
Nonlinear Waves Problem 2.1

Michael Nameika

2.1 Solve the following equations using Fourier transforms.

- (a) $u_t + u_{5x} = 0$, $u(x, 0) = f(x)$.
- (b) $u_t + \int K(x - \xi)u(\xi, t)d\xi = 0$, $u(x, 0) = f(x)$, $\hat{K}(k) = e^{-k^2}$.
- (c) $u_{tt} + u_{4x} = 0$, $u(x, 0) = f(x)$, $u_t(x, 0) = g(x)$.
- (d) $u_{tt} - c^2 u_{xx} - m^2 u(x, t) = 0$, $u(x, 0) = f(x)$, $u_t(x, 0) = g(x)$. Contrast this solution with that of the standard Klein-Gordon equation.

Soln. Throughout, let $\hat{f}(k) = \mathcal{F}(f(x))$, $\hat{g}(k) = \mathcal{F}(g(x))$, and $\hat{u}(k, t) = \mathcal{F}(u(x, t))$. We also use Leibniz's Integral rule which allows us to interchange the temporal derivative with the integral.

- (a) Taking the Fourier transform of the differential equation yields

$$\begin{aligned}\mathcal{F}(u_t) + \mathcal{F}(u_{5x}) &= 0 \\ \frac{\partial}{\partial t} \hat{u} + (ik)^5 \hat{u} &= 0 \\ \frac{\partial \hat{u}}{\partial t} + ik^5 \hat{u} &= 0.\end{aligned}$$

Note that this is an ODE in t , and separating variables gives

$$\begin{aligned}\int \frac{\partial \hat{u}}{\hat{u}} &= -ik^5 \int dt \\ \implies \log(\hat{u}) &= -ik^5 t + C_0(k) \\ \implies \hat{u}(k, t) &= C_1(k)e^{-ik^5 t}.\end{aligned}$$

Where $C_0(k)$ is a function of integration independent of t . Setting $t = 0$ in the above equality yields

$$\begin{aligned}\hat{u}(k, 0) &= C_1(k) \\ \implies \hat{f}(k) &= C_1(k).\end{aligned}$$

Thus

$$\begin{aligned}\hat{u}(k, t) &= \hat{f}(k)e^{-ik^5 t} \\ \implies u(x, t) &= \frac{1}{2\pi} \int_{-\infty}^{\infty} \hat{f}(k)e^{i(kx - k^5 t)} dk.\end{aligned}$$

- (b) Taking the Fourier transform of the differential equation and invoking the convolution theorem gives

$$\begin{aligned}\mathcal{F}(u_t) + \mathcal{F}\left(\int_{-\infty}^{\infty} K(x - \xi)u(\xi, t)d\xi\right) &= 0 \\ \frac{\partial \hat{u}}{\partial t} + \hat{K}(k)\hat{u}(k) &= 0 \\ \frac{\partial \hat{u}}{\partial t} + e^{-k^2} \hat{u}(k) &= 0.\end{aligned}$$

As in part (a), separation of variables gives

$$\begin{aligned}\int \frac{d\hat{u}}{\hat{u}} &= - \int e^{-k^2} dt \\ \implies \log(\hat{u}) &= -e^{-k^2} t + C_0(k) \\ \implies \hat{u}(k, t) &= C_1(k) e^{-te^{-k^2}}\end{aligned}$$

and setting $t = 0$ yields

$$\begin{aligned}\hat{u}(k, 0) &= C_1(k) \\ \implies \hat{f}(k) &= C_1(k)\end{aligned}$$

so that

$$\begin{aligned}\hat{u}(k, t) &= \hat{f}(k) e^{-te^{-k^2}} \\ \implies u(x, t) &= \frac{1}{2\pi} \int_{-\infty}^{\infty} \hat{f}(k) e^{ikx - te^{-k^2}} dk.\end{aligned}$$

(c) Taking the Fourier transform of the differential equation yields

$$\begin{aligned}\mathcal{F}(u_{tt}) + \mathcal{F}(u_{xx}) &= 0 \\ \frac{\partial^2 \hat{u}}{\partial t^2} + (ik)^4 \hat{u} &= 0 \\ \implies \frac{\partial^2 \hat{u}}{\partial t^2} + k^4 \hat{u} &= 0.\end{aligned}$$

From here, we seek solutions of the form

$$\hat{u}(k, t) = h(k) e^{rt}$$

so that, putting this ansatz into the above differential equation gives

$$\begin{aligned}r^2 h(k) e^{rt} + k^4 h(k) e^{rt} &= 0 \\ \implies r &= \pm i k^2.\end{aligned}$$

