Notes and derivations not given explicitly in referenced documents.

Derivatives of regular Laplace expansions, assuming normalized Legendre functions and unit normalization constants ( $\alpha$  and  $\beta$  in G&D). Identities for Legendre functions come from CS-TR-4264.

$$\partial_z = \mu \frac{\partial}{\partial r} + \frac{1 - \mu^2}{r} \frac{\partial}{\partial \mu},\tag{1}$$

$$\partial_{xy} = \frac{\partial}{\partial x} + i \frac{\partial}{\partial y} \tag{2}$$

$$= \frac{e^{i\phi}}{r(1-\mu^2)^{1/2}} \left[ (1-\mu^2) \left( r \frac{\partial}{\partial r} - \mu \frac{\partial}{\partial \mu} \right) + i \frac{\partial}{\partial \phi} \right], \tag{3}$$

$$\overline{\partial_{xy}} = \frac{\partial}{\partial x} - i \frac{\partial}{\partial y} \tag{4}$$

$$= \frac{e^{-i\phi}}{r(1-\mu^2)^{1/2}} \left[ (1-\mu^2) \left( r \frac{\partial}{\partial r} - \mu \frac{\partial}{\partial \mu} \right) - i \frac{\partial}{\partial \phi} \right], \tag{5}$$

$$\partial_x = (\partial_{xy} + \overline{\partial_{xy}})/2,\tag{6}$$

$$\partial_y = -\mathrm{i}(\partial_{xy} - \overline{\partial_{xy}})/2. \tag{7}$$

$$R_n^m = r^n N_n^m P_n^{|m|} e^{im\phi}, \tag{8}$$

$$\frac{\partial R_n^m}{\partial z} = \left[ \mu n r^{n-1} P_n^{|m|} + r^{n-1} (1 - \mu^2) \frac{\partial P_n^{|m|}}{\partial \mu} \right] e^{im\phi} N_n^m, \tag{9}$$

$$= r^{n-1} N_n^m P_{n-1}^{|m|} e^{im\phi}, (10)$$

$$=\frac{N_n^m}{N_{n-1}^m}R_{n-1}^m, (11)$$

$$= \left[ \frac{2n+1}{2n-1} (n^2 - m^2) \right]^{1/2} R_{n-1}^m. \tag{12}$$

$$S_n^m = r^{-n-1} N_n^m P_n^{|m|} e^{im\phi}, (13)$$

$$\partial_z S_n^m = -\left[\mu(n+1)r^{-n-2}P_n^{|m|} + r^{-n-2}(1-\mu^2)\frac{\partial P_n^{|m|}}{\partial \mu}\right] e^{im\phi} N_n^m, \quad (14)$$

$$= -r^{-n-2}N_n^m(n-|m|+1)P_{n+1}^{|m|}e^{im\phi}, (15)$$

$$= -(n - |m| + 1) \frac{N_n^m}{N_{n+1}^m} S_{n+1}^m, \tag{16}$$

$$= -\left\{\frac{2n+1}{2n+3}\left[(n+1)^2 - m^2\right]\right\}^{1/2} S_{n+1}^m,\tag{17}$$

$$\partial_{xy} S_n^m = \left\{ \left( 1 - \mu^2 \right)^{1/2} \left[ -(n+1) P_n^{|m|}(\mu) - \mu \frac{\partial P_n^{|m|}}{\partial \mu} \right] - \right\}$$

$$\frac{m}{(1-\mu^2)^{1/2}} P_n^{|m|}(\mu) \right\} N_n^m e^{i(m+1)\phi} r^{-n-2}, \tag{18}$$

$$= \left[ \frac{2n+1}{2n+3} (n+m+1)(n+m+2) \right]^{1/2} S_{n+1}^{m+1}, \tag{19}$$

$$\overline{\partial_{xy}}S_n^m = \left\{ \left(1 - \mu^2\right)^{1/2} \left[ -(n+1)P_n^{|m|}(\mu) - \mu \frac{\partial P_n^{|m|}}{\partial \mu} \right] + \right\}$$

$$\frac{m}{(1-\mu^2)^{1/2}} P_n^{|m|}(\mu) \right\} N_n^m e^{i(m-1)\phi} r^{-n-2}, \tag{20}$$

$$= -\left[\frac{2n+1}{2n+3}(n-m+1)(n-m+2)\right]^{1/2} S_{n+1}^{m-1}.$$
 (21)

We store rotation coefficients  $H_n^{\nu m}$  for  $\nu=-n,\ldots,n,\ m=0,\ldots,n.$  By symmetry (G&D 5.5.2):

$$H_n^{\nu,-m} = H_n^{-\nu,m}$$

## References

[1] Nail A. Gumerov and Ramani Duraiswami. Recursions for the computation of multipole translation and rotation coefficients for the 3-D Helmholtz equation. *SIAM Journal on Scientific Computing*, 25(4):1344–1381, 2003.