

# The Heat Egn. (Pt. 3)

$$\partial_t u = \Delta u$$

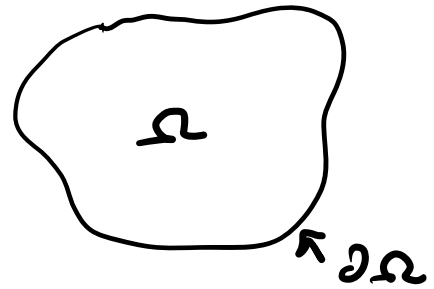
$$u|_{t=0} = g,$$

initial condition

$$u|_{\partial\Omega} = f$$

boundary condition  
(Dirichlet)

→ alt. Neumann, Robin, mixed



Step 1) Solve  $\Delta e_k = \lambda_k e_k$  w/ homogeneous B.C.'s  
replace  $f \rightarrow 0$

⇒ homogeneous soln is  $u_h = \sum_k e^{\lambda_k t} \underbrace{\langle e_k, g - u_* \rangle}_{= e^{\Delta t} (g - u_*)}$

Step 2) Solve  $\Delta u_* = 0$  w/ inhomogeneous B.C.'s

⇒ equilibrium soln is  $u_*$

Step 3) Full soln to Heat Egn. w/ inhomogeneous B.C.'s

$$u = u_* + \sum_k e^{\lambda_k t} \langle e_k, g - u_* \rangle e_k$$

B.C.'s

$$u|_{\partial\Omega} = u_*|_{\partial\Omega} = f + u_h|_{\partial\Omega} = 0 \quad \checkmark$$

I.C.

$$u|_{t=0} = u_* + g - u_* = g \quad \checkmark$$

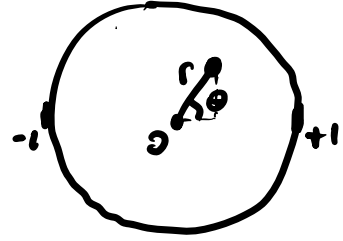
Egn.

$$\begin{aligned} \partial_t u &= \underbrace{\partial_t u_*}_{=0} + \partial_t u_h = \Delta u_h \\ &= \Delta u_h + \underbrace{\partial_t u_*}_{=0} = \Delta u \quad \checkmark \end{aligned}$$

# Diffusion in a disk

$$\partial_t u = \Delta u$$

$$u|_{t=0} = g \quad u|_{\partial\Omega} = f$$



Step 1: Solve  $\Delta u_k = 0$  with  $u_k|_{\partial\Omega} = f$

see  
lecture 1

$$u(r, \theta) = \frac{1}{\sqrt{2\pi}} \sum_{k=-\infty}^{\infty} \hat{f}_k r^{|k|} e^{ik\theta} \quad \hookrightarrow \quad \hat{f}_k = \frac{1}{\sqrt{2\pi}} \int_0^{2\pi} e^{-ik\theta} f(\theta) d\theta$$

Step 2: Solve  $\Delta e_k = \lambda_k e_k$  with  $e_k|_{\partial\Omega} = 0$ .

↓

$$\Rightarrow \partial_r^2 e_k + r^{-1} \partial_r e_k + r^{-2} \partial_\theta^2 e_k = \lambda_k e_k$$

Separation  
of  
variables

$$e_k = R(r) \Theta(\theta)$$

(drop  $k$   
for brevity)

$$\Rightarrow \underbrace{\frac{1}{R} [\partial_r^2 R + r \partial_r R - r^2 \lambda R]}_{\alpha} + \underbrace{\frac{1}{\Theta} \partial_\theta^2 \Theta}_{-\alpha} = 0$$

ODE for  $\Theta$ :

$$\partial_\theta^2 \Theta = -\alpha \Theta \quad \Theta \text{ is } 2\pi\text{-periodic}$$

$$\Rightarrow \Theta_k(\theta) = \frac{1}{\sqrt{2\pi}} e^{ik\theta} \quad -\alpha = -k^2$$

