

# Finite differences in higher dimensions

Laplace and Poisson equations

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18.303 Linear Partial Differential Equations: Analysis and Numerics

# 1d Poisson equation

## Poisson's equation with Dirichlet boundaries

$$\frac{d^2 u(x)}{dx^2} = f(x),$$
$$u(0) = u_0, \quad u(L) = u_{N+1}.$$

Here  $x \in (0, L)$ . We discretize the space  $x$  as before:  $x_k = k\Delta x$  and  $u_k = u(x_k)$ , where  $k = 1, 2, \dots, N$  and  $\Delta x = L/(N+1)$ .

We define the Dirichlet Laplacian

$$D^{(x)} = \frac{1}{\Delta x^2} \begin{pmatrix} -2 & 1 & & & & & \\ 1 & -2 & 1 & & & & \\ & & \ddots & \ddots & \ddots & & \\ & & & & 1 & -2 & 1 \\ & & & & & 1 & -2 \end{pmatrix} \in \mathbb{R}^{N \times N}.$$

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$$\frac{1}{\Delta x^2}(u_0 - 2u_1 + u_2) = f_1$$

while the last equation gives

$$\frac{1}{\Delta x^2}(u_{N-1} - 2u_N + u_{N+1}) = f_N.$$

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We can move the known constants to the RHS of the equation. Now we get

$$D^{(x)}\mathbf{u} = \mathbf{f}_B,$$

where

$$\mathbf{f}_B = \mathbf{f} - \mathbf{b}.$$

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The first entry of  $\mathbf{b} = u_0/\Delta x^2$  and the last one is  $u_{N+1}/\Delta x^2$ . Apart from that, it's zero.

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We can simply set  $\mathbf{f} = 0$  and since  $D^{(x)}$  is invertible, the equation can be solved.

Now we have

$$\Delta u(x, y) = f(x, y),$$

$$u(x, 0) = u^{(b)}(x), u(x, L_y) = u^{(t)}(x),$$

$$u(0, y) = u^{(l)}(y), u(L_x, y) = u^{(r)}(y).$$

Now we have

$$\begin{aligned}\Delta u(x, y) &= f(x, y), \\ u(x, 0) &= u^{(b)}(x), u(x, L_y) = u^{(t)}(x), \\ u(0, y) &= u^{(l)}(y), u(L_x, y) = u^{(r)}(y).\end{aligned}$$

We discretize the space:  $x_i = x(i\Delta x)$ ,  $y_j = y(j\Delta y)$ ,  $u_{i,j} = u(x_i, y_j)$  and so on. Here  $i = 1, 2, \dots, N$  and  $j = 1, 2, 3, \dots, M$ .  $\Delta x = L_x/(N + 1)$  and  $\Delta y = L_y/(M + 1)$ .

We can write the Laplace equation in indices as

$$\frac{1}{\Delta x^2}(u_{i-1,j} - 2u_{i,j} + u_{i+1,j}) + \frac{1}{\Delta y^2}(u_{i,j-1} - 2u_{i,j} + u_{i,j+1}) = f_{i,j}.$$

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We solve them in a similar way. The boundary vector is in indices

$$\frac{1}{\Delta x^2}(\delta_i^1 u_{0,j} + \delta_i^N u_{N+1,j}) + \frac{1}{\Delta y^2}(\delta_j^1 u_{i,0} + \delta_j^M u_{i,N+1}) = \frac{1}{\Delta x^2}(\delta_i^1 u_j^{(l)} + \delta_i^N u_j^{(r)}) + \frac{1}{\Delta y^2}(\delta_j^1 u_i^{(b)} + \delta_j^M u_i^{(t)})$$

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This will give a boundary matrix  $b_{i,j}$  that has the discretized boundary function values on the boundary. Now the source for the equation  $f^{(B)} = f - b$ .

We can regroup the LHS of the equation

$$\frac{1}{\Delta x^2}(u_{i-1,j} - 2u_{i,j} + u_{i+1,j}) - \frac{2u_{i,j}}{\Delta y^2} + \frac{1}{\Delta y^2}(u_{i,j-1} + u_{i,j+1}) = f_{i,j}^{(B)}.$$

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Let us denote the  $j$ th column of  $u$  as  $\mathbf{u}_j$  and the same for  $f^{(B)}$ . Now we have

$$D^{(x)}\mathbf{u}_j - \frac{2}{\Delta y^2}\mathbf{u}_j + \frac{1}{\Delta y^2}(\mathbf{u}_{j-1} + \mathbf{u}_{j+1}) = \mathbf{f}_j^{(B)}.$$

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Let us define  $l_y = 1/\Delta y^2$ . Now we have

$$l_y\mathbf{u}_{j-1} + \underbrace{\left(D^{(x)} - 2l_y\right)}_{=B}\mathbf{u}_j + l_y\mathbf{u}_{j+1} = \mathbf{f}_j^{(B)}.$$

$$l_y \mathbf{u}_{j-1} + B \mathbf{u}_j + l_y \mathbf{u}_{j+1} = \mathbf{f}_j^{(B)}.$$

Doesn't this look familiar?

$$I_y \mathbf{u}_{j-1} + B \mathbf{u}_j + I_y \mathbf{u}_{j+1} = \mathbf{f}_j^{(B)}.$$

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We can make a matrix of matrices  $A$  called a block matrix. The equation will look like this

$$AU = F^{(B)},$$

where

$$A = \begin{pmatrix} B & I_y & & & \\ I_y & B & I_y & & \\ & \ddots & \ddots & \ddots & \\ & & I_y & B & I_y \\ & & & I_y & B \end{pmatrix}$$

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$$AU = F^{(B)}.$$

How should we interpret  $U$  and  $F^{(B)}$ ? If we write  $A$  by just filling in the matrices inside it, it will give  $A$  the dimensions  $NM \times NM$ . We can then express  $U$  as

$$U = \begin{pmatrix} u_{1,1} \\ u_{2,1} \\ \vdots \\ u_{N,1} \\ u_{1,2} \\ \vdots \\ u_{N,M} \end{pmatrix}.$$

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The vector  $F^{(B)}$  is flattened in the same way.

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- Finite difference methods are easy to implement for simple geometries.
- The discrete linear equations can be solved efficiently using existing packages.
- Can be extended to non-uniform meshes but that requires quite a bit of effort.