

Excitation of TAE modes by an electromagnetic antenna using the global gyrokinetic code ORB5

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ORB5 eqs. of motion

Consider evolution of reduced velocity distribution function $f(\mathbf{R}, \mathbf{v}_{||}, \mu; t)$

$$\frac{\partial \delta f_{s}}{\partial t} + \dot{\mathbf{R}} \cdot \frac{\partial \delta f_{s}}{\partial \mathbf{R}} \Big|_{\mathbf{v}_{||}} + \dot{\mathbf{v}}_{||} \frac{\partial \delta f_{s}}{\partial \mathbf{v}_{||}} = -\dot{\mathbf{R}}^{(1)} \cdot \frac{\partial F_{0s}}{\partial \mathbf{R}} \Big|_{\epsilon} - \dot{\epsilon}^{(1)} \frac{\partial F_{0s}}{\partial \epsilon} . \tag{1.1}$$

The 0th/1st-order gyrocenter characteristics follow

$$\dot{\mathbf{R}}^{(0)} = v_{||} \mathbf{b}^* + \frac{1}{qB_{||}^*} \mathbf{b} \times \mu \nabla B,
\dot{v}_{||}^{(0)} = -\frac{\mu}{m} \mathbf{b}^* \cdot \nabla B,
\dot{\mathbf{R}}^{(1)} = \frac{\mathbf{b}}{B_{||}^*} \times \nabla \langle \phi - v_{||} A_{||}^{(s)} - v_{||} A_{||}^{(h)} \rangle - \frac{q_s}{m_s} \langle A_{||}^{(h)} \rangle \mathbf{b}^*,
\dot{v}_{||}^{(1)} = -\frac{q_s}{m_s} \left[\mathbf{b}^* \cdot \nabla \langle \phi - v_{||} A_{||}^{(h)} \rangle + \frac{\partial \langle A_{||}^{(s)} \rangle}{\partial t} \right] - \mu \frac{\mathbf{b} \times \nabla B}{B_{||}^*} \cdot \nabla \langle A_{||}^{(s)} \rangle,$$

and particle energy $\dot{\epsilon}^{(1)} = v_{||}\dot{v}_{||}^{(1)} + \mu \dot{R}^{(1)} \cdot \nabla B$,

where $\mu = v_{\perp}^2/(2B)$, $B_{||}^* = \boldsymbol{b} \cdot \nabla \boldsymbol{A}^*$, $\boldsymbol{b}^* = \nabla \times \boldsymbol{A}^*/B_{||}^*$, and $\boldsymbol{A}^* = \boldsymbol{A} + \frac{m_s v_{||}}{q_s} \boldsymbol{b}$.

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 ϕ is computed from the quasineutrality equation

$$-\nabla \cdot \left[\left(\sum_{s=i,f} \frac{q_s^2 n_s}{T_s} \rho_s^2 \right) \nabla_{\perp} \phi \right] = \sum_{s=i,e,f} q_s n_{1,s}$$
 (1.2)

where $n_{1,s}=\int \delta f_s \delta({m R}+{m
ho}-{m x}) d^6 Z$, $\rho_s=\sqrt{m_s T}/(q_s B)$, and $d^6 Z=B_{||}^* d{m R} \ dv_{||} \ d\mu \ d\alpha$.

Given the decomposition $A_{||} = A_{||}^{(s)} + A_{||}^{(h)}$, symplectic part is computed from Ohm's law

$$\frac{\partial}{\partial t} A_{||}^{(s)} + \boldsymbol{b} \cdot \nabla \phi = 0, \tag{1.3}$$

and the Hamiltonian part from Ampere's law

$$\left(\sum_{s=i,e,f} \frac{\beta_s}{\rho_s^2} - \nabla_{\perp}^2\right) A_{||}^{(h)} = \mu_0 \sum_{s=i,e,f} j_{||,1s} + \nabla_{\perp}^2 A_{||}^{(s)}, \qquad (1.4)$$

and
$$j_{||,1s} = \int v_{||} \delta f_s \delta(\mathbf{R} + \boldsymbol{\rho} - \mathbf{x}) d^6 Z.$$
 (1.5)

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Pullback scheme

Initialize markers in the phase space;

while $t < t_{\rm final}$ do

for k = 1, ..., 4 step of Runge-Kutta scheme **do**

- Compute ϕ , $A_{||}^{(s)}$ and $A_{||}^{(h)}$;
- Push particles according to equations of motion;
- Apply boundary conditions;

end

- Update mixed-variable δf with $\delta f_s^{(m)} \leftarrow \delta f_s^{(m)} + \frac{q_s \langle A_{||}^{(h)} \rangle}{m_s} \frac{\partial F_{0s}}{\partial v_{||}}$;
- Update $A_{||}$ decomposition, i.e, $A_{||}^{(s)} \leftarrow A_{||}^{(s)} + A_{||}^{(h)}$ and $A_{||}^{(h)} \leftarrow 0$;
- $t = t + \Delta t$;

end

Algorithm 1: The δf solution algorithm used in ORB5 within the pullback scheme framework in linear setting

We devise antenna as an electrostatic potential with profile

$$\phi_{\text{ant}}(s,\theta,\varphi;t) = \text{Re}\left[\sum_{I\in\mathcal{T}} h_I(s) A_I e^{\hat{i}(I_1\theta + I_2\varphi + \Phi_I)} (c_1 + c_2 e^{\hat{i}w_{\text{ant}}t})\right]$$
(1.6)

where the set $\mathcal{T} = \{(m_1, n_1), ...\}$ includes the target mode numbers. Electromagnetic potential of antenna from Ohm's law is computed

$$\frac{\partial}{\partial t} A_{\text{ant},||}^{(s)} + \boldsymbol{b} \cdot \nabla \phi_{\text{ant}} = 0.$$
 (1.7)

Eqs. of motion are adapted by setting $\phi_{\rm plasma} \to \phi_{\rm plasma} + \phi_{\rm ant}$.

