

## 3.2 Finite well and harmonic oscillator

Slides: Video 3.2.2 The finite potential well

Text reference: Quantum Mechanics  
for Scientists and Engineers

Section 2.9





# Particles in potential wells



## The finite potential well

Quantum mechanics for scientists and engineers

David Miller

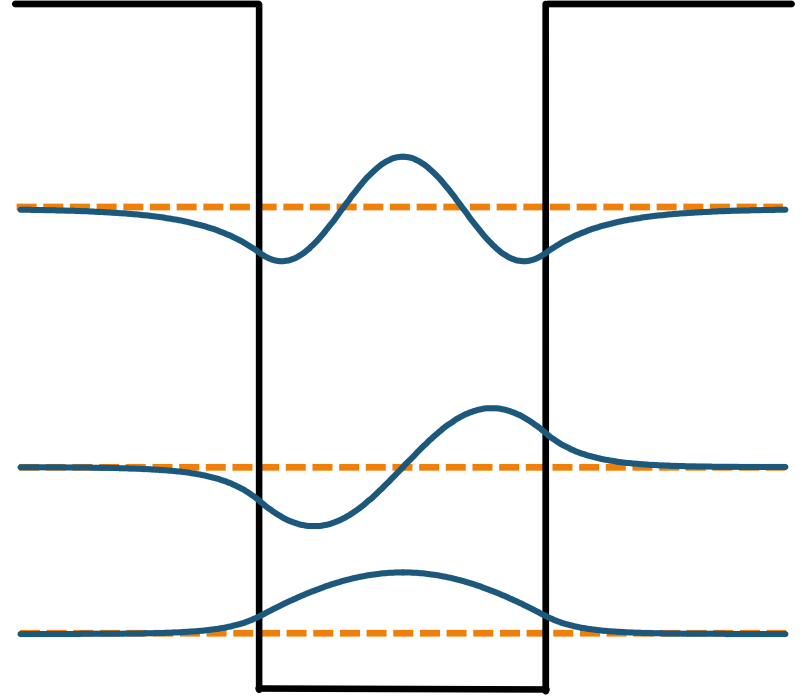
# Finite potential well

Insert video here (split screen)

## Lesson 7

Particles in potential wells

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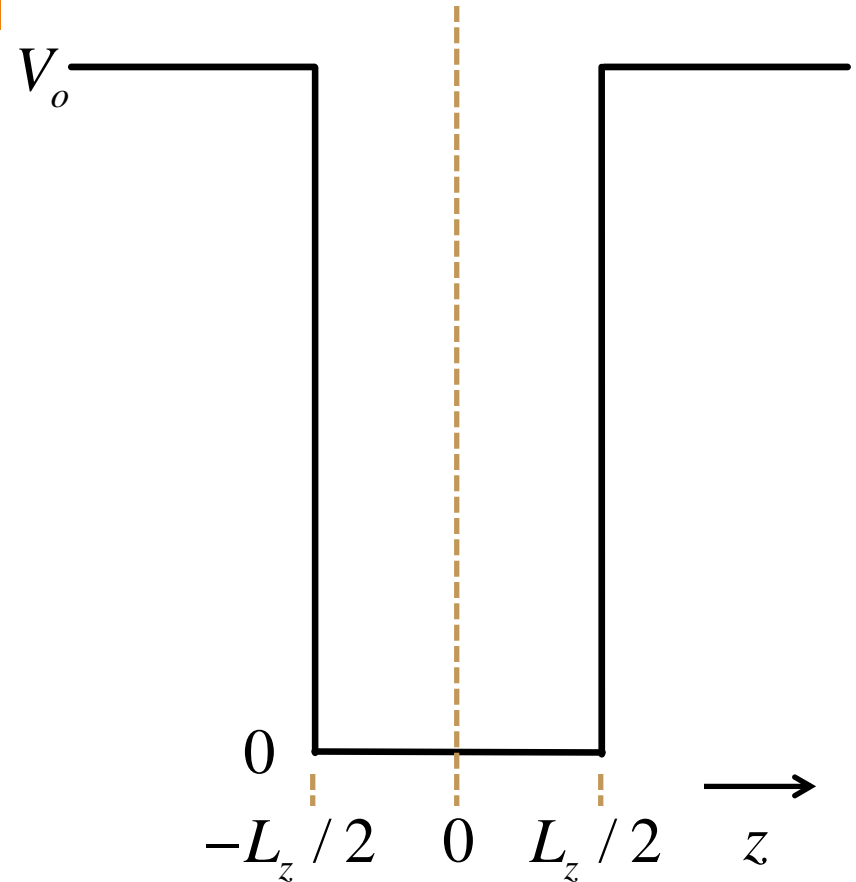


