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Designing a Common Platform for Modelling Soot in Flames and Reactors

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Nomenclature

Acronyms

CFD	Computational Fluid Dynamics
DEM	Discrete Element Modelling
HACA	Hydrogen abstraction carbon addition
MPBM	Monodisperse Population Balance Model
SPBM	Sectional Population Balance Model

Constants

Av	Avogadro's number	$6.02214076 \times 10^{23}$	1/mol
k_B	Boltzmann constant	$1.3806488 \times 10^{-23}$	$\text{m}^2\text{kg/s}^2\text{--K}$

English symbols

C_{tot}	Totla carbon content of soot particles (per section)	mol/kg
D	Diffusion coefficient of particles	m^2/s
d	diameter	m
d_c	Collision diameter	m
d_g	Gyration diameter	m
D_H	Duct hydraulic diameter	—
d_m	Mobility diameter	m
d_p	Primary particle diameter	m
f	Friction factor	—
H_{tot}	Totla hydrogen content of soot particles (per section)	mol/kg
I	Source terms for soot variables	$\text{mol}/(\text{kg} \cdot \text{s})$
k	reaction rate constant	$\text{m}^3/(\text{mol} \cdot \text{s})$
Kn	Knudsen number	—
m	mass	kg
n_p	Number of primary particles of per agglomerate	—
N_{agg}	Number density of agglomerates	mol/kg
N_{pri}	Number density of primary particles	mol/kg
R_H	Duct hydraulic radius	—
SF	Sectional spacing factor	—
T	Temperature	K
W	Molecular weight	kg

Greek symbols

β	Collision frequency	m^3/s
δ_a	Mean distance of particles	m
λ	Mean free path	m
λ_a	Mean stopping distance of particles	m
ν	Gas kinematic viscosity	m^2/s
ρ	Gas density	kg/m^3

ρ_{soot}	Soot density	kg/m ³
φ	Soot volume fraction	m ³ /m ³

Subscripts

ads	Adsorption
agg	Agglomerate
coag	Coagulation
cont	Continuum
f	forward
fm	Free molecular
g	Gyration
grow	Surface growth
inc	Inception
m	Mobility
ox	Oxidation
pri	Primary particle
r	reverse
reac	reactive

Chapter 1

Introduction

1.1 Motivation

Carbonaceous nanoparticles, such as Carbon Black (CB) and soot, are widely encountered in nature and engineering. Every year, nearly 9.5 megatons of soot (black carbon) is emitted into the atmosphere from anthropogenic activities and natural sources such as wildfires and volcanoes [2]. The wide light absorption range of soot reduces the albedo of snow-covered area and alters the radiative forcing balance in the atmosphere making soot the third strongest contributor to climate change after methane and carbon dioxide [2]. Also, exposure to combustion-generated soot could promote respiratory and cardiovascular diseases [3]. So, strict regulations are targeting combustion engines to limit environmental and health risks from soot formation [4]. Accurate and affordable models are essential for the prediction of soot composition and morphology in combustion devices and the reduction of emissions in engines. They can also provide better understanding of soot optical properties that are major indicator of soot environmental effects.

On the other hand, CB with a similar synthesis process and structure to soot but higher elemental carbon to hydrogen ratios (>97%) [5] is commercially produced and sold in large scales. In fact, CB is the largest industrially produced nanomaterial by value and volume (~15 megatons per year with a value of \$17B) with applications as a reinforcing agent in rubber and tire industries [6], and conductive additive in lithium-ion batteries [7]. CB is primarily manufactured by the so-called furnace process in which about 50% of heavy fuel oil is partially combusted to convert the rest of it into CB [8]. This process suffers from low mass yield and excessive emission, generating 4 tons of CO₂ per each ton of product on average [9]. Plasma reactor is an emerging alternative production method with distinct advantages over flame-based methods: They can achieve 100% carbon yields with no direct CO₂ emission or other pollutants [10], and the energy required for pyrolysis is supplied by an electric arc that does not depend on the feedstock composition. Controlling CB properties such as its specific surface area (or primary particle diameter), hard agglomerate size (or gyration diameter of agglomerates with primary particles connected to each other by strong chemical bonds), and composition (or particle carbon to hydrogen ratio) is important to make the process economical and to achieve specific grades of CB for different target applications. However, this is a challenging task because of the complexity of CB formation and mass growth processes and its coupling with gas phase chemistry, dependence on local temperature and pressure. This requires accurate process design and optimization tools built on physics of CB formation and evolution to inform manufacturers' decisions for production of CB with desired grades [11].

The term "soot" usually refers to the unwanted particulate matter formed during incomplete combustion of any carbon-containing material from jet and diesel fuel to wood, heavy oil, and plastics with variable organic content and a large variation in C/H ratio [5], but this research focuses on soot particles generated under controlled laboratory conditions from fuels with known compositions. The mature soot formed in methane and ethylene premixed flame can reach 95% elemental C/H ratio [12], which is close to CB composition. The comparison of transmission electron microscopy (TEM) images of industrially produced CB [13] with soot sampled from diesel fuel [14, 15] shown in

Fig.1.1 indicates similarity of their morphology and structure. Hereafter, soot will be used to collectively refer to carbonaceous nanoparticles produced in flame/reactor during combustion/pyrolysis processes.

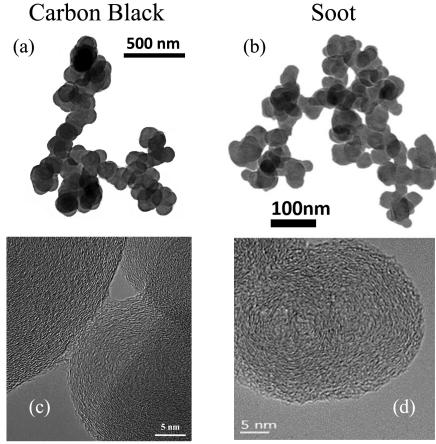


Figure 1.1: The TEM images of Carbon black (a & c) [13] and soot (b & d) [14, 15] that shows soot and CB has similar morphology and structure

1.2 Background

1.2.1 Soot inception and surface growth

The physics of soot formation involves concurrent processes with different time and length scales starting from the breakdown of hydrocarbon molecules to intermediate species and radicals that transition to first soot particles and grow via heterogeneous reactions on the surface as well as coagulation [16]. The TEM analysis of soot sampled from flames [15], reactors [17], and engines [18] with different fuels and process conditions revealed a common fractal-like morphology, often characterized as agglomerates of primary particles. High resolution transmission electron microscopy (HRTEM) of soot primary particles showed clusters of precondensed Polycyclic Aromatic Hydrocarbons (PAHs) [19] pointing to PAH as main soot precursors. This hypothesis was also supported by thermodynamic stability of PAHs enabling them to resist dissociation at high flame temperatures [20].

However, the transition PAHs (soot precursors) to incipient soot, known as *soot inception* has not been well understood at the level of pathways and elementary reactions [21] primarily due to uncertainties in PAH (precursor) chemistry. The thermodynamics of PAH growth from acetylene (C_2H_2) as the dominant hydrocarbon species in fuel pyrolysis and an intermediate of molecular growth does not experience a significant enthalpy release or entropy increase, so the path is driven by a gradual reduction in Gibbs free energy. As a result, the kinetics of PAH growth into subsequent soot particles can be highly reversible hence sensitive to local temperature, pressure and intermediate species concentration.

The growth of PAHs beyond first ring (benzene) is predominantly driven by so-called hydrogen abstraction carbon (acetylene) addition (HACA) mechanism [22] where a hydrogen radical abstracts an hydrogen atom at the edge of PAH, providing a reactive site for acetylene addition. The kinetic reversibility of HACA opens the door for competing pathways such as chain reactions of resonance-stabilized radical (RSR). Propargyl is a prominent example of these radicals, whose combination is known as a major contributor to benzene formation [23]. Built on this hypothesis, Johansson et al. [24] proposed a radical-driven growth mechanism by addition of vinyl (C_2H_3) starting from cyclopentadienyl (C_5H_5) to larger hydrocarbon radicals that can survive long enough in high temperatures to react with other radicals, PAHs and unsaturated aliphatic species, through radical

chain reactions. However, low concentrations of these radicals limit the growth rate through RSR pathways [25]. Moreover, some of the intermediate steps for radical regeneration such as formation of vinylcyclopentadienyl were shown to be kinetically unfavorable [25] compared to HACA. Regardless of their mechanistics, these sequential growth mechanisms, termed as *chemical growth* cannot account for rapid soot formation [22] and its nanostructure [26].

There is abundant but mostly indirect experimental and computational evidence pointing to an alternative collision-based mechanism for soot inception. HRTEM images of nascent and mature soot shows disordered PAH clusters highlighting the role of PAH collisions. The bimodality of particle size distribution (PSD) of nascent soot particles in premixed flames [27] indicates that the kinetics of inception is second order in precursor concentration [28]. The time-of-flight mass spectrometry (TOFMS) experiments in a 13 kPa acetylene-oxygen flame showed a series of peaks with a periodicity of 500 amu [29].

It is not clear how PAH clusters form and what forces allow the binding occur and resist dissociation at flames temperatures (≥ 1600 K). Frenklach [22] characterized the clustering as a physical process where the sticking of PAHs upon collision forms dimers held together by Van der Waals (vdW) forces without involving chemical reactions. In fact, Herdman and Miller [30] found that the binding energy of PAHs dimers due to dispersive and electorstatic forces increases linearly with molecular mass and reaches the limit of exfoliation energy for graphite. However, the entropy barrier of dimerization increases with PAH size making them unfavorable under equilibrium conditions. So, PAHs as large as circumcoronene ($C_{54}H_{18}$) can only form dimer to survive flame temperatures [21]. However, the concentration of large PAHs are too low to account for the observed number density of soot particles in flames [31].

There is also the possibility of PAH clustering governed by non-equilibrium kinetics. PAH molecules can form rovibrationally excited dimer with long enough life-time [32] to react with H atoms forming covalent bonds. There is another possibility for these clusters to be joined by aliphatic linkages. Micro-FT-IR spectroscopy analysis by Cain et al. [33] showed the ratio of aliphatic-to-aromatic C-H bonds can exceed unity at the flame temperature. Using these findings, they suggested that the aliphatic components in the form of alkyl and alkenyl can covalently bound to aromatic units in soot particles. Although such a mechanism is viable near the flame region, it cannot explain persistent soot inception in the post-flame zone [34] where the H atom concentration is too low to initiate H abstraction reactions that produce those radicals.

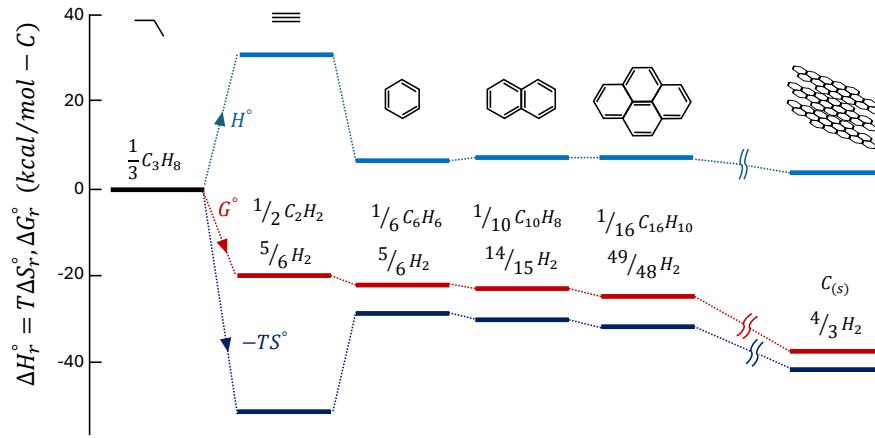


Figure 1.2: Standard enthalpy (ΔH) and entropy ($T\Delta S$) contributions to Gibbs function of reaction ΔG at 1600 K for carbon formation from propane(reprinted from Ref. [21])

There are other complicating factors that hinders the fundamental understanding of soot inception such as lack of a decisive criterion to distinguish gaseous molecules from particles [16], and the overlap of inception with surface growth and agglomeration [35]. Measurement techniques have limited capability in soot detection due to short time and length scales of soot inception and surface

growth ranging from pico- to milliseconds and micro- to millimeter [36]. Fig. 1.3 demonstrates the length and time scales relevant to different stages of soot formation from PAH precursors to incipient, nascent and mature soot in flames. Moreover, the collected data from measurements might not be true representatives of soot formed at the probed location of the studied process. For example, intrusive diagnostic methods based on thermophoresis or dilution have a sampling probe that can perturb flow dynamics or alter structure and composition of soot [37]. Non-intrusive techniques such as optical methods also rely on assumptions about light absorption and scattering of soot [38] and its morphology [39].

Despite the gaps in fundamental understanding of soot formation and limitations of diagnostics methods, models have been developed describe the soot inception and growth. These models have been formulated as a set of clear pathways that explain soot inception based on collisions of PAH molecules. They have to be consistent with current knowledge of soot physics, feasible to be coupled with chemistry and particle dynamics models, and able to predict soot mass, PSD and morphology observed in flames and reactors.

The classic description of soot inception relies on PAH dimerization where collision of two PAH molecules (monomers in this context) forms a dimer held together by Van der Waals forces [40]. The dimerization is a irreversible process with an efficiency that accounts for the reversibility or dissociation of dimers. The theory postulates that PAH growth continues by sequential addition of a monomer (PAH molecule) forming stacks of dimers, trimers, tetramers and so on to reach a certain mass threshold that marks the emergence of incipient soot [40], but for practical purposes, a dimer is usually considered as incipient soot. Here, we call this model *Irreversible Dimerization*. Irreversible Dimerization has been used to predict soot formation in burner-stabilized premixed [41, 42], counterflow diffusion flames [43, 44], coflow diffusion flames [45, 46]. A collision efficiency factor ranging between 10^{-6} to 1 is also employed to adjust the inception flux and PAH adsorption rates to achieve desired soot mass and size distribution. PAHs of moderate sizes such as pyrene (4 rings) to coronene (7 rings) have been considered as the starting point of inception due to their thermodynamic instability that justifies the irreversibility at high temperatures [40]. However, the theoretical calculations [47] and experiments [48] indicated that PAH dimerization is highly reversible in flame conditions.

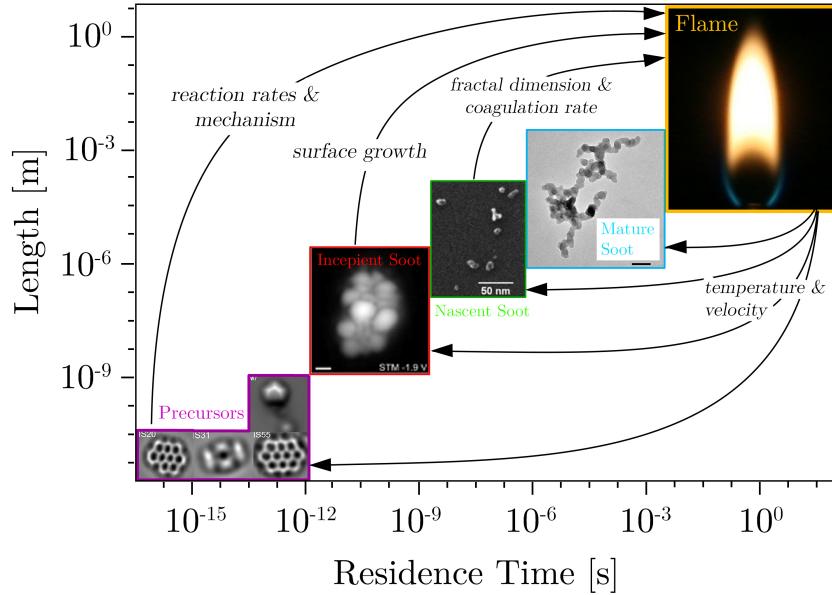


Figure 1.3: The range of time and length scales of the processes involved in soot formation from molecular reactions to particle-fluid interaction in flames

The inception flux of irreversible dimerization is mainly controlled by PAH concentration due

to weak temperature dependence, so it produces new particles at low temperatures (even less than 500 K) [49] despite experimental evidence for termination of inception below 1200 K [50, 51]. Also, the arbitrary selection of efficiency factors alters the distribution of mass between inception of surface growth could significantly change soot mass, PSD, and morphology [52]. Miller [53] used equilibrium constant for PAH dimerization to calculate the net dimerization rate and demonstrated that the collision of PAHs larger than circumovalene (~ 800 amu) could last long enough grow into incipient soot. However, the concentration of PAHs drops rapidly with size [21]. The entropy barrier of dimerization is significant for larger PAHs [54].

Eaves et al. [55] relaxed the irreversibility assumption, and developed a reversible clustering model to simulate inception using an array of PAHs from naphthalene to benzo-pyrene. Building on that work, Kholghy et al. [56] emphasized on the necessity of chemical bond formation after physical PAH clustering for accurate prediction of volume fraction, primary particle diameter and PSD in ethylene coflow diffusion flames. Later, Kholghy et al. [57] proposed the "*Reactive Dimerization*" model which starts with reversible collision of PAHs leading to physical dimers held with vdW forces that are graphitized and form chemically-bonded dimers that serve as soot nuclei grow via surface reactions. They also performed a systematic analysis on the contribution of different PAHs, and concluded that one- and two-ring aromatics account for almost all of inception flux in the so-called "*sooting flame*" [42]. However, Frenklach and Mebel [25] pointed out that an inception model that initiated with a highly reversible step similar to Reactive Dimerization [57] cannot produce sufficient flux of particles to match measurements of the benchmark burner-stabilized stagnation flame [28]. Instead, they proposed a HACA-driven mechanism where addition of monomer molecule to its radical activated by hydrogen abstraction for a stable dimer via an E-Bridge bond formation, and this sequential process continues to form trimers, tetramers, and larger PAH clusters.

The gas-phase chemistry of aromatics can be extended to account for chemical growth of incipient soot via surface reactions [22]. This hypothesis, known as "chemical similarity" postulates that the reactions occurring on the soot surface are similar to those involving large molecules of PAHs in the gas phase. It also provides means to describe the rates of surface growth and particle oxidation in terms of elementary chemical reactions. In other words, it is assumed that the surface of soot particles is made up of lateral faces of large PAHs covered with C-H bonds. This is the basis for HACA mechanism [40, 58] that assumes the soot surface to consist of hydrogenated sites with a predefined density. Mass growth on soot surface requires H-abstraction to form a radical site, followed by acetylene attack similar to growth of PAH molecules in the gas-phase. The reactivity of these sites changes with time and temperature [59, 60], described as soot aging. For modelling purposes, a temperature-dependent multiplier, usually represented by α , was introduced to account for these effects. Appel et al. [58] showed α changes with temperature and particle size.

Adsorption of PAHs on the surface of soot particles is also a viable growth mechanism [40], more specifically called physisorption or chemisorption depending on the mechanisms driving the adsorption process [61]. There is still debate over the stability of adsorbed PAH molecules on soot surface [62]. Following the hypothesis that PAHs are building blocks of soot particles, a mechanism similar to inception is often used to describe PAH-soot growth.

1.3 Coagulation and agglomeration

In typical soot formation processes such as flames and reactor, soot particles are formed at high concentrations (10^{12} 1/cm³), and inception and surface growth are relatively short compared to the total residence of soot particles. As a result, coagulation becomes dominant rapidly attaining both [63] self-preserving size distribution (SPSD) [64] and asymptotic fractal-like structure [65]. The evolving fractal-like structure of agglomerates quantified by their mobility diameter normalized by primary particle, d_m/d_p , and gyration, d_m/d_g , diameters can be described with power laws derived from mesoscale simulations [66]. The collision frequency of agglomerates depends on their evolving fractal-like morphology. Also, polydisperse agglomerates collide more frequently than monodisperse ones. The enhancement in their collision frequency reaches an asymptotic value of 35% [63] or 82% [67] in the free molecular or transition regimes, respectively at SPSD regardless of the poly-

dispersity in their constituent primary particles. Particle morphology formed by inception, surface growth and agglomeration can be tracked precisely by mesoscale simulations, such as Discrete Element Modeling (DEM) [68]. However, they are computationally expensive and interfacing them with chemical kinetics in computational fluid dynamics (CFD) simulations is not trivial [69]. This limits their application. So, sectional population balance models (SPBM) are often used to track agglomerate and primary particle size distribution [70], morphology [71], and composition [45] in complex laminar [45] and turbulent flows [72]. Using the SPBMs coupled with relations for agglomerate fractal-like structure [73] and collision frequency [36], particle size distribution, morphology and composition can be tracked accurately. However, the computational cost of SPBMs increases exponentially with the number of sections [31] and particle properties [45] tracked. Thus, one property (e.g. agglomerate mass) is typically tracked with SPBMs to reduce computational cost. This does not allow to account for agglomerate fractal-like structure [74, 75] which limits SPBM accuracy in predicting surface growth and coagulation rate of agglomerates and their size distribution.

Alternatively, particle dynamics can be tracked by the method of moments (MOM) [76] or monodisperse population balance models (MPBM) [77]. Such models only track average particle properties (e.g. moment ratios) and their accuracy could be limited if unrealistic assumptions (e.g. approximating agglomerates as monodisperse and perfect spheres) are used. However, when inception and surface growth are short [78] and high particle (number) concentrations are formed [66], they lead to rapid attainment of self-preserving size distributions (SPSD) and agglomerates having asymptotic structure [63]. In this case a MPBM or MOM can be assembled on a firm scientific basis with accuracy on par with DEM [68], SPBM [79] and experimental data [80, 81, 82]. Such models can be readily interfaced with CFD simulations [83] without significant computational cost, making them ideal for three-dimensional and even turbulent flame simulations.

The MOM tracks moments of the PSD and estimates average particle properties such as mass [84], surface area [85], the number of constituent primary particles per agglomerate, n_p [76], or even particle composition [86] using the ratio of the moments. The MOM with four equations was used to describe synthesis of optical fibers by simultaneous reaction, diffusion, coagulation and thermophoresis of SiO_2 in laminar flow reactors assuming a lognormal PSD [87]. The MOM with interpolative closure (MOMIC) was developed to predict simultaneous nucleation, surface growth and coagulation of soot agglomerates and estimate its PSD with six equations [76]. To calculate source terms of the transported moments, additional moments that are not tracked are needed preventing the closure of the system of differential equations with the MOM [84, 88]. Thus, often the PSD shape is assumed a priori [84] or extra equations are solved to estimate it [77].

The MPBMs do not have the closure problem and calculate average particle properties by tracking their total concentration, mass [77] and area [89, 90]. Kruis et al. [77] used a 2-equation MPBM (known as the semi-empirical model) to track soot concentration and mass in (non-premixed) flames assuming spherical particles. Good agreement was achieved for measured soot mass. However, the specific surface area [90] and coagulation frequency of spheres are significantly smaller compared to that of agglomerates with the same mass underestimating their oxidation rate [79] and overestimating their concentration [40]. Kruis et al. [77] proposed a 3-equation MPBM to account for the fractal-like structure of nanoparticle agglomerates during coagulation and sintering. Agglomerate volume and area were used to obtain their equivalent primary particle diameter, d_p , and n_p . Then, agglomerate collision diameter, i.e. d_g , was calculated by D_f , d_p and n_p to account for their fractal-like structure that affects their collision frequency. Tsantilis and Pratsinis [89] extended the MPBM to predict hard-(chemically-bonded) and soft- (physically-bonded) agglomerates during synthesis of SiO_2 and TiO_2 [91] nanoparticles with simultaneous reaction, surface growth, coagulation and sintering. Such a MPBM applies best at high concentrations when inception and surface growth are short [78] resulting in the dominance of coagulation where particles rapidly reach their SPSD and asymptotic fractal-like structure. This is often the case for soot emitted from a variety of combustion devices or CB reactors where inception and surface growth are limited to only a few milliseconds when temperature is very high (i.e. $T \geq 1500\text{K}$) [57].

1.4 Soot maturity and its optical properties

The maturity level of soot is described as the evolution of physical and chemical properties from incipient to graphite-like mature soot [92]. It involves the growth of the graphitic crystallite fine structure of soot within, and perpendicular to, the aromatic layers [93, 94], known as graphitization, followed by increase in the size of the crystallite-layer planes and decrease of the interlayer spacing [19]. This process is accompanied by the pyrolytic conversion of hydrocarbon species and substituted hydrocarbons toward elemental carbon and increase in the carbon-to-hydrogen (C/H) ratio, known as carbonization. Soot maturity is closely associated with reactivity of surface sites [95]. Fig. 1.4 compares nanostructure of nascent soot composed of disoriented PAH clusters with that of mature soot with a core-shell pattern where disordered core is surrounded by concentrically oriented graphitic layers [29]. The core-shell nanostructure of mature soot strongly depends on process conditions such as pressure, fuel identity, temperature and residence time of the particles.

Evolution of soot maturity and morphology impacts its optical properties. Incipient and nascent soot absorb shorter wave lengths ($\lambda < 600 \mu\text{m}$) as opposed to mature soot particles that are broad-band light absorbers [96]. Non-intrusive optical diagnostic methods such as light extinction (absorption and scattering) [97] and Laser Induced Incandescence (LII) [38] are widely utilized to measure soot volume fraction, f_v using extinction coefficient and the absorption function, $E(m)$ that depends on soot refractive index, m , soot composition and morphology [98].

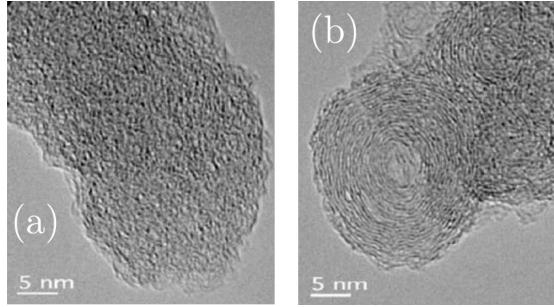


Figure 1.4: The HRTEM image of (a) a nascent soot primary particle with disordered internal nanostructure juxtaposed with that of (b) a mature soot primary particle with well-organized cluster near the shell sampled from soot generated by pyrolysis of ethanol at 1250 and 1650, respectively. Reprinted from Ref.[99]

However, the effects of soot composition, and morphology on its refractive index, and absorption function are not fully understood yet. $E(m)$ of soot at $\lambda = 1064 \text{ nm}$ measured by two-color LII measurements increases from 0.193 to 0.349, and 0.226 to 0.340 for ethylene premixed flames with equivalence ratio of $\phi=2.1$ and $\phi=2.3$, respectively when height above burner (HAB) changes from 8 to 14 mm [100]. During acetylene pyrolysis in a shock tube, $E(m)$ of soot particles increases from 0.05 to 0.25 as their primary particle diameter, d_p , grows to 20 nm within 1.6 ms [101]. Using m and $E(m)$ of mature spherical soot [102] at different HABs neglecting the impact of soot morphology and composition on m [103] can overestimate soot volume fraction by 100% [104]. So, accurate estimation of the evolving optical properties of soot is essential to close the carbon mass balance in the measurements [104], and analyze reaction kinetics for soot inception [57] surface growth [58] and oxidation [105] that are essential for development and validation of aerosol dynamics models coupled with chemistry to predict soot formation.

The dependence of soot m on its morphology and composition can be quantified by soot optical band gap, E_g [106]. The optical band gap concept was originally proposed by Tauc et al. [107] for semi-conductors and later developed for amorphous carbon [108], and it can be used to describe the crystalline character of PAH clusters in soot [108]. Incipient flame-made soot with diameters less than 20 nm exhibit quantum dot behavior [109] with $0.7 < E_g < 2 \text{ eV}$ that is larger than E_g of graphitized soot, 0.12 eV [109]. Optical band gap of organic carbon coated on soot remains nearly constant (1.8-1.9 eV) over soot evolution [110]. In contrast, Russo et al. [111] measured soot optical

band gap using ex-situ and in-situ methods and showed that E_g drops from 0.7 eV at HAB= 8 mm to 0.2 eV at HAB=14 mm in an ethylene premixed flame with $\phi=2.3$ as particles grow and become more mature by carbonization. Kelesidis and Pratsinis [112] correlated the evolving refractive index of soot with its E_g obtained from quantum confinement theory (QCT) [109] for wavelengths of $\lambda= 532$ and 1064 nm [112] by linear interpolations between that of nascent and mature soot. The proposed relations were then used with soot morphologies obtained from Discrete Element Modeling (DEM) simulations to obtain the evolving Mass Absorption Cross-section (MAC), $E(m)$, and the absorption function ratio of soot ($E(\lambda=532\text{nm})/ E(\lambda=1024\text{nm})$) and compared with measurements [113, 114, 115]. Employing the relation in laser diagnostics to reprocess light extinction measurement data resulted in accurate prediction of f_v in moderate ($\phi=2.34$) [104] and rich ($\phi > 3$) ethylene premixed flames [116], and f_v along the centerline of a coflow ethylene diffusion flame [117] which is necessary to close the carbon mass balance in soot generating processes.

1.5 Oxidation

Soot oxidation is described as the removal of soot mass by reaction of molecular oxygen (O_2), oxygen radical (O), and hydroxyl radical (OH) from soot surface. Because of its heterogeneous reaction kinetics and mechanisms, oxidation is expected to be sensitive to surface structure and composition. So, oxidation mechanisms depend on soot maturity and temperature-time history. In near stoichiometric and fuel-rich conditions, the contribution of OH radical to soot oxidation is predominant [118] compared to that of O radicals [119]. Femimore and Jones [120] investigated soot oxidation rate for low oxygen partial pressures and temperatures from 1530 to 1890 K, highlighted the importance of OH as a major oxidation agent, and attributed the faster rates compared to those predicted by Lee et al. [121], to OH oxidation. One approach for describing OH oxidation is by introducing collision efficiency representing the fraction of collisions of OH with soot particles that resulted in the removal of a carbon atom [118]. Scanning Mobility Particle Sizer (SMPS) with a two-stage burner also showed a collision efficiency of 0.13 [122].

The empirical relation of Nagle and Strickland-Constable (NSC) [123] originally developed for oxidation of pyrolytic graphite have been widely used to describe soot oxidation by O_2 . It describes O_2 oxidation rates based on partial pressure of O_2 and the fraction of reactive edge sites to less reactive basal planes using the graphite analogy for soot. The dominance of edge sites in typical combustion conditions accounts for the relatively small reactivity of soot and basal planes become reactive at high temperature (>2500 K) [119]. O_2 oxidation kinetics was also described with a power-law kinetics [121], and shown to be first order in oxygen concentration [124]. Alternatively, soot oxidation can be explained based on chemical similarity using HACA mechanism assuming that active site are attacked by O_2 and OH leading to loss of carbon and release of CO.

A different oxidation regime has been identified for soot at low temperatures ($T<1000$ K) where O_2 diffuses and reacts with bulk soot accounting for most of soot mass consumption [81]. Internal oxidation compacts the pore network of soot resulting in hollow CB [125], diesel [126] and biodiesel [127] soot particles and increases their SSA up to a factor of four [126]. Oxidation can cause fragmentation of soot particles affecting their morphology. The increase in number concentration and the change of soot morphology in lean premixed flames [128] and in the oxidation region of diffusion flames [129] was attributed to fragmentation. Such a change in agglomerate morphology has not been observed in fuel-rich conditions, so fragmentation was linked to O_2 oxidation.

1.6 Research objectives

The purpose of this project is developing a process design tool to predict the yield, morphology, composition and size distribution of soot particles in industrial flames and reactors under different temperatures, pressures and residence times. Integrating gas chemistry, all existing inception and surface growth models and particle dynamics in a single package provides a fully coupled soot description from breakdown of fuel to agglomeration and oxidation of particles that enables sensitivity analysis to shed light on key pathways for soot inception and surface growth in various targets. This

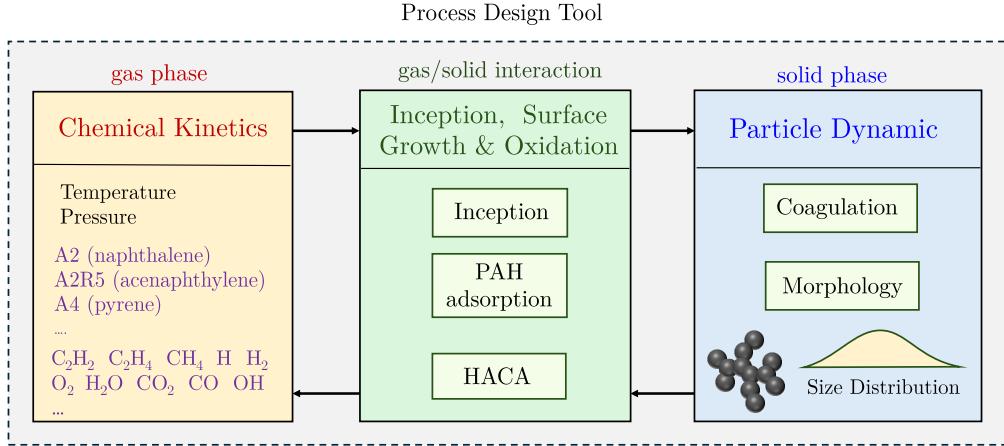


Figure 1.5: The conceptual structure of developed process design tool that account for gas mixture properties, and soot chemistry involving inception, surface growth and oxidation as well as its coagulation leading to the fractal-like morphology of soot particles

tool can also be used for rate constant optimization and mechanism reduction under soot forming conditions. The conceptual structure of this tool is shown in Fig 1.5 that includes a chemical kinetics unit coupled with Cantera [130] to compute thermal and physical properties of gas mixture such as temperature, pressure, density, and enthalpy based on ideal equation of state for the gas mixture. The chemistry of gas mixture is quantified using detail reaction mechanisms that enables tracking concentration of fuel, intermediate species such as C₂H₂ and H₂ essential to surface growth, and PAHs such as benzene (A1), naphthalene (A2) and pyrene (A4) known as building blocks of soot. The second unit accounts for soot inception and surface growth via PAH adsorption and HACA as well as surface oxidation. It accommodates the four inception models widely used in literature. The third unit deals with particle dynamics using a MPBM and SPBM to describe the coagulation of soot particles that results in their fractal-like structure of particle with an evolving size distribution. A multi-step validation procedure is followed to assess reliability of each sub-model that entails the comparison of collision frequency kernels and agglomerate morphology of different particle dynamics models with those of detailed DEM simulations and establishing energy and elemental carbon, and hydrogen balances for all combinations of particle dynamics and inception models with reactors and flames. This tool can be applied to a variety of experimental targets with different temperature- and pressure-time-histories, residence times, and fuel compositions.

As the first step, methane pyrolysis in shock-tube will be simulated using a constant volume reactor. The species measurements will be used to assess the reaction mechanism and examine carbon conversion flux from fuel to smaller intermediates and larger hydrocarbons. The performance of inception models will be analyzed by comparing the predicted soot volume fraction, f_v and primary particle diameter, d_p , with data collected from extinction and TEM measurements. A sensitivity analysis will be conducted on model parameters to identify the determining factors that control both f_v and d_p in the framework of shock-tube simulations with short residence time ($t \approx 2$ ms) and high temperatures (≥ 2000 K). Then, the optimization of inception models will be performed by adjusting the rate constants within their physical limit (maximum collision rate of molecules) to minimize the prediction error for yield and morphology. This shows the range of inception flux that enables the prediction of soot yield and morphology represented by f_v and d_p , respectively in good agreement with measurements. Then, the optimized rates will be applied to shock-tube in a wide temperature range, and the f_v predictions will be compared with data. A similar investigation will be conducted for flow reactors using plug flow reactor model of omnisoot at lower temperatures and longer residence times and then the range of expected inception flux will be compared with those of shock-tubes.

Chapter 2

Theoretical Model

In this chapter, the mathematical basis for *omnisoot* is explained in the top-to-bottom hierarchical order. The highest level is the reactors and flames that include the transport equations of gas mixture and "soot variables". Soot formation source terms are handled by the particle dynamics model that mainly addresses particle size distribution (PSD), morphology and coagulation rate. The "PAH growth model" computes the contribution of inception and adsorption to source terms based on PAHs designated as precursors. Similarly, the "surface reactions" model obtains the surface growth and oxidation rate by HACA mechanism and passes them to the particle dynamics model.