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Integrating electrostatic antenna in the eqs. of motion

$$\dot{\mathbf{R}}^{(1)} = \dot{\mathbf{R}}_{\text{plasma}}^{(1)} + \frac{\mathbf{b}}{B_{\parallel}^*} \times \nabla \langle \phi_{\text{ant}} \rangle, \tag{1.8}$$

$$\dot{\mathbf{v}}_{||}^{(1)} = \dot{\mathbf{v}}_{||,\mathrm{plasma}}^{(1)} - \frac{q_s}{m_s} \mathbf{b}^* \cdot \nabla \langle \phi_{\mathsf{ant}} \rangle, \tag{1.9}$$

$$\dot{\epsilon}^{(1)} = \dot{\epsilon}_{\text{plasma}}^{(1)} - \frac{q_s}{m_s} \left[\frac{v_{||} \mathbf{B}}{B_{||}^*} + m_s \mu \frac{\mathbf{b} \times \nabla B}{q_s B_{||}^*} + \frac{m_s v_{||}^2}{q_s B_{||}^*} \nabla \times \mathbf{b} \right] \cdot \nabla \langle \phi_{\text{ant}} \rangle . \tag{1.10}$$

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Integrating electromagnetic antenna in the eqs. of motion

$$\dot{\mathbf{R}}^{(1)} = \dot{\mathbf{R}}_{\text{plasma}}^{(1)} + \frac{\mathbf{b}}{B_{||}^{*}} \times \nabla \langle \phi_{\text{ant}} \rangle - \frac{\mathbf{b}}{B_{||}^{*}} \times \nabla \langle v_{||} A_{||,\text{ant}}^{(s)} \rangle, \qquad (1.11)$$

$$\dot{v}_{||}^{(1)} = \dot{v}_{||,\text{plasma}}^{(1)} - \frac{q_{s}}{m_{s}} (\mathbf{b}^{*} - \mathbf{b}) \cdot \nabla \langle \phi_{\text{ant}} \rangle - \frac{\mu}{B_{||}^{*}} \mathbf{b} \times \nabla B \cdot \nabla \langle A_{||,\text{ant}}^{(s)} \rangle, \qquad (1.12)$$

and
$$\dot{\epsilon}^{(1)} = \dot{\epsilon}_{\text{plasma}}^{(1)} - \frac{q_s}{m_s} \left[m_s \mu \frac{\boldsymbol{b} \times \nabla B}{q_s B_{||}^*} + \frac{m_s v_{||}^2}{q_s B_{||}^*} \nabla \times \boldsymbol{b} \right] \cdot \nabla \langle \phi_{\text{ant}} \rangle.$$
 (1.13)

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ITPA-TAE test case

Uniform profile of density/temperature for ion and electron. Ad hoc magnetic field.

 $q \in [1.71, 1.87], \mathrm{and} \ \beta = 0.000448$. kinetic ion/electron.

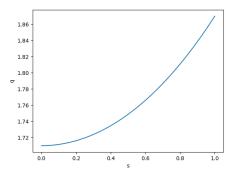
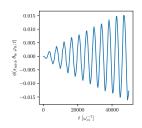
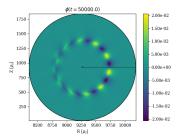
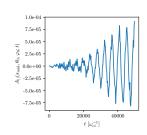


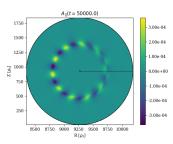
Figure: q profile

Linear simulations: n=6, m=-11,-10

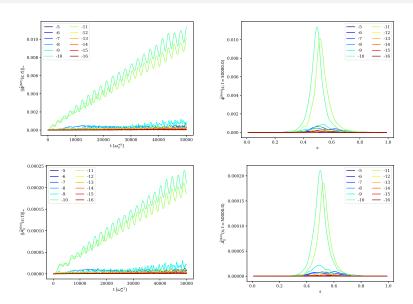








Linear simulations: n=6, m=-10,-11



Resonance frequency .vs. fast-particle simulation result

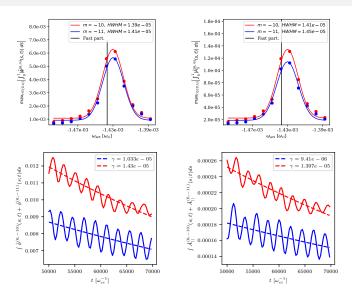
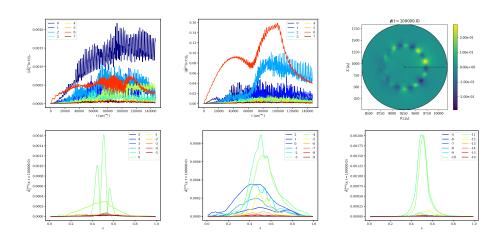
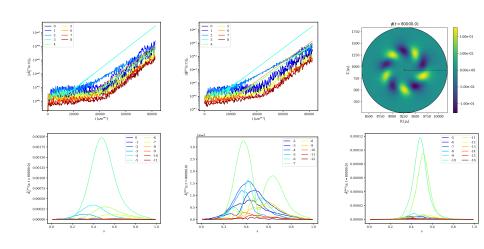


Figure: Damping rates for $\omega_{\rm ant} = -0.00144$ (blue) and 0.00144 (red).

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Conclusion and outlook

- An electromagnetic antenna was devised to excite TAE modes.
- Reasonable agreement with fast particle simulations in linear simulations.
- Damping rate can be measured by switching off antenna.
- First nonlinear simulations are shown.

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