

4D-TExS: A Novel 4D Lattice-based QCD Equation of State with Extended Density Range

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QCD Phase Diagram

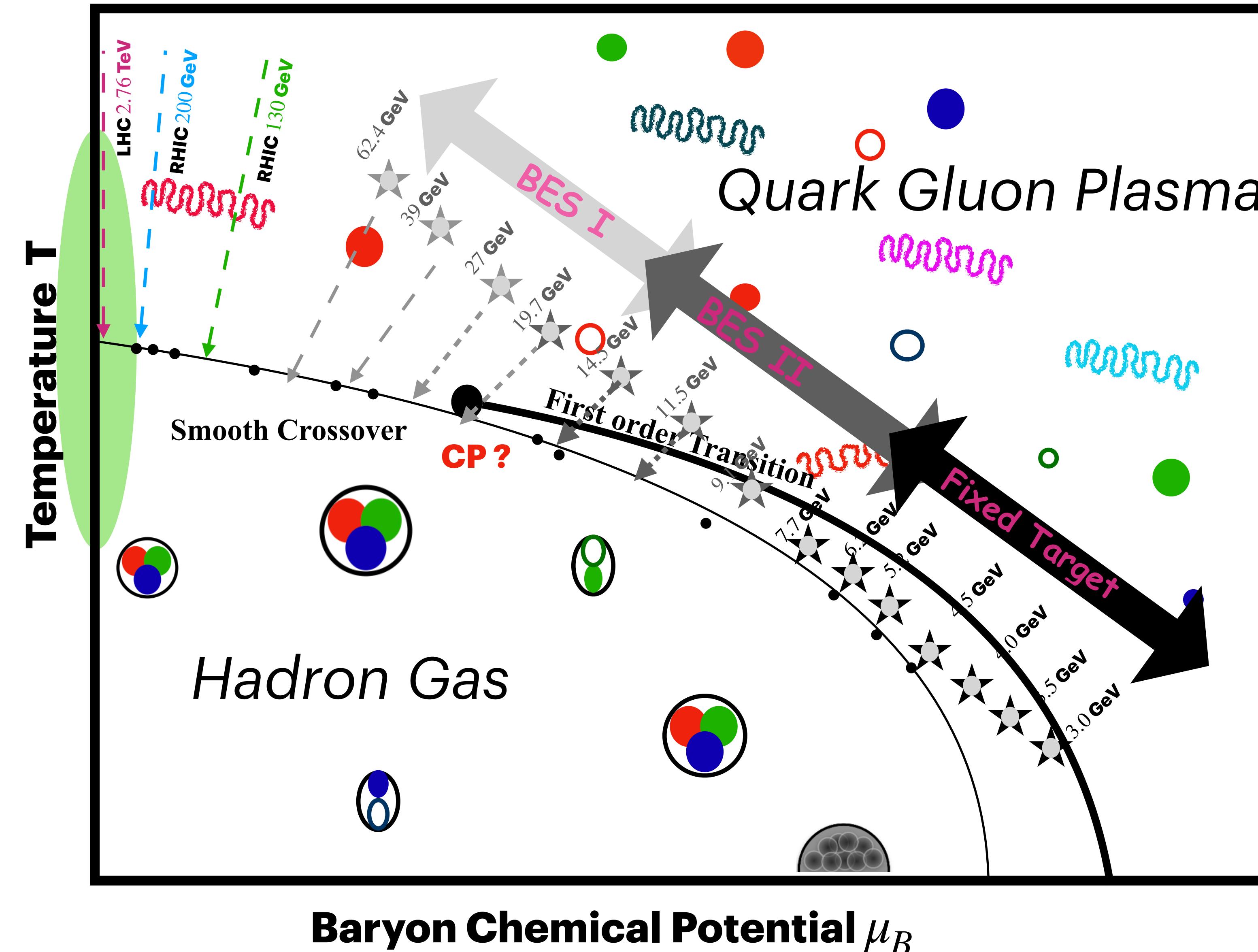
Experiments

- Finite density physics is achieved by lowering the $\sqrt{S_{NN}}$ in **BES program**
- Other experiments FAIR, NICA, J-PARC

Theoretical interpretation

- Transport models (Hydro)** model fireball evolution in heavy-ion collisions and matter dynamics in neutron star mergers.
- The **Equation of State (EoS)** is essential in hydro, to describe thermodynamics and evolution.

