

Room temperature tunable gate frequency InGaAs/InP single-photon detector based on ADC sampling method

Huan Chen,¹ Musheng Jiang,² Shihai Sun,¹ Guangzhao Tang,¹ and Linmei Liang^{1,3,*}

¹Department of Physics, National University of Defense Technology, Changsha, 410073 Peoples Republic of China

²Zhengzhou Information Science and Technology Institute, Zhengzhou 450004, China

³State Key Laboratory of High Performance Computing,

National University of Defense Technology, Changsha 410073, People's Republic of China

The available high speed InGaAs/InP-based single-photon avalanche detectors (SPAD) are normally worked at fixed gate repetition frequency. Here, we present a high speed InGaAs/InP-based SPAD working at telecom wavelength of 1.55 μ m with tunable gate frequency from 900MHz to 1100MHz. In the SPAD, a high speed, 14 bits resolution analog to digital convertor (ADC) is used to sample one voltage of the APD output at every gate. By comparing the voltage we can discriminate an avalanche signal from the noise. With the high resolution of the ADC, a dark count probability of 3×10^{-5} and 0.4% afterpulse probability with 1ns dead time at 10% detection efficiency at 1GHz frequency was achieved at room temperature. The wide tunable gate frequency makes the SPAD very suitable for practical use and commercial producing.

Single-photon detector is a key device in many applications, such as quantum key distribution (QKD)[1–3], quantum entanglement research[4, 5], few-photon imaging[6] and ultrasensitive spectroscopy[7]. Many kinds of infrared single photon detector have been developed, such as upconversion detectors[8], superconducting detectors[9–12] and semiconductor detectors. At near infrared telecom wavelengths, while upconversion with periodically poled LiNbO₃ (PPLN) suffers from high sophisticated optical alignment for noise rejection[13] and superconducting device requires cryogenic cooling to temperatures of around a few kelvin[9], InGaAs avalanche photodiode (APDs) are widely used in many practical applications, such as QKD systems[14] and photon pair detections[15] due to their compact and cryogen-free practicality.

High speed detectors based on APDs are mostly operated in gated mode and the most commonly used are pulsed gate[16] and sinusoidal gate[17, 18]. The single photon avalanche detector (SPAD) is gated-on within the extraordinary short time gate, such as 500ps, detecting photons only within the given time interval. This technique is particularly useful in applications of QKD where photons are coming in sequence. For the pulsed gate, self-differencing (SD) method is used to restrict the pulse noise[16, 19, 20]. SD technique relies on differentially balancing the spikes shifted by a repetitive cycle, allowing sufficient spike cancellation to discriminate the avalanche signals at high gating frequencies. However, the SD method limits the SPAD to only work at fixed repetition frequency [16, 19, 20] or narrow tunable repetition frequency ranging from 0.987GHz to 1.033 GHz using tunable delay which can tune 45ps[21]. For the sinusoidal gate, band elimination filter (BEF) with central frequency of ω_g is used to filter the noise[17, 22, 23], where ω_g is the frequency of the sinusoidal gate. The energy of

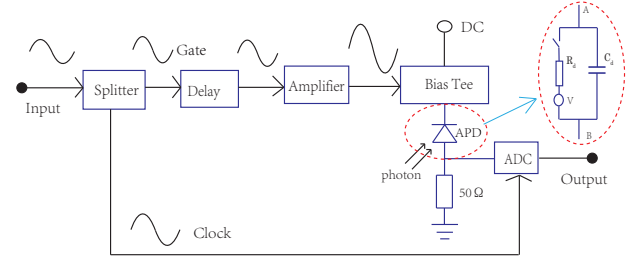


FIG. 1. Diagram of DS-SPAD. The sinusoidal signal is split into two portion, one is used as the gate of the APD and the other is used as the ADC sampling clock. The photon is fiber coupled to the APD.

the eliminated signal concentrated at the ω_g frequency components. In order not to disturb the avalanche signal, the BEF is designed with an extremely narrow band width. Significant improvements have been made using filtering method[22–24]. However, the SPADs using sinusoidal gates also work at fixed frequency because once the BEF is designed, the center frequency is difficult to tune. Alternative scheme has been tried using ADC logic circuit to process the signals[25], and it also worked at fix gate repetition frequency of 1GHz.

