Heat Transfer: Convection

Forced Convection in Internal Flows – HTC in laminar fully developed flows

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Learning goals

Forced convection in internal flows:

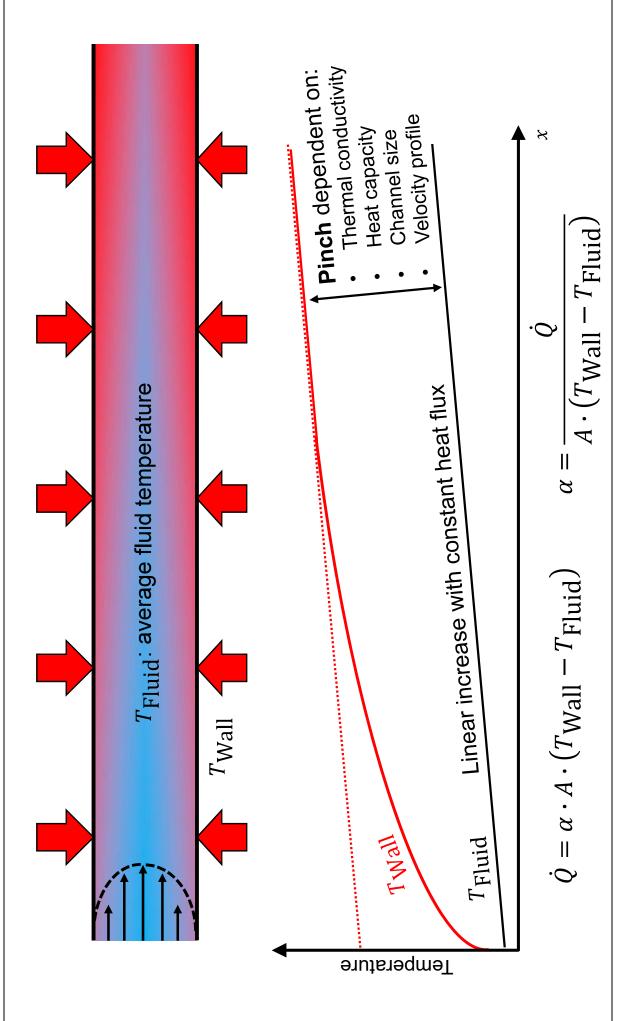
- Ability to calculate the heat transfer coefficient in laminar flows under fully developed conditions
- Ability to distinguish between different flow configurations and to choose the proper correlation for the HTC







