

# Hydrogen Bonding with Hydridic Hydrogen—Experimental Low-Temperature IR and Computational Study: Is a Revised Definition of Hydrogen Bonding Appropriate?

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Cite This: *J. Am. Chem. Soc.* 2023, 145, 8550–8559



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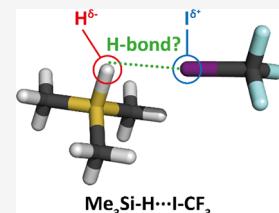
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**ABSTRACT:** Spectroscopic characteristics of  $\text{Me}_3\text{Si}-\text{H}\cdots\text{Y}$  complexes ( $\text{Y} = \text{ICF}_3$ ,  $\text{BrCN}$ , and  $\text{HCN}$ ) containing a hydridic hydrogen were determined experimentally by low-temperature IR experiments based on the direct spectral measurement of supersonically expanded intermediates on a cold substrate or by the technique of argon-matrix isolation as well as computationally at harmonic and one-dimensional anharmonic levels. The computations were based on DFT-D, MP2, MP2-F12, and CCSD(T)-F12 levels using various extended AO basis sets. The formation of all complexes related to the redshift of the Si–H stretching frequency upon complex formation was accompanied by an increase in its intensity. Similar results were obtained for another 10 electron acceptors of different types, positive  $\sigma$ -,  $\pi$ -, and p-holes and cations. The formation of  $\text{HBe}-\text{H}\cdots\text{Y}$  complexes, studied only computationally and again containing a hydridic hydrogen, was characterized by the blueshift of the Be–H stretching frequency upon complexation accompanied by an increase in its intensity. The spectral shifts and stabilization energies obtained for all presently studied hydridic H-bonded complexes were comparable to those in protonic H-bonded complexes, which has prompted us to propose a modification of the existing IUPAC definition of H-bonding that covers, besides the classical protonic form, the non-classical hydridic and dihydrogen forms.



## 1. INTRODUCTION

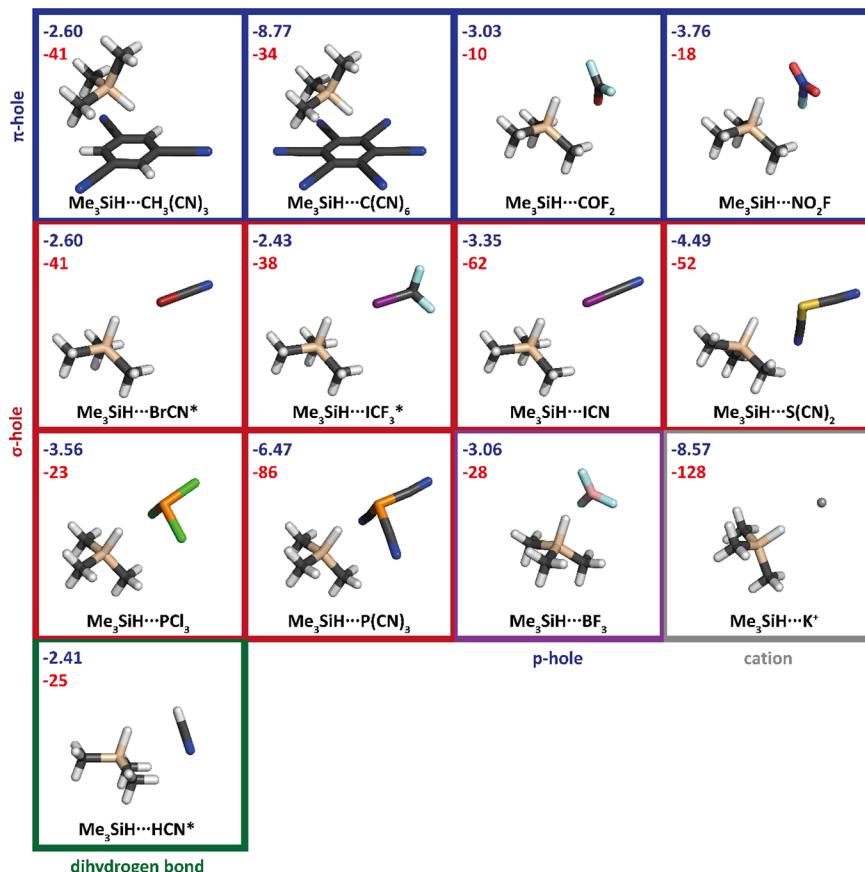
Most of the elements in the periodic table have lower electronegativity than hydrogen (2.2), and only a few of them (C, N, O, F, S, Cl, Se, Br, and I) have it higher. Covalent bonds of hydrogen with the more electronegative atom X are characterized by the polarization of the X–H bond and the formation of a partial positive charge on hydrogen (protonic hydrogen). The molecule thus acts as a Lewis acid, and when it interacts with the electron donor Y (Lewis base), a hydrogen bond (H-bond)  $\text{X}-\text{H}\cdots\text{Y}$  is formed.<sup>1</sup> H-bonds are among the strongest types of non-covalent interactions and are the most common. The specificity of H-bonding is related to its easily detectable spectroscopic manifestation, which originates in different masses of H and X atoms. The formation of the  $\text{X}-\text{H}\cdots\text{Y}$  H-bond, accompanied by a charge transfer from the Lewis base to the Lewis acid, results in a significant change of the X–H covalent bond. NBO orbital analysis<sup>2</sup> has revealed a charge transfer from the lone pair of Y to the X–H  $\sigma^*$  antibonding orbital. The increase in electron density in the  $\sigma^*$  antibonding orbital weakens the X–H covalent bond and lowers the X–H stretching frequency (redshift).<sup>3</sup> Our theoretical work at the end of the last century<sup>4</sup> showed that the formation of the hydrogen bond could also be accompanied by a blueshift of the X–H stretching frequency. When our predictions were proven experimentally,<sup>5</sup> it was obvious that there was a new type of H-bonding, for which we suggested a new name—blueshifting H-bonding. Extensive discussion in the computational and spectroscopic community led to the proposal of a new definition

covering both types of H-bonding. The new IUPAC definition<sup>6</sup> of the  $\text{X}-\text{H}\cdots\text{Y}$  hydrogen bond thus includes both the weakening and strengthening of the X–H covalent bond, leading to red- and blueshifts of the X–H stretching frequency. Nevertheless, the characterization of atom X (X is defined as more electronegative than hydrogen) remained the same. As shown above, however, most of the elements in the periodic table are less electronegative (more electropositive). In these complexes, the interaction scheme is the same, namely,  $\text{X}-\text{H}\cdots\text{Y}$ , but the hydrogen carries a negative charge (hydridic hydrogen) and the molecule acts as a Lewis base, whereas the Y atom carries a positive charge and the molecule acts as a Lewis acid. Both interaction types can be schematically represented as protonic,  $\text{X}-\text{H}^{\delta+}\cdots\text{Y}^{\delta-}$ , and hydridic,  $\text{X}-\text{H}^{\delta-}\cdots\text{Y}^{\delta+}$ . It should be emphasized that a very important feature of hydrogen bonding that makes its detection easy, namely, the position of light hydrogen between two much heavier atoms, is retained. This leads to an important question: does the second scheme correspond to H-bonding or is it another interaction type with a different definition and name? Many complexes containing

Received: January 20, 2023

Published: April 10, 2023





**Figure 1.** Geometries of the complexes investigated optimized at the MP2/cc-pwCVTZ level (cc-pwCVTZ-PP for Br and I): Si – beige, C – black, H – white, I – violet, F – light blue, Br – burgundy, N – dark blue, P – orange, Cl – green, S – yellow, O – red, K – silver, and B – rose. Dark blue is used for the total interaction energy (in kcal/mol) at the MP2/cc-pwCVTZ level for each complex, and red is used for the shift of the Si–H stretching frequency ( $\text{cm}^{-1}$ ). The complexes with the experiment are marked with asterisks.

