

18.369: Mathematical Methods in Nanophotonics

overview lecture slides
(don't get used to it: most lectures are blackboard)

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Maxwell's Equations

$$\nabla \cdot \mathbf{B} = 0$$

Gauss: $\nabla \cdot \mathbf{D} = \rho$ constitutive relations:

Ampere: $\nabla \times \mathbf{H} = \frac{\partial \mathbf{D}}{\partial t} + \mathbf{J}$ $\epsilon_0 \mathbf{E} = \mathbf{D} - \mathbf{P}$

Faraday: $\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$ $\mathbf{H} = \mathbf{B}/\mu_0 - \mathbf{M}$

electromagnetic fields:

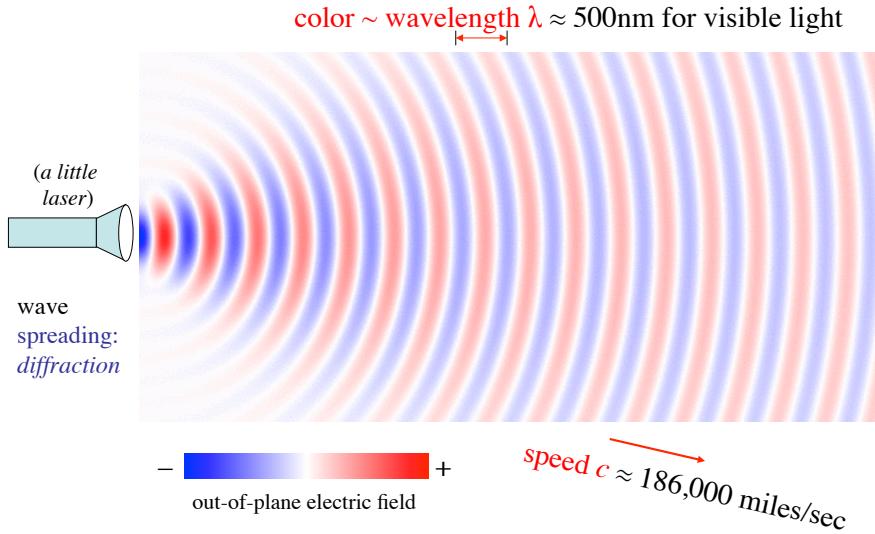
\mathbf{E} = electric field
 \mathbf{D} = displacement field
 \mathbf{H} = magnetic field / induction
 \mathbf{B} = magnetic field / flux density

constants: ϵ_0, μ_0 = vacuum permittivity/permeability
 c = vacuum speed of light = $(\epsilon_0 \mu_0)^{-1/2}$

sources: \mathbf{J} = current density
 ρ = charge density

material response to fields:
 \mathbf{P} = polarization density
 \mathbf{M} = magnetization density

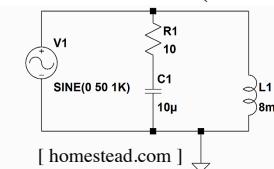
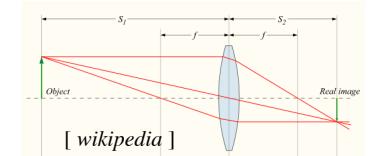
Light is a Wave (and a quantum particle)



When can we solve this mess?

- Very small wavelengths: **ray optics**

- Very large wavelengths:
quasistatics (8.02)
& **lumped circuit models** (6.002)



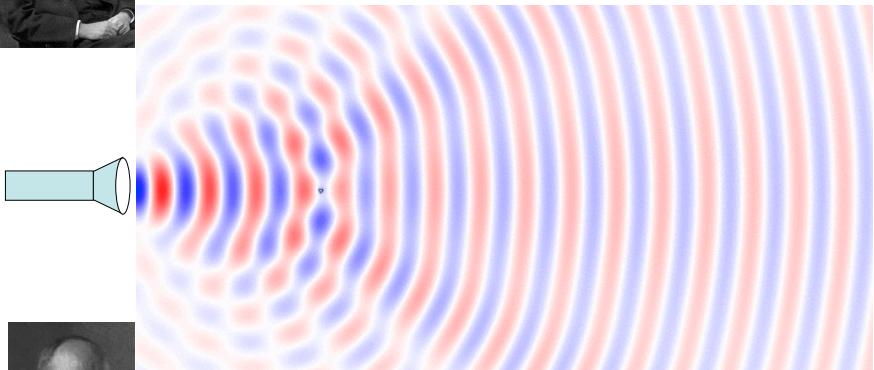
- Wavelengths comparable to geometry?
 - handful of cases can be ~solved analytically:
planes, spheres, cylinders, empty space (8.07, 8.311)
 - everything else just a mess for computer...?



