

rans(eXtreme) Analysis Framework for Compressible Hydrodynamic Simulations

Miroslav Mocák,^{1*}

¹Monash Centre for Astrophysics, School of Physics and Astronomy, Monash University, Clayton, Australia 3800

in preparation for MNRAS

ABSTRACT

ransX ..

Key words: turbulence – analysis – RANS

1 INTRODUCTION

Analysis of hydrodynamic simulations was in past focused either on calculation of simple averaged quantities (like rms vel, Fourier spectra etc.) (citations), mathematical exploration of RANS equations only (citations), or few individual flux calculations (citations). Complete RANS analysis has only been done by Meakin, Viallet, Mocak, Campbell, Hirchi, Arnett. The presented RANS framework called ransX¹ is supposed to promote and streamline such complete RANS analysis of 3D hydrodynamic fully compressible multi-elements simulations of turbulence in stars (no rotation, no magnetic fields).

We obtain our 1D RANS equations by introducing two types of averaging: statistical averaging and horizontal averaging (Besnard et al. 1992; Viallet et al. 2013). In practice, statistical averages are computed by performing a time average (the ergodic hypothesis). Therefore, the combined average of a quantity q is defined as

$$\bar{q}(r, t) = \frac{1}{T\Delta\Omega} \int_{t-T/2}^{t+T/2} q(r, \theta, \phi, t') d\Omega dt' \quad (1)$$

where $d\Omega = \sin\theta d\theta d\phi$ is the solid angle in spherical coordinates, T is the averaging time period, and $\Delta\Omega$ is total solid angle being averaged over.

The flow variables are then decomposed into mean and fluctuation $q = \bar{q} + q'$, noting that $\bar{q'} = 0$ by construction. Similarly, we introduce Favre (or density weighted) averaged quantities by

$$\tilde{q} = \frac{\bar{\rho q}}{\bar{\rho}} \quad (2)$$

which defines a complimentary decomposition of the flow into mean and fluctuations according to $q = \tilde{q} + q''$. Here, q'' is the Favrian fluctuation and its mean is zero when Favre averaged $\tilde{q''} = 0$.

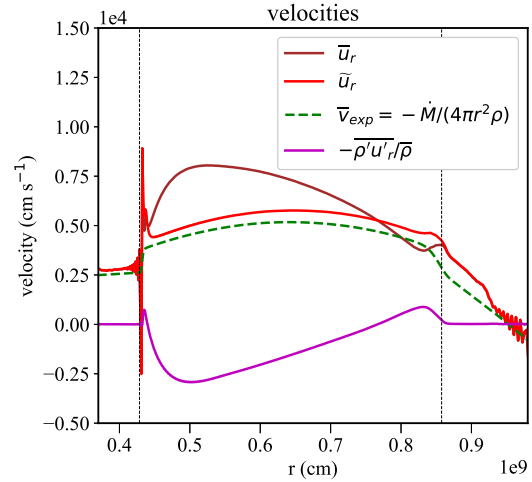


Figure 1. Reynolds and Favre velocities derived from 3D oxygen burning shell model.

Such decomposition gives us the advantage not only to understand compressible hydrodynamics better but to look at some traditional Reynolds mean fields from a new perspective. For example, it allows us to decompose mean Reynolds velocity \bar{u}_r into its background expansion ($\tilde{u}_r = v_{exp}$ how do we prove this mathematically?) and density scaled turbulent mass flux $-\rho'u'_r/\bar{\rho} = u''_r$ (Fig.1).

$$\bar{u}_r = \tilde{u}_r - \overline{u''_r} \quad (3)$$

For a more complete elaboration on the algebra of these averaging procedures we refer the reader to Chassaing et al. (2010).

