

2.5. Applications to Physics and Engineering

2.5.1. Work. Work is the energy produced by a force F pushing a body along a given distance d . If the force is constant, the work done is the product

$$W = F \cdot d.$$

The SI (international) unit of work is the joule (J), which is the work done by a force of one Newton (N) pushing a body along one meter (m). In the American system a unit of work is the foot-pound. Since $1 \text{ N} = 0.224809 \text{ lb}$ and $1 \text{ m} = 3.28084 \text{ ft}$, we have $1 \text{ J} = 0.737561 \text{ ft lb}$.

More generally, assume that the force is variable and depends on the position. Let $F(x)$ be the force function. Assume that the force pushes a body from a point $x = a$ to another point $x = b$. In order to find the total work done by the force we divide the interval $[a, b]$ into small subintervals $[x_{i-1}, x_i]$ so that the change of $F(x)$ is small along each subinterval. Then the work done by the force in moving the body from x_{i-1} to x_i is approximately:

$$\Delta W_i \approx F(x_i^*) \Delta x,$$

where $\Delta x = x_i - x_{i-1} = (b - a)/n$ and x_i^* is any point in $[x_{i-1}, x_i]$. So, the total work is

$$W = \sum_{i=1}^n \Delta W_i \approx \sum_{i=1}^n F(x_i^*) \Delta x.$$

As $n \rightarrow \infty$ the Riemann sum at the right converges to the following integral:

$$W = \int_a^b F(x) dx.$$

2.5.2. Elastic Springs. Consider a spring on the x -axis so that its right end is at $x = 0$ when the spring is at its rest position. According to *Hook's Law*, the force needed to stretch the spring from 0 to x is proportional to x , i.e.:

$$F(x) = kx,$$

where k is the so called *spring constant*.

The energy needed to stretch the spring from 0 to a is then the integral

$$W = \int_0^a kx dx = k \frac{a^2}{2}.$$

2.5.3. Work Done Against Gravity. According to Newton's Law, the force of gravity at a distance r from the center of the Earth is

$$F(r) = \frac{k}{r^2},$$

where k is some positive constant.

The energy needed to lift a body from a point at distance R_1 from the center of the Earth to another point at distance R_2 is given by the following integral

$$W = \int_{R_1}^{R_2} \frac{k}{r^2} dr = \left[-\frac{k}{r} \right]_{R_1}^{R_2} = k \left(\frac{1}{R_1} - \frac{1}{R_2} \right).$$

Example: Find the energy needed to lift 1000 Km a body whose weight is 1 N at the surface of Earth. The Earth radius is 6378 Km.

Answer: First we must determine the value of the constant k in this case. Since the weight of the body for $r = 6378$ Km is 1 N we have $k/6378^2 = 1$, so $k = 6378^2$. Next we have $R_1 = 6378$, $R_2 = 6378 + 1000 = 7378$, hence

$$W = 6378^2 \left(\frac{1}{6378} - \frac{1}{7378} \right) = 864.462 \text{ N Km}.$$

Since 1 Km = 1000 m, the final result in joule is

$$864.462 \text{ N Km} = 864.462 \text{ N} \times 1000 \text{ m} = 864462 \text{ J}.$$

2.5.4. Work Done Filling a Tank. Consider a tank whose bottom is at some height $y = a$ and its top is at $y = b$. Assume that the area of its cross section is $A(y)$. We fill the tank by lifting from the ground ($y = 0$) tiny layers of thickness dy each. Their weight is $dF = \rho A(y) dy$, where ρ is the density of the liquid that we are putting in the tank. The work needed to lift each layer is

$$dW = dF \cdot y = \rho y A(y) dy.$$

Hence, the work needed to fill the tank is

$$W = \int_a^b \rho y A(y) dy.$$

2.5.5. Emptying a Tank. Consider a tank like the one in the previous paragraph. Now we empty it by pumping its liquid to a fix height h . The analysis of the problem is similar to the previous paragraph, but now the work done to pump a tiny layer of thickness dy is

$$dW = dF \cdot (h - y) = \rho(h - y)A(y) dy.$$

Hence the total work needed to empty the tank is

$$W = \int_a^b \rho(h - y)A(y) dy.$$

2.5.6. Force Exerted by a Liquid Against a Vertical Wall.

The pressure p of an homogeneous liquid of density ρ at depth h is

$$p = \rho h.$$

When the pressure is constant, the force exerted by the liquid against a surface is the product of the pressure and the area of the surface. However the pressure against a vertical wall is not constant because it depends on the depth.

Assume that the surface of the liquid is at $y = c$ and we place a vertical plate of width $w(y)$ between $y = a$ and $y = b$. The force exerted at y (so at depth $h = c - y$) against a small horizontal strip of height dy and width $w(y)$ (area = $w(y) dy$) is

$$dF = \rho(c - y)w(y) dy.$$

hence the total force is

$$F = \int_a^b \rho(c - y)w(y) dy.$$

Example: A cylindrical tank of radius 1 m and full of water ($\rho = 9800 \text{ N/m}^3$) is lying on its side. What is the pressure exerted by the water on its (vertical) bottom?