Then

$$\begin{aligned}\hat{u}(k, t) &= h_1(k) e^{ik^2 t} + h_2(k) e^{-ik^2 t} \\ \implies \hat{u}_t(k, t) &= ik^2 h_1(k) e^{ik^2 t} - ik^2 h_2(k) e^{-ik^2 t}.\end{aligned}$$

From our initial conditions, we find

$$\begin{aligned}h_1(k) + h_2(k) &= \hat{f}(k) \\ h_1(k) - h_2(k) &= -\frac{i}{k^2} \hat{g}(k) \\ \implies h_1(k) &= \frac{1}{2} \left(\hat{f}(k) - \frac{i}{k^2} \hat{g}(k) \right) \\ h_2(k) &= \frac{1}{2} \left(\hat{f}(k) + \frac{i}{k^2} \hat{g}(k) \right).\end{aligned}$$

Thus

$$\hat{u}(k, t) = \frac{1}{2} \left[\left(\hat{f}(k) - \frac{i}{k^2} \hat{g}(k) \right) e^{ik^2 t} + \left(\hat{f}(k) + \frac{i}{k^2} \hat{g}(k) \right) e^{-ik^2 t} \right]$$

and taking the inverse Fourier transform gives us the solution to the PDE:

$$u(x, t) = \frac{1}{4\pi} \int_{-\infty}^{\infty} \left(\hat{f}(k) - \frac{i}{k^2} \hat{g}(k) \right) e^{i(kx + k^2 t)} dk + \frac{1}{4\pi} \int_{-\infty}^{\infty} \left(\hat{f}(k) + \frac{i}{k^2} \hat{g}(k) \right) e^{i(kx - k^2 t)} dk$$

(d) Taking the Fourier transform of the PDE gives

$$\begin{aligned}\mathcal{F}(u_{tt}) - c^2 \mathcal{F}(u_{xx}) - m^2 \mathcal{F}(u) &= 0 \\ \frac{\partial^2 \hat{u}}{\partial t^2} + c^2 k^2 \hat{u} - m^2 \hat{u} &= 0 \\ \implies \frac{\partial^2 \hat{u}}{\partial t^2} + (c^2 k^2 - m^2) \hat{u} &= 0.\end{aligned}$$

As in part (c), consider the ansatz $\hat{u}(k, t) = h(k)e^{rt}$. Plugging this ansatz into our above differential equation yields

$$\begin{aligned}h(k)r^2 e^{rt} + (c^2 k^2 - m^2)h(k)e^{rt} &= 0 \\ \implies r &= \pm i \sqrt{c^2 k^2 - m^2}.\end{aligned}$$

Let $\omega(k) = \sqrt{c^2 k^2 - m^2}$. Then The general solution to our ODE in \hat{u} is

$$\hat{u}(k, t) = h_1(k)e^{i\omega(k)t} + h_2(k)e^{-i\omega(k)t}.$$

From our initial conditions, we have

$$\begin{aligned}h_1(k) + h_2(k) &= \hat{f}(k) \\ i\omega(k)h_1(k) - i\omega(k)h_2(k) &= \hat{g}(k) \\ \implies h_1(k) &= \frac{1}{2} \left(\hat{f}(k) - \frac{i}{\omega(k)} \hat{g}(k) \right) \\ h_2(k) &= \frac{1}{2} \left(\hat{f}(k) + \frac{i}{\omega(k)} \hat{g}(k) \right).\end{aligned}$$

Thus

$$\hat{u}(k, t) = \frac{1}{2} \left(\hat{f}(k) - \frac{i}{\omega(k)} \hat{g}(k) \right) e^{i\omega(k)t} + \frac{1}{2} \left(\hat{f}(k) + \frac{i}{\omega(k)} \hat{g}(k) \right) e^{-i\omega(k)t}.$$

Taking the inverse transform yields

$$u(x, t) = \frac{1}{4\pi} \int_{-\infty}^{\infty} \left(\hat{f}(k) - \frac{i}{\omega(k)} \hat{g}(k) \right) e^{i(kx + \omega(k)t)} dk + \frac{1}{4\pi} \int_{-\infty}^{\infty} \left(\hat{f}(k) + \frac{i}{\omega(k)} \hat{g}(k) \right) e^{i(kx - \omega(k)t)} dk$$

with $\omega(k) = \sqrt{c^2 k^2 - m^2}$. Note that if $c^2 k^2 \geq m^2$, $\omega(k) \in \mathbb{R}$, and for $c^2 k^2 < m^2$, $\omega(k) \in \mathbb{C}$. In the latter case, we have exponential behavior. Note that this differs from the Klein-Gordon equation since $\omega(k) \in \mathbb{R}$ for all cases.