2.1 Assumptions and conventions

Here, the main conventions and assumptions used in the derivation of the mathematical model are listed below.

1. The ideal gas law is used to calculate physical, transport, and chemical properties of gas mixture.
2. \dot{s}_k denotes that rate production/consumption of k_{th} species due to soot formation. It is positive when the species is released to gas mixture.
3. Each soot agglomerate consists of spherical monodisperse primary particles in point contact.
4. The word "*particle*" refers to soot both in spherical and agglomerate shape.
5. The density of soot is assumed constant at the value of 1800 kg/m^3 . This density represents an average between particles with large C/H ratio ($\rho = 1600 \text{ kg/m}^3$) and small C/H ratios ($\rho = 2000 \text{ kg/m}^3$) [131].
6. The incipient soot particles are 2 nm in diameter, so no particles could exist with a primary particle diameter smaller than 2 nm. The number of carbon atoms in the incipient soot particle is calculated from the mass of a sphere with the diameter of 2 nm assuming pure carbon content.

$$d_{p,min} = 2 \text{ nm}$$
$$n_{c,min} = \frac{\pi}{6} \rho_{soot} d_{p,min}^3 \frac{1}{MW_c} \approx 378. \quad (2.1)$$

7. The calculation of PAH adsorption and soot oxidation requires "*soot concentration*" which is defined as the number of soot agglomerates per unit volume of gas. The number density of agglomerates, N_{agg} , are tracked per unit mass of gas mixture i.e. #/kg_{gas}. So, soot concentration can be calculated by multiplying agglomerate number density by gas density as:

$$[\text{soot}] = \rho \cdot N_{agg}. \quad (2.2)$$

8. The specific heat, internal energy and enthalpy of soot are approximated by those of pure graphite, and employed to close the energy balance in the system [132].
9. Soot particles and gas are in thermal equilibrium during studied processes.
10. There is no temperature gradient within each agglomerate.
11. *Soot variable* refers to the features/properties of soot particles tracked by the particle dynamics model and used in the soot transport equations.
12. *PAH growth* is a unit of the soot model with a set of pathways that determine the rate of inception and adsorption from PAHs in the gas mixture.
13. *Surface reactions* is a unit of the soot model that describes the addition of acetylene to soot surface, and removal of carbon via oxidation by OH and O₂ in the HACA scheme.
14. The superscript *i* denotes section number of a soot variable or property derived from variables. In the case of the monodisperse model, the section number can be ignored because it is equivalent to the sectional model with one section.
15. The computation of morphological parameters (d_p , d_m , d_g , and n_p) and diffusion coefficient are done similarly by both particle dynamics models, so they are explained separately in standalone sections.

2.2 Flames

This research focuses on axisymmetric stagnation flames that reduces the laminar 3D steady-state equations in the cylindrical coordinates (r-z-θ system) to one-dimensional domain in r-z plane. The derivation of flow equations for a reacting gas is explained in detail in Section 7.2 of [133]. Here, the modified version of equations is provided that takes into account the mass and momentum of particles, the removal/release of species due to soot inception, surface growth and oxidation as well as the formation and sensible energy of soot.

Continuity:

$$\frac{\partial \rho u}{\partial z} = 2\rho V - \sum_j \dot{s}_j W_j, \quad (2.3)$$

Radial Momentum:

$$\rho u \frac{\partial V}{\partial z} + \rho V^2 = -\Lambda + \frac{\partial}{\partial z} \left(\mu \frac{\partial V}{\partial z} \right), \quad (2.4)$$

Energy:

$$[(1 - \varphi)\rho c_p + \varphi \rho_{soot} c_{p,soot}] u \frac{\partial T}{\partial z} = \frac{\partial}{\partial z} \left(K \frac{\partial T}{\partial z} \right) - \sum_k^{n_{sp}} J_k \frac{\partial h_k}{\partial z} - \sum_k^{n_{sp}} h_k W_k (\omega_k + \dot{s}_k), \quad (2.5)$$

Species:

$$\rho u \frac{\partial Y_k}{\partial z} = -\frac{\partial J_k}{\partial z} + W_j (\omega_k + \dot{s}_k), \quad (2.6)$$

where $V = v/r$ is the scaled radial velocity, and Λ is the pressure eigenvalue. J_k denotes the diffusive flux of species k that is usually calculated using mixture averaged formulation.

$$J_K^* = -\rho \frac{W_k}{W_T} D'_{km} \frac{\partial X_k}{\partial z}, \quad (2.7)$$

$$J_K = J_K^* - Y_k \sum_{l=1}^{n_{sp}} J_l^*, \quad (2.8)$$

The details of calculation of mixture-averaged diffusive fluxes can be found in the documentation of `GasTransport::getMixDiffCoeffs()`¹ in Cantera [130]. Soot variables are treated similar to species, and their transport equations can be written as:

$$\rho u \frac{\partial \psi}{\partial z} = -\frac{\partial}{\partial z} \left(D \frac{\partial \psi}{\partial z} \right) + S_\psi, \quad (2.9)$$

where ψ is a generic soot variable that represents any of tracked soot variables.

2.3 Reactors

Zero-dimensional reactors are ideal representation of reactive systems typical in industrial applications. The strong mixing in the volume/cross-section allows description of thermal/chemical/hydrodynamic properties of mixture with a single variable that evolves in time. This section entails the conservation equations of mass, momentum (if applicable), species, energy, and soot for three reactor models.

2.3.1 Constant Volume Reactor

This reactor assumes that the volume of system does not change during the process. In the absence of soot, this leads to gas with constant density. However, soot formation converts part of gaseous species to solid particles thereby affecting its volume and density. Note that, continuity, species and energy transport equations only track gas mixture properties. Figure 2.1 illustrates the control volume over the gas mixture targeted by mass and energy balance equations. Any mass converted to solid soot particles leaves the control volume. Mass and energy passes through the control surface around solid particles by soot formation processes. Reactor volume is the sum of volume of gas mixture and solid particles.

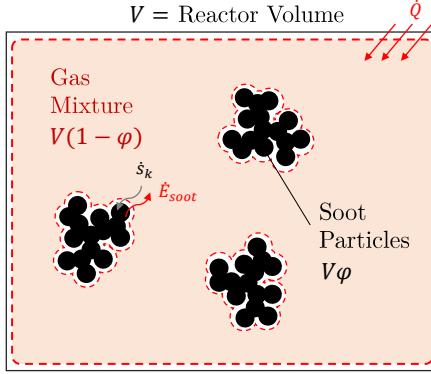


Figure 2.1: The schematics of control volume considered for the constant volume reactor that encompasses the gas mixture and excludes the soot particles. Mass and energy are transferred between gas and soot particles.

The continuity for this reactor can be written as:

$$\frac{d}{dt}(\rho(1 - \varphi)) = (1 - \varphi) \sum_i \dot{s}_i W_i. \quad (2.10)$$

Similarly, the species equation for species k is expressed as:

$$\frac{dY_k}{dt} = \frac{1}{\rho} (\dot{\omega}_k + \dot{s}_k) W_k - \frac{1}{\rho} Y_k \sum_i \dot{s}_i W_i. \quad (2.11)$$

¹https://cantera.org/documentation/docs-3.0/doxygen/html/d8/d58/classCantera_1_1GasTransport.html

The energy balance for the gas mixture can be simplified to the rate change of temperature. An external heat source of \dot{Q} is considered to account for possible heat loss/gain of the reactor.

$$\frac{dT}{dt} = \frac{1}{\rho(1-\varphi)c_v + \rho_{soot}\varphi c_{v,soot}} \left[-(1-\varphi) \sum_k e_k (\dot{\omega}_k + \dot{s}_k) W_k + u_{soot}(1-\varphi) \sum_k \dot{s}_k W_k + \frac{\dot{Q}}{V} \right]. \quad (2.12)$$

The transport equation for a generic soot variable, ψ can be written as:

$$\frac{d\psi}{dt} = S_\psi - \frac{\psi}{\rho} \sum_i \dot{s}_i W_i. \quad (2.13)$$

2.3.2 Perfectly Stirred Reactor

In this reactor, gas enters with a mass flow rate \dot{m}_{in} , composition of Y^* and temperature of T^* , instantaneously mixes and homogeneously reacts with the mixture resident inside the reactor. The reacting gas reaches a spatially uniform temperature and composition described by T , and Y . It is assumed that temperature, composition and soot properties of the outflow are the same as reactor. Figure 2.2 illustrates the schematics of PSR. \dot{m}_{in} and \dot{m}_{out} refer to inflow and outflow gas mass flow rates, respectively. Under no-soot conditions, the inlet and outlet mass flow rates are equal, but the gas mixture loses mass by soot formation, so \dot{m}_{out} is slightly less than \dot{m}_{in} . The pressure of reactor is assumed to stay constant during the process [133]. The nominal residence time of gas mixture in the reactor is defined as:

$$\tau = \frac{\rho V}{\dot{m}_{in}}. \quad (2.14)$$

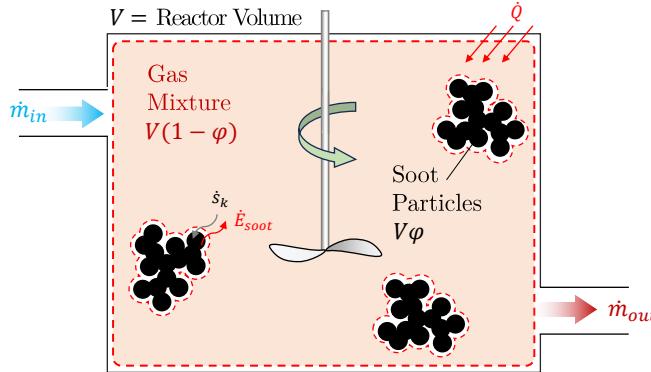


Figure 2.2: The schematics of control volume considered for the perfectly stirred reactor that encompasses the gas mixture and excludes the soot particles. Mass and energy are transferred between gas and soot particles. The inlet flow brings species and enthalpy into the control volume and the outflow discharges them. The outflow mass flow is less inflow mass flow due to soot formation.

The conservation of mass can be written for PSR by considering the mass flux of in- and outflow, and the removal of mass due to soot generation as:

$$\frac{dm}{dt} = \dot{m}_{in} - \dot{m}_{out} + (1 - \varphi) \sum_i \dot{s}_i W_i. \quad (2.15)$$

The density is not determined by solving the continuity equation, but rather from ideal gas law and assuming a constant pressure and the composition from solving the species transport equations as:

$$\frac{dY_k}{dt} = \frac{1}{\tau} (Y_k^* - Y_k) + \frac{1}{\rho} \left[(\dot{\omega}_k + \dot{s}_k) W_k - Y_k \sum_i \dot{s}_i W_i \right]. \quad (2.16)$$

The energy equation for this reactor is written as:

$$\frac{dT}{dt} = \frac{1}{\rho(1-\varphi)c_p + \rho_{soot}c_{p,soot}\varphi} \left[\frac{\dot{m}_{in}}{V} (h^* - h) - \frac{\dot{m}_{in}}{V} \sum_k (Y_k^* - Y_k) h_k \right. \\ \left. - (1-\varphi) \sum_k (\dot{\omega}_k + \dot{s}_k) W_k h_k + (1-\varphi) \sum_i \dot{s}_i W_i h_{soot} + \frac{\dot{Q}}{V} \right]. \quad (2.17)$$

The soot transport equations can also be expressed as:

$$\frac{d\psi}{dt} = \frac{\dot{m}_{in}}{\rho V (1-\varphi)} (\psi^* - \psi) + S_\psi - \frac{1}{\rho} \psi \sum_i \dot{s}_i W_i. \quad (2.18)$$

2.3.3 Plug Flow Reactor

The plug flow reactor (PFR) is an ideal representation of a channel or duct with a constant cross-sectional area where a steady-state one-dimensional flow changes temperature, composition, and soot properties along the channel. There is no spatial gradient over the cross-section due to strong mixing. Diffusion along the channel is negligible. The pressure is assumed constant along the reactor.

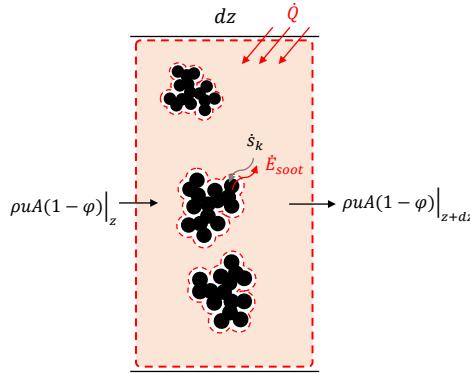


Figure 2.3: The schematics of control volume for a differential element along PFR that includes the gas mixture and excludes the soot particles considering wall heat transfer. Mass and energy are transferred between gas and soot.

The continuity equation for PFR is written as:

$$\frac{d}{dz}(\rho u(1-\varphi)) = (1-\varphi) \sum_i \dot{s}_i W_i. \quad (2.19)$$

The momentum equation can also be established as:

$$u(1-\varphi) \sum_i \dot{s}_i W_i + \rho u(1-\varphi) \frac{du}{dz} = -\frac{d}{dz}(P(1-\varphi)) - \frac{\tau_w}{R_H}. \quad (2.20)$$

where τ_w is the wall shear the can be determined from fraction factor, f as:

$$\tau_w = \frac{1}{2} \rho u^2 f. \quad (2.21)$$

The friction factor, f can be calculated with a good accuracy for the entire range of Reynolds number, Re , from laminar to turbulent flow using the explicit formula given by Haaland [134]:

$$\frac{1}{f^{1/2}} = -1.8 \log \left(\frac{6.9}{Re} + \left[\frac{\epsilon/D_H}{3.7} \right]^{1.11} \right), \quad (2.22)$$

where ϵ is the roughness of reactor wall. R_H and D_H are hydraulic radius and diameter, respectively that can be determined from cross-section geometry of reactor as:

$$D_H = 4R_H = \frac{4A_c}{P_c}, \quad (2.23)$$

A_c and P_c are cross-sectional area and wetted perimeter of the reactor. The species equation can be expressed as:

$$\frac{dY_k}{dz} = \frac{1}{\rho u} \left[(\dot{\omega}_k + \dot{s}_k) W_k - Y_k \sum_i \dot{s}_i W_i \right]. \quad (2.24)$$

The energy equation can be expressed as:

$$\frac{dT}{dz} = \frac{1}{\rho u (1 - \varphi) c_p + \rho_{soot} u \varphi c_{p, soot}} \left[-(1 - \varphi) \sum_k h_k (\dot{\omega}_k + \dot{s}_k) W_k + h_{soot} (1 - \varphi) \sum_k \dot{s}_k W_k + \dot{q}' \right]. \quad (2.25)$$

The soot transport equations can also be written as:

$$\frac{d\psi}{dz} = \frac{S_\psi}{u} - \frac{\psi}{\rho u} \sum_i \dot{s}_i W_i. \quad (2.26)$$

2.4 Particle Dynamics

Population balance models rely on the Eulerian description of particles where bulk properties of particle population such as number density, mass or surface area are treated as continuous quantities and tracked by solving scalar transport equations. These methods are computationally cheaper compared with mesoscale models such as DEM, and can be easily interfaced with chemical kinetics in CFD solvers to simulate soot formation in turbulent configurations. Here, we use two particle dynamics models: a monodisperse population balance model (MPBM) based on four variables leading to 4 transport equations in total, and a fixed sectional population balance model (SPBM) tracking three variables per section. The total number of transport equations in the sectional model is determined by the number of sections and number of equations solved per section. The first two/three variables in the MPBM/SPBM enables description of number, mass, and evolving fractal-like morphology of soot agglomerates that are necessary to accurately predict collision frequency of agglomerates [135] as well as oxidation and surface growth rates [79]. The last variable tracks the number of hydrogen atoms in agglomerates that allows the model to capture the soot composition, thereby its maturity [45], and surface reactivity [86]. The tracked variables are used to address particle dynamics that includes (i) reconstructing particles morphology by determining characteristic diameters from tracked soot variables, (ii) calculating collision frequency and coagulation source term, (iii) combining the contribution of inception, PAH adsorption, surface growth and oxidation into source terms. First, common features of both particle dynamics models are reviewed. As mentioned before, any parameter with superscript i denotes the section number, which can be ignored/dropped for the MPBM that only has one section. For example, d_m^i can be replaced with d_m .

2.4.1 Soot Morphology

The evolving fractal-like structure of agglomerates is quantified by their mobility diameter normalized by primary particle, d_m/d_p , and gyration, d_m/d_g , diameters that can be described with

power-laws derived from mesoscale simulations. Incipient soot is initially a sphere formed of PAHs with constant density that grows in size by surface reactions and forms agglomerates by coagulation. The collision frequency of particles depends on their evolving fractal-like structure [135]. Some simplifying assumptions are made to reconstruct the particle morphology from tracked variables. The primary particles of each agglomerate are similar enough that can be described by mean size and composition. They also stay in point contact during surface growth and agglomeration i.e. the necking is ignored. A universal fractal dimension, $D_f = 1.9$ is used for agglomerates larger than sphere [136]. Mobility and gyration diameters are the diameter of a sphere with the same translational and rotational properties of the agglomerate, respectively. The employed power-laws have been shown to describe the morphology of soot from premixed [80], diffusion [137] flames, and diesel engines [138]. Figure 2.4 illustrates the schematics of a soot agglomerate with 12 primary particles and depicted d_p , d_m , and d_g .

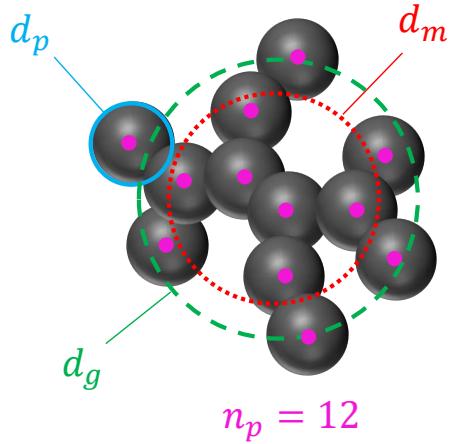


Figure 2.4: The schematics of a soot agglomerate with 12 primary particles ($n_p = 12$). Primary particle (d_p), mobility (d_m), and gyration (d_g) are shown.

n_p^i is the number of primary particles per agglomerate for i^{th} section that can be obtained by dividing the number density of primary particles by the number density of agglomerates of that section as:

$$n_p^i = \frac{N_{pri}^i}{N_{agg}^i}. \quad (2.27)$$

Primary particle diameter, d_p^i , can be obtained from total carbon content and number density of primary particles using

$$d_p^i = \left(\frac{6}{\pi} \frac{C_{tot}^i \cdot W_{carbon}}{\rho_{soot}} \frac{1}{N_{pri}^i \cdot Av} \right)^{1/3}. \quad (2.28)$$

The DEM-derived power-laws [66] relate d_m^i and d_g^i to d_p^i and n_p^i as

$$d_m^i = d_p^i \cdot n_p^i^{0.45}, \quad (2.29)$$

$$d_g^i = \begin{cases} d_m^i / (n_p^i^{-0.2} + 0.4), & \text{if } n_p^i > 1.5 \\ d_m^i / 1.29. & \text{if } n_p^i \leq 1.5 \end{cases} \quad (2.30)$$

The collision diameter, d_c^i is the maximum of d_m^i , d_g^i :

$$d_c^i = \max(d_m^i, d_g^i) \quad (2.31)$$

d_m^i , d_g^i , d_c^i are used to calculate the source terms due to the surface growth, oxidation, PAH adsorption and coagulation. The volume equivalent diameter, d_v^i , is the diameter of the sphere with the same mass as agglomerate, and it is obtained as:

$$d_v^i = d_p^i \cdot n_p^i{}^{1/3} \quad (2.32)$$

The primary particle surface area is calculated from d_p^i assuming spherical primary particles.

$$A_p^i = \pi d_p^i{}^2, \quad (2.33)$$

A_{tot}^i (for each section) is defined as the total surface area of soot particles per unit mass of gas mixture obtained as

$$A_{tot}^i = N_{pri}^i \cdot Av \cdot A_p^i. \quad (2.34)$$

2.4.2 Diffusion of soot particles

The diffusion coefficient of soot particle, D^i , is calculated as

$$D^i = \frac{k_B T}{f^i}, \quad (2.35)$$

where f^i is the friction factor of particles in gas,

$$f^i = \frac{3\pi\mu d_m^i}{C^i(d_m^i)}, \quad (2.36)$$

where C^i is the Cunningham function that corrects the friction factor given a diameter in the continuum regime for transition and free molecular regimes as:

$$C^i(d) = 1 + \frac{2\lambda}{d} \left(1.21 + 0.4 \cdot \exp\left(\frac{-0.78d}{\lambda}\right) \right), \quad (2.37)$$

where λ is the mean free path of gas given as:

$$\lambda = \frac{\mu}{\rho} \sqrt{\frac{\pi W_{gas}}{2k_B AvT}}. \quad (2.38)$$

Note that, λ is a property of the gas mixture that does not depend on particle morphology and size section. The mean velocity, c^i and mean stop distance of particles, λ_a^i can be calculated as:

$$c^i = \sqrt{\frac{8k_B T}{\pi m_{agg}^i}}. \quad (2.39)$$

$$\lambda_a = \frac{8D^i}{\pi c^i}. \quad (2.40)$$

The mean distance of particles are also calculated as:

$$\delta_a^i = \frac{1}{d_c^i \lambda_a^i} \left[(d_c^i + \lambda_a^i)^3 - (d_c^i{}^2 + \lambda_a^i{}^2)^{3/2} \right] - d_{c,j}. \quad (2.41)$$

2.4.3 Soot Composition

The composition of soot is characterized by their elemental carbon to hydrogen ratio (C/H) is a measure of soot maturity and increases from C/H < 2 for incipient soot [139] to 2 < C/H < 10 for nascent soot [140] and C/H > 20 for mature soot [92]. The soot agglomerates are assumed to have pure carbon graphitic core [45] with all hydrogen atoms on the surface [86]. C/H ratio can be obtained from total carbon and hydrogen content as:

$$\left(\frac{C}{H}\right)^i = \frac{C_{tot}^i}{H_{tot}^i}. \quad (2.42)$$

The carbon content of each agglomerate is a predefined parameter in the SPBM (depending on the section the agglomerate is placed), but it can be calculated from dividing C_{tot} by N_{agg} for the MPBM. The hydrogen content of each agglomerate is calculated for both particle dynamics models as:

$$H_{agg}^i = \frac{H_{tot}^i}{N_{agg}^i}. \quad (2.43)$$

2.4.4 Monodisperse Population Balance Model

The MPBM used in this research tracks the number density of primary particles (N_{pri}) and agglomerates (N_{agg}), total carbon (C_{tot}) and hydrogen (H_{tot}) content of soot particles per unit mass of gas mixture. The morphological parameters such as primary particle, mobility and gyration diameters obtained from these soot variables are the average values for the population.

2.4.4.1 Coagulation

Coagulation is the process during which solid and hard soot particles collide and attach at point of contact leading to larger agglomerates. This process conserves the soot mass and composition and number density of primary particles, so coagulation only affects N_{agg} . I_{coag}^N accounts for the decay rate of N_{agg} by the binary collision of soot particles by

$$I_{coag} = -\frac{1}{2}\beta N_{agg}^2, \quad (2.44)$$

where β is the collision frequency of agglomerates for the free molecular ($Kn > 10$) to continuum regime ($Kn < 0.1$). The value of β in the transition regime ($0.1 < Kn < 10$) can be calculated from the harmonic mean of the continuum (β_{cont}) and free molecular (β_{fm}) regime values. Additionally, an enhancement factor of %82 is applied to take into account the effect of polydispersity [67] as:

$$\beta = 1.82 \frac{\beta_{fm} \beta_{cont}}{\beta_{fm} + \beta_{cont}}, \quad (2.45)$$

$$\beta_{fm} = 4 \sqrt{\frac{\pi k_b T}{m_{agg}}} d_c^2, \quad (2.46)$$

$$\beta_{cont} = 8\pi d_c D. \quad (2.47)$$

Alternatively, β can be obtained using Fuchs interpolation [141] as:

$$\beta = \beta_{cont} \left(\frac{d_c}{d_c + 2\sqrt{2}\delta} + \frac{8D}{\sqrt{2}c_r d_c} \right)^{-1}. \quad (2.48)$$

2.4.4.2 Source terms

The source terms of tracked variables combines the effect of the inception, PAH adsorption, surface growth and oxidation and coagulation.

$$S_{N_{agg}} = \frac{I_{N,inc}}{n_{c,min}} + I_{coag}. \quad (2.49)$$

$$S_{N_{pri}} = \frac{I_{N,inc}}{n_{c,min}}. \quad (2.50)$$

$$S_{C_{tot}} = I_{C_{tot},inc} + I_{C_{tot},gr} + I_{C_{tot},ads} + I_{C_{tot},ox}. \quad (2.51)$$

$$S_{H_{tot}} = I_{H_{tot},inc} + I_{H_{tot},gr} + I_{H_{tot},ads} + I_{H_{tot},ox}. \quad (2.52)$$

The partial source terms in Equations 2.49-2.52 denoted by I are determined by PAH growth and surface reaction model explained in Sections 2.6 and 2.5, respectively.

2.4.5 Sectional Population Balance Model

A SPBM with the fixed pivot is used to describe particle dynamics [142]. The mass range of particles are divided into discrete sections each of which includes agglomerates of the same mass. Inception introduces new particles to the first section with the mass corresponding to the incipient particle. The particles of first section can migrate to upper sections by gaining mass via surface growth and coagulation, and return to lower sections when they lose mass through oxidation. The mass of sections is determined by a geometric progression with a scale factor equal to the mass of incipient soot particle, and a common ratio of SF, known as sectional spacing factor. The mass of each section is approximated by the carbon content of agglomerates in moles as:

$$C_{agg}^i = \frac{n_{c,min}}{Av} \cdot SF^{(i-1)}. \quad (2.53)$$

where (i-1) represents the exponent of SF. The mass of hydrogen is ignored in the placement of agglomerates in the sections. The total number density of agglomerates, N_{agg}^i [mol/kg], and primary particles, N_{pri}^i [mol/kg] are tracked for each section. Morphological parameters are determined for each section according to the equations in Section 2.4.1.

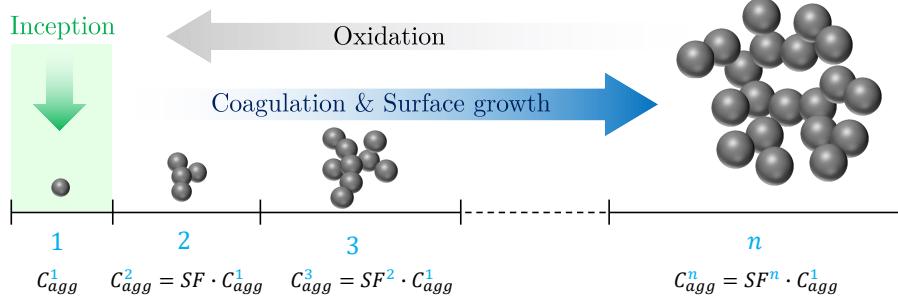


Figure 2.5: The illustration of sections of SPBM. The mass of sections grows progressively by the scale factor of SF. Inception introduces new particles to the first section that propagate to the upper section via coagulation and surface growth and return to lower sections by oxidation

2.4.5.1 Coagulation

In SPBM approach, collisions between particles from every two sections are considered. The new particles formed by coagulation are placed in an upper section with the mass equal to sum of mass of particles involved in the collision. When the mass of yielded particle lies between two consecutive sections, the particles are divided among these sections proportional to their mass. One possible scenario is that the mass of the newly formed particle is greater than the last section, thus leaving tracked mass range. Losing mass is a potential problem with the fixed pivot sectional model, which can be avoided by selecting proper number of sections and spacing factor to ensure the last sections stay empty during the simulation.

The collision frequency between sections j and k can be obtained from the harmonic mean of the values in the continuum and free molecular regimes as:

$$\beta^{jk} = \frac{\beta_{fm}^{jk} \beta_{cont}^{jk}}{\beta_{fm}^{jk} + \beta_{cont}^{jk}}, \quad (2.54)$$

$$\beta_{fm}^{jk} = \sqrt{\frac{\pi k_B T}{2} \left(\frac{1}{m_{agg}^j} + \frac{1}{m_{agg}^k} \right) (d_c^j + d_c^k)^2}, \quad (2.55)$$

$$\beta_{cont}^{ij} = \frac{2k_B T}{3\mu} \left(\frac{C^j}{d_m^j} + \frac{C^k}{d_m^k} \right) (d_c^j + d_c^k)^2. \quad (2.56)$$

The collision frequency can also be determined from the Fuchs interpolation similar to the MPBM as:

$$\beta^{jk} = \beta_{cont}^{ij} \left[\frac{d_c^j + d_c^k}{d_c^j + d_c^k + 2 + \delta_r^{jk}} + \frac{8(D^j + D^k)}{\bar{c}_r^{jk} (d_c^j + d_c^k)} \right]^{-1}, \quad (2.57)$$

where δ_r^{jk} and \bar{c}_r^{jk} are the mean square root of mean distance and velocity of particles, respectively.

$$\delta_r^{jk} = \sqrt{\delta_a^{j2} + \delta_a^{k2}}, \quad (2.58)$$

$$\bar{c}_r^{jk} = \sqrt{c^{j2} + c^{k2}}. \quad (2.59)$$

Coagulation redistributes the total number of agglomerates and primary particles as well as hydrogen atoms among the sections. The partial coagulation source terms for N_{agg}^i , N_{pri}^i and H_{tot}^i can be calculated as:

$$I_{N_{agg}}^i = \sum_{k=1}^{n_{sec}} \sum_{j=k}^{n_{sec}} \left(1 - \frac{\delta_{jk}}{2} \right) \eta_{ijk} \beta^{jk} N_{agg}^j N_{agg}^k - N_{agg}^i \sum_{k=1}^{n_{sec}} \beta^{im} N_{agg}^m. \quad (2.60)$$

$$I_{N_{pri}}^i = \sum_{k=1}^{n_{sec}} \sum_{j=k}^{n_{sec}} \left(1 - \frac{\delta_{jk}}{2} \right) \eta_{p,ijk} \eta_{ijk} \beta^{jk} N_{agg}^j N_{agg}^k - N_{pri}^i \sum_{k=1}^{n_{sec}} \beta^{im} N_{agg}^m. \quad (2.61)$$

$$I_{H_{tot}}^i = \sum_{k=1}^{n_{sec}} \sum_{j=k}^{n_{sec}} \left(1 - \frac{\delta_{jk}}{2} \right) \eta_{h,ijk} \eta_{ijk} \beta^{jk} N_{agg}^j N_{agg}^k - H_{tot}^i \sum_{k=1}^{n_{sec}} \beta^{im} N_{agg}^m. \quad (2.62)$$

where δ_{jk} is the Kronecker delta defined as:

$$\delta_{jk} = \begin{cases} 1, & \text{if } j = k \\ 0, & \text{if } j \neq k \end{cases} \quad (2.63)$$

In Equation (2.60), η_{ijk} assigns newly formed agglomerates to the two consecutive section in order to conserves mass during coagulation [71].

$$\eta_{ijk} = \begin{cases} \frac{C_{agg}^{i+1} - C_{agg}^{jk}}{C_{agg}^{i+1} + C_{agg}^{jk}}, & \text{if } C_{agg}^i \leq C_{agg}^{jk} < C_{agg}^{i+1} \\ \frac{C_{agg}^i - C_{agg}^{jk}}{C_{agg}^i + C_{agg}^{i-1}}, & \text{if } C_{agg}^{i-1} \leq C_{agg}^{jk} < C_{agg}^i \\ 0 & \text{else} \end{cases} \quad (2.64)$$

where $C_{agg}^{jk} = C_{agg}^j + C_{agg}^k$. Similarly, $\eta_{p,ijk}$ in Equation (2.61) and $\eta_{h,ijk}$ in Equation (2.62) adjust the number primary particles and hydrogen atoms added to consecutive sections based on their mass, respectively.

$$\eta_{p,ijk} = \frac{C_{agg}^i}{C_{agg}^{jk}} (n_p^j + n_p^k), \quad (2.65)$$

$$\eta_{h,ijk} = \frac{C_{agg}^i}{C_{agg}^{jk}} (H_{agg}^j + H_{agg}^k), \quad (2.66)$$

2.4.5.2 Source terms

The source terms are split into four parts showing the contribution of different soot formation and evolution factors. The effect of surface growth and PAH adsorption are combined (denoted by the subscript gr,ads) because they are similar mass-gaining mechanisms.

$$S_{N_{agg}} = (S_{N_{agg}})_{inc} + (S_{N_{agg}})_{gr,ads} + (S_{N_{agg}})_{ox} + (S_{N_{agg}})_{coag}, \quad (2.67)$$

$$S_{N_{pri}} = (S_{N_{pri}})_{inc} + (S_{N_{pri}})_{gr,ads} + (S_{N_{pri}})_{ox} + (S_{N_{pri}})_{coag}, \quad (2.68)$$

$$S_{H_{tot}} = (S_{H_{tot}})_{inc} + (S_{H_{tot}})_{gr,ads} + (S_{H_{tot}})_{ox} + (S_{H_{tot}})_{coag}. \quad (2.69)$$

Inception introduces equal number of agglomerates and primary particles to the first section.

$$(S_{N_{agg}})_{inc} = \frac{1}{Av} \frac{I_{N,inc}}{C_{agg}^i}, \quad i = 1. \quad (2.70)$$

$$(S_{N_{pri}})_{inc} = \frac{1}{Av} \frac{I_{N,inc}}{C_{agg}^i}, \quad i = 1. \quad (2.71)$$

$$(S_{H_{tot}})_{inc} = I_{H,inc}, \quad i = 1. \quad (2.72)$$

Surface growth and PAH adsorption increase the (carbon) mass and hydrogen content of agglomerates, and transfer them to upper sections. The removal rate of agglomerates (N_{agg}^i) from the original section due to surface growth and PAH adsorption must be equal to the addition rate of agglomerates to the target section to conserve the mass, and it is calculated by dividing the mass growth rate by the difference of the mass of the adjacent sections.

$$(S_{N_{agg}})_{gr,ads} = \frac{1}{Av} \begin{cases} -\frac{I_{C_{tot},gr}^i + I_{C_{tot},ads}^i}{C_{agg}^{i+1} - C_{agg}^i} & \text{if } i = 1 \\ \frac{I_{C_{tot},gr}^{i-1} + I_{C_{tot},ads}^{i-1}}{C_{agg}^i - C_{agg}^{i-1}} - \frac{I_{C_{tot},gr}^i + I_{C_{tot},ads}^i}{C_{agg}^{i+1} - C_{agg}^i} & \text{if } 1 < i < n_{sec} \\ \frac{I_{C_{tot},gr}^{i-1} + I_{C_{tot},ads}^{i-1}}{C_{agg}^i - C_{agg}^{i-1}} & \text{if } i = n_{sec} \end{cases} \quad (2.73)$$

As agglomerates move up/down through sections, they carry the number of primary particles as well as hydrogen atoms, so the transfer rate of agglomerates is multiplied by n_p^i and H_{agg}^i , respectively.