# Particle in a finite potential well

We will choose the height of the potential barriers as  $V_o$   
with 0 potential energy at the bottom of the well

The thickness of the well is  $L_z$

Now we will choose the position origin in the center of the well

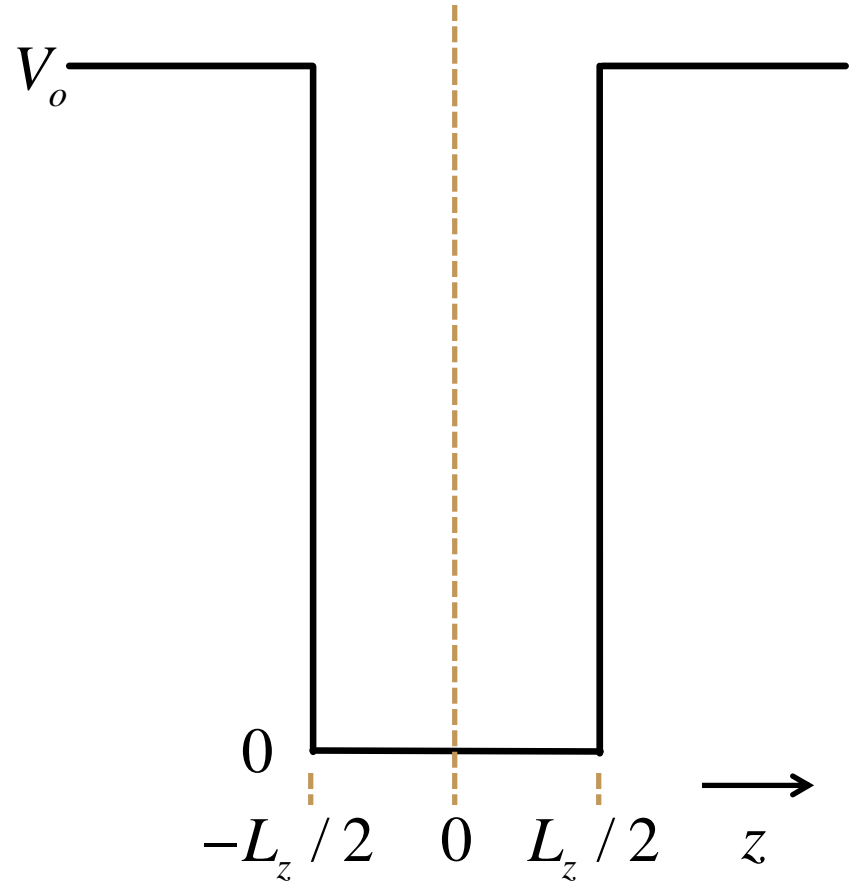


# Particle in a finite potential well

If there is an eigenenergy  $E$  for which there is a solution

then we already know what form the solution has to take

sinusoidal in the middle  
exponentially decaying on either side



# Particle in a finite potential well

For some eigenenergy  $E$

