

## The Heat Equation (pt. 2)

$$\frac{\partial u}{\partial t} = \gamma \frac{\partial^2 u}{\partial x^2}$$

$$\begin{array}{c} u(x,t) \\ \hline -1 \qquad u(-1,t) = u(1,t) \qquad +1 \\ \qquad \partial_x u(-1,t) = \partial_x u(1,t) \end{array}$$

$$u(x,0) = g(x)$$

"initial condition" + "boundary conditions"  
 $\Rightarrow$  Initial Boundary Value Problem

Soln.  
 $\delta = 1$

$$u(x,t) = \frac{1}{\sqrt{2}} \sum_{k=-\infty}^{\infty} \hat{g}_k e^{-(nk)^2 t} e^{inkx}$$

$\uparrow$   
Fourier coeffs

$$\hat{g}_k = \frac{1}{\sqrt{2}} \int_{-1}^{+1} e^{-inkx} g(x) dx$$

$\Rightarrow$  For each  $t \geq 0$ ,  $u(x,t)$  has time-dependent Fourier coeffs that decay exponentially

$$\hat{u}(t) = \hat{g}_k e^{-(nk)^2 t}$$

This means that  $u(x,t)$  is a very smooth function of  $x$ , infinitely differentiable (in fact it is analytic!)

$\Rightarrow$  As  $t \rightarrow \infty$ , the terms with larger  $k$  = "faster oscillation in  $x$ " decay faster than small  $k$ .

Only the  $k=0$  term doesn't decay,

$$\lim_{t \rightarrow \infty} u(x, t) = \frac{1}{\sqrt{2}} \hat{g}_0 = \frac{1}{2} \underbrace{\int_{-1}^1 g(x) dx}_{\text{mean value of } g(x)}$$

The equilibrium ( $t \rightarrow \infty$ ) temperature distribution is uniform at the mean value of the initial temp. dist.

## Operator Exponential

$$\begin{aligned} \frac{\partial u}{\partial t} &= Au, & \langle Af, g \rangle &= \langle f, Ag \rangle \\ u|_{t=0} &= g, & A & \text{ is a self-adjoint diff. op.} \\ & & & \text{w/ eigenvalues } \lambda_1, \lambda_2, \dots \\ & & & \text{eigenfunctions } e_1, e_2, \dots \end{aligned}$$

$$u(t) = e^{At} g = \sum_{k=1}^{\infty} e^{\lambda_k t} \langle e_k, g \rangle e_k$$

E.g.  $\gamma \neq 1$  or heated ring ( $k = \pm 1, \pm 2, \dots$ )

$$\lambda_k = -\gamma (nk)^2$$

$$e_k = \frac{1}{\sqrt{2}} e^{in_k x}$$

$$u(x, t) = \frac{1}{\sqrt{2}} \sum_{k=-\infty}^{\infty} \hat{g}_k e^{-\gamma (nk)^2 t} e^{in_k x}$$

as diffusivity increases,  
high frequency  
decay is  
faster.

## Homogeneous B.C.'s

$$\frac{\partial u}{\partial t} = \gamma \frac{\partial^2 u}{\partial x^2}$$

$$u(x, 0) = g(x)$$

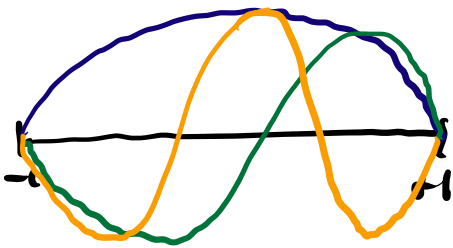
$$\begin{array}{c} u(x, t) \\ \hline -1 \qquad \qquad \qquad 1 \end{array}$$

$$u(-1, t) = u(1, 0) = 0$$

"Fixed temp at endpoints"

Solution is  $u(x, t) = \sum e^{\lambda_k t} \langle e_k, g \rangle e_k$ , so we need to compute eigenvals/eigenfuns of  $A = \frac{d^2 u}{dx^2}$  w/  $u(\pm 1) = 0$

