

Stellar Evolution Differently

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ABSTRACT

Stellar evolution differently ..

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1 INTRODUCTION

Emerging evidence for neglected physics important for stellar evolution based on 3D hydrodynamic compressible simulations of convective-reactive events force us to reconsider our approach to model structure of stars. It looks like there is a whole zoo of missing stellar physics not included in our steller evolutionary calculations, namely (i) explicit background time-dependency (ii) non-local effects and (iii) compressibility effects.

To include all physics present in convective-reactive flow in stars, one has to be able to cope with the following set of hydrodynamic stellar structure equations, where $\mathcal{T}_m, \mathcal{T}_u, \mathcal{T}_\epsilon, \mathcal{T}_T$: describe background structure; $\mathcal{T}_i, \mathcal{T}_{f_i}$: describe nuclear burning and transport of element i ; and $\mathcal{T}_R, \mathcal{T}_{\sigma_T}, \mathcal{T}_{f_T}$: give correct size of convection zone arising naturally from solution (Biferale et al. 2011).

$$\mathcal{T}_m = -\nabla_r f_\rho + 3\bar{\rho}\tilde{u}_r/r + (f_\rho/\bar{\rho})\partial_r\bar{\rho} - \mathcal{C}_m \quad (1)$$

$$\mathcal{T}_u = -\nabla_r\tilde{R} + \bar{\rho}\tilde{g} - \partial_r\bar{P} - \bar{G}_r^M - \bar{\rho}\tilde{u}_r\partial_r\tilde{u}_r \quad (2)$$

$$(1/S)\mathcal{T}_\epsilon = -\nabla_r(f_i + f_{th} + f_K + f_p) + \bar{\rho}\epsilon_{\text{nuc}} - \mathcal{C}_\epsilon^b + W_b - \tilde{R}_{ir}\partial_r\tilde{u}_i + \bar{\rho}\tilde{D}_i\tilde{u}_i\tilde{u}_i/2 - \partial_r\tilde{L} + (\tilde{\epsilon}_t/S)\partial_r4\pi r^2\bar{\rho}\tilde{u}_r \quad (3)$$

$$\mathcal{T}_T = -\nabla_r f_T + \mathcal{C}_T^b + \mathcal{C}_T^t + \epsilon_{\text{nuc}}/c_v + \mathcal{H}_T^T - \bar{u}_r\partial_r\bar{T} \quad (4)$$

$$\mathcal{T}_i = -\nabla_r f_i + \bar{\rho}\tilde{X}_i^{\text{nuc}} - \bar{\rho}\tilde{u}_r\partial_r\tilde{X}_i \quad (5)$$

$$\mathcal{T}_{f_i} = -\nabla_r f_i^r - f_i\partial_r\tilde{u}_r - \tilde{R}_{rr}\partial_r\tilde{X}_i - \mathcal{P}_{f_i} + \mathcal{A}_{f_i}^{\text{nuc}} + \mathcal{G}_i - \bar{\rho}\tilde{u}_r\partial_r f_i/\bar{\rho} \quad (6)$$

$$\mathcal{T}_R = -\nabla_r(2f_k^r + 2f_p) + 2W_b - 2\tilde{R}_{rr}\partial_r\tilde{u}_r + \mathcal{C}_R^t + 2\mathcal{G}_k^r - \bar{\rho}\tilde{u}_r\partial_r\tilde{u}_r'' + \epsilon \quad (7)$$

$$\mathcal{T}_{\sigma_T} = -\nabla_r f_{\sigma_T} - \mathcal{C}_\sigma^o - 2f_T\partial_r\bar{T} - \mathcal{C}_{\sigma_T}^w + \mathcal{C}_{\sigma_T}^e + \mathcal{H}_{\sigma_T}^T + \mathcal{A}_{\sigma_T}^{\text{nuc}} - \tilde{u}_r\partial_r\sigma_T \quad (8)$$

$$\mathcal{T}_{f_T} = -\nabla_r f_T^r - f_T\partial_r\bar{u}_r - \bar{u}_r\partial_r\bar{T} - \mathcal{P}_{f_T} - \mathcal{C}_{f_T}^o - \mathcal{C}_{f_T}^w - \mathcal{C}_{f_T}^t + \mathcal{C}_{f_T}^s + \mathcal{A}_{f_T}^{\text{nuc}} + \mathcal{H}_{f_T}^T - \tilde{u}_r\partial_r f_T + \mathcal{G}_T \quad (9)$$

$$M = \int \rho dV = \int \frac{m}{V} dV = m(r) \ln V(r) \quad ??(check) \quad (10)$$

this is introduction, test citation Sytine et al. (2000)

