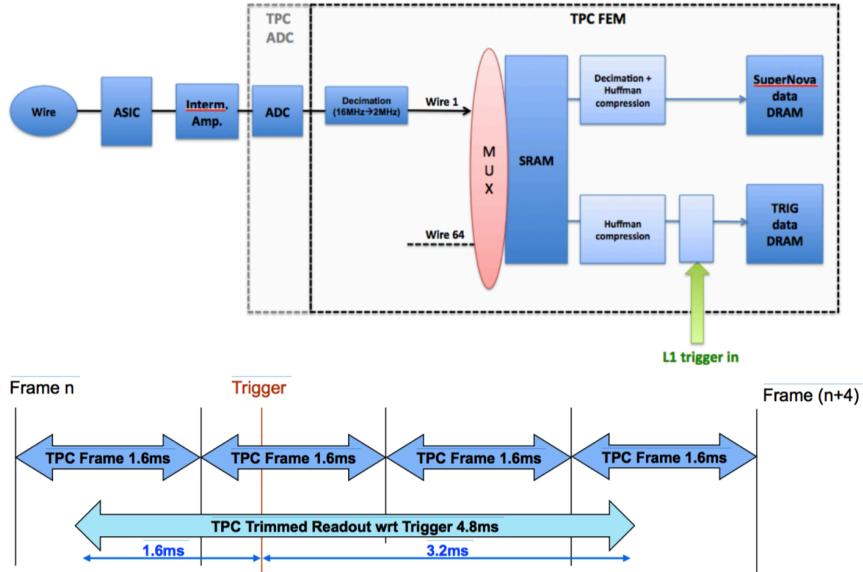


Figure 59. LArTPC trigger readout.

shown in figure 57. A token-passing scheme is utilized to transfer data from each FEM board to the data transmitter module (XMIT) in a controlled way, whereby each FEM, in the order of closest to furthest from the XMIT module, receives a token, transmits its data to the XMIT, and passes the token on to the next FEM in the sequence. For the NU stream, each FEM sends all data associated with a particular trigger number; while for the SN stream, each FEM sends all data associated with a particular frame number. This data transfer is relayed via the otherwise passive crate backplane, and is limited to 512 MB/s. In the XMIT module, the data is buffered temporarily and sent to the DAQ machine through the two streams, SN and NU, which proceed effectively in parallel.

6.3.3 Compression Schemes

In the case of the NU data stream, a lossless Huffman coding scheme implemented in the FEM FPGA compresses the data by approximately a factor of five. Further reduction in the overall rate is achieved by exploiting a PMT trigger in coincidence with the BNB trigger, as described in section 6.5. Huffman coding provides for lossless data compression by taking advantage of the slow variation of the waveform TPC data in any given channel. In particular, this compression scheme relies on the fact that successive data samples on any given wire vary relatively slowly in time. As such, when noise levels are low, any two adjacent data samples either coincide or differ by 1 ADC count. The most frequent values for the difference in ADCs between successive data samples are assigned pre-specified bit patterns with the lowest number of bits possible. Those bit patterns are encoded in the 16-bit data words that would otherwise be used for a single 12-bit ADC sample value. As such, data reduction of up to a factor of 14 is theoretically possible (the uppermost two bits are always reserved for header information, in each 16-bit word). In practice, the data reduction is sensitive to noise levels and LArTPC activity, and is also dependent on the gain setting. The compression factor achieved by MicroBooNE is shown in figure 60.

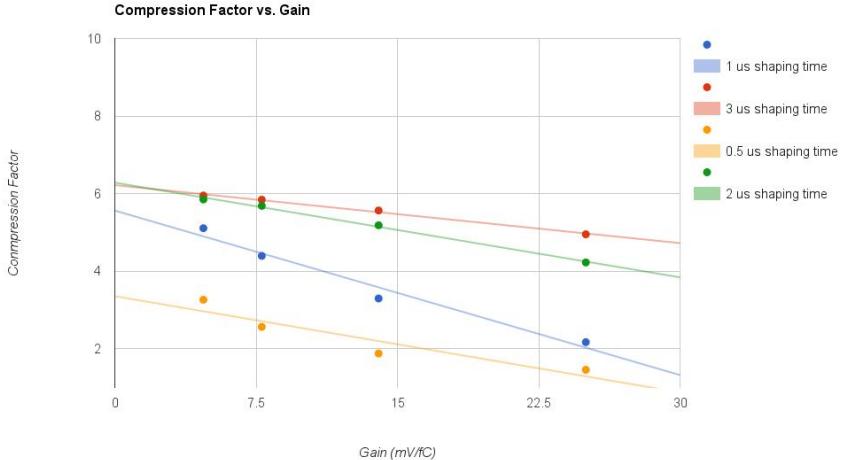


Figure 60. Compression factors achieved on ADC data with Huffman compression.

Because of the low trigger rate (the BNB trigger dictates an upper bound on the trigger rate of 15 Hz), lossless Huffman coding compression proves sufficient for the NU data stream. However, for the continuous SN stream, further compression becomes necessary, resulting in unavoidable data loss. A method called “dynamic decimation” (DD) handles this case. The DD scheme relies on recognizing regions of interest (ROI) in the data stream that contain waveforms corresponding to drift ionization charges. Portions of the data stream not containing ROI contribute to pedestal determination, and ROI are identified as deviations from the continually-updated pedestal, buffered, and read out to disk. At the time of this writing, the MicroBooNE SN stream compression scheme is being finalized and the SN readout stream will be commissioned in September 2016.

6.4 PMT Readout Electronics

The PMT readout electronics are responsible for processing signals from the 32 PMTs described in section 5 and identifying light signatures coincident with the BNB and NuMI beam spills. The coincidences generate PMT triggers that can be later mixed with other triggers in the Trigger Board (TB). Signals from the 4 light paddle guides installed in the LArTPC are also recorded by the PMT readout electronics, but these signals do not participate in the PMT trigger generation.

The stages of signal processing are illustrated in figure 61. First, each PMT signal (with the exception of the light paddle guide signals) is split into two different gains, as described in section 5.2.5, with the HG channel carrying 18% of the PMT signal and the LG channel carrying 1.8% of the PMT signal. Each gain is split once again into HG1 and HG2, and LG1 and LG2, in order to allow different processing of beam-related and beam-unrelated PMT signals. All $32 \times 2 \times 2$ plus 4 signals are pre-amplified and shaped in 16-channel pre-amp/shaper boards (section 6.4.1). Four PMT readout modules receive the analog shaped signals differentially and digitize them (section 6.4.2) at 64 MHz. The PMT readout modules then process the signals in order to prepare them for shipping to a designated DAQ machine and to form a possible PMT trigger (section 6.4.3).

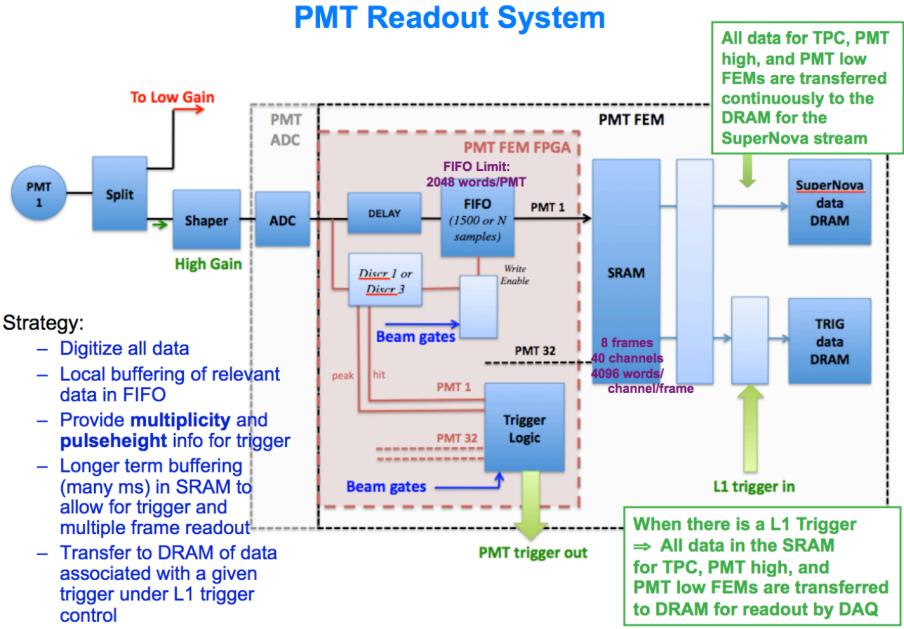


Figure 61. Signal processing in PMT FEM.

1 Each one of the 4 PMT ADC+FEM readout boards used in the PMT readout system handles
 2 one of the following:

- 3 • Readout of and PMT trigger generation using the HG1 PMT signals associated with neutrino
 beam events. The paddle signals are also readout by this board.
- 5 • Readout of and PMT trigger generation using the HG2 PMT signals that are out of beam time
 (i.e. cosmic rays and other cosmogenic backgrounds)
- 7 • Readout of the LG1 PMT signals associated with neutrino beam events
- 8 • Readout of the LG2 PMT signals that are out of beam time

9 After signal processing, the data is sent to a designated DAQ machine via a transmitter (XMIT)
 10 module in the same way as is done for LArTPC data. Two data streams are provided: a NU data
 11 stream associated with event triggers and a SN data stream which a continuous version of the NU
 12 stream readout.

13 6.4.1 Signal Amplification and Shaping

14 The preamp/shaper boards read raw PMT signals from the PMT HV/signal splitters and shape them
 15 into unipolar signals with a 60 ns rise time. The shaped signals are sent to the LArTPC readout
 16 boards differentially via short front-panel cables, in order to minimize noise, where they are digitized
 17 at 64 MHz. The 60 ns peaking time allows digitization of two or three samples on the rising edge.
 18 This in turn enables an accurate determination of the event t_0 needed to determine the x coordinates
 19 of ionization signals along the drift direction. An accurate time measurement also helps reject other
 20 tracks, such as cosmic rays, that cross the detector during the drift time.

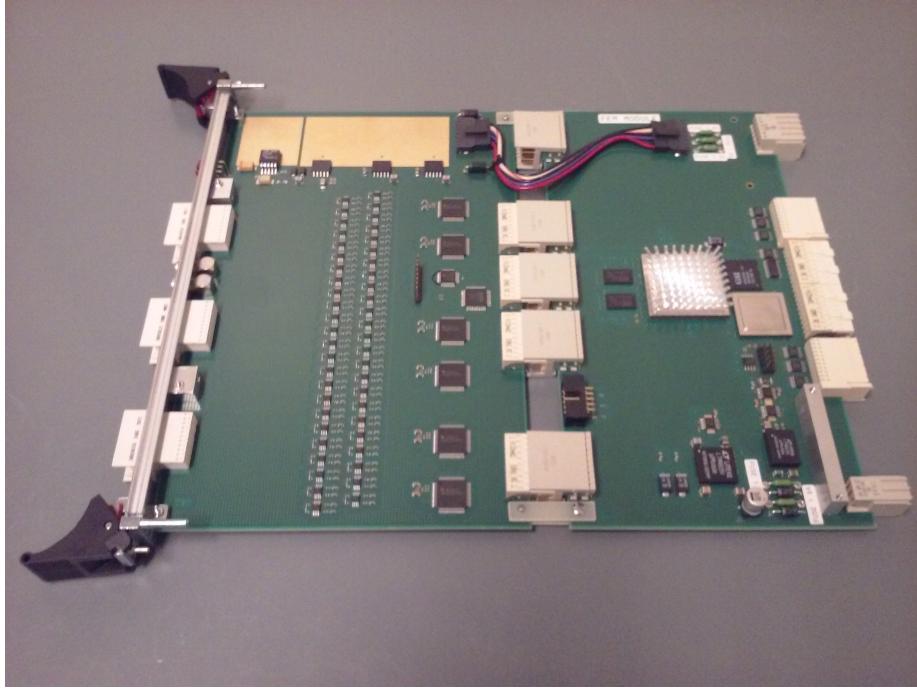


Figure 62. The PMT readout board digitizes 48 input signals.

6.4.2 PMT Data Digitization

The ADC (Texas Instruments, ADS5272) module part of the PMT readout board (figure 62) is responsible for digitization of up to 48 differentially-driven input signals. The differential signals are digitized at 64 MHz. The 64 MHz clock used by the PMT readout is generated starting from the 16 MHz clock that is common to all readout crates (LArTPC and PMT).

In addition to HG and LG waveforms from the PMTs, each readout board also receives, digitizes, and processes beam gate signal markers which arrive 4 μ s before the BNB 1.6 μ s and NuMI 10 μ s beam gates. These gates are used to (a) specially mark regions of interest where PMT data are read out continuously with no compression and (b) look for coincident PMT light signatures for trigger generation.

6.4.3 Data Handling and PMT Trigger Generation

PMT information is recorded in the NU data stream for four 1.6 ms frames associated with an event trigger: the frame containing the (asynchronous) trigger, the frame preceding the trigger frame, and two frames following the trigger frame. To avoid the inordinate amount of data that would be generated at a 64 MHz sampling rate, the FEM applies a zero-suppression immediately after digitization, retaining only samples above a given threshold as well as enough information before and after this useful data to establish a local baseline value; this collection of information is referred to as a PMT readout ROI. An exception is formed for beam-related or other likewise-triggered data where, for example, the 4 μ s-early BNB and NuMI beam gates mentioned in section 6.4.2 instruct readout of 1500 consecutive samples (23.4 μ s) surrounding and including the beam gates regardless of signal activity.

1 Two different discriminators are used: one that is active inside the beam gate(s), and one that
2 is also active outside the beam-gate-surrounding $23.4 \mu\text{s}$. The latter discriminator governs the
3 readout activity due to cosmic rays and other non-beam related activity. The first discriminator
4 enables PMT channels with pulse heights above a configurable threshold (e.g. corresponding to 1
5 photoelectron) to participate in trigger multiplicity and pulse height sum conditions, as described
6 in the following section. The thresholds for those two discriminators are set to different levels and
7 configured with different dead times for the HG and LG signals.

