

Nonlinear Schrödinger Equation with non-zero boundary conditions

1 IST

The focusing nonlinear Schrödinger equation with symmetric NZBC is

$$iq_t + q_{xx} + 2|q|^2 q = 0,$$

where $|q(x, t)| \rightarrow q_{\pm}$ as $x \rightarrow \pm\infty$, and $|q_-| = |q_+|$. Related to this PDE are two ODEs, one of which is the “scattering problem,”

$$\begin{aligned} \phi_x &= (ik\sigma_3 + Q) \phi, \\ \sigma_3 &= \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \quad Q(x, t) = \begin{pmatrix} 0 & q \\ -q^* & 0 \end{pmatrix}. \end{aligned}$$

As $x \rightarrow \pm\infty$, the solutions of the scattering problem are approximated by those of the asymptotic scattering problem $\phi_x = X_{\pm} \phi$, where $X_{\pm} = ik\sigma_3 + Q_{\pm}$ and $Q_{\pm} = \lim_{x \rightarrow \pm\infty} Q(x, t)$. The eigenvalues of X_{\pm} are

$$\lambda(k) = (q_o^2 + k^2)^{1/2}.$$

The eigenvalues are given by a complex square root, so there is a branch cut along $i[-q_o, q_o]$. We write the eigenvector matrix for X_{\pm} (dropping the dependence on t) as

$$Y_{\pm}(k) = I + i\sigma_3 Q_{\pm} / (k + \lambda(k)).$$

The Jost eigenfunctions $\phi_{\pm}(x, k)$ solve the scattering ODE and satisfy the boundary conditions

$$\phi_{\pm}(x, k) = Y_{\pm}(k) e^{i\lambda x \sigma_3} + o(1) \quad \text{as } x \rightarrow \pm\infty.$$

(The Jost eigenfunctions are actually functions of t , as well, but in what follows we only need time $t = 0$.) For all $k \in \mathbb{R} \cup i(-q_o, q_o)$ and for all $x \in \mathbb{R}$, $\phi_{-}(x, k)$ and $\phi_{+}(x, k)$ are two fundamental matrix solutions of the Lax pair, so there is some matrix $S(k)$ (the “scattering matrix”) such that

$$\phi_{+}(x, k) = \phi_{-}(x, k) S(k).$$

It turns out that the columns of ϕ_{\pm} may be extended to different regions of the complex k -plane. The entries of $S(k)$ may be written as Wronskians of columns of the ϕ_{\pm} , so some of them may be extendable as well. In particular, we will care about $s_{11}(k)$, which extends beyond $k \in \mathbb{R} \cup i(-q_o, q_o)$, because the zeros of $s_{11}(k)$ correspond to solitary waves in the solution $q(x, t)$.

2 Piecewise constant potentials

A summary of some results will be given in the case of the IVP with the following piecewise constant, box-like initial conditions:

$$(1) \quad q(x, 0) = \begin{cases} 1, & |x| > L, \\ b e^{i\alpha}, & |x| < L, \end{cases}$$

where $b > 0$. The case $0 < b < 1$ represents a potential well, the case $b > 1$ a potential barrier. The parameter α represents the phase difference between the value of the potential inside the box and its background value.

This is a convenient potential to consider, because the scattering ODE may be solved exactly, and so the Jost solutions at $t = 0$ have a simple representation in certain subsets of the real axis. Namely,

$$(2a) \quad \phi_-(x, 0, k) = \left(I + \frac{i}{k + \lambda} \sigma_3 Q_- \right) e^{i\lambda x \sigma_3} \quad x < -L,$$

$$(2b) \quad \phi_+(x, 0, k) = \left(I + \frac{i}{k + \lambda} \sigma_3 Q_+ \right) e^{i\lambda x \sigma_3} \quad x > L.$$