$$(S_{N_{pri}})_{gr,ads} = \frac{1}{Av} \begin{cases} -\frac{I_{C_{tot},gr}^i + I_{C_{tot},ads}^i}{C_{agg}^{i+1} - C_{agg}^i} & \text{if } i = 1 \\ \frac{I_{C_{tot},gr}^{i-1} + I_{C_{tot},ads}^{i-1}}{C_{agg}^i - C_{agg}^{i-1}} n_p^{i-1} - \frac{I_{C_{tot},gr}^i + I_{C_{tot},ads}^i}{C_{agg}^{i+1} - C_{agg}^i} n_p^i & \text{if } 1 < i < n_{sec} \\ \frac{I_{C_{tot},gr}^{i-1} + I_{C_{tot},ads}^{i-1}}{C_{agg}^i - C_{agg}^{i-1}} n_p^{i-1} & \text{if } i = n_{sec} \end{cases} \quad (2.74)$$

$$(S_{H_{tot}})_{gr,ads} = \frac{1}{Av} \begin{cases} -\frac{I_{C_{tot},gr}^i + I_{C_{tot},ads}^i}{C_{agg}^{i+1} - C_{agg}^i} H_{agg}^i + I_{H_{tot},gr}^i + I_{H_{tot},ads}^i & \text{if } i = 1 \\ \frac{I_{C_{tot},gr}^{i-1} + I_{C_{tot},ads}^{i-1}}{C_{agg}^i - C_{agg}^{i-1}} H_{agg}^{i-1} - \frac{I_{C_{tot},gr}^i + I_{C_{tot},ads}^i}{C_{agg}^{i+1} - C_{agg}^i} H_{agg}^i + I_{H_{tot},gr}^i + I_{H_{tot},ads}^i & \text{if } 1 < i < n_{sec} \\ \frac{I_{C_{tot},gr}^{i-1} + I_{C_{tot},ads}^{i-1}}{C_{agg}^i - C_{agg}^{i-1}} H_{agg}^{i-1} + I_{H_{tot},gr}^i + I_{H_{tot},ads}^i & \text{if } i = n_{sec} \end{cases} \quad (2.75)$$

Similarly, the agglomerates lose (carbon) mass by oxidation, and descend to the lower sections carrying primary particle and hydrogen.

$$(S_{N_{agg}})_{ox} = \frac{1}{Av} \begin{cases} \frac{I_{C_{tot},ox}^{i+1}}{C_{agg}^{i+1} - C_{agg}^i} - \frac{I_{C_{tot},ox}^i}{C_{agg}^i} & \text{if } i = 1 \\ \frac{I_{C_{tot},ox}^{i+1}}{C_{agg}^{i+1} - C_{agg}^i} - \frac{I_{C_{tot},ox}^i}{C_{agg}^i - C_{agg}^{i-1}} & \text{if } 1 < i < n_{sec} \\ -\frac{I_{C_{tot},ox}^i}{C_{agg}^i - C_{agg}^{i-1}} & \text{if } i = n_{sec} \end{cases} \quad (2.76)$$

$$(S_{N_{pri}})_{ox} = \frac{1}{Av} \begin{cases} \frac{I_{C_{tot},ox}^{i+1}}{C_{agg}^{i+1} - C_{agg}^i} n_p^{i+1} - \frac{I_{C_{tot},ox}^i}{C_{agg}^i} & \text{if } i = 1 \\ \frac{I_{C_{tot},ox}^{i+1}}{C_{agg}^{i+1} - C_{agg}^i} n_p^{i+1} - \frac{I_{C_{tot},ox}^i}{C_{agg}^i - C_{agg}^{i-1}} n_p^i & \text{if } 1 < i < n_{sec} \\ -\frac{I_{C_{tot},ox}^i}{C_{agg}^i - C_{agg}^{i-1}} n_p^i & \text{if } i = n_{sec} \end{cases} \quad (2.77)$$

$$(S_{H_{tot}})_{ox} = \frac{1}{Av} \begin{cases} \frac{I_{C_{tot},ox}^{i+1}}{C_{agg}^{i+1} - C_{agg}^i} H_{agg}^{i+1} - \frac{I_{C_{tot},ox}^i}{C_{agg}^i} H_{agg}^i + I_{H_{tot},ox}^i & \text{if } i = 1 \\ \frac{I_{C_{tot},ox}^{i+1}}{C_{agg}^{i+1} - C_{agg}^i} H_{agg}^{i+1} - \frac{I_{C_{tot},ox}^i}{C_{agg}^i - C_{agg}^{i-1}} H_{agg}^i + I_{H_{tot},ox}^i & \text{if } 1 < i < n_{sec} \\ -\frac{I_{C_{tot},ox}^i}{C_{agg}^i - C_{agg}^{i-1}} H_{agg}^i + I_{H_{tot},ox}^i & \text{if } i = n_{sec} \end{cases} \quad (2.78)$$

2.5 Surface reactions model

The heterogeneous surface reactions are described by hydrogen-abstraction-acetylene-addition (HACA) mechanism [40, 58]. The soot growth in HACA scheme is based on a sequential process similar to PAH growth. The hydrogenated arm-chair sites ($C_{soot} - H$) on the edge of aromatic rings are dehydrogenated by H abstraction forming C_{soot}° that bonds with C_2H_2 resulting in an additional aromatic ring with hydrogenated site. These sites can also be attacked by O_2 or OH . The elementary reactions that describe this sequential process are listed in Table 2.1. The rate of mass growth by HACA is obtained from the reaction of C_2H_2 with dehydrogenated sites as:

$$\omega_{gr}^i = \alpha^i k_{f4}[C_2H_2][C_{soot}^\circ]. \quad (2.79)$$

where k_{f4} denotes the forward rate of Reaction 2.1 in Table 2.1, and $[C_{soot}^\circ]$ is obtained by multiplying the surface density of dehydrogenated sites, χ_{soot}° with total surface area of soot (per unit of mass of gas mixture) as:

$$[C_{soot}^i] = \frac{\rho}{Av} A_{tot}^i \cdot \chi_{soot}^i. \quad (2.80)$$

χ_{soot}^i is calculated by assuming the steady-state for $[C_{soot}^i]$ in the system of reactions in Table 2.1:

$$\chi_{soot}^i = \frac{k_f1[H] + k_f2[OH]}{k_{r1}[H_2] + k_{r2}[H_2O] + k_{f3}[H] + k_{f4}[C_2H_2] + k_{f5}[O_2] + k_{f1}[H] + k_{f2}[OH]} \chi_{soot_{CH}}, \quad (2.81)$$

where $\chi_{soot_{CH}}$ is the surface density of hydrogenated sites estimated based on the assumption that soot surface is covered with stacks of benzene rings [58]. Considering the stack spacing of 3.15 Å and 2 C–H bonds per benzene ring length, the surface density of hydrogenated sites, χ_{soot-H} , is calculated to be $0.23 \text{ site}/\text{\AA}^2 = 2.3 \times 10^{19} \text{ site}/\text{m}^2$, which gives the maximum theoretical limit of the reaction sites.

In Equation (2.79), α is the surface reactivity factor between 0 and 1 that represents the decline of reaction sites from the theoretical limit due to particle aging, growth and maturity [143, 144], and it has been observed to depend on temperature time history [145, 60]. The value of α has been described using constant target-specific values as well as empirical equations based on particle size and flame temperature. A detailed review of these can be found in the chapter 4 of [146]. Here, the empirical equations proposed by Appel et al. [58] is used to calculate α :

$$\alpha^i = \tanh \left(\frac{12.56 - 0.00563 \cdot T}{\log_{10} \left(\frac{\rho_{soot} \cdot Av}{W_{carbon}} \frac{\pi}{6} d_p^3 \right)} - 1.38 + 0.00068 \cdot T \right). \quad (2.82)$$

Alternatively, Blanquart and Pitsch [85] related α to the number of surface hydrogen atoms on the soot particles.

$$\alpha^i = \frac{H_{tot}^i}{C_{tot}^i}. \quad (2.83)$$

The contribution of HACA to growth source terms can be computed from HACA rates considering the number of carbon atoms in C_2H_2 and number of arm-chair and zig-zag hydrogenated sites on soot particle [86] using

$$I_{C_{tot},gr}^i = 2\omega_{gr}^i / \rho, \quad (2.84)$$

$$I_{H_{tot},gr}^i = 0.25\omega_{gr}^i / \rho. \quad (2.85)$$

The rate change of C_2H_2 concentration due to mass growth is written as:

$$\left(\frac{d[C_2H_2]}{dt} \right)_{gr} = - \sum_{i=1}^{n_{sec}} \omega_{gr}^i. \quad (2.86)$$

The rate of release of H radicals into the gas mixture due to surface growth is:

$$\left(\frac{d[H]}{dt} \right)_{gr} = 1.75 \sum_{i=1}^{n_{sec}} \omega_{gr}^i. \quad (2.87)$$

The carbons on the surface of soot are oxidized via reaction with O_2 and OH which decreases total carbon of soot and releases CO and H_2 to gas mixture. The oxidation process is described by HACA mechanism. O_2 and OH oxidation rates are calculated as

$$\omega_{ox,O_2}^i = \alpha^i k_{f5}[O_2][C_{soot}^i], \quad (2.88)$$

$$\omega_{ox,OH}^i = \alpha^i k_{f6}[OH][soot^i]. \quad (2.89)$$

The oxidation source term is calculated considering the number of carbon atoms removed from soot through each oxidation pathway by

Table 2.1: Rate coefficients for the various surface reactions in Arrhenius form $k = AT^n \cdot e^{-E/RT}$

No.	Reaction		A $\left[\frac{\text{m}^3}{\text{mol} \cdot \text{s}} \right]$	n	$\frac{E}{R} [\text{K}]$
(R 2.1)	$\text{C}_{\text{soot}-\text{H}} + \text{H} \rightleftharpoons \text{C}_{\text{soot}^\circ} + \text{H}_2$	f	4.17×10^7	0	6542.52
		r	3.9×10^6	0	5535.98
(R 2.2)	$\text{C}_{\text{soot}-\text{H}} + \text{OH} \rightleftharpoons \text{C}_{\text{soot}^\circ} + \text{H}_2\text{O}$	f	10^4	0.734	719.68
		r	3.68×10^2	1.139	8605.94
(R 2.3)	$\text{C}_{\text{soot}^\circ} + \text{H} \longrightarrow \text{C}_{\text{soot}} + \text{H}_2\text{O}$	f	10^4	0.734	719.68
(R 2.4)	$\text{C}_{\text{soot}^\circ} + \text{C}_2\text{H}_2 \longrightarrow \text{C}_{\text{soot}-\text{H}}$	f	80	1.56	1912.43
(R 2.5)	$\text{C}_{\text{soot}^\circ} + \text{O}_2 \longrightarrow 2 \text{CO}$	f	2.2×10^6	0	3774.53
(R 2.6)	$\text{C}_{\text{soot}-\text{H}} + \text{OH} \longrightarrow \text{CO} + \frac{1}{2} \text{H}_2$	f	1.3×10^7	0	0

$$I_{C_{tot},ox}^i = -(2\omega_{ox,O_2}^i + \omega_{ox,OH}^i)/\rho, \quad (2.90)$$

We assume that oxidation does not change the number of surface hydrogen atoms. The rate change of concentration of CO, H and OH by oxidation is calculated as:

$$\left(\frac{d[\text{CO}]}{dt} \right)_{ox} = \sum_{i=1}^{n_{sec}} \omega_{ox,O_2}^i. \quad (2.91)$$

$$\left(\frac{d[\text{O}_2]}{dt} \right)_{ox} = - \sum_{i=1}^{n_{sec}} \omega_{ox,O_2}^i. \quad (2.92)$$

$$\left(\frac{d[\text{OH}]}{dt} \right)_{ox} = - \sum_{i=1}^{n_{sec}} \omega_{ox,OH}^i. \quad (2.93)$$

$$\left(\frac{d[\text{H}_2]}{dt} \right)_{ox} = \frac{1}{2} \sum_{i=1}^{n_{sec}} \omega_{ox,OH}^i. \quad (2.94)$$

2.6 PAH growth models

Here, four different PAH growth models are used to describe the conversion of PAHs to incipient particles and their adsorption on existing agglomerates. As mentioned before, the soot inception and surface growth is not fully understood yet, but there is substantial evidence to support the collision of PAHs as a key step in inception and surface growth [27, 28, 29]. So, global inception models have been developed based PAH collision consisting of different pathways with single- or multi-step reactions. The collision frequency of gaseous species including PAH molecules and polymers depend on their mass and diameter, and it is obtained as:

$$\beta_{dim_{jk}} = 2.2 \cdot d_r^2 \sqrt{\frac{8\pi k_B T}{m_r}} \quad (2.95)$$

where d_r and m_r are reduced diameter and mass for two PAH molecules, respectively.

$$d_r = 2 \frac{d_{PAH_k} \cdot d_{PAH_j}}{d_{PAH_k} + d_{PAH_j}} \quad (2.96)$$

$$m_r = \frac{m_{PAH_k} \cdot m_{PAH_j}}{m_{PAH_k} + m_{PAH_j}} \quad (2.97)$$

The mass of each PAH molecule is obtained from its molecular weight as:

$$m_{PAH_j} = \frac{W_{PAH_j}}{Av} \quad (2.98)$$

The diameter of PAH is calculated from its mass and density.

$$d_{PAH_j} = \left(\frac{6 \cdot m_{PAH_j}}{\pi \cdot \rho_{PAH_j}} \right)^{1/3} \quad (2.99)$$

The density of a PAH molecule is estimated using the relation proposed by Johansson et al. [147].

$$\rho_{PAH_j} = 171943.5197 \frac{W_{carbon} \cdot n_{C,PAH_j} + W_{hydrogen} \cdot n_{H,PAH_j}}{n_{C,PAH_j} + n_{H,PAH_j}} \quad (2.100)$$

The collision frequency of PAH_j and soot agglomerates in each section can be determined for the entire regime by harmonic mean of the collision frequency in the free molecular and continuum regimes as:

$$\beta_{ads_j}^i = \frac{\beta_{fm,ads}^i \cdot \beta_{cont,ads}^i}{\beta_{fm,ads}^i + \beta_{cont,ads}^i} \quad (2.101)$$

$$\beta_{fm,ads_j}^i = 2.2 \sqrt{\frac{\pi k_B T}{2} \left(\frac{1}{m_{agg}^i} + \frac{1}{m_{PAH_j}} \right) (d_g^i + d_{PAH_j})^2} \quad (2.102)$$

$$\beta_{cont,ads_j}^i = \frac{2k_B T}{3\mu} \left[\frac{C^i(d_m)}{d_g^i} + \frac{C^i(d_{PAH_j})}{d_{PAH_j}} \right] (d_g + d_{PAH_j}) \quad (2.103)$$

where C^i is the Cunningham function calculated as Equation 2.37.

2.6.1 Irreversible Dimerization

The irreversible dimerization is based on the irreversible collision of PAHs leading to their clustering/polymerization that forms dimers, trimers, and tetramers until the polymer mass reaches a threshold that can be considered a solid particle. For practical purposes, dimer is usually considered as a incipient particle that grows by surface growth and coagulation. A single-step irreversible collision of two similar PAHs forms a new dimer as:



Similarly, the adsorption of each PAH molecule on soot particles is described by the irreversible collision of soot and PAH_j as:



The forward rate of dimerization, k_{f,dim_j} and adsorption, k_{f,ads_j} in Reactions (R.2.7) and (R.2.8) are calculated from Equations (2.95) and (2.101), respectively.

$$k_{f,dim_j} = \gamma_{inc} \cdot \beta_{jk,PAH} \cdot Av, \quad (2.104)$$

$$k_{f,ads_j}^i = \gamma_{ads_j} \cdot \beta_{j,ads}^i \cdot Av, \quad (2.105)$$

where γ_{inc} and γ_{ads} are the collision efficiencies for dimerization and adsorption, respectively. Their value vary in $[10^{-7}, 1]$ range, and usually chosen to match the predicted soot mass with the experimental data. Here, $\gamma_{inc} = 10^{-4}$ and $\gamma_{ads} = 10^{-3}$. The rate of dimerization and adsorption from PAH_j are calculated accordingly as:

$$w_{dim_j} = k_{f,dim_j} [PAH_j][PAH_j] \quad (2.106)$$

The partial source terms for inception are calculated as:

$$I_{N,inc} = \frac{1}{\rho} \sum_{j=1}^{n_{PAH}} w_{dim_j} 2n_{PAH_j,C} \quad (2.107)$$

$$I_{C_{tot},inc} = \frac{1}{\rho} \sum_{j=1}^{n_{PAH}} w_{dim_j} 2n_{PAH_j,C} \quad (2.108)$$

$$I_{H_{tot},inc} = \frac{1}{\rho} \sum_{j=1}^{n_{PAH}} w_{dim_j} 2n_{PAH_j,H} \quad (2.109)$$

The rate of PAH adsorption for each section is obtained as:

$$w_{ads_j}^i = k_{f,ads_j}^i [\text{soot}][PAH_j] \quad (2.110)$$

The contribution of PAH adsorption to the source terms are expressed as:

$$I_{C_{tot},ads}^i = \frac{1}{\rho} \sum_{j=1}^{n_{PAH}} w_{ads_j}^i n_{PAH,C} \quad (2.111)$$

$$I_{H_{tot},ads}^i = \frac{1}{\rho} \sum_{j=1}^{n_{PAH}} w_{dim_j}^i (n_{PAH,H} - 2) \quad (2.112)$$

Each PAH molecule loses one H atom becoming a radical that forms bonds with a dehydrogenated site on soot surface, so two H atoms are released during the process that is taken into account in Equation (2.112).

The formation of a dimer consumes two PAH molecules, and during adsorption one PAH molecule is removed from the gas mixture, so the total rate of PAH_i removal by the irreversible dimerization is obtained as:

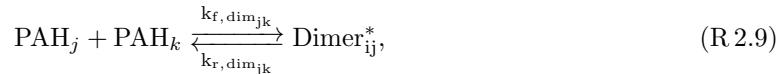
$$\left(\frac{d [PAH_j]}{dt} \right)_{inc} = -2w_{dim_j} - \sum_{i=1}^{n_{sec}} w_{ads_j}^i. \quad (2.113)$$

During the adsorption process one H₂ is released to the gas mixture.

$$\left(\frac{d [H_2]}{dt} \right)_{inc} = \sum_{i=1}^{n_{sec}} w_{ads_j}^i. \quad (2.114)$$

2.6.2 Reactive Dimerization

This model is built on Irreversible Dimerization with two main differences: The first step of dimerization and adsorption is reversible forming physically bonded dimers followed by a irreversible carbonization that leads to chemical bond formation in dimers [57]. This approach allows formation of homo- and heterodimers. The dimerization of PAH_j and PAH_k is described as:



where Dimer_{jk}^{*} and Dimer_{jk} physically and chemically bonded dimers, respectively, from PAH_j and PAH_k. The forward rate of physical dimerization, $k_{f,dim_{jk}}$ is calculated from Equation (2.95) as:

$$k_{f,dim_{jk}} = p'' \cdot \beta_{jk,PAH} \cdot Av, \quad (2.115)$$

where $p'' = 0.1$ accounts for the probability of PAH-PAH collisions in “FACE” configuration that results in successful vdW bond formation [148]. The reverse rate of physical dimerization, $k_{r,dim_{jk}}$ is obtained from the dimerization equilibrium constant [53] as:

$$\log_{10} K_{eq} = a \frac{\epsilon_{jk}}{RT} + b, \quad (2.116)$$

$$k_{r,dim_{jk}} = k_{f,dim_{jk}} 10^{-b} e^{-a\epsilon_{jk}\ln(10)/(RT)}, \quad (2.117)$$

$$\epsilon_{jk} = cW_{jk} - d, \quad (2.118)$$

$$W_{jk} = \frac{W_j \cdot W_k}{W_j + W_k}, \quad (2.119)$$

where $a = 0.115$ (obtained from pyrere dimerization data [48]) and $b=1.8$ [57], $c=933420$ j/kg, and $d=34053$ j/mol [57].

The rate of chemical bond formation, k_{reac} is defined in the Arrhenius form [49] as

$$k_{reac} = 5 \times 10^6 \cdot e^{(-96232/RT)}. \quad (2.120)$$

Assuming a steady state condition for the physical dimers, $\partial[\text{Dimer}_{jk}^*]/\partial t = 0$, the rate of formation of chemically-bonded dimers can be obtained as

$$\omega_{dim_{jk}} = k_{reac} \frac{k_{f,dim_{jk}} [\text{PAH}_j][\text{PAH}_k]}{k_{r,dim_{jk}} + k_{c,dim}}. \quad (2.121)$$

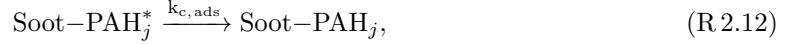
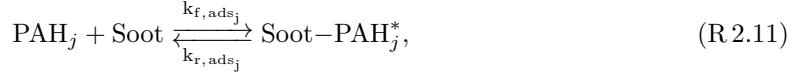
The contribution of dimer formation to partial source terms is expressed by looping over all combinations of PAHs as:

$$I_{N,inc} = \frac{1}{\rho} \sum_{j=1}^{n_{PAH}} \sum_{k=j}^{n_{PAH}} \omega_{dim_{kj}} (n_{PAH_j,C} + n_{PAH_k,C}), \quad (2.122)$$

$$I_{C_{tot},inc} = \frac{1}{\rho} \sum_{j=1}^{n_{PAH}} \sum_{k=j}^{n_{PAH}} \omega_{dim_{kj}} (n_{PAH_j,C} + n_{PAH_k,C}), \quad (2.123)$$

$$I_{H_{tot},inc} = \frac{1}{\rho} \sum_{j=1}^{n_{PAH}} \sum_{k=j}^{n_{PAH}} \omega_{dim_{kj}} (n_{PAH_j,H} + n_{PAH_k,H}), \quad (2.124)$$

Similarly, PAH adsorption is described by a two-step process where the collision of PAH_j with soot agglomerates leads to physically bonded Soot – PAH* that is carbonized and forms chemically-bonded Soot – PAH added to soot surface.



The forward rate of PAH-soot collision is calculated from Equation (2.101), and the reverse rate is determined same as inception (Equation(2.117)).

$$k_{f,ads}^i = \beta_{jk,ads}^i \cdot Av, \quad (2.125)$$

$$k_{r,ads}^i = k_{f,ads}^i \cdot 10^{-b} e^{-a\epsilon_{soot,j}\ln(10)/(RT)}, \quad (2.126)$$

$$\epsilon_{soot,j} = cW_{soot,j} - d, \quad (2.127)$$

a, b, c, d values are the same as those used in inception. Computing $\epsilon_{soot,j}$ also requires "equivalent soot molecular weight", W_{soot} for section i, which is estimated from carbon mass of each agglomerate as:

$$W_{soot}^i = \frac{C_{tot}^i W_{carbon}}{N_{agg}^i} \quad (2.128)$$

The rate constant of carbonization of Soot – PAH_j^{*} is defined as in the Arrhenius form similar to inception (Equation (2.120)). The prefactor is adjusted based on matching the numerical PSD [49] with measurements in the ethylene pyrolysis in a flow reactor [149].

$$k_{c,dim} = 2 \times 10^{10} \cdot e^{(-96232/RT)}. \quad (2.129)$$

The total adsorption rate can be calculated assuming a steady-state concentration for physically adsorbed PAH on soot, $\partial[\text{Soot} - \text{PAH}^*]/\partial t = 0$ calculated in a similar way to inception flux (Equation (2.121)) as

$$\omega_{ads_j}^i = k_{c,ads} \frac{k_{f,ads_j} [\text{Soot}][\text{PAH}_j]}{k_{r,ads_j} + k_{c,ads_j}}, \quad (2.130)$$

The contribution of PAH adsorption rate to partial source terms can be expressed as:

$$I_{C_{tot},ads}^i = \frac{1}{\rho} \sum_{i=1}^{n_{PAH}} \omega_{ads_j}^i n_{C,PAH_j}, \quad (2.131)$$

$$I_{C_{tot},ads}^i = \frac{1}{\rho} \sum_{i=1}^{n_{PAH}} \omega_{ads_j}^i (n_{H,PAH_j} - 2). \quad (2.132)$$

The rate of removal of PAH from gas mixture due to adsorption is given as

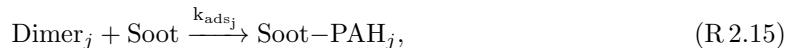
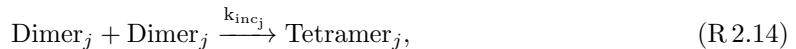
$$\left(\frac{d[\text{PAH}_j]}{dt} \right)_{inc} = - \sum_{k=1}^{n_{PAH}} w_{dim_{jk}} - \sum_{i=1}^{n_{sec}} w_{ads_j}^i. \quad (2.133)$$

During the adsorption process one H₂ is released to the gas mixture.

$$\left(\frac{d[\text{H}_2]}{dt} \right)_{inc} = \sum_{i=1}^{n_{sec}} w_{ads_j}^i. \quad (2.134)$$

2.6.3 Dimer Coalescence

Dimer coalescence model is a multi-step irreversible model proposed by Blanquart and Pitsch [85] where self-collision of PAH molecules form dimers that are intermediate state between gaseous PAH molecules and solid soot particles. The self-coalescence of dimers forms incipient soot particles. Alternatively, dimers can adsorb on the surface of existing soot particles and contribute to their surface growth. The equations are based on implementation of Dimer Coalescence by Sun et al. [150] for a sectional model.



where the rate constant of dimerization, k_{dim_j} and inception k_{inc_j} are calculated from collision rate of PAHs in Equation (2.101).

Table 2.2: The dimerization efficiency, γ_{dim_j} , for different PAH in dimer coalescence model [86]

Species name	Chemical formula	W [kg/mol]	γ_{dim_j}
Naphthalene	C ₁₀ H ₈	0.128	0.002
Acenaphthylene	C ₁₂ H ₈	0.152	0.004
Biphenyl	C ₁₂ H ₁₀	0.154	0.0085
Phenanthrene	C ₁₄ H ₁₀	0.178	0.015
Acephenanthrylene	C ₁₆ H ₁₀	0.202	0.025
Pyrene	C ₁₆ H ₁₀	0.202	0.025
Fluoranthene	C ₁₆ H ₁₀	0.202	0.025
Cyclopenta[cd]pyrene	C ₁₈ H ₁₀	0.226	0.039

$$k_{\text{dim}_j} = \gamma_{\text{dim}_j} \cdot \beta_{jj,PAH} \cdot Av, \quad (2.135)$$

$$k_{inc_j} = \beta_{jj,dimer} \cdot Av, \quad (2.136)$$

where γ_{dim_j} is the dimerization efficiency that is assumed to scale with fourth power of PAH molecular weight [86] as:

$$\gamma_{\text{dim}_j} = C_{N,j} \cdot W_{PAH_j}^4, \quad (2.137)$$

Blanquart and Pitsch [85] estimated the constant $C_{N,j}$ by comparing the profiles of several PAH species with experimental measurements in a single premixed benzene flame [151], and provided a efficiency values for various PAHs that are listed in Table 2.2. The rate of dimer collision is expressed as:

$$w_{\text{dim}_j} = k_{inc_j} [\text{Dimer}_j][\text{Dimer}_j] \quad (2.138)$$

Similarly, the rate of adsorption of dimers on soot particles is obtained as:

$$w_{ads_j}^i = k_{ads_j}^i [\text{soot}]^i [\text{Dimer}_j] \quad (2.139)$$

Assuming fast dimer consumption leads to the steady-state concentration of dimers that can be determined by solving a quadratic equation as:

$$a_{inc_j} [\text{dimer}]^2 + b_{ads_j} [\text{dimer}] = \omega_{\text{dim},j} \quad (2.140)$$

$$[\text{Dimer}_j] = \begin{cases} \frac{-b_{ads_j} + \sqrt{\Delta_j}}{2a_{inc_j}}, & \text{if } \Delta_j \geq 0 \\ 0 & \text{if } \Delta_j < 0 \end{cases} \quad (2.141)$$

$$\Delta_j = b_{ads_j}^2 - 4a_{inc_j}\omega_{\text{dim},j} \quad (2.142)$$

where $a_{inc_j} = k_{inc_j}$ and b_{ads_j} is calculated by summing the adsorption rate of dimer for all sections and dividing it by the dimer concentration.

$$b_{ads_j} = \sum_{i=1}^{n_{sec}} k_{ads_j}^i [\text{soot}]^i \quad (2.143)$$

After determining the concentration of each dimer, the contribution of inception and PAH adsorption to source terms of tracked soot variables can be calculated similar to previous inception models considering the number of carbon and hydrogen atoms involved in the process.

$$I_{N,inc} = \frac{1}{\rho} \sum_{j=1}^{n_{PAH}} 4\omega_{inc_j} n_{PAH_j,C}, \quad (2.144)$$

$$I_{C_{tot},inc} = \frac{1}{\rho} \sum_{j=1}^{n_{PAH}} 4\omega_{inc_j} n_{PAH_j,C}, \quad (2.145)$$

$$I_{H_{tot},inc} = \frac{1}{\rho} \sum_{j=1}^{n_{PAH}} 4\omega_{inc_j} (n_{PAH_j,H} - 2), \quad (2.146)$$

$$I_{C_{tot},ads}^i = \frac{1}{\rho} \sum_{i=1}^{n_{PAH}} 2\omega_{ads_j}^i n_{C,PAH_j}, \quad (2.147)$$

$$I_{C_{tot},ads}^i = \frac{1}{\rho} \sum_{i=1}^{n_{PAH}} 2\omega_{ads_j}^i (n_{H,PAH_j} - 2). \quad (2.148)$$

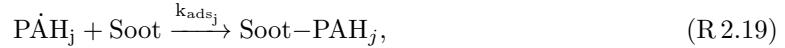
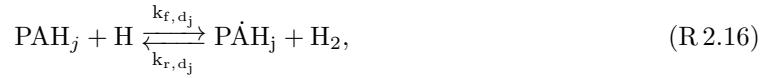
The rate of removal of PAHs and release of H₂ due to inception and PAH adsorption is calculated as:

$$\left(\frac{d [\text{PAH}_j]}{dt} \right)_{inc} = -4 \sum_{k=1}^{n_{PAH}} w_{inc_j} - 2 \sum_{i=1}^{n_{sec}} w_{ads_j}^i. \quad (2.149)$$

$$\left(\frac{d [\text{H}_2]}{dt} \right)_{inc} = 2 \sum_{i=1}^{n_{sec}} w_{ads_j}^i. \quad (2.150)$$

2.6.4 E-Bridge Formation

The E-Bridge Formation was originally proposed by Frenklach and Mebel [25] to describe soot inception using a HACA-like scheme that starts with dehydrogenation of PAH monomers, often pyrene, which forms the monomer radicals and continues with sequential addition of the radicals to PAHs that form dimers, trimers and larger polymers until the PAH structure reaches the mass threshold and the clustering process becomes irreversible. Here, a modified version of EBridge Formation model is used where dimers are considered as incipient soot, and monomer radical are adsorbed on soot agglomerates. This PAH growth model is described using the following set of pathways:



The rate constants of Reactions (R 2.16)&(R 2.17) are listed in Table 2.3 while those of dimer production and adsorption are calculated based on Equations (2.95)&(2.101), respectively. For both steps, it is assumed the all collisions are successful i.e. 100% collision efficiency for radical-monomer and radical-soot.

Table 2.3: Rate coefficients for the monomer de-/hydrogenation reaction of E-bridge formation in Arrhenius form $k = AT^n \cdot e^{-E/RT}$ [25]

Reaction	A	$\left[\frac{m^3}{mol \cdot s} \right]$	n	$\frac{E}{R}[K]$
(R 2.16)	f	$98 \times n_{C,PAH_j}$	1.8	7,563.519
	r	1.6×10^{-2}	2.63	2145.346
(R 2.17)	f	4.8658×10^7	0.13	0.0

$$k_{inc_j} = \beta_{jj,PAH} \cdot Av, \quad (2.151)$$

$$k_{ads_j}^i = \beta_{ads_j}^i \cdot Av, \quad (2.152)$$

The rate of dimer formation and adsorption is calculated as:

$$w_{dim_j} = k_{inc_j} [\text{PAH}_j] [\dot{\text{PAH}}_j] \quad (2.153)$$

$$w_{ads_j}^i = k_{ads_j}^i [\text{Soot}]^i [\dot{\text{PAH}}_j] \quad (2.154)$$

The calculations of rate of inception and PAH adsorption from PAH_j requires the concentration of corresponding monomer radical that can be determined by applying the steady-state assumption for $\dot{\text{PAH}}_j$.

$$\frac{d[\dot{\text{PAH}}_j]}{dt} = 0$$

$$\begin{aligned} k_{f,d_j} [\text{PAH}_j] [\text{H}] - k_{r,d_j} [\dot{\text{PAH}}_j] [\text{H}_2] - k_{f,h_j} [\dot{\text{PAH}}_j] [\text{H}] - k_{inc_j} [\dot{\text{PAH}}_j]^2 \\ - \sum_{i=1}^{n_{sec}} k_{ads_j}^i [\dot{\text{PAH}}_j] [\text{Soot}]^i = 0 \end{aligned}$$

The above equations can be rearranged as a quadratic equation similar to the dimer coalescence.

$$a_{inc_j} [\dot{\text{PAH}}_j]^2 + b_{ads_j} [\dot{\text{PAH}}_j] + c_j = 0, \quad (2.155)$$

$$a_{inc_j} = k_{f,d_j} \quad (2.156)$$

$$b_{ads_j} = k_{r,d_j} [\text{H}_2] + k_{f,h_j} [\text{H}] + \sum_{i=1}^{n_{sec}} k_{ads_j}^i [\text{Soot}]^i \quad (2.157)$$

$$c_{inc_j} = k_{f,d_j} [\text{PAH}_j] [\text{H}] \quad (2.158)$$

Finally, solving the quadratic equation for each PAH results in concentration of the radical using the following equation as:

$$[\dot{\text{PAH}}_j] = \begin{cases} \frac{-b_{ads_j} + \sqrt{\Delta_j}}{2a_{inc_j}}, & \text{if } \Delta_j \geq 0 \\ 0 & \text{if } \Delta_j < 0 \end{cases} \quad (2.159)$$

$$\Delta_j = b_{ads_j}^2 - 4a_{inc_j}c_j \quad (2.160)$$

The contribution of inception and adsorption to the partial source terms for E-Bridge formation can be written as:

$$I_{N,inc} = \frac{1}{\rho} \sum_{j=1}^{n_{PAH}} 2\omega_{inc_j} n_{PAH_j,C}, \quad (2.161)$$

$$I_{C_{tot},inc} = \frac{1}{\rho} \sum_{j=1}^{n_{PAH}} 2\omega_{inc_j} n_{PAH_j,C}, \quad (2.162)$$

$$I_{H_{tot},inc} = \frac{1}{\rho} \sum_{j=1}^{n_{PAH}} 2\omega_{inc_j} (n_{PAH_j,H} - 2), \quad (2.163)$$

$$I_{C_{tot},ads}^i = \frac{1}{\rho} \sum_{i=1}^{n_{PAH}} \omega_{ads_j}^i n_{C,PAH_j}, \quad (2.164)$$

$$I_{C_{tot},ads}^i = \frac{1}{\rho} \sum_{i=1}^{n_{PAH}} \omega_{ads_j}^i (n_{H,PAH_j} - 2). \quad (2.165)$$

The rate of removal of each PAH involved in soot inception and PAH adsorption and release of H₂ to the gas mixture can be expressed as:

$$\left(\frac{d[PAH_j]}{dt} \right)_{inc} = -2 \sum_{k=1}^{n_{PAH}} w_{inc_j} - \sum_{i=1}^{n_{sec}} w_{ads_j}^i. \quad (2.166)$$

$$\left(\frac{d[H_2]}{dt} \right)_{inc} = \sum_{i=1}^{n_{sec}} w_{ads_j}^i. \quad (2.167)$$

2.7 Gas scrubbing rates

The rate of production/destruction of species involved in soot formation must be taken into account to preserve the mass and energy balance in reactive systems. In order to do that, the production rate of gaseous species calculated by Cantera must be corrected for the rate of release/consumption due to PAH growth and surface reaction models.

$$\left(\frac{d[PAH_j]}{dt} \right)_{tot} = \left(\frac{d[PAH_j]}{dt} \right)_{gas} + \left(\frac{d[PAH_j]}{dt} \right)_{inc} + \left(\frac{d[PAH_j]}{dt} \right)_{ads}. \quad (2.168)$$

H₂ is released to the gas mixture due to inception, PAH adsorption as well as oxidation.

$$\left(\frac{d[H_2]}{dt} \right)_{tot} = \left(\frac{d[H_2]}{dt} \right)_{gas} + \left(\frac{d[H_2]}{dt} \right)_{inc} + \left(\frac{d[H_2]}{dt} \right)_{ads} + \left(\frac{d[H_2]}{dt} \right)_{ox}. \quad (2.169)$$

Surface growth consumes C₂H₂ and adds H₂ to the gas mixture.

$$\left(\frac{d[C_2H_2]}{dt} \right)_{tot} = \left(\frac{d[C_2H_2]}{dt} \right)_{gas} + \left(\frac{d[C_2H_2]}{dt} \right)_{gr}. \quad (2.170)$$

$$\left(\frac{d[H]}{dt} \right)_{tot} = \left(\frac{d[H]}{dt} \right)_{gas} + \left(\frac{d[H]}{dt} \right)_{gr}. \quad (2.171)$$

Oxidation uses O₂ and OH to remove carbon from soot particles and generates H₂ and CO.

$$\left(\frac{d[CO]}{dt} \right)_{tot} = \left(\frac{d[CO]}{dt} \right)_{gas} + \left(\frac{d[CO]}{dt} \right)_{ox}. \quad (2.172)$$

$$\left(\frac{d[O_2]}{dt} \right)_{tot} = \left(\frac{d[O_2]}{dt} \right)_{gas} + \left(\frac{d[O_2]}{dt} \right)_{ox}. \quad (2.173)$$

$$\left(\frac{d[OH]}{dt} \right)_{tot} = \left(\frac{d[OH]}{dt} \right)_{gas} + \left(\frac{d[OH]}{dt} \right)_{ox}. \quad (2.174)$$

Chapter 3

Results

3.1 Validation

A set of simulations was performed to ensure the accuracy and reliability of omnisoot for prediction of soot formation. Aerosol dynamics is validated by comparing the results of population balance models implemented in omnisoot with those of DEM simulations from literature. Carbon and hydrogen mass and energy balance is also rigorously evaluated to ensure that residuals fall within the bounds of acceptable numerical error.