with  $k = \sqrt{2mE / \hbar^2}$

and  $\kappa = \sqrt{2m(V_o - E) / \hbar^2}$

for  $z < -L_z / 2$

$$\psi(z) = G \exp(\kappa z)$$

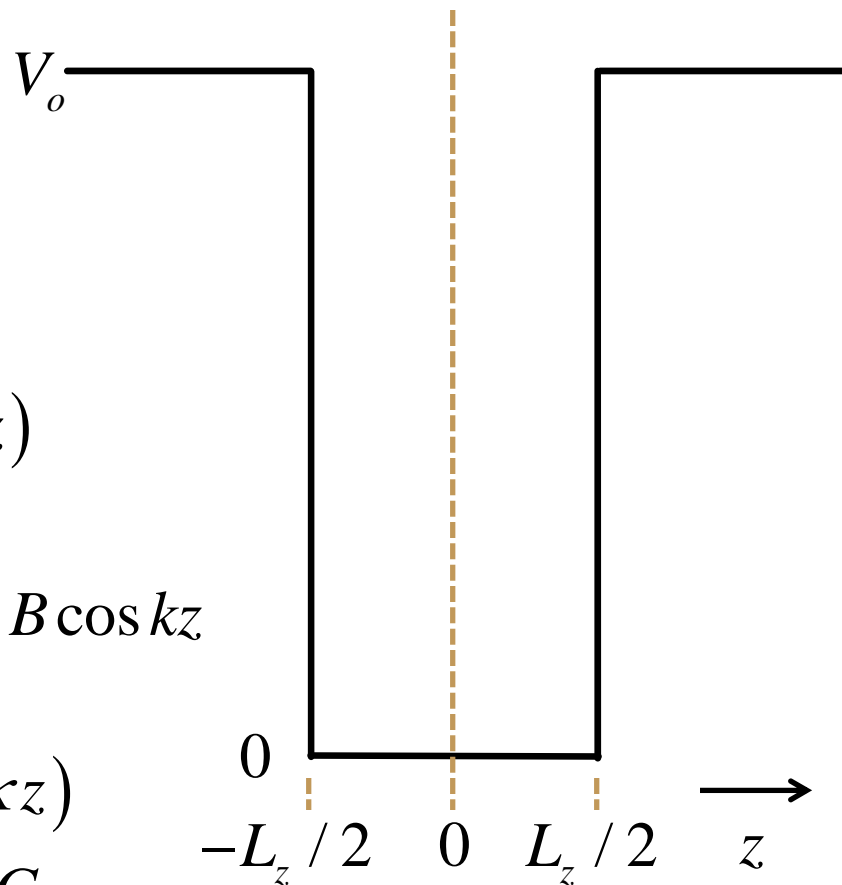
for  $-L_z / 2 < z < L_z / 2$

$$\psi(z) = A \sin kz + B \cos kz$$

for  $z > L_z / 2$

$$\psi(z) = F \exp(-\kappa z)$$

with constants  $A, B, F,$  and  $G$



# Particle in a finite potential well

Now we need to apply the boundary conditions to solve for the unknown coefficients

constants  $A, B, F$ , and  $G$

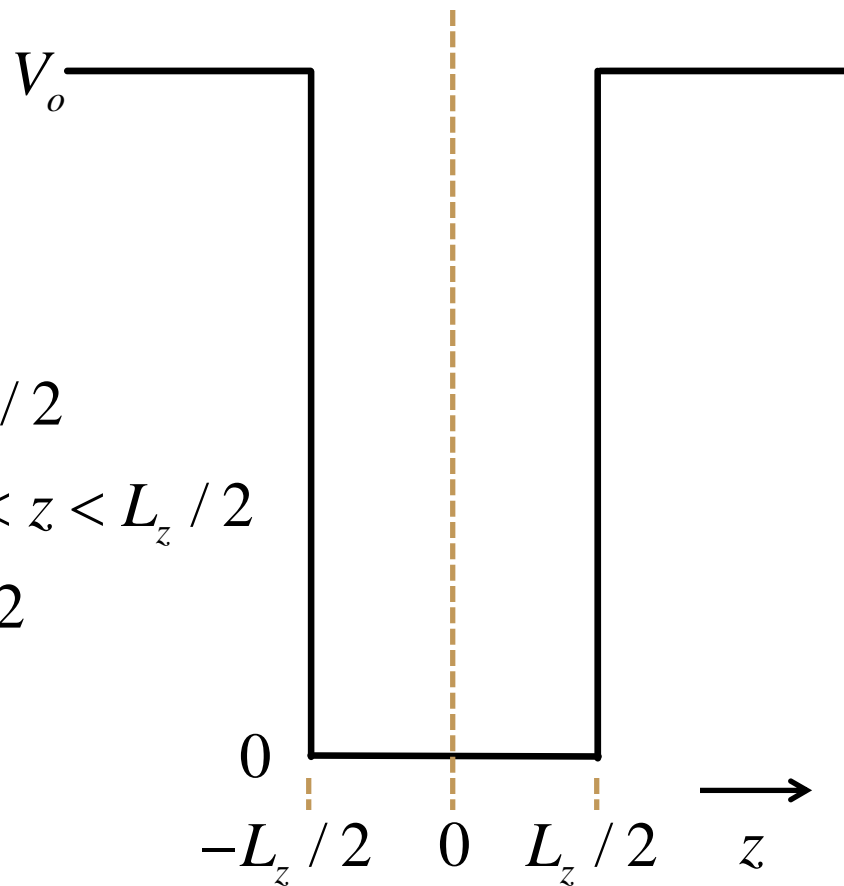
$$\psi(z) = G \exp(\kappa z) \quad z < -L_z / 2$$