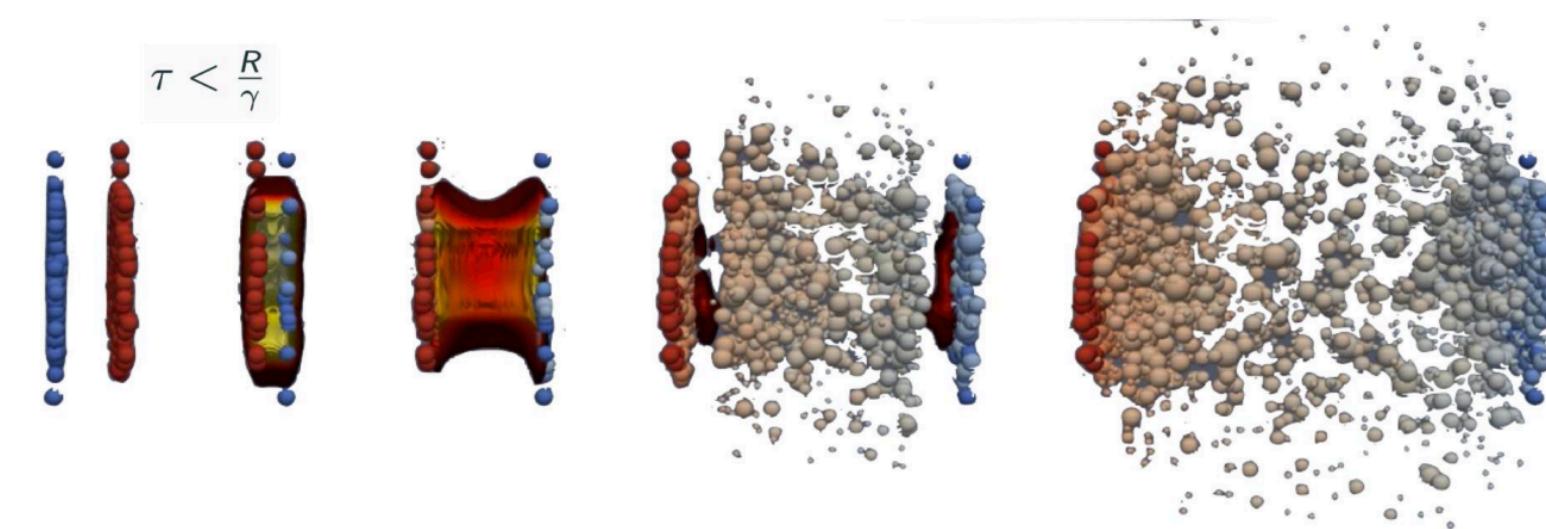
The EoS must capture relevant physics and cover much of the phase diagram



Why we need 4D EoS (T, μ_B, μ_Q, μ_S)

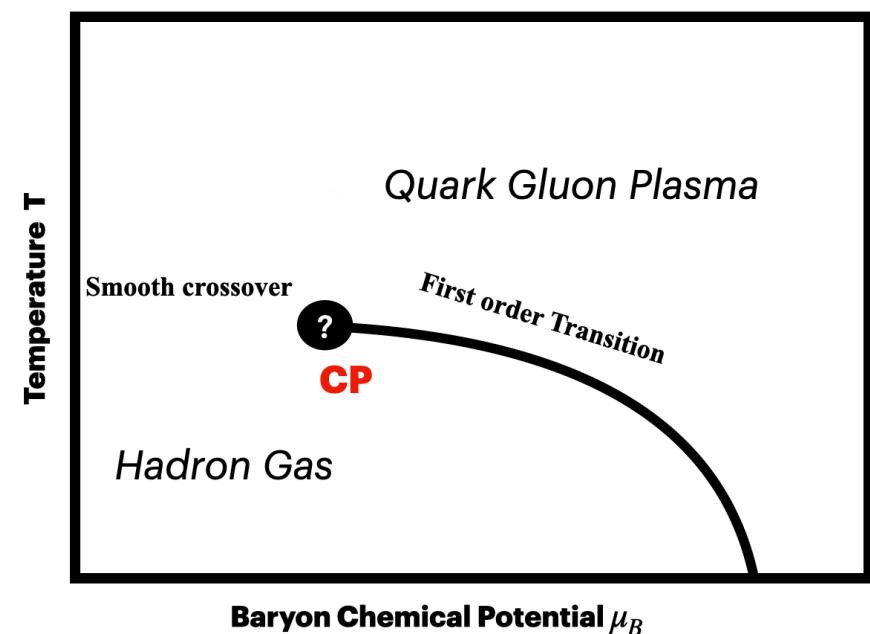
Heavy Ion collision, HI

- Net electric charge Q is non-zero
- To describe **kaons, hyperons** and **strange quark matter**, finite μ_S is necessary



