

K223: Angular Correlation

Bence Mitlasóczy* and Benoît Scholtes†
Rheinische-Friedrich-Wilhelms Universität Bonn

March 11, 2018

This paper analyses the angular correlation between subsequent photon emissions in the γ - γ cascade of decaying ^{60}Co . A fast-slow coincidence circuit with scintillation detectors, single channel analysers, and constant fraction discriminators were used measure the number of γ - γ cascade decays for angular separations from 90° to 270° . The data was corrected for the measured accidental coincidences. The data was in agreement with the theoretical angular correlation function and its coefficients. The fit coefficient values calculated are $B = 0.116 \pm 0.066$ and $C = 0.040 \pm 0.075$. The large uncertainty is to be expected from previous experiments.

1 Introduction

In order to conserve angular momentum, the probability that a nucleus relaxing toward the ground state will emit a photon in a given direction is dependent on the angle between the direction of emission and the spin axis of the nucleus [1]. When studying a free nucleus, the nuclear spin axis is free to rotate in space and thus γ emissions are isotropic. If relaxation proceeds through successive γ emissions however, the two emissions will be angularly correlated if there isn't enough time for the projection of the nuclear spin to be perturbed and rotated by extra-nuclear fields in between the two emissions. The aim of this paper is to verify the angular correlation expected from the γ - γ cascade of decaying ^{60}Co . To do so, a FAST-SLOW coincidence circuit was initially set up and the γ spectrum of a ^{60}Co sample was measured. The number of angularly correlated photons for different angles were then measured. The angular correlation of decaying photons sheds light on the properties of the nucleus and nuclear decay, such as the spins of nuclear levels and the angular momentum carried by radiation.

2 Theory

A detailed treatment of the theory of the angular correlation of nuclear radiation is provided in chapter XIX of [1]. Figure 1 shows the decay scheme of the sample used. The ^{60}Co with half-life of ≈ 5.3 years decays via β^- -radiation into ^{60}Ni with highest probability to the 4^+ angular momentum state. The Ni then relaxes to the ground state with the largest branching ratio being the $4^+ \rightarrow 2^+ \rightarrow 0^+$ decay, producing two γ photons with respective energies of 1.17 and 1.33 MeV. The lifetime of the intermediate 2^+ Ni state is on the order of 1 ps which is too short to be perturbed by extra-nuclear fields and thus allowing for the angular correlation treatment. As the direction of the emission of the first photon is encoded with information on the projection of the spin axis of the nucleus due to angular momentum conservation, the probability of the direction of emission of the second photon depends on the direction of the first emission. Thus, there is an angular correlation between the directions of the two γ emissions, given by the following directional correlation function for the γ - γ cascade herein considered:

$$W(\theta) = 1 + A_{22}P_2(\cos\theta) + A_{44}P_4(\cos\theta), \quad (1)$$

*s6bemitl@uni-bonn.de

†s6bescho@uni-bonn.de

the fast circuit to measure the arrival time of the incident photon. As this procedure distorts the signal shape, which is required for the measurement of the energy of the incident photon, the output pulse cannot be used to measure the energy. Instead, an undistorted signal from an earlier phase of the amplification process is used as the input for the single channel analyzer (**SCA**) in the slow circuit.

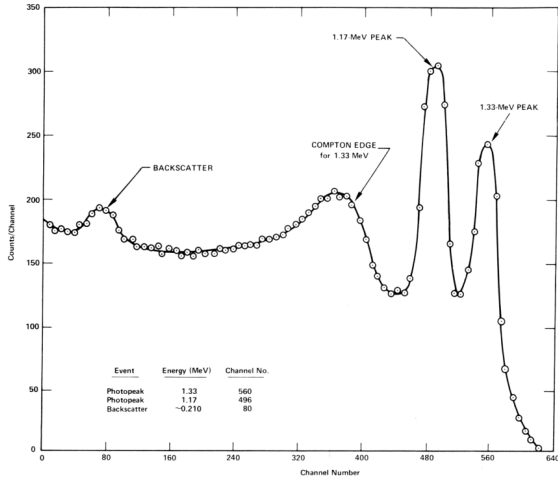


Figure 4: Compton scattering in the ^{60}Co γ - γ cascade spectrum [7]