$$\psi(z) = A \sin kz + B \cos kz \quad -L_z / 2 < z < L_z / 2$$

$$\psi(z) = F \exp(-\kappa z) \quad z > L_z / 2$$

or at least three of them

the fourth could be found  
by normalization



# Particle in a finite potential well

From continuity of the  
wavefunction at  $z = L_z / 2$

$$\begin{aligned}\psi(L_z / 2) &= F \exp(-\kappa L_z / 2) \\ &= A \sin(kL_z / 2) + B \cos(kL_z / 2)\end{aligned}$$

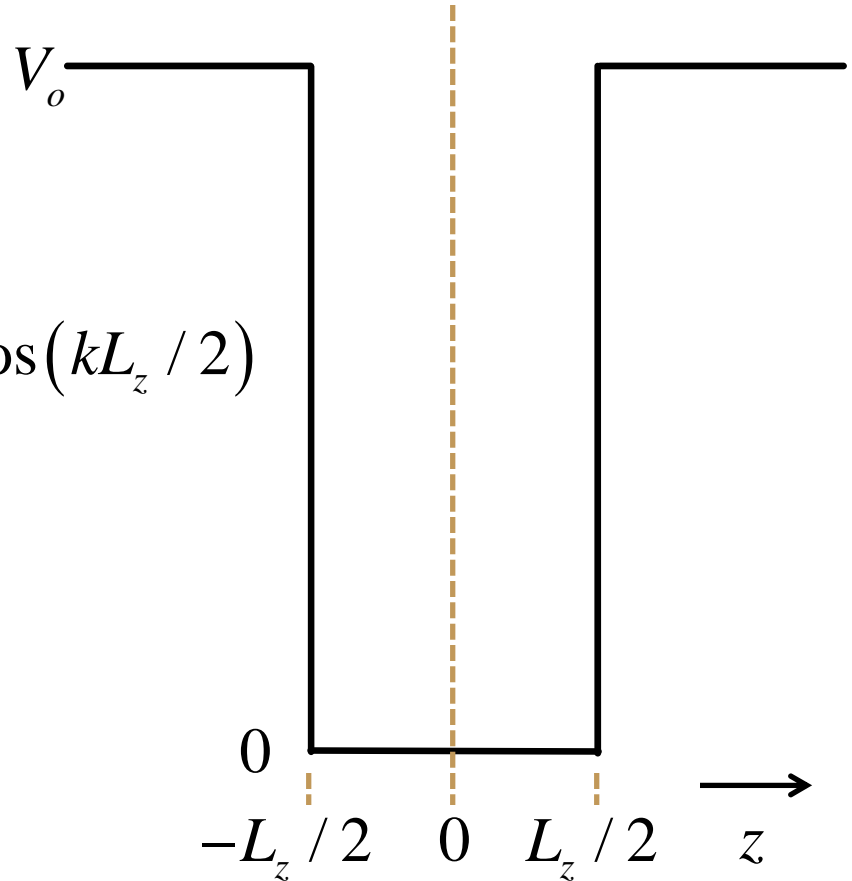
Writing  $X_L = \exp(-\kappa L_z / 2)$

$$S_L = \sin(kL_z / 2)$$