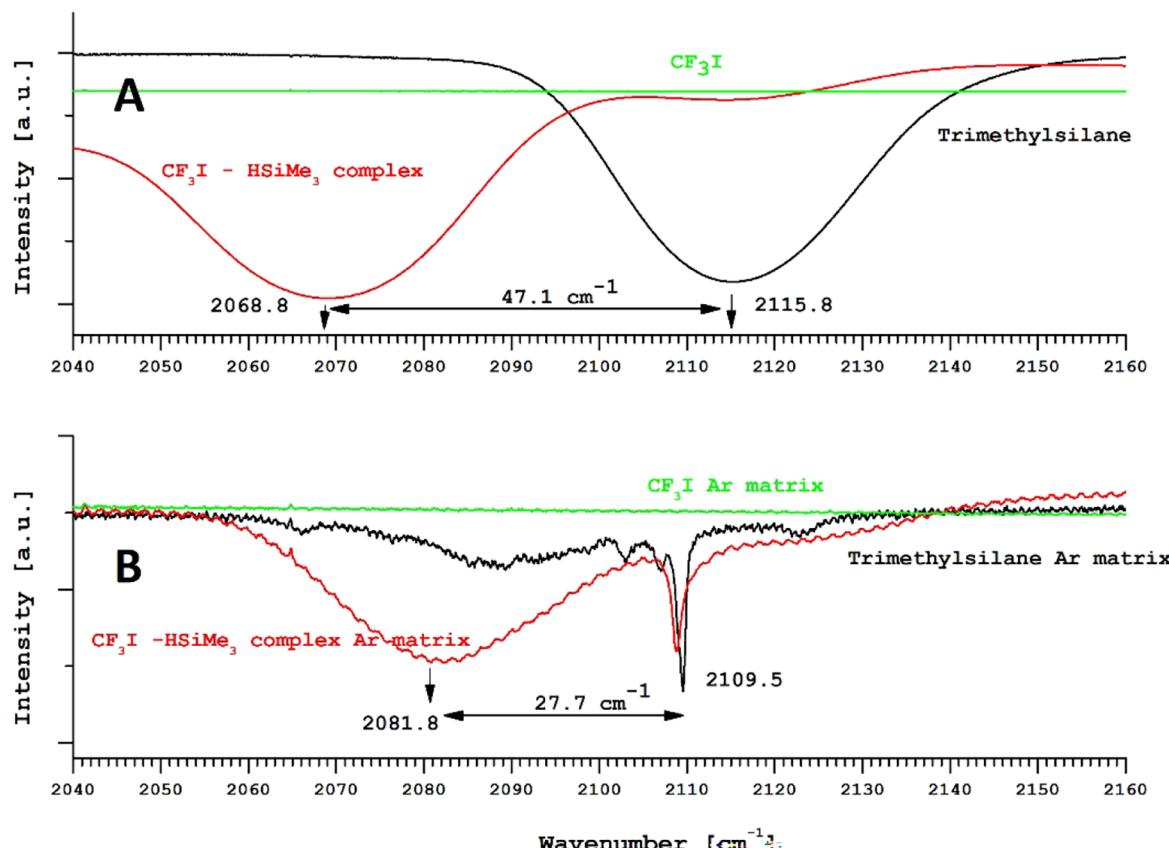
Lewis bases including Si–H, Ge–H, Be–H, Mg–H, Zn–H, Li–H, and Cu–H hydridic bonds and different Lewis acids have been studied computationally in the laboratory of Jabłoński, who has recommended using the name charge-inverted H-bond<sup>7–15</sup> (CIHB). The experimental detection of these non-classical H-bonds is, however, missing; their very existence is thus based on solving the Schrödinger equation only within the rigid rotor–harmonic oscillator–ideal gas approximation. The only exception is one class of CIHB, so-called dihydrogen bonding, where hydridic hydrogen interacts with protonic hydrogen. The structures of several dihydrogen-bonded complexes have been determined by X-ray and neutron diffraction studies; for some of them, theoretical calculations have revealed spectral shifts upon dihydrogen-bond formation.<sup>16</sup> Finally, clusters of phenol and aniline with borane–amines having the B–H…H–X dihydrogen bond have been studied in supersonic jets using electronic and vibrational spectroscopy.<sup>17</sup>

Noncovalent interactions have been studied using many different experimental approaches, starting from classical solution-phase studies involving spectroscopic and thermodynamic measurements. More recently, important, complementary information to that obtained in solutions has been provided by studies in the gas phase, particularly the supersonic beam studies of Flygare, Klemperer, and others.<sup>18–21</sup> The binding energy of weakly bound complexes is usually much smaller than room-temperature thermal energy. For this reason, the low temperature of a gas-phase supersonic jet, a rare-gas solid matrix,

or a liquid helium droplet is the typical temperature in which noncovalent complexes are studied in the laboratory.

Besides the direct high-resolution gas-phase techniques, there are also other very powerful low-temperature methods. The possibility of storing spectroscopically detectable concentrations of reaction intermediates in a solid or rare-gas matrix was first recognized by Whittle, Dows, and Pimentel.<sup>22</sup> These matrix materials are often chemically inert and are optically transparent from the far-infrared range well into the vacuum–ultraviolet region. The early experiments demonstrated that at 20 K, which is a temperature conveniently obtained using liquid hydrogen, solid nitrogen and argon are rigid enough to eliminate molecular diffusion and effectively inhibit subsequent chemical reactions. At the cryogenic temperatures required for studying rare-gas solids, molecules reside in their ground electronic and vibrational states. Since diffusion is inhibited, reaction intermediates do not undergo further reaction, and sufficient concentrations of many of them have been obtained for studies of their electronic and infrared spectra.

The aim of the present paper is to study the X'–H…Y' complexes (X' is more electropositive than hydrogen, X'–H acts as a Lewis base) containing Si–H and Be–H hydridic bonds and various hydridic–hydrogen acceptor Ys (Lewis acids) having a positive  $\sigma$ -,  $\pi$ -, and n-hole or a positive atom (e.g., hydrogen). Spectroscopic characteristics of the  $\text{Me}_3\text{Si}-\text{H}\cdots\text{Y}'$  (Y' =  $\text{ICF}_3$ ,  $\text{BrCN}$  and  $\text{HCN}$ ) complexes obtained not only at the harmonic but also at the more reliable anharmonic level have been verified by low-temperature IR experiments. Notice that the present



**Figure 2.** Solid-state and Ar-matrix spectra of the  $\text{Me}_3\text{Si}-\text{H}\cdots\text{ICF}_3$  complex.

experiments are the first ones to confirm unambiguously the formation of the CIHB. Finally, an attempt has been made to find a new definition of H-bonding that would cover both protonic (red- as well as blueshifting) and hydridic (CIHB) schemes.

