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## 1 Introduction

In order to conserve angular momentum, the probability that a nucleus relaxing toward the ground state will emit a photon in a given direction is dependent on the angle between the direction of emission and the spin axis of the nucleus [1]. When studying a free nucleus, the nuclear spin axis is free to rotate in space and thus  $\gamma$  emissions are isotropic. If relaxation proceeds through successive  $\gamma$  emissions however, the two emissions will be angularly correlated if there isn't enough time for the projection of the nuclear spin to be perturbed and rotated by extra-nuclear fields in between the two emissions. The aim of this paper is to verify the angular correlation expected from the  $\gamma$ - $\gamma$  cascade of decaying  $^{60}\text{Co}$ . To do so, a FAST-SLOW coincidence circuit was initially set up and the  $\gamma$  spectrum of a  $^{60}\text{Co}$  sample was measured. The number of angularly correlated photons for different angles were then measured.[5]

## 2 Theory

A more detailed treatment of the theory of the angular correlation of nuclear radiation is provided in chapter XIX of [1]. Figure 1 shows the decay scheme of the sample used. The  $^{60}\text{Co}$  with half-life of  $\approx 5.3$  years decays via  $\beta^-$ -radiation into  $^{60}\text{Ni}$  with highest probability to the  $4^+$  angular momentum state. The Ni then relaxes to the ground state with the largest branching ratio being the  $4^+ \rightarrow 2^+ \rightarrow 0^+$  decay, producing two  $\gamma$  photons with respective energies of 1.17 and 1.33 MeV. The lifetime of the intermediate  $2^+$  Ni state is on the order of 1 ps which is too short to be perturbed by extra-nuclear fields and thus

allowing for the angular correlation treatment. Decaying to the ground state follows via emitting one or more  $\gamma$  photons. The relaxing process with the largest branching ratio is the  $4^+ \rightarrow 2^+ \rightarrow 0^+$  decay, produces two  $\gamma$  photons. The corresponding lifetimes are short enough for the emissions to be considered coincidental (from the detecting electronics point of view), and for the assumption that extra-nuclear forces do not cause perturbation in the correlation between the photons. As the direction of the emission of

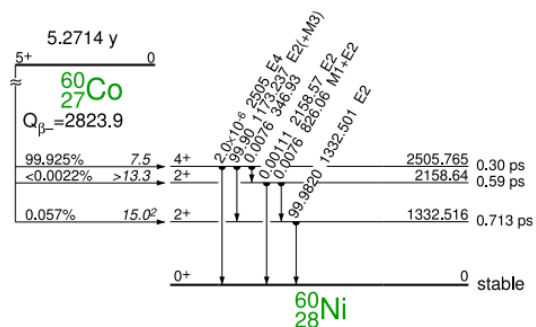


Figure 1:  $^{60}\text{Co}$  decay scheme [3]

the first photon is encoded with information on the projection of the spin axis of the nucleus due to angular momentum conservation, the probability of the direction of emission of the second photon depends on the direction of the first emission. Thus, there is an angular correlation between the directions of the two  $\gamma$  emissions, given by the following directional correlation function for the  $\gamma$ - $\gamma$  cascade herein considered:

$$W(\theta) = 1 + A_{22}P_2(\cos \theta) + A_{44}P_4(\cos \theta),$$

where  $P_2$  and  $P_4$  are Legendre Polynomials,  $A_2 = 0.1020$ , and  $A_4 = 0.0091$  [1].

### 3 Experimental setup

The design of the apparatus is shown in Figure 2.

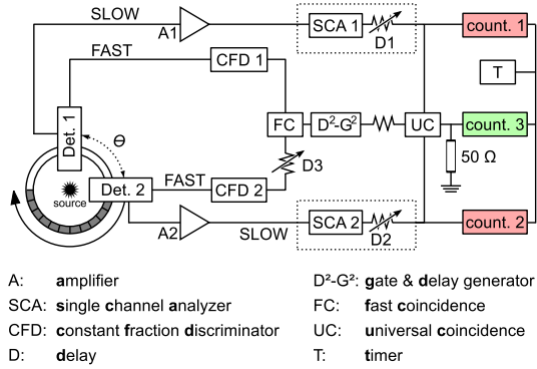


Figure 2: Experimental setup [2]

#### 3.1 Detector

The detectors consist of a crystal scintillator and a photomultiplier. The scintillator absorbs the  $\gamma$ -ray and re-emits visible light in form of scintillation, which induces electron emission in the photomultiplier via the photoelectric effect. The high voltage provided between the (photo-)cathode and the anode accelerate the electron, which induces an avalanche of electrons by colliding with each of the dynodes, as depicted in Figure 3. As this procedure distorts signal shape, which is required for the measurement of the energy absorbed by the scintillator, the signal on one of the first dynodes is also used as the input for the Single Channel Analyzers **SCA1** and **SCA2**.

#### 3.2 The fast coincidence circuit: Constant Fraction Discriminators

The fast coincidence circuit checks whether the two detected photons come from the same decay process.

