

University of California at San Diego – Department of Physics – Prof. John McGreevy
Physics 210B Non-equilibrium Fall 2025
Assignment 5 – Solutions

Due 11:59pm Monday, November 3, 2025

1. **The adjoint and the FP operator.** Write the Fokker-Planck equation $\partial_t P + \vec{\nabla} \cdot \vec{J} = 0$ for an overdamped particle in a potential $U(x)$ as

$$\partial_t P(x, t) = -\hat{L}P(x, t). \quad (1)$$

Assume the fluctuation-dissipation relation $D = \mu T$, so that

$$P_{\text{eq}}(x) = e^{-U(x)/T} / Z \quad (2)$$

is a stationary solution of the FP equation.

- (a) Show that the linear operator

$$\hat{H} = \frac{1}{\sqrt{P_{\text{eq}}}} \hat{L} \sqrt{P_{\text{eq}}} \quad (3)$$

is hermitian with respect to the inner product $\langle f|g \rangle \equiv \int_{-\infty}^{\infty} dx f(x)g(x)$ (the usual L^2 inner product on real functions on \mathbb{R}). (\hat{H} is the linear operator in the FP equation for $P/\sqrt{P_{\text{eq}}}$.)

For one thing, this shows that the eigenvalues of \hat{L} are real, so that, in this problem, there are no oscillations in the approach to equilibrium.

$$\hat{H} = \frac{1}{\sqrt{P_{\text{eq}}}} \hat{L} \sqrt{P_{\text{eq}}} \quad (4)$$

let $\psi \equiv \frac{P}{\sqrt{P_{\text{eq}}}}$

$$\partial_t \psi = -\hat{H} \psi \quad (5)$$

$$\partial_x \sqrt{P_{\text{eq}}} = -\frac{1}{2T} \sqrt{P_{\text{eq}}} U' \quad (6)$$

$$\partial_x^2 \sqrt{P_{\text{eq}}} = \left(\frac{1}{4T^2} U'^2 - \frac{1}{2T} U'' \right) \sqrt{P_{\text{eq}}} \quad (7)$$

Let's apply to a test function $\phi(x)$:

$$\hat{L}(\sqrt{P_{\text{eq}}} \psi) = \sqrt{P_{\text{eq}}} \left[-\mu T \partial_x^2 \phi + \mu \left(\frac{U'^2}{4T} - \frac{U''}{2} \right) \phi \right] \quad (8)$$

So,

$$\hat{H} = D \left[-\partial_x^2 + \frac{U'^2}{4T^2} - \frac{U''}{2T} \right] = D \left[-\partial_x^2 + W^2 - W' \right] \quad (9)$$

This can be factorized into

$$\hat{H} = D (-\partial_x + W) (\partial_x + W) \quad (10)$$

where $W = \frac{U'}{2T}$. Now to show it, we do

$$\langle f | \hat{H} g \rangle = D \int dx f (-\partial_x + W) (\partial_x + W) g = D \int dx f (-g'' + Vg) \quad (11)$$

$$\langle f | \hat{H} g \rangle = D \int dx (f'g' + fVg) = D \int dx (-f'' + Vf) g = \langle \hat{H} f | g \rangle \quad (12)$$

- (b) [Bonus] Show that the adjoint of \hat{L} with respect to this inner product is the operator appearing in the *backwards* FP equation describing the evolution of the *initial* time

$$\partial_{t'} P(x, t | x', t') = -\hat{L}^\dagger P(x, t | x' t'). \quad (13)$$

2. **Stochastic topology.** Suppose you tie your dog to a flagpole (by an infinitely-extendible, massless leash that can pass through itself) in the center of a big round field and go run an errand¹. Your dog performs Brownian motion with diffusion constant D . We would like to find (analytically!) the distribution for the *winding angle*

$$\theta(t) = \int_0^t dt \dot{\theta} = \int_0^t dt \partial_t \arctan(y/x) = \int_0^t dt \frac{x\dot{y} - y\dot{x}}{x^2 + y^2} \quad (14)$$

of your dog's motion around the flagpole.