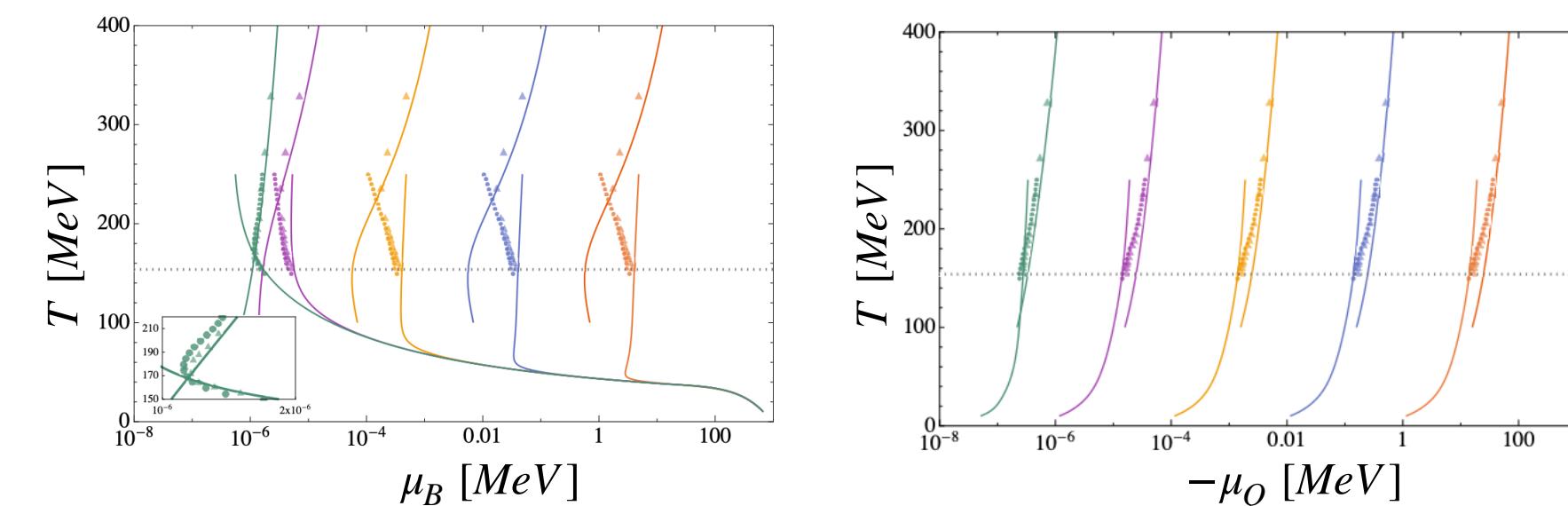
Phase Transitions and Critical point

- QCD phase structure and phase boundary shift differently in μ_B , μ_Q and μ_S (**BES Program**)



Cosmic trajectories

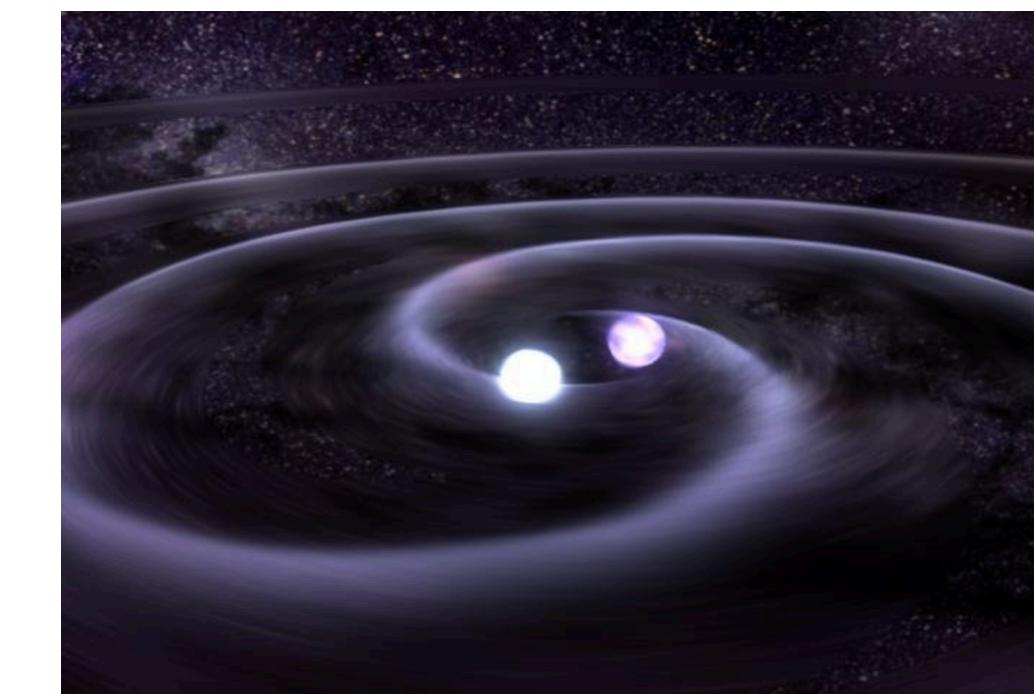
- μ_B , μ_Q and μ_S are important for cosmic trajectories



[Mandy M et al PRL. 121 20, 201302(2018)]

Neutron Star mergers

- To consider Asymmetric matter content



"The 4D EoS is in high demand for describing strongly interacting matter dynamics".

Taylor-expanded Lattice QCD

Limitations

2D Taylor Expansion around $\mu_B = 0$

$$\frac{P(T, \mu_B)}{T^4} = \sum_{n=0}^{\infty} \frac{1}{2n!} \chi_{2n}(T, \mu_B = 0) \left(\frac{\mu_B}{T} \right)^{2n}$$

$$\frac{\chi_n^B(T, \mu_B = 0)}{n!} = \frac{1}{n!} \left(\frac{\partial}{\partial(\mu_B/T)} \right)^n (P/T^4) \Big|_{\mu_B=0}$$

- Currently limited to $\mu_i/T \sim 2.5 - 3$, $i = \{B, Q, S\}$ despite great computational effort.
- Including one more higher-order term does not remove unphysical behavior due to truncation of Taylor series.

