

Fourier Spectral Methods (pt. 2)

$$-u_{xx} = f$$

periodic B.C.s

on $[0, 2\pi]_{\text{per}}$

$$\sim/\text{gauge } g = \frac{1}{2\pi} \int_0^{2\pi} u(x) dx$$

ONB

$$\phi_n(x) = \frac{1}{\sqrt{2\pi}} e^{inx}$$

for $n = 0, \pm 1, \pm 2, \dots$

Soln

$$u(x) \approx \sum_{|n| \leq N} \hat{u}_n \phi_n(x)$$

RHS

$$f(x) \approx \sum_{|n| \leq N} \hat{f}_n \phi_n(x)$$

Algorithm

① Compute $\hat{f}_n = F \frac{f}{\tau}$ values
Fourier coeffs \uparrow \downarrow FFT
on grid
 $x_n = \frac{hn}{2\pi}$

② Solve $\hat{u}_n = \hat{f}_n / \lambda_n$ eigenvalues of $-\partial_x^2$ on $[0, 2\pi]_{\text{per}}$
 and $\hat{u}_0 = g$ (gauge) $n = \pm 1, \pm 2, \dots, \pm N$

③ Compute $\underline{u} = F^{-1} \hat{\underline{u}}$
IFFT

Note: identical to fast FD solver except we use eigenvalues of $-\partial_x^2$ (n^2) instead of $D = \begin{bmatrix} 2 & -1 & & \\ -1 & 2 & & \\ & & \ddots & \\ & & & 2 & -1 \\ & & & -1 & 2 \end{bmatrix}$.

^{2nd order}
 \Rightarrow fastest FD solver is only 2nd order accurate
 because

$$\lambda_m^{\text{FD}} = \left(\frac{2\pi}{N}\right)^2 \left[2 - 2\cos\left(\frac{2\pi m}{N}\right) \right] = m^2 \left(1 + \mathcal{O}\left(\left(\frac{m}{N}\right)^2\right) \right)$$

\uparrow
 Taylor expansion

Only the low lying modes are accurate ($m \ll N$),
 leading to the low-order accuracy of the FD scheme.

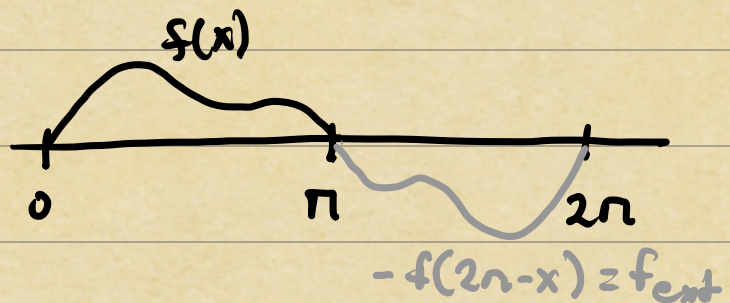
Non-periodic problems

Odd periodic extension

$$-u_{xx} = f$$

$$u(0) = u(\pi) = 0$$

(Dirichlet B.C.s)



We can solve using a sine series, BUT
 convergence rates will depend on smoothness
 of the odd-periodic extension of $f(x)$.

Algorithm

ONB

$$\phi_n(x) = \sqrt{\frac{2}{\pi}} \sin(nx)$$

soln

$$u(x) = \sum_{n=1}^{\infty} \hat{u}_n \phi_n(x)$$

RHS

$$f(x) = \sum_{n=1}^{\infty} \hat{f}_n \phi_n(x)$$

$$\textcircled{1} \hat{f} = \text{DST1}(f)$$

$$\textcircled{2} \hat{u}_n = \hat{f}_n / n^2$$

$$\textcircled{3} u = \text{IDST1}(\hat{u})$$

Note that for, e.g., super-algebraic convergence we need for $k=0,1,2,\dots$

smooth interior $\Rightarrow f^{(k)}(x)$ continuous on $[0, 2\pi]$

smooth extension $\Rightarrow f^{(k)}(0) = f^{(k)}(2\pi)$

$$f_{\text{ext}}(x) = -f(2\pi - x) \Rightarrow f_{\text{ext}}^{(n)}(x) = (-1)^{n+1} f^{(n)}(2\pi - x)$$

for continuity
at $x=\pi$

$$\lim_{x \rightarrow \pi^-} f^{(n)}(x) = \lim_{x \rightarrow \pi^+} (-1)^{n+1} f^{(n)}(2\pi - x)$$

$$\text{for odd } n \Rightarrow f^{(n)}(\pi) = f^{(n)}(\pi)$$

$$\text{for even } n \Rightarrow f^{(n)}(\pi) = -f^{(n)}(\pi)$$

The odd n condition is always satisfied but the even n condition can only be true when $f^{(n)}(\pi) = 0$. Similar calculation at $x=0$ leads to

$$\text{Super-algebraic convergence rate} \Leftrightarrow f^{(n)} \Big|_{x=0, \pi} = 0 \text{ for all even } n$$

