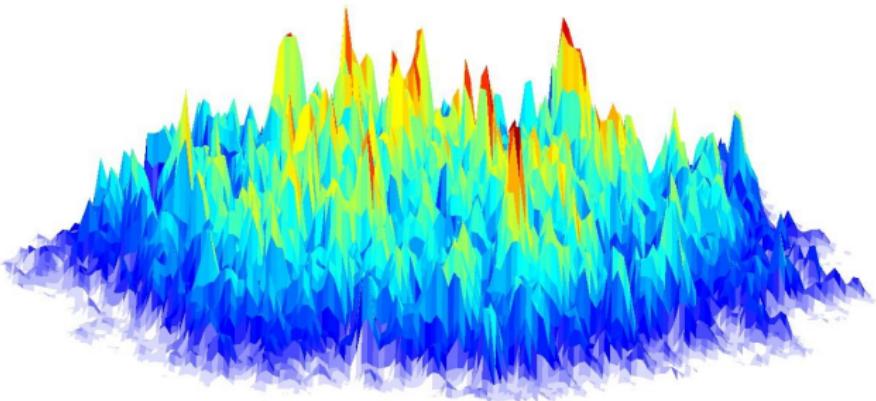




# Imprinting Quantum Fluctuations on Hydrodynamic Initial Conditions

J. Scott Moreland, Z. Qiu; Advisors S. Bass & U. Heinz

- Adding quark and gluon degrees of freedom to Monte Carlo  
Color-Glass Condensate initial conditions



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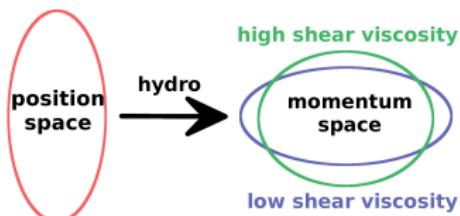
\*Special thanks to U. Heinz and S. Bass for their guidance and to the Krell Institute for funding my future graduate studies.

# Brief background: extracting the QGP shear viscosity

- Hybrid models using viscous hydrodynamics + Boltzmann cascade provide useful tool for extracting QGP shear viscosity  $\eta/s$ .

$$\eta/s \leftrightarrow \frac{v_n}{\epsilon_n} \begin{matrix} \leftarrow \text{anisotropic flow} \\ \leftarrow \text{spatial anisotropy} \end{matrix}$$

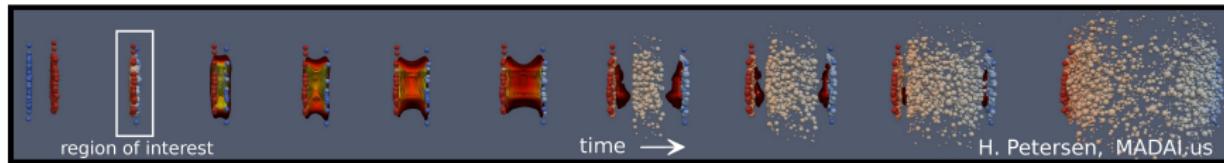
viscosity  $\eta/s$  dampens conversion of spatial anisotropy into flow anisotropy



1. Initial conditions (IC) simulate transverse energy (entropy) density
2. Hydrodynamic transport equations evolve IC down to  $T_c$  (critical temperature)
3. Boltzmann cascade evolves hadron resonance gas from  $T_c$  to freeze-out
4. Particle interactions cease and particles free-stream

(1) Initial Condition

(3) Hadron Resonance Gas (Boltzmann)



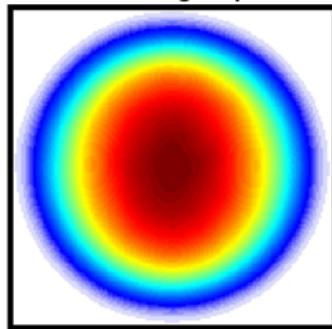
(2) QGP (Hydrodynamics)