8 **6.5 Level-1 Trigger Generation**

9 The TB, which physically resides in the PMT readout crate, issues a “Level-1” trigger in order to
10 flag frames that must be treated differently. In the case of the LArTPC readout, the TB flags the 4
11 frames that must be trimmed and readout through the NU data stream and, in the case of the PMT
12 readout, it flags the 4 frames that must be readout in full through the NU data stream.

13 The inputs to the TB include a BNB trigger input (maximum rate of 15 Hz), a NuMI trigger input
14 (1.25 Hz), a Fake Beam trigger input (configurable), a PMT trigger input, and two calibration trigger
15 inputs, provided by the laser calibration system and the cosmic ray muon telescope, respectively.
16 The TB also has the ability to receive, via the crate controller, DAQ-issued calibration triggers,
17 which are used explicitly for cold electronics and PMT calibration. The various input triggers can
18 be independently pre-scaled, masked, and mixed together (OR or AND) to generate an event trigger.

19 The FPGA firmware in the PMT FEM can generate two different types of PMT triggers based
20 on the PMT signals: a cosmic PMT trigger and a beam gate PMT trigger. Beam gate PMT triggers
21 are configured in the same way for the BNB, NuMI, and Fake Beam. The nominal criteria for these
22 triggers are (1) PMT multiplicity ≥ 1 and (2) summed PMT pulse-height ≥ 2 PE summed over all
23 32 HG1 PMT channels. Both criteria must be met during any 100 ns time interval coincident with
24 the beam spill duration (1.6 μs in the case of the BNB and Fake Beam gates and 10 μs in the case
25 of the NuMI gate), and only channels enabled by the beam gate discriminator can participate in
26 the active pulse-height and multiplicity sums. The criteria for a cosmic PMT trigger are (1) PMT
27 multiplicity ≥ 1 and (2) summed PMT pulse-height ≥ 40 PE summed over any one of 28 preset
28 groups of 5 HG2 PMT channels that are grouped based on their spatial correlation. Again, only
29 channels enabled by the cosmic discriminators can participate in the trigger generation.

30 In addition, a software-based algorithm has been written to mimic the capabilities of the beam
31 gate PMT trigger performed in the FPGA and provide more flexibility in trigger criteria settings.
32 Details of this higher-level software trigger are described in section 6.7.

33 Figure 63 diagrams the PMT readout and trigger logic. Activation or masking of each of the
34 trigger inputs and outputs is DAQ-controlled. The trigger condition and explicit PMT trigger type,
35 if applicable, is available for every event in the NU data stream at both the event-building stage and
36 offline; this information is read out via a dedicated optical data stream, directly from the TB. The
37 trigger number and trigger time are also propagated and available in the NU data streams arriving
38 independently at the assembler DAQ machine from each LArTPC and PMT crate, and can therefore
39 be used to correctly associate data from the same event.

40 The readout control sequence is illustrated in figure 64. When a trigger is generated by the TB it
41 is passed to a fan-out module on a single cable and from there it is distributed to all crate controllers
42 (LArTPC and PMT). Through the crate backplane, the trigger gets propagated to each FEM. An

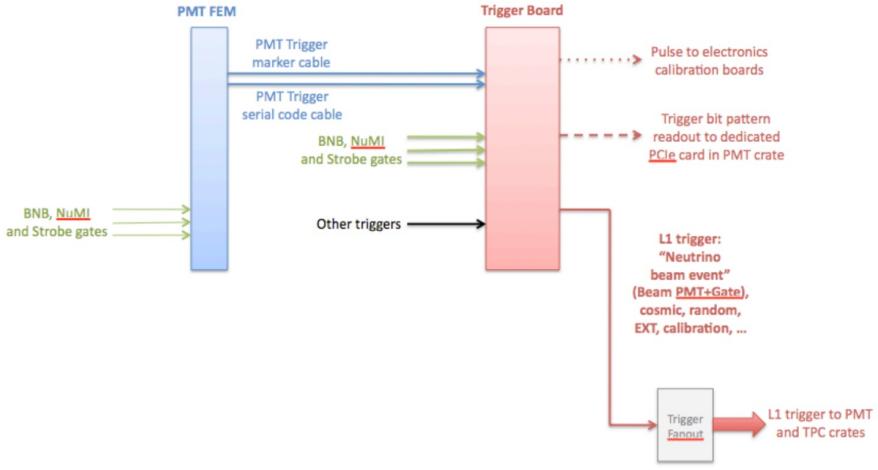


Figure 63. PMT readout and trigger logic.

- 1 FEM that receives a trigger temporarily inhibits the SN stream with its associated decimation and
- 2 initiates the loss-less readout scheme to direct the data to the appropriate readout path. SN readout
- 3 resumes once the XMIT is done sending all NU data associated with an event to the DAQ.

4 6.6 DAQ Design

- 5 The MicroBooNE DAQ system acquires data from the readout electronics, writes data to local disk
- 6 before transferring it to long-term storage, configures and controls the readout electronics during
- 7 data-taking periods, and monitors the data flow and detector conditions. These tasks are performed
- 8 on a network of commodity servers running both custom and open-source software.

9 The data from each crate of the backend electronics is sent to a dedicated server (called the
10 sub-event buffer, or SEB) via an optical fiber, arriving in a card on the SEB's PCIe bus. A real-time
11 application places these data in an internal buffer, collects all segments belonging to an event, and
12 creates a sub-event fragment that may be routed to a specified destination. For the NU stream,
13 in which the data arrives with every trigger, these fragments are sent to a single event-building
14 machine (EVB) over an internal network. Full events are checked for consistency and written to
15 local disk on the EVB before being sent offline for further processing. A high-level software trigger,
16 described in section 6.7, is applied to the data to determine whether events should be written locally
17 or ignored. For the SN stream the data remains on the SEB where it is written to disk and only sent
18 for offline analysis on explicit requests.

19 Data writing to either triggered or SN streams is limited by the RAID6 disk write speeds which
20 are roughly 300 MB/s. This is much less than the network bandwidth bottleneck, which is 10
21 GB/s. The 300 MB/s disk write speed therefore sets the maximum aggregate rate at which all SEB
22 fragments can ship data to the EVB without loss of data. With Huffman compression, which gives
23 a data reduction of approximately a factor of five (figure 60) and the PMT trigger, which reduces
24 the data rates by another factor of > 70, this is more than sufficient for MicroBooNE's maximum
25 15 Hz beam spill rate. MicroBooNE expects a total triggered write rate of around 12 MB/s. The

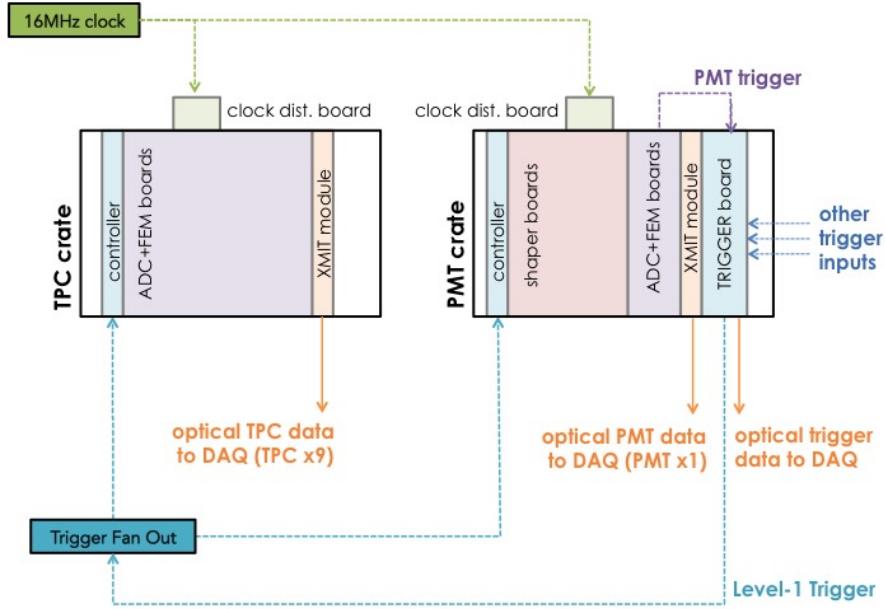


Figure 64. Readout control sequence. The MicroBooNE Level-1 trigger is a hardware trigger which consists of the OR between a BNB, NuMI, and EXT (strobe) trigger. Once received, the Level-1 trigger is propagated to all readout crates and instructs PMT and TPC data readout into ten dedicated sub-event buffer DAQ machines. The data across different DAQ machines is correlated at event building stage by the trigger number and corresponding trigger frame and sample numbers recorded in each data stream, per event. All readout crates are synchronized and correlated to the same 16 MHz clock.

1 SN stream circular buffers, which will be aggressively (non-losslessly) compressed beyond what
 2 the triggered stream experiences, will fill each server's 14 TB in on the order of one day, which is
 3 ample time to respond to a Super Nova Early Warning System (SNEWS) alert [81].

4 After data is written to disk, it is then copied to another server on the internal DAQ network,
 5 where the raw data is further compressed, shipped, and queued to be stored on in tape and disk
 6 cache using the Fermilab central data management system known as SAM. Offline applications then
 7 begin processing the raw data, converting the binary data format into a LArSoft [82] ROOT-based
 8 format which can be used as input for reconstruction algorithms. LArSoft is a common framework
 9 of software tools utilized by many LArTPC experiments at Fermilab and elsewhere. A separate
 10 process collects beam data and, during binary to LArSoft conversion, inserts that data into the built
 11 events. A duplicate copy of the data is also stored offsite at Pacific Northwest National Laboratory
 12 (PNNL). This collection of approximately 15 “projects” and the database which holds and monitors
 13 the state of the data flow is known as the Python/Postgres for MicroBooNE Scripting system
 14 (PUBS), and is patterned after a similar database state machine that the Double Chooz experiment
 15 used for data management. As PUBS pushes the data through this process, the progress of each
 16 project is monitored and viewable via GUI. PUBS can also monitor the state of the SN stream
 17 data, held locally on the SEBs. A separate offline PUBS instance controls the processing of the
 18 data, including applying newly calculated calibration constants as part of data quality management.

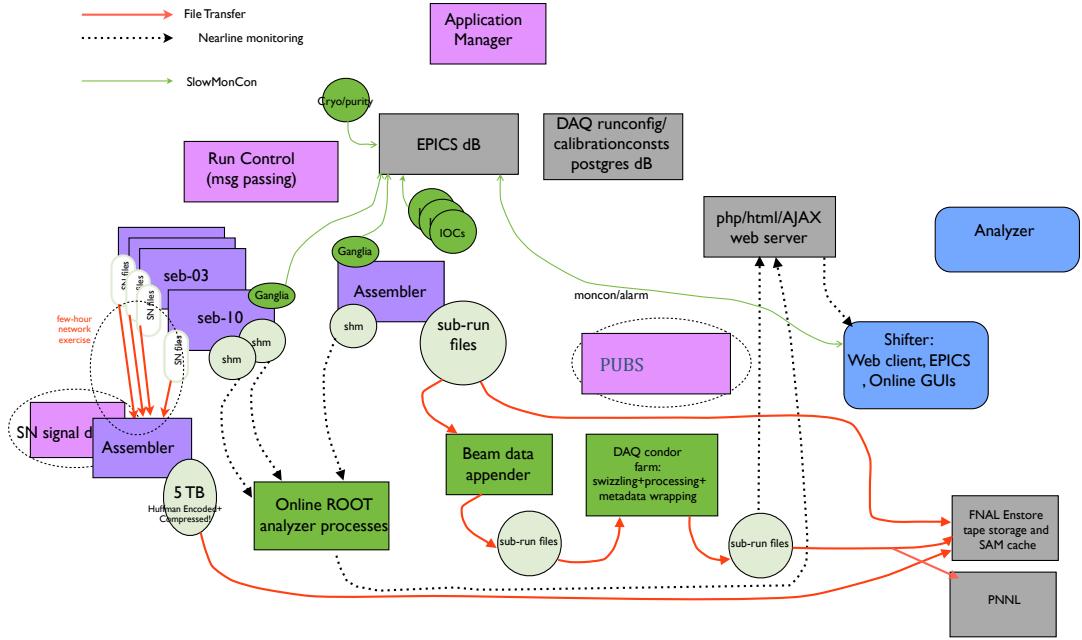


Figure 65. Data flow for MicroBooNE DAQ from raw to processed.

1 Figure 65 schematically depicts the flow of data throughout the MicroBooNE DAQ system.

2 Additional software components handle the management of the main DAQ processes moving
 3 the data. A run control application issues configuration and state-progressing commands to the SEBs
 4 and EVB. Configuration states are stored in a dedicated run configuration database, which allows
 5 for the setting and preserving of configuration information for the DAQ, readout, and additional
 6 components. This database not only allows for creating the large (~200 parameters) intricate
 7 DAQ run configuration files, but also enforces certain conditions which must hold for consistency.
 8 For example, the configuration that initiates the ASICS charge-injection calibration also dials and
 9 captures the settings on the external pulser that drives the calibration signal and assures that ASICS
 10 gains and peaking time parameters are enforced and recorded.

11 Another important aspect of the system is monitoring the health of the DAQ. Monitoring
 12 of DAQ components is accomplished through Ganglia which monitors basic system states (such
 13 as CPU, memory, and network usage) as well as allows use of custom metrics to monitor the
 14 data flow and status of the readout electronics [83]. These metrics are sampled and collected
 15 by the Experimental Physics and Industrial Control System (EPICS) slow monitoring and control
 16 processes, which archives desired quantities and provides alarms when pre-defined thresholds are
 17 exceeded [84]. Some examples of Ganglia metrics that are monitored and alarmed in EPICS are the
 18 rates of growth of the SEB data buffers, the fragment rates leaving each SEB, and fragment arrival
 19 rates at the EVB. Figure 66 shows examples of Ganglia metrics.