Similar scheme based on even extensions and DC71 for Neumann boundary conditions.

Multi-dimensional Fourier Spectral Methods

Fourier, Sine, and Cosine bases diagonalize const coefficient differential operators. Can extend to multi-dimensions using Kronecker products analogous to multi-dim finite difference schemes.

Example: 2D Poisson ^{on $[0, \pi]^2$} w/homogeneous Dirichlet BCs.

ONB $\phi_{nm}(x, y) = \frac{2}{\pi} \sin(nx) \sin(my)$

w/inner product $\langle f, g \rangle = \int_0^\pi \int_0^\pi f(x, y) g(x, y) dx dy$

Soln $u(x, y) = \sum_{n=1}^N \sum_{m=1}^N \hat{u}_{nm} \phi_{nm}(x, y)$

RHS $f(x, y) = \frac{1}{2\pi} \sum_{n=1}^N \sum_{m=1}^N \hat{f}_{nm} \phi_{nm}(x, y)$

Step 1: Compute $\hat{f}_{nm} = \langle \phi_{nm}, f \rangle$

Iterated integral = Sine Transform in x , then y

$$\begin{aligned}
 \hat{f}_{nm} &= \langle \phi_{nm}, f \rangle = \frac{2}{n} \int_0^n \int_0^n \sin(nx) \sin(my) f(x,y) dx dy \\
 &= \sqrt{\frac{2}{n}} \int_0^n \sin(my) \underbrace{\left(\sqrt{\frac{2}{n}} \int_0^n \sin(nx) f(x,y) dx \right)}_{\substack{\text{1D sine transform} \\ \text{in } x \Rightarrow \tilde{f}_n(y)}} dy \\
 &= \sqrt{\frac{2}{n}} \int_0^n \sin(my) \tilde{f}_n(y) dy \quad \text{1D transform in } y
 \end{aligned}$$

Approximating the inner and outer integrals using a trapezoidal rule on an equi-spaced grid leads to a Discrete Sine Transform (Type-1) along both array axes of sample matrices

$$F = \begin{bmatrix} f(x_1, y_1) & \dots & f(x_1, y_N) \\ \vdots & & \vdots \\ f(x_N, y_1) & \dots & f(x_N, y_N) \end{bmatrix}$$

of values of $f(x,y)$ on 2D equispaced grid (x_n, y_m)

(N interior grid points)

$$y_1 = x_1 = \frac{n}{N+1}, \quad y_2 = x_2 = \frac{2n}{N+1}, \quad \dots, \quad y_N = x_N = \frac{2nN}{N+1}$$

Step 2: Solve for \hat{u}_{nm} (diagonal system)

eigenvals of $-\Delta u/208k_1$

$$\langle \phi_{nm}, -\Delta u \rangle = (n^2 + m^2) \hat{u}_{nm} \quad \text{and} \quad \langle \phi_{nm}, f \rangle = \hat{f}_{nm}$$

$$\Rightarrow \hat{u}_{nm} = \hat{f}_{nm} / (n^2 + m^2) \quad \text{for } n, m = 1, \dots, N$$

Step 3: Compute $u(x, y)$ on grid (equispaced)

Evaluating the 2D sine series

$$u(x, y) = \sum_{n=1}^N \sum_{m=1}^N \hat{u}_{nm} \phi_{nm}(x, y)$$

on the equispaced grid (x_n, y_n) above leads to an inverse Discrete Sine Transform (Type-1) along both array axis of the $N \times N$ matrix

$$\hat{u} = \begin{bmatrix} \hat{u}_{11} & \dots & \hat{u}_{1N} \\ \vdots & & \vdots \\ \hat{u}_{N1} & \dots & \hat{u}_{NN} \end{bmatrix}$$