$$C_L = \cos(kL_z / 2)$$

gives

$$FX_L = AS_L + BC_L$$





# Particle in a finite potential well

Similarly at  $z = -L_z / 2$

$$GX_L = -AS_L + BC_L$$

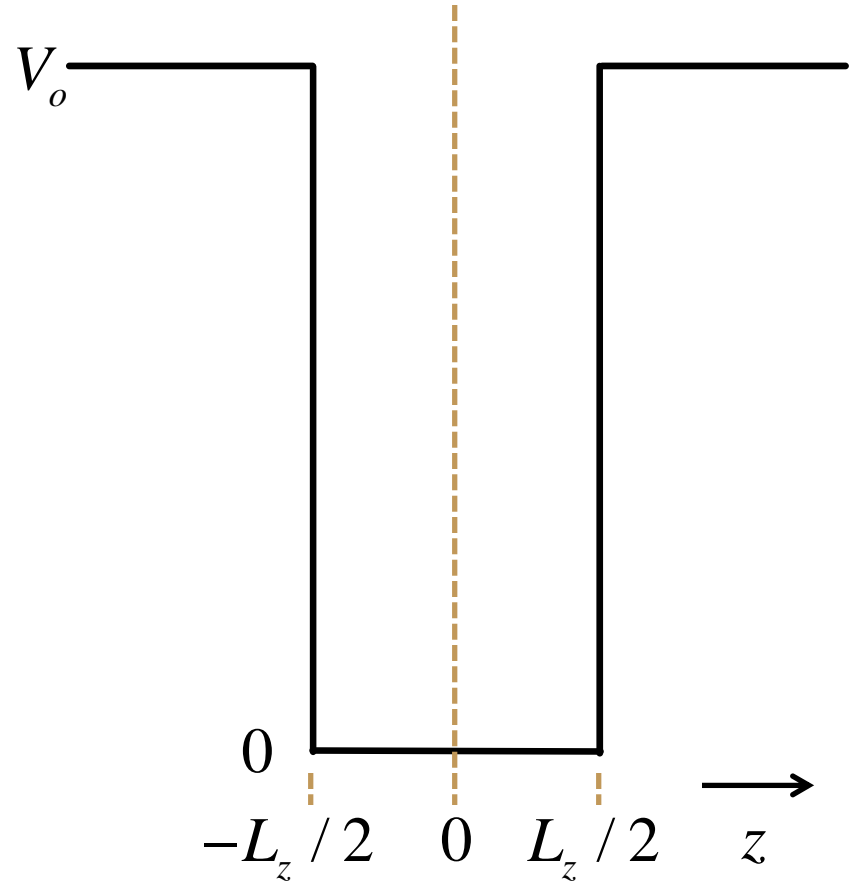
Continuity of the derivative  
gives

at  $z = -L_z / 2$

$$\frac{\kappa}{k}GX_L = AC_L + BS_L$$

at  $z = L_z / 2$

$$-\frac{\kappa}{k}FX_L = AC_L - BS_L$$



# Particle in a finite potential well

So we have four relations

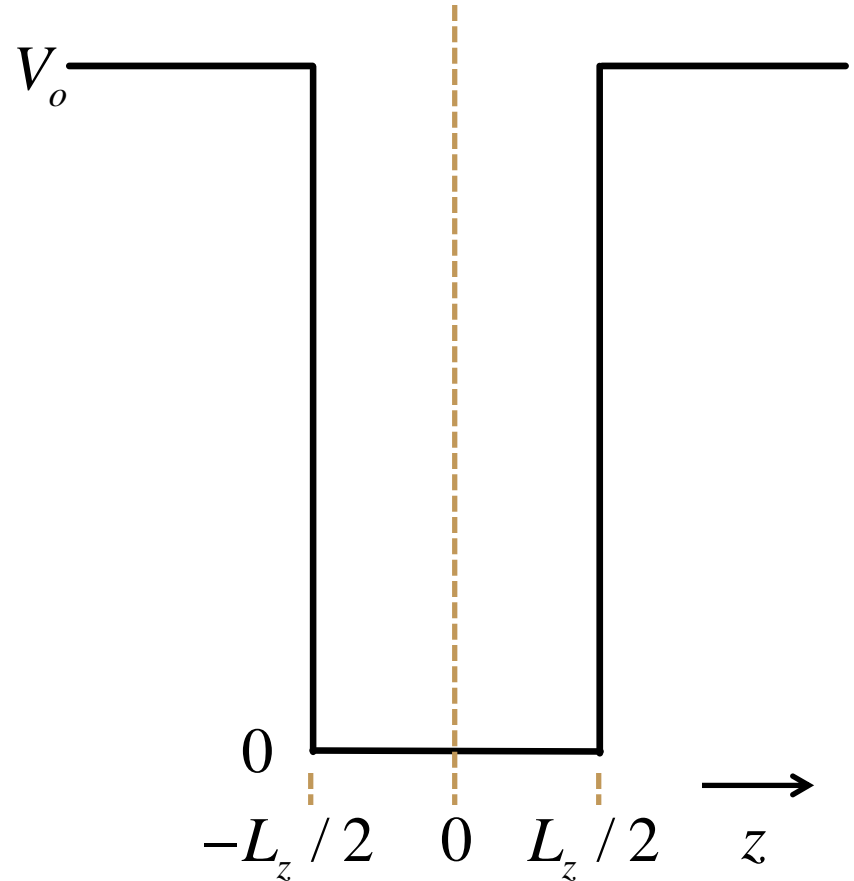
$$GX_L = -AS_L + BC_L$$

$$FX_L = AS_L + BC_L$$