Sytine I. V., Porter D. H., Woodward P. R., H. H. S., Winkler K. H., 2000, Journal of Computational Physics, 158, 224

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Table 1. Definitions:

$R = \overline{\rho u_r'' u_r''}$	$\mathcal{C}_m = \overline{\rho d}$	$\mathcal{H}_\sigma^T = \overline{(2/\rho c_v) T' \nabla \cdot F_{th}}$
$f_i = \overline{\rho X_i'' u_r''}$	$\mathcal{C}_\epsilon^b = \overline{P d}$	$\mathcal{C}_{f_T}^o = (\Gamma_3 - 1) (\overline{T} \overline{u_r' d''})$
$\mathcal{T}_m = \partial_t \overline{m}$	$\mathcal{C}_T^b = (1 - \Gamma_3) \overline{T} \overline{d}$	$\mathcal{C}_{f_T}^w = (\Gamma_3 - 1) (\overline{\tilde{d}} \overline{u_r' T'})$
$\mathcal{T}_u = \overline{\rho \partial_t \tilde{u}_r}$	$\mathcal{C}_T^t = (2 - \Gamma_3) \overline{T' d'}$	$\mathcal{C}_{f_T}^t = (\Gamma_3 - 1) (\overline{u_r' T' d''})$
$\mathcal{T}_\epsilon = \overline{\rho \partial_t \tilde{\epsilon}_t}$	$\mathcal{C}_R^t = 2 \overline{P' \nabla_r u_r''}$	$\mathcal{C}_{f_T}^s = \overline{T' u_r' d''}$
$\mathcal{T}_T = \partial_t \overline{T}$	$f_{\sigma_T} = \overline{(T' T' u_r'')}$	$\mathcal{C}_{f_T}^T = \overline{u_r' \nabla \cdot F_{th} / \rho c_v}$
$\mathcal{T}_i = \overline{\rho \partial_t \tilde{X}_i}$	$f_T = \overline{T' u_r''}$	$\mathcal{A}_{f_T}^{\text{nuc}} = \overline{2 T' \epsilon_{\text{nuc}} / c_v}$
$\mathcal{T}_{f_i} = \overline{\rho \partial_t (f_i / \bar{\rho})}$	$\sigma_T = \overline{T' T'}$	$\mathcal{A}_{\sigma_T}^{\text{nuc}} = \overline{u_r' \epsilon_{\text{nuc}} / c_v}$
$\mathcal{T}_R = \overline{\rho \partial_t u_r'' u_r''}$	$\mathcal{C}_\sigma^o = 2(\Gamma_3 - 1) \overline{T} \overline{T' d''}$	$\mathcal{A}_{f_i}^{\text{nuc}} = \overline{u_r'' \rho \dot{X}_i^{\text{nuc}}}$
$\mathcal{T}_{\sigma_T} = \partial_t \overline{T' T'}$	$\mathcal{C}_\sigma^w = 2(\Gamma_3 - 1) \overline{\tilde{d} \sigma_T}$	$\mathcal{P}_{f_T} = \overline{T' \partial_r P / \rho}$
$\mathcal{T}_{f_T} = \partial_t \overline{T' u_r'}$	$\mathcal{C}_\sigma^e = (3 - 2\Gamma_3) \overline{T' T' d''}$	$\mathcal{P}_{f_i} = \overline{X_i'' \partial_r P}$
$d = \nabla \cdot \mathbf{u}$	$\mathcal{H}_T^T = (1/\rho c_v) \nabla \cdot F_{th}$	$\mathcal{S} = 4\pi r^2$

Table 2. Definitions:

ρ density	g_r radial gravitational acceleration
$m = \rho V = \rho \frac{4}{3} \pi r^3$ mass	$M = \int \rho(r) dV = \int \rho(r) 4\pi r^2 dr$ integrated mass
T temperature	$\mathcal{S} = \rho \epsilon_{\text{nuc}}(q)$ nuclear energy production (cooling function)
P pressure	$\tau_{ij} = 2\mu S_{ij}$ viscous stress tensor (μ kinematic viscosity)
u_r, u_θ, u_ϕ velocity components	$S_{ij} = (1/2)(\partial_i u_j + \partial_j u_i)$ strain rate
$\mathbf{u} = u(u_r, u_\theta, u_\phi)$ velocity	$\tilde{R}_{ij} = \overline{\rho u_i'' u_j''}$ Reynolds stress tensor

(11)

Table 3. Definitions:

$$\begin{aligned}
 R &= \overline{\rho u_r'' u_r''} \\
 f_i &= \overline{\rho X_i'' u_r''} \\
 \mathcal{T}_m &= \partial_t \overline{m} \\
 \mathcal{T}_u &= \overline{\rho \partial_t \tilde{u}_r} \\
 \mathcal{T}_\epsilon &= \overline{\rho \partial_t \tilde{\epsilon}_t} \\
 \mathcal{T}_T &= \partial_t \overline{T} \\
 \mathcal{T}_i &= \overline{\rho \partial_t \tilde{X}_i} \\
 \mathcal{T}_{f_i} &= \overline{\rho \partial_t (f_i / \rho)} \\
 \mathcal{T}_R &= \overline{\rho \partial_t u_r'' u_r''} \\
 \mathcal{T}_{\sigma_T} &= \partial_t \overline{T' T'} \\
 \mathcal{T}_{f_T} &= \overline{\partial_t T' u_r'} \\
 \mathcal{S} &= 4\pi r^2 \\
 \mathcal{C}_m &= \overline{\rho \tilde{d}} \\
 \mathcal{C}_\epsilon^b &= \overline{P \tilde{d}} \\
 \mathcal{C}_T^b &= (1 - \Gamma_3) \overline{T \tilde{d}} \\
 \mathcal{C}_T^t &= (2 - \Gamma_3) \overline{T' d'} \\
 \mathcal{C}_R^t &= \overline{2 P' \nabla_r u_r''} \\
 f_{\sigma_T} &= \overline{(T' T' u_r'')} \\
 f_T &= \overline{T' u_r''} \\
 \sigma_T &= \overline{T' T'} \\
 \mathcal{C}_\sigma^o &= 2(\Gamma_3 - 1) \overline{T T' d''} \\
 \mathcal{C}_\sigma^w &= 2(\Gamma_3 - 1) \overline{d \sigma_T} \\
 \mathcal{C}_\sigma^e &= (3 - 2\Gamma_3) \overline{T' T' d''} \\
 \mathcal{H}_T^T &= (1/\rho c_v) \nabla \cdot F_{th} \\
 \mathcal{H}_\sigma^T &= \overline{(2/\rho c_v) T' \nabla \cdot F_{th}} \\
 \mathcal{C}_{f_T}^o &= (\Gamma_3 - 1) \overline{(T u_r' d'')} \\
 \mathcal{C}_{f_T}^w &= (\Gamma_3 - 1) \overline{(\tilde{d} u_r' T')} \\
 \mathcal{C}_{f_T}^t &= (\Gamma_3 - 1) \overline{(u_r' T' d'')} \\
 \mathcal{C}_{f_T}^s &= \overline{T' u_r' d''} \\
 \mathcal{C}_{f_T}^T &= \overline{u_r' \nabla \cdot F_{th} / \rho c_v} \\
 \mathcal{A}_{f_T}^{\text{nuc}} &= \overline{2 T' \epsilon_{\text{nuc}} / c_v} \\
 \mathcal{A}_{\sigma_T}^{\text{nuc}} &= \overline{u_r' \epsilon_{\text{nuc}} / c_v} \\
 \mathcal{A}_{f_i}^{\text{nuc}} &= \overline{u_r'' \rho \dot{X}_i^{\text{nuc}}} \\
 \mathcal{P}_{f_T} &= \overline{T' \partial_r P / \rho} \\
 \mathcal{P}_{f_i} &= \overline{X_i'' \partial_r P}
 \end{aligned}$$

(12)