2 GENERAL FORM OF RANS EQUATIONS

The RANS equations terms can be split in general into:

* E-mail:miroslav.mocak@gmail.com

¹ <https://github.com/mmicromegas/ransX>

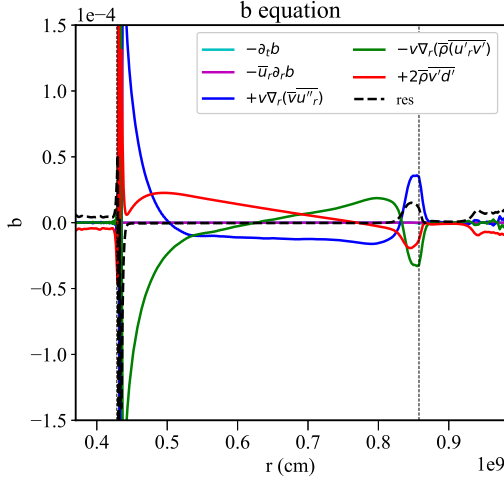


Figure 2. Example RANS equation for density-specific volume covariance.

- time-dependence \mathcal{T}
- redistribution (the div terms) \mathcal{R}
- production/destruction \mathcal{P}, \mathcal{D}
- geometry terms \mathcal{G}
- numerical residuals \mathcal{N}

General form of RANS equations is:

$$\mathcal{T} = \mathcal{R} + \mathcal{P} + \mathcal{D} + \mathcal{G} + \mathcal{N} \quad (4)$$

3 TIME-DEPENDENCE TERMS

4 REDISTRIBUTION TERMS

5 PRODUCTION/DESTRUCTION TERMS

6 GEOMETRY TERMS

7 NUMERICAL RESIDUALS

8 SOFTWARE IMPLEMENTATION

All presented equations (including their Cartesian geometry equivalents) were implemented into a software framework done in Python (<https://github.com/mmimromegas/ransX>), which relies on a specific space-time averaged field output currently programmed into hydrodynamic PROMPI code only, but which can easily be implemented in any other hydrodynamic code too. Using the ransX framework is then just a matter of working with the specific hydrodynamic output.

REFERENCES

- Besnard D., Harlow F., Rauen Zahn R., Zemach C., 1992, Technical report, Turbulence transport equations for variable-density turbulence and their relationship to two-field models. Los Alamos National Lab., NM (United States)
- Chassaing P., Antonia R., Anselmet F., Joly L., Sarkar R., 2010, Variable Density Fluid Turbulence. Kluwer Academic Publishers

Viallet M., Meakin C., Arnett D., Mocák M., 2013, *ApJ*, **769**, 1

This paper has been typeset from a $\text{\TeX}/\text{\LaTeX}$ file prepared by the author.

[t]

Table 1. 1D RANS equations in Lagrangian form.

$$\tilde{D}_t \bar{\rho} = -\bar{\rho} \tilde{d} + \mathcal{N}_\rho \quad (1)$$

$$\bar{\rho} \tilde{D}_t \tilde{u}_r = -\nabla_r \tilde{R}_{rr} - \overline{G_r^M} - \partial_r \bar{P} + \bar{\rho} \tilde{g}_r + \mathcal{N}_{ur} \quad (2)$$

$$\bar{\rho} \tilde{D}_t \tilde{u}_\theta = -\nabla_r \tilde{R}_{\theta r} - \overline{G_\theta^M} - (1/r) \partial_\theta \bar{P} + \mathcal{N}_{u\theta} \quad (3)$$

$$\bar{\rho} \tilde{D}_t \tilde{u}_\phi = -\nabla_r \tilde{R}_{\phi r} - \overline{G_\phi^M} + \mathcal{N}_{u\phi} \quad (4)$$

$$\bar{\rho} \tilde{D}_t \tilde{\epsilon}_I = -\nabla_r (f_I + f_T) - \bar{P} \bar{d} - W_P + \mathcal{S} + \mathcal{N}_{\epsilon I} \quad (5)$$

$$\bar{\rho} \tilde{D}_t \tilde{\epsilon}_K = -\nabla_r (f_k + f_P) - \tilde{R}_{ir} \partial_r \tilde{u}_i + W_b + W_P + \bar{\rho} \tilde{D}_t (\tilde{u}_i \tilde{u}_i / 2) + \mathcal{N}_{\epsilon k} \quad (6)$$

$$\bar{\rho} \tilde{D}_t \tilde{\epsilon}_t = -\nabla_r (f_I + f_T + f_k + f_P) - \bar{P} \bar{d} - \tilde{R}_{ir} \partial_r \tilde{u}_i + W_b + \mathcal{S} + \bar{\rho} \tilde{D}_t (\tilde{u}_i \tilde{u}_i / 2) + \mathcal{N}_{\epsilon t} \quad (7)$$