$$\frac{\kappa}{k}GX_L = AC_L + BS_L$$

$$-\frac{\kappa}{k}FX_L = AC_L - BS_L$$

Now we need to find what solutions are compatible with these



# Particle in a finite potential well

Adding  $GX_L = -AS_L + BC_L$

$$FX_L = AS_L + BC_L$$

gives

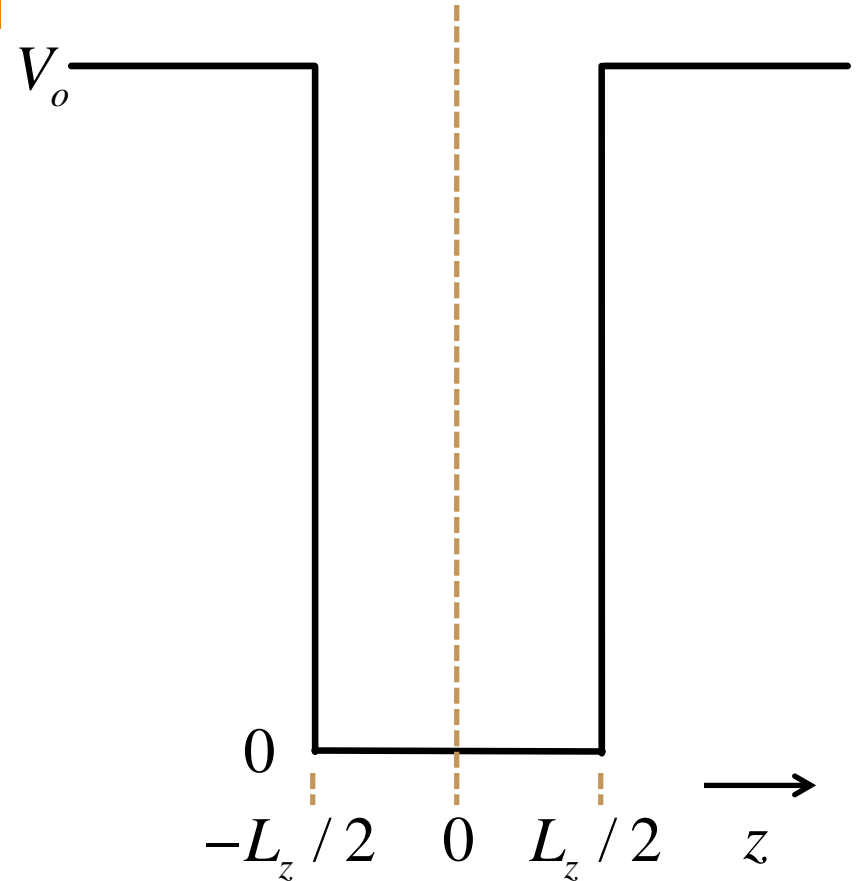
$$2BC_L = (F + G)X_L$$

Subtracting

$$-\frac{\kappa}{k}FX_L = AC_L - BS_L$$

from  $\frac{\kappa}{k}GX_L = AC_L + BS_L$

gives  $2BS_L = \frac{\kappa}{k}(F + G)X_L$



# Particle in a finite potential well

As long as  $F \neq -G$