One should note that Compton scattering also occurs in the scintillator when incident photons scatter inelastically inside the crystal instead of being absorbed. Scintillation photons are then produced with lower energy than those produced via absorption. This creates a Compton plateau spectrum bordered by two peaks on either side, as seen in Figure 4. The Compton edge is caused by the photons backscattering in the crystal, depositing the maximum of energy in the crystal from Compton scattering. The backscattered peak is instead caused by photons scattering back into the detector from the lead casing of the detector, with a large loss of energy, and then being absorbed by the crystal [8]. Though Compton scattering distorts the measure energy spectrum, the peaks from the absorbed photons should still be easily discernible at higher frequencies, visible in Figure 4.

3.2 Constant Fraction Discriminator

The constant fraction discriminators modify the signals as seen in Figure 5: an attenuated inverted copy of the input is added to the delayed input signal. The resulting shape crosses the 0 V line ("zero crossing point"). This point serves as a time stamp for the signal which does not depend on the amplitude of the signal. The discriminator has two out-

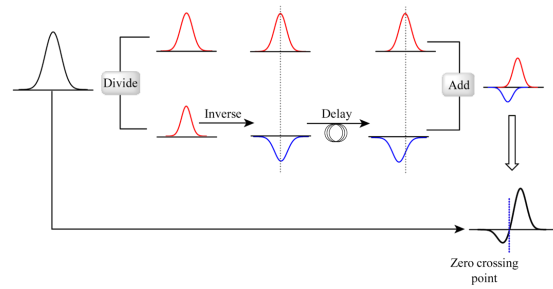


Figure 5: Scintillation detector with photomultiplier [5]

puts: the "fast" output which is a negative digital pulse and the "slow" output which is a positive one. The latter is used by the fast coincidence unit **FC**.

3.3 Single Channel Analyzers

The weak, undistorted detector signal coming from the **PMT** is fed through the amplifiers **A1** and **A2** first. The analyzers **SCA1** and **SCA2** determine whether the amplitudes of the signals, which are proportional to the energy of the detected γ photons, fall into an interval set by the **SCAs**. The interval is over a range of so-called channels of the **SCA** which are set either with upper and lower limit channels or with a lower limit and a window size.

3.4 Universal coincidences

The output signals from the **SCAs** and the **FC** are sent to the universal coincidence unit **UC** to count the total number of coincident photons within a time span set by the timer. The signals from the **SCAs**

are first sent through delays to make sure the signals incident to the **UC** correspond to photons hitting both detectors at approximately the same time. The output from the **FC** is fed to a gate & delay generator **D²-G²** which delays the fast signal to ensure that the fast coincidences are correctly assigned to the same photons measured in the slow circuit when fed into the **UC**. Thus, the **UC** counts the total number of pairs of photons which have the right energy, and arrive within a short enough time frame, which is expected from γ - γ cascade events.

4 Procedure

4.1 Solid Angle Corrections

[1] provides solid angle correction factors Q_2 and Q_4 for detectors placed 5, 7, and 10 cm from the ^{60}Co sample. These factors correct for the assumption made by the angular correlation function equation 1 that the detectors are point-like though indeed they are not. The closer the detectors are placed to the source, the larger the solid angle of the detectors become and thus the uncertainty introduced from the fact that they do not simply measure counts for a single angle is increased. That said, the further away the detectors are placed from the source, the smaller the photon count rate becomes for a given time span which increases the relative uncertainty of the count rate (as photon emissions follow Poisson statistics). Rough measurements were taken for different detector placements to estimate the expected uncertainty introduced by photon counts for the angular correlation measurements. It was decided to place the detectors 7 cm away from the source in order to minimise uncertainty, though whether this is indeed the best placement cannot be known due to the unknown nature of the uncertainty introduced by the solid angle correction. As the solid angle correction factors depend on the energy of the incident photons and [1] only provides correction factors for energies of 0.5, 1.0, and 1.5 MeV, linear interpolation was utilised to obtain factor values for our photons which have energies of 1.17 and 1.33 MeV and thus a mean of

$E_{\text{avg}} = 1.252$ MeV. The factor values obtained are,

$$\begin{aligned} Q_2(E_{\text{avg}}) &\equiv Q_2^* = 0.9623, \\ Q_4(E_{\text{avg}}) &\equiv Q_4^* = 0.8782. \end{aligned}$$

This method reduces the uncertainty caused by the energy dependence of the solid angle correction factors. That said, there is still a small uncertainty which remains caused by the different energies of the two cascade photons which would require individual correction factors, though we do not distinguish between the two different cascade photons during measurements in this experiment.