[Bollweg, D. et al Phys.Rev.D 108 1, 014510 (2023)]

[Borsanyi , S et al Phys.Rev.D 110 1, L011501 (2024)]

4D Taylor Expansion around $\hat{\mu}_B = \mu_B/T = 0$

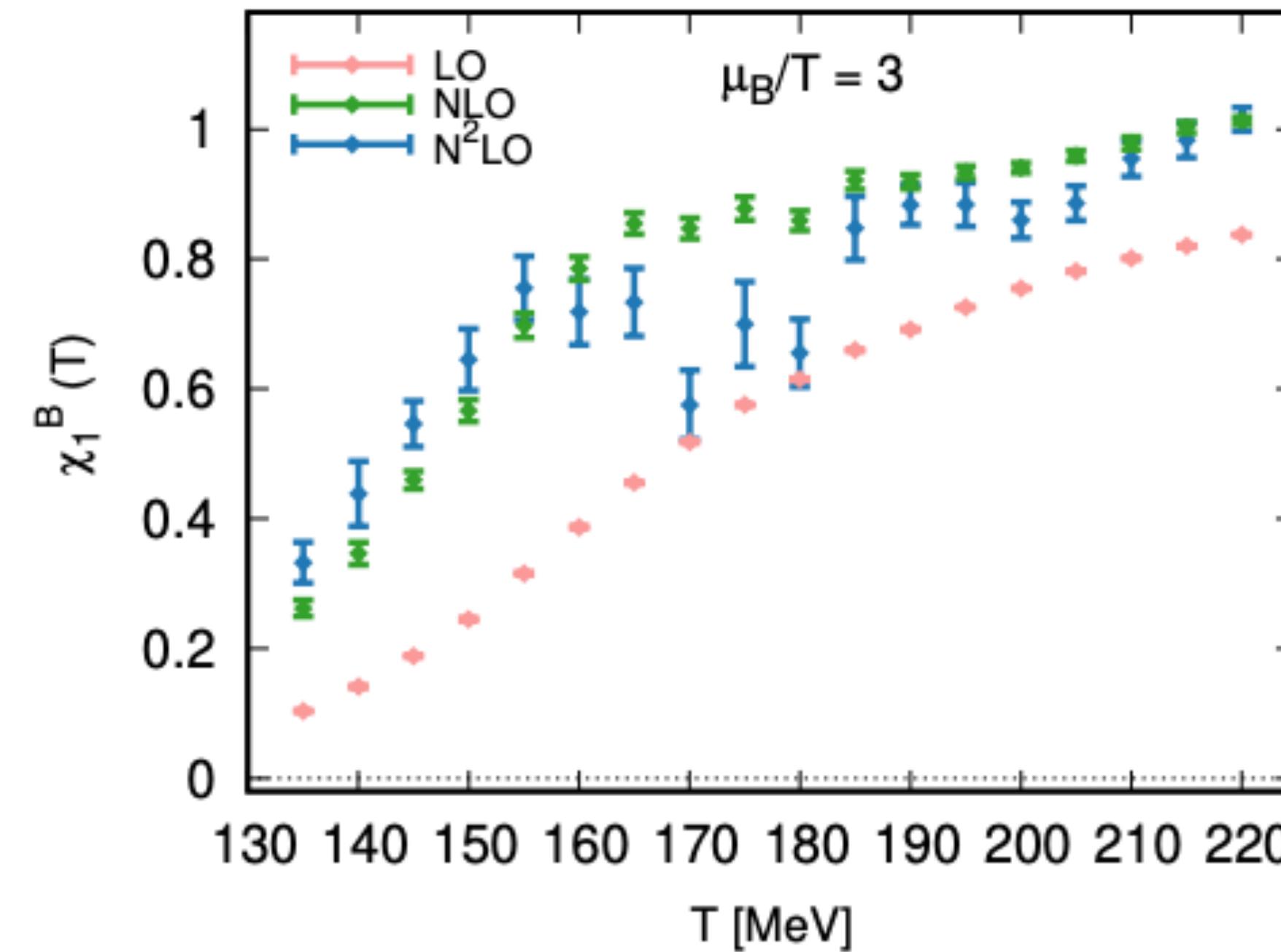
$$\frac{P(T, \hat{\mu}_B, \hat{\mu}_S, \hat{\mu}_Q)}{T^4} = \sum_{i,j,k} \frac{1}{i!j!k!} \chi_{ijk}^{BSQ}(T) \hat{\mu}_B^i \hat{\mu}_Q^j \hat{\mu}_S^k$$

$$\chi_{ijk}^{BQS}(T) = \frac{\partial^{i+j+k} (P/T^4)}{\partial \hat{\mu}_B^i \partial \hat{\mu}_Q^j \partial \hat{\mu}_S^k} \Big|_{\hat{\mu}_B, \hat{\mu}_Q, \hat{\mu}_S=0}$$

[J. Noronha-Hostler et al Phys.Rev.C 100 6, 064910 (2019)]

[A Monnai et al Phys.Rev.C 100 2, 024907 (2019)]

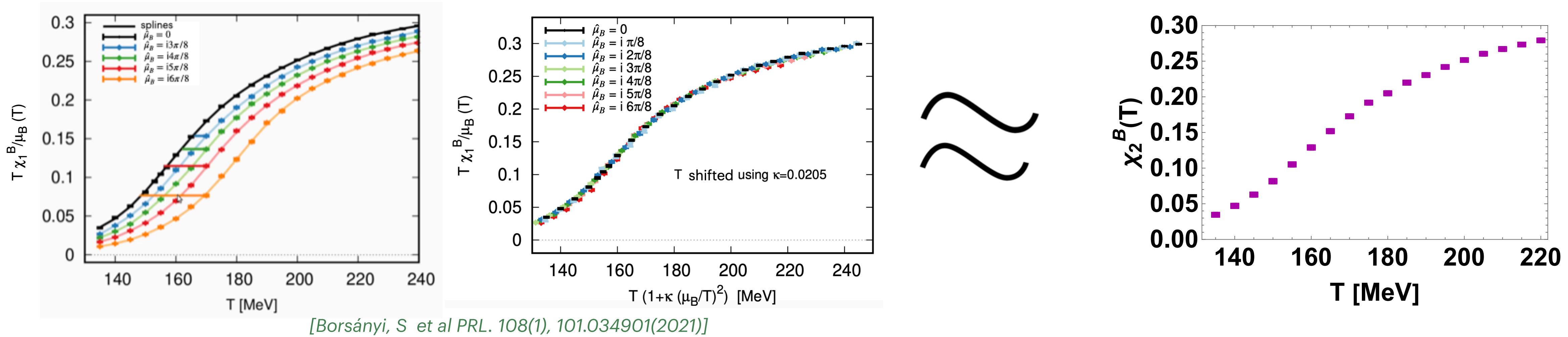
[A Monnai et al Phys.Rev.C 110 4, 044905 (2024)]