In this work, we present a high digital sampling single photon avalanche detector (DS-SPAD), which can work at tunable gate repetition frequency using an ADC circuit at room temperature. Instead of restricting the spike noise by self-differencing or filtering, we use an ADC to sampling the output voltage of the APD. This detector implements continuous tunable gate repetition frequency which is under test from 900MHz to 1100MHz while theoretically it can be worked at a wider range of repetition frequency.

The diagram of our DS-SPAD is shown in Figure 1. In the diagram, an external sinusoidal clock is power split into two sinusoidal waves, i.e. *Gate* and *Clock*. *Gate* is used as the gate of the APD and *Clock* is used as the sampling clock of the ADC. The *Delay* is used to adjust the time difference between the *Clock* and the *Gate*.

* nmliang@nudt.edu.cn

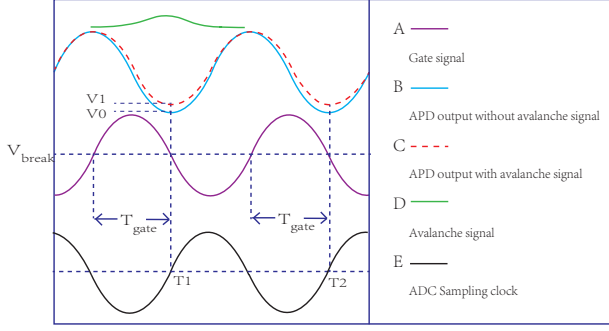


FIG. 2. Time relation between the APD gate signal, ADC sampling clock and the output signal. V_{break} is the breakdown voltage of the APD. $T1$ and $T2$ are the sampling time of the ADC. T_{gate} is the gate width.

Then the *Gate* will be amplified by a tunable amplifier. Following the amplifier, the appropriate amplified *gate* will be coupled with a DC voltage by a bias tee module. Finally the output of the APD is connected to the input of the ADC. A digitalized value will be outputted by the ADCs at every period clock.

The APD is simply equivalent to the circuit shown as the inset of Figure 1[26]. With the charge and discharge effect of the C_d , when a sinusoidal signal shown as *A* in Figure 2 is input onto the APD, a response noise shown as *B* will be generated on it. In this case, when a photon comes in the gate and triggers an avalanche, the output signal of the APD will be a superimpose of the response noise and the weak avalanche signal, shown as *C*. Here the avalanche signal shown as *D* is much weaker than the response noise. The output voltage of the APD with avalanche signal superimposed will be slightly higher than the voltage without avalanche signal, shown as *V1* and *V0* at time *T1*. Since the frequency of the ADC sampling clock and the APD gate are exactly the same and synchronized, then by adjusting the *delay* we can sample the signal exactly at point *T1* in every period and obtain one sampling value of the output signal at every gate. By comparing the sampling value we can judge whether there is an avalanche occurs or not. Since there is no frequency strongly dependent devices, the sinusoidal gate frequency can be quite broad theoretically instead of a fix frequency, which makes the gate repetition frequency of the detector tunable.

The ADC used here is ADQ7 with 14bits resolution. It will sample a value when a rise edge of the sampling clock occurs and then digitalize it. The ENOB (Equivalent Number of Bits) of the device at 2.5GHz clock is 7.3, $2^{7.3} \approx 158$, that means the resolution ratio is about 0.6%. In the experiment, the ADC outputs are saved into 16bits of binary format files for processing. The ADC also provides a real time processing module FPGA, which is essential for practical use. The InGaAs APD used in our experiment is commercially available from Princeton Lightwave. The beakdown voltage is -71V at 25°C. In the experiment, the InGaAs APD worked at room temperature of 21°C. We use a homemade pulsed