4.2 Amplifiers, CFDs

The amplifiers **A₁**, **A₂** were then adjusted. We changed the gain using an oscilloscope such that the peaks corresponding to different detected photon energy levels don't hit the 9 V output ceiling of the amplifiers. We set the maximum to approximately 8 V. The result of the adjustment is shown in Figure 6. Next we set the threshold of constant-fraction discriminators above the noise level by using the output of the **CFD** as the trigger for the amplifier signal of the same detector, increasing the threshold until the bright line at 0 V showing no photon detection (e.g. noise) disappeared.

4.3 Prompt Curve

We connected the positive outputs of the **CFDs** to the fast coincidence unit **FC**, one of them through a delay unit **D3**. While keeping the resolution time of the **FC** at 15 ns, which was a value approximately expected, we changed the delay of **D3** to measure the coincidental detections as a function of the delay. We measured the counts for (10.0 ± 0.1) s for each delay value which was sufficient to plot an informative prompt curve, given in Figure 7.

From the figure, the count rate approximately plateaus from a delay of approximately 4 ns to 18 ns. This is comparable to the resolution time of the **FC** set at 15 ns, although a more careful investigation follows below. We set the delay of **D3** to 11 ns, roughly

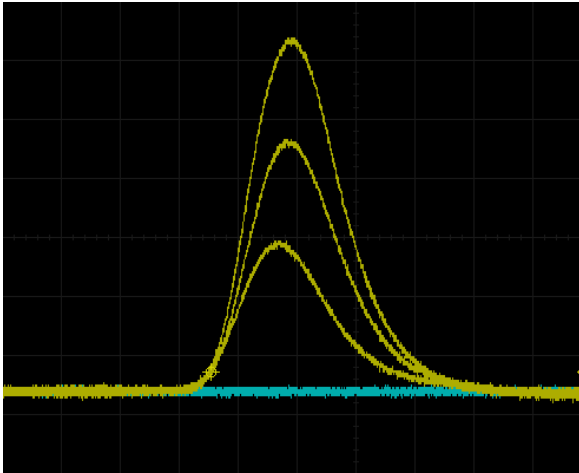


Figure 6: Well-adjusted signal of the first amplifier. A too high gain would cause the highest peaks to flatten due to saturation before the maximum.

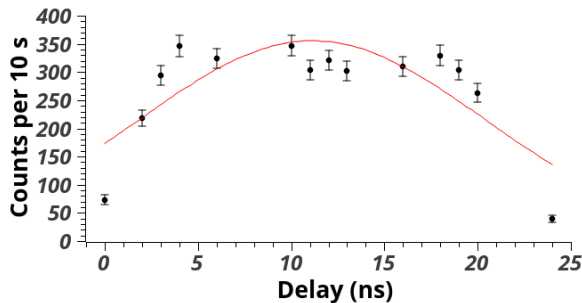


Figure 7: The measured prompt curve and a fitted Gaussian to quantify the width. The parameters are $\mu = (11.1 \pm 0.9)$ ns (mean), $\sigma = (9.3 \pm 1.3)$ ns (deviation), $A = 8259.8 \pm 881.8$ (scaling factor). Two measured values at 61 ns (17 ± 4.1 counts) and 63 ns (29 ± 5.4) are not shown.