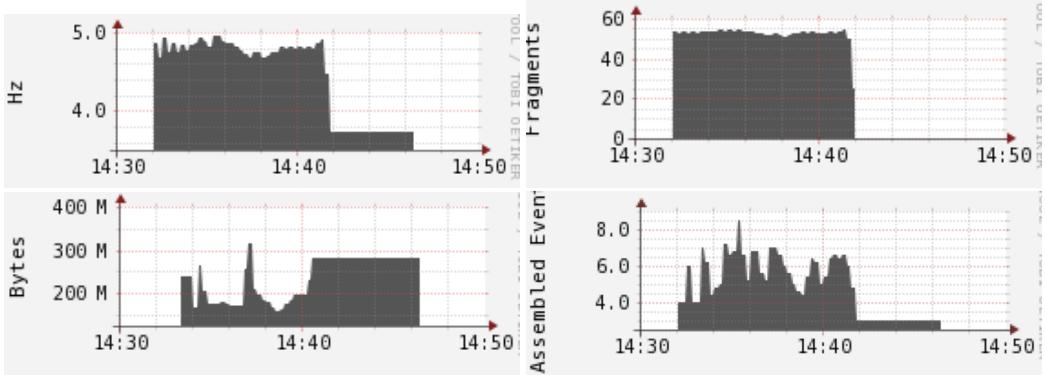


Figure 66. Ganglia metrics showing data flow on the EVB machine during a 5 Hz test run. Shown are the external trigger rate (top-left), the received fragment rate (top-right), the disk data write rate (bottom-left), and the number of assembled queued-up events waiting to be written (bottom-right). (11 fragments constitute a complete event in the MicroBooNE DAQ.)

Additional online monitoring exists to check data quality in more detail, through both programmed checks and visual checks including a real-time event display. The online monitoring takes snapshots into shared memory segments on the SEBs and the EVB, and thus provides the desired low latency checks of newly-arriving data. It continually walks through these ≈ 150 MB snap-shotted events and outputs histograms of occupancies and rates which are saved in ROOT files [85]. The histograms are then displayed in a web-based monitoring system that is easily accessible by the shift crew. Channels are aggregated in a variety of formats, including the order in which they appear in crates or across the wires and PMTs themselves. In this way, potential problems across connectors or crates, for example, may be more readily identified. Noisy, quiet, and unresponsive channels are easily marked and displayed to the shift crew. Figure 67 shows an example of available online monitoring information.

6.7 High-Level Software Trigger

After events are collected on the DAQ event-builder server, a suite of trigger algorithms are applied to the data. Currently, these algorithms mimic the PMT readout electronics' FPGA-based beam gate trigger algorithms, described in section 6.5. A search over the PMT digitized waveforms from the beam-spill period is performed, and if there is a significant amount of light in coincidence with the expected arrival time of the neutrinos, the event is marked and saved. This selection is performed on data that passes the Level-1 trigger, which typically includes data from the BNB and NuMI beams, and randomly selected off-beam data from an “external” trigger. A fraction of the data from each of these Level-1 trigger input streams is also retained via a random prescale, which provides a selection of data that has not been biased by the trigger. The high-level trigger algorithms take approximately 10 ms to return a result, a latency that is well-below the event-taking rate and so does not impact data-taking performance. The average pass rate for data-events in the PMT beam gate trigger algorithm is roughly 5%.

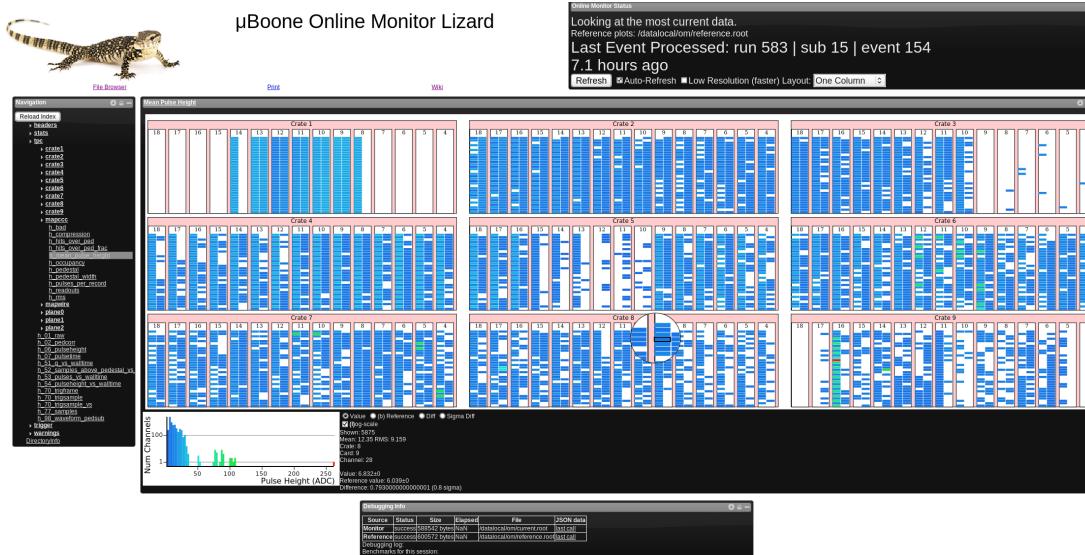


Figure 67. The online monitoring GUI (Lizard). The pedestal-subtracted ADC values for all 8256 wires in one minimum bias event are shown. Each box is a channel. Not all crates are fully populated with electronics.

7 Infrastructure and Monitoring Systems

MicroBooNE is housed at LArTF, which is located on-axis to the BNB. Complete knowledge of the electrical and cryogenic systems housed within LArTF is necessary to maintain acceptable operating conditions for the experiment. Continuous monitoring of the beam being delivered to LArTF is also necessary for subsequent physics analyses. This section describes the details of infrastructure within LArTF, as well as monitoring of the experiment and beam conditions.

7.1 Electronics Infrastructure at LArTF

This section describes the electronics infrastructure of the experiment, which is essential for proper operation of the electronics that control the functioning of and extract data from the LArTPC, PMTs, and other detector systems. Figure 68 shows a diagram of this system at LArTF, depicting racks located on a platform directly above the cryostat that house electronics for: LArTPC control and readout, light collection system, drift high voltage, purity monitor, calibration laser, trigger, and cryogenic control systems. Additional server racks containing the data acquisition, beam timing, and external network electronics are located in a separate computer room above and adjacent to the above-cryostat platform. The distribution of power, data, and network connections to and from all of these racks are also presented in figure 68, and is described in more detail in the following sections, as are electronics safety systems and interlocks.

7.1.1 AC Power Distribution and Grounding for Low-Noise LArTF Data-taking

As with any large detector operating with a high dynamic range, prevention of electromagnetic interference and its attendant effects on MicroBooNE data is an essential aspect of detector design. MicroBooNE's strategy for producing a low-noise environment for the LArTPC and associated readout electronics can be largely summarized in a few key points. AC power distribution-related

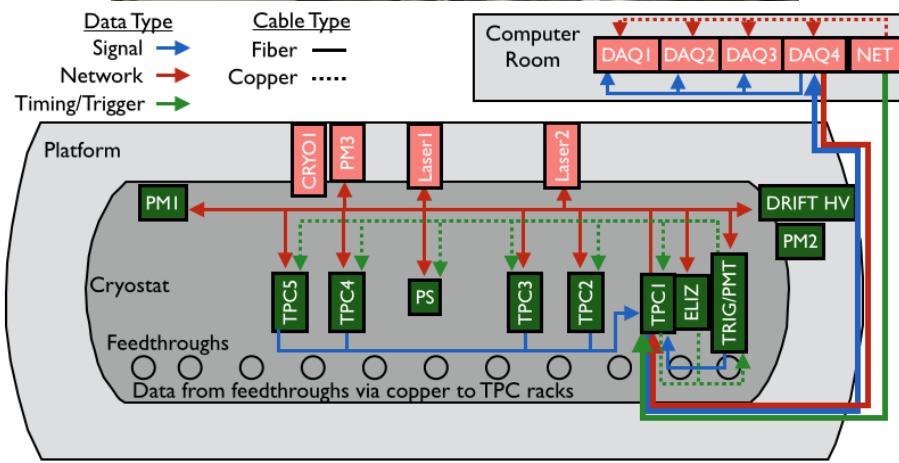
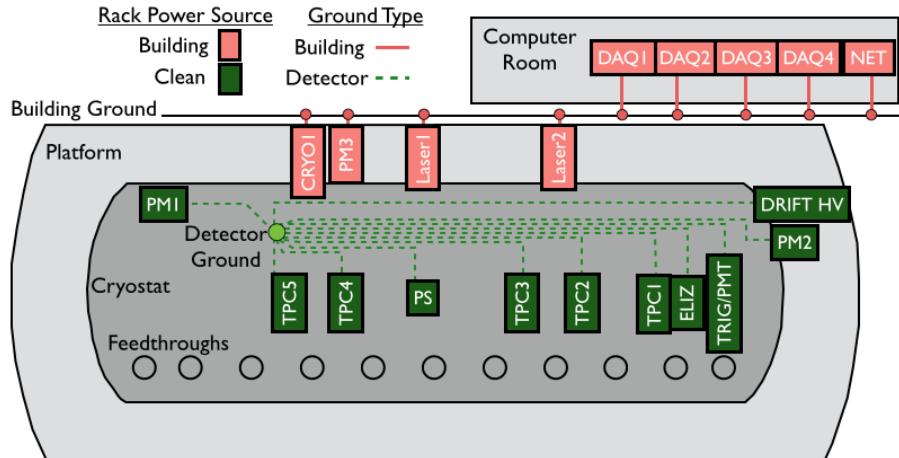


Figure 68. Top: Diagram illustrating location of deployed electronics racks and separation of detector and building grounds and clean and building power for differing racks. Middle: Photograph of installed electronics racks on the LArTF platform. Bottom: A diagram illustrating the general scheme of signal, network, and timing signal cabling in the LArTF computer room and platform.

1 items will be described here, while cabling, connections, and shielding will be described in a
2 following section.

- 3 • “Clean power,” or AC power electrically isolated from AC power for the rest of LArTF
4 (“building power”), is supplied to all sensitive electronics via an isolation transformer.
- 5 • Highly sensitive electronics are housed inside the Faraday cage provided by the detector
6 cryostat or inside Faraday cages directly grounded to the cryostat.
- 7 • The detector cryostat is grounded to a “detector ground,” which is physically and electrically
8 isolated from the ground provided to all other LArTF power circuits, or “building ground.”
- 9 • No direct electrical connections are present between detector ground and building ground.
10 This is accomplished through the use of insulating platform and cryostat saddle materials,
11 insulating cable trays and cables, and by inserting insulating “breaks” (i.e. fiber data links
12 or insulating cryo pipe sections) when connections between sensitive and potentially noisy
13 detector components are necessary.
- 14 • Indirect pickup on clean signals through capacitive coupling to adjacent noise sources is
15 minimized through use of detector-grounded shielding and electrically-insulating cable trays.
- 16 • Ground loops on detector ground are avoided wherever possible by connecting all electronics
17 racks directly to the cryostat and by minimizing direct electrical connections between racks.
- 18 • Direct or capacitive couplings between building and detector ground are constantly monitored
19 during installation and operation with a custom-designed impedance monitor.

20 A line drawing describing the production of clean power and clean ground is shown in figure 69.
21 A pair of 200 A clean power circuits produced at isolation transformers are used to power all sensitive
22 racks, which are indicated in figure 68. All racks containing LArTPC readout electronics are placed
23 on one circuit, while all other sensitive equipment is placed on the alternate circuit. On the platform,
24 all racks utilize clean power with the exception of the calibration laser and in-line purity monitor
25 racks, which either contain noise-producing elements or support building-grounded components.
26 All racks in the LArTF computer room utilize building power.

27 208-3 phase power is distributed to each individual electronics rack. For racks with significant
28 power requirements or a large number of components, this power is delivered to a Fermilab-
29 designed “AC switch box,” which distributes power to an Eaton Power Distribution Unit (PDU)
30 only upon receiving an interlock signal from a smoke detection system in each rack, which will
31 be described in more detail below. Rack components then receive power from one of the three
32 phases on this PDU. For racks with fewer requirements, power is supplied to components directly
33 from an interlocked simplified AC switch box or SurgeX SX-1120-RT PDU. Racks with sensitive
34 electronics are grounded to the cryostat via copper sheeting running throughout insulated cable
35 trays above the cryostat. Sensitive components within each rack are connected to a tin-plated
36 copper grounding bar electrically connected to the rack bottom and running the height of the rack.
37 Mechanical attachments to the rack provide grounding for less sensitive rack components. As
38 mentioned before, any unintentional direct connection between building and detector ground can

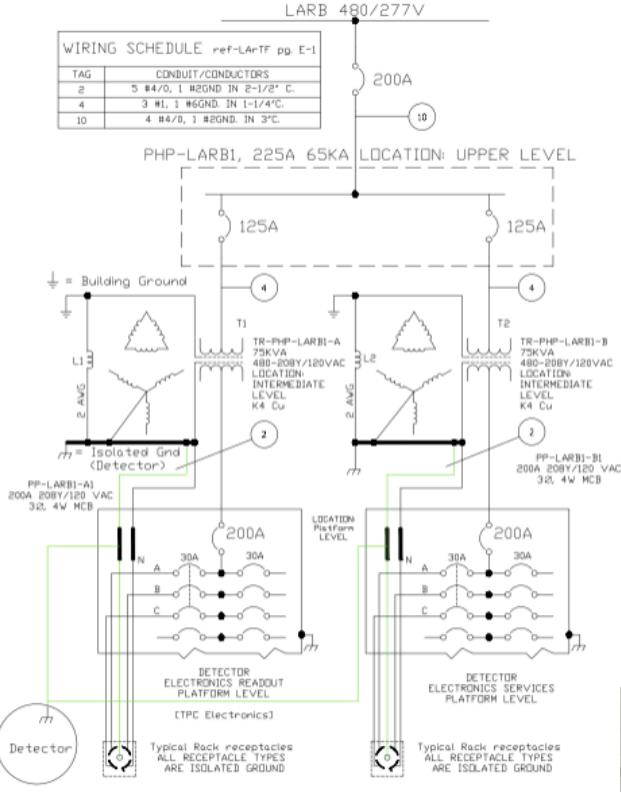


Figure 69. Line drawing of clean power generation and distribution and connections to detector ground.

be quickly identified by the impedance monitor located on the LArTF platform. Figure 70 shows pictures of this equipment.