## 2. INSTRUMENTATION AND SPECTRAL MEASUREMENT

In our low-temperature experiments, two different approaches have been applied (for more details, see the SI).

The first of them (A) is the technique of the direct spectral measurement of a supersonically expanded mixture of reaction intermediates on a cold substrate (solid-phase complex, SPC) and the second (B) is the technique of noble-gas matrix isolation (MI).

(A) The procedure of low-temperature experiments has already been described elsewhere.<sup>23,24</sup> The gas mixtures are usually deposited onto a cooled KBr substrate of the cryostat at a temperature of 4–20 K. The relative concentrations of the products have been monitored using the integrated absorption intensities of selected infrared bands. The intermediates forming part of the low-temperature complex have been supersonically expanded into a high vacuum ( $10^{-6}$  Torr) on a cold substrate (of 18 K, which is the minimum attainable temperature) inside a Leybold cryostat chamber. The spectra were obtained using a Bruker Vertex spectrometer with KBr optics, a HgCdTe detector, and a KBr beam splitter. The broad spectral region was cut by optical interference filters with transparency in the range of 700–5000  $\text{cm}^{-1}$ . The KBr entry window of the spectrometer was used. The unapodized spectral resolution was  $0.06 \text{ cm}^{-1}$ . Between 30 and 100 scans, depending on the sample,

were used to obtain a reasonable signal-to-noise ratio. The observed wavenumbers were calibrated using  $\text{CO}_2$ -absorption rotation-vibration lines.

(B) Matrix isolation is a well-known technique frequently used for the measurement of unstable species such as ions, radicals, and low-temperature-existing molecular complexes in a cold matrix of noble gas (Ng). Like in the procedure A, a mixture of reaction intermediates mixed together with argon gas (molar ratio 1:1000) was expanded through a pulse nozzle onto the cold (18 K) KBr substrate, and the spectra were recorded using the Bruker Vertex spectrometer.

## 3. SYSTEM CONSIDERED

$\text{Me}_3\text{Si}-\text{H}$  and  $\text{HBe}-\text{H}$  systems containing a hydridic hydrogen and different electron acceptors ( $\text{ICF}_3$ ,  $\text{BrCN}$ ,  $\text{HCN}$ ,  $\text{K}^+$ ,  $\text{C}_6(\text{CN})_6$ ,  $\text{C}_6\text{H}_5(\text{CN})_3$ ,  $\text{BF}_3$ ,  $\text{P}(\text{CN})_3$ ,  $\text{PCl}_3$ ,  $\text{S}(\text{CN})_2$ ,  $\text{NO}_2\text{F}$ ,  $\text{CO}_2\text{F}$ , and  $\text{ICN}$ ) forming CIHBs have been considered. Lewis bases contain a positive  $\sigma$ -,  $\pi$ -, and p-hole or a positive hydrogen. The latter complexes are also known as complexes with a dihydrogen bond. All complexes considered are depicted in Figure 1.

In this paper, we use both techniques A and B. It is well known that the method of matrix isolation gives narrow lines, rare-gas atoms isolate the molecules from mutual interaction, and subsequent chemical reactions are effectively inhibited. On the contrary, the A arrangement provides the possibility to study the basic energy characteristic of the molecular complexes, such as thermodynamic stability, on the board temperature scale.

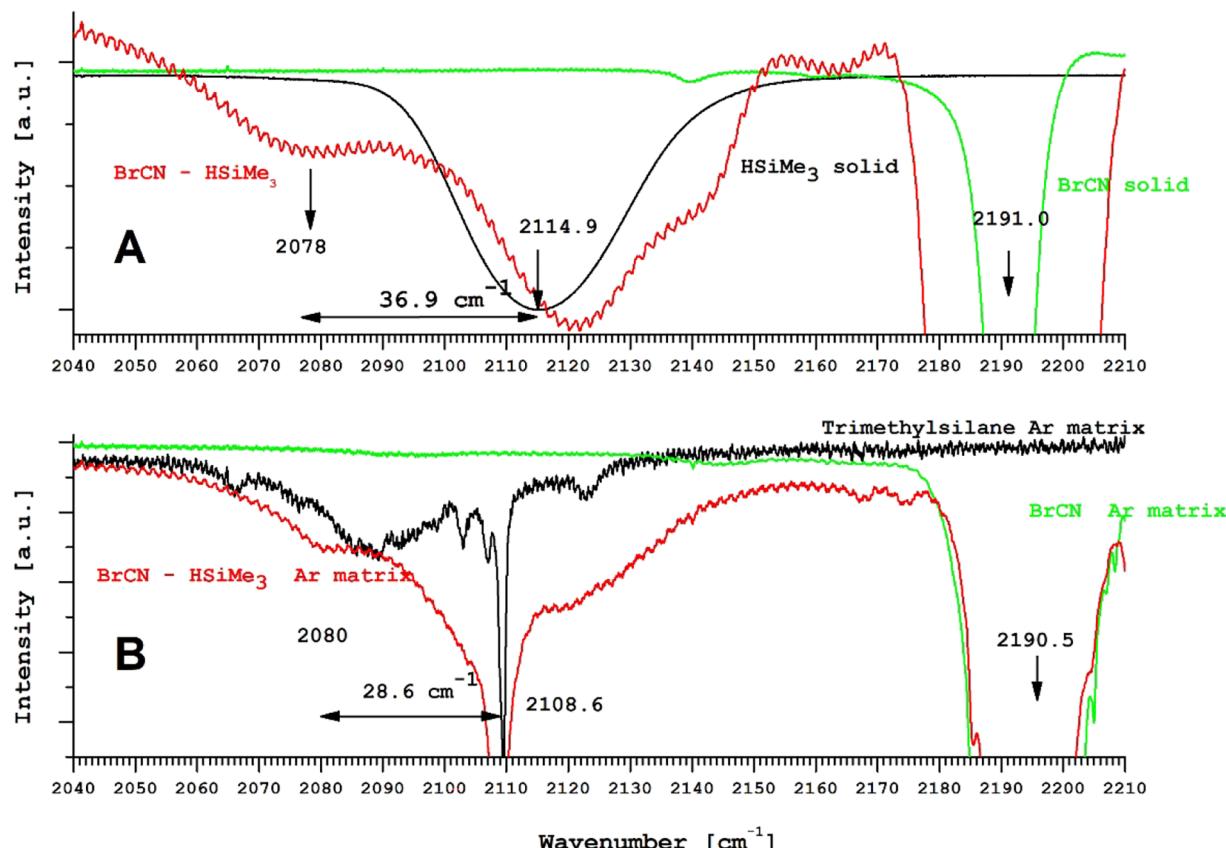


Figure 3. Solid-state and Ar-matrix spectra of the  $\text{Me}_3\text{Si}-\text{H}\cdots\text{BrCN}$  complex.