small particles:
Lord Rayleigh (1871)
why the sky is blue

Waves Can Scatter

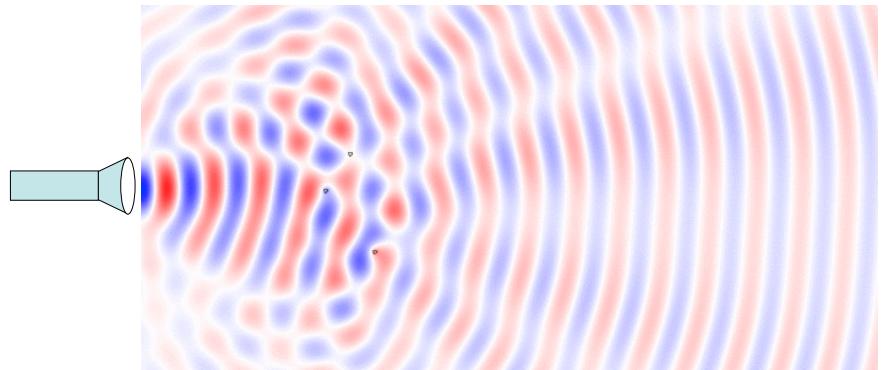
here: a little circular speck of silicon



checkerboard pattern: interference of waves
traveling in different directions
scattering by spheres:
solved by Gustave Mie (1908)

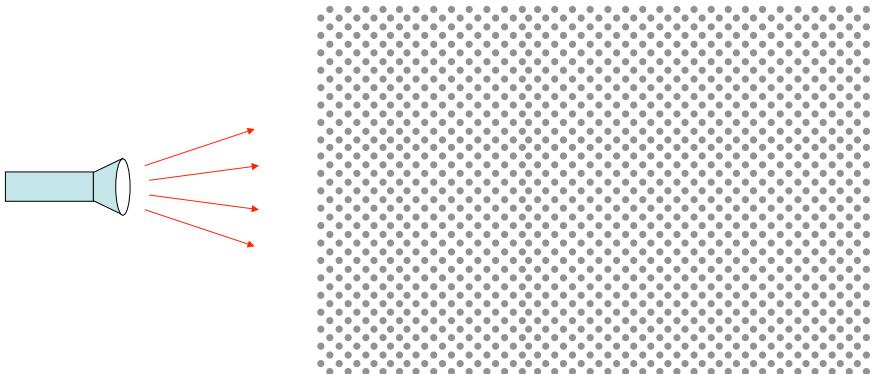
Multiple Scattering is Just Messier?

here: scattering off three specks of silicon



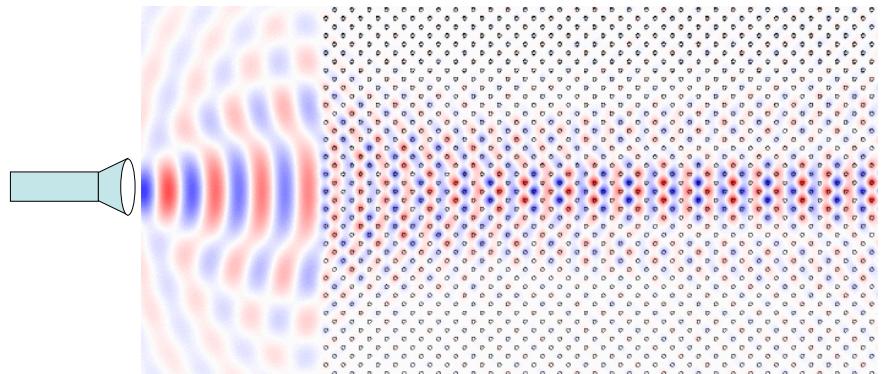
can be solved on a computer, but not terribly interesting...

An even bigger mess? zillions of scatterers



Blech, light will just scatter like crazy
and go all over the place ... how boring!

Not so messy, not so boring...



the light seems to form several *coherent beams*
that propagate *without scattering*
... and *almost without diffraction* (*supercollimation*)

...the magic of symmetry...