(4) Kinetic Freezeout

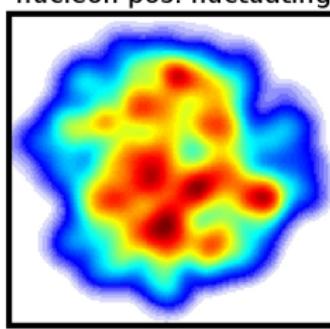
# Concentrating on initial conditions

- ▶ At present, largest source of error for extracting  $\eta/s$  is uncertainty in initial conditions:  
MC-Glauber? DIPSY? AMPT? MC-KLN? MC-rcBK? NeXuS? UrQMD? IP-Sat?  
IP-Glasma?... It's a zoo.
- ▶ Current models not yet converging, only discovering new sources of uncertainty
- ▶ Let's concentrate on developments in fluctuations  
Classic Example: Optical nucleus → Monte Carlo nucleus w/finite nucleons

event averaged profile



nucleon pos. fluctuating



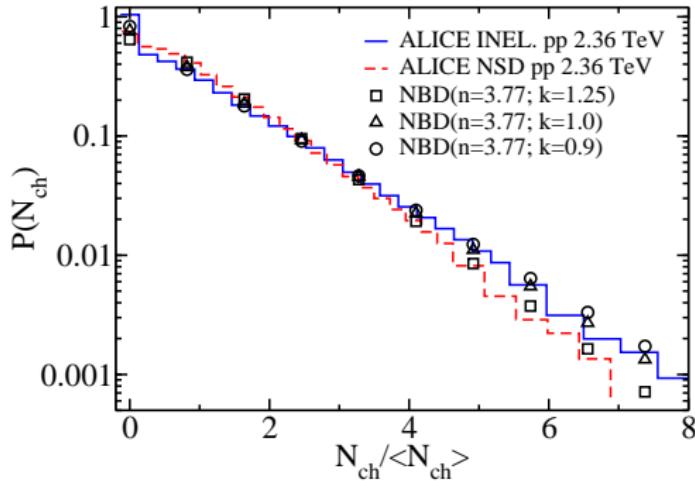
sub-nucleonic d.o.f.



# What's next? Event-by-event gluon-field fluctuations

- ▶ Large event-by-event fluctuations in pp charged particle multiplicity  $N_{ch}$  evidence for sub-nucleonic initial state fluctuations

K. Aamodt et al. (ALICE Collaboration) Eur.Phys.J. C68 (2010) 89-108

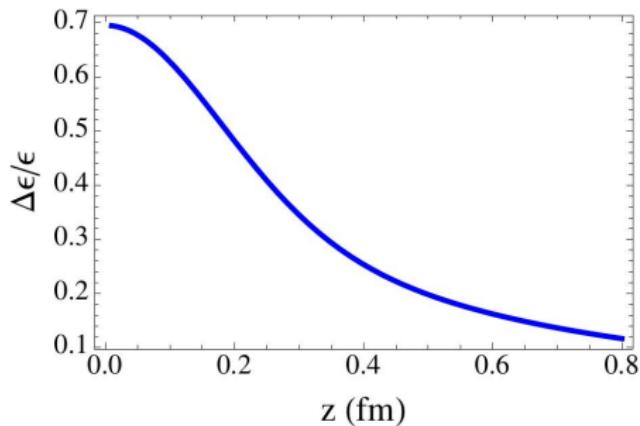


- ▶ pp multiplicity fluctuations **cannot be explained** using Glauber model for pp interaction where,  $P(b) = 1 - \exp[-\sigma_{gg} T_{pp}(b)]$
- ▶ Sub-nucleonic fluctuations clearly exist in the initial state, but where to begin?