we can divide

$$2BS_L = \frac{\kappa}{k}(F + G)X_L$$

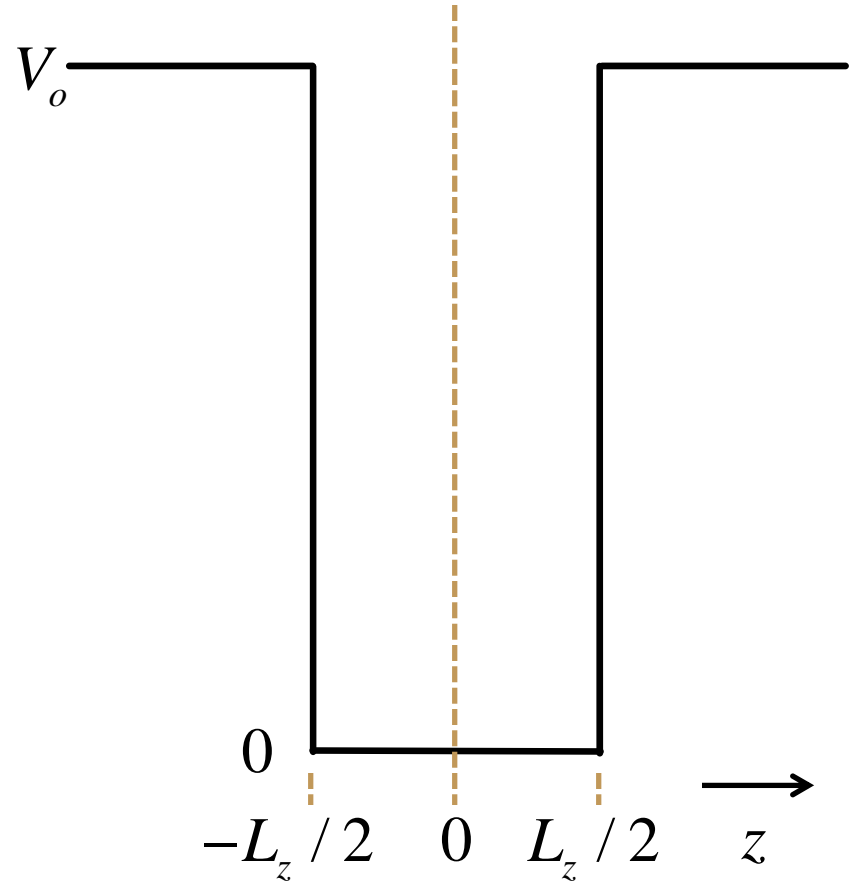
by

$$2BC_L = (F + G)X_L$$

to obtain

$$\tan(kL_z / 2) = \kappa / k$$

This relation is effectively a condition for eigenvalues



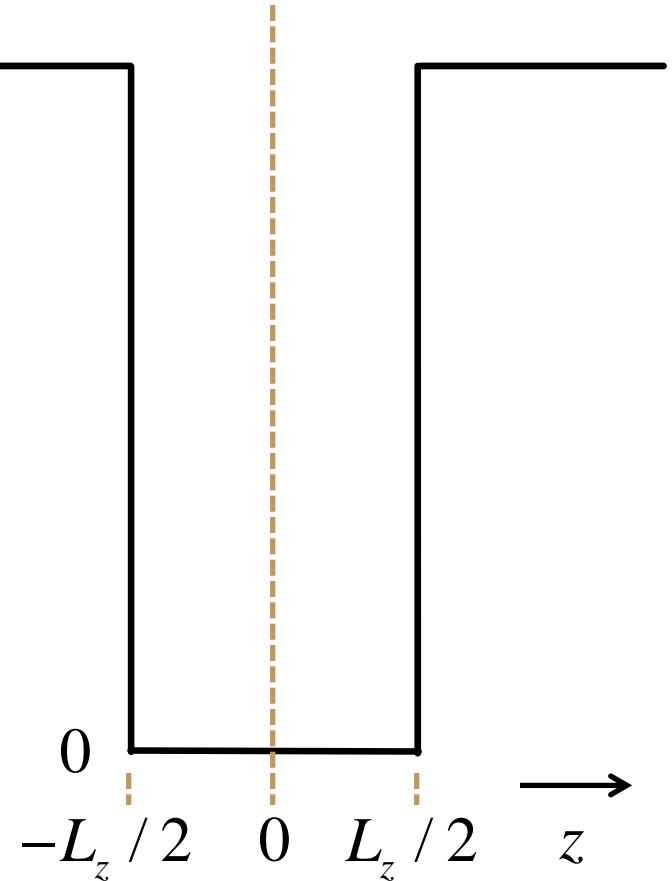
# Particle in a finite potential well

Subtracting  $GX_L = -AS_L + BC_L$   
from  $FX_L = AS_L + BC_L$   
gives  $2AS_L = (F - G)X_L$

Adding  $-\frac{\kappa}{k}FX_L = AC_L - BS_L$

and  $\frac{\kappa}{k}GX_L = AC_L + BS_L$