7.1.2 DC Power Distribution to the MicroBooNE Detector

DC power is provided to the LArTPC and readout electronics by power supplies in clean-powered, detector-grounded racks for a variety of purposes:

- Holding the LArTPC cathode plane at voltage to produce the desired ionization electron drift speed
- Holding the two ungrounded anode planes at the proper constant voltage to ensure the planes are transparent to drifting electrons
- Operating the light collection system PMTs
- Powering the cold electronics located inside the LArTPC
- Powering the warm electronics located in the LArTPC readout electronics racks
- Powering auxiliary systems, such as LArTPC temperature sensors and purity monitors

Table 7 summarizes the required voltages or currents for each of these purposes as well as the power supply make and model utilized in each case. Power supplies are located in the relevant subsystem's electronics rack, with power and grounding connections as dictated by that rack.

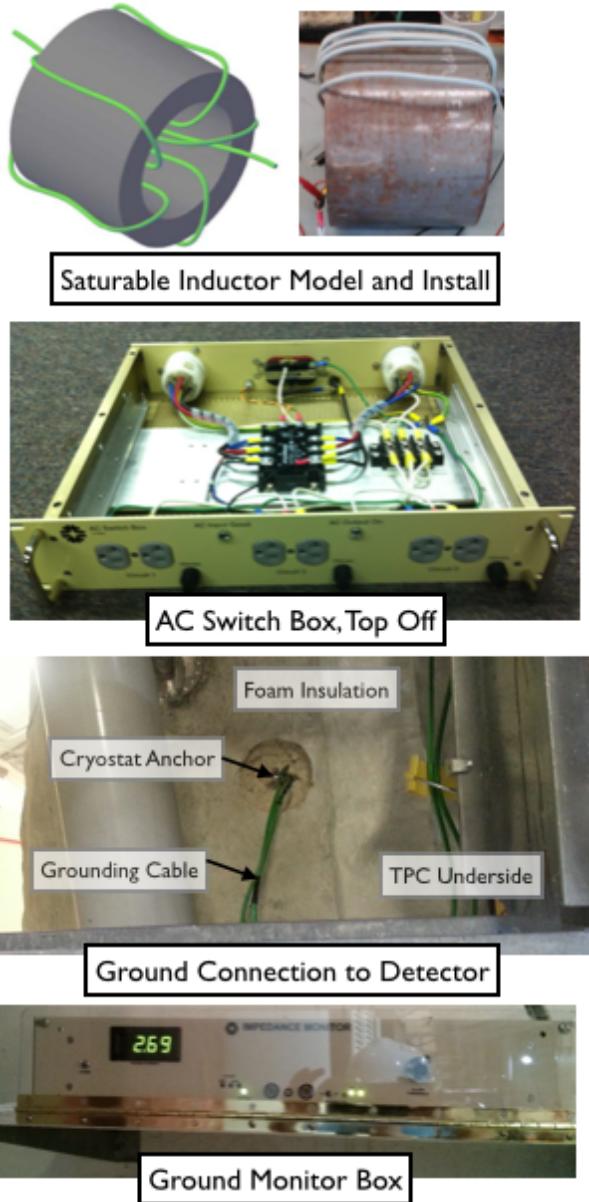


Figure 70. Pictures of the installed saturable inductor (top), AC switch box (top middle), detector ground strap and connection (bottom middle), and impedance monitor (bottom).

DC supply power consumption is minimal in most cases, with the exception of the warm LArTPC electronics PL-508 supplies. Care was taken to distribute the AC power load by limiting the number of high-draw PC supplies per electronics rack.

7.1.3 Network, Timing, and Data Distribution for Low-Noise LArTF Data-taking

Network, timing, and data connections must be made between the detector, building-ground, and detector-ground racks to properly read out MicroBooNE data. However, as described above, these connections must be made while maintaining strict detector-building electrical isolation. Deployed

Table 7. Overview of MicroBooNE DC power distribution. Delivered voltages or currents are listed, along with power supply makes and models and whether each supply utilizes clean or building power.

System	Supplied Value	Supply Make/Model	Clean Power?
Drift High Voltage	120 kV	Glassman LX150N12	Y
Anode Plane Voltage	1 kV / 8 mA	Wiener MPOD	Y
PMT High Voltage	2 kV / 3 mA	BIRA T4	Y
Cold Electronics Power	8 V / 10 A	MPOD MPV 800I	Y
Warm LArTPC/PMT Electronics	+/-5 V MDH 2-8 V/25 A +12 V MEH 8-15 V/92 A +3.3 V MEH 2-7 V/115 A	Wiener PL508	Y
Auxiliary Systems Power	Various	Various	Laser: N Inline PM: N Cryostat PM: Y

¹ interconnections meeting both of these requirements are displayed in figure 68.

² Timing and LArTF-external network signals are brought into LArTF via electronics in the computer room, where all racks are building grounded and powered. These signals are distributed and processed in the computer room via copper cable, while network and processed timing signals to be sent to the platform are converted onto fiber cables and aggregated into a central fiber termination box. A fiber trunk line then delivers these signals to the platform, where another fiber termination box on a detector-ground rack is used to fan out these signals. Network connections are fanned out via fiber to a network switch in each platform rack, while timing signals are re-converted to copper and further processed for use by the trigger system on a different detector-grounded rack.
¹⁰ All rack-to-rack cables are run in insulating cable trays beneath the platform.

¹¹ PMT and LArTPC data are transferred from each detector feedthrough to readout crates in detector-ground racks via insulated copper cable whose shield is tied to detector ground. Digitized crate output is then sent to the aforementioned platform fiber termination box, where these signals are sent via fiber trunk line to the computer room. In the computer room, these fibers are then fanned out to the appropriate DAQ computer. Readout crate and cold electronics control commands are transmitted in the opposite direction utilizing a similar scheme, with crate controls delivered directly via fiber, and cold electronics commands delivered via fiber to a copper fanout in a detector-ground rack.

¹⁹ Clock and trigger signals must also be sent from a central trigger rack to all detector-ground LArTPC/PMT readout racks. These signals are transmitted via copper connections, and represent the only source of ground loops on detector ground. To further reduce the possible impact of induced noise in these and all copper cables mentioned above, all insulating cable trays beneath the platform are lined with copper sheeting grounded to the detector. As an additional precaution, all LArTPC signal copper cables are run in separate cable trays from power and auxiliary cabling beneath the platform as well as inside every rack.

²⁶ All cables between all detector components have been uniquely labelled with serial number, source, and destination to allow for ease of replacement and reconnection. Ample fiber and copper spares for every major cable type are also installed along with the production cables to allow for

1 quick replacement of any failed cable.

2 **7.1.4 Interlocks and Safety Systems**

3 All electronics racks contain smoke-sensing and temperature-monitoring systems, which, when
4 interlocked with AC and DC power transmission in each rack, constitute a rack protection system
5 (RPS) designed to meet Fermilab safety requirements and reduce the risk of fire and related damage
6 in LArTF and to individual rack components.

7 The RPS principally consists of a smoke sensor connected to a Fermilab-designed rack pro-
8 tection box. This box produces and outputs a 12 V interlock signal when the rack protection box is
9 on and receiving a “no-smoke” signal from the smoke sensor. This 12 V signal can be sent to the
10 AC distribution box located in each rack, as described above, to allow AC transmission to all rack
11 components only if the RPS is on and not detecting smoke. A similar 12 V “RPS Status” signal
12 is also produced by the RPS box for input into the MicroBooNE slow control box, which will be
13 described in following sections. Alternate contacts are available on the rear of the RPS box for
14 coupling the status of additional subsystems, such as the DAQ and calibration laser uninterruptible
15 power supply (UPS), to smoke sensor or rack power status.

16 Temperature sensors deployed in two or three locations in each electronics rack sample air
17 temperature within each rack. Temperatures at each sensor are read out and recorded in the slow-
18 control database by the slow-control monitoring box. In addition, the box also produces a 5 V
19 interlock signal if all sampled temperatures are within pre-programmed thresholds. In electronics
20 racks distributing PMT- or LArTPC-related DC power these temperature interlock signals are input
21 into each relevant power supply, allowing DC power distribution only when this interlock signal is
22 present, for safety purposes.

23 Additional hardware interlocks ensure the non-simultaneous operation of particular systems.
24 In particular, the PMT system is disabled when cryogenic system liquid level sensors detect a
25 level below that of the highest PMT bases, or when the UV laser system is active. The former
26 requirement is enforced with a dry-contact hardware interlock, while the latter is enforced with a
27 software interlock in the MicroBooNE online software.

28 **7.1.5 Performance Measurements**

29 The proper operation of each production electronics rack’s AC and DC distribution and RPS systems
30 has been tested prior to installation at LArTF. Furthermore, test stands exercising functionality of
31 DAQ, PMT and LArTPC electronics, trigger, and drift HV systems have successfully incorporated
32 and tested various aspects of these same AC and DC distribution and RPS systems. Impedances
33 between detector and building grounds were recorded throughout the installation of the rack infras-
34 tructure at LArTF using the impedance monitor located on the LArTF platform.

35 **7.2 Slow monitoring and control system**

36 MicroBooNE uses EPICS for controlling and monitoring most devices and conditions important
37 to the experiment. These include power supply controls, temperatures, fan speeds, rack protection
38 interlock status, and various environmental conditions. The DAQ, cryogenics systems, and beam
39 data collection systems operate independently of the EPICS slow monitoring, but export data

1 which are imported into EPICS for archiving and status displays. Applications from the Control
2 System Studio software collection [86] are used for providing displays, alarm notifications, and data
3 archiving.

4 An EPICS system consists of any number of server programs implementing the EPICS Channel
5 Access (CA) protocol [87] to provide client programs access to any number of process variables,
6 where each process variable represents a quantity being controlled (an output) or measured (an
7 input). The EPICS base distribution provides a standard type of channel access server called an
8 Input/Output Controller (IOC), which can be extended to support specific hardware as desired.

9 Most power supplies are controllable over the network through the NetSNMP protocol [88].
10 Several EPICS driver modules are available for SNMP, and MicroBooNE utilizes one written at
11 NSCL [89]. An IOC with this SNMP module runs on a central computer and contacts the power
12 supplies over a private network for monitoring and control. The photomultiplier power supplies are
13 reused from the DØ experiment and have custom IOCs running in their own controllers. The main
14 high voltage power supply has only a simple RS-232 serial interface; control and monitoring for it is
15 provided by a nearby computer running an IOC with the EPICS asynDriver [90] and StreamDevice
16 [91] modules.

17 MicroBooNE has a number of racks in various positions above the detector and in an adjacent
18 server room. Each is equipped with a rack-protection system and multiple digital temperature
19 sensors, and most contain one or two fan packs, each containing 6 fans. To monitor and control
20 these devices, each rack has an 1U rack-mount enclosure containing an ARM-based single-board
21 computer (SBC) running Linux. An off-the-shelf GESBC-9G20 from Glomation Inc. [92] is utilized
22 for the SBC. A custom interface board [93] connects the SBC to front panel LEDs, temperature
23 probes, fan packs, and rack-protection-status input. The temperature sensors are DS1621 chips,
24 controlled and read out over an I2C bus by the SBC's I2C controller. The DS1621 also has a
25 thermostat output with programmable trip and reset temperatures, which are connected via the
26 interface board to outputs that can be used to interlock devices in the racks, such as power supplies.
27 The fans provide pulse-per-rotation outputs, which are monitored by a 12-channel tachometer
28 implemented via a PIC16F887 microcontroller, and also read out by the I2C bus. An EPICS IOC
29 runs in each SBC, with custom device drivers for reading all status information and controlling the
30 heartbeat LED and temperature sensor trip and reset points.

31 Data are imported into EPICS channels from a number of external sources. The primary
32 reason for duplicating these data in EPICS is to integrate displays and warnings into one system
33 for the experiment operators, and to provide integrated archiving for sampled data in the archived
34 database. An IOC running on a central computer provides “soft” process-variables channels for these
35 data. The data acquisition system provides many metrics describing its operation via the Ganglia
36 system[83, 94], which makes the data available in an XML format easily read by a Python script,
37 which in turn writes to EPICS using the PyEPICS module [95]. The hardware and system status of
38 the DAQ computers is monitored through the industry standard Intelligent Platform Management
39 Interface (IPMI); rather than writing a script to import data from IPMI directly into EPICS, a IPMI-
40 to-Ganglia interface provided by the FreeIPMI’s “ipmi-sensors” package [96] is used, allowing
41 data to be imported via the same mechanism used for the DAQ metrics. Separate Python scripts
42 periodically retrieve data about outside weather conditions from various sources, cryogenics system
43 data from a file retrieved non-intrusively from the IFIX cryogenics control system, and beam data

¹ from Fermilab’s Intensity Frontier Beam Database (IFDB) [97].