# Müller & Schäfer calculate transverse energy density fluctuations in Color-Glass Condensate model

- ▶ Müller-Schäfer paper predicts structure of quantum fluctuations!
- ▶ Paper calculates mean normalized covar.  
 $\text{Cov}[\epsilon(r)/\epsilon_0] = (\Delta\epsilon(r)/\epsilon_0)^2$  for  
collision of two infinite sheets of  
nuclear matter
- ▶ Gaussian color distribution (color  
content contribution from each  
nucleon is independent).
- ▶ Correlation length *proportional to*  
 $1/Q_s \approx 0.14$  fm (relevant  $Q_s$   
at RHIC)
- ▶ Use Müller-Schäfer covariance of gluon-field energy density fluctuations to texture  
MC-KLN initial conditions

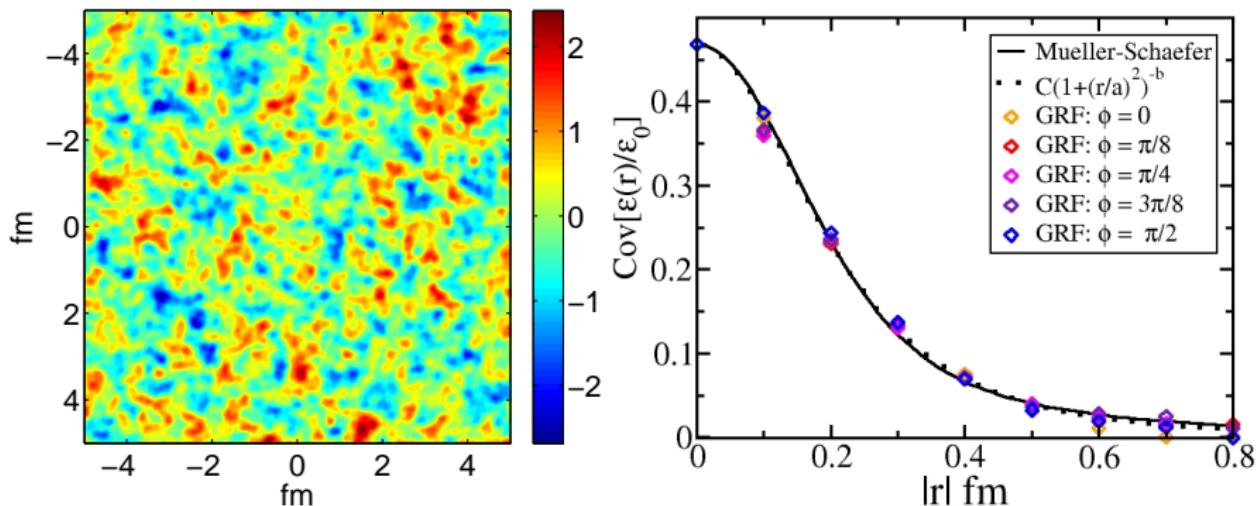
B. Müller, A. Schäfer, Phys.Rev.D 85, 114030 (2012)



# Fitting the Müller-Schäfer Gaussian random field (GRF)

We use GRF Turning Band **SIMulator** TBSIM developed by X. Emery and C. Lantuéjoul to generate a large 4000x4000 cell GRF w/ the Müller-Schäfer covariance

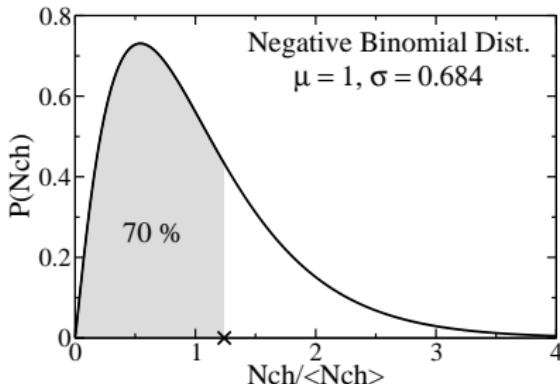
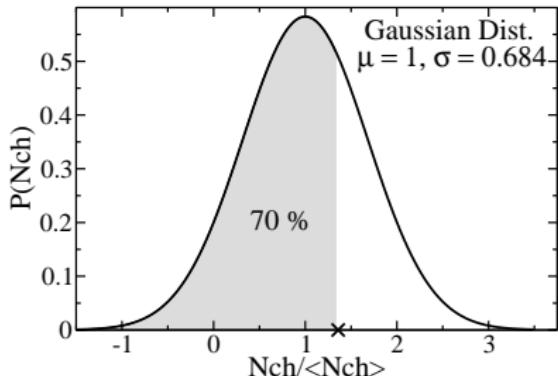
X. Emery and C. Lantuéjoul, Computers and Geosciences 32, 1615 (2006)



Fit to Müller-Schäfer covariance is nearly exact. One small problem: GRF allows for fluctuations  $\Delta\varepsilon/\varepsilon_0 < -1$ , negative energy densities which are, of course, non-physical.