In addition, a fundamental matrix solution of the scattering problem in the region $|x| < L$ is

$$(2c) \quad \phi_0(x, 0, k) = \left(I + \frac{i}{k + \mu} \sigma_3 Q_0 \right) e^{i\mu x \sigma_3} \quad |x| < L,$$

where

$$(3) \quad \mu^2 = k^2 + b^2.$$

The expression of all these eigenfunctions beyond these domains is obtained by requiring continuity at the boundary. Thus, there exist matrices $S_+(k)$ and $S_-(k)$, independent of x , such that

$$(4a) \quad S_+(k) = \phi_0^{-1}(x, 0, k) \phi_+(x, 0, k),$$

$$(4b) \quad S_-(k) = \phi_0^{-1}(x, 0, k) \phi_-(x, 0, k).$$

The scattering matrix $S(k) = \phi_0^{-1}(x, 0, k) \phi_+(x, 0, k)$ is given by

$$S(k) = S_-^{-1}(k) S_+(k).$$

Evaluating (??) at $x = \pm L$ (respectively),

$$s_{1,1}(k) = e^{-2iL(\mu-\lambda)} \left[e^{4iL\mu} (b^2(k+\lambda)^2 - 2b \cos \alpha (k+\lambda)(k+\mu) + (k+\mu)^2) \right. \\ \left. + b^2 + 2b \cos \alpha (k+\lambda)(k+\mu) + (k+\lambda)^2(k+\mu)^2 \right] \\ / ((k+\lambda)^2 + 1)((k+\mu)^2 + b^2).$$

Note that $s_{1,1}(k)$ depends on $\mu(k)$, which is defined by (??) up to a sign. After some algebra,

$$(5) \quad s_{1,1}(k) = e^{2iL\lambda} \left(\cos(2L\mu) - \frac{i(b \cos \alpha + k^2)}{\lambda\mu} \sin(2L\mu) \right).$$

2.1 Results

1. If $b < 1$ and $\cos \alpha > b$ there are no discrete eigenvalues (that is, zeros of s_{11}). (So there is a class of perturbations of the constant background that don't have any solitons.)
2. If $b > 1$ and $\cos \alpha > 1/b$, there is always at least one eigenvalue. The eigenvalues are always purely imaginary, in the interval $i(1, b)$. More precisely, there are n discrete eigenvalues in $i(1, b)$ whenever $(n-1)\pi/(2\sqrt{b^2-1}) < L < n\pi/(2\sqrt{b^2-1})$, and these eigenvalues accumulate near ib (on the imaginary axis) as L increases. (This means that there are
3. If $b < 1$ and $\cos \alpha < b$, we don't have analytical results. Numerical evidence suggests there are no zeros for $L < (2n+1)\pi/(4\sqrt{b^2-b_a^2})$. As L increases, an eigenvalues appear on the continuous spectrum and accumulate near ib .

4. If $b > 1$ and $\cos \alpha < 1/b$, $s_{1,1}$ has a zero on the continuous spectrum at ib_α if $L = (2n+1)\pi/(4\sqrt{b^2 - b_\alpha^2})^3$. Numerical evidence suggests that as L increases, the eigenvalues travel towards the imaginary axis, bifurcating at points in $i(1, b)$. The eigenvalues that travel down disappear at i ; those that travel up accumulate on the imaginary axis near ib .

3 Numerical Questions

I have been wondering what any of these solutions look like. In particular,

1. What happens to $q(x, t)$ in the case when there are no solitons? That is, what does $q(x, t)$ look like if $b < 1$ and $\cos \alpha > b$?
2. Is there a difference between the $q(x, t)$ with a purely imaginary eigenvalue and the $q(x, t)$ with a complex eigenvalue, or a purely real “eigenvalue?” (Technically, if a zero of s_{11} is real, then it is on the continuous spectrum, and so is not actually an eigenvalue.)
3. What happens at the L at which eigenvalues bifurcate? That is, in a case where $L - \epsilon$ has one eigenvalue, and $L + \epsilon$ has 2 eigenvalues, what happens to $q(x, t)$ right at L ?