

Transport Of Density In Stars

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ABSTRACT

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Key words: turbulence – convection zone – stellar evolution

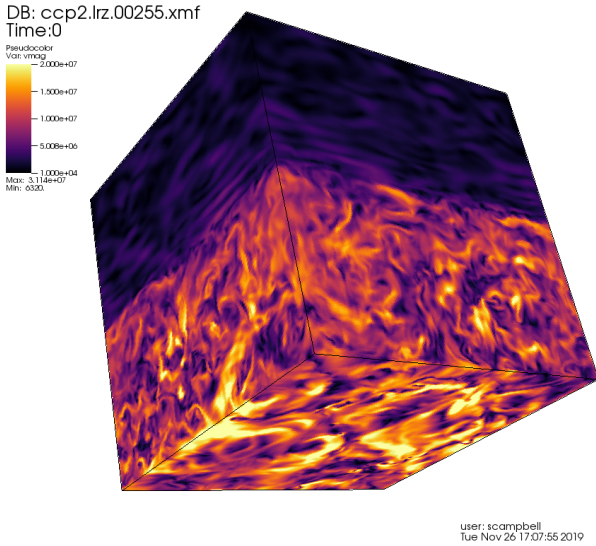


Figure 1. Visualization of turbulence in a 3D box under study.

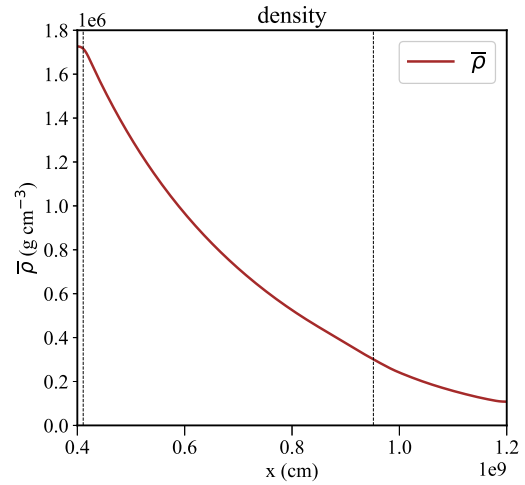


Figure 2. Density radial profile of the 3D model.

1 INTRODUCTION

What does RANS transport equation for density (Mocák et al. 2014) tell us about convection zones in stars?

Time-dependency $\partial_t \bar{\rho}$, compressibility $\bar{\rho} \nabla_x \tilde{u}_x$ and non-local physics $\nabla_x f_\rho$ play important role in convection zone.

The following relation between the dilatation terms and non-local physics can be deduced from Fig.3:

$$-\bar{\rho} \nabla_x \tilde{u}_x \sim -\nabla_x f_\rho + f_\rho / \bar{\rho} \partial_r \bar{\rho} - \bar{\rho} \nabla_x \tilde{u}_x \quad (1)$$

Linearized continuity equation leads to incomplete picture of density physics in stars (see Sect.2).

2 LINEARIZED CONTINUITY EQUATION

$$\partial_t \rho + \nabla \cdot \rho \mathbf{u} = 0 \quad (2)$$

Let us assume, that $\rho(\mathbf{r}, t) = \rho_0(\mathbf{r}) + \rho'(\mathbf{r}, t)$ where $\rho_0(\mathbf{r})$ is time-independent background density state around which we'll linearize the continuity equation. We get:

$$\partial_t [\rho_0(\mathbf{r}) + \rho'(\mathbf{r}, t)] + \nabla \cdot [\rho_0(\mathbf{r}) + \rho'(\mathbf{r}, t)] \mathbf{u}(\mathbf{r}, t) = 0 \quad (3)$$

$$\cancel{\partial_t \rho_0(\mathbf{r})} + \partial_t \rho'(\mathbf{r}, t) + \nabla \cdot [\rho_0(\mathbf{r}) \mathbf{u}(\mathbf{r}, t)] + \nabla \cdot [\rho'(\mathbf{r}, t) \mathbf{u}(\mathbf{r}, t)] = 0 \quad (4)$$

And because we also assume that $\rho' \ll \rho_0$, we have that $\nabla \cdot [\rho_0(\mathbf{r}) \mathbf{u}(\mathbf{r}, t)] \gg \nabla \cdot [\rho'(\mathbf{r}, t) \mathbf{u}(\mathbf{r}, t)]$ and get:

$$\partial_t \rho' + \nabla \cdot (\rho_0 \mathbf{u}) \sim 0 \quad (5)$$

Furthermore, we know that $\mathbf{u} \sim \mathbf{u}'$ and therefore we can write:

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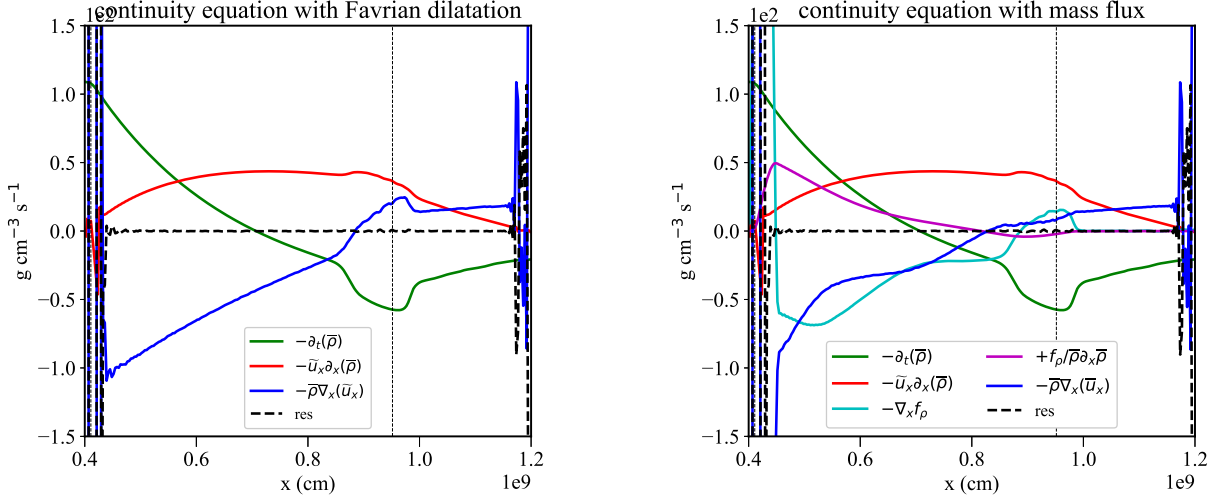


Figure 3. Transport density equation with Left: Favrian dilatation $\nabla_r \tilde{u}_x$ and Right: turbulent mass flux f_ρ

$$\partial_t \rho' + \nabla \cdot (\rho_0 \mathbf{u}') \sim 0 \quad (6)$$

This is equation 23 from (Viallet et al. 2013).

3 MEAN LINEARIZED CONTINUITY EQUATION

$$\partial_t \rho' + \nabla \cdot (\rho_0 \mathbf{u}) \sim 0 \quad (7)$$

After space-time averaging of this linearized continuity equation, we get:

$$\overline{\partial_t \rho'} + \overline{\nabla \cdot (\rho_0 \mathbf{u})} \sim 0 \quad (8)$$

$$\overline{\partial_t \rho'} + \overline{\nabla \cdot (\rho_0 \mathbf{u})} \sim 0 \quad (9)$$

$$\nabla_r \rho_0 \bar{u}_r \sim 0 \quad (10)$$

$$(11)$$

But because $\rho_0 \equiv \bar{\rho}$, we can write:

$$\nabla_r \bar{\rho} \bar{u}_r \sim 0 \quad (12)$$

$$\bar{\rho} \nabla_r \bar{u}_r + \bar{u}_r \partial_r \bar{\rho} \sim 0 \quad (13)$$

$$\bar{\rho} \nabla_r \bar{u}_r + (\bar{u}_r'' + \tilde{u}_r) \partial_r \bar{\rho} \sim 0 \quad (14)$$

$$\bar{\rho} \nabla_r \bar{u}_r + \bar{u}_r'' \partial_r \bar{\rho} + \tilde{u}_r \partial_r \bar{\rho} \sim 0 \quad (15)$$

$$\bar{\rho} \nabla_r \bar{u}_r + \bar{\rho}' \bar{u}_r' / \bar{\rho} \partial_r \bar{\rho} + \tilde{u}_r \partial_r \bar{\rho} \sim 0 \quad (16)$$

$$\bar{\rho} \nabla_r \bar{u}_r - f_\rho / \bar{\rho} \partial_r \bar{\rho} + \tilde{u}_r \partial_r \bar{\rho} \sim 0 \quad (17)$$

But from my latest analysis of the full (non-linearized) continuity equation, it turns out that:

$$\bar{\rho} \nabla_r \bar{u}_r - f_\rho / \bar{\rho} \partial_r \bar{\rho} + \tilde{u}_r \partial_r \bar{\rho} \neq 0 \quad (18)$$

Or, when multiply this by -1:

$$-\bar{\rho} \nabla_r \bar{u}_r + f_\rho / \bar{\rho} \partial_r \bar{\rho} - \tilde{u}_r \partial_r \bar{\rho} \neq 0 \quad (19)$$

Linearization relies on time-independent background state around which you can linearize, which turns out not to be our case and $\partial_t \rho_0$ or in other words $\partial_t \bar{\rho} \neq 0$ (green curve). Also it gets rid of the transport of turbulent density field $\nabla_r f_\rho$ (cyan curve) (Fig.3):

REFERENCES

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Viallet M., Meakin C., Arnett D., Mocák M., 2013, *ApJ*, **769**, 1

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