² 7.3 Beam Monitoring

³ The primary source of neutrinos for the MicroBooNE experiment is the BNB. The primary beamline
⁴ is lined with instrumentation including toroids which indicate beam intensity, “multiwires” showing
⁵ beam profile in the horizontal and vertical planes, and beam position monitors measuring the mean
⁶ beam position. Data from these monitors are stored on a spill-by-spill basis in the IFDB. Many of
⁷ MicroBooNE’s physics analyses require that beam data are recorded for each spill and matched to
⁸ detector events.

⁹ Primary beam monitoring in MicroBooNE is done using a “dashboard” interface to IFDB.
¹⁰ By using the IFBD instead of the accelerator control system, the experiment can also verify that
¹¹ data are being acquired by the IFBD. The dashboard is accessible over the network using a web
¹² browser. The final monitoring step includes a post-data-merge check, ensuring that beam data are
¹³ successfully matched with detector data for all beam spills. This is done once the detector DAQ
¹⁴ binary data file is closed.

¹⁵ The dashboard presents a graphic representation of the data, allowing for easy error identification,
¹⁶ as shown in figure 71. The experiment monitors: two toroids which indicate beam intensity;
¹⁷ three multiwires, each of which shows beam profile in each plane; and beam position monitors
¹⁸ along the beamline which show the vertical and the horizontal position. Parameters pertaining to
¹⁹ the target and horn, such as cooling air temperature and horn current, can also be monitored. The
²⁰ dashboard allows the experiment to easily add additional devices if experience demonstrates the
²¹ need for their monitoring.

²² Data are monitored in near real-time. A reasonable history is also kept so that changes are
²³ easily identified. The accelerator control system provides detailed diagnostics tools to experts and
²⁴ can be used in case of any problems.

Booster Neutrino Beamline Status Display (e,1d events) - Historical data for 2015-07-01 11:15:03

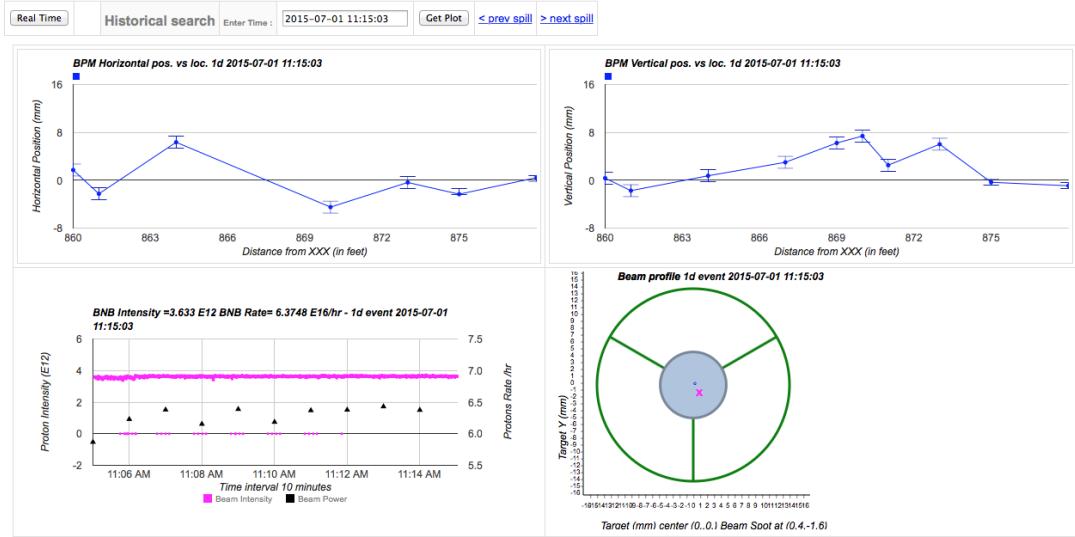


Figure 71. The BNB dashboard, showing graphical representation of beam instrumentation data, is used to monitor the beam. The top box shows the timestamp of the beam spill and indicates if data is stale by changing the color. The two top plots show the primary proton beam position along the BNB. The bottom left plot shows the recent beam spill intensity and rate. The bottom right plot shows the beam as projected onto the Beryllium target (the grey circle in the middle with radius of 0.5 cm). The dashboard is accessible via web page providing both real time updates and the review of past data. The page can be easily extended to monitor additional beam devices.

8 UV Laser System

The knowledge of the electric field inside the drift volume of a LArTPC is a necessary aspect for performing subsequent event reconstruction. Distortions of particle tracks due to field non-uniformities affect the accuracy of the particle momentum reconstruction based on multiple scattering. Deviations from a uniform drift field may arise mainly due to accumulation of positive argon ions in the drift volume. These ions are produced by ionizing particles from neutrino interactions, as well as by cosmic rays. While electrons produced by ionizing particles are quickly (within few milliseconds) swept towards the readout system, ions have significantly lower mobility. Their drift velocity in the MicroBooNE detector at nominal drift field is of the order of 0.8 cm/s. The rate of cosmic muons in the LArTPC volume is calculated to be 11,000 muons/s within the active volume (assuming no overburden, and a cosmic rate of 200 muons/m²/s through a horizontal plane at the earth's surface and 63 muons/m²/s through a vertical plane), traversing a combined length of 1.9×10⁴ m through the liquid argon [98, 99]. Assuming that cosmic muons are minimum-ionizing (2.1 MeV/cm) and produce 23.6 eV per ion pair, positive ion charge is produced at a rate of 2.8×10⁻⁸ C/s in the MicroBooNE LArTPC. These ions are continuously neutralized at the cathode. The resulting charge distribution in equilibrium is shown in figure 72. Such accumulated space charge leads to noticeable distortion of the drift field and, consequently, to deviations of reconstructed track coordinates by up to 10 cm (see figure 73). The ion drift velocity is comparable to local argon flow velocities,

1 produced by global argon recirculation flow and thermal convection. Therefore the distribution
 2 of positive space charge inside the drift volume may be not only nonuniform (figure 74), but also
 3 non-stationary.

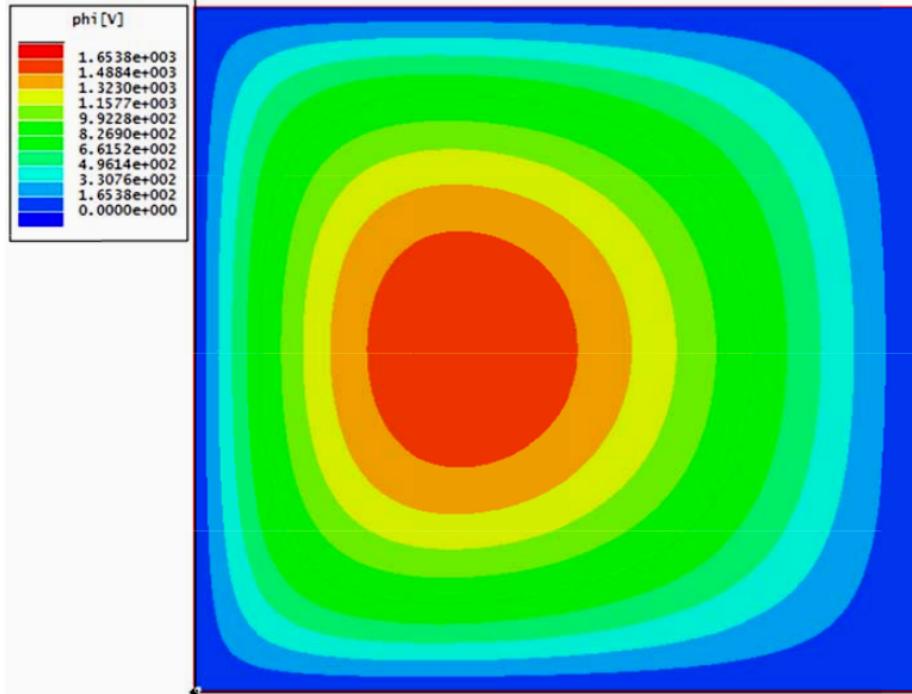


Figure 72. Distorting potential distribution due to positive space charge in equilibrium in the MicroBooNE detector.

4 A nonuniform drift field in the LArTPC leads to the apparent bending of truly straight tracks
 5 of high-momentum ionizing particles. In principle, a set of events from such particles allows for
 6 the reconstruction of the field in any small region of the LArTPC drift volume, using the systematic
 7 apparent curvature of tracks at different angles passing through that region. In practice, the rate
 8 of such events from cosmic muons is too low to acquire sufficient statistics in reasonable time.
 9 A method to generate straight ionization tracks at a defined location in liquid argon is described
 10 in [100]. A collimated photon beam from a pulsed UV laser with $\lambda=266$ nm can ionize liquid
 11 argon via multi-photon absorption. The resulting ionization track is straight, characterized by low
 12 electron density, therefore featuring little charge recombination loss, unlike cosmic muon tracks.
 13 Laser tracks are also free from δ -electrons, which complicate track reconstruction in the case of
 14 muons. The method was successfully exploited in the Argontube long drift LArTPC [101–103] to
 15 derive the non-uniformity of the electric field along its 5 m long drift volume [104].

16 The MicroBooNE LArTPC requires a set of such tracks in order to cover the whole sensitive
 17 volume to reconstruct field distortions. This is the purpose of the laser calibration system, which
 18 features two UV lasers, one on either end of the LArTPC. Tracks are generated by steering a pulsed
 19 laser beam, introduced to the cryostat via a custom-designed opto-mechanical feedthrough [105],
 20 through the LArTPC active volume.

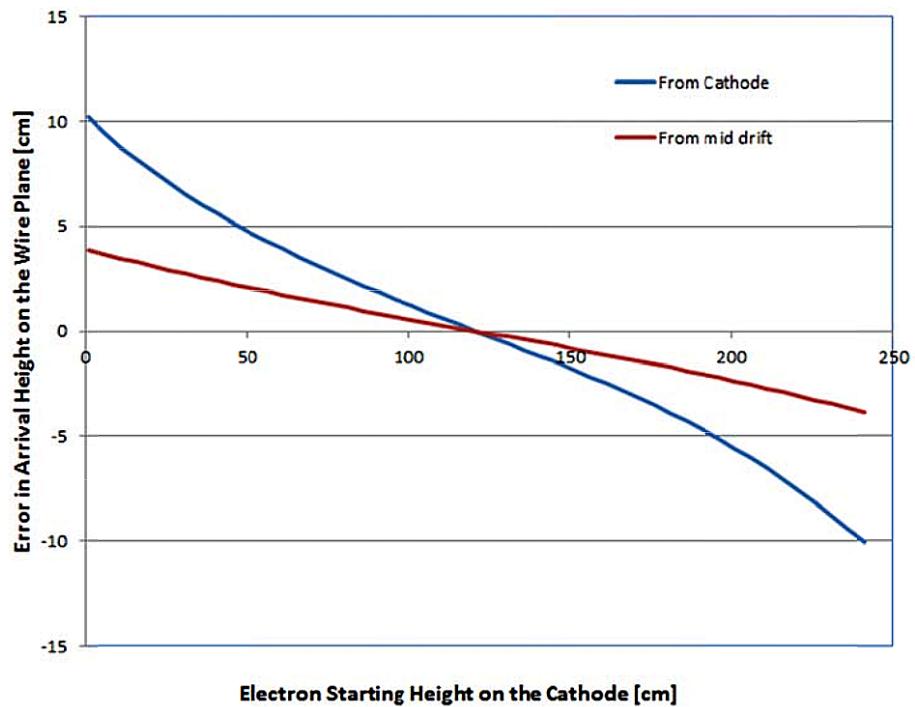


Figure 73. Deviation of a crossing track from its true coordinates due to positive ion space charge.

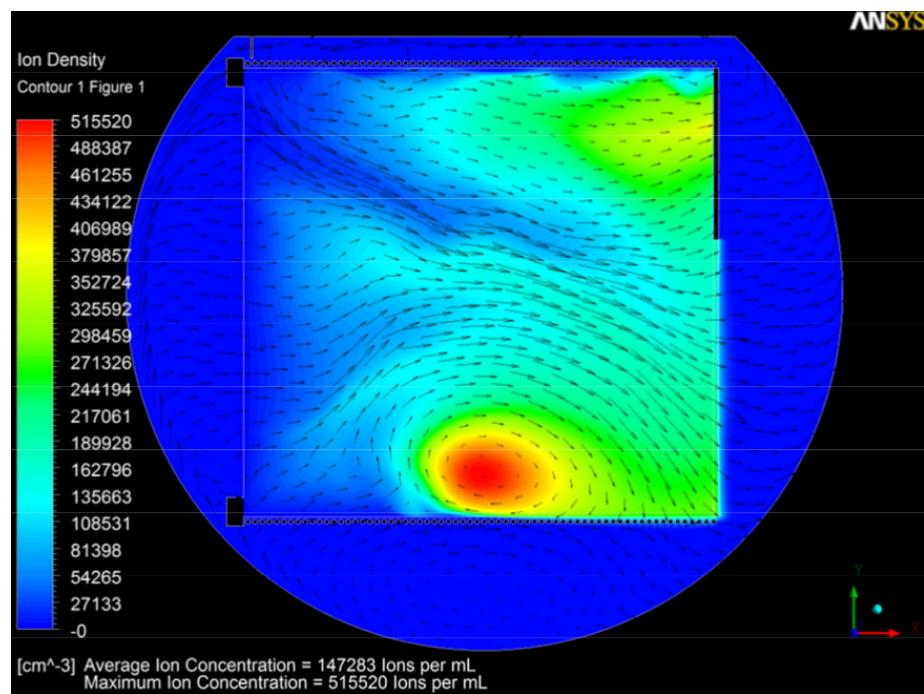


Figure 74. Distribution of the positive space charge in presence of argon circulation.