[Borsányi, S et al PRL. 108(1), 101.034901(2021)]

T' - Expansion scheme (T ExS)

Simulating at Imaginary μ_B



$$T \frac{\chi_1^B(T, \mu_B)}{\mu_B} = \chi_2^B(T, 0)$$

$$T'(T, \mu_B) = T \left[1 + \kappa_2^{BB}(T) \left(\frac{\mu_B}{T} \right)^2 + \kappa_4^{BB}(T) \left(\frac{\mu_B}{T} \right)^4 + \mathcal{O} \left(\frac{\mu_B}{T} \right)^6 \right]$$

T' -TExS with corrected SB-limit

$$\frac{\chi_1(T, \hat{\mu}_B)}{\chi_1^{SB}(\hat{\mu}_B)} = \frac{\chi_2(T, 0)}{\chi_2^{SB}(0)}$$

$$T'(T, \mu_B) = T [1 + \lambda_2(T) \hat{\mu}^2 + \dots]$$

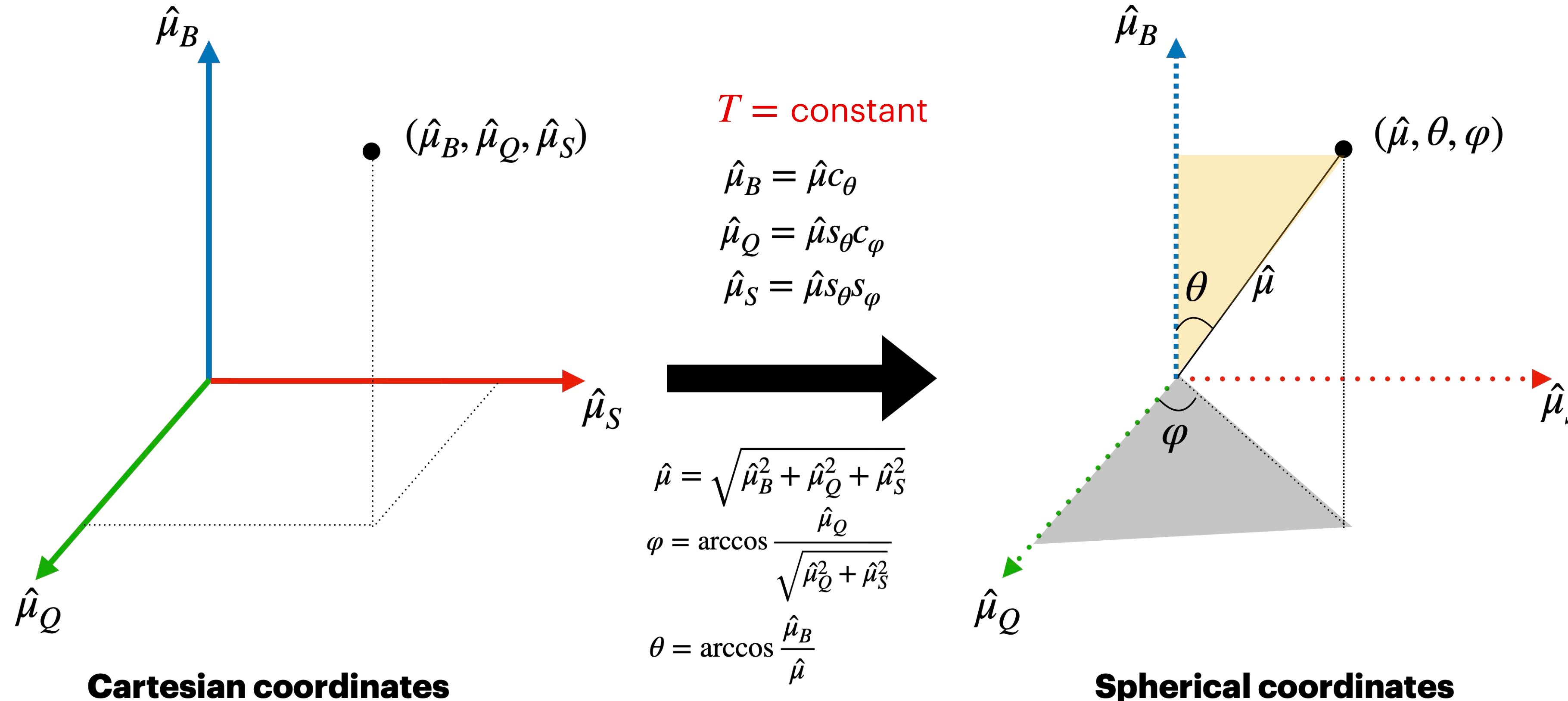
$$\lambda_2(T \rightarrow \infty) = 0$$

Pros

- Uses few expansion terms.
- μ_B dependence is captured in T' -rescaling.
- Trusted up to $\mu_B/T=3.5$ in the region where Critical point is expected.