$$\bar{\rho} \tilde{D}_t \tilde{X}_\alpha = -\nabla_r f_\alpha + \bar{\rho} \tilde{X}_\alpha^{\text{nuc}} + \mathcal{N}_\alpha \quad (8)$$

$$\bar{\rho} \tilde{D}_t \tilde{s} = -\nabla_r f_s + (\nabla \cdot F_T) / T + \mathcal{S} / T + \mathcal{N}_s \quad (9)$$

$$\bar{\rho} \tilde{D}_t \tilde{A} = -\nabla_r f_A - \overline{\rho A^2 \Sigma_\alpha (\dot{X}_\alpha^{\text{nuc}} / A_\alpha)} + \mathcal{N}_A \quad (10)$$

$$\bar{\rho} \tilde{D}_t \tilde{Z} = -\nabla_r f_Z - \overline{\rho Z A \Sigma_\alpha (\dot{X}_\alpha^{\text{nuc}} / A_\alpha)} + \overline{\rho A \Sigma_\alpha (Z_\alpha \dot{X}_\alpha^{\text{nuc}} / A_\alpha)} + \mathcal{N}_Z \quad (11)$$

$$\bar{D}_t \bar{P} = -\nabla_r f_P - \Gamma_1 \bar{P} \bar{d} + (1 - \Gamma_1) W_P + (\Gamma_3 - 1) \mathcal{S} + (\Gamma_3 - 1) \nabla_r f_T + \mathcal{N}_P \quad (12)$$

$$\bar{D}_t \bar{T} = -\nabla_r f_T + (1 - \Gamma_3) \bar{T} \bar{d} + (2 - \Gamma_3) \bar{T}' \bar{d}' + (\nabla \cdot F_T) / \rho c_v + (\tau_{ij} \partial_i u_j) / \rho c_v + \epsilon_{\text{nuc}} / c_v + \mathcal{N}_T \quad (13)$$

$$\bar{\rho} \tilde{D}_t \tilde{k} = -\nabla_r (f_k + f_P) - \tilde{R}_{ir} \partial_r \tilde{u}_i + W_b + W_P + \mathcal{N}_k \quad (14)$$

$$\bar{\rho} \tilde{D}_t \overline{u_r''} = -(\overline{\rho' u_r' u_r' / \bar{\rho}}) \partial_r \bar{\rho} + (\tilde{R}_{rr} / \bar{\rho}) / \partial_r \bar{\rho} - \bar{\rho} \nabla_r (\overline{u_r'' u_r''}) + \nabla_r \overline{\rho' u_r' u_r'} - \overline{\rho u_r'' \nabla_r u_r} + \overline{\rho u_r' d''} - b \partial_r \bar{P} + \overline{\rho' v \partial_r P'} + \mathcal{G}_a + \mathcal{N}_a \quad (15)$$

$$\bar{D}_t b = + \bar{v} \nabla_r \overline{\rho u_r''} - \bar{\rho} \nabla_r (\overline{u_r' v'}) + 2 \bar{\rho} \overline{v' d'} + \mathcal{N}_b \quad (16)$$

$$\bar{\rho} \tilde{D}_t (f_I / \bar{\rho}) = -\nabla_r f_I^r - f_I \partial_r \tilde{u}_r - \tilde{R}_{rr} \partial_r \tilde{\epsilon}_I - \overline{\epsilon_I'' \partial_r \bar{P}} - \overline{\epsilon_I'' \partial_r P'} - \overline{u_r'' (P d)} + \overline{u_r'' (\mathcal{S} + \nabla \cdot F_T)} + \mathcal{G}_I + \mathcal{N}_{fI} \quad (17)$$

$$\bar{\rho} \tilde{D}_t (f_s / \bar{\rho}) = -\nabla_r f_s^r - f_s \partial_r \tilde{u}_r - \tilde{R}_{rr} \partial_r \tilde{s} - \overline{s'' \partial_r \bar{P}} - \overline{s'' \partial_r P'} + \overline{u_r'' (\mathcal{S} + \nabla \cdot F_T) / T} + \mathcal{G}_s + \mathcal{N}_{fs} \quad (18)$$