From lecture 6,  $\lambda_k = -\left(\frac{k\pi}{2}\right)^2 \quad k = 1, 2, 3, \dots$



$$e_k(x) = \begin{cases} \cos\left(\frac{k\pi x}{2}\right) & k = 1, 3, 5, \dots \\ \sin\left(\frac{k\pi x}{2}\right) & k = 2, 4, 6, \dots \end{cases}$$

Solution:

$$u(x, t) = \sum_{k=\text{odd}} e^{-\left(\frac{k\pi}{2}\right)^2 t} a_k \cos \frac{k\pi x}{2} + \sum_{k=\text{even}} e^{-\left(\frac{k\pi}{2}\right)^2 t} b_k \sin \frac{k\pi x}{2}$$

$$a_k = \int_{-1}^1 \cos \frac{k\pi x}{2} g(x) dx$$

$$b_k = \int_{-1}^1 \sin \frac{k\pi x}{2} g(x) dx$$

Note that  $u(x,t)$  can be written as a standard Fourier series (complex exp. form) by plugging in  $\cos \frac{k\pi x}{2} = \frac{1}{2}(e^{inkx} + e^{-inkx})$

$$\sin \frac{k\pi x}{2} = \frac{1}{2i}(e^{inkx} - e^{-inkx})$$

What happens as  $t \rightarrow \infty$ ?

Inhomogeneous B.C.s

$$\frac{\partial u}{\partial t} = \gamma \frac{\partial^2 u}{\partial x^2}$$

$$\begin{array}{c} u(x,t) \\ \hline -1 \qquad \qquad \qquad +1 \end{array}$$

$$u(x,0) = g(x)$$

$$u(-1,t) = T_1, \quad u(1,t) = T_2$$

"Fixed temp at endpoints"

Idea: Split  $u(x,t) = u_s(x) + \tilde{u}(x,t)$

↑  
"equilibrium"

$$\frac{\partial u_s}{\partial t} = 0$$

$$u_s(-1) = T_1$$

$$u_s(+1) = T_2$$

↑  
"time-dependent"

$$\frac{\partial \tilde{u}}{\partial t} = \frac{\partial^2 \tilde{u}}{\partial x^2}$$

$$\tilde{u}(\pm 1, t) = 0$$

$$\tilde{u}(x,0) = g - u_s$$

Then,  $u(x,t)$  solves heat equation w/ inhom. B.C.'s.

Egn  $\frac{\partial u}{\partial t} = \frac{\partial u_s}{\partial t} + \frac{\partial \tilde{u}}{\partial t} = \frac{\partial^2 \tilde{u}}{\partial x^2} = \frac{\partial^2 u}{\partial x^2} + \frac{\partial u_s}{\partial x^2} = \frac{\partial^2 u}{\partial x^2}$  ✓

B.C.1

$$u(-1,0) = \underbrace{\tilde{u}(-1,0)}_0 + \underbrace{u_{\#}(-1)}_{T_1} = T_1 \quad \checkmark$$

B.C.2

$$u(+1,0) = \underbrace{\tilde{u}(+1,0)}_0 + \underbrace{u_{\#}(+1)}_{T_2} = T_2 \quad \checkmark$$

I.C.

$$u(x,0) = \tilde{u}(x,0) + u_{\#} = g(x) \quad \checkmark$$

So we need to solve for  $u_{\#}$  and  $\tilde{u}$  to construct full solution to heat eqn w/inhom. BCs.

Soln to  
Homogen.  
BVP w/  
Inhom. BCs

$$\Rightarrow u_{\#}(x) = T_1 + T_2 \left( \frac{1+x}{2} \right)$$

$$\tilde{u}(x,t) = \sum_{k=\text{odd}} e^{-\left(\frac{\pi k}{2}\right)^2 t} a_k \cos \frac{\pi k x}{2} + \sum_{k=\text{even}} e^{-\left(\frac{\pi k}{2}\right)^2 t} b_k \sin \frac{\pi k x}{2}$$

$\uparrow$   
Soln to IBVP w/inhomogeneous BCs

$$\text{Full Soln.} \Rightarrow u(x,t) = u_{\#}(x) + \tilde{u}(x,t)$$

Notice similarity to solution decomposition for BVP (Poisson) w/inhomogeneous B.C.'s.