[MK et al arXiv:2408.04588[nucl-th] (2024)]

Generalization to 4D- T' expansion scheme



2D T' Expansion

$$\frac{\chi_1(T, \hat{\mu}_B)}{\chi_1^{SB}(\hat{\mu}_B)} = \frac{\chi_2(T', 0)}{\chi_2^{SB}(0)}$$

4D T' Expansion

$$\frac{X_1^{\theta, \varphi}(T, \hat{\mu})}{\bar{X}_1^{\theta, \varphi}(\hat{\mu})} = \frac{X_2^{\theta, \varphi}(T'^{\theta, \varphi}(T, \hat{\mu}), 0)}{\bar{X}_2^{\theta, \varphi}(0)}$$

$$T'^{\theta, \varphi}(T, \hat{\mu}) = T \left(1 + \lambda_2^{\theta, \varphi}(T) \hat{\mu}^2 + \dots \right)$$

Apply 2D T' -expansion along a given trajectory at fixed θ and φ .

Generalization to 4D- T' expansion scheme

Generalized Susceptibilities

$$X_n^{\theta,\varphi}(T) = \frac{\partial^n p(T, \hat{\mu}, \theta, \varphi)/T^4}{\partial \hat{\mu}^n} \Bigg|_{\hat{\mu}=0}^{\theta,\varphi}$$

- Linear combination of χ_{ijk}^{BQS} (**Lattice data + HRG**)

$$X_2^{\theta,\varphi}(T) = c_\theta^2 \chi_2^B(T) + s_\theta^2 c_\varphi^2 \chi_2^Q(T) + s_\theta^2 s_\varphi^2 \chi_2^S(T) + \dots$$

$$X_4^{\theta,\varphi}(T) = c_\theta^4 \chi_4^B(T) + s_\theta^4 c_\varphi^4 \chi_4^Q(T) + s_\theta^4 s_\varphi^4 \chi_4^S(T) + \dots$$

Expansion parameter λ_2

$$\lambda_2^{\theta,\varphi}(T) = \frac{1}{6T(X_2^{\theta,\varphi}(T))'} \left(X_4^{\theta,\varphi}(T) - \frac{\bar{X}_4^{\theta,\varphi}(0)}{\bar{X}_2^{\theta,\varphi}(0)} X_2^{\theta,\varphi}(T) \right)$$

Generalized Charge density

$$X_1^{\theta,\varphi}(T, \hat{\mu}) = \frac{\bar{X}_1^{\theta,\varphi}(\hat{\mu})}{\bar{X}_2^{\theta,\varphi}(0)} X_2^{\theta,\varphi}(T'^{\theta,\varphi}(T, \hat{\mu}), 0)$$

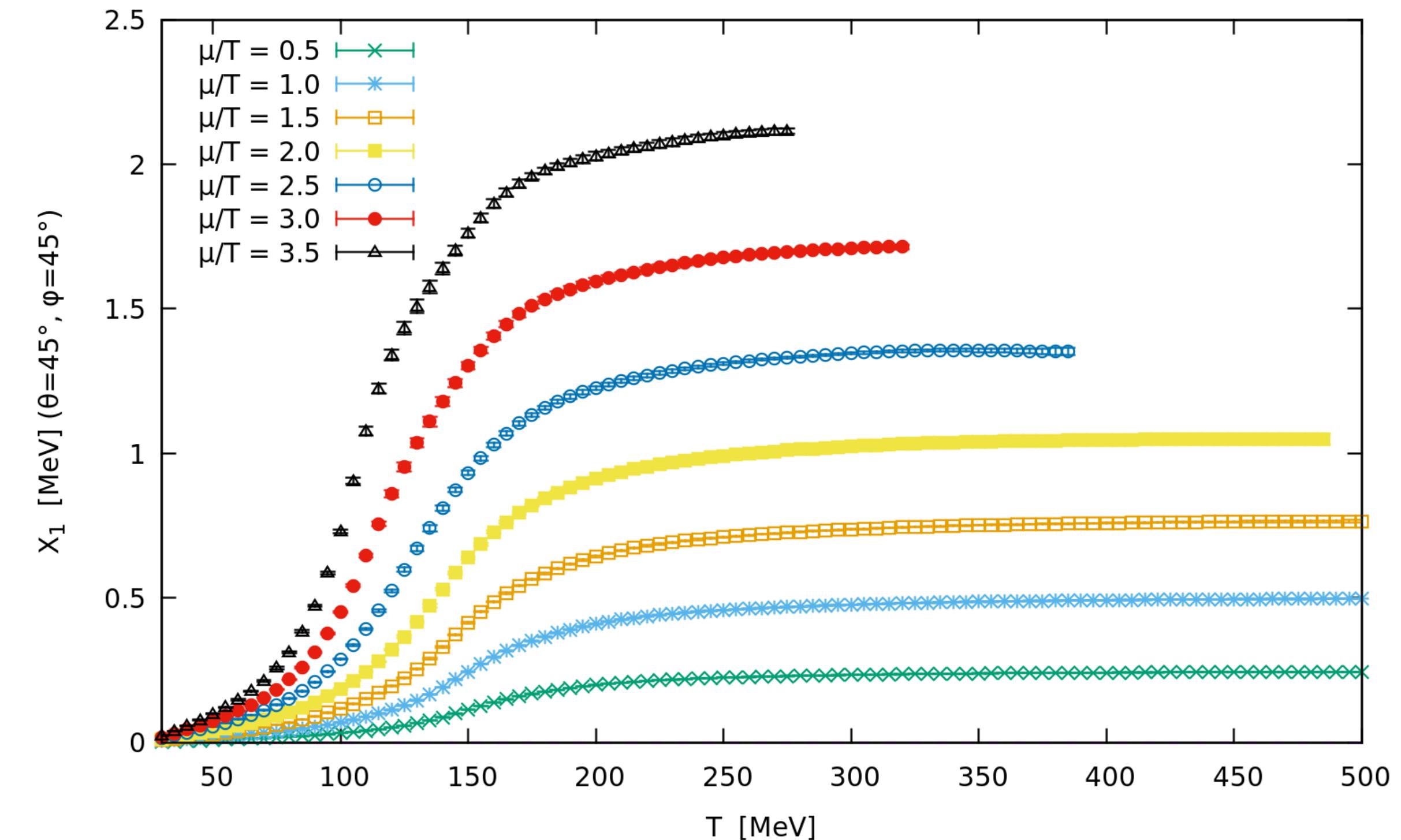
Example:

For $(\theta = 45^\circ, \varphi = 45^\circ)$

4D T' Expansion

$$\frac{X_1^{\theta,\varphi}(T, \hat{\mu})}{\bar{X}_1^{\theta,\varphi}(0)} = \frac{X_2^{\theta,\varphi}(T'^{\theta,\varphi}(T, \hat{\mu}), 0)}{\bar{X}_2^{\theta,\varphi}(0)}$$

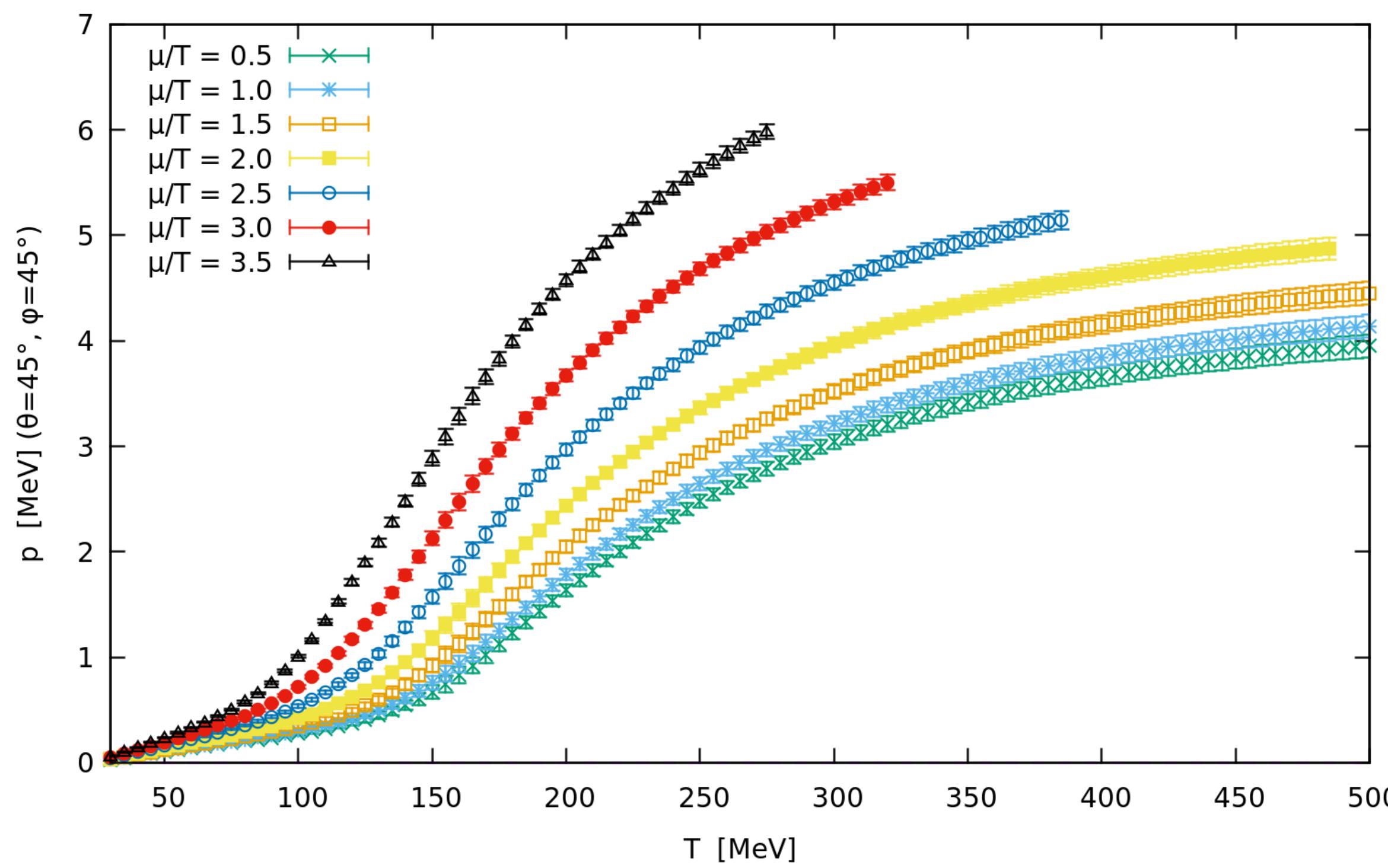
$$T'^{\theta,\varphi}(T, \hat{\mu}) = T \left(1 + \lambda_2^{\theta,\varphi}(T) \hat{\mu}^2 + \dots \right)$$