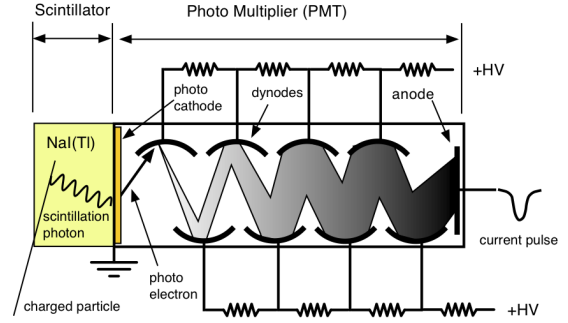


Figure 3: Scintillation detector with photomultiplier [4]

The constant fraction discriminators modify the signals as seen in Figure 4: an attenuated inverted copy of the input is added to the delayed input signal. The resulting shape crosses the 0 V line ("zero crossing point"). This point may serve as a time stamp for further processing because the result does not depend on the amplitude.

The discriminators have two outputs: while the "fast" output is a negative pulse, the "slow" output is a positive one, hence the latter is used by the fast coincidence unit **FC**, which does the actual coincidence checking. The output is fed to a gate & delay generator **D<sup>2</sup>-G<sup>2</sup>**, which delays the gate signal ("trigger" signal): this is required to match the output of the slow coincidence circuit, described in the next section.

#### 3.3 The slow coincidence circuit: Single-Channel Analyzers

The weak detector signal coming from the earlier dynode is fed through the amplifiers **A1** and **A2** first. The analyzers **SCA1**, **SCA2** determine whether the amplitudes of the signals — which are proportional to the energy of the  $\gamma$ -rays — fall into an interval with adjustable upper and lower limits.

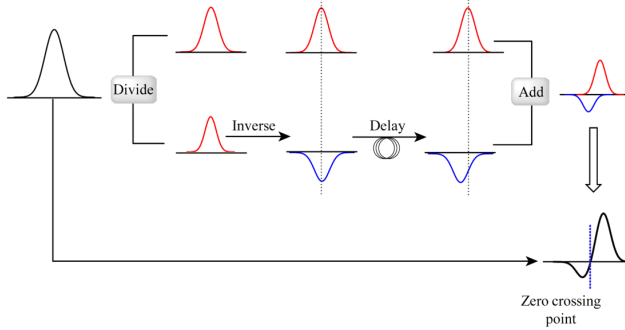


Figure 4: Scintillation detector with photomultiplier [5]

## 4 Procedure

### 4.1 Amplifiers, constant-fraction discriminators

First the amplifiers  $A_1$ ,  $A_2$  were adjusted. We changed the gain using an oscilloscope such that the peaks corresponding to different detected photon energy levels don't hit the 9 V output ceiling of the amplifiers. The result of the adjustment is shown on Figure 5. Next we set the threshold of constant-fraction discriminators above the noise level by using the output of the CFD as the trigger for the amplifier signal of the same detector, increasing the threshold until the bright line at 0 V showing no photon detection (e.g. noise) disappeared.

### 4.2 Fast coincidence circuit, prompt curve

We connected the CFDs' positive outputs to the fast coincidence unit **FC**, one of them through a delay unit **D3**. Keeping the resolution time of **FC** at 15 ns, we used one of the counter units and changing the delay of **D3** to measure the coincidental detections as the function of the delay. We measured for 10 s for each delay value.

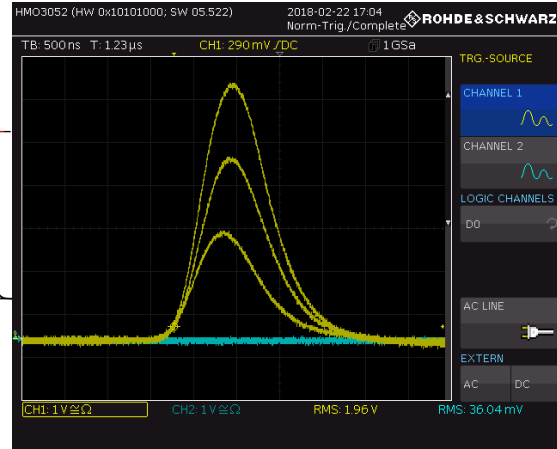


Figure 5: Well-adjusted signal of the amplifiers on channel 1. A too high gain would cause the highest peaks to flatten as saturation happens before reaching the maximum.

## References

- [1] K. Siegbahn, ALPHA-, BETA-, AND GAMMA-RAY SPECTROSCOPY, Vol. 2, North Holland Publishing Company, Amsterdam (1965).
- [2] Booklet.
- [3] R. B. Firestone, Table of Isotopes 8<sup>th</sup> edition (Wiley, New York, 1996)
- [4] [http://wanda.fiu.edu/teaching/courses/Modern\\_lab\\_manual/scintillator.html](http://wanda.fiu.edu/teaching/courses/Modern_lab_manual/scintillator.html)
- [5] [https://en.wikipedia.org/wiki/Constant\\_fraction\\_discriminator#/media/File:Operation\\_of\\_a\\_CFD.png](https://en.wikipedia.org/wiki/Constant_fraction_discriminator#/media/File:Operation_of_a_CFD.png)