8.1 UV Laser Calibration

In order to unambiguously reconstruct a drift field vector at any point within the detector fiducial volume, a minimum of two ionization tracks are required to cross in the region of interest. The total number of crossing points is determined by the required reconstruction granularity. In MicroBooNE the initial scenario is to acquire 100 tracks from each direction, producing a reconstructed 3-D map with voxels that are approximately $10 \times 10 \times 50 \text{ cm}^3$ in volume. This map provides a rough picture of the space charge distribution. Depending on the results of this measurement, areas of interest can be studied in more detail. Repetitive study of small volumes may reveal dynamics in the space charge distribution due to turbulent circulation, and should further inform an optimized scenario for a standard UV laser calibration procedure.

An algorithm of drift field calibration utilizes an input array of detector events with one straight ionization track in each element of the array. The result of the calibration is a coordinate correction map, to be applied to each track, which converts apparently curved track images back to the true coordinate system where they are straight. The algorithm is iterative with an optimizable iteration step and required accuracy. An example of simulated reconstruction in 2-D space is depicted in figure 75, showing that the magnitude of the field distortions can potentially be reduced from 10 cm down to several millimeters in 99% of the detector volume.

8.2 Laser Source and Optics

A Nd:YAG laser emitting light at a wavelength of 1024 nm is used as the primary light source [106]. Inside the laser head, nonlinear crystals are installed in the beam line for frequency doubling and summing, resulting in a wavelength of 266 nm needed for ionization of liquid argon. For this wavelength, the Nd:YAG laser is specified to produce an output energy of 60 mJ for each 4 to 6 ns long pulse and a horizontal polarization. The maximal repetition rate is 10 Hz; the beam has a divergence of 0.5 mrad.

Figure 76 depicts the arrangement of the laser system. An optical table, mounted on the cryostat, was developed to accommodate the necessary parts in a stable and compact environment. The components were chosen to be accessible remotely where necessary. The emitted laser beam contains not only ultraviolet light but also all other harmonics generated in the crystal and the primary light of 1024 nm. Dichroic Mirrors optimized to reflect only wavelengths in the UV region are used to filter higher wavelengths out. To absorb the transmitted wavelength behind the mirrors, glass-ceramic plates are installed. The beam leaving the laser head is reflected by the first 45° -mirror into an attenuator. For optical adjustment and verification of the non-visible UV-beam, a green alignment laser is placed behind this mirror and adjusted such that its path is coincident with the UV-laser beam. In the attenuator (Altechna Wattpilot) a turnable $\lambda/2$ -plate enables rotation of the orientation of the laser beam polarization. Behind the attenuator two parallel plates are installed such that the angle of the incident beam matches the Brewster angle of the reflector. Modulating the polarization of the beam adjusts the intensity of the reflected beam. An aperture is placed in the optical path of the beam after the attenuator to control the beam diameter. The last part in the beam line is a remotely-controllable mirror mount (Zaber T-OMG), which directs the beam to the laser feedthrough on the cryostat. A photodiode (Thorlabs DET10A/M), which is sensitive in the ultraviolet region, detects the scattered light when a laser pulse is fired. Its signal is then used as a

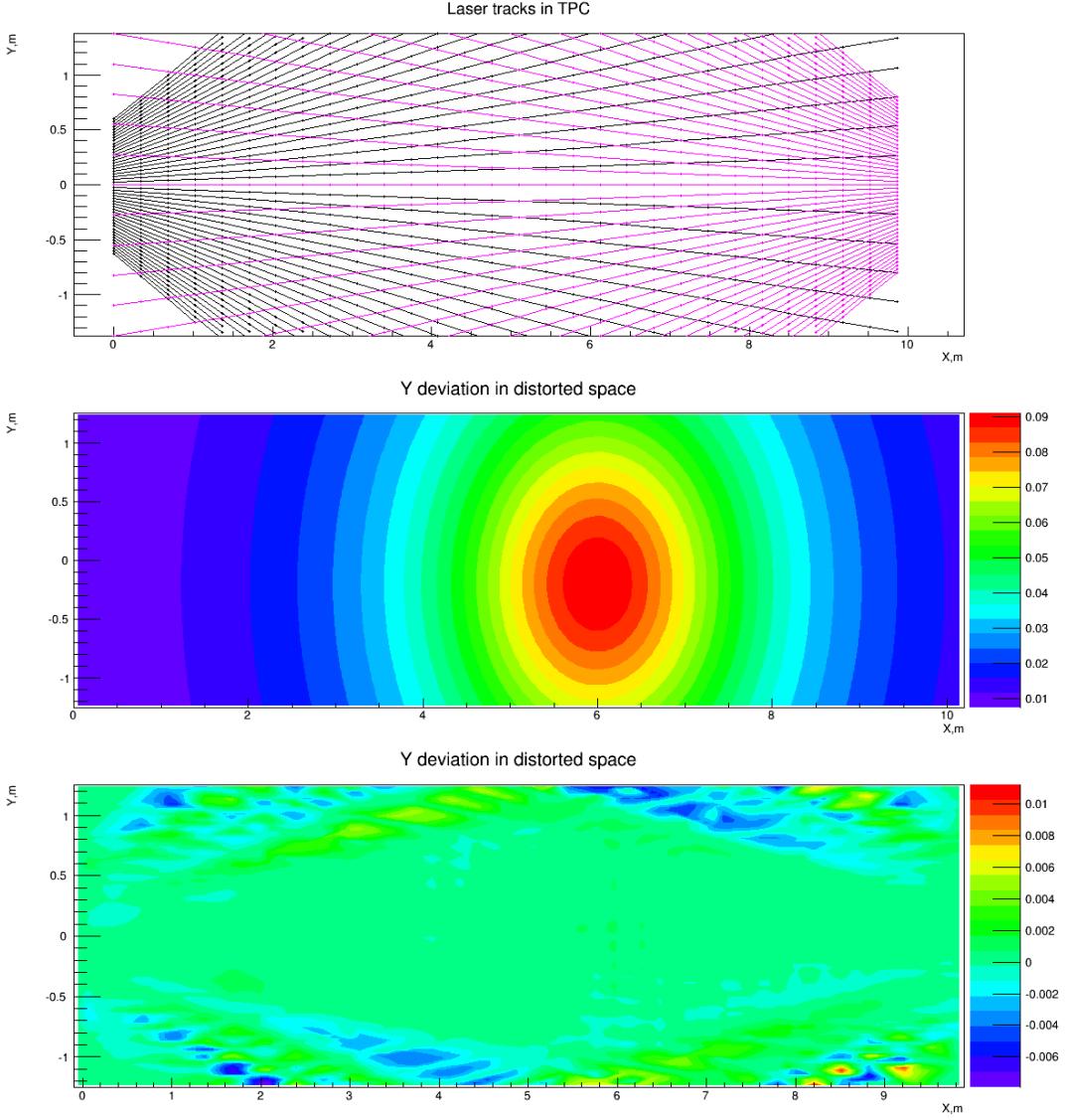


Figure 75. Top: Simulated true laser beam trajectories in the detector; Middle: map of Y coordinate of track deviation under influence of non-uniform electric field; Bottom: residual map of reconstructed track Y-coordinates versus true ones.

1 trigger for data taking. Both the UV-laser head and the optical table are mounted on a 15 mm thick
 2 aluminum plate.

3 8.3 Steering System

4 One of the main challenges of the laser calibration system is the introduction of a steerable laser beam
 5 into the detector. Earlier, an evacuated quartz-glass [107] was utilized to introduce a laser beam into
 6 liquid argon, however this beam had a fixed path through the detector. For the purpose of scanning
 7 the full detector a fully steerable mirror in liquid argon is necessary. In the MicroBooNE detector,
 8 this is achieved by mounting a mirror on a horizontally-rotatable support structure. A rack and

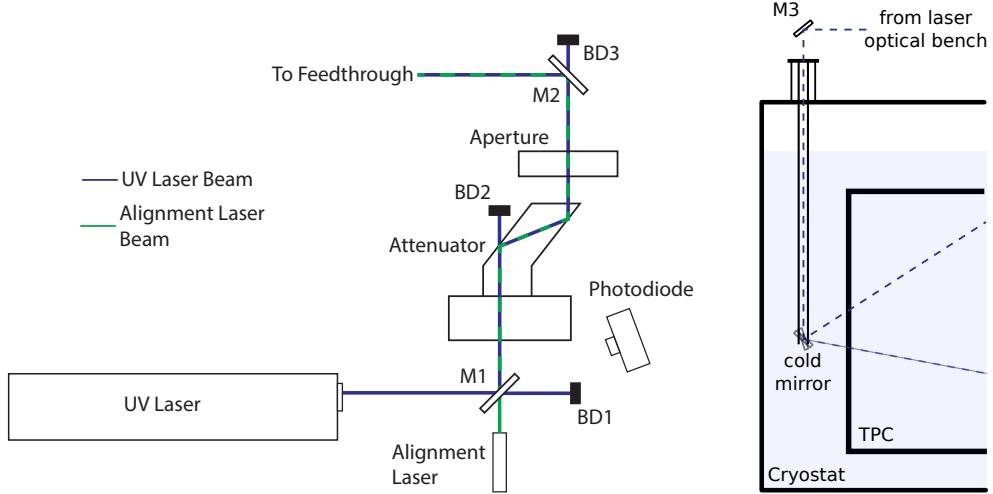


Figure 76. Left: A schematic drawing (not to scale) of the components used for laser beam configuration. An alignment laser (visible light) is introduced along the UV-laser path at the first dichroic mirror (M1), such that the paths overlap. In the attenuator the UV-laser beam intensity can be adjusted to the desired level, and the diameter of the beam is controlled by an aperture. A motorized mirror (M2) deflects the light into the direction of the feedthrough. Beam dumps (BD) are installed behind all mirrors to absorb the non-reflected laser light. Right: Side view of the cryostat indicating the mirror support structure with respect to the LArTPC.

1 pinion construction, where the mirror is mounted on the frontside of a half gear (pinion), provides the
 2 necessary freedom for the vertical movement (see figure 77 right). To steer the horizontal movement
 3 from outside the cryostat, a commercial differentially-pumped rotational feedthrough is deployed
 4 (see figure 77 left). The rack and pinion construction is attached to a linear feedthrough. Both
 5 feedthroughs are motorized to allow for remote control and automation of the mirror movement.
 6 The mirror support structure was fabricated out of polyamide-imide (Duratron T4301 PAI), which
 7 has a very low outgassing rate and low thermal expansion coefficient, and is certified for operation
 8 at 87 K. To minimize the probability of discharges due to the close location of the feedthrough to
 9 the field cage structure in MicroBooNE, no conductive parts were used in the support structure.
 10 The support structure has a total length of 2.5 m in MicroBooNE. Both feedthroughs are equipped
 11 with high precision position encoders from Heidenhain. The accuracy of the encoders is chosen
 12 such that a position accuracy of 2 mm for the laser beam spot over 10 m distance is achieved. An
 13 external interface box controls the encoders and records a position reading upon receiving a trigger
 14 signal from the photodiode. The DAQ computer accesses the position information over an ethernet
 15 connection. The same computer is also used for steering the two motors via a motor driver system
 16 (over a RS232 interface).

17 **8.4 Performance Tests and Initial Operation**

18 A full performance test of the laser calibration system identical to the one installed in the Micro-
 19 BooNE LArTPC was performed prior to the final installation. Apart from the general proof-of-
 20 principle of the laser calibration system, several operationally relevant parameters were identified.
 21 These include scanning speed, positioning accuracy, positioning limits, optimal laser beam in-

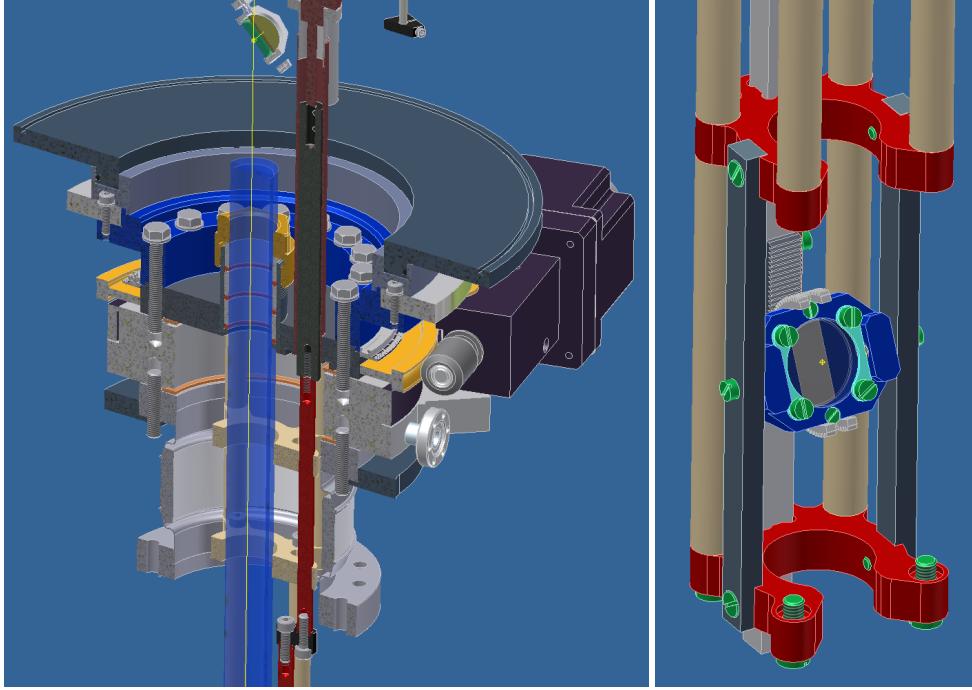


Figure 77. Left: CAD cutaway drawing of the feedthrough construction is shown. The yellow line indicates the laser path. Right, CAD drawing of the cold mirror including the support structure.

tensity, beam diameter, and the minimal achievable field distortion which can be resolved. The test system consists of a LArTPC equipped with 64 readout channels and an active area of about 400 cm^2 , with a drift distance of 40 cm (see [108] for further details).