$$\bar{\rho} \tilde{D}_t (f_\alpha / \bar{\rho}) = -\nabla_r f_\alpha^r - f_\alpha \partial_r \tilde{u}_r - \tilde{R}_{rr} \partial_r \tilde{X}_\alpha - \overline{X_\alpha'' \partial_r \bar{P}} - \overline{X_\alpha'' \partial_r P'} + \overline{u_r'' \rho \dot{X}_\alpha^{\text{nuc}}} + \mathcal{G}_\alpha + \mathcal{N}_{f\alpha} \quad (19)$$

$$\bar{\rho} \tilde{D}_t (f_A / \bar{\rho}) = -\nabla_r f_A^r - f_A \partial_r \tilde{u}_r - \tilde{R}_{rr} \partial_r \tilde{A} - \overline{A'' \partial_r \bar{P}} - \overline{A'' \partial_r P'} - \overline{u_r'' \rho A^2 \Sigma_\alpha (\dot{X}_\alpha^{\text{nuc}} / A_\alpha)} + \mathcal{G}_A + \mathcal{N}_{fA} \quad (20)$$

$$\bar{\rho} \tilde{D}_t (f_Z / \bar{\rho}) = -\nabla_r f_Z^r - f_Z \partial_r \tilde{u}_r - \tilde{R}_{rr} \partial_r \tilde{Z} - \overline{Z'' \partial_r \bar{P}} - \overline{Z'' \partial_r P'} - \overline{u_r'' \rho Z A \Sigma_\alpha (\dot{X}_\alpha^{\text{nuc}} / A_\alpha)} - \overline{u_r'' \rho A \Sigma_\alpha (Z_\alpha \dot{X}_\alpha^{\text{nuc}} / A_\alpha)} + \mathcal{G}_Z + \mathcal{N}_{fZ} \quad (21)$$

$$\tilde{D}_t f_{pr} = -\nabla_r f_{pr}^r - f_{pr} \partial_r \tilde{u}_r + \overline{u_r'' u_r'' \partial_r \bar{P}} + \Gamma_1 \overline{u_r'' P d} + (\Gamma_3 - 1) \overline{u_r'' \rho \epsilon_{\text{nuc}}} + \overline{P' u_r'' d''} - \overline{P' G_r^M / \rho} - \overline{P' \partial_r P / \rho} + \mathcal{N}_{fpr} \quad (22)$$

$$\bar{\rho} \tilde{D}_t (f_h / \bar{\rho}) = -\nabla_r f_h^r - f_h \partial_r \tilde{u}_r - \tilde{R}_{rr} \partial_r \tilde{h} - \overline{h'' \partial_r \bar{P}} - \overline{h'' \partial_r P'} - \Gamma_1 \overline{u_r'' (P d)} + \Gamma_3 \overline{u_r'' (\mathcal{S} + \nabla \cdot F_T)} + \mathcal{G}_h + \mathcal{N}_h \quad (23)$$

$$\bar{\rho} \tilde{D}_t (f_s / \bar{\rho}) = -\nabla_r f_s^r - f_s \partial_r \tilde{u}_r - \tilde{R}_{rr} \partial_r \tilde{s} - \overline{s'' \partial_r \bar{P}} - \overline{s'' \partial_r P'} + \overline{u_r'' (\mathcal{S} + \nabla \cdot F_T) / T} + \mathcal{G}_s + \mathcal{N}_{fs} \quad (24)$$

$$\bar{\rho} \tilde{D}_t \sigma_\alpha = -\nabla_r (\overline{\rho X_\alpha'' X_\alpha'' u_r''}) - 2 f_\alpha \partial_r \tilde{X}_\alpha + \overline{2 X_\alpha'' \rho \dot{X}_\alpha^{\text{nuc}}} + \mathcal{N}_{\sigma_\alpha} \quad (25)$$

Definitions:

ρ density	g_r gravitational acceleration
T temperature	$\mathcal{S} = \bar{\rho}\epsilon_{\text{nuc}}(q)$ nuclear energy production (cooling function)
P pressure	$\tau_{ij} = 2\mu S_{ij}$ viscous stress tensor (μ kinematic viscosity)
u_r, u_θ, u_ϕ velocity components	$S_{ij} = (1/2)(\partial_i u_j + \partial_j u_i)$ strain rate
\mathbf{u} velocity	$\tilde{R}_{ij} = \bar{\rho} \widetilde{u_i'' u_j''}$ Reynolds stress tensor
$d = \nabla \cdot \mathbf{u}$ dilatation	$\tilde{k} = (1/2) \widetilde{u_i'' u_i''}$ turbulent kinetic energy
ϵ_I specific internal energy	$\tilde{k}^r = (1/2) \widetilde{u_r'' u_r''} = (1/2) \tilde{R}_{rr} / \bar{\rho}$ radial turbulent kinetic energy
$\epsilon_K = (1/2) u_i u_i$ specific kinetic energy	$\tilde{k}^\theta = (1/2) \widetilde{u_\theta'' u_\theta''} = (1/2) \tilde{R}_{\theta\theta} / \bar{\rho}$ angular turbulent kinetic energy
ϵ_t specific total energy	$\tilde{k}^\phi = (1/2) \widetilde{u_\phi'' u_\phi''} = (1/2) \tilde{R}_{\phi\phi} / \bar{\rho}$ angular turbulent kinetic energy
s specific entropy	$\tilde{k}^h = \tilde{k}^\theta + \tilde{k}^\phi$ horizonatal turbulent kinetic energy
$v = 1/\rho$ specific volume	$f_k = (1/2) \bar{\rho} \widetilde{u_i'' u_i'' u_r''}$ turbulent kinetic energy flux
X_α mass fraction of isotope α	$f_k^r = (1/2) \bar{\rho} \widetilde{u_r'' u_i'' u_i''}$ radial turbulent kinetic energy flux
$\dot{X}_\alpha^{\text{nuc}}$ rate of change of X_α	$f_k^\theta = (1/2) \bar{\rho} \widetilde{u_\theta'' u_i'' u_i''}$ angular turbulent kinetic energy flux
A_α number of nucleons in isotope α	$f_k^\phi = (1/2) \bar{\rho} \widetilde{u_\phi'' u_i'' u_i''}$ angular turbulent kinetic energy flux
Z_α charge of isotope α	$f_k^h = f_k^\theta + f_k^\phi$ horizontal turbulent kinetic energy flux
A mean number of nucleons per isotope	$W_p = \overline{P' d''}$ turbulent pressure dilatation
Z mean charge per isotope	$W_b = \bar{\rho} \widetilde{u_r'' g_r}$ buoyancy energy

Table 2. Definitions (continued):