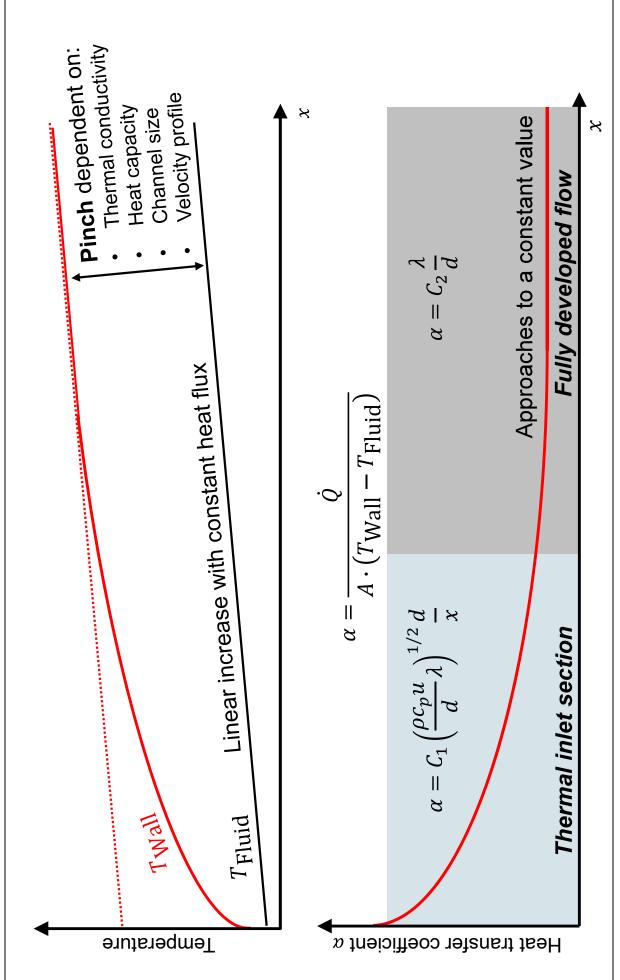










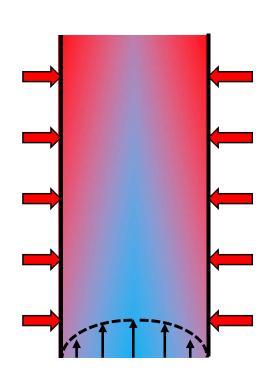








Heat Transfer: HTC in laminar fully developed flows



$$\alpha = \frac{Q}{A \cdot (T_{\text{Wall}} - T_{\text{Fluid}})}$$

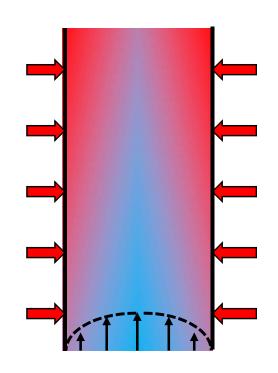
- ► Heat flux is constant
- $ightharpoonup T_{
 m Fluid}$ need to be determined

Procedure:

- Determine the radial temperature profile in the pipe flow
 - Calculate the caloric mean fluid temperature $T_{
 m Fluid}$







- ► Heat flux is constant
- → Caloric mean temperature increases linearly in axial direction
- → Temperature increases linearly with the same slope at any radial location
- → Temperature gradient in r-direction needs to be constant

$$ightarrow rac{\partial T}{\partial x} = rac{\dot{q}''\pi D}{\dot{m}c_p}$$
 at any location r with $\dot{m} = \rho U\pi R^2$

U = mean velocity

Equation:

(steady state, advection in x-direction, conduction in r-direction): Governing energy conservation equation

$$u\rho c_{p} \frac{\partial T}{\partial x} = \frac{\lambda}{r} \frac{\partial}{\partial r} \left(r \frac{\partial T}{\partial r} \right)$$





1) Temperature profile in a fully-developed pipe with constant heat flux

Equation:

$$u\rho c_p \frac{\partial T}{\partial x} = \frac{\lambda}{r} \frac{\partial}{\partial r} \left(r \frac{\partial T}{\partial r} \right)$$

$$\frac{\partial T}{\partial z} = \frac{\dot{q}'' 2\pi}{\dot{q}'' \dot{q}}$$

with
$$\frac{\partial T}{\partial x} = \frac{\dot{q}'' 2\pi}{\rho U \pi R c_p}$$

And the velocity profile in a laminar pipe flow: $u=2U\left(1-rac{r^2}{R^2}
ight)$

Mathematical steps (1):

$$2U\left(1 - \frac{r^2}{R^2}\right)\frac{\dot{q}''\rho c_p\pi^2}{\rho U\pi R c_p} = \frac{\lambda}{r}\frac{\partial}{\partial r}\left(r\frac{\partial T}{\partial r}\right)$$

U, c_p , π and ρ cancel out

$$4\left(1 - \frac{r^2}{R^2}\right)\frac{\dot{q}''}{R\lambda} = \frac{1}{r}\frac{\partial}{\partial r}\left(r\frac{\partial T}{\partial r}\right)$$

$$4\left(r - \frac{r^3}{R^2}\right)\frac{\dot{q}''}{\lambda R} = \frac{\partial}{\partial r}\left(r\frac{\partial T}{\partial r}\right)$$

First integration:

$$\left(r\frac{\partial T}{\partial r}\right) = 4\left(\frac{1}{2}r^2 - \frac{r^4}{4R^2} + C_1\right)\frac{\dot{q}''}{\lambda R}$$







1) Temperature profile in a fully-developed pipe with constant heat flux

Radial temperature gradient:

$$\left(r\frac{\partial T}{\partial r}\right) = 4\left(\frac{1}{2R}r^2 - \frac{r^4}{4R^3} + C_1\right)\frac{\dot{q}''}{\lambda}$$