in the middle. These values were used for the angular correlation measurements. The Gaussian fitted post factum confirmed our choice for the delay (μ), and provided the experimental resolving time of the fast coincidence circuit $\text{FWHM} = 2\sqrt{2\ln 2}\sigma = (21.9 \pm 3.0)$ ns[6]. Here one must consider the numerous error sources. First, the theoretical prompt curve would be a rectangle (see [10]), and the time

jitter manifesting in the smoothened edges are attributed to the non-idealness of the electronic components (not perfectly rectangular generated pulses, for example). If one wants to read off the resolution time, defined as the FWHM for the prompt curve, one needs to take this jitter into account; thus even though the Gaussian is not the theoretical shape, it can be used more easily and yields a more precise result than taking the length of the plateau. Secondly, there exists a constant background, caused by the accidental coincidences (two atoms radiate at the same time). This is clearly visible if one considers the fitted Gaussian at > 60 ns values, where it gives a count of $< 10^{-4}$, practically 0, while the measurements yield a substantially different result. They could be used to estimate the accidental coincidences; we only consider later the case when the two photons also have the relevant energy, see section 4.5.

4.4 Energy Spectra

The energy spectra from both detectors was plotted by scanning the entire relevant channels of both **SCAs**. The detectors were placed at 180° separation where the largest number of coincidence counts are theoretically expected. We again measured the counts for (10.0 ± 0.1) s, which was sufficient to obtain adequate spectra. A window size of 150 channels for both **SCAs** was used. Measurement were taken until the full spectrum expected from Figure 4 was obtained for both detectors. The spectra are given in Figure 8. The two γ - γ cascade photon peaks are clearly visible for both detectors at the upper end of the spectra, as well as the backscatter peaks at the lower end of the spectra and Compton edges around channel 450. The channel windows chosen for **SCA1** and **SCA2** were 520–720 and 540–740 respectively, in order to include both γ - γ cascade photon peaks.

We can assume that the difference between the two spectra is a constant factor due to the different instruments of the two detector systems (not investigated during the experiment). A rough estimate follows. Using the channel number as units, the first peak appears at 565 ± 15 (detector 1, black) and 595 ± 15 (detector 2). From this, $\alpha = 0.950 \pm 0.035$ results as the value for detector 1 divided by the value for

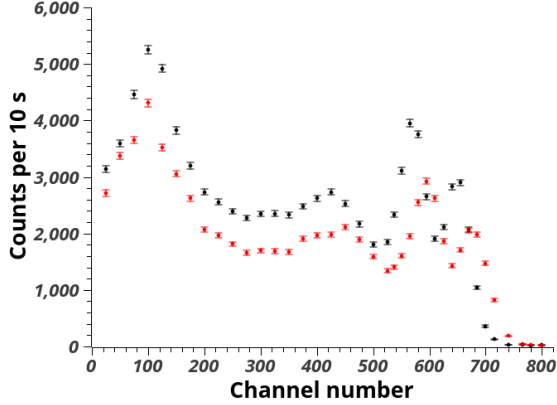


Figure 8: Measured energy spectra of both detectors 1 (fixed location, black) and 2 (red)

detector 2. Similarly, we detect the second peak for detectors 1 and 2 at 655 ± 15 and 670 ± 15 , respectively, yielding $\alpha = 0.978 \pm 0.031$.

4.5 Accidental Coincidences

Two photons might be detected as coincidental yet come from two different atoms, we need to take this noise into account. We measured the accidental coincidences overnight, and during the length of 59 820 s, 21855 ± 148 coincidences were recorded, which corresponds to 146.14 ± 0.99 counts over 400 s, or 0.3653 ± 0.0025 counts each second. The individual **SCA** units measured 260560000 ± 16142 and 243720000 ± 15612 incoming photons within the correct energy interval, or

$$N_{\text{SCA } 1} = 4355.73 \pm 0.27 \quad \text{and} \quad N_{\text{SCA } 2} = 4074.22 \pm 0.26$$

each second. Using the time resolution $\sigma = (21.9 \pm 3.0)$ ns of the **FCC** (see Figure 7), we can calculate the expected accidental coincidences (see [10] for a detailed explanation) as

$$N_{\text{acc}} = \sigma \cdot N_{\text{SCA } 1} \cdot N_{\text{SCA } 2} = 0.3886 \pm 0.0532$$

counts per second, which agrees with the experimental results.