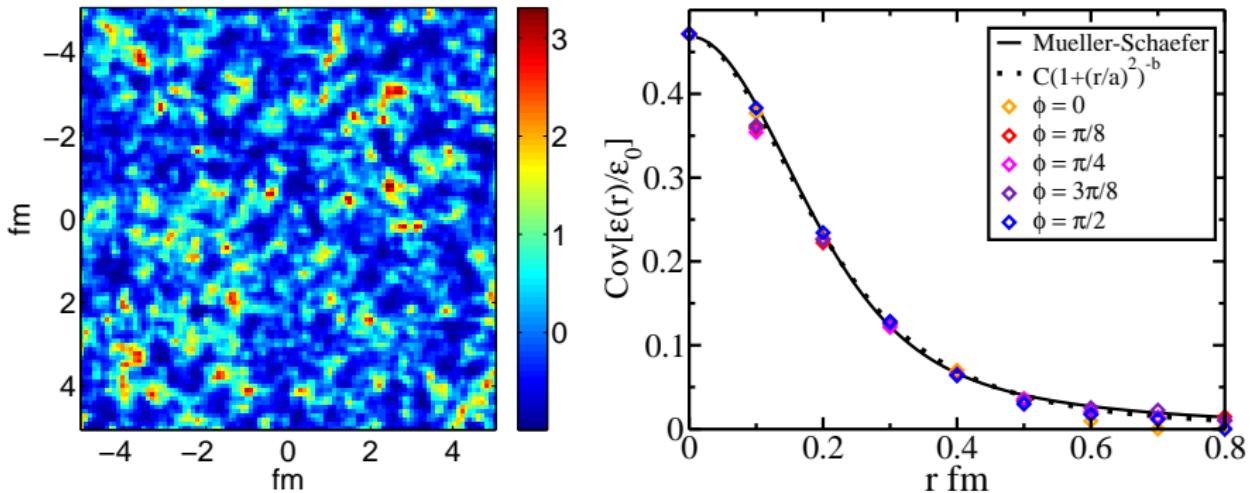
# Mapping to positive definite Neg. Binom. random field

- ▶ 7.2% chance Gaussian fluctuations are negative (non-physical)
- ▶ Send Gaussian  $\rightarrow$  Negative Binomial Distribution (NBD) motivated by p-p multiplicity fluctuations
- ▶  $NBD(\bar{n}, k; n) = \frac{\Gamma(k+n)}{\Gamma(k)\Gamma(n+1)} \frac{\bar{n}^n k^k}{(\bar{n}+k)^{n+k}}$  2-parameter function on  $\mathbb{Z}$
- ▶ Generate single NBD with mean and standard deviation equal to Gaussian
- ▶ Equate random variables via their cumulative distribution functions (CDF's)  
i.e. iff  $\int_{-\infty}^x \text{Gaussian}(x)dx = \int_{-\infty}^y \text{NBD}(y)dy$  then  $x \leftrightarrow y$