3.1.1 Collision Frequency

The collision frequency function determines the rate at which two particles collide, which results in the reduction of total number of agglomerates and increase in size. In the absence of strong flow shear or external forces, Brownian motion is the main driving force for particle coagulation. As explained in Sections 2.4.4.1 & 2.4.5.1, omnisoot employs harmonic mean and Fuchs interpolations to calculate collision frequency of agglomerates from free-molecular to continuum regimes based on gas mean free path, and particle morphology.

The test case for validation of collision frequency is based on the DEM simulation of 2000 monodisperse spherical particles with the density of 2200 kg/m³ in a cubic cell with the constant temperature of 298 K and pressure of 1 atm [152]. Figure 3.1 depicts the collision frequency plotted against Knudsen number ($\text{Kn} = 2\lambda/d_m$) obtained by omnisoot using harmonic mean (red solid line) and Fuchs interpolation (green dashed line) and DEM results of Goudeli et al. [152]. The Fuchs interpretation perfectly matches DEM data over the free-molecular ($\text{Kn} < 10$) to the continuum ($\text{Kn} > 10$) range. Harmonic mean is also in good agreement with the DEM results in the free-molecular and continuum regime, but slightly underpredicts the collision frequency in the transition regime ($0.1 \leq \text{Kn} \leq 10$) with relative errors less than 16%.

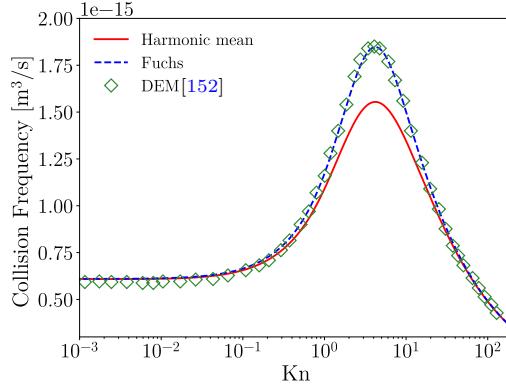


Figure 3.1: The comparison of collision frequency, β , obtained by omnisoott using harmonic mean (red solid line) and Fuchs interpolation (green dashed line) with DEM results (symbols) [152]

3.1.2 Coagulation

This test case was designed and conducted to validate the coagulation sub-unit of both particle dynamics models, MPBM and SPBM, by comparing the results of omnisoott with those of DEM [153]. The constant volume reactor was used for this test case, but it will be applicable to other reactors and flame models as long as the particle residence time matches with the values obtained by DEM. An adiabatic reactor with the volume of 1 m^3 is initialized with 2.6261×10^{18} spherical particles 2 nm in diameter. The initial conditions are indicated in Table 3.1. The particles are allowed to coagulate in the free molecular regime and grow in size while no inception, PAH adsorption and surface growth occur. Figure 3.3 demonstrates the number density of agglomerates (N_{agg}) and primary particles (N_{pri}), and mobility (d_m) and gyration (d_g) diameters of particle obtained by omnisoott that are in good agreement with DEM results. N_{pri} is conserved during coagulation resulting in identical flat lines for both particle dynamics models, but N_{agg} declines over time with the higher decay rate for SPBM because it accounts for the polydispersity of agglomerates that results in larger collision frequency compared to MPBM. Therefore, mean mobility and gyration diameter of SPBM (red lines in Fig. 3.3-b) are slightly larger than those of MPBM (blue lines of the same figure).

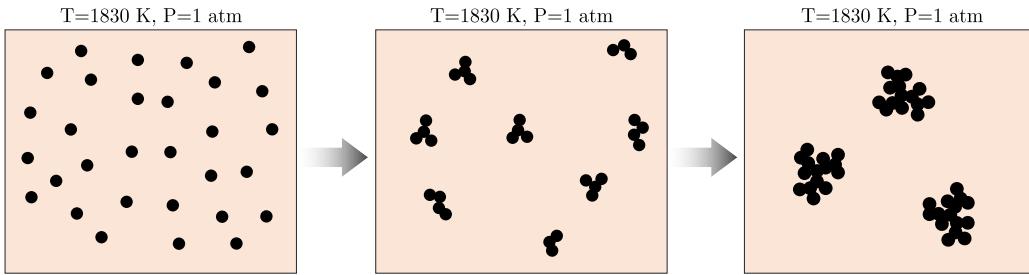


Figure 3.2: The schematic of agglomeration process in the coagulation test cases where initially spherical particle collide and form agglomerate

MPBM model cannot resolve PSD because of the monodispersity assumption. In contrast, SPBM tracks the number concentration of particles in separate sections that can be used to construct evolving PSD and calculate mean properties and determine the spread of size distribution of particles during coagulation. Figure 3.4-a shows the standard deviation of mobility diameter, σ_g predicted by SPBM in close agreement with DEM results. σ_g starts from unity indicating a monodisperse population at the beginning of simulation, and it finally reaches 2.03 that is the signature standard deviation of the free molecular regime [154]. Fig. 3.4 demonstrates the evolution of non-dimensional PSD from $t=1 \text{ ms}$ to 677 ms . The PSD is plotted for the normalized concentration, $\Psi = \bar{v}n_{agg}(v, t)/N_{agg,inf}$

Table 3.1: The simulations conditions of the coagulation test case [153]

Property	Value
Composition	CH ₄ :0.425, O ₂ :0.435, N ₂ :0.14
T	1830 K
P	1 atm
N_{agg}^1	3.514×10^{-5} mol/kg
N_{pri}^1	3.514×10^{-5} mol/kg
d_p^1	2 nm

and dimensionless volume, $\eta = v/\bar{v}$, where $n_{agg}(v, t)$ is the size distribution function of agglomerate, v particle volume, \bar{v} mean particle volume, $N_{agg,inf}$ total number concentration of agglomerates. For short residence times, $t \approx 4$ ms, the PSD resembles a half bell curve because the majority of particles has sizes close to $d_0 = 2\text{nm}$ with the average volume close to the minimum volume, so the particles with $\eta \approx 1$ has the largest concentration. As particles grow by coagulation, the PSD rapidly transitions to a full bell-curve ($t \geq 22\text{ms}$) and does not change for longer residence times, $t \geq 447\text{ms}$ marking the attainment of SPSD in a good agreement with DEM results. This confirms the capability of SPBM implemented in omnisoott to capture SPSD for soot agglomerates as a signature of Brownian-driven particle coagulation.

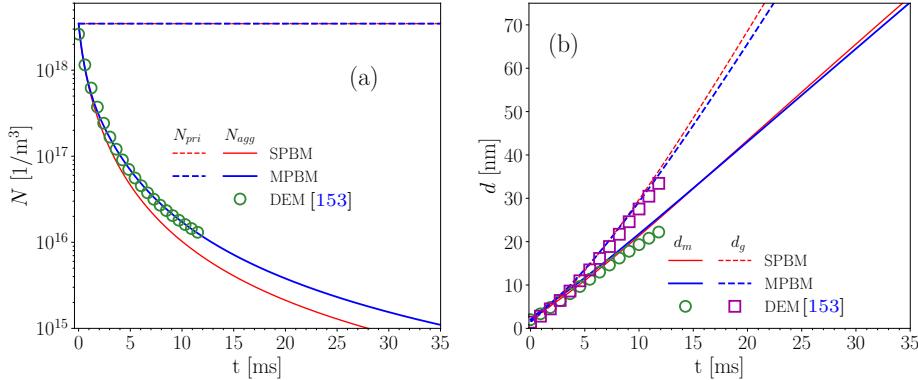


Figure 3.3: The total number concentration of agglomerates and primary particles (a), and mobility and gyration diameter (b) obtained with omnisoott using MPBM and SPBM that are in close agreement with the DEM results [153] indicating the validation of coagulation sub-model

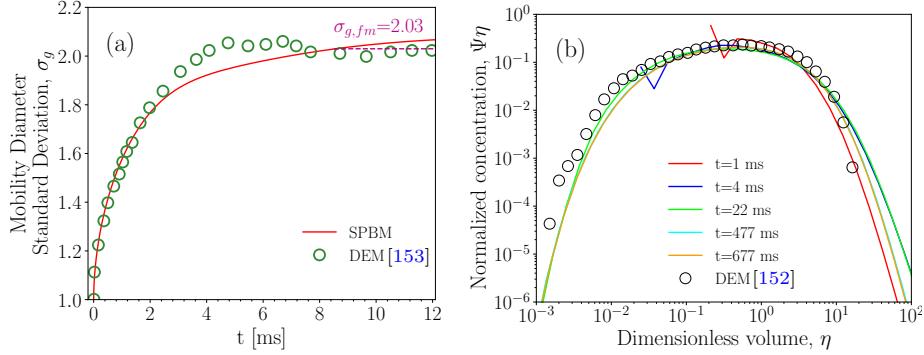


Figure 3.4: The standard deviation (residual) of mobility diameter, σ_g obtained with SPBM in close agreement with DEM results [153] (left pane) that reaches $\sigma_{g,fm} = 2.03$ characteristic of the free molecular regime [154]; the particle size distribution (normalized number concentration of agglomerates is plotted against non-dimensional volume in the right pane) at different residence times that overlaps after initial transient phase marking the attainment of self-preserving size distribution in good agreement with DEM results [152]

3.1.3 Constant Volume Reactor

The pyrolysis of 30% CH₄ diluted in N₂ with the initial temperature and pressure of 2455 K and 3.47 atm, respectively, was simulated using the constant volume reactor model in the residence time of 40 ms. The combination of available PAH growth and particle dynamics models leads to eight different cases that were simulated to ensure the conservation of mass and energy. Here, we focus on the total elemental balance of carbon and hydrogen because they are involved in soot processes. Figure 3.5 demonstrates the relative error of total carbon, hydrogen and energy of system for different PAH growth and particle dynamics models in the constant volume that falls below 10^{-10} for all parameters confirming the validity of model in satisfying the mass and energy balance in the constant volume reactor using all models.

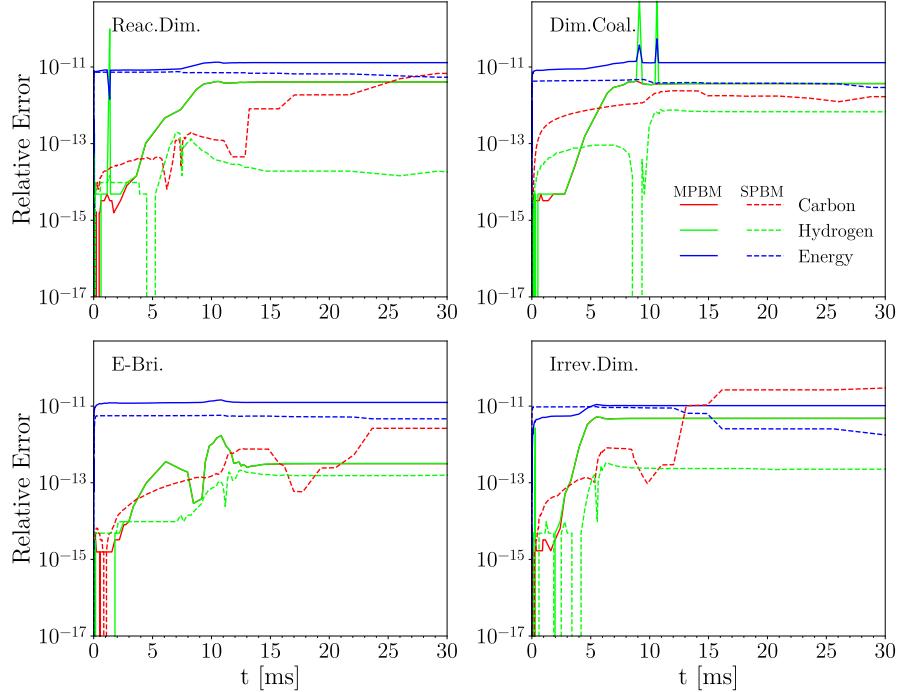


Figure 3.5: The relative error (residual) of total carbon (red line) and hydrogen (green line) mass, and total internal energy residual of gas and soot (blue line) plotted against residence time during pyrolysis of 30% $\text{CH}_4\text{-N}_2$ at 2455 K and 3.47 atm in the constant volume reactor simulated using different PAH growth models along with MPBM (solid line) and SPBM (dashed line)

3.1.4 Plug Flow Reactor

Methane pyrolysis in an adiabatic flow reactor is used to check elemental carbon and hydrogen, and energy balance in the PFR model. The inlet flow enters the reactor at the composition of 30% CH_4 diluted in N_2 , and $T=2100$ K and $P=1$ atm. Figure 3.6 shows the residual of total elemental carbon and hydrogen, and energy up to 40 cm of the reactor length using all PAH growth and particle dynamics model. The residuals are in the order of 10^{-11} and start to grow at the beginning of the reactor by pyrolysis of CH_4 and the formation of intermediate species such as C_2H_2 and PAHs. This initiates soot inception of surface growth affecting the gas chemistry and energy that ends near $x=10$ cm, and then the coagulation of particles is dominant with no affect on mass and energy of particles. As a result, PFR model of omnisoot satisfied the conservation of the mass and energy.

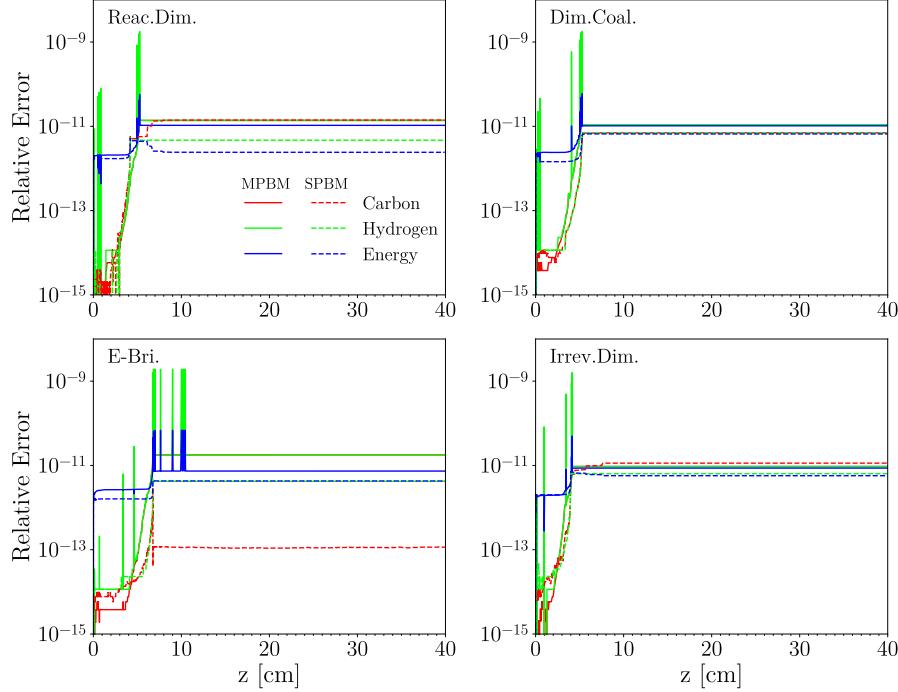


Figure 3.6: The relative error (residual) of total carbon (red line) and hydrogen (green line) mass, and total internal energy residual of gas and soot (blue line) plotted against reactor length (cm) in the adiabatic flow reactor during pyrolysis of 30% $\text{CH}_4\text{-N}_2$ at 2100 K and 1 atm simulated using different PAH growth models and MPBM (solid line) and SPBM (dashed line)

3.2 Soot yield and morphology during methane pyrolysis in shock-tubes

In this section, we focus on simulation of soot formation in shock-tubes using the constant volume reactor of omnisoott in order to (i) assess the performance of reaction mechanisms in prediction of the species mole fraction measured by laser diagnostics, and (ii) evaluate the accuracy of inception models to predict soot yield and morphology.

Temperature ranges achieved in shock-tube experiments ($T > 2000$ K) are relevant for plasma pyrolysis of methane (CH_4) as an emerging alternative production method for co-generation of CB and hydrogen [155]. Achieving specific grades of CB for different target applications requires a good control over its properties such as its morphology (quantified by primary particle, d_p and mobility diameter, d_m), composition (elemental carbon to hydrogen ratio), specific surface area (inversely proportional to d_p) and internal nanostructure (composed of aligned graphitic shell and disordered core). However, this is challenging because of the complexity of soot inception and mass growth processes and its coupling with gas phase chemistry, dependence on local temperature and pressure, and short time scales of soot inception and surface growth [36]. Therefore, robust particle dynamics models coupled with detailed chemical mechanisms are essential to better understand methane-to-CB conversion process under different feedstock compositions, and temperature and pressure time-histories.

While methane combustion has been extensively studied [156] due to the need for a detailed chemical description of natural gas combustion, the formation of carbonaceous particles such as soot and CB from methane pyrolysis are not well understood yet. Significant differences exist between reaction mechanisms in prediction of CH_4 and intermediate species such as C_2H_2 , and C_2H_4 in flame and reactors [156]. These uncertainties are amplified for larger hydrocarbons, especially polycyclic

aromatic hydrocarbons (PAH) due to multiplicity of pathways that are not identified in available reaction mechanisms.

Shock tubes can be effective for kinetic studies and measurements of methane pyrolysis and the subsequent soot inception and surface growth due to the uniform and instant heating of the mixture behind the reflected shock to the desired temperature and pressure as well as the absence of complicating factors such as diffusion and mixing. The shock-wave propagation and reflection creates nearly isothermal and isobaric conditions ideal for kinetic studies of pyrolysis and particle formation. However, these conditions are limited to short residence times usually 1-3 ms, so they can only be used to study early stages of soot formation such as inception and surface growth, and not for processes occurring at longer residence times such as coagulation, graphitization and carbonization.

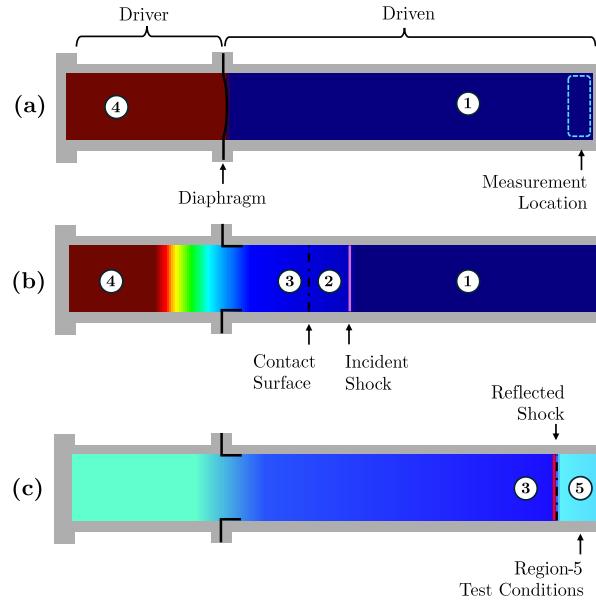


Figure 3.7: Different stages of shock wave propagation during shock-tube test from (a) before rupture of diaphragm leading to (b) propagation of incident shock towards the end of shock-tube and (c) its reflection from the end wall that increases pressure and temperature of test gas (Reprinted from Ref. [157])

Different stages of shock-wave propagation is illustrated in Fig.3.7. A shock tube consists of a driver and a driven section separated by a single-use diaphragm. The driver section contains usually an inert gas with a low molecular weight such as He, and the driven section is filled with the studied test gas. The driver section is pressurized to the desired pressure until the rupture of the diaphragm. This creates a shock wave that propagates through the driven section causing a pressure and temperature jump in the test gas.

After reaching the end wall, the shock wave reflects back towards the driver end of the shock-tube resulting in a secondary compression and heating of the mixture, and ultimately stagnating the test gas. In nearly 10 μs , this shock-heating process can bring the test gas from room-temperature to temperatures upwards of 10,000 K, and pressures more than 1000 atm.

The passage of reflected shock, referred to as region 5 with T_5 and P_5 , is the reference for calculation of residence time and the starting point for simulations. The delay in appearance of soot (often quantified by a threshold for LII or extinction signals) is known as induction time, τ_{ind} . There has been a lot of research in the literature on induction time (similarly on ignition delay time) in shock tubes [158], but it is not the focus of this work. Instead, we mainly focus on species concentrations and soot characteristics during the pyrolysis of methane, at atmospheric and higher pressures which can be used for the design and optimization of CB in plasma reactors [155].

3.2.1 Experimental setup and data collection

The experiments on CH₄ pyrolysis were conducted by Hanson Research Group at Stanford University. The data has not been published yet (at the time of writing the document) and were provided through the collaboration with Monolith Materials and Stanford University.

Mole fraction time histories of methane (CH₄), acetylene (C₂H₂), and ethylene (C₂H₄) were captured using continuous wave (CW) laser absorption at 3.365 μm, 2.998 μm, and 10.532 μm, respectively [159, 160, 161]. Soot volume fraction was measured using laser light extinction at $\lambda=633$ nm and 1064 nm with absorption function, E(m) of 0.174 and 0.203 [162], respectively. Additionally, soot samples were collected onto imaging stubs mounted in the shock-tube end-wall. A schematic of the experimental setup for laser diagnostics in the shock tube is shown in Fig. 3.8. Samples were extracted from the interior surface of the shock tube endwall after each experiment to allow for imaging and analysis of the particulates. TEM images were recorded with a FEI Tecnai G2 F20 X-TWIN microscope and Gatan SC200 camera for the test case of P₅=4.5 atm and T₅=2217 K.

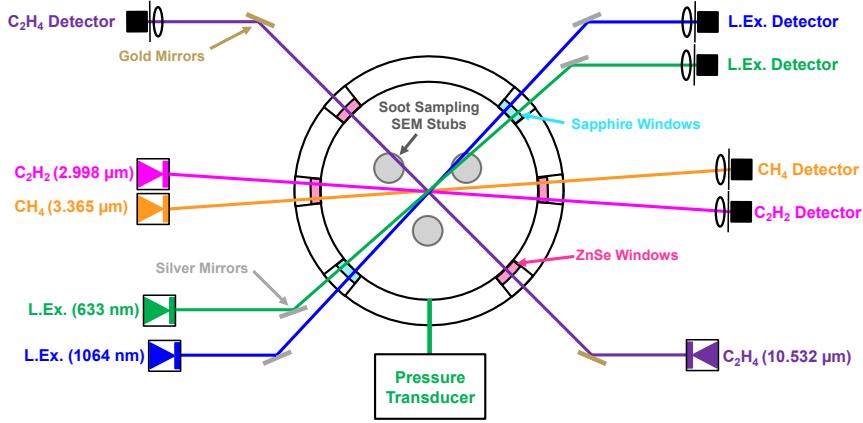


Figure 3.8: Layout of the laser diagnostics and shock-tube setup. Spatial and spectral filtering is employed to ensure high-quality absorbance measurements to infer species mole fractions and soot volume fraction. Soot samples are collected at the shock-tube endwall. All lasers are aligned in a plane 1 cm from the endwall

3.2.2 30% Methane Data-set

The first data set includes eight data points with the fuel loading of 30% CH₄-Ar at P=4±0.5 atm in the temperature range of 1800-2500 K. Table 3.2 lists the test conditions including pressure, temperature and composition of all data points.

Table 3.2: The pressure, temperature and composition of simulation data points for 30% CH₄

	Datapoints							
	(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)
T [K]	1861	1917	2030	2155	2184	2313	2375	2455
P [atm]	4.12	3.74	3.56	3.93	3.62	3.58	3.75	3.47
Composition	CH ₄ : 0.3, Ar: 0.7							

The time history of CH₄, C₂H₄ and C₂H₂ mole fraction as well as soot yield and volume fraction were reported up to 0.5 ms. The constant volume reactor of omnisoot was used with all PAH growth and particle dynamics models. The performance of reaction mechanisms are assessed by

comparing the mole fraction of species captured by laser diagnostics, CH₄, C₂H₄, and C₂H₂ with model predictions using Caltech [163], KAUST [156] and ABF [58] mechanisms. The analysis of simulation results (see Figs. B.2-B.2) showed that soot formation described by different PAH growth and particle dynamics models have negligible impact on prediction the mole fraction of measured species in the studied temperature range. Therefore, the assessment of reaction mechanisms for prediction of major species is performed without soot to simplify the analysis. Figs. 3.9-3.11 compares the mole fraction of CH₄, C₂H₄, C₂H₂ predicted by Caltech, KAUST, and ABF with data from laser diagnostics. Note that only a subset of simulations are shown for brevity, and results for the whole temperature range can be found in Figs. B.8-B.10. CH₄ mole fraction starts with a sharp drop from 0.3 at the beginning (corresponding to fast decomposition of methane) followed by a more gradual decrease. The initial decomposition rate increases with temperature. All mechanisms underpredict CH₄ mole fraction, but ABF mechanism predictions are closer to the measurements than KAUST and Caltech. In other words, KAUST and Caltech overestimate CH₄ conversion rate. The mole fraction of C₂H₄ (Fig.3.10) starts with a quick jump after methane decomposition, reaches the maximum value and gradually decreases towards its final value. The simulations extended beyond test time of shock-tube calculations showed that C₂H₄ mole fraction predicted by KAUST and Caltech mechanisms continues to decrease to the equilibrium values, but ABF predicts an increase at longer residence times indicating reversal of C₂H₄ decomposition which is more pronounced at high temperatures. The C₂H₂ mole fraction increases in two steps, a rapid increase at the beginning followed by a gradual one. ABF mechanism predictions are in good agreement with measurements, but Caltech and KAUST underpredict the C₂H₂ mole fraction. The difference between mechanisms are more pronounced at equilibrium conditions. While ABF predicts a plateau beyond 1.5 ms, KAUST and Caltech predict a decrease towards a lower value indicating conversion of acetylene to larger hydrocarbons and PAHs. The carbon mass fraction (CMF) of measured species, CH₄, C₂H₄ and C₂H₂ combined is shown in Fig. 3.12. The CMF of different mechanisms decreases up to a certain time that shortens with temperature. Beyond that, ABF mechanism predicts an increasing trend indicating that carbon returns to major species, but KAUST and Caltech continuously transfer carbon from measured species to larger hydrocarbons. This is consistent with the CMF of soot precursors shown in Fig. 3.13, which is significantly larger for KAUST and Caltech compared with ABF by two to three orders of magnitude, so it is expected that ABF mechanism cannot provide enough precursors for soot model to predict volume fraction comparable with the measurements.

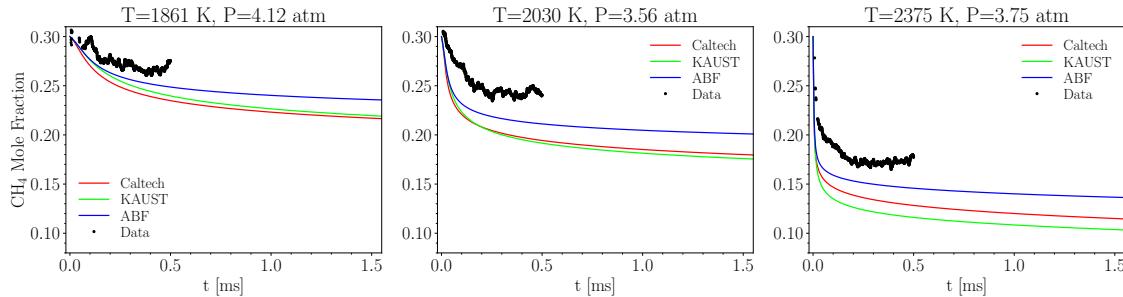


Figure 3.9: The time history of carbon mass fraction of CH₄ of 30% CH₄ pyrolysis at T=1861, 2083, and 2375 K using Caltech, KAUST, and ABF mechanisms compared with laser diagnostics data

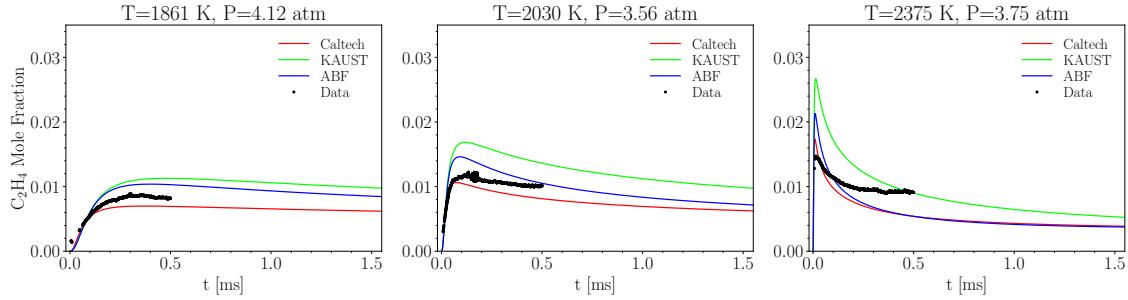


Figure 3.10: The time history of carbon mass fraction of C₂H₄ of 30% CH₄ pyrolysis at T=1861, 2083, and 2375 K using Caltech, KAUST, and ABF mechanisms compared with laser diagnostics data

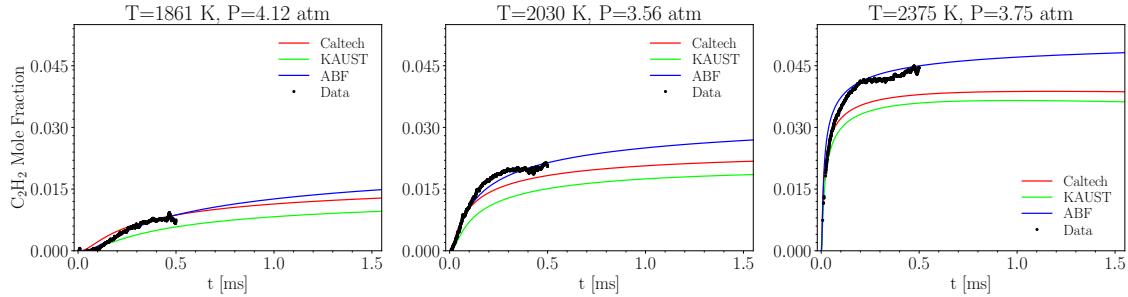


Figure 3.11: The time history of carbon mass fraction of C₂H₂ of 30% CH₄ pyrolysis at T=1861, 2083, and 2375 K using Caltech, KAUST, and ABF mechanisms compared with laser diagnostics data

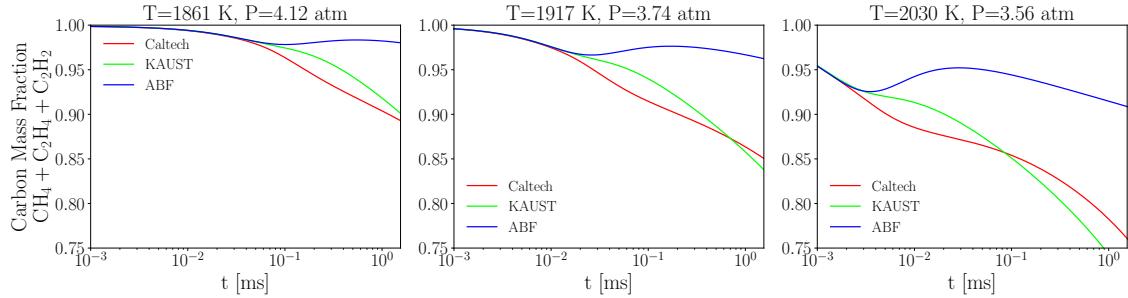


Figure 3.12: The time history of carbon mass fraction of CH₄, C₂H₄, C₂H₂ of 30% CH₄ pyrolysis at T=1861, 2083, and 2375 K using Caltech, KAUST, and ABF

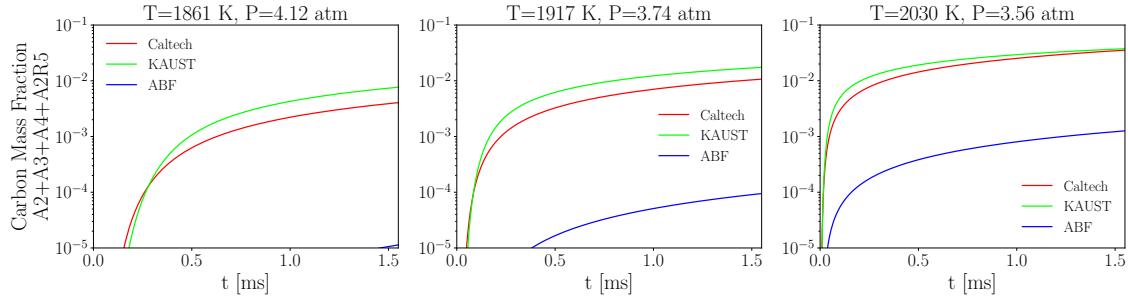


Figure 3.13: The time history of carbon mass fraction of soot precursors, A2, A3, A4 and A2R5 of 30% CH_4 pyrolysis at $T=1861$, 2083, and 2375 K using Caltech, KAUST, and ABF mechanisms

Comparing f_v predicted by three reaction mechanisms at $T=2375$ K and $P=3.75$ atm with the data from extinction measurements confirms the above-mentioned conclusion. As shown in Fig. 3.14, the f_v by ABF follows a similar trend as data, but the values are nearly two orders of magnitude lower indicating small PAH concentrations. Although both Caltech and KAUST predict f_v comparable with data, KAUST will be used for detailed soot simulations as its combination with E-Bridge Modified accurately captures soot volume fraction. Having said that, any conclusions drawn using KAUST can be easily generalized to Caltech as well.