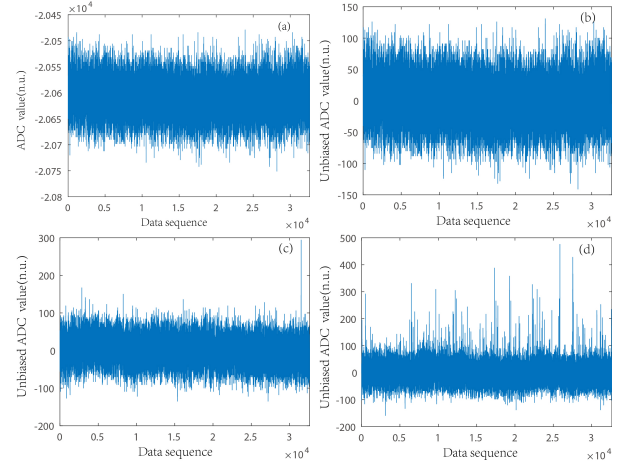


FIG. 3. (a) the direct digital output of ADC with no avalanche signal, (b)unbiased output of (a), (c)unbiased output with one dark count avalanche signal, (d)unbiased output with 0.1 photon per pulse on average incident. One data represents one gate and the gate frequency is 1GHz.

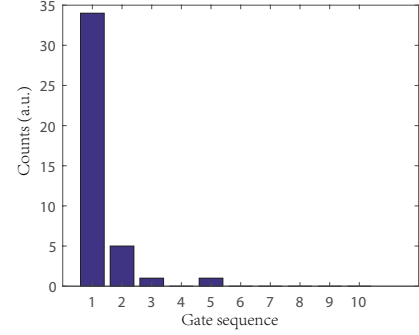


FIG. 4. The accumulation of detected counts of Figure 3(d) at every 10th gate. For example, the first bar shows the accumulation count of gate 1, gate 11, gate 21, and so forth.

laser diode at 1.55 μ m with a pulse duration of 50ps to act as single photon source. The laser is attenuated to 0.1 photon per pulse on average by the optical attenuator before coupling into the APD fiber pigtail. In addition, the laser pulse is 1/10th synchronized to the gate frequency which means if the gate frequency is 1GHz, then the laser pulse frequency is 100MHz. Figure 3 (a) is the direct digital output of ADC with laser switched off and no dark count occurs. The vertical axis value of Figure 3 (a) is the digital converted value of the sampling data. For instance, the value -2.06×10^4 represents that the sampling point voltage value is $\frac{FS}{2} \times (-2.06 \times 10^4)/2^{15}$, where FS is the Full-scale voltage of the ADC and is 1V for ADQ7. So the average voltage of the response noise at the sampling point shown in Figure 3 (a) is -315.9mV and the fluctuation of the response noise at the sampling point is $1V \times 200/2^{16} \approx 3.1mV$. When applying different gate amplitude, the average voltage of the response noise would be different. It would be inconvenient to set the threshold value. Here we subtract the mean value of Figure 3 (a) and only get the fluctuation value shown in Figure 3 (b). Figure 3 (c) is a processed output with laser switched off and one dark count occurs. Figure 3 (d) is a processed output with laser switched on. We can