#### 4. METHODS

**4.1. Geometry Optimization and Thermodynamic Properties.** Geometries of all subsystems and complexes were optimized at the RI-MP2/cc-pwCVTZ<sup>25</sup> level. For heavy halogens (Br, I), basis sets with pseudopotentials (cc-pwCVTZ-PP<sup>26</sup>) were used. Thermodynamic properties as well as harmonic vibration frequencies, determined using the rigid rotor–harmonic oscillator–ideal gas approximation, were evaluated at the same theoretical level. Geometry optimization and harmonic vibrational analysis on selected complexes ( $\text{ICF}_3$ ,  $\text{BrCn}$ , and  $\text{HCN}$  complexes with  $\text{Me}_3\text{SiH}$ ) were also performed using the more reliable explicitly correlated RI-MP2-F12<sup>27</sup> method in cc-pVTZ-F12<sup>28</sup> basis sets (cc-pVTZ-PP-F12 for Br and I). All calculations were performed with MOLPRO 2022.<sup>29,30</sup> The potential energy surfaces for the anharmonic calculations of Si–H frequencies were obtained on the MP2/cc-pVTZ level (cc-pwCVTZ-PP for Br and I) using the Cuby4<sup>31</sup> framework and the TURBOMOLE 7.5<sup>32</sup> program.

**4.2. Single-Point Energy Calculations.** CCSD(T)-F12<sup>33</sup> energies were determined with cc-pVTZ-F12 (cc-pVTZ-PP-F12 for Br and I) basis sets implemented in the MOLPRO 2022 package on geometries obtained from the MP2-F12 method. All interaction energies were systematically corrected using the Boys and Bernardi counterpoise technique.<sup>34</sup> Energy decomposition analysis was performed using the SAPT2 + 3<sup>35</sup>/aug-cc-pwCVTZ (aug-cc-pwCVTZ-PP for Br and I) method implemented in the PSI4 package.

**4.3. NBO Analysis.** NBO analysis was performed at the  $\omega\text{B}97\text{X-D}/\text{cc-pwCVTZ}$  level (cc-pwCVTZ-PP for Br and I) using the NBO<sup>2</sup> program implemented in Gaussian 16.<sup>37</sup>

**4.4. One-Dimensional Harmonic and Anharmonic Vibrational Analysis.** Like in our previous publications,<sup>38,39</sup> the evaluation of the sought vibrational Si–H stretching frequencies relies on the deep adiabatic separability of the Si–H mode ( $s$ ) from the rest of the vibrational degrees of freedom of the probed compounds and on the use of a HBJ non-rigid reference configuration of the atomic nuclei that essentially follows the Si–H motion.<sup>40</sup> The vibrational energies can

then be obtained by solving the Schrödinger equation for the following Hamiltonian:

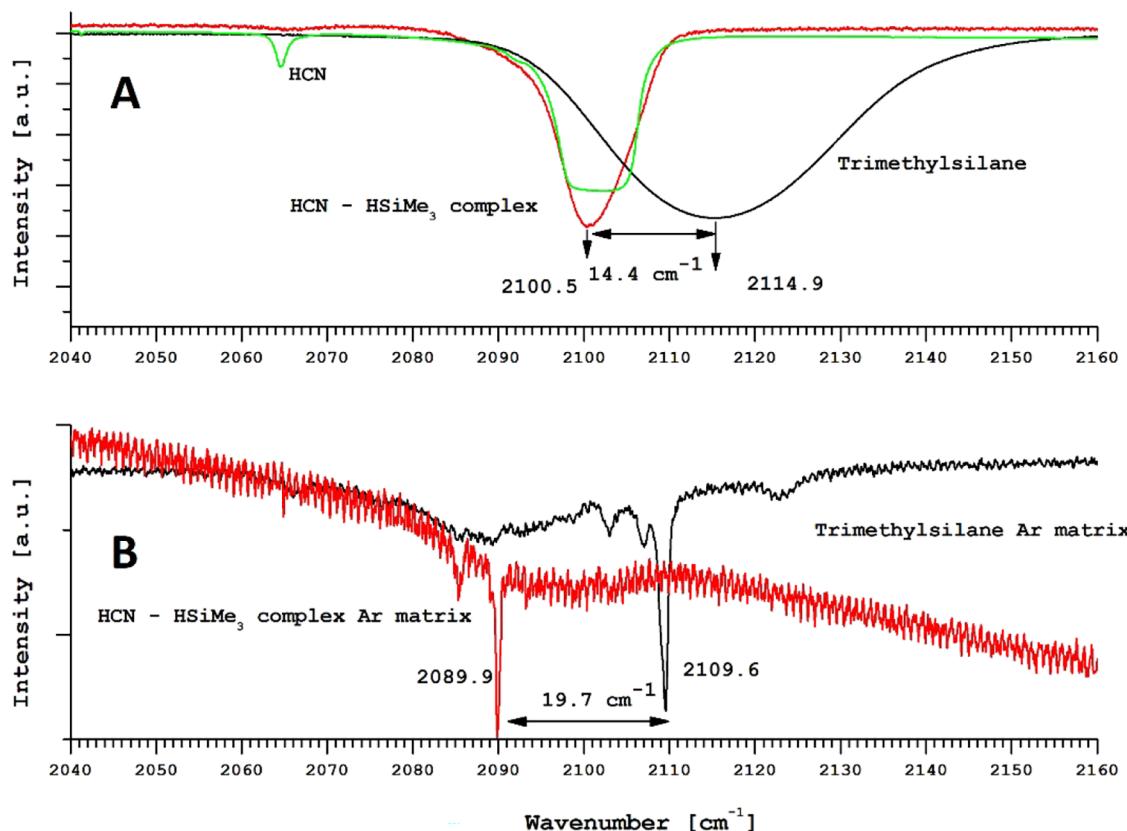
$$H_s = \frac{-1}{2} \mu_{ss} J_s^2 + \frac{1}{2} (J_{ss} \mu_{ss}) J_s + V_{\text{pseudo}}(s) + V(s)$$

where  $J_s = -i\hbar(d/ds)$ ,  $\mu_{ss}$  is the Si–H stretching component of the tensor that is the inverse of the  $4 \times 4$  generalized molecular inertia tensor,  $\mu$  is the determinant of the matrix  $[\mu_{\alpha\beta}]$  ( $\alpha, \beta = x, y, z, s$ ), with  $x$ ,  $y$ , and  $z$  being the Cartesian atomic coordinates in the molecule-fixed axis system,  $V_{\text{pseudo}}(s)$  is a mass-dependent kinematic pseudopotential, and  $V(s)$  is the Si–H minimum-energy-path stretching potential.

Obviously, mainly for the practical impossibility to account reliably for the aggregation effects of the molecular environments used, the adopted “isolated-molecule” theory may seem inadequate. Nevertheless, if one assumes a purely linear dependence of the atomic coordinates of the probed compounds (deposited on cold substrates or trapped in rare-gas matrices) on the stretching Si–H distortion, the stretching-reduced mass  $\mu_{ss}$  becomes constant and can thus be used as a single scaling parameter. Interestingly, as illustrated in Figure S1, the assumption of “linearity” seems to hold at least for the minimum-energy-path relaxation effects in isolated molecules (note that the  $x$  and  $y$  coordinates exhibit even much lower relaxation dependence than their  $z$ -counterparts). As shown in Figure S2, the “complete” scaling factors corresponding to the observed spectral shifts  $\Delta\nu$  obtained using the argon matrix-isolation technique acquire rather coinciding values, thus explicitly justifying the use of the “linearity” assumption for the rationalization of the Ar-matrix-isolation data. The anharmonic calculation thus appears to be a useful complement to the standard normal coordinate analysis.

#### 5. RESULTS AND DISCUSSION

**5.1. Experiment.** All complexes were studied in a solid phase on a cold (18 K) KBr substrate and simultaneously in the Ar matrix.