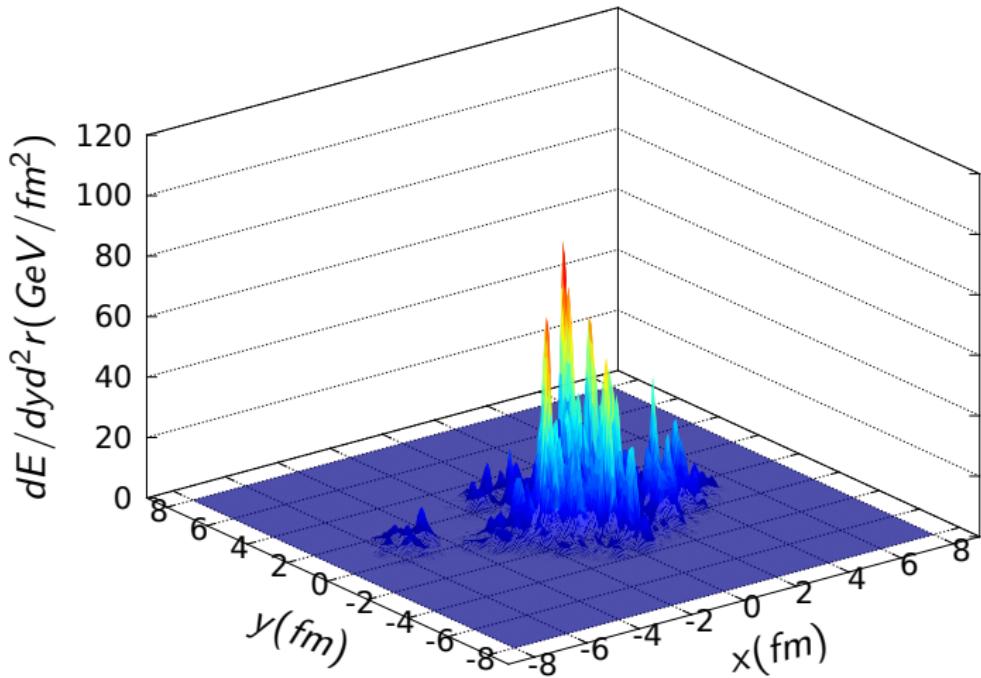
# Did this transformation destroy our covariance?

- ▶ Nope. Almost indistinguishable from GRF fit to Müller-Schäfer covariance

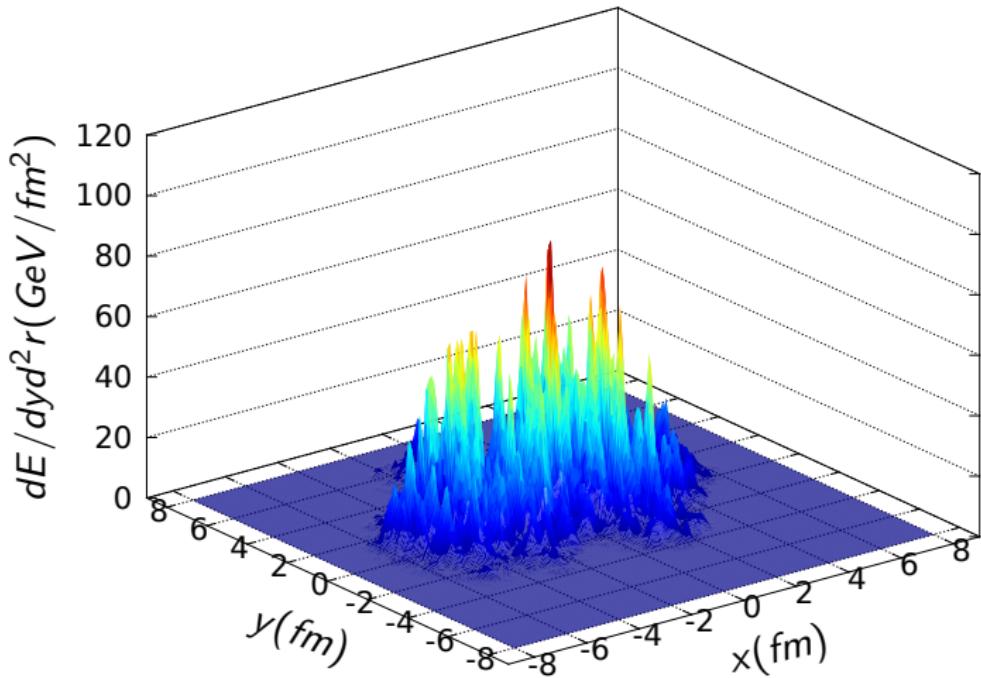


- ▶ Bulk of profile is pulled down and hot spots are more sharply peaked
- ▶ Texture can be thought of as percent fluctuation about **mean gluon-field energy density**

