

# Finite element methods

Making use of the weak formulation

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18.303 Linear Partial Differential Equations: Analysis and Numerics

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$$\mathcal{L}u(\mathbf{x}) = f(\mathbf{x}),$$

on some domain  $\Omega$  subject to a boundary condition

$$\mathcal{G}u(\mathbf{x}) = b(\mathbf{x})$$

on the boundary  $\mathbf{x} \in \partial\Omega$ . The weak (variational) formulation reads

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In order to find an approximate solution to this problem on a computer we choose a suitable basis  $\{\varphi_i\}$  for  $\varphi$  and  $u$ . The space for  $u$  is typically slightly bigger than the space for  $\varphi$  because of the boundaries but we'll get back to that later.

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Let's write

$$u(\mathbf{x}) = \sum_j \hat{u}_j \varphi_j(\mathbf{x}).$$

After some calculation we get

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We write  $L_{i,j} = \langle \varphi_i, \mathcal{L}\varphi_j \rangle$  and  $\hat{f}_i = \langle \varphi_i, f \rangle$  giving us a linear system

$$L\hat{u} = \hat{f},$$

where  $\hat{u}$  is a vector with the expansion coefficients for  $u$  in the basis  $\{\varphi_i\}$  and  $\hat{f}$  is a vector with  $f$  projected on the respective basis functions  $\varphi_j$ .

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**Finite element methods** (FEMs) constitute a group of techniques where the basis functions are chosen in such a way that they have a small support in real space. We'll still have to deal with the boundary conditions and some other technical details but this might be best illuminated by an example.

## Poisson equation with FEM

Let's consider the good old Poisson equation

$$\frac{d^2 u(x)}{dx^2} = f(x)$$

on the interval  $x \in (0, 1)$  with the boundary conditions  $u(0) = u_-$  and  $u(1) = u_+$ .



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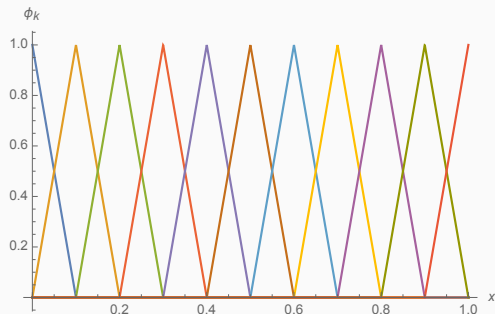
Let's choose the basis  $\varphi_k$  to be piecewise linear such that

$$\varphi_i(x_i) = 1$$

and

$$\varphi_i(x_i \pm \Delta x) = \varphi_i(x_{i \pm 1}) = 0.$$

Here  $x_i = i\Delta x$  and  $\Delta x = 1/(N + 1)$ .



The basis functions  $\varphi_k$  for  $N = 9$ .

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For the right hand side of the equation we calculate

$$\langle \varphi_i, f \rangle = \int_0^1 \varphi_i(x) f(x) dx = f_i,$$

where  $i = 1, 2, \dots, N$ .

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Now we have to evaluate the matrix

$$L_{i,j} = \langle \varphi_i, \frac{d^2}{dx^2} \varphi_j \rangle = \int_0^1 \varphi_i(x) \varphi_j''(x) dx = \left( \varphi_i(x) \varphi_j'(x) \right)_0^1 - \int_0^1 \varphi_i'(x) \varphi_j'(x) dx.$$

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Now, two functions  $\varphi_i$  and  $\varphi_j$  are overlapping only if they are the same or if  $i = j \pm 1$ . Otherwise this integral will evaluate to zero i.e.  $L$  will be tridiagonal. Furthermore, the boundary term will vanish since  $\varphi_i(0) = \varphi_i(1) = 0$ , when  $i = 1, 2, 3, \dots, N$ .

Let's calculate the diagonals first assume  $i = j$ . We have

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The function  $\varphi_i$  goes from 0 to 1 when  $x$  goes from  $x_{i-1}$  to  $x_i$ . This behavior is reversed when going from  $x_i$  to  $x_{i+1}$ . Hence we have

$$\varphi_i'(x) = \begin{cases} 1/\Delta x, & x \in (x_{i-1}, x_i) \\ -1/\Delta x, & x \in (x_i, x_{i+1}) \\ 0, & \text{otherwise.} \end{cases}$$

This gives

$$L_{i,i} = - \int_{x_i - \Delta x}^{x_i + \Delta x} \frac{1}{\Delta x^2} dx = -\frac{2}{\Delta x}.$$

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The integrand will always be negative because of the alternating behavior of the derivative of  $\varphi_i$ . We get

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We can carry out the same calculation for  $j = i + 1$  giving

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What about the boundaries?

In a similar way as before, we have a linear system for the *interior points*  $x_i$ , where  $i = 1, 2, 3, \dots, N$  but the boundaries will enter these equations.

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We require that the weak formulation holds for *all* test functions for which  $\varphi(0) = \varphi(1) = 0$ . Hence  $i$  goes only from 1 to  $N$ . However, the function  $u$  itself is not in this space of functions but includes the boundary functions  $\varphi_0$  and  $\varphi_{N+1}$ .