## ODE for R:

$$(*) \quad r^2 \partial_r^2 R + r \partial_r R - (r^2 \lambda + k^2) R = 0$$

This ODE eigenvalue problem (remember, we need to find both  $\lambda$  and  $R$ !) is closely related to a famous ODE called *Bessel's equation*, and its solutions are related to *Bessel functions*  $J_n(x)$ .

$$\text{With } \tilde{r} = \sqrt{-\lambda} r \quad \text{and} \quad R(r) = R(\tilde{r}/\sqrt{-\lambda}) = \tilde{R}(\tilde{r})$$

$$(*) \Rightarrow \tilde{r} \partial_{\tilde{r}}^2 \tilde{R} + \tilde{r} \partial_{\tilde{r}} \tilde{R} + (\tilde{r}^2 - k^2) \tilde{R} = 0$$

Bounded solution in disk is *first-kind Bessel function of order*  $n = k^2$ ,

$$n = k^2 \quad J_n(\tilde{r}) = \frac{1}{\pi} \int_0^\pi \cos(nz - \tilde{r} \sin z) dz$$

and our solution is  $J_n(\tilde{r}) = \tilde{R}(\tilde{r}) = R(r)$

$$R(r) = J_n(\sqrt{-\lambda} r) \quad \text{for } n = k^2$$

What about  $\lambda$ ? To find  $\lambda$ , apply B.C.'s

$$u|_{\partial\Omega} = 0 \Rightarrow 0 = R(1) = J_n(\sqrt{\lambda})$$

$\sqrt{\lambda}$  must be a root

So our eigenvalues  $\lambda$  are determined by zeros of Bessel functions! For each  $n \geq k^2$ , we get a countable set of distinct positive roots:

$$0 \leq s_{n,1} < s_{n,2} < \dots < s_{n,m} < \dots$$

The eigenfunctions / eigenvalues are therefore

$$R_{n,m}(r) = J_n(s_{n,m}r) \quad \text{and} \quad \lambda_{n,m} = -s_{n,m}^2.$$

Putting together the results with  $k = 0, \pm 1, \pm 2, \dots$   
 $m = 1, 2, 3, \dots$

$$e_{k,m}(r, \theta) = R_{k^2,m}(r) \Theta_k(\theta)$$

$$= J_{k^2}(s_{k^2,m}r) e^{ik\theta}$$

$$\lambda_{k,m} = -s_{k^2,m}^2$$

The operator exponential  $e^{\Delta t}$  on the disk is

$$u_h(r, \theta, t) = \sum_{k=-\infty}^{+\infty} \sum_{m=1}^{\infty} \tilde{g}_{k,m} e^{-(s_{k,m})^2 t} \tilde{J}_k(s_{k,m} r) e^{ik\theta}$$

Here, expansion coeffs  $\tilde{f}_{k,m}$  are such that

$$f(r, \theta) - u_x(r, \theta) = \sum_{k=-\infty}^{\infty} \sum_{m=1}^{\infty} \tilde{g}_{k,m} \tilde{J}_k(s_{k,m} r) e^{ik\theta}$$

so that  $u_h(r, \theta, 0) = g(r, \theta) - u_x(r, \theta)$  as required.

Step 3 Finally, we can stitch together the full solution from steps 1 & 2.

$$u(r, \theta, t) = u_x(r, \theta) + u_h(r, \theta, t)$$

$$= \sum_{k=-\infty}^{\infty} \hat{f}_k r^{|k|} e^{ik\theta}$$

$$+ \sum_{k=-\infty}^{\infty} \sum_{m=1}^{\infty} \tilde{g}_{k,m} e^{-(s_{k,m})^2 t} \tilde{J}_k(s_{k,m} r) e^{ik\theta}$$

$$= \sum_{k=-\infty}^{\infty} \left[ \hat{f}_k r^{|k|} + \sum_{m=1}^{\infty} \tilde{g}_{k,m} e^{-(s_{k,m})^2 t} \tilde{J}_k(s_{k,m} r) \right] e^{ik\theta}$$