

HW 1.3

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PHYS 513

a.) Show that E from a finite line of length $2L$ with charge density λ is the same as Gauss's when taking the Taylor approx of the E field.

Gauss's Law

$r \gg L$

$$\oint \vec{E} \cdot d\vec{a} = \frac{Q_{enc}}{\epsilon_0}$$

$$E \cdot 4\pi r^2 = \frac{2\lambda L}{\epsilon_0}$$

$$E = \frac{2\lambda L}{4\pi\epsilon_0 r^2} \hat{r} \rightarrow \frac{q}{4\pi\epsilon_0 r^2} \hat{r}$$



Starting with E field from line charge

$$\vec{E}_z(z) = \frac{1}{4\pi\epsilon_0} \frac{2\lambda L}{z \sqrt{z^2 + L^2}} \hat{z} \quad \leftarrow \text{from Griffith's example 2.2}$$

$$= \frac{2\lambda}{4\pi\epsilon_0} \frac{L}{z \cdot z \sqrt{1 + (\frac{L^2}{z^2})}} \hat{z}$$

$$= \frac{2\lambda}{4\pi\epsilon_0} \frac{L/z^2}{\sqrt{1 + (L^2/z^2)}} \cdot \frac{L}{L} \hat{z}$$

$$\vec{E}_z(z) = \frac{2\lambda}{4\pi\epsilon_0 L} \frac{L^2/z^2}{\sqrt{1 + (L^2/z^2)}} \cdot \hat{z}$$

define $f(x) = \frac{x}{\sqrt{1+x}}$ so $x = \frac{L^2}{z^2}$ & $\frac{L^2}{z^2} \ll 1$

$$f(x) = f(0) + x \cdot f'(0) + O(x^2)$$

$$f(0) = \frac{0}{\sqrt{1+0}} = 0$$

$$f'(x) = \frac{(x+2)}{2(1+x)^{3/2}} \rightarrow f'(0) = \frac{0+2}{2(1+0)^{3/2}} = 1$$

\leftarrow wolfram alpha

$$f(x) = x + O(x^2)$$

$$\vec{E}_z(z) = \frac{2\lambda}{4\pi\epsilon_0 L} f(x) \hat{z} \rightarrow \frac{2\lambda}{4\pi\epsilon_0 L} \cdot x \hat{z}$$

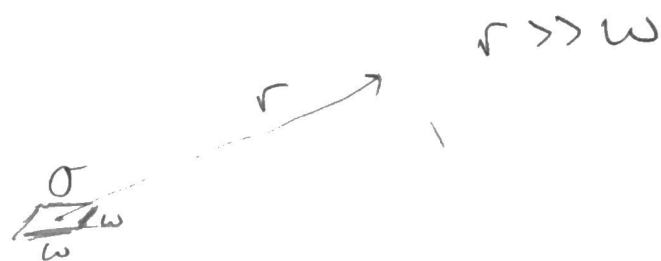
$$\vec{E}_z(z) = \frac{2\lambda}{4\pi\epsilon_0 L} \frac{L^2}{z^2} \rightarrow \boxed{\vec{E}_z(z) = \frac{2\lambda L}{4\pi\epsilon_0 z} \hat{z}}$$

(2)

b.) Show that E from sheet-charge of sides w and w with charge density σ is the same as Gauss's law when taking the Taylor approx of the E field.

Gauss's Law

$$\oint \vec{E} \cdot d\vec{a} = \frac{Q_{\text{enc}}}{\epsilon_0}$$



$$E \cdot 4\pi r^2 = \frac{\sigma w^2}{\epsilon_0}$$

$$\vec{E} = \frac{\sigma w^2}{4\pi \epsilon_0 r^2} \hat{r} \rightarrow \frac{q}{4\pi \epsilon_0 r^2}$$

Start $\vec{E}_z(z) = \frac{\sigma_0}{\pi \epsilon_0} \tan^{-1} \left[\frac{W^2}{4z} \frac{1}{\sqrt{z^2 + \frac{W^2}{2}}} \right] \hat{z}$

$$= \frac{\sigma}{\pi \epsilon_0} \tan^{-1} \left[\frac{W^2}{4z} \frac{1}{z \sqrt{1 + \frac{W^2}{2z^2}}} \right] \hat{z}$$

$$\vec{E}_z(z) = \frac{\sigma}{\pi \epsilon_0} \tan^{-1} \left[\frac{1}{2} \frac{W^2}{z^2} \frac{1}{\sqrt{1 + \frac{W^2}{2z^2}}} \right] \hat{z}$$

define $f(x) = \tan^{-1} \left[\frac{1}{2} \frac{x}{\sqrt{1+x}} \right]$

so $x = \frac{W^2}{2z^2}$ & $\frac{W^2}{2z^2} \ll 1$

$$f(x) = f(0) + x \cdot f'(0) + O(x^2)$$

$$f(0) = \tan^{-1} \left[\frac{1}{2} \frac{0}{0} \right] = 0$$

$$f'(x) = \frac{1}{(x+2)\sqrt{1+x}} \quad f'(0) = \frac{1}{(0+2)\sqrt{1+0}} = \frac{1}{2}$$

← wolfram alpha

$$f(x) = 0 + x \frac{1}{2} + O(x^2)$$

$$\vec{E}_z(z) = \frac{\sigma}{\pi \epsilon_0} f(x) \hat{z} \rightarrow \frac{\sigma}{\pi \epsilon_0} \cdot \frac{1}{2} x \hat{z}$$

$$\vec{E}_z(z) = \frac{\sigma}{\pi \epsilon_0} \frac{1}{2} \frac{W^2}{2z^2} \hat{z}$$

$$\vec{E}_z(z) = \frac{W^2 \sigma}{4 \pi \epsilon_0 z^2} \hat{z}$$

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