Several tests of the motorized feedthrough were performed under warm conditions before cold tests were conducted. One crucial parameter for the quality of the electric field calibration is the resolution at which laser tracks can be aimed in the detector. For the rotational axis this angle is directly measured on a circular scale. For the vertical movement the linear displacement of the bellow is translated into a rotation inside the cryostat, as can be seen in the CAD drawing in figure 77. This construction introduces uncertainties to the measurement position and backlash. The backlash can be compensated by always approaching positions from the same direction. For the translation of the linear movement ΔL into a rotation $\Delta\phi$ the translation ratio s according to $\Delta L = s \cdot \Delta\phi$ was measured with a Bosch GPL3 laser alignment device. The obtained ratio is $s = 0.3499 \pm 0.0002 \text{ mm/}^\circ$. The dominant uncertainty in the vertical position measurement is the accuracy of the encoder $\sigma_{\text{linear}} = \pm 1 \mu\text{m}$, which translates into a vertical rotation measurement accuracy $\sigma_{\text{vertical}} = 0.050 \text{ mrad}$.

Horizontal movement limitations arise from the construction of the feedthrough system, namely the warm mirror support structure. This limitation originates in the manner the laser table is mounted relative to the feedthrough on the MicroBooNE cryostat. Vertically the mirror can be rotated more than 45° relative to the horizon in both directions. In an upward looking configuration, no limitations arise which would affect the coverage of the detector with the beam. When the mirror faces the opposite downward direction and the laser is properly aligned onto the center of the mirror, only the

laser diameter and the size of the mirror limit the achievable coverage. However slight misalignment will affect this, since the beam will not be in the optimal spot anymore. In warm tests a maximal downward angle of the beam of 52.5° with respect to the horizon was achieved. During the cold tests the horizontal and vertical movement speeds were set to $2.6^\circ/\text{s}$ and $1^\circ/\text{s}$, respectively, and horizontal and vertical angles of 81° and 22° , respectively, were covered. Warm tests of the fully expanded setup showed vibrations if the chosen movement speed was too large. The vibrations are expected to be damped with a more stable installation on the detector, and with the immersion of the setup in liquid argon.

Modulation of the beam energy with respect to the vertical alignment of the cold mirror was found to be crucial for obtaining sufficient ionization in the detector. Investigations showed that the reflectivity of the selected dielectric mirrors, which were optimized for 45° in air, are very sensitive to the angle of incidence in liquid argon. Therefore during a calibration run, the beam energy has to be controlled. The emitted UV-laser beam has a diameter of 6 mm and will spatially diverge during propagation. A beam with this diameter will produce an ionization signal larger than the wire spacing, which will limit the capabilities of the full system. With the aperture a small as possible diameter of the laser was selected to enter the detector. Measurements of the diameter were performed with thermal paper (used for thermal printing) on which the selected beam spot burns in. The minimal achieved diameter was 1 mm.

9 Conclusion

The MicroBooNE detector is the culmination of several years of development and construction. It is the largest LArTPC ever constructed in the U.S., and innovations such as custom cold electronics and filling without evacuation represent major technological advances that future experiments will build upon. Operations of MicroBooNE began in the summer of 2015, and a cosmic ray tagger system was added in the fall of 2016. Future publications will be dedicated to describing the performance of the detector systems introduced in this paper.

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1 References

- [1] T. Doke, *A historical view on the r&d for liquid rare gas detectors*, *Nuclear Instruments and Methods in Physics Research Section A* **327** (1993), no. 1 113 – 118.
- [2] W. Willis and V. Radeka, *Liquid-argon ionization chambers as total-absorption detectors*, *Nuclear Instruments and Methods* **120** (1974), no. 2 221 – 236.
- [3] C. Rubbia, *The liquid-argon time projection chamber: A new concept for neutrino detectors*, EP Internal Report 77-8, CERN, May, 1977.
- [4] P. Benetti et al., *A three-ton liquid argon time projection chamber*, *Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment* **332** (1993), no. 3 395 – 412.
- [5] P. Cennini et al., *Performance of a three-ton liquid argon time projection chamber*, *Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment* **345** (1994), no. 2 230 – 243.
- [6] F. Arneodo et al., *The ICARUS 50-l LAr TPC in the CERN ν beam*, [hep-ex/9812006](#).
- [7] S. Amerio et al., *Design, construction and tests of the ICARUS T600 detector*, *Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment* **527** (2004), no. 3 329 – 410.
- [8] **ArgoNeuT** Collaboration, C. Anderson et al., *The ArgoNeuT Detector in the NuMI Low-Energy beam line at Fermilab*, *JINST* **7** (2012) P10019, [[arXiv:1205.6747](#)].
- [9] **ArgoNeuT** Collaboration, R. Acciarri et al., *A study of electron recombination using highly ionizing particles in the ArgoNeuT Liquid Argon TPC*, *JINST* **8** (2013) P08005, [[arXiv:1306.1712](#)].
- [10] **ArgoNeuT** Collaboration, C. Anderson et al., *First Measurements of Inclusive Muon Neutrino Charged Current Differential Cross Sections on Argon*, *Phys.Rev.Lett.* **108** (2012) 161802, [[arXiv:1111.0103](#)].
- [11] **ArgoNeuT** Collaboration, R. Acciarri et al., *Measurements of Inclusive Muon Neutrino and Antineutrino Charged Current Differential Cross Sections on Argon in the NuMI Antineutrino Beam*, *Phys.Rev.* **D89** (2014) 112003, [[arXiv:1404.4809](#)].
- [12] **ArgoNeuT** Collaboration, R. Acciarri et al., *First Measurement of Neutrino and Antineutrino Coherent Charged Pion Production on Argon*, *Phys. Rev. Lett.* **113** (2014), no. 26 261801, [[arXiv:1408.0598](#)].
- [13] **LAr1-ND, ICARUS-WA104, MicroBooNE** Collaboration, M. Antonello et al., *A Proposal for a Three Detector Short-Baseline Neutrino Oscillation Program in the Fermilab Booster Neutrino Beam*, [arXiv:1503.0152](#).
- [14] **DUNE** Collaboration, R. Acciarri et al., *Long-Baseline Neutrino Facility (LBNF) and Deep Underground Neutrino Experiment (DUNE)*, [arXiv:1512.0614](#).
- [15] **MiniBooNE** Collaboration, A. A. Aguilar-Arevalo et al., *Unexplained Excess of Electron-Like Events From a 1-GeV Neutrino Beam*, *Phys. Rev. Lett.* **102** (2009) 101802, [[arXiv:0812.2243](#)].
- [16] C. Adams et al., *LAr1-ND: Testing Neutrino Anomalies with Multiple LArTPC Detectors at Fermilab*, [arXiv:1309.7987](#).
- [17] A. Marchionni, *Status and New Ideas Regarding Liquid Argon Detectors*, *Ann.Rev.Nucl.Part.Sci.* **63** (2013) 269–290, [[arXiv:1307.6918](#)].

- 1 [18] **MiniBooNE** Collaboration, A. A. Aguilar-Arevalo et al., *The Neutrino Flux prediction at*
 2 *MiniBooNE*, *Phys. Rev.* **D79** (2009) 072002, [[arXiv:0806.1449](https://arxiv.org/abs/0806.1449)].
- 3 [19] P. Adamson et al., *The NuMI Neutrino Beam*, *Nucl. Instrum. Meth.* **A806** (2016) 279–306,
 4 [[arXiv:1507.0669](https://arxiv.org/abs/1507.0669)].
- 5 [20] **MiniBooNE** Collaboration, A. Aguilar-Arevalo et al., *The Neutrino Flux prediction at MiniBooNE*,
 6 *Phys. Rev.* **D79** (2009) 072002, [[arXiv:0806.1449](https://arxiv.org/abs/0806.1449)].
- 7 [21] E. Aprile, K. Giboni, and C. Rubbia, *A study of ionization electrons drifting large distances in liquid*
 8 *and solid argon*, *Nuclear Instruments and Methods in Physics Research Section A: Accelerators,*
 9 *Spectrometers, Detectors and Associated Equipment* **241** (1985), no. 1 62 – 71.
- 10 [22] C. Tegeler, R. Span, and W. Wagner, *A new equation of state for argon covering the fluid region for*
 11 *temperatures from the melting line to 700 k at pressures up to 1000 mpa*, *Journal of Physical and*
 12 *Chemical Reference Data* **28** (1999), no. 3 779–850.
- 13 [23] R. L. Amey and R. H. Cole, *Dielectric constants of liquefied noble gases and methane*, *The Journal*
 14 *of Chemical Physics* **40** (1964), no. 1 146–148.
- 15 [24] C. Amsler et al., *Review of particle physics*, *Physics Letters B* **667** (2008), no. 1–5 1 – 6.
 16 <http://pdg.lbl.gov/2008/reviews/rpp2008-rev-passage-particles-matter.pdf>.
- 17 [25] E. Shibamura, A. Hitachi, T. Doke, T. Takahashi, S. Kubota, and M. Miyajima, *Drift velocities of*
 18 *electrons, saturation characteristics of ionization and W-values for conversion electrons in liquid*
 19 *argon, liquid argon-gas mixtures and liquid xenon*, *Nuclear Instruments and Methods* **131** (Dec.,
 20 1975) 249–258.
- 21 [26] M. Miyajima, T. Takahashi, S. Konno, T. Hamada, S. Kubota, H. Shibamura, and T. Doke, *Average*
 22 *energy expended per ion pair in liquid argon*, *Phys. Rev. A* **9** (Mar, 1974) 1438–1443.
- 23 [27] E. Shibamura, T. Takahashi, S. Kubota, and T. Doke, *Ratio of diffusion coefficient to mobility for*
 24 *electrons in liquid argon*, *Phys. Rev. A* **20** (Dec, 1979) 2547–2554.
- 25 [28] S. Derenzo, *Electron diffusion and positive ion charge retention in liquid-filled high-resolution*
 26 *multistrip ionization-mode chambers*, *Lawrence Berkeley Laboratory, Group A Physics Note No.*
 27 786 (1974).
- 28 [29] V. Atrazhev and I. Timoshkin, *Transport of electrons in atomic liquids in high electric fields*, *IEEE*
 29 *Transactions on Dielectrics and Electrical Insulation* **5** (June, 1998) 450 –457.
- 30 [30] M. Adamowski et al., *The Liquid Argon Purity Demonstrator*, *JINST* **9** (2014) P07005,
 31 [[arXiv:1403.7236](https://arxiv.org/abs/1403.7236)].
- 32 [31] *ASME boiler and pressure vessel code: an international code. I - Rules for construction of power*
 33 *boiler*. American Society of Mechanical Engineers, New York, NY, 2010.
- 34 [32] B. Jones, C. Chiu, J. Conrad, C. Ignarra, T. Katori, et al., *A Measurement of the Absorption of Liquid*
 35 *Argon Scintillation Light by Dissolved Nitrogen at the Part-Per-Million Level*, *JINST* **8** (2013)
 36 P07011, [[arXiv:1306.4605](https://arxiv.org/abs/1306.4605)].
- 37 [33] Barber-Nichols Inc., 6325 West 55th Avenue, Arvada, CO 80002 USA.
- 38 [34] Sigma-Aldrich, P.O. Box 14508, St. Louis, MO 63178 USA.
- 39 [35] BASF Corp., 100 Park Avenue, Florham Park, NJ 07932 USA.
- 40 [36] Aspen Aerogels, Inc., 30 Forbes Road, Building B, Northborough, MA 01532 USA.