Answer: We assume the center of the tank is at $y = 0$, so the top of the liquid is at $y = 1$, and its bottom is at $y = -1$. On the other hand we obtain geometrically $w(y) = 2\sqrt{1 - y^2}$. Hence the total force is:

$$F = \int_{-1}^1 9800 \cdot (1 - y)2\sqrt{1 - y^2} dy = 19800 \cdot \frac{\pi}{2} = 30787.6 \text{ N}.$$

2.5.7. Moments, Center of Mass. Here we will find a point in a solid in which it can be balanced. That point is called *center of mass*.

In the particular case when the solid consists of just two particles, we can use the Law of the Lever, discovered by Archimedes. If two masses m_1 and m_2 are attached to a rod of negligible mass on opposite sides of a fulcrum at distance d_1 , d_2 from the fulcrum, the rod will balance if

$$m_1 d_1 = m_2 d_2 .$$

If the particles lie along the x axis at coordinates x_1 and x_2 respectively (with $x_1 < x_2$), and \bar{x} is the coordinate of the center of mass, then $d_1 = \bar{x} - x_1$, $d_2 = x_2 - \bar{x}$, so

$$\begin{aligned} m_1(\bar{x} - x_1) &= m_2(x_2 - \bar{x}) \\ m_1\bar{x} + m_2\bar{x} &= m_1x_1 + m_2x_2 \\ \bar{x} &= \frac{m_1x_1 + m_2x_2}{m_1 + m_2} . \end{aligned}$$

The numbers m_1x_1 and m_2x_2 are called *moments* of the masses m_1 and m_2 . So the center of masses is found by dividing the total moment by the total mass. More generally, for a solid made up of n particles of masses m_i placed on the x -axis at points x_i :

$$\bar{x} = \frac{\sum_{i=1}^n m_i x_i}{\sum_{i=1}^n m_i} = \frac{M}{m} ,$$

where $M = \sum_{i=1}^n m_i x_i$ is the total moment of the solid, and $m = \sum_{i=1}^n m_i$ is its total mass.

If we now consider particles distributed in a plane with coordinates (x_i, y_i) , then the moment of the system about the y -axis is

$$M_y = \sum_{i=1}^n m_i x_i ,$$

and the moment of the system about the x -axis is

$$M_x = \sum_{i=1}^n m_i y_i .$$

M_y measures the tendency of the solid to rotate about the y -axis, and M_x measures the tendency of the solid to rotate about the x -axis.

The coordinates of the center of mass (\bar{x}, \bar{y}) are now

$$\bar{x} = \frac{M_y}{m}, \quad \bar{y} = \frac{M_x}{m}.$$

Now consider a flat plate of uniform density ρ occupying the plane region under a curve $y = f(x)$ between $x = a$ and $x = b$. Slices parallel to the y axis of width Δx and length $f(x)$ have mass $\rho f(\bar{x}) \Delta x$, and their moment about the y -axis will be $x \rho f(x) \Delta x$. The sum of their moments in the limit as $\Delta x \rightarrow 0$ is the integral:

$$M_y = \rho \int_a^b x f(x) dx.$$

The moment about the x -axis of a slice can be found taking into account that by symmetry its center of mass is at distance $f(x)/2$ from the x -axis, so the moment is $\frac{1}{2} f(x) \rho f(\bar{x}) \Delta x = \rho \frac{1}{2} [f(x)]^2 \Delta x$. Adding and taking the limit we get the integral

$$M_x = \rho \int_a^b \frac{1}{2} [f(x)]^2 dx.$$

The total mass of the plate is its density times its area:

$$m = \rho \int_a^b f(x) dx.$$

Hence, the coordinates of its center of mass are:

$$\begin{aligned} \bar{x} &= \frac{M_y}{m} = \frac{\rho \int_a^b x f(x) dx}{\rho \int_a^b f(x) dx} = \frac{\int_a^b x f(x) dx}{\int_a^b f(x) dx} \\ \bar{y} &= \frac{M_x}{m} = \frac{\rho \int_a^b \frac{1}{2} [f(x)]^2 dx}{\rho \int_a^b f(x) dx} = \frac{\int_a^b \frac{1}{2} [f(x)]^2 dx}{\int_a^b f(x) dx}. \end{aligned}$$

Example: Find the center of mass of a semicircular plate of radius r .

Answer: We use coordinates so that the plate occupies the region under the graph of $y = \sqrt{r^2 - x^2}$, $-r \leq x \leq r$. The area of the semicircle is $A = \pi r^2/2$. By symmetry $\bar{x} = 0$, so we only need to

find \bar{y} :

$$\begin{aligned}\bar{y} &= \frac{1}{A} \int_{-r}^r \frac{1}{2} \{f(x)\}^2 dx = \frac{1}{\pi r^2/2} \cdot \frac{1}{2} \int_{-r}^r (\sqrt{r^2 - x^2})^2 dx \\ &= \frac{2}{\pi r^2} \int_0^r (r^2 - x^2) dx = \frac{2}{\pi r^2} \left[r^2 x - \frac{x^3}{3} \right]_0^r dx \\ &= \frac{2}{\pi r^2} \left(r^3 - \frac{r^3}{3} \right) = \frac{2}{\pi r^2} \frac{2r^3}{3} = \boxed{\frac{4r}{3\pi}}.\end{aligned}$$