Boundary condition (symmetry):
$$\frac{\partial T}{\partial r}(r=0)=0 \rightarrow C_1=0$$

$$\frac{\partial T}{\partial r} = 4\left(\frac{1}{2R}r^1 - \frac{r^3}{4R^3}\right)\frac{\dot{q}''}{\lambda} \qquad \Longrightarrow \qquad \frac{\partial T}{\partial r} = \left(\frac{2}{R}r^1 - \frac{r^3}{R^3}\right)\frac{\dot{q}''}{\lambda}$$

$$\frac{\partial T}{\partial r} = \left(\frac{2}{R}r^1 - \frac{r^3}{R^3}\right)\frac{\dot{q}''}{\lambda}$$

Radial temperature distribution:

Second integration:
$$T(r) = \frac{\dot{q}''}{\lambda} \left(\frac{1}{R} r^{\lambda} \right)$$

$$T(r) = \frac{\dot{q}''}{\lambda} \left(\frac{1}{R} r^2 - \frac{r^4}{4R^3} + C_2 \right)$$

Boundary condition (wall temperature):
$$T(r = R) = T_{\rm w}$$





1) Temperature profile in a fully-developed pipe with constant heat flux

Radial temperature distribution:

$$T(r) = \frac{\dot{q}''}{\lambda} \left(\frac{1}{R} r^2 - \frac{r^4}{4R^3} + C_2 \right)$$

$$C_2 = T_W \frac{\lambda c_p}{\dot{q}''} - \frac{3}{4}R$$

$$T(r) = T_W + \frac{\dot{q}''}{\lambda} \left(\frac{1}{R} r^2 - \frac{r^4}{4R^3} - \frac{3}{4} R \right)$$

$$C_2 = T_W \frac{ncp}{q''} - \frac{2}{4}$$

$$T(r) = T_W - \frac{\dot{q}''R}{\lambda} \left(\frac{3}{4} - \frac{1}{R^2} r^2 + \frac{r^4}{4R^4} \right)$$

2) Caloric mean temperature in the pipe flow:

► Determine mean temperature of the fluid (mass flow averaged temperature):

$$T_{\rm m} = \frac{\int u(r)A(r)T(r)\rho c_p}{\int u(r)A(r)\rho c_p}$$

$$A(r) = 2\pi r dr$$







2) Mass flow averaged temperature in the pipe flow

Mathematical calculations:

$$T_{\rm m} = \frac{\int \left({\rm T_w} - \frac{{\dot q'''R}}{\lambda} \left(\frac{3}{4} - \frac{1}{R^2} \, r^2 + \frac{r^4}{4R^4} \right) \right) \ 2 U \left(1 - \frac{r^2}{R^2} \right) \, 2 \pi r dr}{\int \ 2 U \left(1 - \frac{r^2}{R^2} \right) \, 2 \pi r dr}$$

$$T_{\rm m} = \frac{\int \left({\rm T_w} - \frac{{\dot q''}R}{\lambda} \left(\frac{3}{4} - \frac{1}{R^2} r^2 + \frac{r^4}{4R^4} \right) \right) \left(r - \frac{r^3}{R^2} \right) dr}{\int \left(r - \frac{r^3}{R^2} \right) dr} = \frac{\int \left(r - \frac{r^3}{R^2} \right) dr}{\int \left(r - \frac{r^3}{R^2} \right) dr} - \frac{{\dot q''}R}{\lambda} \frac{\int_0^R \left(\frac{3}{4} - \frac{1}{R^2} r^2 + \frac{r^4}{4R^4} \right) \left(r - \frac{r^3}{R^2} \right) dr}{\int \left(r - \frac{r^3}{R^2} \right) dr}$$

Integration boundaries: r = 0 and r = R

$$T_{\rm m} = T_{\rm w} - \frac{11}{24} \frac{\dot{q}''}{\lambda}$$

$$_{\rm n} = T_{\rm w} - \frac{11}{24} \frac{\dot{q}'' \, k}{\lambda}$$







Heat transfer coefficient and Nusselt number

Final steps:

$$T_{\rm m} = T_{\rm w} - \frac{11}{24} \frac{\dot{q}''R}{\lambda}$$

$$\dot{q}'' = \frac{24 \, \lambda}{11 \, R} (T_{\rm w} - T_{\rm m})$$

$$\alpha = \frac{24 \ \lambda}{111 R} = 2.18 \frac{\lambda}{R}$$

$$Nu = \frac{\alpha D}{\lambda} = 4.36$$

The mean temperature is directly coupled to the wall temperature through the heat flux