There are some choices to make in modeling this situation (which actually make a difference for the late time behavior):

- The flagpole is pointlike or has a finite radius a .
- The field is bounded by $r < R$ (and your dog can't leave) or goes on forever.

Pick an option for both items. The case $a = 0$ and an unbounded region is the most interesting. As a bonus problem, compare multiple options for the boundary conditions.

¹I've never had a dog, so I don't know if this is an OK thing to do, but let's say that the dog doesn't mind.

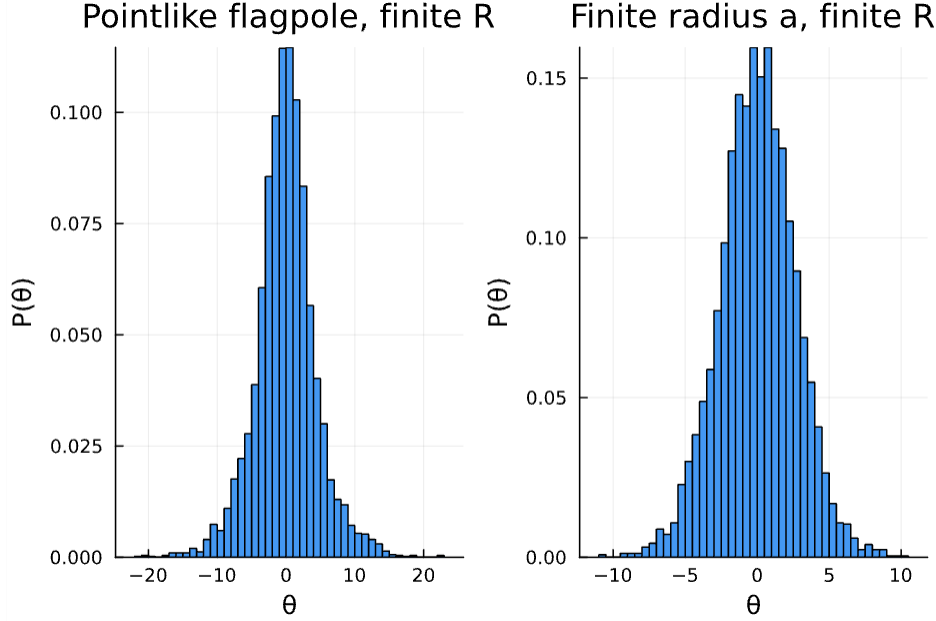


Figure 1: Winding angle distribution at $t = 1$ for pointlike flagpole with finite R (left) and finite radius a with finite R (right).

- (a) Simulate the problem with the various choices mentioned above and make histograms of the winding angle distribution at late times. Warning: the case where the region is unbounded requires some care.

Please see the figures 1 and 2. In the case with a pointlike flagpole, the dog is able to jump through and make large angles.

- (b) Here is a much easier but much less interesting version of the problem that may help get you started on the next part. Suppose we have a particle diffusing on a circle, S^1 (or a torus $\equiv (S^1)^d$). What is the probability distribution for the winding angle $\theta(t)$ as a function of time? By the winding angle, I mean we should keep track of the total change in angle from its initial position, not just mod 2π :

$$\theta \equiv \int_0^t ds \dot{\theta} \quad (15)$$

(or the discrete version of this formula).

Angular diffusion with constant D

$$d\psi(t) = \sqrt{2D} dW_t, \theta(0) = 0 \quad (16)$$

Now this is equivalent to a 1D Brownian motion with diffusion constant D .