**Figure 4.** Solid-state and Ar-matrix spectra of the  $\text{Me}_3\text{Si}-\text{H}\cdots\text{HCN}$  complex.

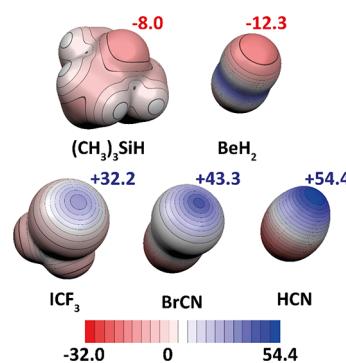
**5.1.1.  $\text{Me}_3\text{Si}-\text{H}\cdots\text{ICF}_3$  Complex.** Panel (A) of Figure 2 depicts the spectrum of trimethylsilane (black) on a cold (18 K) KBr substrate together with the spectrum of the trimethylsilane- $\text{CF}_3\text{I}$  1:1 molecular complex ( $2068.8 \text{ cm}^{-1}$ ) at 18 K. The obtained redshift of the Si–H bond was estimated to be about  $47.1 \text{ cm}^{-1}$ . Panel (B) shows the Ar-matrix spectrum of pure trimethylsilane (black) and the spectrum of the  $\text{Me}_3\text{Si}-\text{H}\cdots\text{CF}_3\text{I}$  complex at 18 K (red). The measured Si–H bond shift inside the argon matrix was  $27.7 \text{ cm}^{-1}$ .

**5.1.2.  $\text{Me}_3\text{Si}-\text{H}\cdots\text{BrCN}$  Complex.** Panel (A) in Figure 3 shows the spectrum of trimethylsilane (black) on a cold (18 K) KBr substrate together with the absorption spectrum of BrCN at 18 K (green) and the  $\text{Me}_3\text{Si}-\text{H}\cdots\text{BrCN}$ –silane 1:1 molecular mixture at 18 K. If we accept the formation of the  $\text{Me}_3\text{Si}-\text{H}\cdots\text{BrCN}$  molecular complex (the broad peak around  $2078 \text{ cm}^{-1}$ ), the obtained redshift of the Si–H bond gives the value of  $36.9 \text{ cm}^{-1}$  (see Figure 3A). Panel (B) shows the Ar-matrix spectrum of pure trimethylsilane (black) and the spectrum of the Ar– $\text{Me}_3\text{Si}-\text{H}\cdots\text{BrCN}$  mixture (2000,1) at 18 K (red) and the argon-matrix spectrum of BrCN (green). From the argon-matrix spectra (Figure 3B), we attributed the broad peak around  $2080 \text{ cm}^{-1}$  to the molecular  $\text{Me}_3\text{Si}-\text{H}\cdots\text{BrCN}$  complex, with the resulting redshift of the H–Br bond being  $28.6 \text{ cm}^{-1}$ .

**5.1.3.  $\text{Me}_3\text{Si}-\text{H}\cdots\text{HCN}$  Complex.** Panel (A) in Figure 4 depicts the spectrum of trimethylsilane (black) on a cold (18 K) KBr substrate together with the HCN molecular absorption band at 18 K (green) and the  $\text{Me}_3\text{Si}-\text{H}\cdots\text{HCN}$  1:1 molecular complex ( $2100.5 \text{ cm}^{-1}$ ) at 18 K. The obtained redshift of the Si–H bond was estimated to be about  $14.4 \text{ cm}^{-1}$ . Panel (B) shows the Ar-matrix spectrum of pure trimethylsilane (black) and the spectrum of the  $\text{Me}_3\text{Si}-\text{H}\cdots\text{HCN}$  complex at 18 K. The

measured Si–H bond shift inside the argon matrix was  $19.7 \text{ cm}^{-1}$ .

**5.2. Calculations.** **5.2.1. Subsystems.** Electrostatic potentials including the  $V_{s,\max}$  and  $V_{s,\min}$  values for the optimized structures of selected Lewis bases and acids (for which experimental measurements have been performed) are visualized in Figure 5. The most basic hydridic hydrogen (i.e.,



**Figure 5.** MEP calculated at the  $\omega\text{B97X-D}/\text{def2-QZVPP}$  level for the monomers studied. The ESP scale is in kcal/mol. The  $V_{s,\max}$  on the hydrogen atom from the  $\text{Me}_3\text{SiH}$  molecule is depicted in red. The  $V_{s,\min}$  in the center of the  $\sigma$ -hole or on top of the hydrogen atom is in blue.

the strongest electron donor) has been found in  $\text{BeH}_2$ , followed by  $\text{Me}_3\text{Si}-\text{H}$ , while the strongest electron acceptor has been detected in HCN.

**5.2.2. Complexes.** The intramolecular X’–H (X’–H) and intermolecular H...Y’ distances for  $\text{Me}_3\text{Si}-\text{H}\cdots\text{Y}'$  complexes are shown in Table 1, whereas these characteristics for  $\text{BeH}_2$  complexes are summarized in Table S1. In comparison with

**Table 1. Intermolecular Distances and Changes in Intramolecular Bond Lengths (in Å) of Selected Bonds in  $\text{Me}_3\text{SiH}\cdots\text{Y}'$  Complexes ( $\text{Y}' = \text{ICF}_3, \text{BrCN}, \text{HCN}$ ) Calculated at the MP2-F12/cc-pVTZ-F12 Level<sup>b</sup>**

|                | $\Delta r(\text{Si}-\text{H})$ | $\Delta r(\text{Si}-\text{C})$ | $\Delta r(\text{I}-\text{C})$  | $\text{H}\cdots\text{I}/\text{vdW}^{\text{a}}$  |
|----------------|--------------------------------|--------------------------------|--------------------------------|---|
| $\text{ICF}_3$ | 0.007                          | -0.003/-0.001/-0.001           | 0.002                          | 2.874/-0.573                                    |
|                | $\Delta r(\text{Si}-\text{H})$ | $\Delta r(\text{Si}-\text{C})$ | $\Delta r(\text{Br}-\text{C})$ | $\text{H}\cdots\text{Br}/\text{vdW}^{\text{a}}$ |
| $\text{BrCN}$  | 0.008                          | -0.003/-0.002/-0.002           | 0.005                          | 2.705/-0.585                                    |
|                | $\Delta r(\text{Si}-\text{H})$ | $\Delta r(\text{Si}-\text{C})$ | $\Delta r(\text{H}-\text{C})$  | $\text{H}\cdots\text{H}/\text{vdW}^{\text{a}}$  |
| $\text{HCN}$   | 0.005                          | -0.003/-0.003/0                | 0                              | 2.729/+0.320                                    |

<sup>a</sup>The difference between the intermolecular distance and the sum of the vdW radii. <sup>b</sup>All values are in Å.

the sum of the vdW radii,<sup>41</sup> the intermolecular distances in complexes with  $\text{ICF}_3$  and  $\text{BrCN}$  are significantly shorter (by 0.573 and 0.585 Å); in the complex with  $\text{HCN}$ , it is longer. Table 1 further shows that all intramolecular Si–H distances are systematically elongated upon complex formation. Notice that the elongation of the Si–H bond is smaller than that of the X–H bonds in the X–H···Y H-bonded systems. As expected, the Si–C bonds are contracted, but the respective changes are smaller. Finally, upon complex formation, the I–C and Br–C intramolecular distances are elongated, whereas the H–C distance is not changed. Qualitatively similar results have also been obtained for  $\text{HBeH}\cdots\text{Y}'$  complexes.