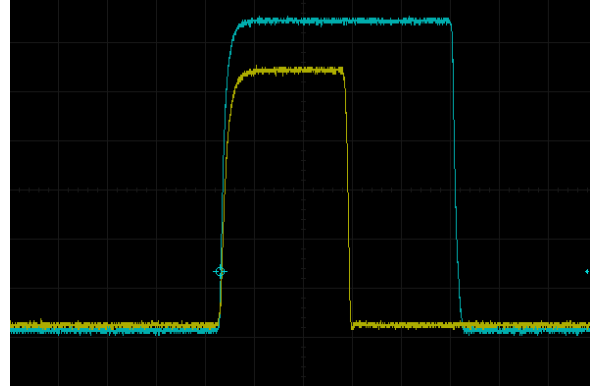


Figure 9: Adjusting the delay of the SCAs such that the leading edges of the signals from one SCA (yellow) and the gate and delay generator coincide.

5 Results

5.1 Angular Correlation

Table 1 presents the angular correlation measurements. We took measurements from 90° to 270° to span a large range of angles to obtain a more accurate weighted fit and also to test whether the angular correlation relation is indeed symmetric about 90° . If not, it would suggest that the ^{60}Co is not exactly in the centre of both detectors (or that the relation is incorrect). Spanning a larger range of angles can partially offset errors introduced from a sample not perfectly in the centre. The theoretical function is adapted with the solid angle corrections from Equation 1 to:

$$W(\theta) = A(1 + B(Q_2^*)^2 \cos^2 \theta + C(Q_4^*)^2 \cos^4 \theta), \quad (2)$$

with theoretical B and C values [1]

$$B_{\text{th}} = \frac{1}{8} = 0.125,$$

$$C_{\text{th}} = \frac{1}{24} = 0.041\bar{6}.$$

A is a simply scaling factor which depends on the measuring time. As the function is a linear combination of \cos^2 and \cos^4 , a more accurate fit is obtained

| Angle | UC | UC-norm. | UC-acc. | Unc. |
|-------|-------|----------|---------|------|
| 90 | 5,576 | 3717 | 3,571 | 61 |
| 110 | 3,837 | 3837 | 3,691 | 62 |
| 130 | 3,892 | 3892 | 3,746 | 62 |
| 150 | 4,102 | 4102 | 3,956 | 64 |
| 170 | 4,222 | 4222 | 4,076 | 65 |
| 180 | 6,331 | 4221 | 4,075 | 65 |
| 200 | 4,165 | 4165 | 4,019 | 65 |
| 220 | 4,085 | 4085 | 3,939 | 64 |
| 240 | 3,807 | 3807 | 3,661 | 62 |
| 260 | 3,797 | 3797 | 3,651 | 62 |
| 190 | 4,306 | 4306 | 4,160 | 66 |

Table 1: Data for the total universal counts **UC** at each angle for their respective measuring times mentioned below, normalised to 400 seconds **UC-norm.**, and lastly corrected for accidental counts **UC-acc.** with propagated uncertainty.

from more precise values on and adjacent to the expected peaks and troughs, while making sure there are sufficient values measured in between them. As a result, the angular correlation counts at the trough and peak expected at 90° and 180° respectively were measured longer for 600 seconds each while other measurements were taken approximately every 20° for 400 seconds. Table 1 provides the data normalised to 400 seconds. The table also includes the total coincidence counts corrected by the accidental coincidences by simply taking them away. As photons follows Poisson statistics, the uncertainty in all counts measured is simply the square root of the count. The uncertainty in the angles was 1° as a result of the imprecision of the apparatus. We fitted the measured data to Equation 2 using the least-squares-fit method, which produced

$$\begin{aligned}
A &= 3602 \pm 40, \\
B &= 0.116 \pm 0.066, \\
C &= 0.040 \pm 0.075.
\end{aligned}$$

These values are in agreement with the theoretical values expected, given above. That said, the uncertainties in the fit values are quite large, especially for C , though this was also found in previous experi-