$f_P = \overline{P'u_r'}$ acoustic flux	$F_T = \chi \partial_r T$ full heat flux (χ thermal conductivity)
$f_T = \overline{\chi \partial_r T}$ turbulent heat flux	
$f_I = \overline{\rho \epsilon_I'' u_r''}$ internal energy flux	$f_\alpha = \overline{\rho X_\alpha'' u_r''}$ X_α flux
$f_s = \overline{\rho s'' u_r''}$ specific entropy flux	$b = \overline{v' \rho'}$ density-specific volume covariance
$f_A = \overline{\rho A'' u_r''}$ A flux	\mathcal{N}_ρ numerical effect
$f_Z = \overline{\rho Z'' u_r''}$ Z flux	$\mathcal{N}_{ur}, \mathcal{N}_{u\theta}, \mathcal{N}_{u\phi}$ numerical effect
$f_\tau = f_\tau^r + f_\tau^\theta + f_\tau^\phi$ viscous flux	$\mathcal{N}_{\epsilon I} = +\varepsilon_k$ numerical effect
$f_\tau^r = -\overline{\tau_{rr}' u_r'}$ viscous flux	$\mathcal{N}_{\epsilon k} = -\varepsilon_k$ numerical effect
$f_\tau^\theta = -\overline{\tau_{\theta r}' u_\theta'}$ viscous flux	$\mathcal{N}_{\epsilon t}$ numerical effect
$f_\tau^\phi = -\overline{\tau_{\phi r}' u_\phi'}$ viscous flux	\mathcal{N}_α numerical effect
$f_\tau^h = f_\tau^r + f_\tau^\theta + f_\tau^\phi$ viscous flux	$\mathcal{N}_s = \overline{-\varepsilon_k/T}$ numerical effect
$f_I^r = \overline{\rho \epsilon_I'' u_r'' u_r''}$ radial flux of f_I	\mathcal{N}_A numerical effect
$f_s^r = \overline{\rho s'' u_r'' u_r''}$ radial flux of f_s	\mathcal{N}_Z numerical effect
$f_\alpha^r = \overline{\rho \alpha'' u_r'' u_r''}$ radial flux of f_α	$\mathcal{N}_k = -\nabla_r f_\tau - \varepsilon_k$ numerical effect
$f_A^r = \overline{\rho A'' u_r'' u_r''}$ radial flux of f_A	$\mathcal{N}_{kr} = -\nabla_r f_\tau^r - \varepsilon_k^r$ numerical effect
$f_Z^r = \overline{\rho Z'' u_r'' u_r''}$ radial flux of f_Z	$\mathcal{N}_{kh} = -\nabla_r f_\tau^h - \varepsilon_k^h$ numerical effect
$\mathcal{G}_k^r = -(1/2) \overline{G_{rr}^R} - \overline{u_r'' G_r^M}$	$\mathcal{N}_a = -\varepsilon_a$ numerical effect
$\mathcal{G}_k^\theta = -(1/2) \overline{G_{\theta\theta}^R} - \overline{u_\theta'' G_\theta^M}$	\mathcal{N}_b numerical effect
$\mathcal{G}_k^\phi = -(1/2) \overline{G_{\phi\phi}^R} - \overline{u_\phi'' G_\phi^M}$	$\mathcal{N}_{fI} = -\nabla_r (\epsilon_I'' \tau_{rr}') + \overline{u_r'' \tau_{ij} \partial_i u_j} - \varepsilon_I$ numerical effect
$\mathcal{G}_k^h = +\mathcal{G}_k^r + \mathcal{G}_k^\theta + \mathcal{G}_k^\phi$	$\mathcal{N}_{fs} = -\nabla_r (s'' \tau_{rr}') + \overline{u_r'' \tau_{ij} \partial_i u_j / T} - \varepsilon_s$ numerical effect
$\mathcal{G}_a = +\rho' v \overline{G_r^M}$	$\mathcal{N}_{f\alpha} = -\nabla_r (\overline{X_\alpha'' \tau_{rr}'}) - \varepsilon_\alpha$ numerical effect
$\mathcal{G}_I = -\overline{G_r^I} - \overline{\epsilon_I'' G_r^M}$	$\mathcal{N}_{fA} = -\nabla_r (\overline{A'' \tau_{rr}'}) - \varepsilon_A$ numerical effect
$\mathcal{G}_\alpha = -\overline{G_r^\alpha} - \overline{\epsilon_\alpha'' G_r^M}$	$\mathcal{N}_{fZ} = -\nabla_r (\overline{Z'' \tau_{rr}'}) - \varepsilon_Z$ numerical effect
$\mathcal{G}_A = -\overline{G_r^A} - \overline{A'' G_r^M}$	$\mathcal{G}_Z = -\overline{G_r^Z} - \overline{Z'' G_r^M}$
$\varepsilon_k^r = \overline{\tau_{rr}' \partial_r u_r''} + \overline{\tau_{r\theta}' (1/r) \partial_\theta u_r''} + \overline{\tau_{r\phi}' (1/r \sin \theta) \partial_\phi u_r''}$	$\overline{G_r^M} = -\overline{\rho u_\theta u_\theta / r} - \overline{\rho u_\phi u_\phi / r}$