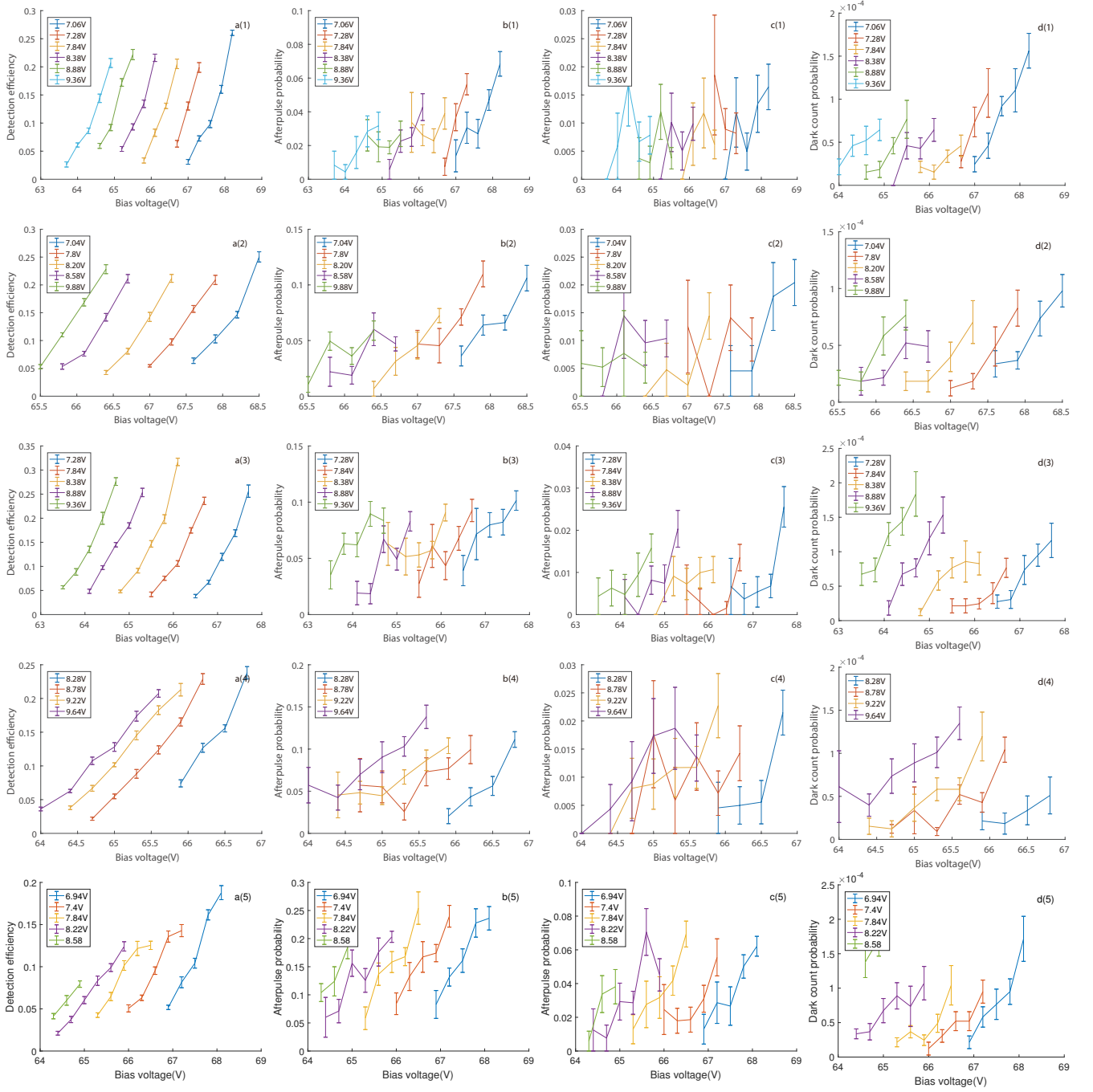


FIG. 5. The figures in the first row are measured at 900MHz gate repetition frequency and the second row at 950MHz, the third at 1000MHz, the fourth row at 1050MHz and fifth row at 1100MHz. The first column a(1)-a(5) show the detection efficiencies. The second column b(1)-b(5) show the afterpulse probabilities with no additional deadtime. The third column c(1)-(c5) show the afterpulse probabilities with 1 gate additional deadtime. The fourth column d(1)-d(5) show the dark count probabilities. The legend values in each figure are the peak-peak voltage of the sinusoidal gate, and the horizontal axis values are the bias voltage added to the sinusoidal gate.

clearly see the avalanche signal the highest avalanche signal voltage is about $1V \times 500/2^{16} \approx 7.6mV$. For each measurement, the ADC samples 32634 data. As only 1/10th of the gate is with photon input, we can plot the detected counts respected to gates sequence into Figure 4. The vertical axis is the accumulated counts of detected photons. The horizontal axis 'Gate sequence' means that

Gate sequence 1 is the accumulation of gate 1, gate 11 of ADC gate and so on, Gate sequence 2 is the accumulation of gate 2, gate 12 of ADC and so on, and the same for Gate sequence 3 to Gate sequence 10. Gate sequence 1 are the gates with photon incident, while others are with no photon incident. Here we define R_i to be the accumulation count at i th Gate sequence.