- 1 [37] G. Carugno et al., *Electron lifetime detector for liquid argon*, *Nuclear Instruments and Methods in*
 2 *Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment* **292**
 3 (1990) 580.
- 4 [38] L. F. Bagby et al., *Breakdown voltage of metal-oxide resistors in liquid argon*, *JINST* **9** (2014),
 5 no. 11 T11004, [[arXiv:1408.4013](#)].
- 6 [39] **MicroBooNE** Collaboration, R. Acciarri et al., *Liquid Argon Dielectric Breakdown Studies with the*
 7 *MicroBooNE Purification System*, *JINST* **9** (2014), no. 11 P11001, [[arXiv:1408.0264](#)].
- 8 [40] J. Asaadi, J. M. Conrad, S. Gollapinni, B. J. P. Jones, H. Jostlein, J. M. St. John, T. Strauss,
 9 S. Wolbers, and J. Zennamo, *Testing of High Voltage Surge Protection Devices for Use in Liquid*
 10 *Argon TPC Detectors*, *JINST* **9** (2014) P09002, [[arXiv:1406.5216](#)].
- 11 [41] R. Acciarri et al., *Construction and Assembly of the Wire Planes for the MicroBooNE Time*
 12 *Projection Chamber*, [arXiv:1609.0616](#).
- 13 [42] W. Buescher, *Spectrum Laboratory*, . <http://www.qsl.net/dl4yhf/spectra1.html>.
- 14 [43] B. Rebel et al., *Results from the Fermilab materials test stand and status of the liquid argon purity*
 15 *demonstrator*, *J.Phys.Conf.Ser.* **308** (2011) 012023.
- 16 [44] “Citranox acid detergent.”
 17 http://www.alconox.com/resources/StandardDocuments/MSDS/msds_citranox_english_ghs.pdf.
- 18 [45] “Simple Green detergent.” http://simplegreen.com/pdfs/MSDS_EN-US_AllPurposeCleaner.pdf.
- 19 [46] “Elma Clean 65.” [http://www.elma-](http://www.elma-ultrasonic.com/fileadmin/downloads/CleaningAgents/SafetyDataSheets/GB/S_EC%2065_GB.pdf)
 20 [ultrasonic.com/fileadmin/downloads/CleaningAgents/SafetyDataSheets/GB/S_EC%2065_GB.pdf](http://www.elma-ultrasonic.com/fileadmin/downloads/CleaningAgents/SafetyDataSheets/GB/S_EC%2065_GB.pdf).
- 21 [47] **MicroBooNE** Collaboration, T. Katori, *The MicroBooNE light collection system*, *JINST* **8** (2013)
 22 C10011, [[arXiv:1307.5256](#)].
- 23 [48] L. Bugel, J. M. Conrad, C. Ignarra, B. J. P. Jones, T. Katori, T. Smidt, and H. K. Tanaka,
 24 *Demonstration of a Lightguide Detector for Liquid Argon TPCs*, *Nucl. Instrum. Meth. A* **640** (2011)
 25 69–75, [[arXiv:1101.3013](#)].
- 26 [49] **MicroBooNE** Collaboration, T. Katori, *MicroBooNE, A Liquid Argon Time Projection Chamber*
 27 (*LaTPC*) *Neutrino Experiment*, *AIP Conf. Proc.* **1405** (2011) 250–255, [[arXiv:1107.5112](#)].
- 28 [50] C. Chiu, C. Ignarra, L. Bugel, H. Chen, J. Conrad, et al., *Environmental Effects on TPB*
 29 *Wavelength-Shifting Coatings*, *JINST* **7** (2012) P07007, [[arXiv:1204.5762](#)].
- 30 [51] B. Baptista, L. Bugel, C. Chiu, J. Conrad, C. Ignarra, et al., *Benchmarking TPB-coated Light Guides*
 31 *for Liquid Argon TPC Light Detection Systems*, [arXiv:1210.3793](#).
- 32 [52] T. Briese, L. Bugel, J. Conrad, M. Fournier, C. Ignarra, et al., *Testing of Cryogenic Photomultiplier*
 33 *Tubes for the MicroBooNE Experiment*, *JINST* **8** (2013) T07005, [[arXiv:1304.0821](#)].
- 34 [53] B. Jones, T. Alexander, H. Back, G. Collin, J. Conrad, et al., *The Effects of Dissolved Methane upon*
 35 *Liquid Argon Scintillation Light*, *JINST* **8** (2013) P12015, [[arXiv:1308.3658](#)].
- 36 [54] Z. Moss, L. Bugel, G. Collin, J. Conrad, B. Jones, et al., *Improved TPB-coated Light Guides for*
 37 *Liquid Argon TPC Light Detection Systems*, [arXiv:1410.6256](#).
- 38 [55] **MicroBooNE** Collaboration, J. Conrad, B. Jones, Z. Moss, T. Strauss, and M. Toups, *The*
 39 *Photomultiplier Tube Calibration System of the MicroBooNE Experiment*, *JINST* **10** (2015), no. 06
 40 T06001, [[arXiv:1502.0415](#)].

- [56] Z. Moss, M. Toups, L. Bugel, G. Collin, and J. Conrad, *Anode-Coupled Readout for Light Collection in Liquid Argon TPCs*, [arXiv:1507.0199](https://arxiv.org/abs/1507.0199).
- [57] B. J. P. Jones, *Sterile Neutrinos in Cold Climates*. PhD thesis, MIT, 2015.
<http://inspirehep.net/record/1389805>.
- [58] A. Hitachi, T. Takahashi, N. Funayama, K. Masuda, J. Kikuchi, and T. Doke, *Effect of ionization density on the time dependence of luminescence from liquid argon and xenon*, *Phys. Rev. B* **27** (May, 1983) 5279–5285.
- [59] WArP Collaboration, R. Acciarri et al., *Effects of Nitrogen contamination in liquid Argon*, *JINST* **5** (2010) P06003, [[arXiv:0804.1217](https://arxiv.org/abs/0804.1217)].
- [60] M. Suzuki and S. Kubota, *Mechanism of proportional scintillation in argon, krypton and xenon*, *Nuclear Instruments and Methods* **164** (1979), no. 1 197 – 199.
- [61] N. Ishida, M. Chen, T. Doke, K. Hasuike, A. Hitachi, M. Gaudreau, M. Kase, Y. Kawada, J. Kikuchi, T. Komiyama, K. Kuwahara, K. Masuda, H. Okada, Y. H. Qu, M. Suzuki, and T. Takahashi, *Attenuation length measurements of scintillation light in liquid rare gases and their mixtures using an improved reflection suppresser*, *Nuclear Instruments and Methods in Physics Research A* **384** (Feb., 1997) 380–386.
- [62] Hamamatsu, “Large area photomultipier tubes.” Specification for R5912 and R5912-02 photomultiplier tubes:
http://www.hamamatsu.com/resources/pdf/etd/LARGE_AREA_PMT TPMH1286E05.pdf.
- [63] E. H. Bellamy, G. Bellettini, F. Gervelli, M. Incagli, D. Lucchesi, C. Pagliarone, F. Zetti, Yu. Budagov, I. Chirikov-Zorin, and S. Tokar, *Absolute calibration and monitoring of a spectrometric channel using a photomultiplier*, *Nucl. Instrum. Meth. A* **339** (1994) 468–476.
- [64] H. O. Meyer, *Dark Rate of a Photomultiplier at Cryogenic Temperatures*, [arXiv:0805.0771](https://arxiv.org/abs/0805.0771).
- [65] H. Meyer, *Spontaneous electron emission from a cold surface*, *EPL* (2010), no. 89 58001.
- [66] Hamamatsu Photonics, *Photomultiplier Tubes, Basics and Applications, Third Edition, Chapter 13.9*, . https://www.hamamatsu.com/resources/pdf/etd/PMT_handbook_v3aE.pdf.
- [67] W. Burton and B. Powell, *Fluorescence of tetraphenyl-butadiene in the vacuum ultraviolet.*, *Appl Opt.* **12** (1973), no. 1 87–9.
- [68] C. M. Ignarra, *Sterile Neutrino Searches in MiniBooNE and MicroBooNE*. PhD thesis, MIT, Cambridge, 2014. <http://inspirehep.net/record/1316612>.
- [69] V. Gehman, S. Seibert, K. Rielage, A. Hime, Y. Sun, et al., *Fluorescence Efficiency and Visible Re-emission Spectrum of Tetraphenyl Butadiene Films at Extreme Ultraviolet Wavelengths*, *Nucl.Instrum.Meth. A* **654** (2011) 116–121, [[arXiv:1104.3259](https://arxiv.org/abs/1104.3259)].
- [70] R. Francini, R. Montereali, E. Nichelatti, M. Vincenti, N. Canci, et al., *VUV-Vis optical characterization of Tetraphenyl-butadiene films on glass and specular reflector substrates from room to liquid Argon temperature*, *JINST* **8** (2013) P09006.
- [71] B. J. P. Jones, *A simulation of the optical attenuation of TPB coated light-guide detectors*, *JINST* **8** (2013) C10015, [[arXiv:1307.6906](https://arxiv.org/abs/1307.6906)].
- [72] Sundance Supply, “Lexan Thermoclear Multiwall Polycarbonate Sheet.”
<https://www.sundancesupply.com/Lexan.pdf>.
- [73] “Stycast 2850.” <http://www.ellsworth.com/manufacturers/emerson-cuming>.

- 1 [74] J. B. Birks and K. N. Kuchela, *Energy transfer in fluorescent plastic solutions*, *Disc. Faraday. Soc* **27**
 2 (1959) 57–63.
- 3 [75] J. B. Birks and K. N. Kuchela, *Energy transfer in organic systems ii: Solute-solute transfer in liquid*
 4 *solutions*, *Proc. Phys. Soc* **77** (1961) 1083.
- 5 [76] S. Hanagodimath, B. Siddlingeshwar, J. Thipperudrappa, S. Kumar, and B. Hadimani,
 6 *Fluorescence-quenching studies and temperature dependence of fluorescence quantum yield, decay*
 7 *time and intersystem crossing activation energy of tpb*, *J. Lumin* **129** (2008), no. 4 335–339.
- 8 [77] L. Liu, “Time resolved fluorescence studies of dye-polymers as excited by laser and beta radiation.”
 9 PhD Thesis at Texas Tech University -
 10 <http://repositories.tdl.org/ttu-ir/bitstream/handle/2346/13744/31295012500558.pdf?sequence=1>.
- 11 [78] Hamamatsu, “Photomultiplier tube r7723, r7724, r7725.”
 12 <http://www.hamamatsu.com/resources/pdf/etd/R7723 TPMH1315E01.pdf>.
- 13 [79] **ICARUS** Collaboration, A. Ankowski et al., *Energy reconstruction of electromagnetic showers from*
 14 *pi0 decays with the ICARUS T600 Liquid Argon TPC*, *Acta Phys. Polon.* **B41** (2010) 103–125,
 15 [[arXiv:0812.2373](https://arxiv.org/abs/0812.2373)].
- 16 [80] A. Ereditato Private Communication.
- 17 [81] K. Scholberg, *The SuperNova Early Warning System*, *Astron. Nachr.* **329** (2008) 337–339,
 18 [[arXiv:0803.0531](https://arxiv.org/abs/0803.0531)].
- 19 [82] “LArSoft.” <http://larsoft.org>.
- 20 [83] M. Massie et al., *Monitoring with Ganglia: Tracking Dynamic Host and Application Metrics at*
 21 *Scale*. O'Reilly Media, 2012.
- 22 [84] “Experimental physics and industrial control system.” <http://www.aps.anl.gov/epics/>.
- 23 [85] R. Brun and F. Rademakers, *Root — an object oriented data analysis framework*, *Nuclear*
 24 *Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and*
 25 *Associated Equipment* **389** (1997), no. 1 81 – 86.
- 26 [86] “Control system studio.” <http://controlsystemstudio.org/>; developer’s manual:
 27 <http://cs-studio.sourceforge.net/docbook/>.
- 28 [87] *Channel access protocol specification*, September, 2014. <http://www.aps.anl.gov/epics/docs/ca.php>.
- 29 [88] “Net-snmp.” <http://www.net-snmp.org>.
- 30 [89] J. Priller, *Epics snmp device support module: Release 1.0.0*, April, 2014.
 31 <https://groups.nscl.msu.edu/controls/>.
- 32 [90] M. Rivers, *asyndriver: Asynchronous driver support*, September, 2014.
 33 <http://www.aps.anl.gov/epics/modules/soft/asy/>.
- 34 [91] D. Zimoch, *Epics streamdevice*, September, 2014.
 35 <http://epics.web.psi.ch/software/streamdevice/doc/>.
- 36 [92] *Glomation, inc., gesbc-9g20 user’s manual, version 0.2*, June, 2009.
 37 <http://www.glamationinc.com/GESBC-9G20-user-manual.pdf>.
- 38 [93] D. Huffman, *Peripheral interface: Fan and temperature probe*, September, 2014.
 39 http://www-ppd.fnal.gov/EEDOffice-W/Infrastructure_group/Huffman/WEB/uboone/pia.html.
- 40 [94] *Ganglia*, September, 2014. <http://ganglia.info/>.

- 1 [95] M. Newville, *Pyepics: Epics channel access for python*, September, 2014.
 2 http://cars9.uchicago.edu/software/python/pyepics3/.
- 3 [96] *Gnu freeipmi*, September, 2014. http://www.gnu.org/software/freeipmi/.
- 4 [97] I. Mandrichenko, A. Norman, A. Petrov, and V. Podstavkov, *Implementation of beam conditions*
 5 *database for intensity frontier experiments*, June, 2012.
 6 http://cd-docdb.fnal.gov/cgi-bin>ShowDocument?docid=4781.
- 7 [98] K. McDonald and H. Jostlein, “Path Length of Muons Traversing an Arbitrary Volume.”
 8 http://physics.princeton.edu/mcdonald/examples/muonpath.pdf.
- 9 [99] S. Palestini and K. McDonald, “Space Charge in Ionization Detectors.”
 10 http://physics.princeton.edu/mcdonald/examples/spacecharge.pdf.
- 11 [100] I. Badhrees et al., *Measurement of the two-photon absorption cross-section of liquid argon with a*
 12 *time projection chamber*, *New J.Phys.* **12** (2010) 113024, [[arXiv:1011.6001](#)].
- 13 [101] I. Badhrees et al., *Argontube: An R&D Liquid Argon Time Projection Chamber*, *JINST* **7** (2012)
 14 C02011.
- 15 [102] A. Ereditato et al., *Design and operation of ARGONTUBE: a 5 m long drift liquid argon TPC*,
 16 *JINST* **8** (2013) P07002, [[arXiv:1304.6961](#)].
- 17 [103] M. Zeller et al., *First measurements with ARGONTUBE, a 5m long drift Liquid Argon TPC*,
 18 *Nucl.Instrum.Meth.* **A718** (2013) 454–458.
- 19 [104] A. Ereditato et al., *Measurement of the drift field in the ARGONTUBE LAr TPC with 266 nm pulsed*
 20 *laser beams*, [arXiv:1408.6635](#).
- 21 [105] A. Ereditato et al., *A steerable UV laser system for the calibration of liquid argon time projection*
 22 *chambers*, [arXiv:1406.6400](#).
- 23 [106] “Continuum Lasers.” <http://www.continuumlasers.com>.
- 24 [107] I. Badhrees, A. Ereditato, I. Kreslo, M. Messina, U. Moser, B. Rossi, M. S. Weber, M. Zeller,
 25 C. Altucci, S. Amoruso, R. Buzzese, and R. Velotta, *Measurement of the two-photon absorption*
 26 *cross-section of liquid argon with a time projection chamber*, *New Journal of Physics* **12** (2010),
 27 no. 11 113024.
- 28 [108] A. Ereditato, I. Kreslo, M. Lüthi, C. R. von Rohr, M. Schenk, T. Strauss, M. Weber, and M. Zeller, *A*
 29 *steerable uv laser system for the calibration of liquid argon time projection chambers*, *Journal of*
 30 *Instrumentation* **9** (2014), no. 11 T11007.