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$$\langle \varphi_1, \varphi_0 \rangle \hat{u}_0 = \frac{u_-}{\Delta x}.$$

For the inner product we get  $1/\Delta x$  for the same reasons as before. Now this term will enter the first equation just as when we talked about finite difference methods.

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Similarly we will have the term

$$\langle \varphi_N, \varphi_{N+1} \rangle \hat{u}_{N+1} = \frac{u_+}{\Delta x}$$

for the last row.

Now we have the equation

$$L\hat{\mathbf{u}} + \mathbf{b} = \mathbf{f},$$

where  $\mathbf{b}$  has the boundary conditions i.e.  $b_1 = u_-/\Delta x$  and  $b_N = u_+/\Delta x$  and otherwise  $b_k = 0$ .

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How does  $L$  look like?

$$L = \frac{1}{\Delta x} \begin{pmatrix} -2 & 1 & & & \\ 1 & -2 & 1 & & \\ & \ddots & \ddots & \ddots & \\ & & 1 & -2 & 1 \\ & & & 1 & -2 \end{pmatrix}.$$

Looks familiar?

Furthermore. Assume we are given  $f(x_k)$ . We have

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This integral just evaluates the area of the triangle with height 1 and width  $2\Delta x$ . Now the discretized equation is exactly the same as for our basic finite difference case.

## Some remarks

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- Notice however that we couldn't have used zeroth order piecewise functions.
- In practice the basis functions are chosen to be polynomials for which the order and other specifics are dictated by the problem at hand. The unifying theme is that they have small support so that matrix  $L$  will be sparse.

# Higher dimensions

Consider the Laplace equation with  
Dirichlet boundaries

$$\Delta u(\mathbf{x}) = 0,$$

when  $\mathbf{x} \in \Omega$  and

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on the boundary  $\mathbf{x} \in \partial\Omega$ .

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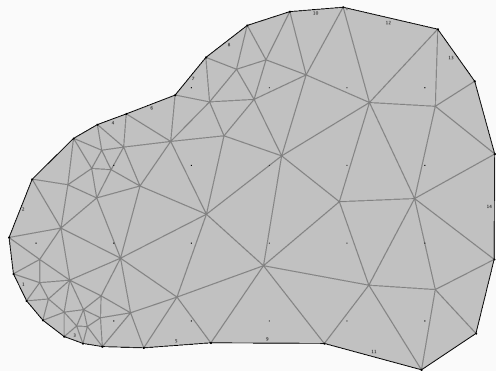
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We can define piecewise linear functions on triangles shown on the right.



Triangular mesh for the Laplace equation.

Solving the 2d equation is very similar.

- The matrix elements  $L_{i,j} = \langle \varphi_i, \varphi_j \rangle$  will only have nearest neighbor terms.



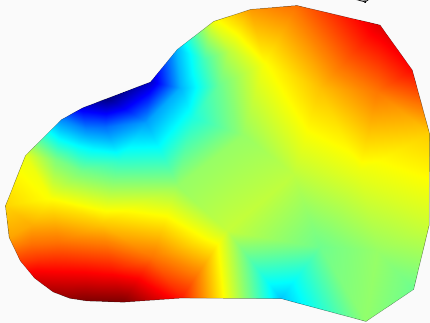
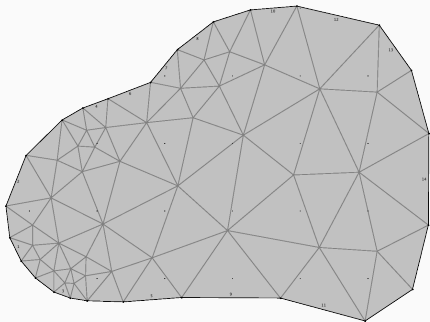
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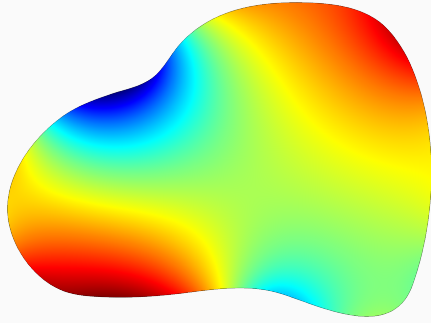
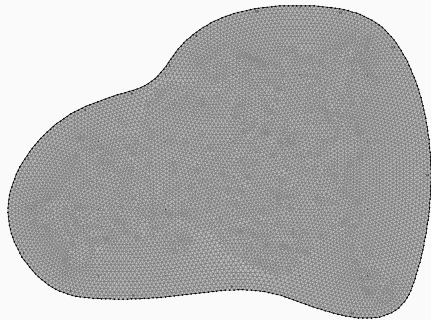
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- We will have to solve for the sparse system

$$L\hat{\mathbf{u}} = -\mathbf{b}.$$



Loose mesh.



More triangles.