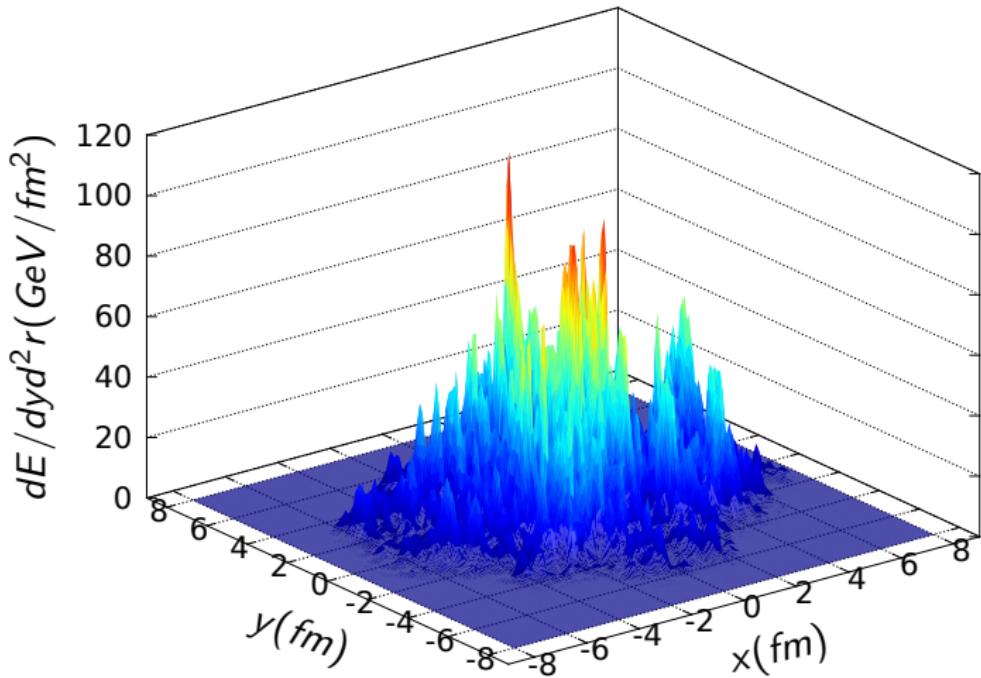
Fluctuated MC-KLN, Au-Au, 200 GeV; Npart=100



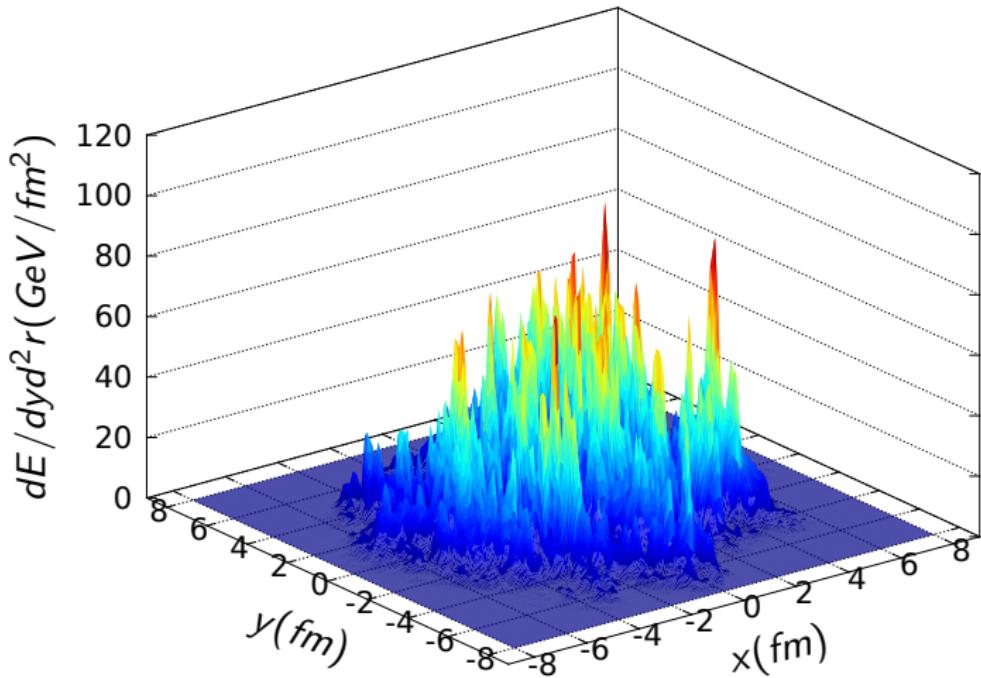
# Fluctuated MC-KLN, Au-Au, 200 GeV; Npart=200