**5.2.3. Energies.** Table 2 contains calculated energy characteristics for  $\text{Me}_3\text{SiH}\cdots\text{Y}'$  complexes, specifically MP2, MP2–F12,

**Table 2. Energy Characteristics of the  $\text{Me}_3\text{SiH}\cdots\text{Y}'$  Complexes ( $\text{Y}' = \text{ICF}_3, \text{BrCN}, \text{HCN}$ )<sup>a</sup>**

| $\text{Y}'$    | $\Delta E^{\text{MP2}}/\Delta E^{\text{MP2-F12}}/\Delta E^{\text{CCSD(T)-F12}}$ | $\Delta G^{\text{MP2}}(18 \text{ K})/\Delta G^{\text{MP2-F12}}(18 \text{ K})$ |
|----------------|---|---|
| $\text{ICF}_3$ | -2.43/-3.27/-2.70   | -1.85/-2.42   |
| $\text{BrCN}$  | -2.60/-3.02/-2.76   | -2.07/-2.77   |
| $\text{HCN}$   | -2.41/-2.59/-2.24   | -1.79/-1.97   |

<sup>a</sup>The total interaction energy calculated at the MP2/cc-pwCVTZ level ( $\Delta E^{\text{MP2}}$ ) and at the MP2-F12/cc-pVTZ-F12 level ( $\Delta E^{\text{MP2-F12}}$ ), intrinsic interaction energy with counterpoise correction at the CCSD(T)-F12/cc-pVTZ-F12 level on the MP2-F12/cc-pVTZ-F12 geometries ( $\Delta E^{\text{CCSD(T)-F12}}$ ) and Gibbs free energy at MP2/cc-pwCVTZ ( $\Delta G^{\text{MP2}}$ ) and MP2-F12/cc-pVTZ-F12 ( $\Delta G^{\text{MP2-F12}}$ ) levels. All values are in kcal/mol.

and CCSD(T)-F12 interaction energies and binding free energies at 18 K based on the MP2 and MP2–F12 characteristics. Corresponding results obtained for  $\text{BeH}_2$  complexes are collected in Table S2.

The MP2 stabilization energies of all CIHB complexes are comparable, whereas the more reliable MP2–F12 and mainly CCSD(T)–F12 energies of the first two complexes are similar and are higher than that of the third (dihydrogen-bonded) one. Further, all MP2–F12 and CCSD(T)–F12 stabilization energies lie between 2.6 and 3.3 kcal/mol and between 2.2 and 2.8 kcal/mol, respectively, and are thus well comparable to those of classical H-bonded complexes. These findings contradict the results from the previous subchapter (Subsystems), showing that the strongest electron acceptor is  $\text{HCN}$ , followed by  $\text{BrCN}$  and  $\text{ICF}_3$ . Therefore, the complex with  $\text{HCN}$  was expected to be the strongest. The fact that the opposite is true, i.e., this complex is the weakest, is caused by dispersion energy. The SAPT2 + 3<sup>35</sup> dispersion and total interaction energies for

the complexes of  $\text{Me}_3\text{SiH}$  with  $\text{CF}_3\text{I}$ ,  $\text{BrCN}$ , and  $\text{HCN}$  amount to -4.35, -3.79, and -3.05 kcal/mol and -3.34, -3.20, -2.50 kcal/mol, respectively. Clearly, the smallest dispersion energy for the last complex is responsible for its smallest stabilization energy. Notice that the SAPT2 + 3 interaction energies for all three complexes agree surprisingly well with the MP2–F12 ones. The negative binding free energies calculated for all the complexes indicate their formation at 18 K. The MP2 and CCSD(T) interaction energies of  $\text{BeH}_2$  complexes are systematically smaller, with the largest values found for the  $\text{HBeH}\cdots\text{BrCN}$  complex. The binding free energies of all complexes are smaller than those of  $\text{Me}_3\text{SiH}$ , but they are still negative, which ensures their formation at 18 K.

Table 3 shows the difference in the orbital occupancies (of monomers and complexes) of occupied and unoccupied Si–H

**Table 3. Orbital–Occupation Difference between Monomers and in  $\text{Me}_3\text{SiH}\cdots\text{Y}'$  Complexes Using NBO Analysis Calculated at the  $\omega\text{B97X-D/aug-cc-pwCVTZ}$  Level on MP2-F12/cc-pVTZ-F12 Geometries**

| $\text{Y}'$    | $\sigma/\sigma^* \text{ Si-H}$ | $\sigma/\sigma^* \text{ I-C}$  | $\Sigma \text{LP I}$  |
|----------------|--------------------------------|--------------------------------|-----------------------|
| $\text{ICF}_3$ | -0.010/0.004                   | -0.003/0.005                   | -0.005                |
| $\text{Y}'$    | $\sigma/\sigma^* \text{ Si-H}$ | $\sigma/\sigma^* \text{ Br-C}$ | $\Sigma \text{LP Br}$ |
| $\text{BrCN}$  | -0.010/0.004                   | -0.001/0.009                   | -0.002                |
| $\text{Y}'$    | $\sigma/\sigma^* \text{ Si-H}$ | $\sigma/\sigma^* \text{ H-C}$  | $\Sigma \text{LP N}$  |
| $\text{HCN}$   | -0.001/0.002                   | -0.001/0                       | -0.001                |

$\sigma$ -orbitals (the corresponding values for the  $\text{BeH}_2$  complexes are included in Table S3). Changes in orbital occupancies are related to changes in bond lengths; the biggest changes have been found for the Si–H bond. A decrease in the occupancy of the Si–H  $\sigma$ -orbitals and its increase in the case of the  $\sigma^*$ -orbitals, found for all three complexes, lead to a weakening of the Si–H bond, which is manifested by the elongation of the Si–H bond (cf. Table 1). Very similar results were obtained for the I–C, Br–C, and H–C bonds in the  $\text{ICF}_3$ ,  $\text{BrCN}$ , and  $\text{HCN}$  electron acceptors, respectively. In these cases, the elongation of the respective bonds was smaller and, in the case of  $\text{HCN}$  even, equal to zero.

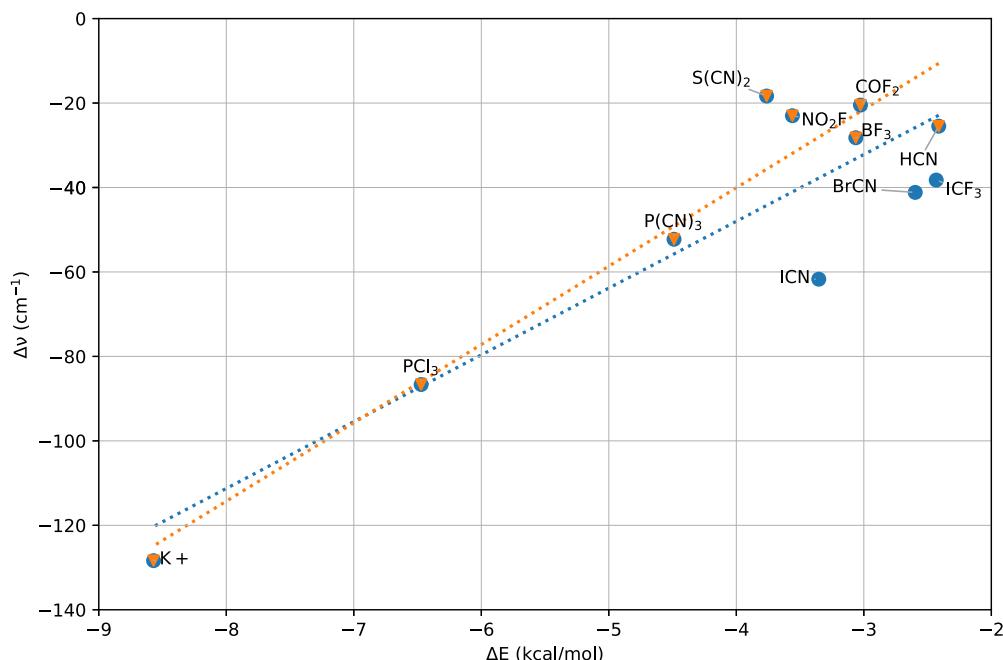
Other occupancy changes and the respective bond-length changes are less pronounced and are not discussed below.