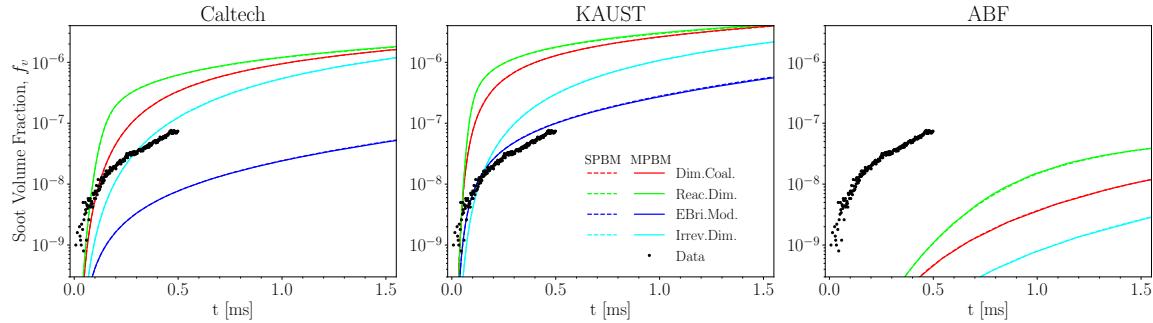


Figure 3.14: The time history of soot volume fraction of 30% CH_4 pyrolysis at $T=2375$ K and $P=3.75$ atm using Caltech, KAUST, and ABF mechanism and different PAH growth and particle dynamics models

Fig. 3.15 shows that the time variation of soot volume fraction in $T=2155$, 2313, and 2455 K. The measurement data was not available (or too noisy) below $T=2155$ K due to lack of enough soot particles leading to weak extinction signals (see Fig. B.13). SPBM and MPBM results coincide indicating that taking polydispersity into account by the particle dynamics model does not affect soot mass during 1.5 ms of simulation. f_v by EBridge Modified is in good agreement with data and the lowest among inception models across the temperature range. On the other hand, Reactive Dimerization and Dimer Coalescence yields the largest f_v nearly two orders of magnitude larger than data (and EBridge Modified). There is a remarkable similarity between trend of f_v and CMF of soot precursors (Fig. 3.13) highlights that major role of inception and PAH adsorption in soot mass growth. Temperature promotes soot formation in the studied shock tube. Fig. 3.16 depicts that f_v computed at 1.5 ms increases with temperature in a nearly linear fashion expect for EBridge Modified.

Fig. 3.17 shows the time history of primary particle diameter, d_p at different temperatures. As expected, EBridge Formation and Irreversible Dimerization has the lowest d_p due to low soot mass growth rate leading to small f_v values. Dimer Coalescence and Reactive Dimerization generate a relatively close f_v values (corresponding to similar soot mass growth), but Reactive Dimerization predicts a larger d_p indicating lower inception flux and stronger PAH adsorption rate of this model

compared to Dimer Coalescence. There is a noticeable difference between sectional and monodisperse models in the d_p predicted by Reactive Dimerization, but other inception models does not exhibit such a sensitivity to particle dynamics model. Both models predict the initial rapid rise in d_p . While MPBM predicts a gradual increase to final value, d_p by SPBM decreases after the initial rise due to stronger inception rate by SPBM that generates particle with $d_p=2$ nm bringing down the average d_p . Fig. 3.18 depicts temperature dependence of d_p at 1.5 ms where SPBM predicts a decreasing trend as opposed to MPBM.

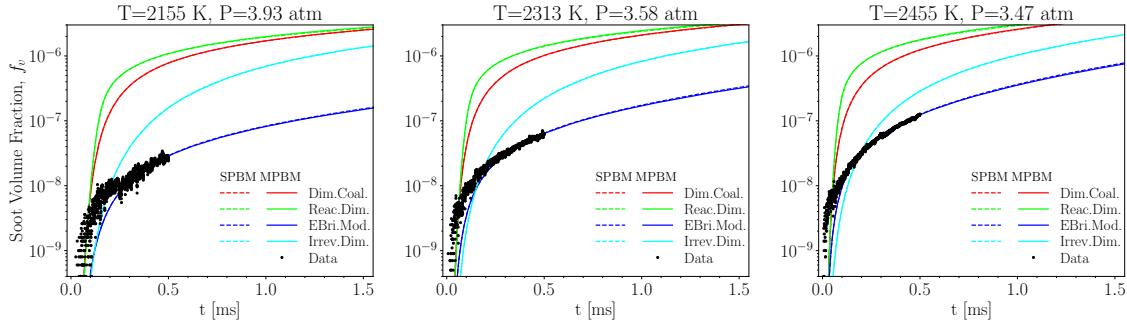


Figure 3.15: The time history soot volume fraction, f_v of 30% CH_4 pyrolysis for $T=2155$, 2313, and 2455 K using KAUST with different combinations of PAH growth and particle dynamics models compared with extinction measurements

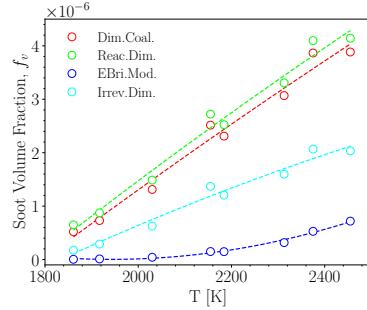


Figure 3.16: The soot volume fraction, f_v at 1.5 ms increasing with temperature during 30% CH_4 pyrolysis using KAUST and SPBM. The dashed lines are added to guide the eye

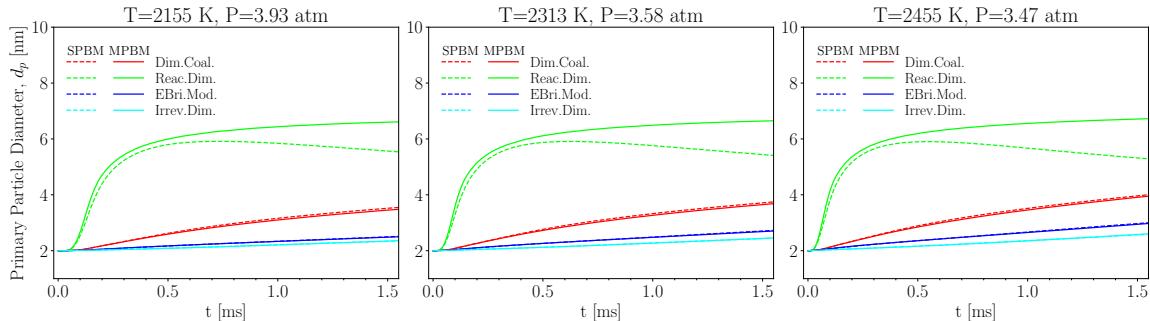


Figure 3.17: The time history of primary particle diameter, d_p of 30% CH_4 pyrolysis for $T=2155$, 2313, and 2455 K using KAUST with different combinations of PAH growth and particle dynamics models

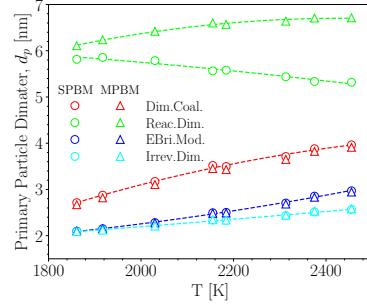


Figure 3.18: The primary particle diameter, d_p at 1.5 ms increasing with temperature during 30% CH_4 pyrolysis using KAUST with MPBM and SPBM. The dashed lines are added to guide the eye

As shown in Fig. 3.19, the inception flux of Reactive Dimerization starts with a quick rise followed by a quick drop until $t \approx 0.25$ ms which is nearly the same for both particle dynamics models, but beyond that SPBM predicts an increasing inception flux. This stems from a lower PAH adsorption rate by SPBM (with Reactive Dimerization) directing more PAHs to inception. All models predict an initial rapid rise due to PAH production by CH_4 decomposition. Dimer Coalescence reaches the highest peak and continuously decreases afterward because of consumption of PAHs. However, Irreversible Dimerization and EBridge Modified nearly stay constant towards the end of simulation time.

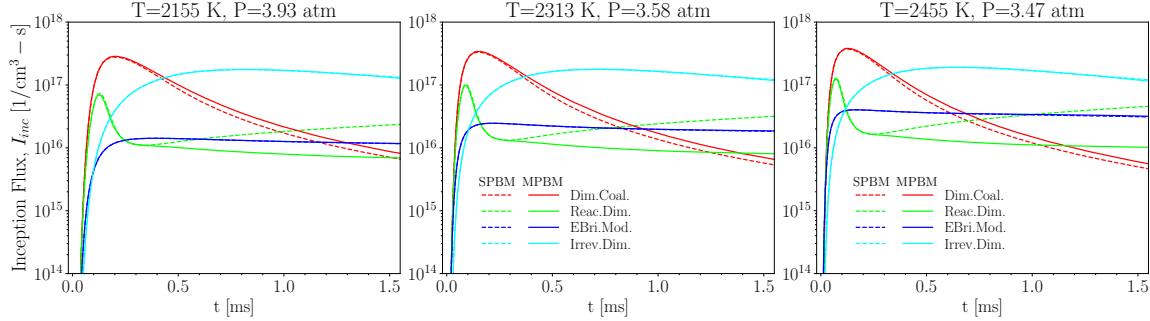


Figure 3.19: The time history of inception flux, I_{inc} of 30% CH_4 pyrolysis for $T=2155$, 2313, and 2455 K using KAUST with different combinations of PAH growth and particle dynamics models

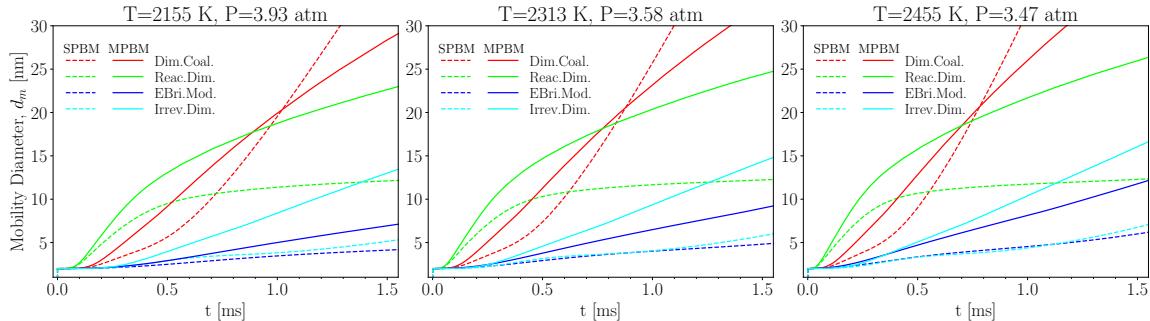


Figure 3.20: The time history of mobility diameter, d_m of 30% CH_4 pyrolysis for $T=2155$, 2313, and 2455 K using KAUST with different combinations of PAH growth and particle dynamics models

As shown Fig. 3.20, d_m exhibits a stronger sensitivity to the employed particle dynamics model compared to d_p because MPBM cannot capture the initial transition of particle population before the

Table 3.3: The pressure, temperature and composition of simulation data points for 10% CH₄

	Datapoints									
	(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)
T [K]	1873	1927	2074	2030	2348	1776	1897	2085	2277	2406
P [atm]	1.09	1.37	1.27	1.22	1.13	4.76	4.49	4.85	4.43	3.89
Composition	CH ₄ : 0.1, CO ₂ : 0.01, Ar: 0.7									

attainment of self-preserving size distribution. As a result, morphology (quantified by characteristic diameters) predicted by MPBM diverges from the SPBM results that are closer to actual polydisperse morphology of soot. Overall, Dimer Coalescence predicts a larger mean mobility diameter because of stronger inception flux at the beginning of simulation time (see Fig. 3.19) leading to more primary particles and higher coagulation rate that generates larger agglomerates. Fig. 3.21 shows the carbon addition rate by PAH adsorption is comparable and greater than that of HACA. This difference is more pronounced for Reactive Dimerization where PAH adsorption is larger by more than one order of magnitude.

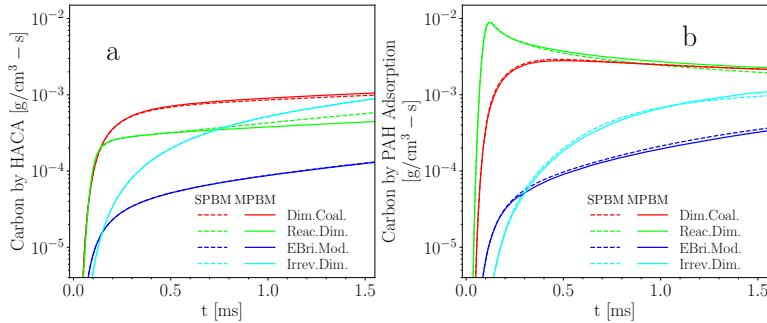


Figure 3.21: The time history of rate of carbon mass addition by (a) HACA and (b) PAH adsorption at T=2313 K using KAUSt with different combinations of PAH growth and particle dynamics models

3.2.3 10% Methane Data-set

The second data set from Stanford's group includes the measurements on 10% CH₄ pyrolysis at P=1±0.5 atm and 4±0.5 atm in the temperature range of 1800-2500 K. The first step is the examination of chemistry by comparing the predicted species mole fraction using different reaction mechanisms with data from laser diagnostics. Then, we look into soot volume fraction and morphology for a single data point at which TEM measurements were conducted, analyze the collected TEM images, identify main factors contributing to the discrepancies in volume fraction and morphology, optimize the PAH growth models to minimize the prediction error, and apply optimized model to the rest of data points. We limit our discussion to the 4-atm cases as extinction measurements showed negligible soot in 1-atm cases. The results of whole pressure and temperature ranges can be found in Section B.1.2.

3.2.3.1 Reaction Mechanism Assessment

As shown in Fig. 3.22, similar to 30% dataset, KAUSt yields the highest CH₄ conversion followed by Caltech and ABF mechanism. While KAUSt accurately predicts CH₄ mole fraction over the studied temperature range, Caltech and ABF overpredict it indicating that less carbon is directed toward intermediates and larger hydrocarbons.

Fig. 3.23 shows that ABF predicts the initial jump and gradual decrease towards equilibrium C_2H_4 mole fraction at $T=1897$ K in good agreement with the data, but underpredicts it at higher temperatures. KAUST mechanism overpredicts C_2H_4 production at the beginning leading to higher peak as well as its later consumption resulting in lower mole fraction towards the end of simulation compared to the measurements. Fig. 3.24 depicts underprediction of C_2H_2 mole fraction by all mechanism over the whole temperature range, which is more pronounced for KAUST mechanism. Fig. 3.25 shows that carbon mass fraction in measured species decreases for all mechanisms showing conversion of carbon from major species to larger hydrocarbons, but similar to 30% simulation, ABF predicts an increase in total carbon mass fraction. As shown in Fig. 3.26, the carbon mass fraction of designated soot precursors predicted by ABF are lower by more than one order of magnitude compared to KAUST and Caltech.

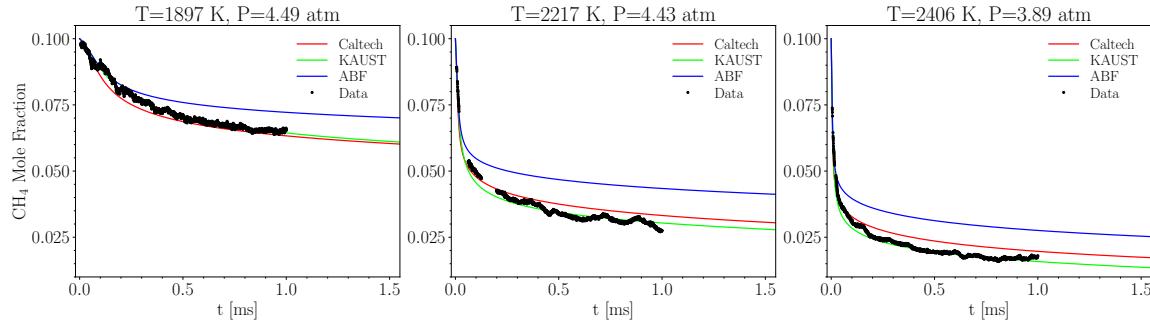


Figure 3.22: The time history of mole fraction of CH_4 of 10% CH_4 pyrolysis at $T=1897$, 2217, and 2406 K using Caltech, KAUST, and ABF mechanisms compared with laser diagnostics data

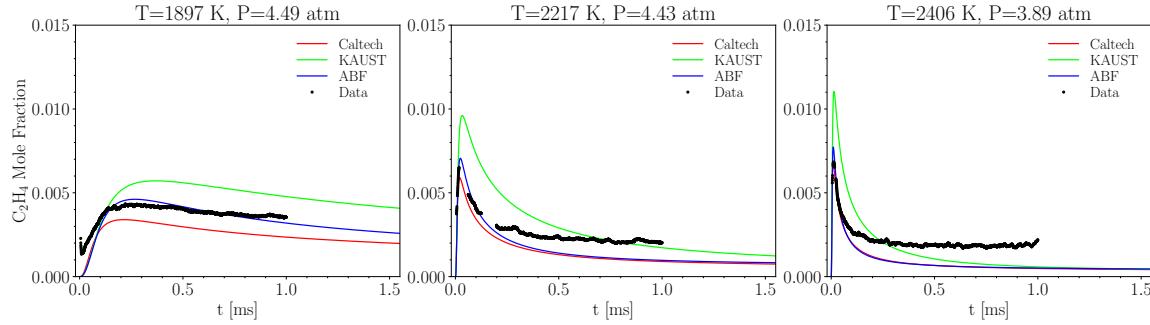


Figure 3.23: The time history of mole fraction of C_2H_4 of 10% CH_4 pyrolysis at $T=1897$, 2217, and 2406 K using Caltech, KAUST, and ABF mechanisms compared with laser diagnostics data

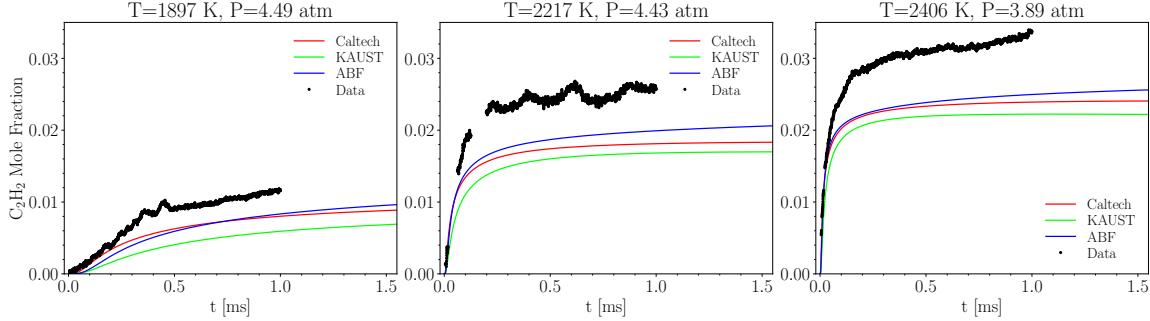


Figure 3.24: The time history of mole fraction of C₂H₂ of 10% CH₄ pyrolysis at T=1897, 2217, and 2406 K using Caltech, KAUST, and ABF mechanisms compared with laser diagnostics data

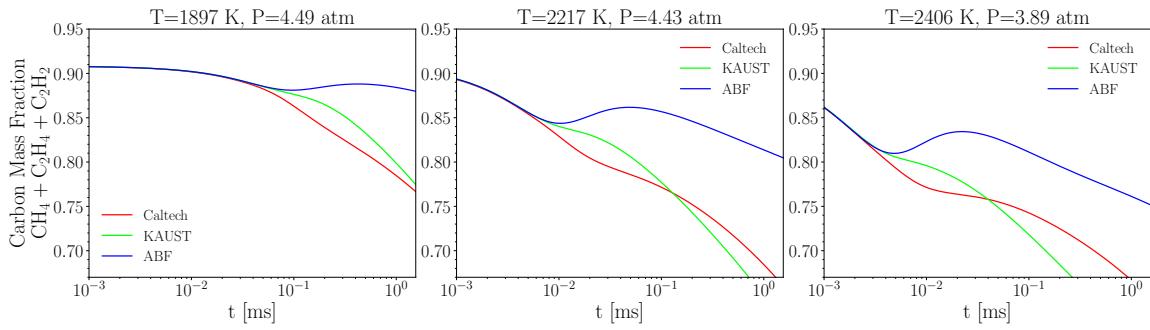


Figure 3.25: The time history of carbon mass fraction of CH₄, C₂H₄, C₂H₂ of 10% CH₄ pyrolysis at T=1897, 2217, and 2406 K using Caltech, KAUST, and ABF mechanisms compared with laser diagnostics data

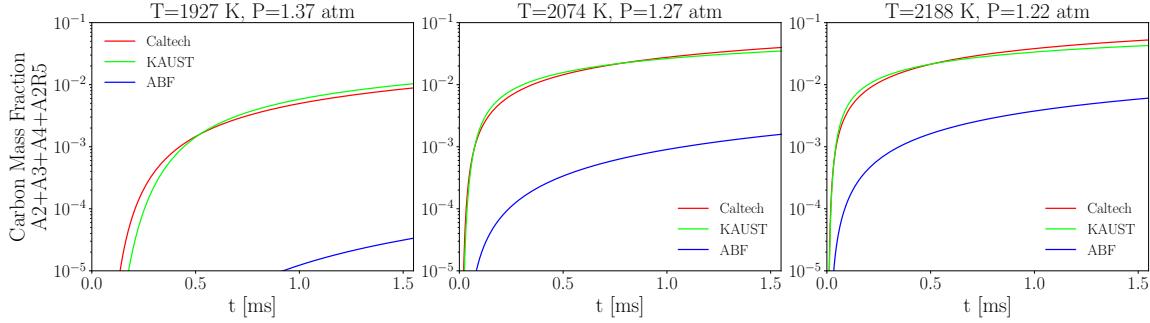


Figure 3.26: The time history of carbon mass fraction of A₂, A₃, A₄, A₂R₅ of 10% CH₄ pyrolysis at T=1897, 2217, and 2406 K using Caltech, KAUST, and ABF mechanisms compared with laser diagnostics data

3.2.3.2 Characterization of soot morphology in methane pyrolysis shock-tube

The TEM images of 10% CH₄ was provided by Stanford group for a single data point of T=2217 K and P=4.5 atm that can be used to characterize soot morphology from d_p and d_m. The images in addition to soot volume fraction enables assessing the performance of omnisoott in the description of soot generated during methane pyrolysis in shock-tube. As reported by Stanford team, the expansion wave traverses the shock tube around 2 ms reducing the temperature to nearly 1000 K freezing the chemical reactions that contribute to the surface growth, but the coagulation might continue until the

collection of particles leading to larger agglomerates. As a result, d_p estimated from the TEM images closely represents d_p of agglomerates at the end of process, but d_m estimated from TEM images could be larger the actual values due to post-expansion-wave growth of particles. *atems* package [164] was used for characterizing soot aggregates recorded in TEM images with different methods for evaluating the aggregate projected area, perimeter, and primary particle diameter. It creates a binary map from the raw TEM image where white and black pixels represent the agglomerates and background, respectively, and feeds the map to a agglomerate segmentation algorithm to detect individual agglomerates. Fig.3.27 shows a sample from TEM images and segmented agglomerates from generated map.

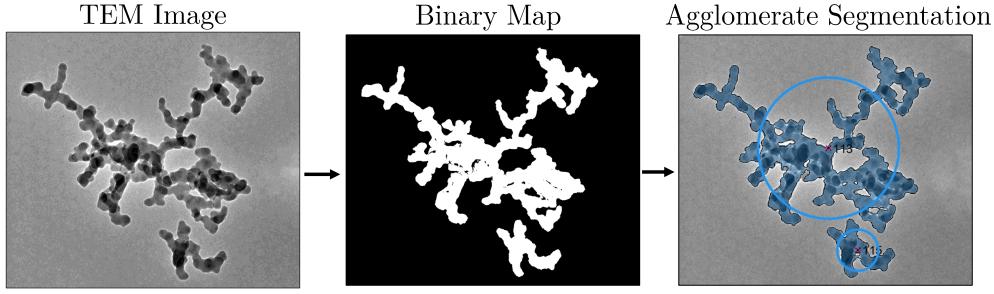


Table 3.4: The morphological characteristics of agglomerates quantified by atems using KMC and PCM

Property	Arithmetic mean	Median
d_m [nm]	97	69
d_p [nm]	22	18

from various sources. Soot morphology is mainly governed by coagulation leading to self-similar structures in which d_m scales with d_p and n_p . Olfert and Rogak [166] analyzed soot particles from flares [167], inverted burners [168], compression ignition engines [169] and other sources to support the external mixing hypothesis and related agglomerate size characterized by d_m to d_p using the following power-law:

$$d_p = d_{p,100} \left(\frac{d_m}{100\text{ nm}} \right)^{D_{\text{TEM}}}, \quad (3.2)$$

where $d_{p,100}$ is the average primary particle diameter for a 100 nm aggregate, and D_{TEM} is the exponent. Both quantities were obtained by fitting Eq. (3.2) to soot sampled from different sources. Fig 3.29 demonstrates the scatter plot of d_p against d_m compared with power-law curves of Eq. (3.2) using two sets of prefactor and exponent values, i) $D_{\text{TEM}}=0.32$, $d_{p,100}=20.6$ nm, and ii) $D_{\text{TEM}}=0.39$, $d_{p,100}=20.2$ nm taken from average values reported in summary of fit parameters in Table 1 of Olfert and Rogak [166]. The power-law show a good agreement with quantities that d_p and d_m computed by atems can be good representative of soot agglomerates with minimal agglomerate overlap.

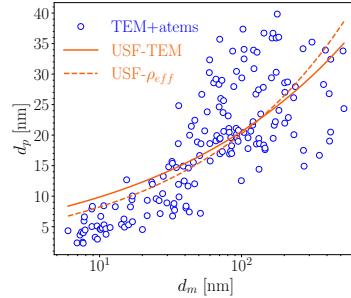


Figure 3.29: The d_m as a function of d_p from TEM image obtained by atems [164] (symbols) compared with the power law in Eq. (3.2)

The number of primary particles, N_{pri} was estimated from f_v and d_p assuming spherical primary particles with fixed size during the process. Using minimum, maximum, and mean of d_p from TEM measurements provides a range of expected N_{pri} can be obtained.

$$N_{pri} = \frac{f_v}{\pi d_p^3 / 6}. \quad (3.3)$$

3.2.3.3 Modelling soot yield and morphology

Fig. 3.30 compares predicted f_v and d_p with data. The horizontal error bars denotes the uncertainty in the residence time corresponding to the TEM measurement. EBridge Modified and Irreversible Dimerization predicted soot volume fraction in close agreement with measurements, but Reactive Dimerization and Dimer Coalescence overpredict f_v by a factor of 2-3. While the model has an overall accurate description of soot yield, the primary particle diameter, d_p is significantly underestimated by all PAH growth models. Reactive Dimerization yields the largest d_p nearly 6 nm due to stronger PAH adsorption rate.

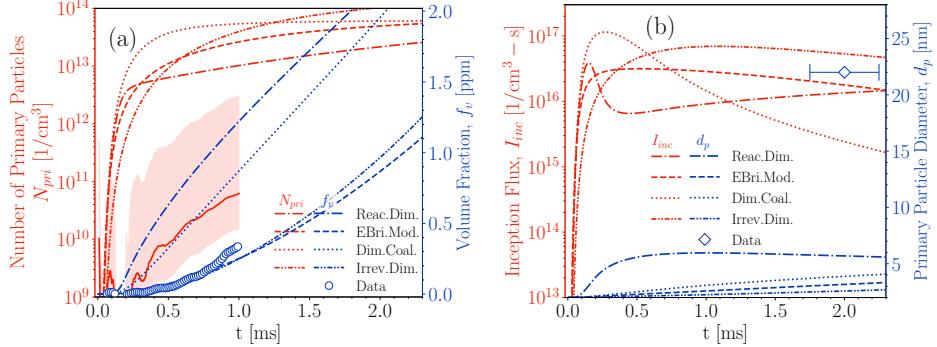


Figure 3.30: The soot volume fraction, f_v and number of primary particles, N_{pri} (left pane), and , the soot inception flux, I_{inc} and the primary particle diameter, d_p predicted by KAUST mechanism, SPBM, and different PAH growth models at $T=2230$ K, $P=4.5$ atm that corresponds to shock-tube conditions of TEM measurement

The underprediction of d_p by the PAH growth models despite producing enough or even more soot carbon mass compared with the measurement motivated a sensitivity analysis for soot yield and morphology. It should be noted that the sensitivity analysis is only focused on soot model to identify the sub-model(s) with the dominant effect on yield and morphology, and it does not include reactions in the gas phase.

First, the effect of HACA rate on soot yield and morphology is examined for the data point with $T=2230$ K and $P=4.5$ atm. To this purpose, the surface reactivity, α , in HACA formulation is modified by introducing a damping factor, ζ in Eq. (2.79) as:

$$\alpha^i = \tanh \left(\frac{12.56 - 0.00563 \cdot \zeta T}{\log_{10} \left(\frac{\rho_{soot} \cdot Av}{W_{carbon}} \frac{\pi}{6} d_p^3 \right)} - 1.38 + 0.00068 \cdot \zeta T \right). \quad (3.4)$$

Note that, there are many adjustable parameters in HACA scheme such as rate constants. We picked ζ to modify α for a number of reasons: First, α was initially introduced as a tuning parameter as a function of local temperature and primary particle diameter to control surface growth rate, and it has been usually adjusted specifically for each flame [170, 128] to match the predicted volume fraction with the measurements. Second, the global empirical relation of Appel et al. [58] (used in omnisoott to quantify α) developed by fitting parameters of α to minimize the prediction error of volume fraction for various premixed flames, but shock tubes generally have larger temperature ranges compared with premixed flames, which could excessively reduce surface reactivity and HACA growth rates. Finally, larger volume fractions (in the order of few ppm) were underpredicted by the global α relation (Fig. 9 of [58]). So, the low values of α in the shock tube at $T=2230$ K corresponding to process conditions of TEM measurements can contribute to underprediction of surface growth rates and d_p . We introduce ζ to reduce the damping effect of temperature on surface reactivity that promote HACA growth rates resulting in larger primary particle diameters, and possibly lessen the discrepancy between predicted and measured d_p . Fig. 3.31 demonstrates the variation of α with temperature for $\zeta=0.6, 0.8, 1$ and primary particle diameter 2 and 6 nm. α decreases with temperature, and it is inversely proportional to ζ .

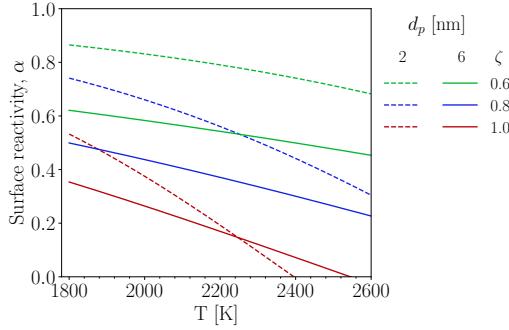


Figure 3.31: The variation of surface reactivity, α , as a function of temperature for different values of ζ and primary particle size of 2 nm (dashed lines) and 6 nm (solid lines)

A series of simulations was performed to study the damping effect of temperature on soot mass and morphology by varying ζ from 1 to 0.8 and 0.6. Note that, ζ of 1 represent the original α formulation. The reaction mechanism and particle dynamic model were set to KAUST and SPBM, respectively, but all PAH growth were used in the simulations. Fig. 3.32 depicts the volume fraction variation for each PAH growth model using different ζ values. As expected, lowering ζ increases the soot volume fraction with a maximum of 29% at $t=2$ ms for Dimer Coalescence. However, as shown in Fig. 3.33 ζ has negligible effect on d_p with all PAH growth models. As a result, HACA growth rate cannot account for the observed gap between measure and predicted primary particle within 2 ms of the simulation of shock tube.

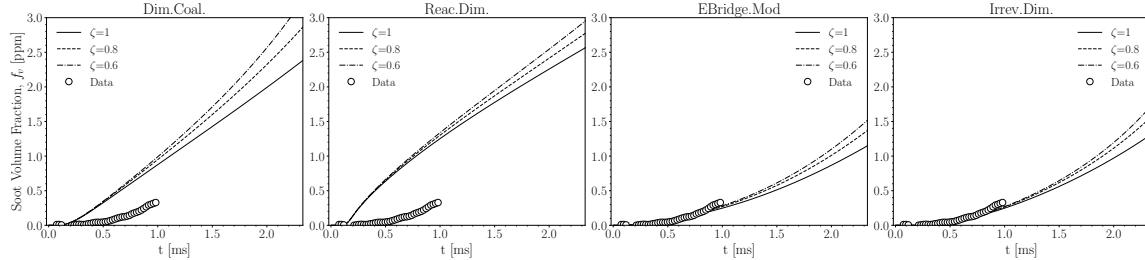


Figure 3.32: The effect of reducing ζ leading to larger HACA rates on soot volume fraction, f_v using KAUST, SPBM and different PAH growth model

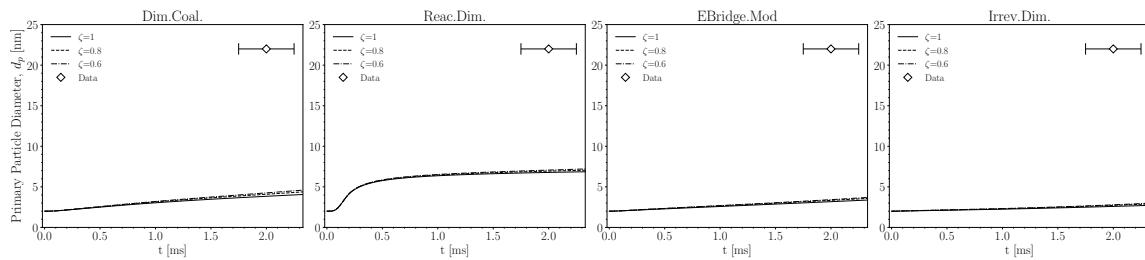


Figure 3.33: The effect of reducing ζ leading to larger HACA rates on primary particle diameter, d_p using KAUST and MPBM and different PAH growth model

There is evidence supporting the coalescence (or sintering) of soot particles at initial stages of formation. The TEM images collected thermophoretically [171] near the inception zone of a burner-stabilized laminar premixed ethylene–oxygen–argon flame ($\phi = 2.5$) showed particles with apparent sizes larger than those measured by Scanning Mobility Particle Sizer (SMPS) [172] suggesting that

the particles flatten on the TEM grid upon impact. In other words, incipient and nascent soot particles may have a liquid-like state that quickly transforms to a fully solid particles due to soot aging and maturity[45]. At intial stages, these particle can coalesce (merge) upon collision forming larger primary particles. There are a few experimental studies on soot coalescence. Ono et al. [17] reheated certain size groups of soot particles formed during pyrolysis of ethylene and benzene in a secondary reactor and observed the morphological changes by comparing the TEM images. Matsukawa et al. [173] proposed a relation based on solid sintering mechanism of metallic nanoparticles to describe coalescence characteristic time of soot as:

$$\tau_c^i = A(d_p^i)^4 \cdot T \cdot \exp\left(\frac{E_a}{RT}\right), \quad (3.5)$$

where $A = 3.0110^{22}[\text{s}/\text{m}^4 - \text{K}]$ and $E_a = 1.4010^5[\text{J/mol}]$ are the prefactor, and activation energy, respectively, that are determined by curve-fitting to the change in primary particle diameter of soot after reheating in a flow reactor. Coalescence does not affect the agglomerate mass, but rather increase the surface area by merging primary particles. Coalescence is described in omnisoott by adding another partial source term to Eq. 2.68 as:

$$(S_{N_{pri}})_{coal} = -\frac{3}{\tau_c^i} \left(n_p^i - (n_p^i)^{2/3} \right) N_{agg}^i \quad (3.6)$$

Fig.3.34 depicts the decrease in f_v by considering coalescence because of reduction in the surface area of agglomerates hence the number of active sites for HACA growth. As shown in Fig.3.35, coalescence increases d_p with a maximum of 64% for Dimer Coalescence, but it cannot account for the significant underprediction of primary particle diameter by the model.