TABLE I. A comparison of our results with the other reported high-speed SPDs. P_e is the detection efficiency, P_a is the afterpulse probability, P_d is the dark count probability.

	temperature	frequency	P_e	P_a	P_d	deadtime	working mode
sinusoidal gate[23]	-50 °C	1.25GHz	10%	11.7%	4.66×10^{-6}	0 ^a	fixed gating
sinusoidal gate[22]	-60 °C	1.5GHz	10.8%	2.8%	6.3×10^{-7}	50 μ s	fixed gating
self-differencing[27]	-30 °C	921MHz	9.3%	3.4%	2.8×10^{-6}	10ns	fixed gating
self-differencing[24]	20 °C	1GHz	10%	$\sim 2\%$	$\sim 1.2 \times 10^{-5}$	0	fixed gating
self-differencing[28]	20 °C	1GHz	10%	$\sim 1.7\%$	$\sim 2 \times 10^{-5}$	— ^b	fixed gating
	20 °C	1GHz	25%	$\sim 2.8\%$	5.9×10^{-5}	—	fixed gating
self-differencing[21]	-30 °C	2GHz	11.8%	1.43%	3.79×10^{-6}	0	fixed gating
	-30 °C	1GHz	10%	2.5%	6.6×10^{-6}	0	tunable gating(0.987 ~ 1.033GHz)
NFAD[29]	-110 °C	-	10%	2.2%	1Hz	20 μ s	free running
id210	-	-	10%	-	3.3KHz	50 μ s	free running
	-	-	10%	-	2.04×10^{-6}	10 μ s	fixed gating
this work	21 °C	1GHz	10%	0.4%	3×10^{-5}	1ns	tunable gating(0.9 ~ 1.1GHz)
	21 °C	1GHz	20%	0.7%	5.6×10^{-5}	1ns	tunable gating(0.9 ~ 1.1GHz)

^a 0 dead time represents dead time free

^b — represents no data

We tested the detector at different gate repetition frequencies range from 900MHz to 1100MHz. The results are shown in Figure 5. The values in the legends represent the gate peak-peak voltages. The dark count probability shown in Figure 5 d(1)-d(5) are measured by switching off the incident laser and counting the ADC outputs which exceed the threshold value. The dark count probability is defined as: $P_d = \frac{\sum_{i=1}^{10} R_i}{R_{gate}}$, where R_{gate} is the total count of gate. The detection efficiency is calculated by: $P_e = 100 \times R_1/R_{gate}$ and shown in Figure 5 a(1)-a(5). Figure 5 b(1)-b(5) show the first gate afterpulse probability which is defined as: $P_{a1} = R_{12}/R_1$, here R_{12} is a conditional count that when there is a count in gate 1 and then also gets a count in the next gate, i.e. gate2. Similarly, Figure 5 c(1)-c(5) show the second gate afterpulse probability which is defined as: $P_{a2} = R_{13}/R_1$, where R_{13} is a conditional count that when there is a count in gate 1 and then also gets a count in gate3. The afterpulse of third and other gates are not shown here, because they are decreased to below 0.1%. Combine Figure 5 a(1)-a(5) and Figure 5 d(1)-d(5), we can see that for each gate repetition frequency, there is an optimum peak-peak voltage to achieve optimum dark count probability. The optimum peak-peak voltages are 7.84V, 7.8V, 7.84V, 8.28V and 7.84V at 900MHz, 950MHz, 1000MHz, 1050MHz and 1100MHz gate repetition frequency respectively. The optimum peak-peak voltages show a very nice consistency. To get a more clear comparison of the performance at different repetition frequency, we select the datasets of detection efficiency, afterpulse probability and dark count probability measured at 7.84V, 7.8V, 7.84V, 8.28V and 7.84V peak-peak voltage at 900MHz 950MHz, 1000MHz, 1050MHz and 1100MHz respectively. Then we do the linear fit to the data. Figure 6 shows the mean value with error bars and fitting lines of the detection efficiency, dark count probability and afterpulse probability with 0, 1 and 2 gates of dead time. From the fitted line, we obtain the afterpulse probability and dark count

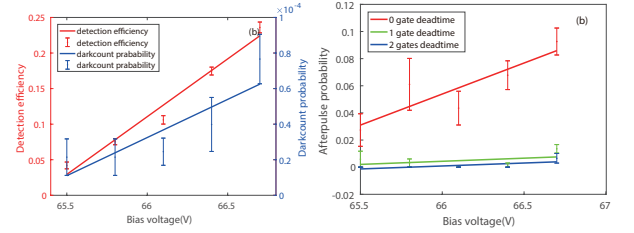


FIG. 6. Specifics at 1000MHz gate repetition frequency with 7.84V peak-peak gate voltage. (a) Detection efficiency and dark count probability, (b) afterpulse probability with 0, 1 and 2 gates of dead time.