**5.2.4. Vibration Frequencies.** Table 4 presents the experimental and calculated shifts of Si–H stretching frequencies upon complex formation. First, both experimental techniques predict larger shifts for the first two complexes; the same trend appears in all calculated shifts. Second, the calculated MP2–F12 harmonic shifts are systematically larger than the

**Table 4. IR Shift ( $\Delta\nu$  in  $\text{cm}^{-1}$ ) and Change of Intensity ( $\Delta I$  in  $\text{km/mol}$ ) of the Si–H Band upon the Formation of  $\text{Me}_3\text{SiH}\cdots\text{Y}'$  Complexes**

| $\text{Y}'$    | $\Delta\nu^{\text{a}}$                  | $\Delta I^{\text{b}}$ |
|----------------|---|-----------------------|
| $\text{ICF}_3$ | -47/-28/-31(-52)/-38(-49 <sup>c</sup> ) | +136                  |
| $\text{BrCN}$  | -37/-29/-36(-55)/-41(-50)               | +138                  |
| $\text{HCN}$   | -14/-20/-20(-35)/-25(-29)               | +37                   |

<sup>a</sup>Experimental – cold substrate/experimental – Ar-matrix//calculated anharmonic scaled (calculated anharmonic unscaled)/calculated harmonic at the MP2/cc-pwCVTZ level (calculated harmonic at the MP2-F12/cc-pVTZ-F12 level). <sup>b</sup>Calculated harmonic at the MP2/cc-pwCVTZ level. <sup>c</sup>Geometry minimum with one negative frequency.



**Figure 6.** Correlation between the total interaction energy ( $\Delta E$ ) and the shift of the Si–H stretching frequency ( $\Delta \nu$ ) for selected complexes of  $\text{Me}_3\text{SiH}$  with various electron acceptors. All values are calculated at the MP2/cc-pwCVTZ level. The blue correlation includes all electron acceptors; in the orange one, the electron acceptors exhibiting  $\sigma$ -holes (ICN, BrCN, and ICF<sub>3</sub>) have been omitted.

MP2 harmonic ones and are similar to the anharmonic unscaled ones. Finally, the agreement between the experimental and calculated frequencies is very good, which supports the reliability and suitability of all calculated techniques. In summary, the agreement between the experimental and calculated frequency shifts is satisfactory and all Si–H frequencies are redshifted upon complex formation. As expected, the magnitude of the redshifts correlates with the complex stability. Finally, the intensities of the Si–H stretch in all three complexes have increased upon complex formation. Notice here that the experimental technique adopted has not made it possible to detect intensity changes.

In the case of Be–H, the lower-energy symmetric stretching vibrations are forbidden; thus, only the antisymmetric vibrations, which are higher in energy, are discussed. As observed for the above-discussed Si–H frequencies, the intensities of the Be–H stretching frequencies (see Table S4 for the calculated harmonic Be–H stretching frequencies and their intensities) upon complexes formation also significantly increase. Frequency shifts, on the other hand, provide a different picture, with a blueshift in HBeH···BrCN, redshift in HBeH···HCN, and only negligible shift in HBeH···ICF<sub>3</sub>.

**5.2.5. Characterization and Classification of the Complexes Studied.** The classification of  $\text{Me}_3\text{Si}-\text{H}\cdots\text{Y}'$  ( $\text{Y}' = \text{ICF}_3$ , BrCN, HCN) complexes, studied both experimentally and computationally, is not unique. The first two complexes can be viewed either as CIHB or as halogen-bonded (XB) ones. Their classification is not easy because both forms are involved—the stabilization of the complex comes from Si–H···Y' CIHB as well as from the  $\text{H}^{\delta-} \rightarrow \text{I}(\text{Br})$   $\sigma$ -hole halogen bond. The dominant role of the former contribution is supported by the following evidence:

- (i) The CIHB is accompanied by a charge transfer (CT) from Y' to  $\text{Me}_3\text{Si}-\text{H}$ , whereas the CT accompanying the XB is reversed, i.e., from  $\text{Me}_3\text{Si}-\text{H}$  to Y'. Because of the

partial compensation of both contributions, the total CT between the subsystems should be small, which is fully supported by the calculated value (0.007 and 0.003 electrons transferred from  $\text{Me}_3\text{SiH}$  to BrCN and ICF<sub>3</sub>, respectively). This is further verified by the NBO E2 charge-transfer energies ( $\text{Me}_3\text{Si}-\text{H} \rightarrow \text{Y}'$  and  $\text{Y}' \rightarrow \text{Me}_3\text{Si}-\text{H}$ ), which are almost equal for both complexes. Table 3 shows that electron density in these two complexes is transferred predominantly to the  $\sigma$ - and  $\sigma^*$  Si–H orbitals, which leads to the weakening and elongation of this bond, accompanied by a redshift of the Si–H stretching frequency. In the reverse case, the electron density in the X–C  $\sigma$ - and  $\sigma^*$  orbitals of ICF<sub>3</sub> and BrCN also decreases and increases, resulting again in the weakening and elongation of the bond accompanied by a redshift of the respective stretching frequencies. The redshifts of X–C stretching frequencies are smaller than that in that in the previous case (Si–H) due to the presence of two heavy atoms (I,C and Br,C). Shifts of the Si–H bond stretching vibration frequencies are thus the most visible presentation of the complex formation.