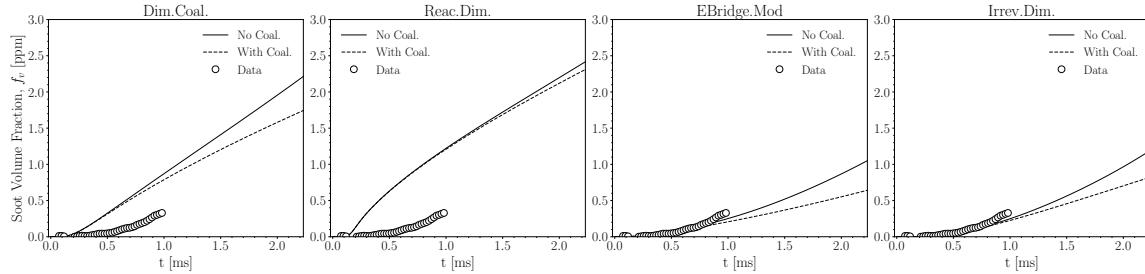


Figure 3.34: The effect of coalescence onsoot volume fraction, f_v using KAUST, SPBM and different PAH growth model

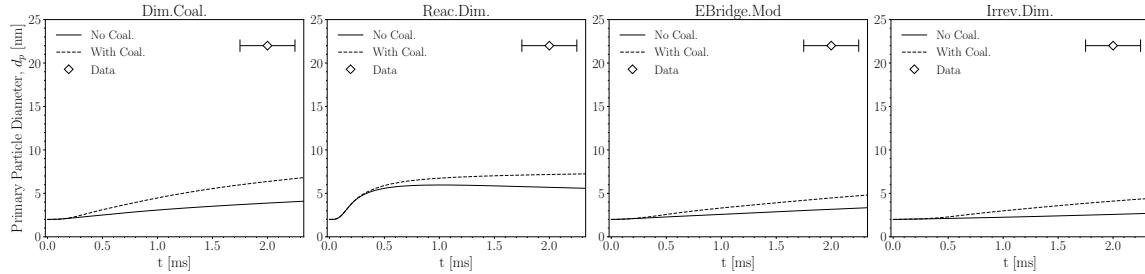


Figure 3.35: The effect of coalescence on primary particle diameter, d_p using KAUST and SPBM and different PAH growth model

The next step is examining the effect on inception and PAH adsorption on predicted soot volume fraction and primary particle diameter. This is achieved by introducing two prefactors, δ and η to adjust inception flux and PAH adsorption rate, respectively. Nine different simulation cases are designed by assigning 1, 0.01, 0.001 to δ and 1, 10, 100 to η where the case with $\delta, \eta = 1$

corresponds to the default PAH growth parameters. Fig. 3.36 and 3.37 shows modifying δ and η values in Reactive Dimerization model can significantly change the f_v and d_p . Modified primary particle diameter values increased large enough to cover the gap between TEM measurements and simulations. By introducing the same prefactors in the rest of PAH growth model, we noticed that d_p exhibits a similar sensitivity with respect to δ and η values. Therefore, the inception flux and increasing the PAH adsorption rate are the major factors that determine primary particle size and soot yield in shock-tube simulations.

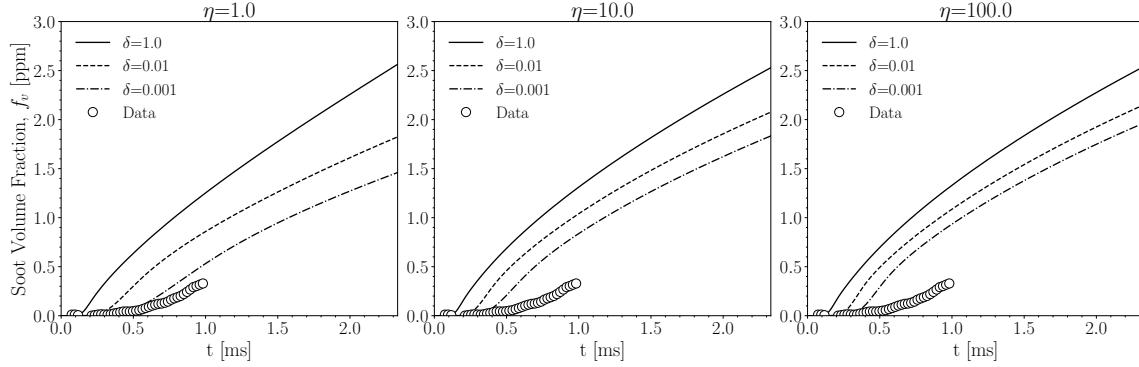


Figure 3.36: The effect of inception flux and PAH adsorption rate on soot volume fraction, f_v , using KAUST, Reactive Dimerization and SPBM

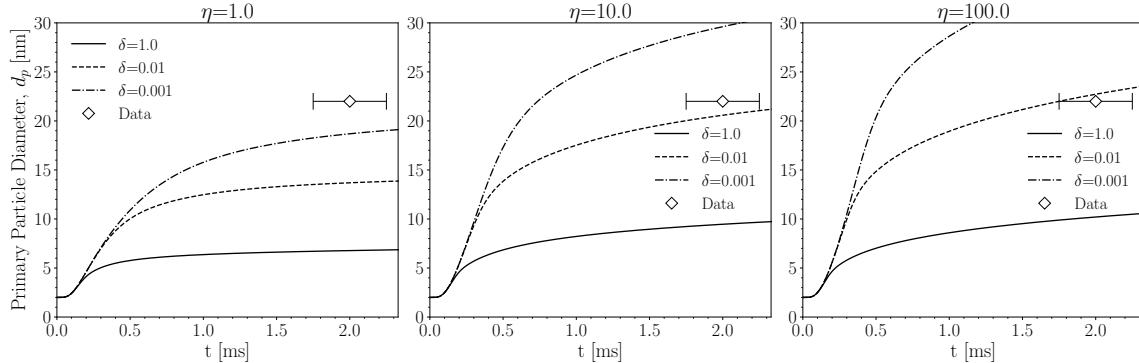


Figure 3.37: The effect of inception flux and PAH adsorption rate on primary particle diameter, d_p using KAUST, Reactive Dimerization and SPBM

Now that the importance of inception and PAH adsorption is demonstrated, a systematic analysis will be conducted for each PAH growth model by altering δ and η to determine the rate constants that minimize the prediction error of both f_v and d_p . We employ a two-step grid search for all PAH growth models. In the first step, both δ and η are divided and multiplied by a series factors ranging 1- 10^4 , respectively, leading to 81 simulations in total. The simulation case with the lowest error is used as the starting point for the second grid search by perturbing δ and η between -50% and +50%, leading to 45 simulations. It is important to note that the purpose of optimization is not necessarily finding the global minimum of cost function (the combined error of d_p and f_v), but rather to show that it is possible to predict both f_v and d_p in good agreement with data by adjusting the inception flux and PAH adsorption rate ensuring that the rate constants stay within their physical limit. The agreement between of data and predictions using optimized PAH growth models shown in Fig. 3.38 confirms the above-mentioned conclusion. Interestingly, the predicted N_{pir} by all PAH growth models reach a plateau of $10^{11} [1/cm^3]$ and fall in the range inferred from extinction and TEM measurements. The inception flux, I_{inc} of PAH growth models shows similar

behavior spanning over the range of $\approx 3 \times 10^{12} [1/cm^3 - s]$ to $\approx 3 \times 10^{14} [1/cm^3 - s]$. But, they also have noticeable differences. While Dimer Coalescence starts with a quick rise in inception flux followed by a decrease toward the end of test time, Reactive Dimerization predicts a more uniform inception over the simulated time.

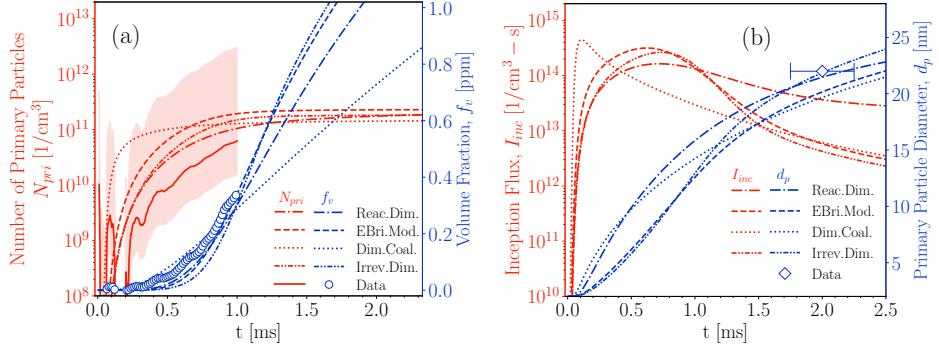


Figure 3.38: The soot volume fraction, f_v and number of primary particles, N_{pri} (left pane), the soot inception flux, I_{inc} and the primary particle diameter, d_p predicted by KAUST mechanism, SPBM, and different PAH growth models with optimized rate at $T=2230$ K, $P=4.5$ atm that corresponds to shock-tube conditions of TEM measurement

The PAH growth models with optimized rates were applied to the whole temperature and pressure ranges (see Fig. B.23 and B.24). As shown in Fig. 3.39, the measured f_v demonstrates a relatively strong sensitivity to temperature best captured by EBridge Formation. This model start with a temperature-dependent dehydrogenation step, but other model are initialed with PAH collision that weakly depends on temperature. On the other hand, Dimer Coalescence is least sensitive to temperature. While it slightly underestimates f_v at $T=2085$ K, it predicts soot volume fraction four time larger than data at $T=2406$ K. The evolution of primary particle diameter, d_p seems to not be considerably affected with temperature in $2000\text{ K} < T < 2400$ K. The d_p using optimized PAH growth models is calculated at $t=2$ ms and plotted for the studied T_5 range in Fig. 3.41 that shows a bell-like temperature dependency with peak around 2200 K for all PAH growth models.

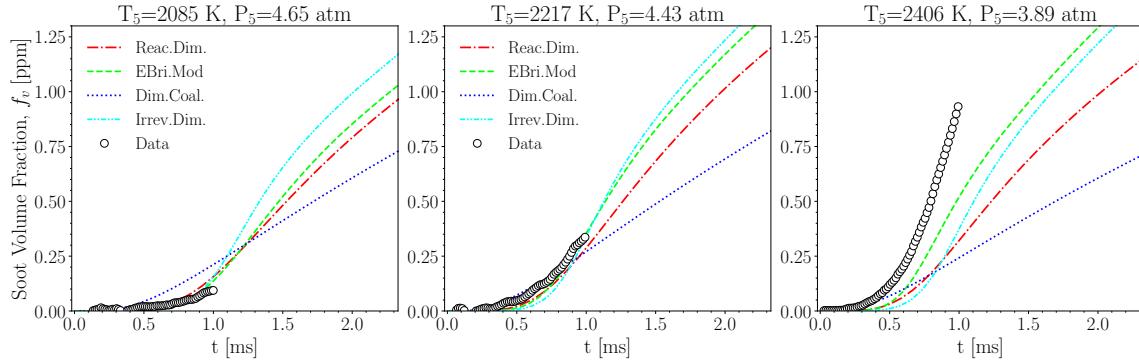


Figure 3.39: The soot volume fraction, f_v over $2000\text{ K} < T < 2400$ K, using KAUST, SPBM, and optimized inception models

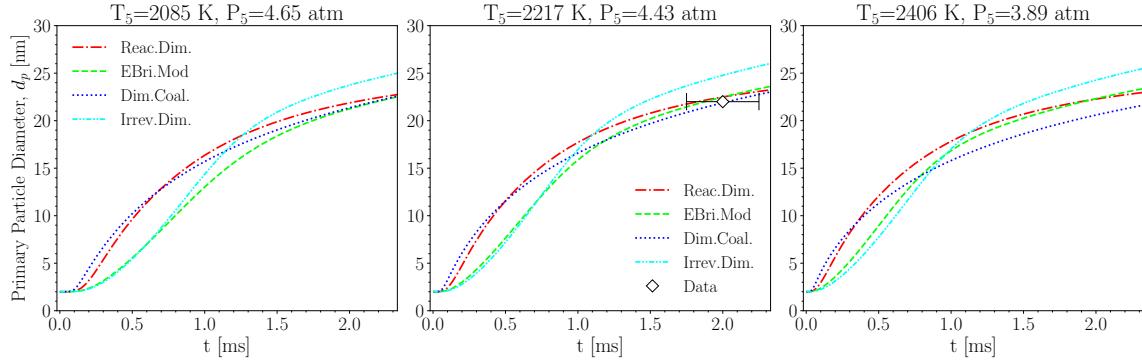


Figure 3.40: The primary particle diameter, d_p over $2000\text{ K} < T < 2400\text{ K}$, using KAUST, SPBM, and optimized inception models

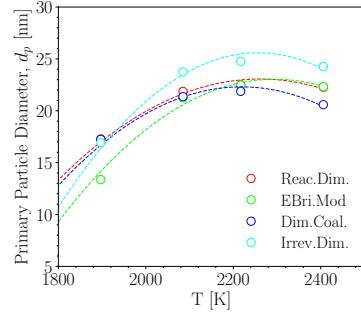


Figure 3.41: The primary particle diameter, d_p at $t=2\text{ ms}$ over the studied temperature range using KAUST, SPBM, and optimized inception models. The dashed line is added to guide the eye.

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Appendix A

Derivation of Transport Equations

This chapter explains the derivation of transport equations from first principles for mass, species, energy and soot variables for the reactors employed in omnisooot. The conventional reactor transport equation have been derived for a reacting gas without solid particle [133], so we provide a step by step derivation of equations that describes transport/evolution of mass, species and energy of gas as well as those of soot variables.

A.1 Constant Volume Reactor

A.1.1 Continuity

The rate change of mass of gas mixture is equal to the production rate of soot.

$$\frac{d}{dt} (m) = V (1 - \varphi) \sum_i \dot{s}_i W_i,$$

where V is the reactor volume that stays constant during the process, and gas occupies a fraction of volume reactor without soot, so $m = \rho V(1 - \varphi)$.

$$\begin{aligned} \frac{d}{dt} (\rho V(1 - \varphi)) &= V (1 - \varphi) \sum_i \dot{s}_i W_i \\ &\Downarrow \\ \frac{d}{dt} (\rho(1 - \varphi)) &= (1 - \varphi) \sum_i \dot{s}_i W_i \end{aligned} \tag{A.1}$$

A.1.2 Species

The rate change of mass of each species can be described as:

$$\frac{d}{dt} (m_k) = V (1 - \varphi) (\dot{\omega}_k + \dot{s}_k) W_k$$

The mass of each species is defined as the total gas mass multiplied by the mass fraction as

$$\frac{dY_k}{dt} \rho V(1 - \varphi) + Y_k \frac{d}{dt} (\rho V(1 - \varphi)) = V (1 - \varphi) (\dot{\omega}_k + \dot{s}_k) W_k$$

where the colored term can be substituted from Equation (A.1).

$$\frac{dY_k}{dt} \rho V(1 - \varphi) + Y_k \cancel{V(1 - \varphi)} \sum_i \dot{s}_i W_i = V (1 - \varphi) (\dot{\omega}_k + \dot{s}_k) W_k$$

$$\Downarrow \times \frac{1}{\rho V(1-\varphi)}$$

$$\frac{dY_k}{dt} = \frac{1}{\rho} \left[(\dot{\omega}_k + \dot{s}_k) W_k - Y_k \sum_i \dot{s}_i W_i \right] \quad (\text{A.2})$$

A.1.3 Energy

In the constant volume reactor, the internal energy is favorable since the external boundaries of the control volume are not changed, and the work from soot volume is neglected, so formulating the energy balance based on the internal energy is preferred. The energy passes through the boundaries of the control volume via soot generation and external heat source.

$$\begin{aligned} \frac{dE}{dt} &= -\frac{dE_{soot}}{dt} + \dot{Q} \\ \Downarrow E &= m \sum Y_k e_k \\ \frac{dm}{dt} \sum_k Y_k e_k + m \sum_k \frac{dY_k}{dt} e_k + m \sum_k Y_k \frac{de_k}{dt} &= -\frac{dE_{soot}}{dt} + \dot{Q} \\ \Downarrow \text{from continuity} \\ V(1-\varphi) \sum_i \dot{s}_i W_i \sum_k Y_k e_k + m \sum_k \frac{dY_k}{dt} e_k + m \sum_k Y_k \frac{de_k}{dt} &= -\frac{dE_{soot}}{dt} + \dot{Q} \\ \Downarrow \text{from species} \\ V(1-\varphi) \sum_i \dot{s}_i W_i \sum_k Y_k e_k + m \sum_k \frac{1}{\rho} \left[(\dot{\omega}_k + \dot{s}_k) W_k - Y_k \sum_i \dot{s}_i W_i \right] e_k + m \sum_k Y_k \frac{de_k}{dt} &= -\frac{dE_{soot}}{dt} + \dot{Q} \\ \Downarrow Y_k \frac{de_k}{dt} &= Y_k c_{c,k} \frac{dT}{dt} \\ \Downarrow \sum_k Y_k c_{c,k} &= c_p \\ V(1-\varphi) \sum_i \dot{s}_i W_i \sum_k Y_k e_k + m \sum_k \frac{1}{\rho} \left[(\dot{\omega}_k + \dot{s}_k) W_k - Y_k \sum_i \dot{s}_i W_i \right] e_k + m c_v \frac{dT}{dt} &= -\frac{dE_{soot}}{dt} + \dot{Q} \\ \Downarrow \frac{dE_{soot}}{dt} &= \left(\rho_{soot} c_{v,soot} \varphi \frac{dT}{dt} - (1-\varphi) \sum_i \dot{s}_i W_i e_{soot} \right) V \\ V(1-\varphi) \sum_i \dot{s}_i W_i \sum_k Y_k e_k + m \sum_k \frac{1}{\rho} \left[(\dot{\omega}_k + \dot{s}_k) W_k - Y_k \sum_i \dot{s}_i W_i \right] e_k + m c_v \frac{dT}{dt} &= \\ - \left(\rho_{soot} c_{v,soot} \varphi \frac{dT}{dt} - (1-\varphi) \sum_i \dot{s}_i W_i e_{soot} \right) V + \dot{Q} & \end{aligned}$$

$$\Downarrow m = \rho V(1-\varphi)$$

$$\begin{aligned}
& V(1-\varphi) \sum_i \dot{s}_i W_i \sum_k Y_k e_k + V(1-\varphi) \sum_k (\dot{\omega}_k + \dot{s}_k) W_k e_k - V(1-\varphi) \sum_k Y_k e_k \sum_i \dot{s}_i W_i \\
& + \rho V(1-\varphi) c_v \frac{dT}{dt} = - \left(\rho_{soot} c_{v, soot} \varphi \frac{dT}{dt} - (1-\varphi) \sum_i \dot{s}_i W_i e_{soot} \right) V + \dot{Q} \\
& \Downarrow \text{Eliminating } V(1-\varphi) \sum_i \dot{s}_i W_i \sum_k Y_k e_k \\
& \Downarrow \times 1/V \\
& (1-\varphi) \sum_k (\dot{\omega}_k + \dot{s}_k) W_k e_k + \rho(1-\varphi) c_v \frac{dT}{dt} = - \rho_{soot} c_{v, soot} \varphi \frac{dT}{dt} + \\
& (1-\varphi) \sum_i \dot{s}_i W_i e_{soot} + \dot{Q} \\
& \Downarrow \\
& (\rho(1-\varphi) c_v + \rho_{soot} c_{v, soot} \varphi) \frac{dT}{dt} = - (1-\varphi) \sum_k e_k (\dot{\omega}_k + \dot{s}_k) W_k + (1-\varphi) \sum_i \dot{s}_i W_i e_{soot} + \frac{\dot{Q}}{V} \\
& \Downarrow \\
& \frac{dT}{dt} = \frac{1}{\rho(1-\varphi) c_v + \rho_{soot} c_{v, soot} \varphi} \left[- (1-\varphi) \sum_k e_k (\dot{\omega}_k + \dot{s}_k) W_k + (1-\varphi) \sum_i \dot{s}_i W_i e_{soot} + \frac{\dot{Q}}{V} \right] \quad (\text{A.3})
\end{aligned}$$

A.1.4 Soot Variables

The derivation of equations for soot variables can be accomplished similar to species evolution.

$$\frac{d}{dt} (m\psi) = \rho V(1-\varphi) S_\psi$$

\Downarrow

$$\frac{d}{dt} (\rho V(1-\varphi\psi)) = \rho V(1-\varphi) S_\psi$$

$\Downarrow \times 1/V$

$$\rho(1-\varphi) \frac{d\psi}{dt} + \psi \frac{d}{dt} (\rho(1-\varphi)) = \rho(1-\varphi) S_\psi$$

\Downarrow from continuity

$$\begin{aligned}
\rho(1-\varphi) \frac{d\psi}{dt} + \psi V(1-\varphi) \sum_i \dot{s}_i W_i &= \rho(1-\varphi) S_\psi \\
\Downarrow \times \frac{1}{\rho(1-\varphi)} \\
\frac{d\psi}{dt} &= S_\psi - \frac{1}{\rho} \left[\psi \sum_i \dot{s}_i W_i \right]
\end{aligned} \tag{A.4}$$

A.2 Plug Flow Reactor

A.2.1 Continuity

The gas flow that passes through the cross-section of a differential element along the flow reactor is described as:

$$\begin{aligned}
\rho u A (1-\varphi) |_{z+dz} - \rho u A (1-\varphi) |_z &= Adz(1-\varphi) \sum_i \dot{s}_i W_i \\
\Downarrow \\
\rho u (1-\varphi) A &= Adz(1-\varphi) \sum_i \dot{s}_i W_i \\
\Downarrow \\
\frac{d}{dz} (\rho u (1-\varphi)) &= (1-\varphi) \sum_i \dot{s}_i W_i
\end{aligned} \tag{A.5}$$

A.2.2 Momentum

$$\begin{aligned}
\frac{d}{dz} (\dot{m} u) &= -\frac{d}{dz} (P(1-\varphi)) - \tau_w \frac{P_c}{A} \\
\Downarrow \\
u \frac{d}{dz} (\dot{m}) + \dot{m} \frac{d}{dz} (u) &= -\frac{d}{dz} (P(1-\varphi)) - \tau_w \frac{1}{R_H} \\
\Downarrow \\
u(1-\varphi) \sum_i \dot{s}_i W_i + \rho u (1-\varphi) \frac{d}{dz} (u) &= -\frac{d}{dz} (P(1-\varphi)) - \frac{\tau_w}{R_H}
\end{aligned} \tag{A.6}$$

A.2.3 Species

$$\frac{d}{dz} (m_k) = (1 - \varphi) (\dot{\omega}_k + \dot{s}_k) W_k$$

↓

$$\frac{d}{dz} (\rho u (1 - \varphi) Y_k) = (1 - \varphi) (\dot{\omega}_k + \dot{s}_k) W_k$$

↓

$$\rho u (1 - \varphi) \frac{dY_k}{dz} + Y_k \frac{d}{dz} (\rho u (1 - \varphi)) = (1 - \varphi) (\dot{\omega}_k + \dot{s}_k) W_k$$

↓ from Equation (A.5)

$$\rho u (1 - \varphi) \frac{dY_k}{dz} + Y_k (1 - \varphi) \sum_i \dot{\omega}_i W_i = (1 - \varphi) (\dot{\omega}_k + \dot{s}_k) W_k$$

↓ × 1/(ρ(1 - φ))

$$\frac{dY_k}{dz} = \frac{1}{\rho u (1 - \varphi)} \left[(1 - \varphi) (\dot{\omega}_k + \dot{s}_k) W_k - Y_k (1 - \varphi) \sum_i \dot{s}_i W_i \right] \quad (\text{A.7})$$

A.2.4 Energy

$$\frac{d}{dz} (\rho u (1 - \varphi) h) + \frac{d}{dz} (\rho u \varphi h_{soot}) = \dot{q}'$$

↓

$$\frac{dh}{dz} (\rho u (1 - \varphi)) + h \frac{d}{dz} (\rho u (1 - \varphi)) + \frac{d}{dz} (\rho u \varphi) h_{soot} + \frac{dh_{soot}}{dz} \rho u \varphi = \dot{q}'$$

$$\downarrow h = \sum_k Y_k h_k$$

$$\rho u (1 - \varphi) \frac{d}{dz} \left(\sum_k Y_k h_k \right) + h \frac{d}{dz} (\rho u (1 - \varphi)) + \frac{d}{dz} (\rho u \varphi) h_{soot} + \frac{dh_{soot}}{dz} \rho u \varphi = \dot{q}'$$

↓

$$\rho u (1 - \varphi) \sum_k \frac{dY_k}{dz} h_k + \rho u (1 - \varphi) \sum_k \frac{dh_k}{dz} Y_k + h \frac{d}{dz} (\rho u (1 - \varphi)) + \frac{d}{dz} (\rho u \varphi) h_{soot} + \frac{dh_{soot}}{dz} \rho u \varphi = \dot{q}'$$

\Downarrow from Equation (A.5)&(A.7)

$$\rho u (1 - \varphi) \sum_k \frac{1}{\rho u} \left[(\dot{\omega}_k + \dot{s}_k) W_k - Y_k \sum_i \dot{s}_i W_i \right] h_k + \rho u (1 - \varphi) \sum_k \frac{dh_k}{dz} Y_k + \\ h \frac{d}{dz} \left((1 - \varphi) \sum_i \dot{s}_i W_i \right) + \frac{d}{dz} (\rho u \varphi) h_{soot} + \frac{dh_{soot}}{dz} \rho u \varphi = \dot{q}'$$

$$\Downarrow \frac{dh_k}{dz} = c_{p,k} \frac{dT}{dz}$$

$$\Downarrow \sum_k Y_k c_{p,k} = c_p$$

$$\Downarrow \frac{dh_{soot}}{dz} = c_{p,soot} \frac{dT}{dz}$$

$$(1 - \varphi) \sum_k \left[(\dot{\omega}_k + \dot{s}_k) W_k - Y_k \sum_i \dot{s}_i W_i \right] h_k + \rho u (1 - \varphi) c_p \frac{dT}{dz} + h (1 - \varphi) \sum_i \dot{s}_i W_i \\ + \frac{d}{dz} (\rho u \varphi) h_{soot} + c_{p,soot} \frac{dT}{dz} \rho u \varphi = \dot{q}'$$

\Downarrow

$$(1 - \varphi) \sum_k [(\dot{\omega}_k + \dot{s}_k) W_k] h_k + \rho u (1 - \varphi) c_p \frac{dT}{dz} - (1 - \varphi) \sum_i \dot{s}_i W_i h_{soot} + c_{p,soot} \frac{dT}{dz} \rho u \varphi = \dot{q}'$$

\Downarrow

$$[\rho u (1 - \varphi) c_p + \rho u \varphi c_{p,soot}] \frac{dT}{dz} + (1 - \varphi) \sum_k [(\dot{\omega}_k + \dot{s}_k) W_k] h_k - (1 - \varphi) \sum_i \dot{s}_i W_i h_{soot} = \dot{q}'$$

\Downarrow

$$\frac{dT}{dz} = \frac{1}{\rho u (1 - \varphi) c_p + \rho u \varphi c_{p,soot}} \left[- (1 - \varphi) \sum_k (\dot{\omega}_k + \dot{s}_k) W_k h_k + (1 - \varphi) \sum_i \dot{s}_i W_i h_{soot} + \dot{q}' \right]$$

A.3 Perfectly Stirred Reactor

A.3.1 Continuity

The mass conservation for the partially stirred reactor can be written as:

$$\frac{dm}{dt} = \dot{m}_{in} - \dot{m}_{out} + V (1 - \varphi) \sum_i \dot{s}_i W_i \quad (\text{A.8})$$

A.3.2 Species

$$\frac{d}{dt} (m_k) = \dot{m}_{in} Y_k^* - \dot{m}_{out} Y_k + V (1 - \varphi) (\dot{\omega}_k + \dot{s}_k) W_k$$

$$\Downarrow m_k = m Y_k$$

$$Y_m \frac{dm}{dt} + m (Y_k) = \dot{m}_{in} Y_k^* - \dot{m}_{out} Y_k + V (1 - \varphi) (\dot{\omega}_k + \dot{s}_k) W_k$$

\Downarrow from Equation A.8

$$\left(\dot{m}_{in} - \dot{m}_{out} + V (1 - \varphi) \sum_i \dot{s}_i W_i \right) Y_k + \frac{dY_k}{dt} m = \dot{m}_{in} Y_k^* - \dot{m}_{out} Y_k + V (1 - \varphi) (\dot{\omega}_k + \dot{s}_k) W_k$$

$$\Downarrow$$

$$V (1 - \varphi) Y_k \sum_i \dot{s}_i W_i + \frac{dY_k}{dt} m = \dot{m}_{in} (Y_k^* - Y_k) + V (1 - \varphi) (\dot{\omega}_k + \dot{s}_k) W_k$$

$$\Downarrow$$

$$\frac{dY_k}{dt} \rho V (1 - \varphi) = \dot{m}_{in} (Y_k^* - Y_k) + V (1 - \varphi) (\dot{\omega}_k + \dot{s}_k) W_k - V (1 - \varphi) Y_k \sum_i \dot{s}_i W_i$$

$$\Downarrow$$

$$\frac{dY_k}{dt} = \frac{\dot{m}_{in}}{\rho V (1 - \varphi)} (Y_k^* - Y_k) + \frac{1}{\rho} (\dot{\omega}_k + \dot{s}_k) W_k - \frac{1}{\rho} Y_k \sum_i \dot{s}_i W_i \quad (\text{A.9})$$

A.3.3 Energy

$$\frac{dH}{dt} = \dot{m}_{in} h^* - \dot{m}_{out} h - \frac{dH_{soot}}{dt} + \dot{Q}$$

$$\Downarrow H = mh$$

$$\frac{dm}{dt} h + m \frac{dh}{dt} = \dot{m}_{in} h^* - \dot{m}_{out} h - \frac{dH_{soot}}{dt} + \dot{Q}$$

\Downarrow from Equation A.8

$$\left(\dot{m}_{in} - \dot{m}_{out} + V (1 - \varphi) \sum_i \dot{s}_i W_i \right) h + m \frac{dh}{dt} = \dot{m}_{in} h^* - \dot{m}_{out} h - \frac{dH_{soot}}{dt} + \dot{Q}$$

$$\Downarrow$$

$$V(1-\varphi)h \sum_i \dot{s}_i W_i + m \frac{dh}{dt} = \dot{m}_{in}(h^* - h) - \frac{dH_{soot}}{dt} + \dot{Q}$$

↓

$$m \frac{dh}{dt} = \dot{m}_{in}(h^* - h) - V(1-\varphi)h \sum_i \dot{s}_i W_i - \frac{dH_{soot}}{dt} + \dot{Q}$$

$$\Downarrow h = \sum Y_k h_k$$

$$m \sum \frac{dY_k}{dt} h_k + m \sum \frac{dh_k}{dt} Y_k = \dot{m}_{in}(h^* - h) - V(1-\varphi)h \sum_i \dot{s}_i W_i - \frac{dH_{soot}}{dt} + \dot{Q}$$

↓ from Equation A.9

$$m \sum \left(\frac{\dot{m}_{in}}{\rho V(1-\varphi)} (Y_k^* - Y_k) + \frac{1}{\rho} (\dot{\omega}_k + \dot{s}_k) W_k - \frac{1}{\rho} Y_k \sum_i \dot{s}_i W_i \right) h_k + m \sum c_{p,k} \frac{dT}{dt} Y_k = \\ \dot{m}_{in}(h^* - h) - V(1-\varphi)h \sum_i \dot{s}_i W_i - \frac{dH_{soot}}{dt} + \dot{Q}$$

$$\Downarrow m = \rho V(1-\varphi)$$

$$\dot{m}_{in} \sum (Y_k^* - Y_k) h_k + V(1-\varphi) \sum (\dot{\omega}_k + \dot{s}_k) W_k h_k + mc_p \frac{dT}{dt} = \dot{m}_{in}(h^* - h) - \frac{dH_{soot}}{dt} + \dot{Q}$$

↓

$$\dot{m}_{in} \sum (Y_k^* - Y_k) h_k + V(1-\varphi) \sum (\dot{\omega}_k + \dot{s}_k) W_k h_k + mc_p \frac{dT}{dt} = \dot{m}_{in}(h^* - h) \\ - \left(\rho_{soot} c_{p,soot} \varphi V \frac{dT}{dt} - (1-\varphi) V \sum_i \dot{s}_i W_i h_{soot} \right) + \dot{Q}$$

↓

$$\dot{m}_{in} \sum (Y_k^* - Y_k) h_k + V(1-\varphi) \sum (\dot{\omega}_k + \dot{s}_k) W_k h_k + mc_p \frac{dT}{dt} = \dot{m}_{in}(h^* - h) \\ - \left(\rho_{soot} c_{p,soot} \varphi V \frac{dT}{dt} - (1-\varphi) V \sum_i \dot{s}_i W_i h_{soot} \right) + \dot{Q}$$

$$\Downarrow m = \rho V(1-\varphi)$$

$$\Downarrow \times 1/V$$

$$\begin{aligned}
& \dot{m}_{in} \sum (Y_k^* - Y_k) h_k + V (1 - \varphi) \sum (\dot{\omega}_k + \dot{s}_k) W_k h_k + \rho V (1 - \varphi) c_p \frac{dT}{dt} = \\
& \dot{m}_{in} (h^* - h) - \left(\rho_{soot} c_{p,soot} \varphi V \frac{dT}{dt} - (1 - \varphi) V \sum_i \dot{s}_i W_i h_{soot} \right) + \dot{Q} \\
& \Downarrow \\
& (\rho (1 - \varphi) c_p + \rho_{soot} c_{p,soot} \varphi) \frac{dT}{dt} = \frac{\dot{m}_{in}}{V} (h^* - h) - \frac{\dot{m}_{in}}{V} \sum (Y_k^* - Y_k) h_k \\
& - (1 - \varphi) \sum (\dot{\omega}_k + \dot{s}_k) W_k h_k + (1 - \varphi) \sum_i \dot{s}_i W_i h_{soot} + \frac{\dot{Q}}{V} \\
& \Downarrow \\
& \frac{dT}{dt} = \frac{1}{\rho (1 - \varphi) c_p + \rho_{soot} c_{p,soot} \varphi} \left[\frac{\dot{m}_{in}}{V} (h^* - h) - \frac{\dot{m}_{in}}{V} \sum (Y_k^* - Y_k) h_k - (1 - \varphi) \sum (\dot{\omega}_k + \dot{s}_k) W_k h_k \right. \\
& \quad \left. + (1 - \varphi) \sum_i \dot{s}_i W_i h_{soot} + \frac{\dot{Q}}{V} \right] \tag{A.10}
\end{aligned}$$

Appendix B

Additional Results

B.1 Methane pyrolysis data from Hanson Research Group

B.1.1 30% CH₄ dataset

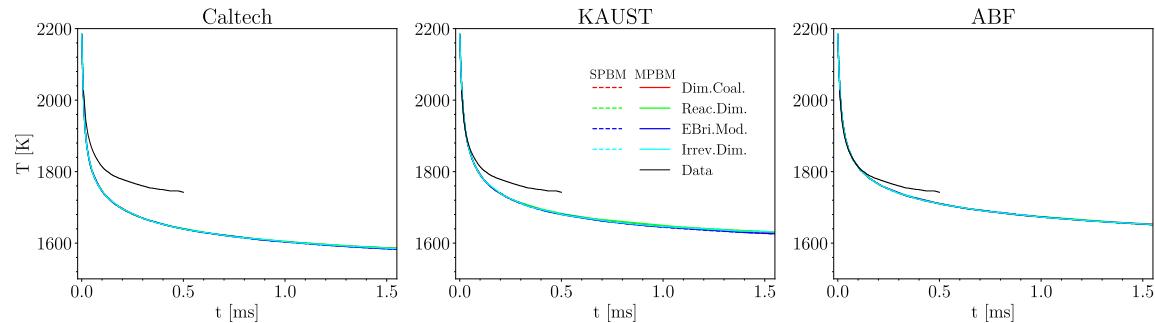


Figure B.1: The time history of temperature of 30% CH₄ pyrolysis at T=2184 K and P=3.62 atm using Caltech, KAUST, and ABF mechanisms and different PAH growth and particle dynamics models compared with the simulation results reported by Hanson Research Group