# Fluctuated MC-KLN, Au-Au, 200 GeV; Npart=300



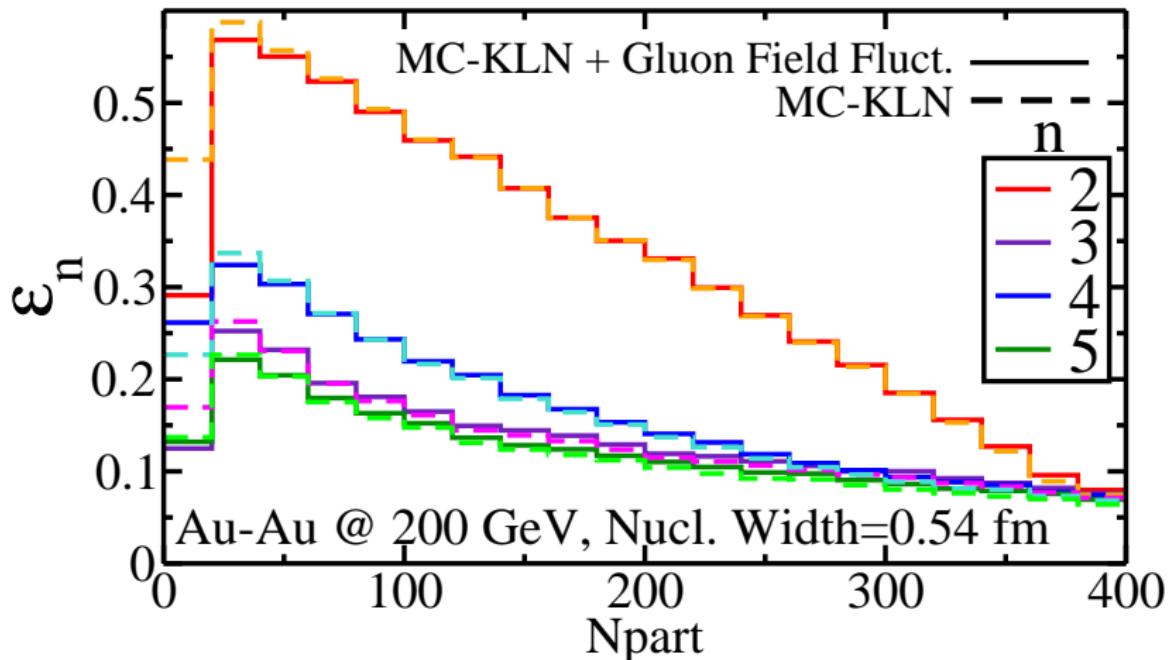
Fluctuated MC-KLN, Au-Au, 200 GeV; Npart=380