[Emmy Noether, 1915]

Noether's theorem:
symmetry = conservation laws

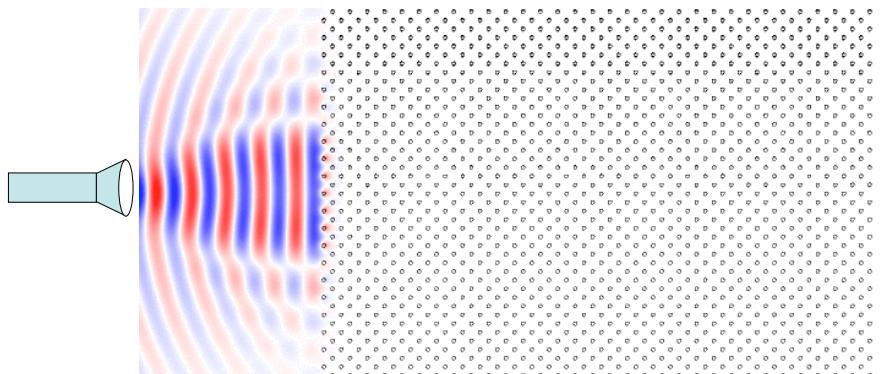
In this case, periodicity
= conserved "momentum"
= wave solutions without scattering
[Bloch waves]



Felix Bloch
(1928)

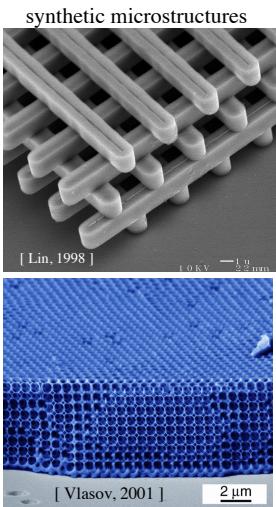
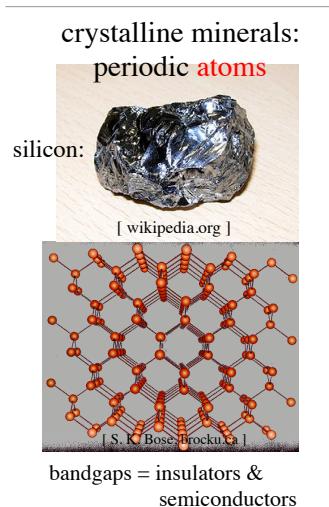
Mathematically, use *structure* of the equations, not explicit solution:
linear algebra, group theory, functional analysis, ...

A slight change? Shrink λ by 20%
an "optical insulator" (*photonic bandgap*)



light **cannot** penetrate the structure at this wavelength!
all of the scattering destructively interferes

Photonic Crystals: periodic structures for light



Lots of math to
understand this...

...linear algebra:

$$\nabla \times \frac{1}{\epsilon} \nabla \times \vec{H} = \left(\frac{\omega}{c} \right)^2 \vec{H}$$

(eigenproblem)

...symmetry \Rightarrow group
representation theory

...computational
methods...

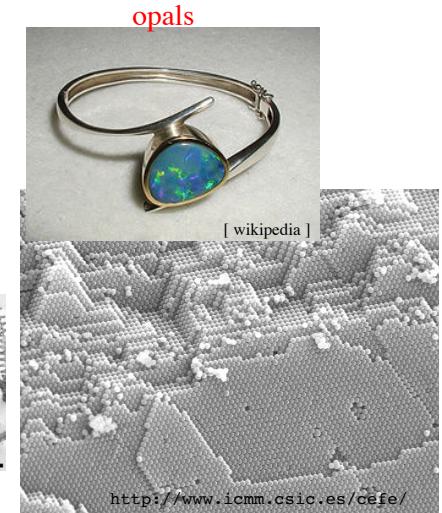
...many open
questions...