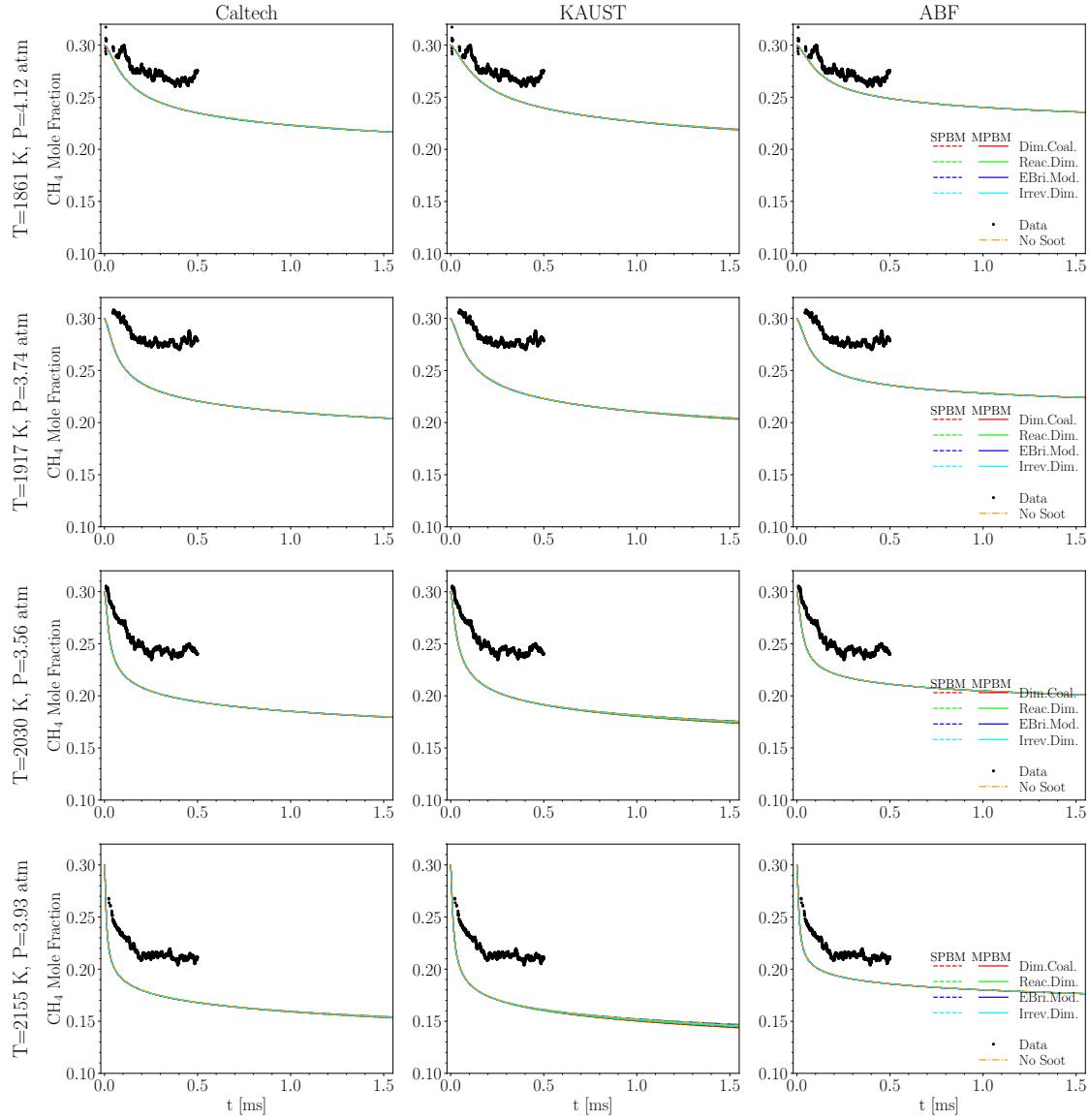


Figure B.2: The time history of CH₄ mole fraction of 30% CH₄ pyrolysis for 1861 K < T < 2155 K using Caltech, KAUST, and ABF mechanisms without soot (only chemistry) and with soot modeling described by different PAH growth and particle dynamics models compared with laser diagnostic data

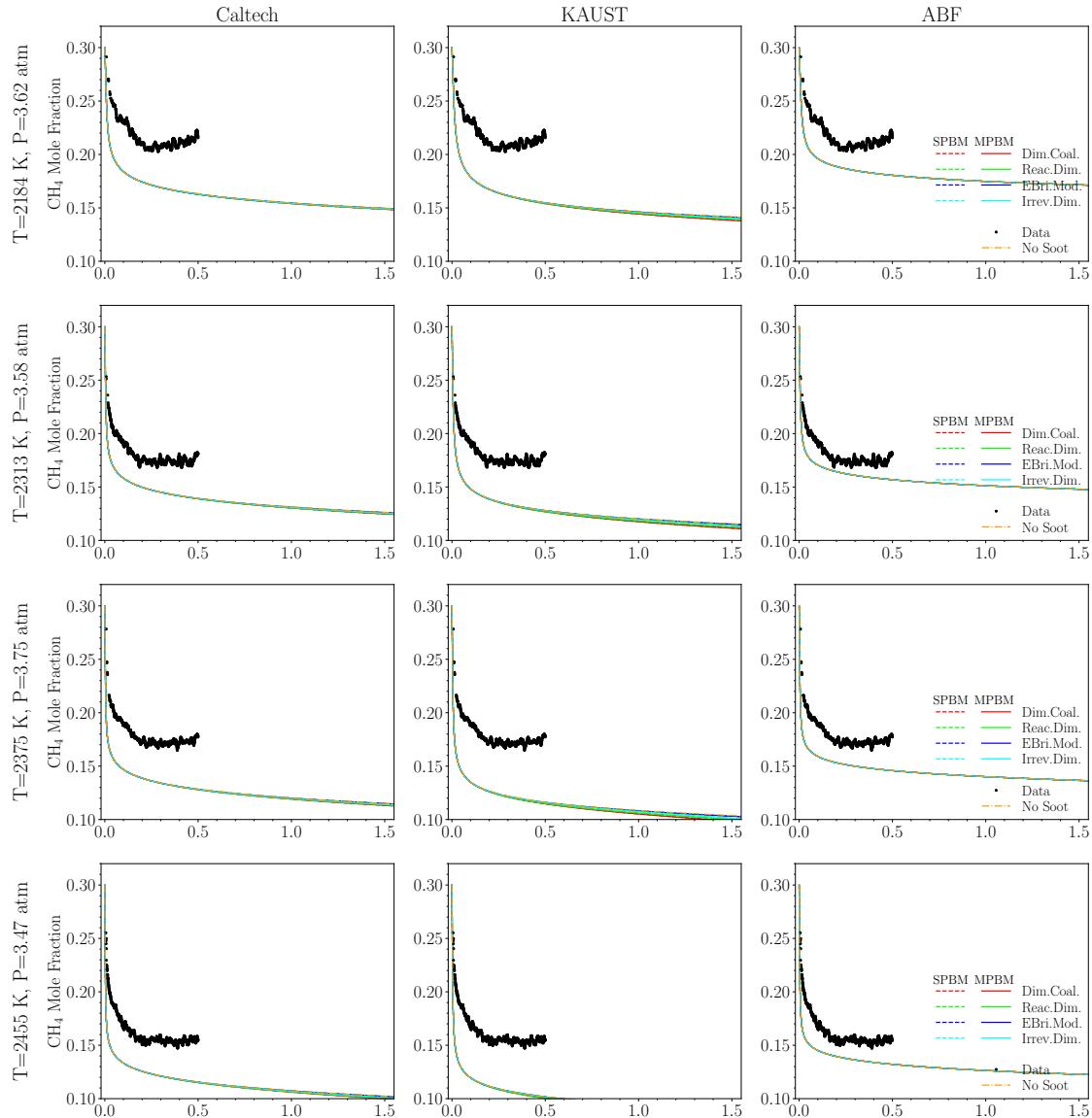


Figure B.3: The time history of CH_4 mole fraction of 30% CH_4 pyrolysis for $2184 \text{ K} < T < 2455 \text{ K}$ using Caltech, KAUST, and ABF mechanism without soot (only chemistry) and with soot modeling described by different PAH growth and particle dynamics models compared with laser diagnostic data

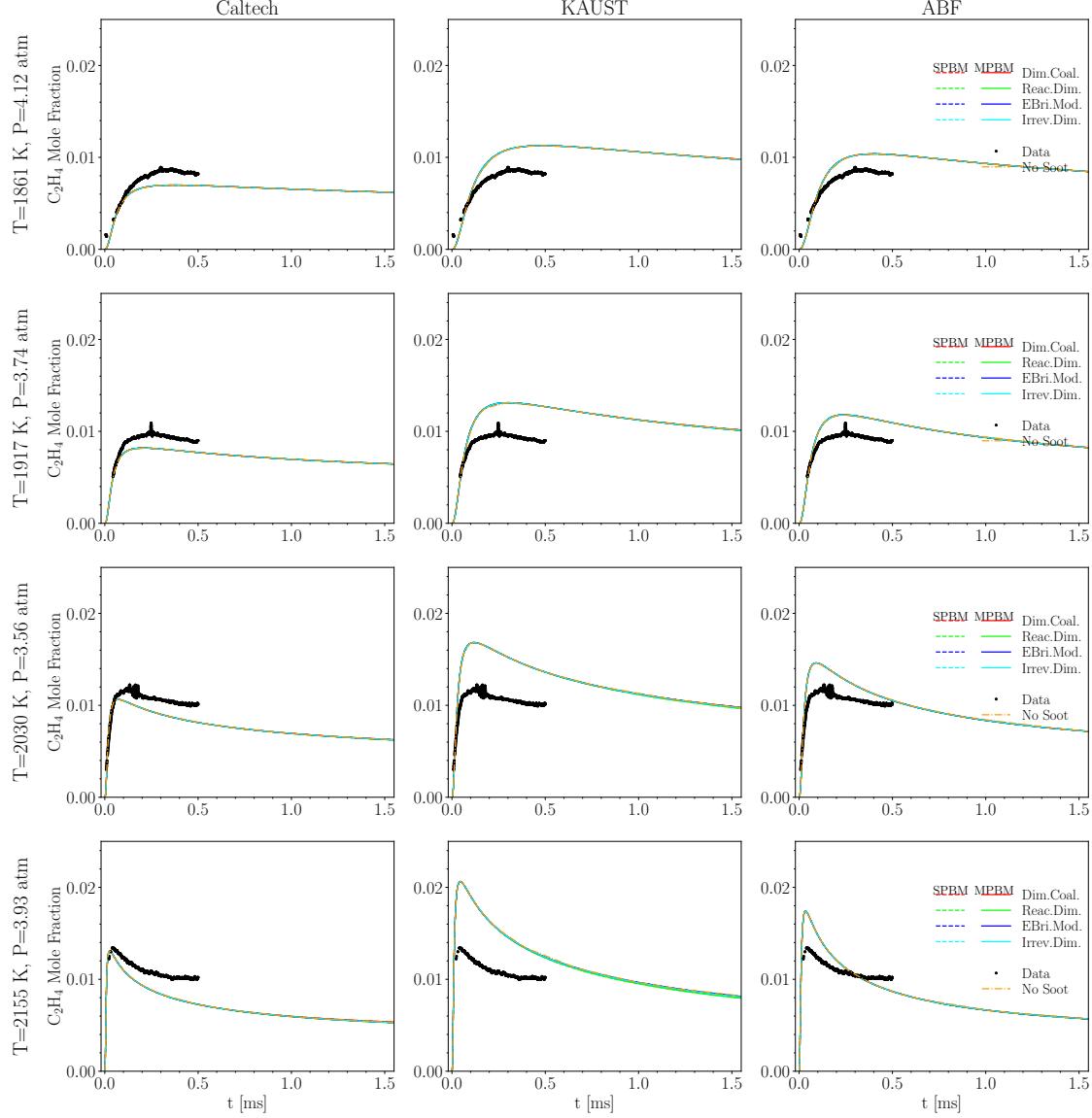


Figure B.4: The time history of C₂H₄ mole fraction of 30% CH₄ pyrolysis for 1861 K < T < 2155 K using Caltech, KAUST, and ABF mechanisms without soot (only chemistry) and with soot modeling described by different PAH growth and particle dynamics models compared with laser diagnostic data

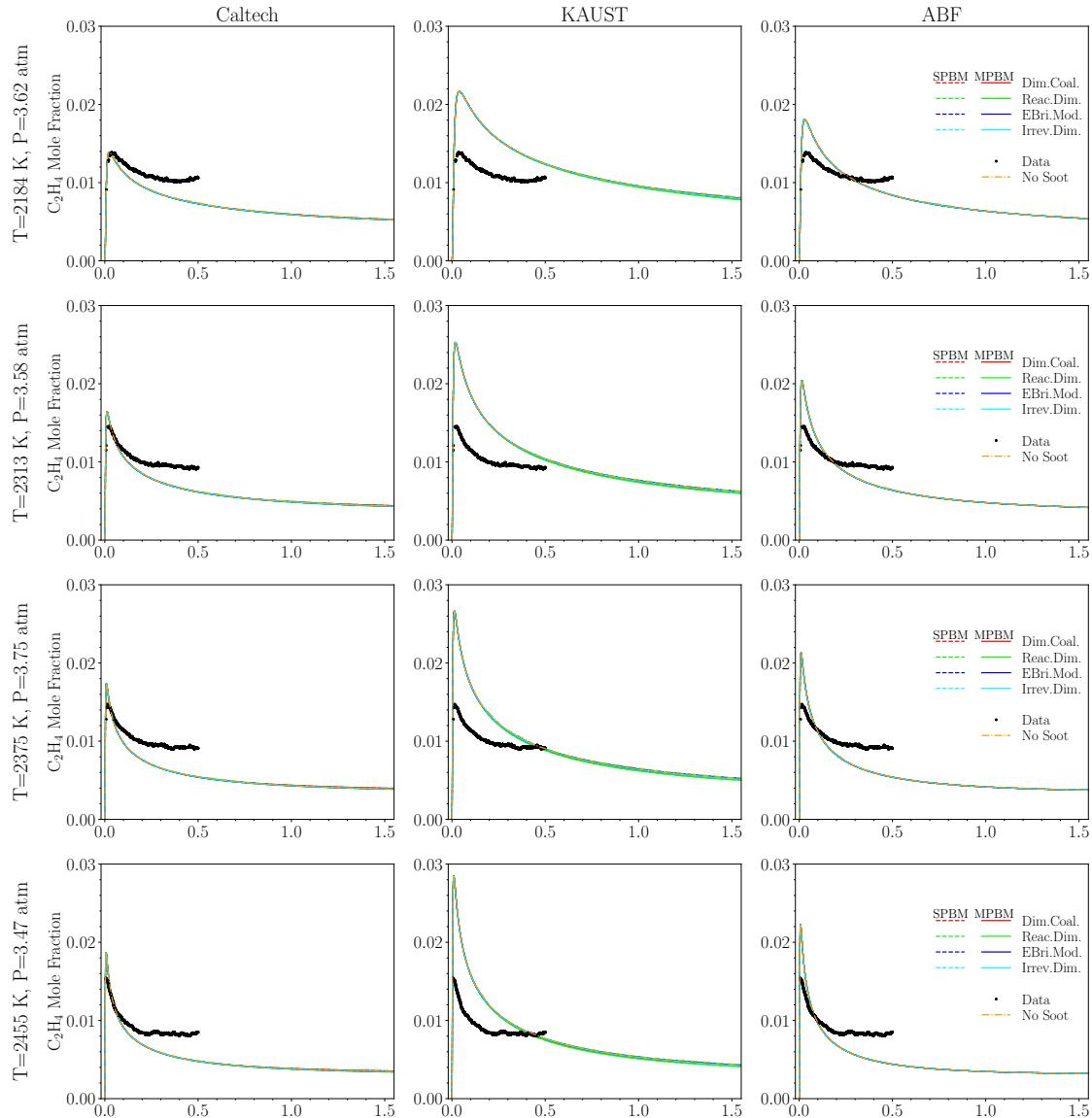


Figure B.5: The time history of C_2H_4 mole fraction of 30% CH_4 pyrolysis for $2184 \text{ K} < T < 2455 \text{ K}$ using Caltech, KAUST, and ABF mechanism without soot (only chemistry) and with soot modeling described by different PAH growth and particle dynamics models compared with laser diagnostic data

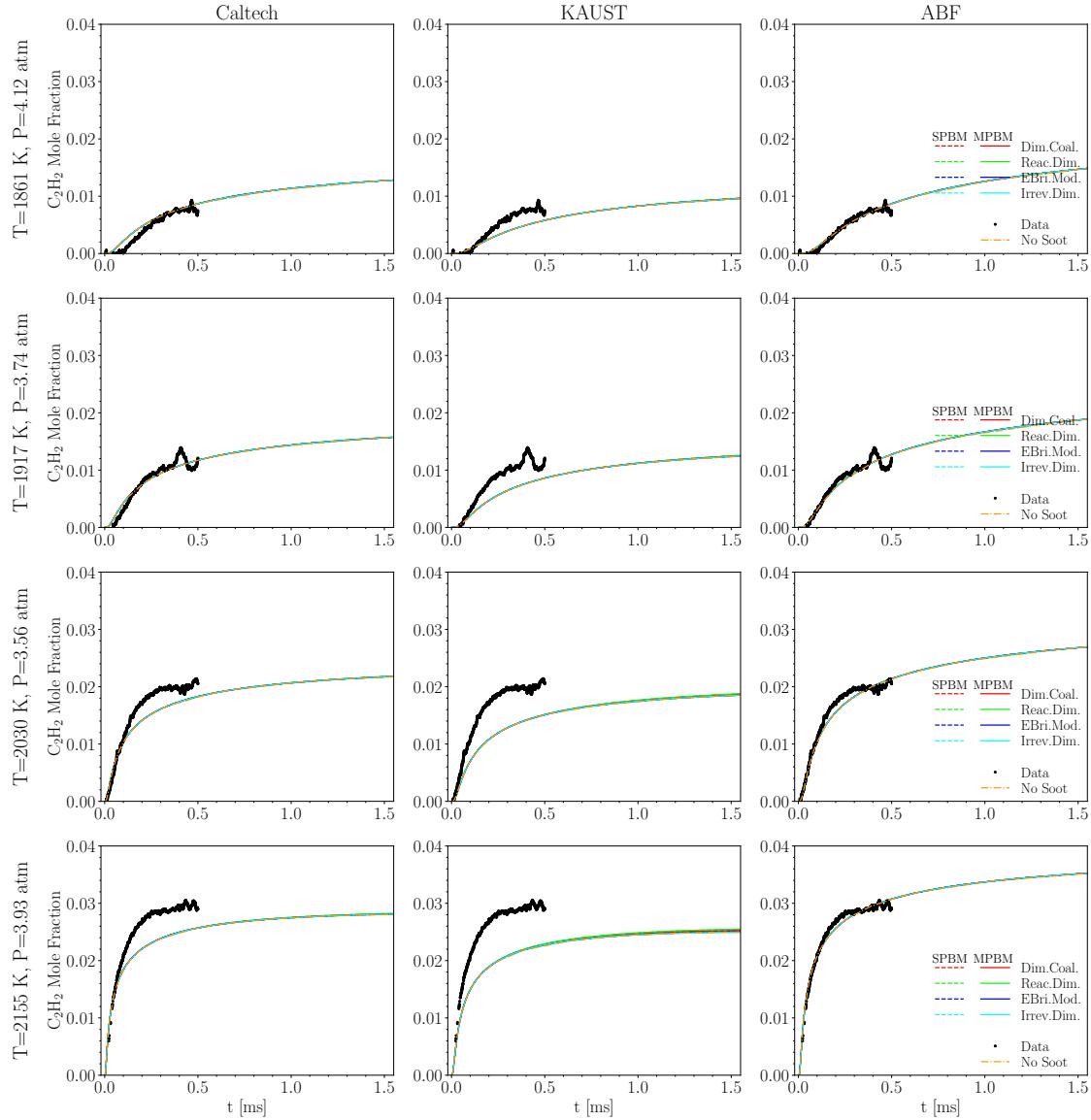


Figure B.6: The time history of C_2H_2 mole fraction of 30% CH_4 pyrolysis for $1861 \text{ K} < T < 2155 \text{ K}$ using Caltech, KAUST, and ABF mechanisms without soot (only chemistry) and with soot modeling described by different PAH growth and particle dynamics models compared with laser diagnostic data

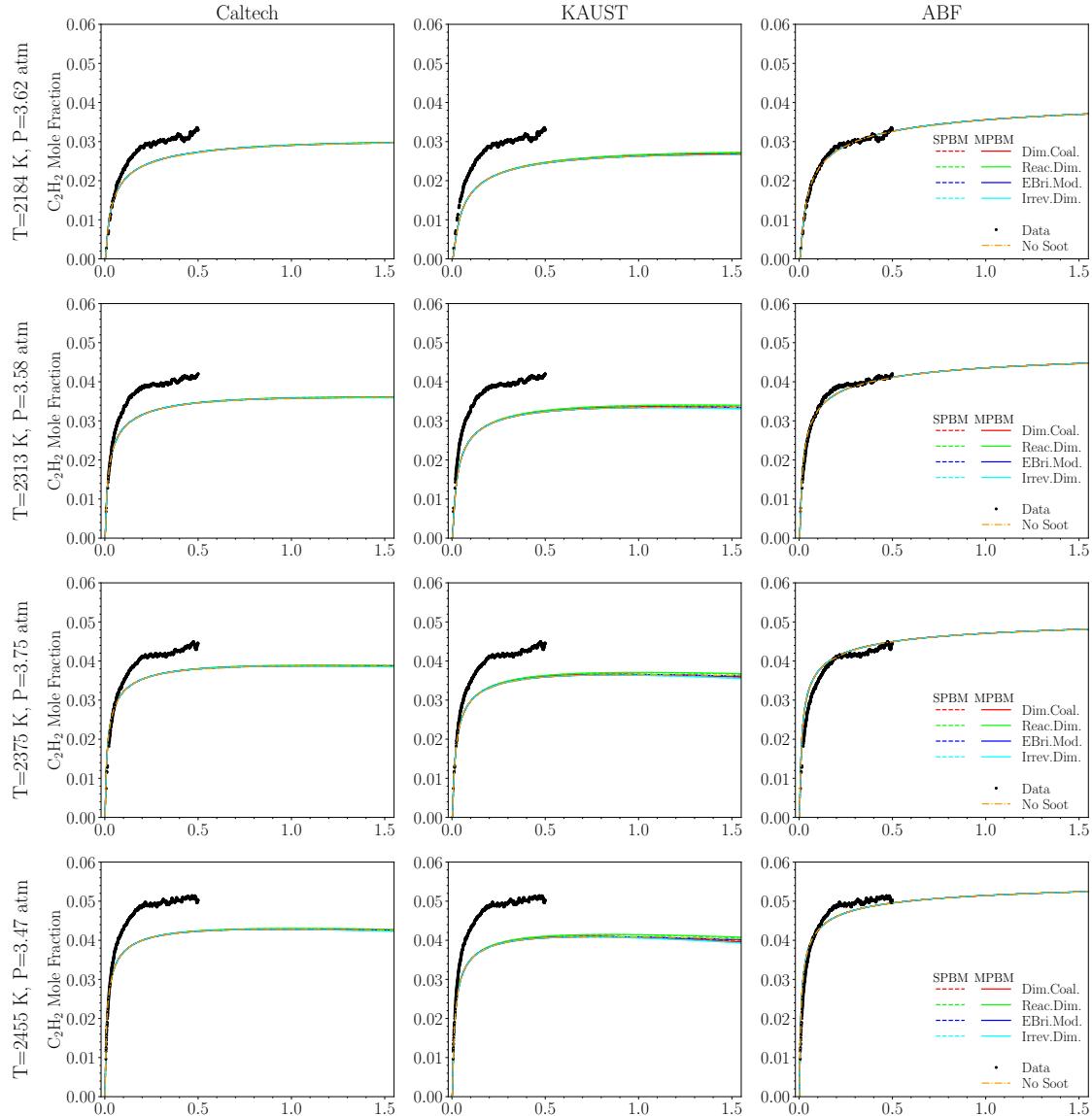


Figure B.7: The time history of C_2H_2 mole fraction of 30% CH_4 pyrolysis for $2184 \text{ K} < T < 2455 \text{ K}$ using Caltech, KAUST, and ABF mechanism without soot (only chemistry) and with soot modeling described by different PAH growth and particle dynamics models compared with laser diagnostic data

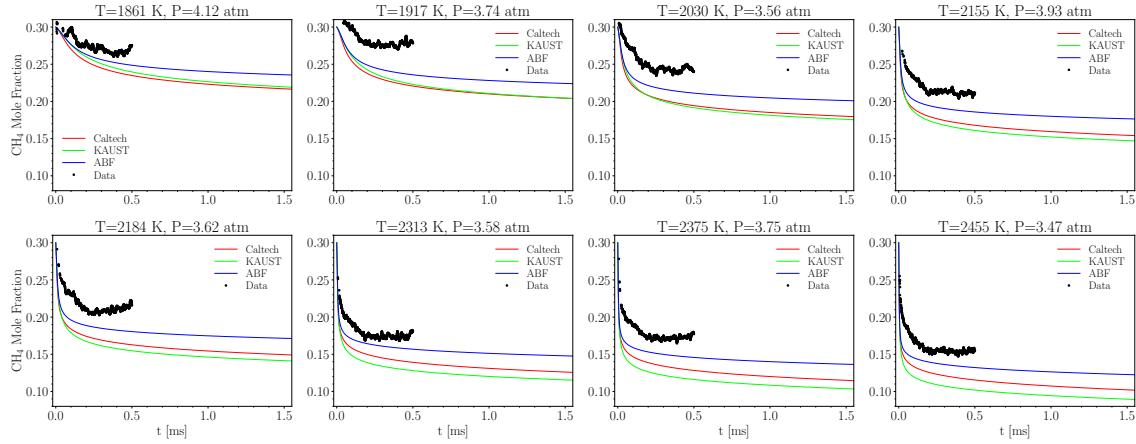


Figure B.8: The time history of CH_4 mole fraction of 30% CH_4 pyrolysis for $1861 \text{ K} < T < 2455 \text{ K}$ using Caltech, KAUST, and ABF mechanisms without soot compared with laser diagnostic data

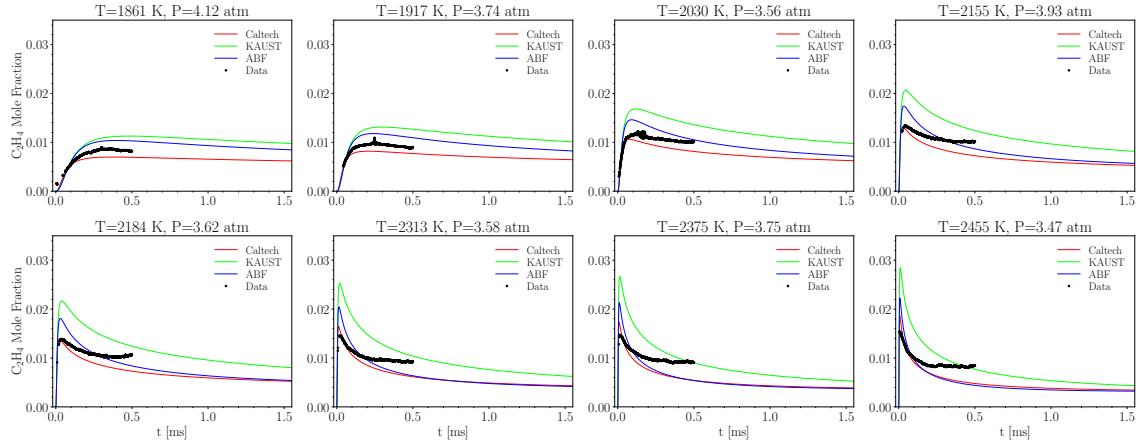


Figure B.9: The time history of C_2H_4 mole fraction of 30% CH_4 pyrolysis for $1861 \text{ K} < T < 2455 \text{ K}$ using Caltech, KAUST, and ABF mechanisms without soot compared with laser diagnostic data

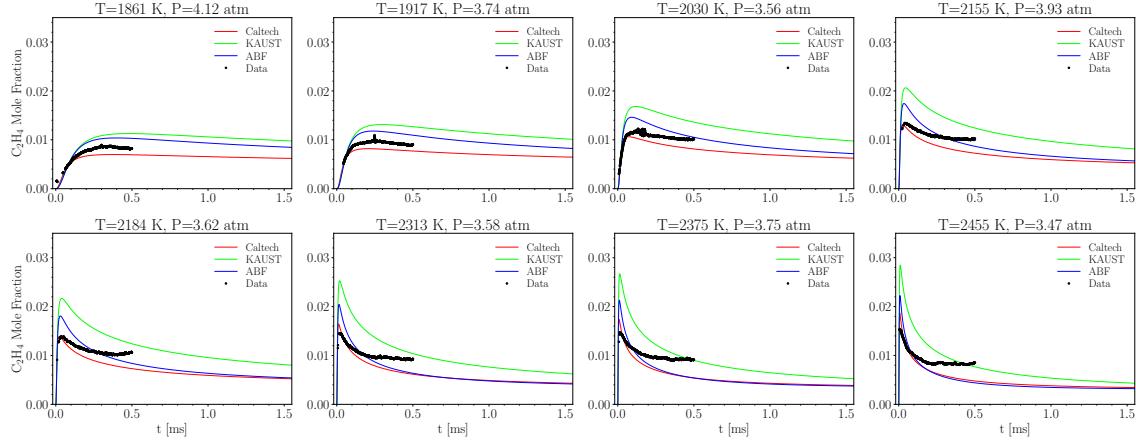


Figure B.10: The time history of C_2H_2 mole fraction of 30% CH_4 pyrolysis for $1861 \text{ K} < T < 2455 \text{ K}$ using Caltech, KAUST, and ABF mechanisms without soot compared with laser diagnostic data

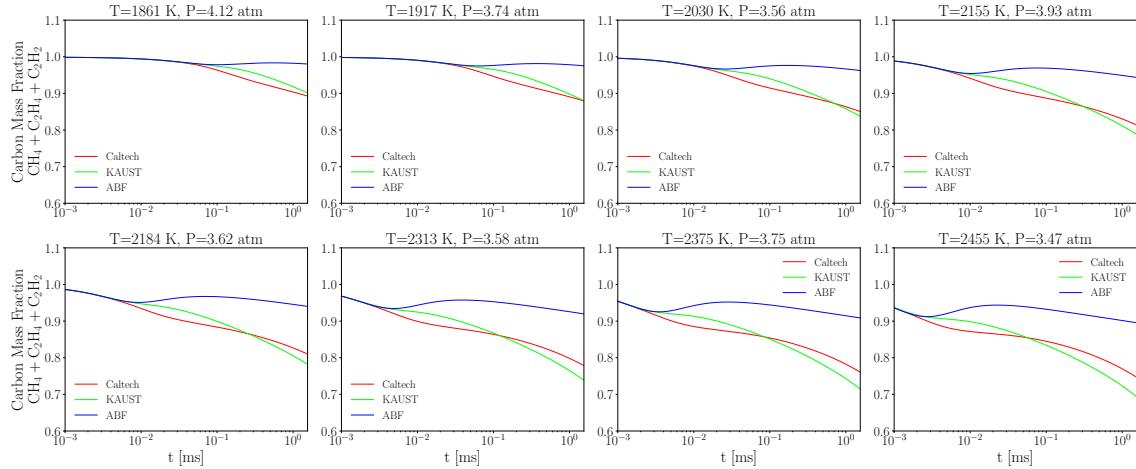


Figure B.11: The time history of carbon mass fraction of CH_4 , C_2H_4 , C_2H_2 of 30% CH_4 pyrolysis for $1861 \text{ K} < T < 2455 \text{ K}$ using Caltech, KAUST, and ABF mechanisms without soot

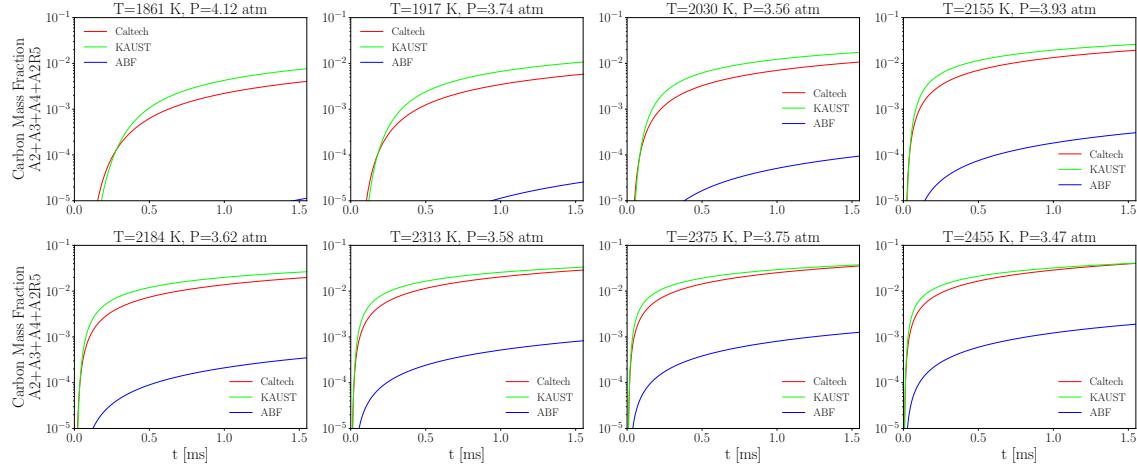


Figure B.12: The time history of carbon mass fraction of soot precursors, A2, A3, A4, and A2R5 of 30% CH_4 pyrolysis for $1861 \text{ K} < T < 2455 \text{ K}$ using Caltech, KAUST, and ABF mechanisms without soot

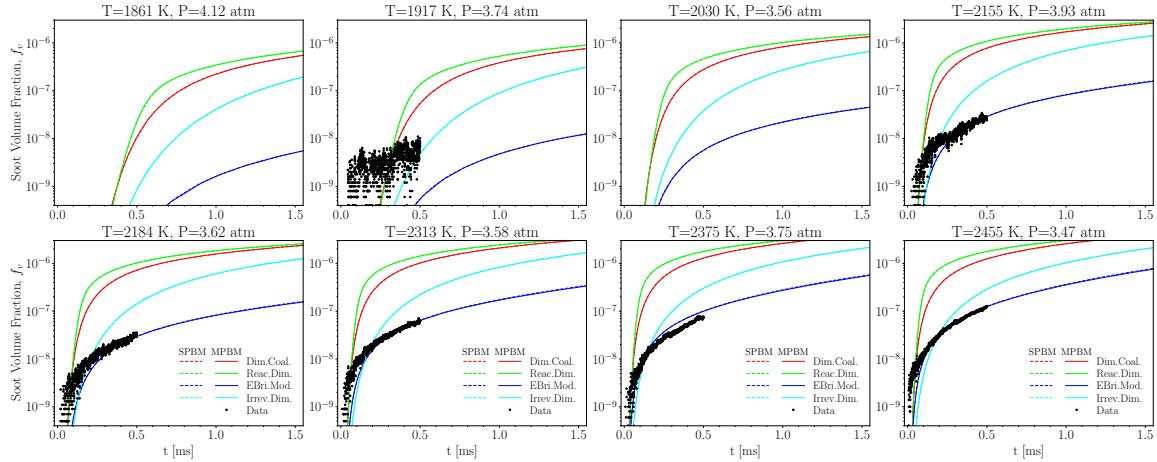


Figure B.13: The time history of soot volume fraction, f_v of 30% CH_4 pyrolysis for $1861 \text{ K} < T < 2455 \text{ K}$ using KAUST with different combinations of PAH growth and particle dynamics models compared with extinction measurements