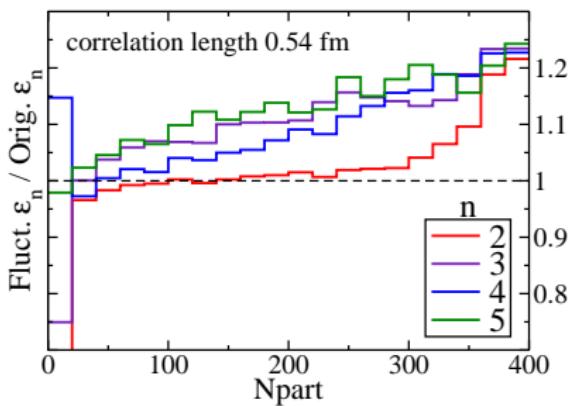
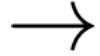
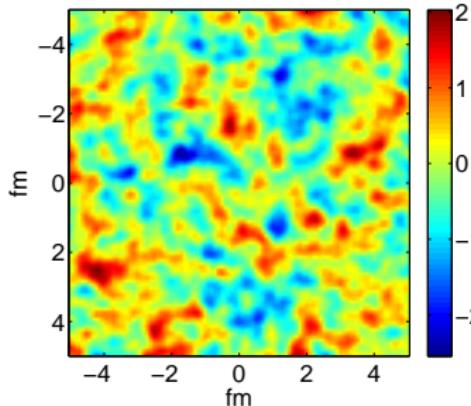
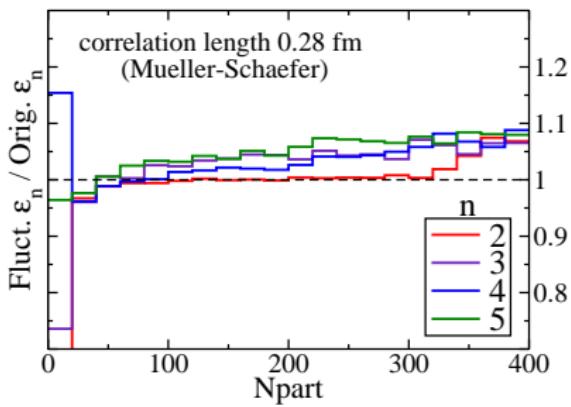
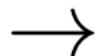
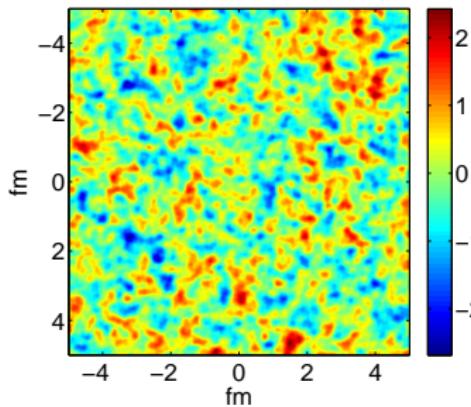
## Analysis: Calculating the eccentricity harmonics

Using definition:  $\epsilon_n e^{in\Psi_n^{PP}} = -\frac{\int dx dy r^2 e^{in\phi} \rho(x,y)}{\int dx dy r^2 \rho(x,y)}$

Quantum Fluct. w/Mueller-Schaefer Covar.



# Analysis: Dependence on correlation length



## Comparison with similar studies

### Related Publications:

- ▶ A. Dumitru & Y. Nara  
“KNO scaling of fluctuations in pp and pA, and eccentricities in heavy-ion collisions” Phys. Rev. C 85, 034907 (2012)  
[arXiv:1201.6382 [nucl-th]].
- ▶ B. Schenke, P. Tribedy & R. Venugopalan  
“Fluctuating Glasma initial conditions and flow in heavy ion collisions” arXiv:1202.6646 [nucl-th].

### Strengths of the present analysis:

- ▶ Realistic covariance applicable in central (hot) regions of the fireball
- ▶ Fluctuations can be turned on and off for comparison (factorized) and applied to any suitable Color-Glass Condensate model

## Concluding remarks

- ▶ We have implemented gluon-field fluctuations in MC-KLN initial conditions with a uniform Müller-Schäfer covariance.
- ▶ We see a small increase in  $\epsilon_2 - \epsilon_5$  in central collisions  $\approx 5 - 10\%$ .
- ▶ Increasing the correlation length to the width of a nucleon (as may happen for smaller  $Q_s$  values) results in  $\approx 20 - 25\%$  increase of  $\epsilon_n$  in central collisions.
- ▶ Increases in  $\epsilon_2$  vanish for  $N_{part} < 300$ .

## Looking forward

- ▶ Refine model to better account for fluctuations in dilute regions, i.e. fitting pp multiplicity fluctuations
- ▶ Converge with parallel studies (B. Schenke, P. Tribedy, R. Venugopalan & A. Dumitru, Y. Nara)
- ▶ Run fluctuated profiles through hydro

Special thanks to Berndt Müller, Adrian Dumitru,  
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