$$\partial_t P = D \frac{\partial^2 P}{\partial \theta^2} \quad (17)$$

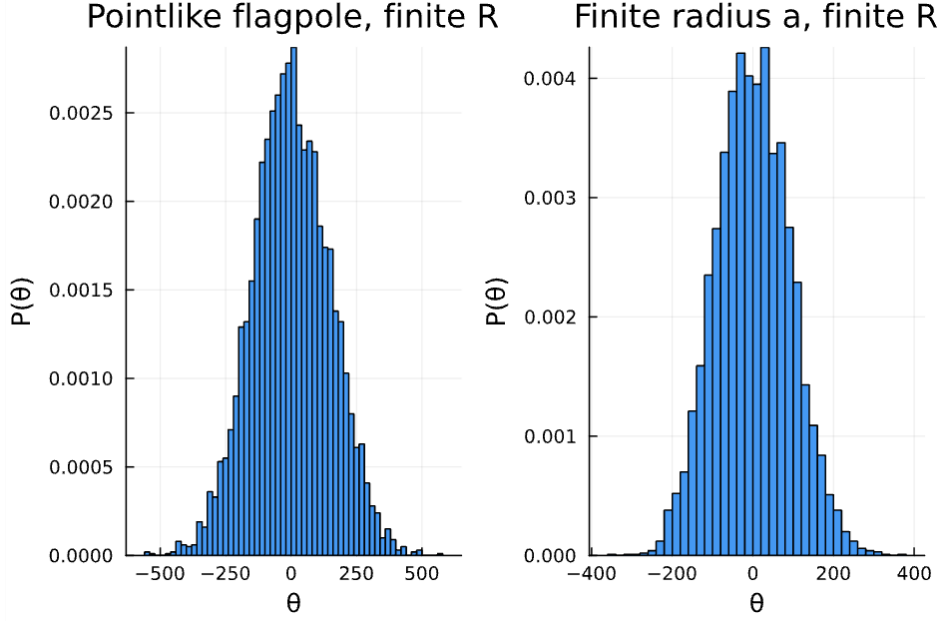


Figure 2: Winding angle distribution at $t = 1200$ for pointlike flagpole with finite R (left) and finite radius a with finite R (right).

$$P(\theta, 0) = \delta(\theta) \quad (18)$$

The solution is

$$P(\theta, t) = \frac{1}{\sqrt{4\pi Dt}} e^{-\frac{\theta^2}{4Dt}} \quad (19)$$

Thus, the distribution of the winding angle is gaussian.

- (c) Returning to the 2d story of the dog, here is the real problem. Find analytically the probability distribution $P(\theta, t)$ for the winding angle at late times. Suppose the initial distribution is localized to $r = r_0$.

Here are some hints for the case with unbounded domain.

- The following integral is useful:

$$\int_0^\infty dk k J_{|m|}(kr_0) J_{|m|}(kr) e^{-k^2 t} = \frac{1}{2t} e^{-\frac{r^2 + r_0^2}{4t}} I_{|m|}\left(\frac{rr_0}{2t}\right). \quad (20)$$

- if you try to do the r integral under the integral over eigenvalues, you'll get nonsense.

Here are some hints for the cases with a bounded domain $R < \infty$:

- Let $j_{m,n}$ be the location of the n th zero of the ordinary Bessel function $J_m(x)$. For small $m \neq 0$, it has the Taylor expansion: $(j_{|m|,1})^2 \approx cm^2 + \mathcal{O}(m^4)$ for some constant c .

- Similarly, let $j'_{m,n}$ be the location of the n th zero of $\partial_x J_m(x)$. For small $m \neq 0$, it has the Taylor expansion: $(j'_{m,1})^2 \approx c'm^2 + \mathcal{O}(m^4)$ for some constant c' .
- (d) [Bonus] Here is another wrinkle to consider: suppose that $r = R$ is an absorbing barrier (say that if $r > R$, the leash breaks and your dog runs home). What's the distribution of the winding angle *given* that the leash didn't break?
- (e) [Bonus] Suppose that there is a wind spiraling around the flagpole counter-clockwise, which applies a drift force on your dog of the form $\vec{F} = r\Omega(r)\hat{\theta}$ with $\Omega(r) = \frac{\beta}{r^2}$. Show that the FP equation is instead

$$\partial_t P = -\Omega(r)\partial_\theta P + D\vec{\nabla}^2 P. \quad (21)$$

Now what is the late-time distribution of θ ?