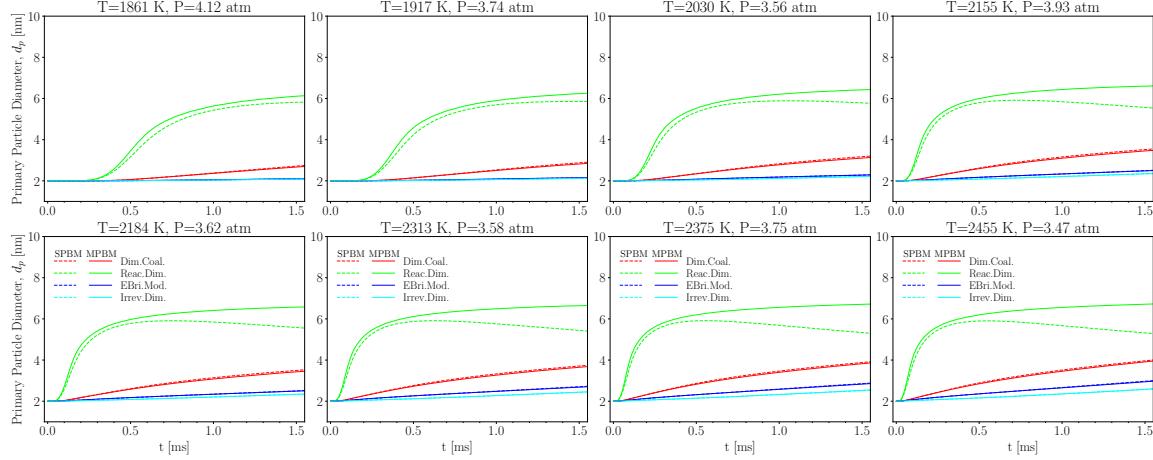


Figure B.14: The time history of primary particle diameter, d_p of 30% CH_4 pyrolysis for $1861 \text{ K} < T < 2455 \text{ K}$ using KAUST with different combinations of PAH growth and particle dynamics models compared with extinction measurements

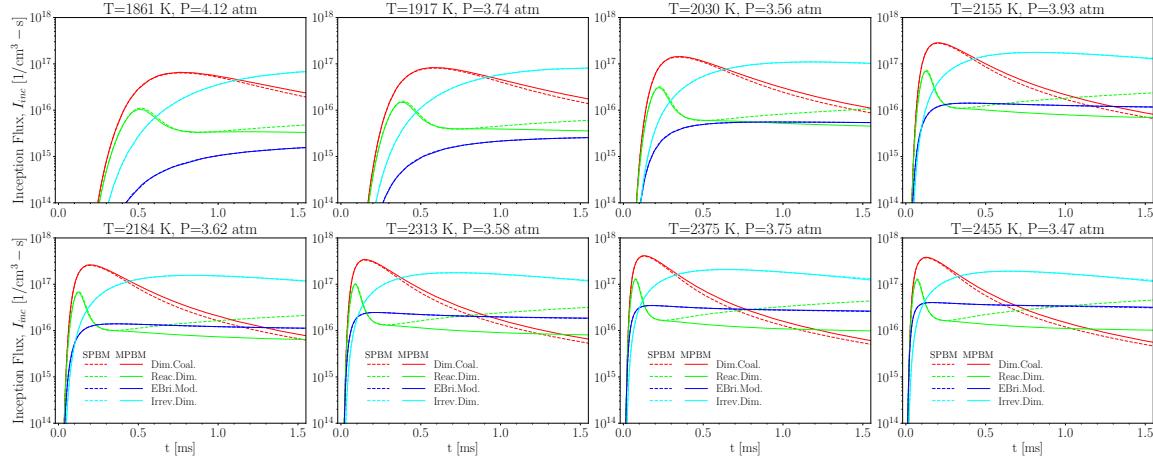


Figure B.15: The time history of inception flux, I_{inc} of 30% CH_4 pyrolysis for $1861 \text{ K} < T < 2455 \text{ K}$ using KAUST with different combinations of PAH growth and particle dynamics models compared with extinction measurements

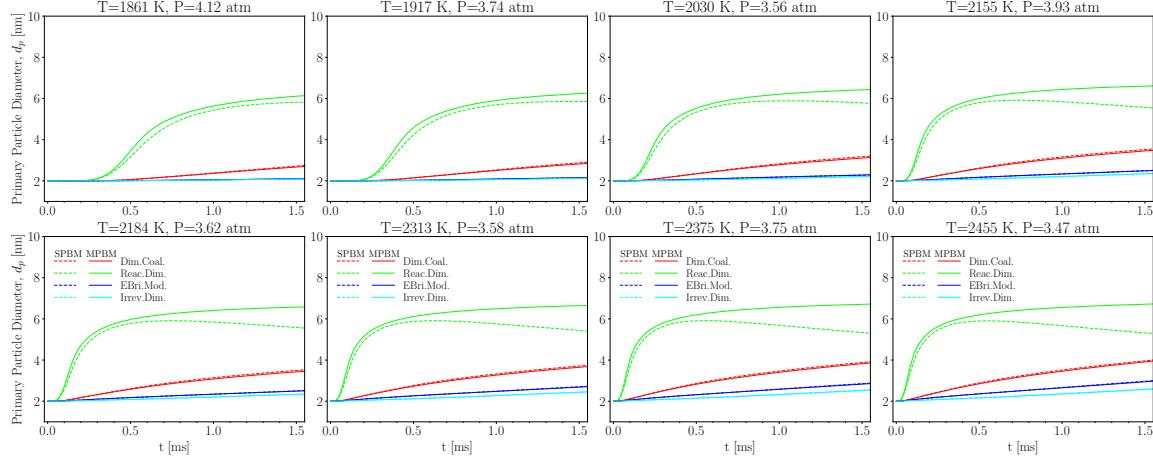


Figure B.16: The time history of mobility diameter, d_m of 30% CH_4 pyrolysis for $1861 \text{ K} < T < 2455 \text{ K}$ using KAUST with different combinations of PAH growth and particle dynamics models compared with extinction measurements

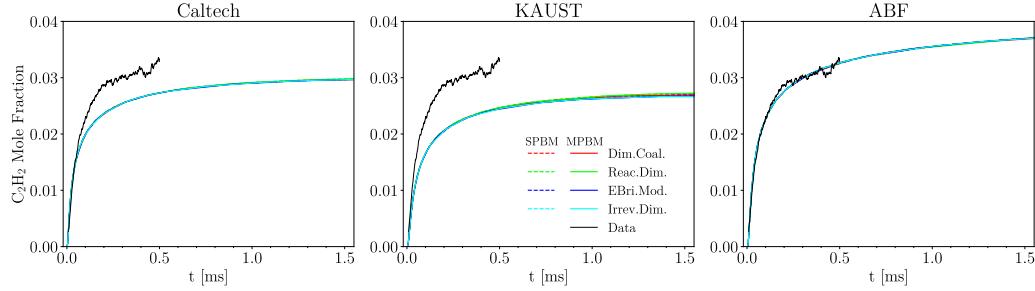


Figure B.17: The time history of C_2H_2 mole fraction of 30% CH_4 pyrolysis at $T=2184 \text{ K}$ and $P=3.62 \text{ atm}$ using Caltech, KAUST, and ABF mechanism and different PAH growth and particle dynamics models compared with laser diagnostic data

B.1.2 10% CH_4 dataset

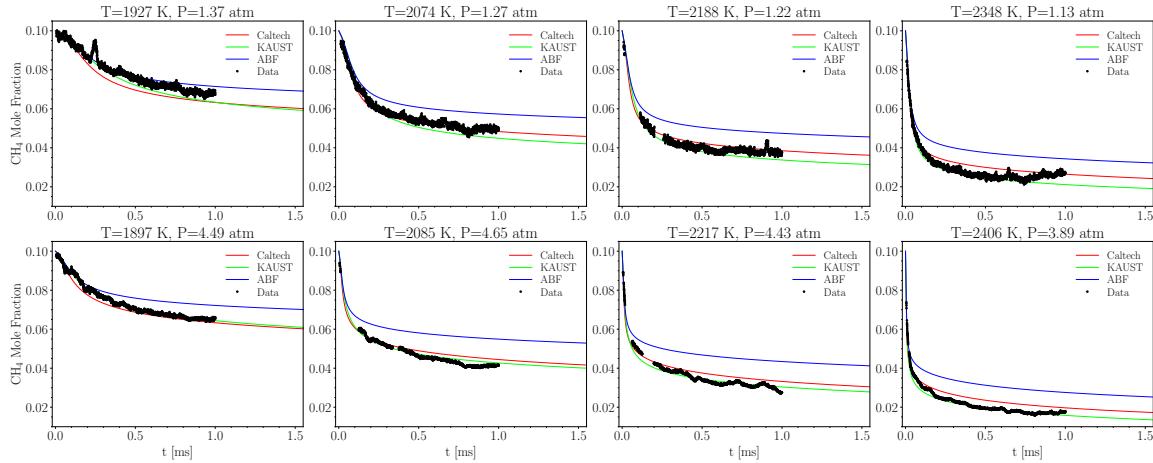


Figure B.18: The time history of CH_4 mole fraction of 10% CH_4 pyrolysis for $1861 \text{ K} < T < 2455 \text{ K}$ using Caltech, KAUST, and ABF mechanisms without soot compared with laser diagnostic data

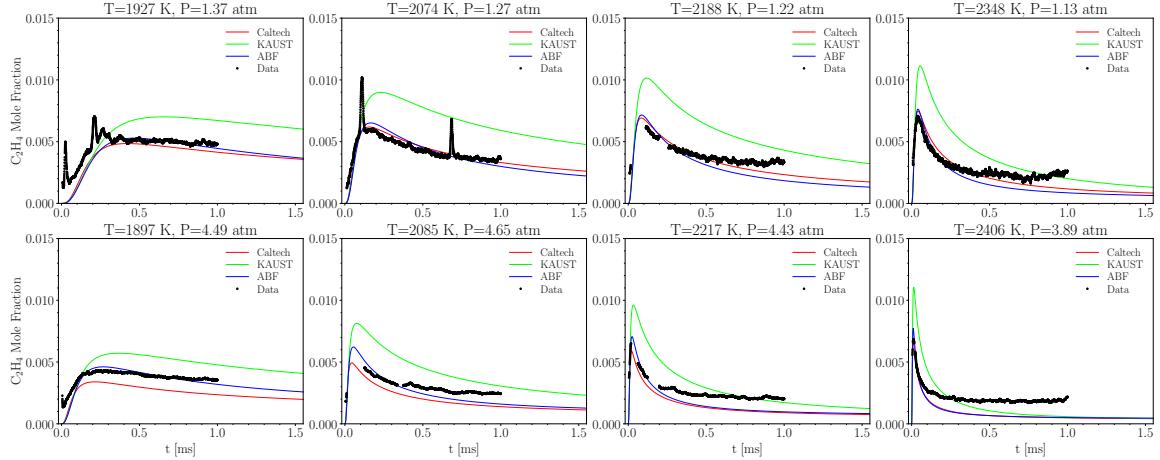


Figure B.19: The time history of C_2H_4 mole fraction of 10% CH_4 pyrolysis for $1861 \text{ K} < T < 2455 \text{ K}$ using Caltech, KAUST, and ABF mechanisms without soot compared with laser diagnostic data

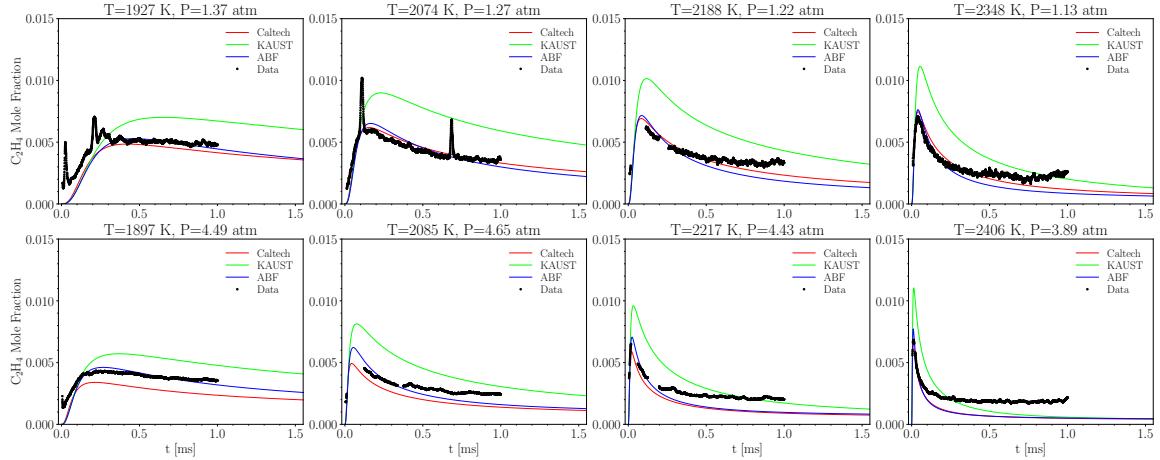


Figure B.20: The time history of C_2H_2 mole fraction of 10% CH_4 pyrolysis for $1861 \text{ K} < T < 2455 \text{ K}$ using Caltech, KAUST, and ABF mechanisms without soot compared with laser diagnostic data

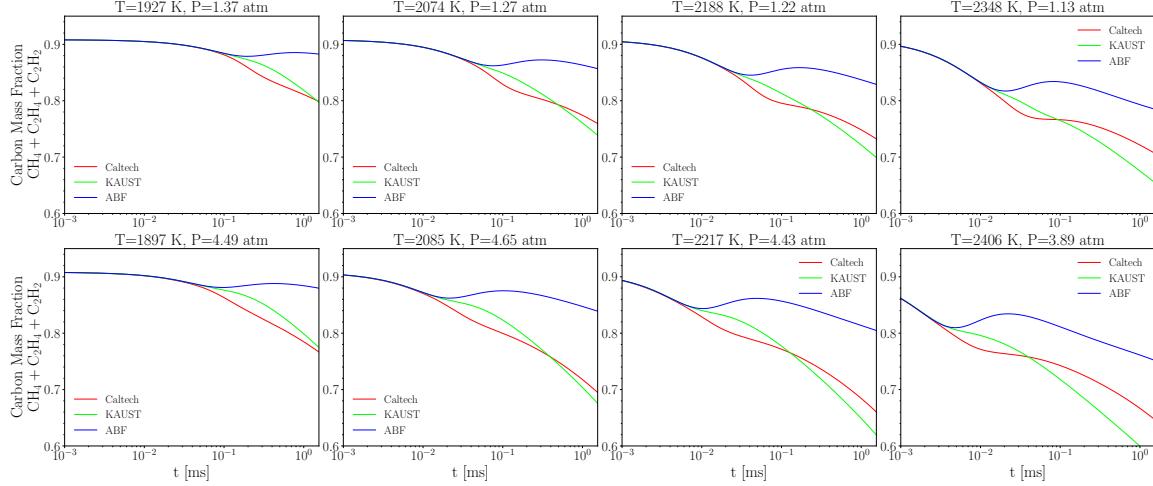


Figure B.21: The time history of carbon mass fraction of CH_4 , C_2H_4 , C_2H_2 of 10% CH_4 pyrolysis for $1861 \text{ K} < T < 2455 \text{ K}$ using Caltech, KAUST, and ABF mechanisms without soot

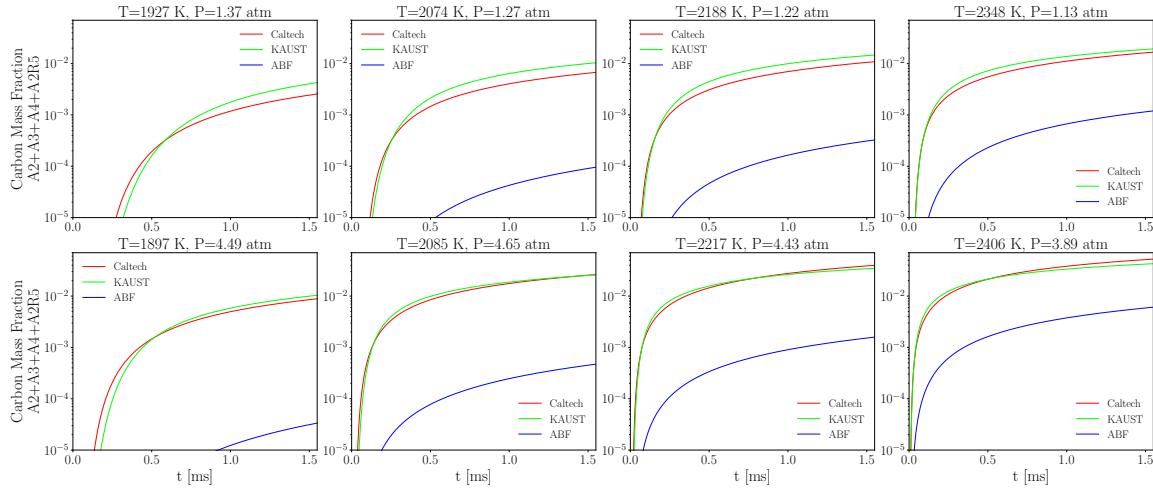


Figure B.22: The time history of carbon mass fraction of soot precursors, $\text{A}2$, $\text{A}3$, $\text{A}4$, and $\text{A}2\text{R}5$ of 10% CH_4 pyrolysis for $1861 \text{ K} < T < 2455 \text{ K}$ using Caltech, KAUST, and ABF mechanisms without soot

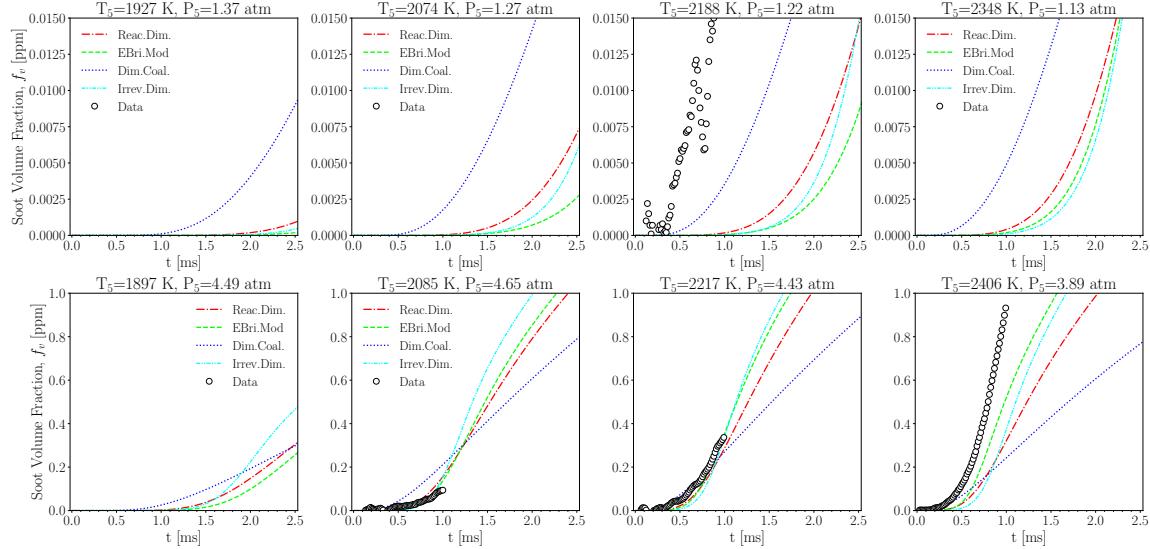


Figure B.23: The soot volume fraction, f_v over $1800 \text{ K} < T < 2400 \text{ K}$ and P near 1 and 4 atm for 10% CH_4 pyrolysis using KAUST, SPBM and optimized PAH growth models that minimizes f_v and d_p at $T=2217$ and $P=4.43$ atm

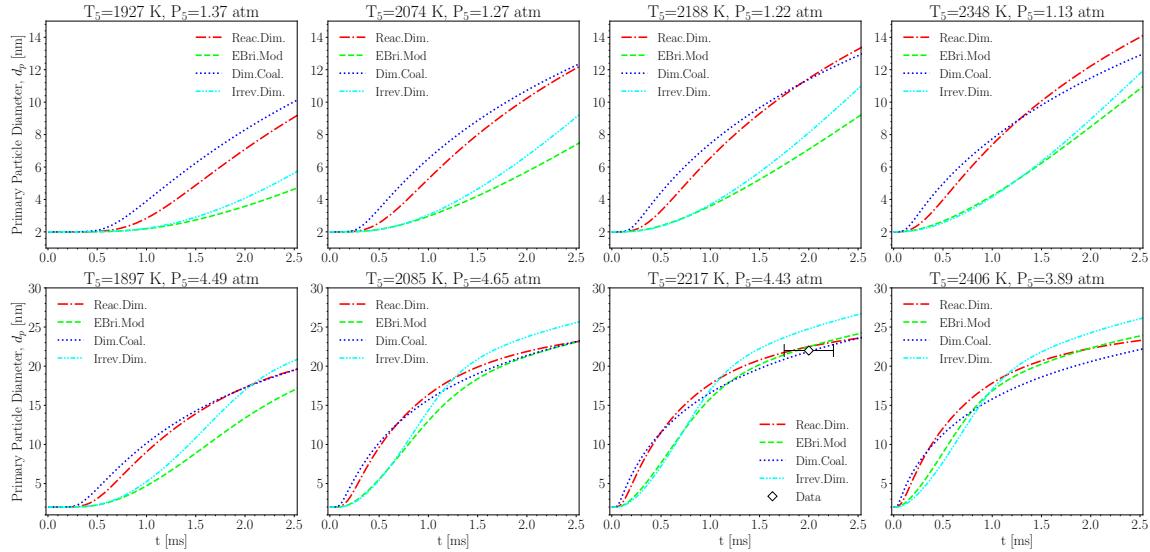


Figure B.24: The primary particle diameter, d_p over $1800 \text{ K} < T < 2400 \text{ K}$ and P near 1 and 4 atm for 10% CH_4 pyrolysis using KAUST, SPBM and optimized PAH growth models that minimizes f_v and d_p at $T=2217$ and $P=4.43$ atm

B.2 Methane pyrolysis data of Agafonov et al. [1]

The pyrolysis of 5% and 10% CH_4 diluted with Ar is investigated for the post-shock temperature range of 1800-3000 K and the pressure range of 4.7-7.1 bar. We assume the pressure linearly increase by temperature across the simulation cases. The obtained soot yields were compared with the soot yield measured by Agafonov et al. [1] using a dual-beam absorption-emission technique. Agafonov et al. [1] reported yield $\times E(m)$ at $\lambda=632 \text{ nm}$, and yield data was retrieved using $E(m)=0.37$ suggested therein.

Figure B.25 depicts soot yield (SY) at $t=1.5$ ms for 5% and 10% CH_4 over the temperature range of 1800-3000 K using MPBM and SPBM and different inception model and Caltech [163] mechanism. Note that, the simulations were also performed using ABF [58] and KAUST [156] mechanisms, and the results were included in Appendix ***. The soot yield predictions exhibit a close agreement with measurements [1] considering the uncertainties in residence time from experiments. The model also successfully captures the bell-shape dependence of soot yield on temperature that has been observed in a variety of hydrocarbons in shock tube [174, 175]. The particle dynamics model has minimal effect on the predicted soot yield. In the vicinity of peak soot yield, SPBM results in slightly lower yield than MPBM, but they are indistinguishable in the rest of the temperature range.

In 5% CH_4 , the SY peak temperature obtained from the model is slightly shifted towards higher temperatures ($2300 \text{ K} < T_{peak} < 2400 \text{ K}$) compared to the measurements $T_{peak} = 2200 \text{ K}$. There are noticeable differences in the behavior of PAH growth model depending on the shock-tube initial temperature. When $T < 2100 \text{ K}$, Reactive Dimerization and EBridge Formation have the highest and lowest soot yield, respectively with Dimer Coalescence predicting soot yields that always falls between Reactive and Irreversible Dimerization. However, The soot yield of EBridge Formation rapidly rises with temperature and exceeds that of Reactive Dimerization and stays higher for the rest of temperature range.

The SY noticeably increases for higher initial CH_4 mole fraction of 10% because more PAH and C_2H_2 are formed leading to stronger carbon conversion rate to soot via inception and surface growth. The peak SY obtained from the model occurs in a higher temperature range 2600-2700 K compared to 5% CH_4 . Soot yield trends can be better understood by examining carbon mass fraction of species directly contributing to soot mass. Figure B.26 depicts the bell-shape distribution of the carbon mass fraction of soot precursors (A2, A2R5, A3, and A4) over the studied temperature range. In low ($T=1800 \text{ K}$) and high ($T=2300 \text{ K}$) end of distribution, a low amount of precursors are formed resulting in low inception rates, particle number concentration and surface growth sites that reduces the soot yield. Additionally, 10% CH_4 has a wider spread with peaks at higher temperature compared to 5% CH_4 which explains the shift in peak yield temperature in Figure B.25. The effect of particle dynamics model is only noticeable in Reactive Dimerization where SPBM results in higher precursor CMF (lower consumption) due to its lower PAH adsorption rates.

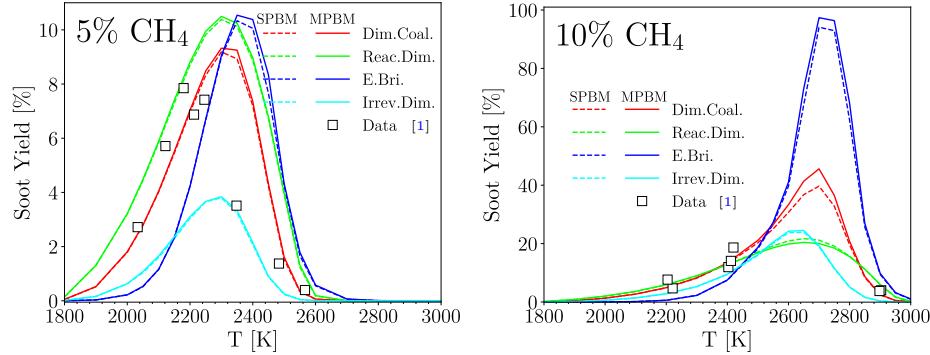


Figure B.25: The temperature dependence of soot yield during pyrolysis of 5% CH_4 -Ar (left pane) and 10% CH_4 -Ar (right pane) at $P = 4.5\text{--}6.7$ bar obtained using Caltech mechanism and different inception models compared with measurements at 1.5ms [1] where the absorption function of $E(m)=0.37$ is used to estimate soot yield.

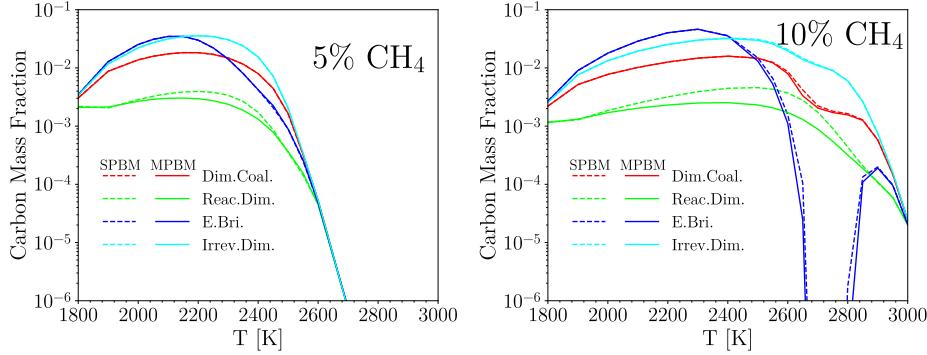


Figure B.26: The temperature dependence of carbon mass fraction of soot precursors (A2, A2R5, A3, and A4) combined during pyrolysis of 5% CH₄-Ar (left pane) and 10% CH₄-Ar (right pane) at P = 4.5–6.7 bar obtained using Caltech mechanism and different inception models at t=1.5 ms

EBridge Formation exhibits a similar behavior in %10 CH₄ simulations where it starts with the lowest SY at T<2500 K and then quickly increases reaching its peak at 100% which is significantly larger than other PAH growth models. The remarkable drop in carbon mass fraction of precursors with EBridge Formation corresponds to %100 yield meaning that all gaseous carbon including the precursors are directed towards soot particles. The higher precursor CMF with Irreversible Dimerization near the peak yield temperature region (2200–2400 K for %5 CH₄ and 2600–2800 K for %10 CH₄) indicates less consumption of precursors via inception and PAH adsorption.

Figure B.27 shows the CMF of C₂H₂ that has an overall increasing trend in the temperature range, but it reaches a plateau for %5 CH₄. There is also a remarkable drop in C₂H₂ CMF in 2600–2800 K due to strong mass growth rate of soot particles that drains C₂H₂ from the gas mixture leading to high soot yield ≈ 100%.

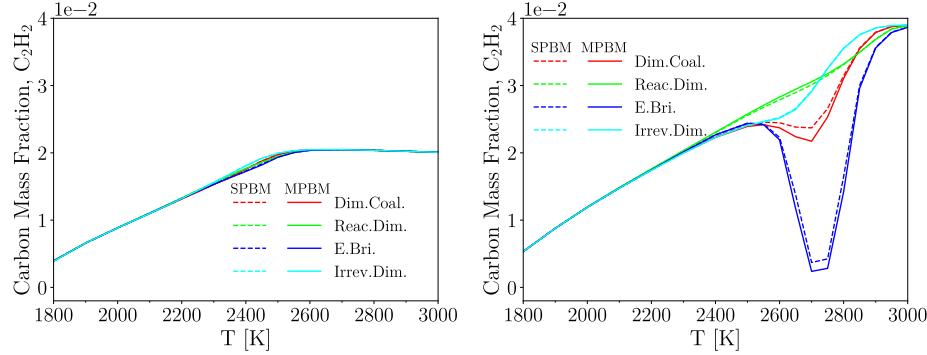


Figure B.27: The temperature dependence of carbon mass fraction of C₂H₂ during pyrolysis of 5% CH₄-Ar (left pane) and 10% CH₄-Ar (right pane) at P = 4.5–6.7 bar obtained using Caltech mechanism and different inception models at t=1.5 ms

Although soot yield, and precursor and acetylene CMF are not sensitive to particle dynamics model, there is a significant difference in agglomerate morphology between SPBM and MPBM predictions. Fig. B.28 shows the average number of primary particles per agglomerate, which is larger for SPBM by the maximum factor of 5 for EBridge Formation and Dimer Coalescence in 10% CH₄. SPBM predicts larger agglomerates due to accounting for polydispersity of particles that results in higher overall collision rate and faster growth by coagulation. n_p follows a bell-shape trend similar to soot yield (Fig. B.25). The n_p values and the difference between the particle dynamics models reaches their maximum in 2200–2400 K for %5 CH₄ and 2600–2800 K for %10 CH₄ because of a stronger inception flux leading to larger number concentration of particles and higher coagulation rate. n_p is larger for Dimer Coalescence and EBridge Formation reaching the maximum of nearly

100 and 1000 in %5 and %10 CH₄, respectively using SPBM. The effect of particle dynamics is minimum for Reactive and Irreversible Dimerization at 5% CH₄ over the whole temperature range.

Fig.B.29 shows the standard geometric deviation of mobility diameter, σ_g obtained by SPBM that reaches the maximum of 5 and 10 for %5 and %10 CH₄, respectively indicating a significant degree of polydispersity in the generated particles at t=1.5ms.

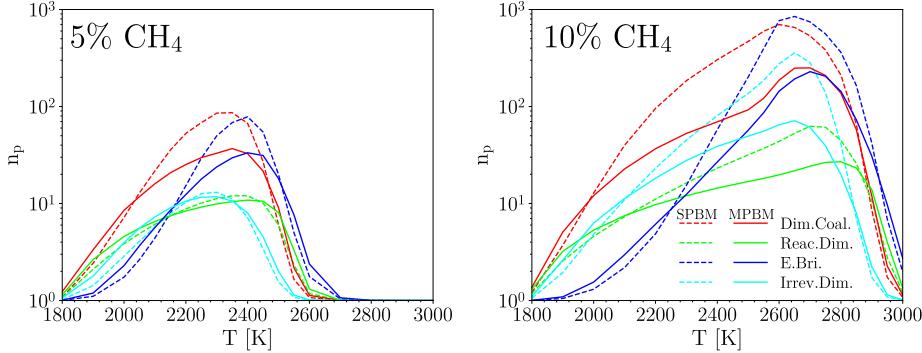


Figure B.28: The temperature dependence of average number of primary particle per agglomerate, n_p during pyrolysis of 5% CH₄-Ar (left pane) and 10% CH₄-Ar (right pane) at P = 4.5–6.7 bar obtained using Caltech mechanism and different inception models at t=1.5 ms

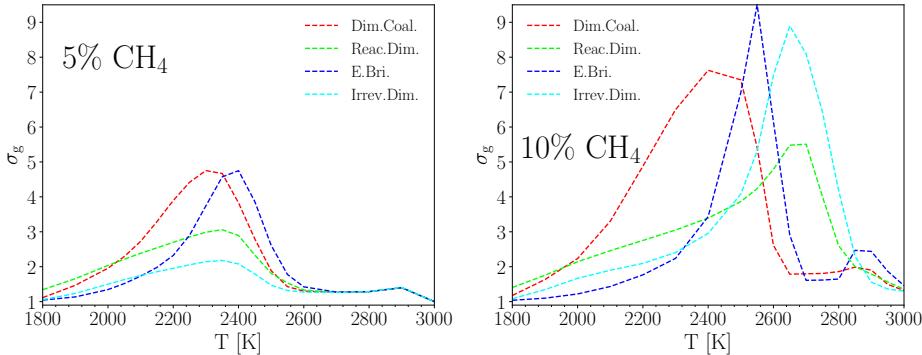


Figure B.29: The time variation of standard geometric deviation of mobility diameter, σ_g during pyrolysis of 5% CH₄-Ar (left pane) and 10% CH₄-Ar (right pane) at P = 4.5–6.7 bar obtained using Caltech mechanism and different inception models at t=1.5 ms

The σ_g values from the SMPS measurements of soot particles at t≈45ms in a benchmark burner-stabilized premixed are close to 1.1, which is significantly lower than values observed here. So, the evolution of σ_g in the studied shock tube is examined in an extended time frame up to 4 msf for 10% CH₄ at T=2500 and 2700 K. As shown in Fig.B.30, for all PAH growth models σ_g rises in the beginning due to simultaneous inception and coagulation that increases polydispersity and rapidly drops and approaches 1.5 before t=3ms when coagulation becomes dominant and inception weaken due to consumption of precursors by PAH adsorption and HACA.

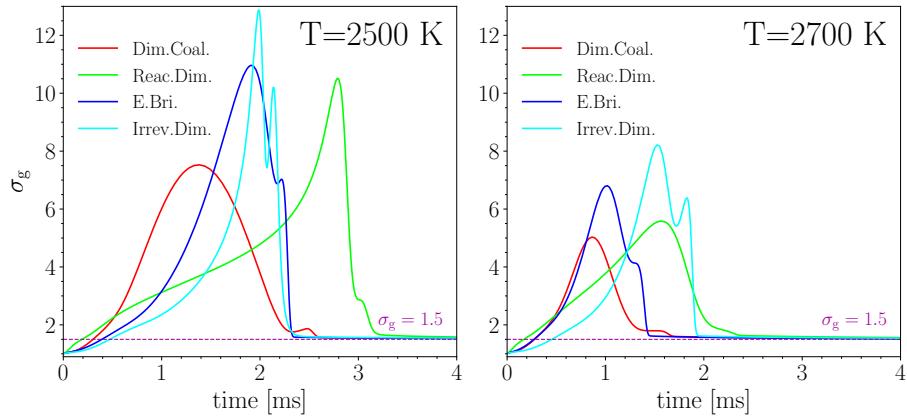


Figure B.30: The temperature dependence of standard geometric deviation of mobility diameter, σ_g during pyrolysis of 10% CH₄-Ar at T=2500 K (left pane) and T=2700 K (right pane)