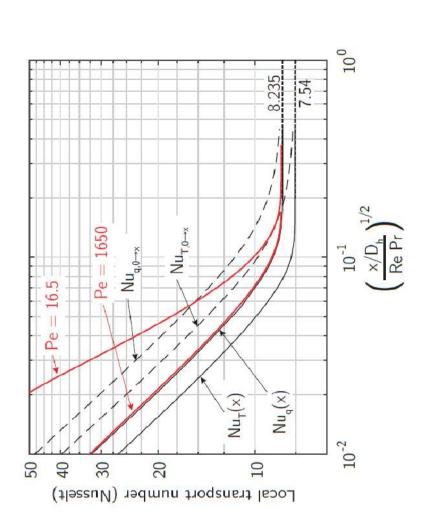
Rearrangement yields for the local heat flux

The fully developed velocity and temperature profile gives a constant Nusselt number!





Dimensionless heat transfer laws for a flow between two parallel plates



Dimensionless heat transfer laws:

$$Nu = \frac{\alpha d}{\lambda}$$

$$Re = \frac{\rho ud}{\eta}$$

$$Pr = \frac{\eta c_p}{\lambda}$$

$$Pe = Re \cdot Pr$$

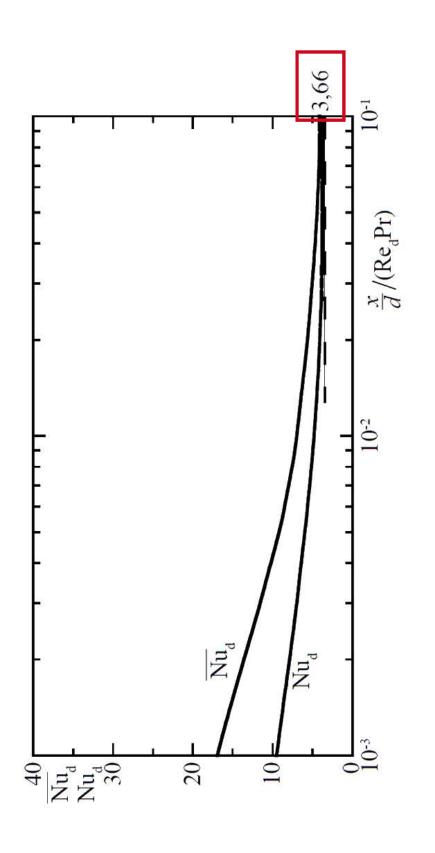
Necessary assumptions:

- Constant fluid properties
- → no shear thinning
- → no temperature dependencies





Laminar pipe flow with hydrodynamically developed flow at the beginning of the heated/cooled pipe section and isothermal surface



$$\overline{\mathrm{Nu_d}} = \left(3.66 + \frac{0.19 \left(\mathrm{Re_d Pr} \frac{d}{L} \right)^{0.8}}{1 + 0.117 \left(\mathrm{Re_d Pr} \frac{d}{L} \right)^{0.467}} \right) \left(\frac{\eta}{\eta_{\mathrm{w}}} \right)^{0.14}$$
 (HTC. 12)







Laminar pipe flow with hydrodynamically developed flow at the beginning of the heated/cooled pipe section and constant heat flux

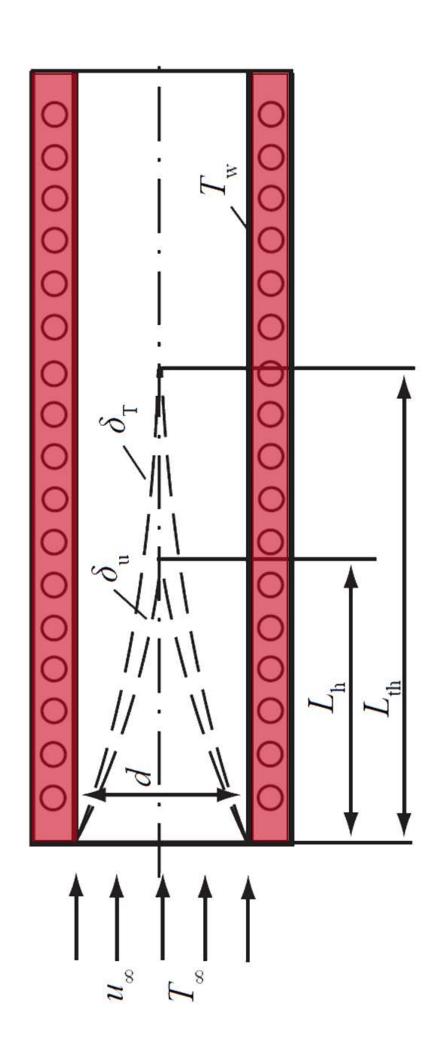
- If the heat flow density at the wall is kept constant instead of the wall temperature, the heat transfer coefficients are about 20% higher
- ▶ The final value for long pipes in this case is $Nu_{d,\infty} = 4.36 \left(\frac{\eta}{\eta_W}\right)^{0.14}$