- (ii) The strength of the XB in  $\text{Me}_3\text{Si}-\text{H}\cdots\text{Y}'$  ( $\text{Y}' = \text{ICF}_3$ , BrCN) complexes can only be estimated indirectly. The stabilization energy of the XB complexes with identical electron acceptors and a different electron donor, namely,  $\text{H}_3\text{N}\cdots\text{BrCN}$  and  $\text{H}_3\text{N}\cdots\text{ICF}_3$ , where stabilization comes exclusively from the halogen bond, is considerably higher (6.0 and 5.4 kcal/mol, respectively) than that of parent complexes with  $\text{Me}_3\text{Si}-\text{H}$  (cf. Table 2). This finding corroborates the fact that XB in  $\text{Me}_3\text{Si}-\text{H}\cdots\text{Y}'$  ( $\text{Y}' = \text{ICF}_3$ , BrCN) complexes is only weak.
- (iii) The most conclusive evidence comes from the correlation between the stabilization energy and shift of the Si–H stretching frequency upon  $\text{Me}_3\text{Si}-\text{H}\cdots\text{Y}'$  complex formation. Thirteen different electron acceptors have been considered here: first, those having pronounced  $\sigma$ - (ICF<sub>3</sub>,

BrCN, ICN, P(CN)<sub>3</sub>, PCl<sub>3</sub>, S(CN)<sub>2</sub>, and  $\pi$ - (CH<sub>3</sub>(CN)<sub>3</sub>, C(CN)<sub>6</sub>, COF<sub>2</sub>, NO<sub>2</sub>F) holes, allowing the formation of the XB; next, an electron acceptor with a p-hole (BF<sub>3</sub>); and, finally, electron acceptors with a positive charge (K<sup>+</sup>, HCN). The respective structures, stabilization energies, and shifts of the Si–H stretching frequencies are shown in Figure 1. The formation of CIHB is systematically manifested by the redshift of the Si–H stretching frequency. As expected, this shift is proportional to the stabilization energy of the complex (cf. Figure 6). The figure depicts two correlation lines: the blue one shows correlation for all electron acceptors, while in the case of the orange one, the electron acceptors exhibiting  $\sigma$ -holes (ICF<sub>3</sub>, BrCN, and ICN) have been omitted. Evidently, both correlations are very similar, which indicates that the nature of stabilization in the last three complexes is not significantly different from those complexes where XB character is absent.

Finally, the third complex for which experimental results exist, Me<sub>3</sub>Si–H...HCN, can be characterized as the CIHB or the dihydrogen one. In this case, the classification is easier because we can directly compare shifts of Si–H and C–H vibration stretching frequencies in both H-bonds as both possess hydrogen. Upon complex formation, the computed Si–H stretching frequency is redshifted (by 25 cm<sup>-1</sup>), whereas the C–H stretching frequency is blueshifted. The blueshift is, however, considerably smaller (7 cm<sup>-1</sup>). This means that the CIHB character of the complex is dominant over the dihydrogen-bonded one, which is further supported by correlation between the stabilization energy and vibration shifts, discussed in previous paragraph (Me<sub>3</sub>Si–H...HCN is very close to the ideal correlation line).

We can thus conclude that in all three complexes studied experimentally as well as computationally, the CIHB form is dominant over the XB and dihydrogen forms. This is reflected in the moderate redshift of Si–H stretching frequency, which is the largest among the stretching frequencies. These shifts are well comparable to redshifts in complexes with protonic H-bonds.

## 6. CONCLUSIONS AND REVISED DEFINITION OF H-BONDING

The complexation of Me<sub>3</sub>Si–H and HBe–H molecules (containing a hydridic hydrogen) with different electron acceptors is accompanied by red- and blueshifts of the Si–H and Be–H stretching frequencies, which are comparable with those in protonic H-bonded complexes. Both the red- or blueshifts were accompanied by an increase in intensity (the experiment did not make it possible to detect the intensity changes during complex formation), comparable to intensity changes of protonic H-bonded complexes.

An important question arises as to what to call these interactions with hydridic–hydrogen participation. Jabłoński, who studied these interactions intensely, introduced a new name for them, namely, charge-inverted H-bonding. In our opinion, this name does not capture the nature of the interaction, and “hydridic H-bonding” seems more appropriate. The wide use of this term is, however, connected with another problem. If we characterize the type of hydrogen bond by the charge on the participating hydrogen, we should consistently use “protonic H-bond” instead of “H-bond”. The undesirable inflation of new names in this case can easily be avoided by changing the current IUPAC definition of hydrogen bonding. The current IUPAC

definition requires “X (from X–H) to be more electronegative than H.” We prefer this avenue because the relevant change in the definition (mainly concerning the relative electronegativities of X and H) is only marginal and, besides hydridic and protonic H-bonding, it will also cover the dihydrogen bond. Below, we present the very first draft of this modification including a part of the original IUPAC definition, with the proposed corrections in italics and with the part of the definition to be removed underscored. The revised version covers three different interaction schemes where hydrogen plays a dominant role: X–H<sup>δ+</sup>...Y<sup>δ-</sup> (protonic H-bond), X'–H<sup>δ-</sup>...Y<sup>δ+</sup> (hydridic H-bond), and X'–H<sup>δ-</sup>...H<sup>δ+</sup>–Z (dihydrogen bond). We realize that any change of existing definition is difficult and cannot be rushed and requires an extended, broad, and in-depth discussion within the scientific community. Despite that, the benefit is obvious: the simplification and clarification of the nomenclature of one of the most important type of non-covalent interactions.

Our proposed definition of the hydrogen bond is as follows: The hydrogen bond is an attractive interaction between a hydrogen atom from a molecule or a molecular fragment X–H in which X is more or less electronegative than H and an atom or a group of atoms in the same or a different molecule in which there is evidence of bond formation. A typical hydrogen bond may be depicted as X–H...Y–Z, where the three dots denote the bond.

Depending on the electronegativity of X, the hydrogen carries a positive charge (protonic hydrogen) and thus acts as the hydrogen-bond donor (Lewis acid) or a negative charge (hydridic hydrogen) and thus acts as the hydrogen-bond acceptor (Lewis base). The atom Y in the molecule in the former case may be an electron-rich region such as, but not limited to, a lone electron pair of Y or a  $\pi$ -bonded pair of Y–Z, while in the second case, it may be an electron-deficient region such as  $\sigma$ -, n-, or  $\pi$ -hole or a positively-charged atom including hydrogen or an ion or a fragment of a molecule. X–H represents the hydrogen bond donor. The acceptor may be an atom or an anion Y, or a fragment or a molecule Y–Z, where Y is bonded to Z. In some cases, X and Y are the same. In more specific cases, X and Y are the same and X–H and Y–H distances are the same as well leading to symmetric hydrogen bonds. In any event, the acceptor is an electron rich region such as, but not limited to, a lone pair of Y or  $\pi$ -bonded pair of Y–Z.

The other parts of the original IUPAC definition remain unchanged.

## ■ ASSOCIATED CONTENT

### SI Supporting Information

The Supporting Information is available free of charge at <https://pubs.acs.org/doi/10.1021/jacs.3c00802>.

BeH<sub>2</sub> complexes, figures, and MP2-optimized geometries (PDF)

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## Notes

The authors declare no competing financial interest.

## ACKNOWLEDGMENTS

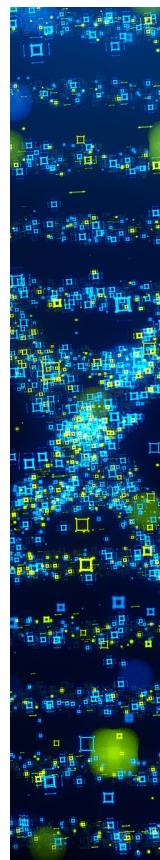
This work was supported by the ERDF/ESF “Centre of Advanced Applied Sciences” (no. CZ.02.1.01/0.0/0.0/16\_019/0000778) (S.C.); by Palacký University through the Internal Grant Association, project IGA\_PrF\_2023\_018 (M.L.); by the Ministry of Education, Youth, and Sports from the Large Infrastructures for Research, Experimental Development and Innovations project “e-Infrastructure CZ-LM2018140” (M.Š.); and by the Czech Science Foundation, project 19-27454X (P.H.).

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