Structural Color in Nature

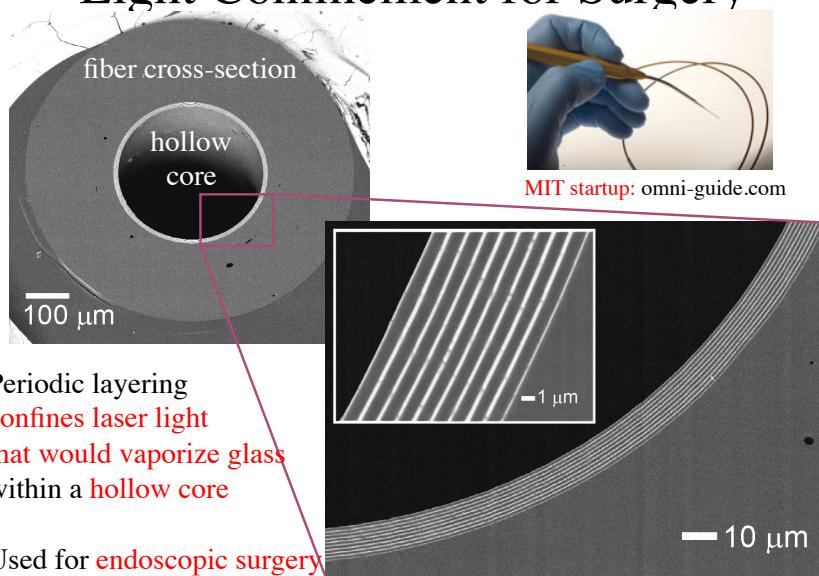
bandgap = wavelength-selective mirror = bright iridescent colors



wing scale:



Light Confinement for Surgery



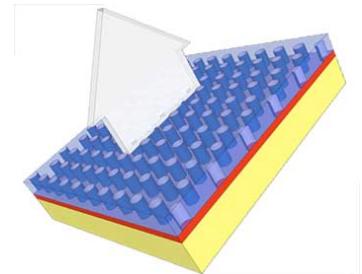
Molding Diffraction for Lighting

[another MIT startup (by a colleague): Luminus.com]



ultra-bright/efficient LEDs

periodic pattern gathers & redirects it in one direction



new projection TVs,
pocket projectors,
lighting applications,

...

Back to Maxwell, with some simplifications

- *source-free* equations (propagation of light, not creation): $\mathbf{J} = 0, \rho = 0$
- *Linear, dispersionless* (instantaneous response) materials:

$$\mathbf{P} = \epsilon_0 \chi_e \mathbf{E} \quad \Rightarrow \quad \mathbf{D} = \epsilon_0 (1 + \chi_e) \mathbf{E} = \epsilon_0 \epsilon_r \mathbf{E}$$

$$\mathbf{M} = \chi_m \mathbf{H} \quad \Rightarrow \quad \mathbf{B} = \mu_0 (1 + \chi_m) \mathbf{H} = \mu_0 \mu_r \mathbf{H}$$

(nonlinearities very weak
in EM ... we'll treat later)
(dispersion can be negligible
in narrow enough bandwidth)

where $\epsilon_r = 1 + \chi_e$ = relative permittivity
(drop r subscript) or dielectric constant
 $\mu_r = 1 + \chi_m$ = relative permeability
 $\epsilon\mu = (\text{refractive index})^2$

- *Isotropic* materials: $\epsilon, \mu = \text{scalars}$ (not matrices)
- *Non-magnetic* materials: $\mu = 1$ (true at optical/infrared)
- *Lossless, transparent materials*: ϵ real, > 0 (< 0 for metals...bad at infrared)

Simplified Maxwell

$$\nabla \cdot \mathbf{H} = 0 \quad \nabla \cdot \epsilon \mathbf{E} = 0$$

$$\nabla \times \mathbf{H} = \epsilon_0 \epsilon(\mathbf{x}) \frac{\partial \mathbf{E}}{\partial t} \quad \nabla \times \mathbf{E} = -\mu_0 \mu \frac{\partial \mathbf{H}}{\partial t} = -\mu_0 \frac{\partial \mathbf{H}}{\partial t}$$

- *Linear, time-invariant* system:
⇒ look for *sinusoidal solutions* $\mathbf{E}(\mathbf{x},t) = \mathbf{E}(\mathbf{x})e^{-i\omega t}$, $\mathbf{H}(\mathbf{x},t) = \mathbf{H}(\mathbf{x})e^{-i\omega t}$

(i.e. Fourier transform)

$$\nabla \times \mathbf{H} = -i\omega \epsilon_0 \epsilon(\mathbf{x}) \mathbf{E} \quad \nabla \times \mathbf{E} = i\omega \mu_0 \mathbf{H}$$

[note: real materials have
dispersion: ϵ depends on ω
= non-instantaneous response]

... these, we can work with