Laminar pipe flow with simutaneous hydrodynamic and thermal inlet and isothermal surface

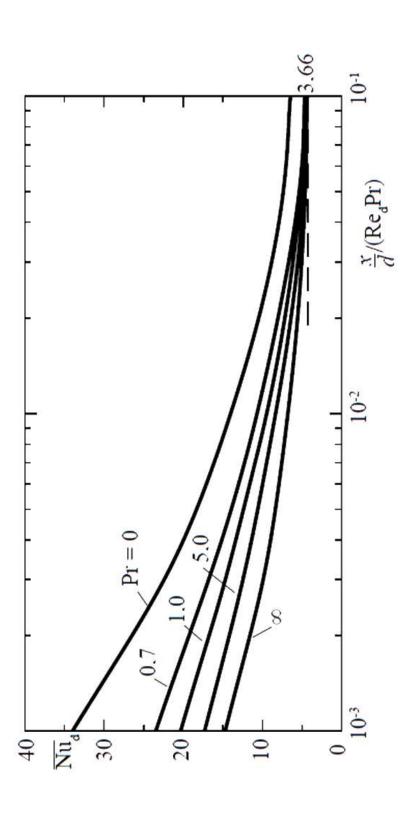








<u>Laminar pipe flow with simutaneous hydrodynamic and thermal inlet and isothermal</u> surface



$$\overline{\text{Nu}_{d}} = \left(\frac{0.0677 \left(\text{Re}_{d} \text{Pr} \frac{d}{L}\right)^{1.33}}{1 + 0.1 \left(\text{Re}_{d} \text{Pr} \frac{d}{L}\right)^{0.83}}\right) \left(\frac{\eta}{\eta_{W}}\right)^{0.14}$$

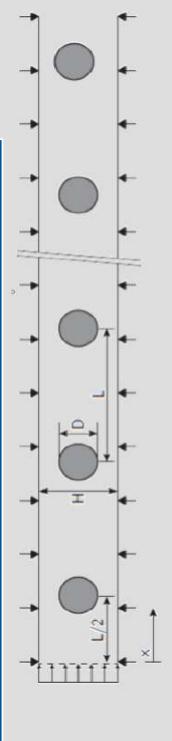




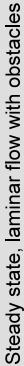


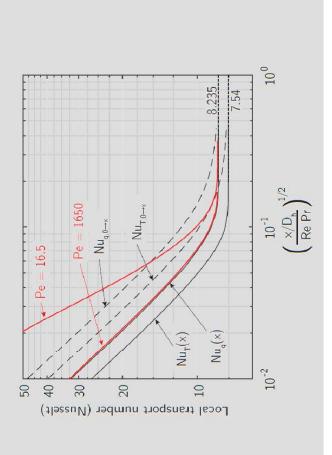
Influence of complex obstacles on the heat transfer of laminar internal flow

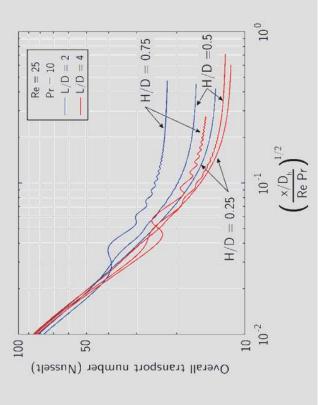
Simulation of a long internal flow - detection of inlet length effects:



Steady state, laminar conditions without obstacles





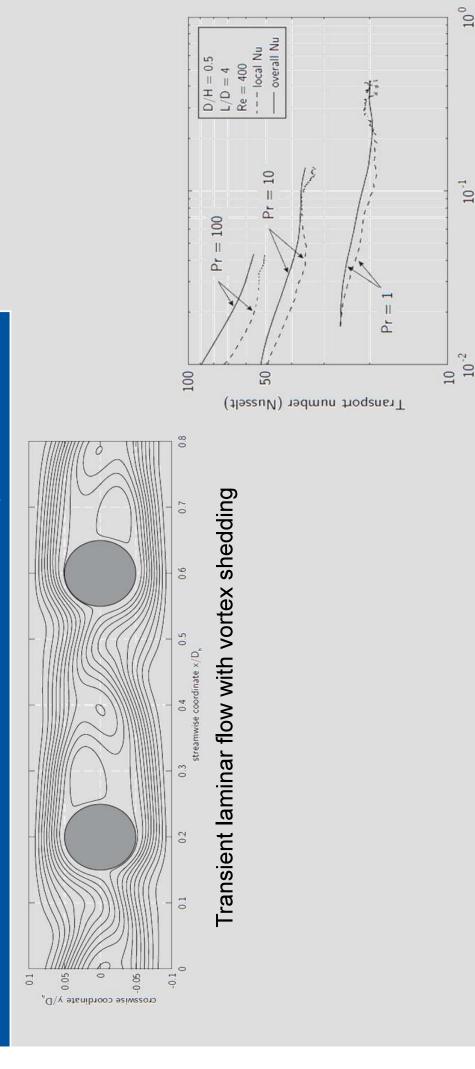








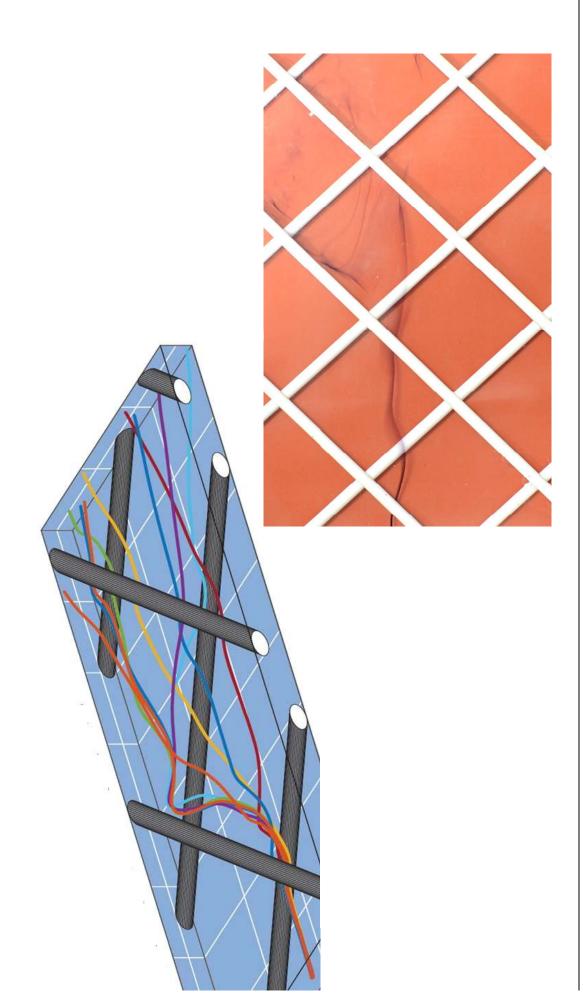
Transient flow behavior and the loss of self-similarity:











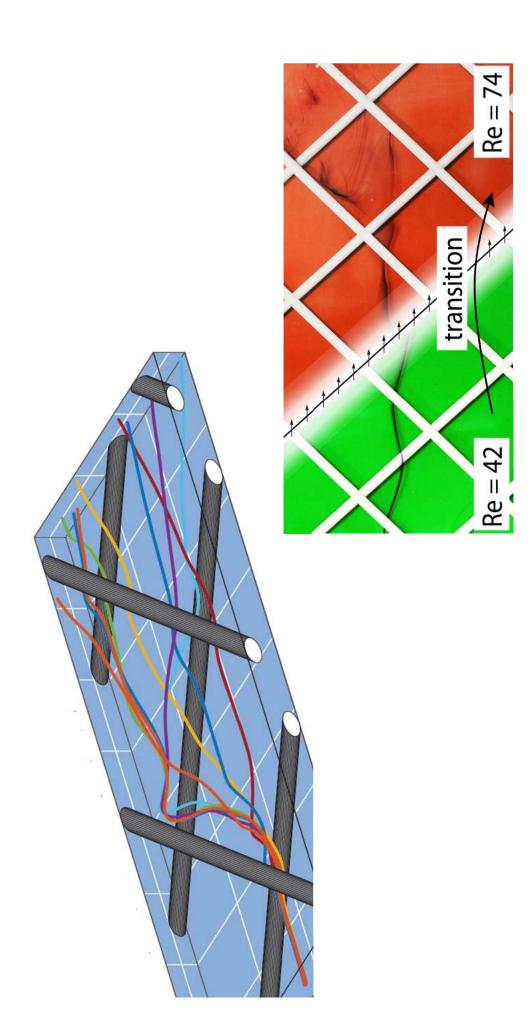








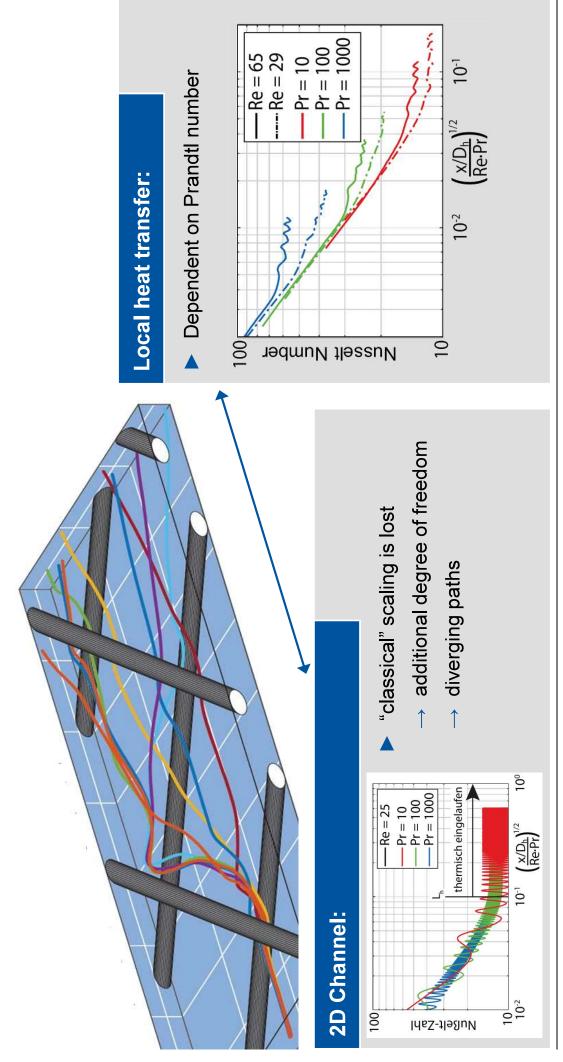
Heat Transfer: HTC in laminar fully developed flows















Turbulent pipe flow at isothermal wall temperature

$$\overline{\mathrm{Nu_d}} = 0.0235 \left(\mathrm{Re}_d^{0.8} - 230\right) (1.8 \mathrm{Pr}^{0.3} - 0.8) \left(1 + \left(\frac{d}{L}\right)^{\frac{2}{3}}\right) \left(\frac{\eta}{\eta_\mathrm{w}}\right)^{0.14} \tag{HTC. 14}$$

For
$$0.6 < Pr < 500$$
 and $Re_d > 2300$

▶ In many cases, instead of the equation (HTC.14), the simpler relationship is sufficient:

$$\overline{\mathrm{Nu_d}} = 0.027 \mathrm{Re_d^{0.8} Pr^{\frac{1}{3}}} \left(\frac{\eta}{\eta_\mathrm{w}}\right)^{0.14}$$
 (HTC. 15)

For
$$\text{Re}_d > 10^5$$
; $0.7 \le \text{Pr} \le 16,700$; $\frac{L}{d} > 10$



Heat Transfer: HTC in laminar fully developed flows

Comprehension questions

Why is the HTC constant in the fully developed region of an internal flow?

What are the major steps to calculate the HTC in the fully developed region?

What can result in a loss of self-similarity of the heat transfer behavior?

Extra tasks:

Proof that the Nusselt number for a laminar flow between two parallel plates with a constant heat flux boundary condition is Nu = 8.235.

Think about another geometry/flow configuration for which you can determine a laminar velocity profile analytically and calculate the Nusselt number.