Thermodynamic Observable

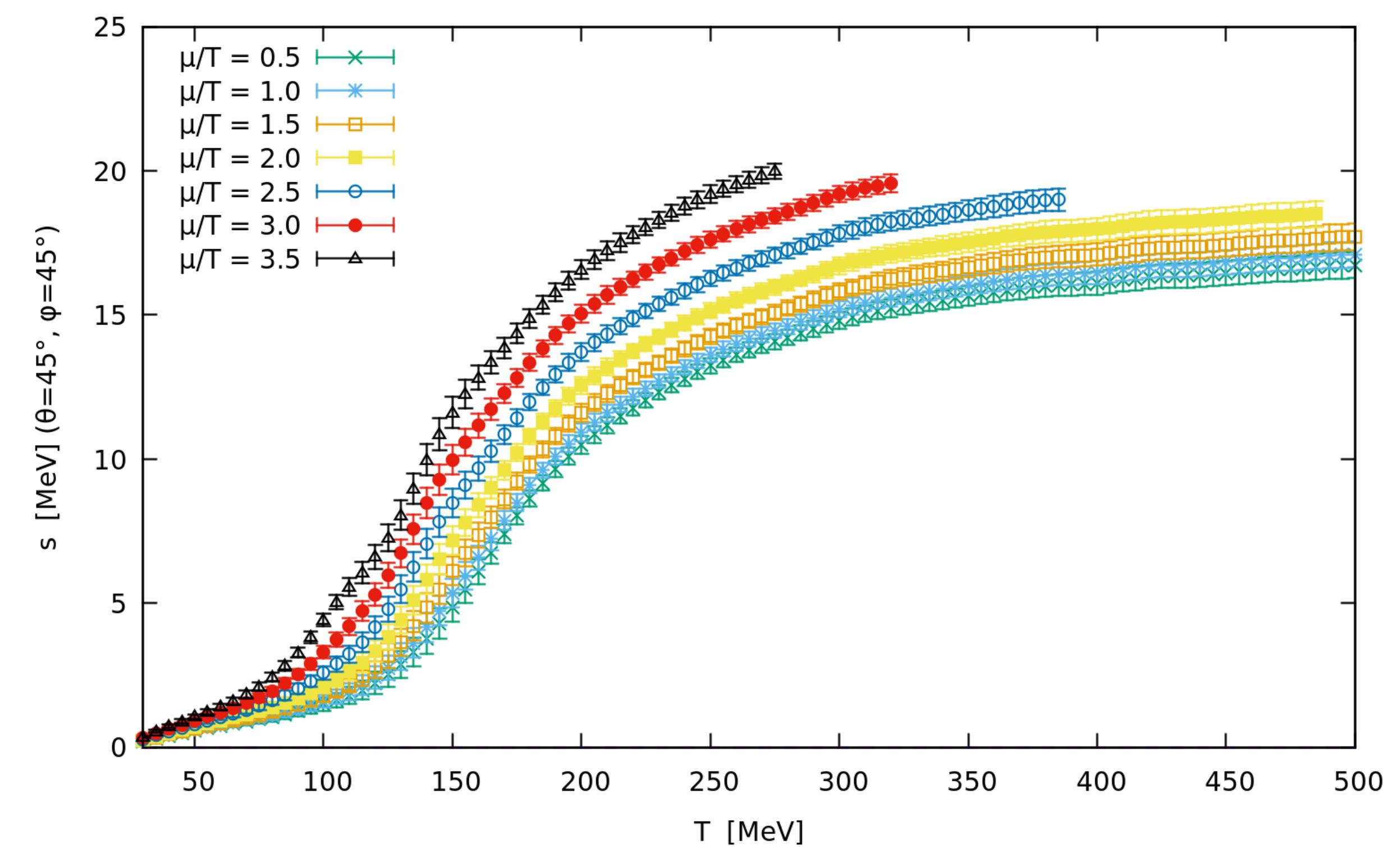
Pressure

$$\hat{p}(T, \hat{\mu}) = \hat{p}(T, \hat{\mu} = 0) + \int_0^{\hat{\mu}} d\hat{\mu}' X_1(T, \hat{\mu}')$$



Entropy

$$\hat{s}(T, \hat{\mu}) = \left. \frac{\partial \hat{p}(T, \hat{\mu})}{\partial T} \right|_{\hat{\mu}}$$



Conclusion and Outlook

- We have generalized the Lattice T' -Expansion Scheme to give 4D $(T, \hat{\mu}_B, \hat{\mu}_Q, \hat{\mu}_S)$ EoS.
- This construction includes a comprehensive and well-propagated **error analysis** based on new lattice data.
- We have constructed a **4D lattice EoS (BQS-TExS)** that show no unphysical behavior at **higher densities**.

Outlook

- Implement 4D-TExS in MUSES Calculation Engine



Thank you!