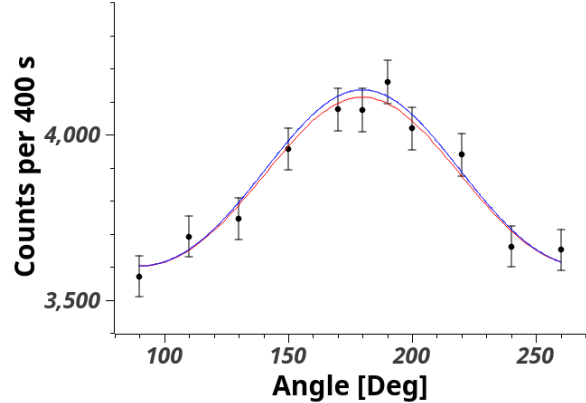


Figure 10: Angular correlation with fitted (red) and theoretical function (blue).

ments [9]. A plot of the data with the weighted fit is given in Figure 10.

6 Conclusion

We set up an experiment to analyze the angular correlation of the γ - γ cascade in a $^{60}_{27}\text{Co}$ sample. First we measured the prompt curve, where despite the low measurement times, we were able to easily find the optimal delay, confirmed by later analysis. Next, we found the relevant energy window of the individual detector energy spectra by scanning through a wide range of the single channel analyzers. We found a different interval for the two detectors, stemming from the presumable differences in the two systems, but we found the spectra to match our expectations, including the Compton edge and the two peaks. Then we proceeded to investigate the angular correlation. We achieved a somewhat good agreement with the theoretical prediction, and the main source of the errors can be easily identified: a longer measurement time would have allowed for smaller relative uncertainties, as visible in the case of the overnight recording of the accidental coincidences.

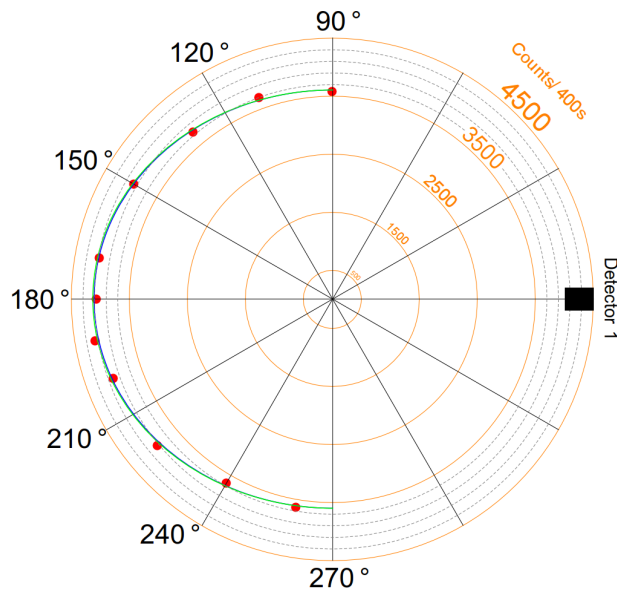


Figure 11: Polar plot of the counts against angle. Uncertainties are not displayed for the measurements. The weighted fit is given in blue, the theoretical result in green.

References

- [1] Siegbahn, K.: Alpha-, beta-, and gamma-ray spectroscopy, Vol. 2, (North Holland Publishing Company, Amsterdam, 1965).
- [2] Unspecified author: Advanced Laboratory Course (physics601): Description of Experiments, (University of Bonn, 2018).
- [3] R. B. Firestone, Table of Isotopes 8th edition, (Wiley, New York, 1996).
- [4] http://wanda.fiu.edu/teaching/courses/Modern_lab_manual/scintillator.html
- [5] https://en.wikipedia.org/wiki/Constant_fraction_discriminator#/media/File:Operation_of_a_CFD.png
- [6] Nakhostin: Signal Processing for Radiation Detectors, Wiley (2018)
- [7] <https://www.nuclear-power.net/nuclear-power/reactor-physics/interaction-radiation-matter/interaction-gamma-radiation-matter/compton-scattering/compton-edge>
- [8] https://www.phys.ufl.edu/courses/phy4803L/group_I/gamma_spec/gamspec.pdf
- [9] A. C. Melissinos & J. Napolitano, Experiments in Modern Physics, 2nd edition (Academic Press, San Diego, 2003), pp 419-21.
- [10] William R. Leo: Techniques for Nuclear and Particle Physics Experiments (Springer-Verlag, 1987)