$\varepsilon_k^\theta = \overline{\tau_{\theta r}' \partial_r u_\theta''} + \overline{\tau_{\theta\theta}' (1/r) \partial_\theta u_\theta''} + \overline{\tau_{\theta\phi}' (1/r \sin \theta) \partial_\phi u_\theta''}$	$\overline{G_\theta^M} = +\overline{\rho u_\theta u_r / r} - \overline{\rho u_\phi u_\phi / (r \tan \theta)}$
$\varepsilon_k^\phi = \overline{\tau_{\phi r}' \partial_r u_\phi''} + \overline{\tau_{\phi\theta}' (1/r) \partial_\theta u_\phi''} + \overline{\tau_{\phi\phi}' (1/r \sin \theta) \partial_\phi u_\phi''}$	$\overline{G_\phi^M} = +\overline{\rho u_\phi u_r / r} + \overline{\rho u_\phi u_\theta / (r \tan \theta)}$
$\varepsilon_k = (1/2)(\varepsilon_k^r + \varepsilon_k^\theta + \varepsilon_k^\phi)$	$\overline{G_{rr}^R} = -\overline{\rho u_\theta'' u_\theta'' / r} - \overline{\rho u_\theta'' u_r'' / r} - \overline{\rho u_\phi'' u_\phi'' / r} - \overline{\rho u_\phi'' u_r'' / r}$
$\varepsilon_k^h = (1/2)(\varepsilon_k^r + \varepsilon_k^\theta + \varepsilon_k^\phi)$	$\overline{G_{\theta\theta}^R} = +\overline{\rho u_\theta'' u_\theta'' / r} + \overline{\rho u_\theta'' u_r'' / r} - \overline{\rho u_\phi'' u_\phi'' / (r \tan \theta)} - \overline{\rho u_\phi'' u_r'' / (r \tan \theta)}$
$\varepsilon_a = \overline{\rho' v \nabla_r \tau_{rr}'}$	$\overline{G_{\phi\phi}^R} = +\overline{\rho u_\phi'' u_\phi'' / r} + \overline{\rho u_\phi'' u_\theta'' / (r \tan \theta)} + \overline{\rho u_\phi'' u_r'' / r} + \overline{\rho u_\phi'' u_\theta'' / (r \tan \theta)}$
$\varepsilon_I = \overline{\tau_{rr}' \partial_r \epsilon_I''} + \overline{\tau_{r\theta}' (1/r) \partial_\theta \epsilon_I''} + \overline{\tau_{r\phi}' (1/r \sin \theta) \partial_\phi \epsilon_I''}$	$\overline{G_r^I} = -\overline{\rho \epsilon_I'' u_\theta'' / r} - \overline{\rho \epsilon_I'' u_\phi'' / r}$
$\varepsilon_s = \overline{\tau_{rr}' \partial_r s''} + \overline{\tau_{r\theta}' (1/r) \partial_\theta s''} + \overline{\tau_{r\phi}' (1/r \sin \theta) \partial_\phi s''}$	$\overline{G_r^s} = -\overline{\rho s'' u_\theta'' / r} - \overline{\rho s'' u_\phi'' / r}$
$\varepsilon_\alpha = \overline{\tau_{rr}' \partial_r X_\alpha''} + \overline{\tau_{r\theta}' (1/r) \partial_\theta X_\alpha''} + \overline{\tau_{r\phi}' (1/r \sin \theta) \partial_\phi X_\alpha''}$	$\overline{G_r^\alpha} = -\overline{\rho X_\alpha'' u_\theta'' / r} - \overline{\rho X_\alpha'' u_\phi'' / r}$
$\varepsilon_A = \overline{\tau_{rr}' \partial_r A''} + \overline{\tau_{r\theta}' (1/r) \partial_\theta A''} + \overline{\tau_{r\phi}' (1/r \sin \theta) \partial_\phi A''}$	$\overline{G_r^A} = -\overline{\rho A'' u_\theta'' / r} - \overline{\rho A'' u_\phi'' / r}$
$\varepsilon_Z = \overline{\tau_{rr}' \partial_r Z''} + \overline{\tau_{r\theta}' (1/r) \partial_\theta Z''} + \overline{\tau_{r\phi}' (1/r \sin \theta) \partial_\phi Z''}$	$\overline{G_r^Z} = -\overline{\rho Z'' u_\theta'' / r} - \overline{\rho Z'' u_\phi'' / r}$
	Differential operators
$\nabla(\cdot) = \nabla_r(\cdot) + \nabla_\theta(\cdot) + \nabla_\phi(\cdot) = \frac{1}{r^2} \partial_r(r^2 \cdot) + \frac{1}{r \sin \theta} \partial_\theta(\sin \theta \cdot) + \frac{1}{r \sin \theta} \partial_\phi(\cdot)$	