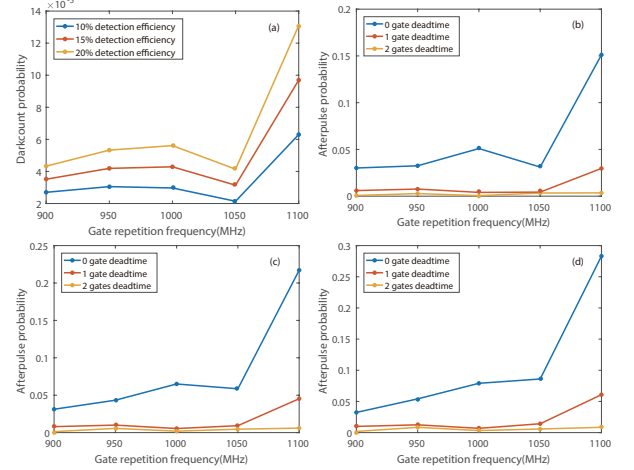


FIG. 7. (a) Dark count probability at different gate repetition frequency. (b) Afterpulse probability at 10% detection efficiency. (c) Afterpulse probability at 15% detection efficiency. (d) Afterpulse probability at 20% detection efficiency.

probability at 10%, 15% and 20% detection efficiency, which are shown in Figure 7.

Figure 7(a) shows the dark count probability respected to the gate repetition frequency. At 10% detection efficiency, the dark count probability is around 3×10^{-5} . The dark count probability increases with the gate rep-

etition frequency. Figure 7 (b) shows that the afterpulse probability with no additional gate dead time is quite high, ranging from 3% to 15% at 10% detection efficiency. However, if 1 additional gate software dead time is used, the afterpulse will be greatly decreased to 0.6% to 3%. What's more, if 2 additional gate software dead time is used, the afterpulse will be greatly decreased to 0.08% to 0.35%. At 15% detection efficiency and 20 % detection efficiency also have the same decrease trend as shown in Figure 7 (c) and Figure 7 (d).

Table I shows a comparison of our results at 1GHz gate repetition frequency with the other reported high-speed SPDs. At both 10% and 20% detection efficiency, we achieves a very low afterpulse probability with 1 additional gate dead time, which means the detection rate can reach as high as 500MHz. Also the dark count probabilities are at the same level with other SPADs which also working at room temperature.

In conclusion, we report a single photon avalanche detector DS-SPD which can work at continuous tunable

frequency from 900MHz to 1100MHz at sinusoidal gate mode. An ADC is used to sample the output of the APD without suppressing the response noise. By using this technique, we can discriminate a very weak avalanche signal from the response noise. As a result, a dark count probability of 3×10^{-5} and afterpulse probability of 0.4% with 1ns software dead time working at room temperature at 10% detection efficiency and 1GHz gate repetition frequency was achieved. Furthermore, the circuits of the detector is quite simple and works at room temperature make it very robust and portable. The most important is that the APD can work in high speed at broad frequency of the sinusoidal gate instead of a fixed or narrow frequency used in SD or filtering method. With the tunable gate frequency, the SPAD can be commercially produced and exploited for different and innovative applications.

The authors thank Queentest for providing ADQ7 product. This work is supported by the National Natural Science Foundation of China, Grants No.61072071 and No.11304391. L.-M.L. is supported by the Program for New Century Excellent Talents.

-
- [1] H. Lo, M. Curty, and K. Tamaki, *Nature Photon.* **8**, 595 (2014).
 - [2] G.-Z. Tang, S.-H. Sun, H. Chen, C.-Y. Li, and L.-M. Liang, *Chinese Physics Letters* **33**, 120301 (2016).
 - [3] G.-Z. Tang, S.-H. Sun, F. Xu, H. Chen, C.-Y. Li, and L.-M. Liang, *Physical Review A*, **94** (2016).
 - [4] H. Jayakumar, A. Predojevi, T. Kauten, I. Huber, Tobias, G. S. Solomon, and G. Weihs, *Nature Communications* **5**, 5251 (2014).
 - [5] J. Yin, Y. Cao, Y. H. Li, S. K. Liao, L. Zhang, J. G. Ren, W. Q. Cai, W. Y. Liu, B. Li, and H. Dai, *Science* **356**, 1140 (2017).
 - [6] H. Zhou, Y. He, L. You, S. Chen, W. Zhang, J. Wu, Z. Wang, and X. Xie, *Opt. Express* **23**, 14603 (2015).
 - [7] R. E. Warburton, A. McCarthy, A. M. Wallace, S. Hernandez-Marin, R. H. Hadfield, S. W. Nam, and G. S. Buller, *Opt. Lett.* **32**, 2266 (2007).
 - [8] H. Pan, H. Dong, H. Zeng, and W. Lu, *Appl. Phys. Lett.* **89**, 191108 (2006).
 - [9] F. Marsili, V. Verma, and S. W. Nam, *Nature Photon.* **7**, 210 (2013).
 - [10] A. E. Lita, A. J. Miller, and S. W. Nam, *Opt. Express* **16**, 3032 (2002).
 - [11] C. M. Natarajan, M. G. Tanner, and R. H. Hadfield, *Superconductor Science and Technology* **25**, 063001 (2012).
 - [12] W. Zhang, L. You, H. Li, J. Huang, C. Lv, L. Zhang, X. Liu, J. Wu, Z. Wang, and X. Xie, *Science China Physics, Mechanics & Astronomy* **60**, 120314 (2017).
 - [13] M. A. Albota and F. N. C. Wong, *Opt. Lett.* **29**, 1449 (2004).
 - [14] A. Rubenok, J. A. Slater, P. Chan, I. Lucio-Martinez, and W. Tittel, *Phys. Rev. Lett.* **111**, 130501 (2013).
 - [15] T. Gnathner, B. Pressl, K. Laiho, J. Geler, S. Hfling, M. Kamp, C. Schneider, and G. Weihs, *Journal of Optics* **17**, 125201 (2015).
 - [16] Z. L. Yuan, B. E. Kardynal, A. W. Sharpe, and A. J. Shields, *Appl. Phys. Lett.* **91**, 041114 (2007).
 - [17] N. Namekata, S. Sasamori, and S. Inoue, *Opt. Express* **14**, 10043 (2006).
 - [18] W.-H. Jiang, J.-H. Liu, Y. Liu, G. Jin, J. Zhang, and J.-W. Pan, *Opt. Lett.* **42**, 5090 (2017).
 - [19] L. L. Xu, E. Wu, X. R. Gu, Y. Jian, G. Wu, and H. P. Zeng, *Appl. Phys. Lett.* **94**, 161106 (2009).
 - [20] S.-B. Cho and S.-K. Kang, *Opt. Express* **19**, 18510 (2011).
 - [21] Z. L. Yuan, A. W. Sharpe, J. F. Dynes, A. R. Dixon, and A. J. Shields, *Appl. Phys. Lett.* **96**, 071101 (2010).
 - [22] N. Namekata, S. Adachi, and S. Inoue, *Opt. Express* **17**, 6275 (2009).
 - [23] X. Liang, J. Liu, Q. Wang, D. Du, J. Ma, G. Jin, Z.-B. Chen, J. Zhang, and J.-W. Pan, *REV. SCI. INSTRUM.* **83**, 083111 (2012).
 - [24] L. C. Comandar, B. Frhlich, J. F. Dynes, A. W. Sharpe, M. Lucamarini, Z. L. Yuan, R. V. Penty, and A. J. Shields, *J. Appl. Phys.* **117**, 083109 (2015).
 - [25] S. Takahashi, A. Tajima, and A. Tomita, 13th Microoptics Congerence (MOC'07), Sunport Takamatsu, Kagawa, Japan (2007).
 - [26] S. Cova, M. Ghioni, A. Lacaita, C. Samori, and F. Zappa, *Appl. Opt.* **35**, 1956 (1996).
 - [27] J. Zhang, R. Thew, C. Barreiro, and H. Zbinden, *Appl. Phys. Lett.* **95**, 091103 (2009).
 - [28] L. C. Comandar, B. Frhlich, M. Lucamarini, K. A. Patel, A. W. Sharpe, J. F. Dynes, Z. L. Yuan, R. V. Penty, and A. J. Shields, *Appl. Phys. Lett.* **104**, 175 (2014).
 - [29] B. Korzh, N. Walenta, T. Lunghi, N. Gisin, and H. Zbinden, *Appl. Phys. Lett.* **104**, 081108 (2014).