gives  $2AC_L = -\frac{\kappa}{k}(F - G)X_L$



# Particle in a finite potential well

Similarly, as long as  $F \neq G$

we can divide

$$2AC_L = -\frac{\kappa}{k}(F - G)X_L$$

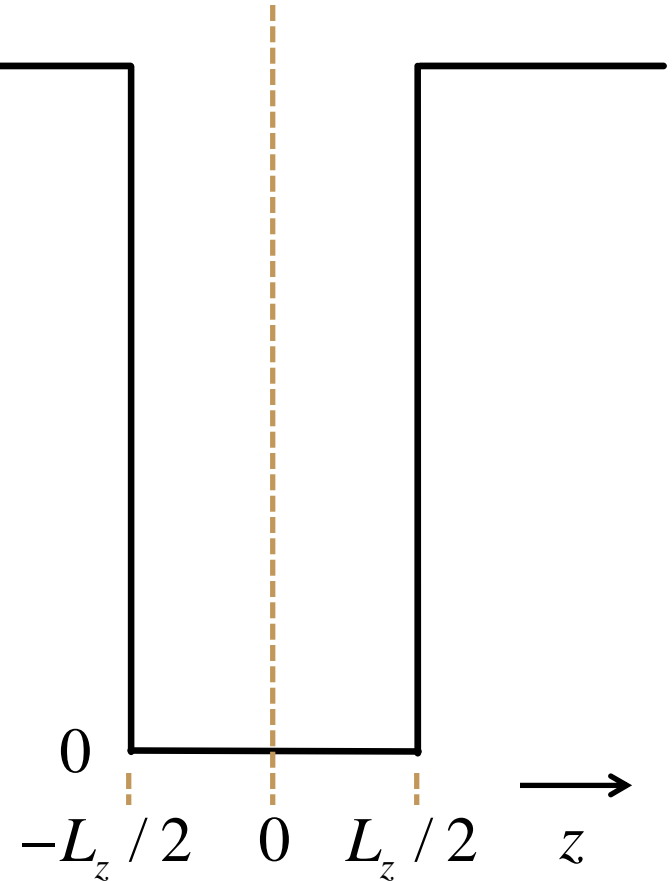
by

$$2AS_L = (F - G)X_L$$

to obtain

$$-\cot(kL_z / 2) = \kappa / k$$

This relation is also effectively a condition for eigenvalues



# Particle in a finite potential well

For any case other than  $F = G$

which leaves  $\tan(kL_z / 2) = \kappa / k$

but not  $-\cot(kL_z / 2) = \kappa / k$

or  $F = -G$

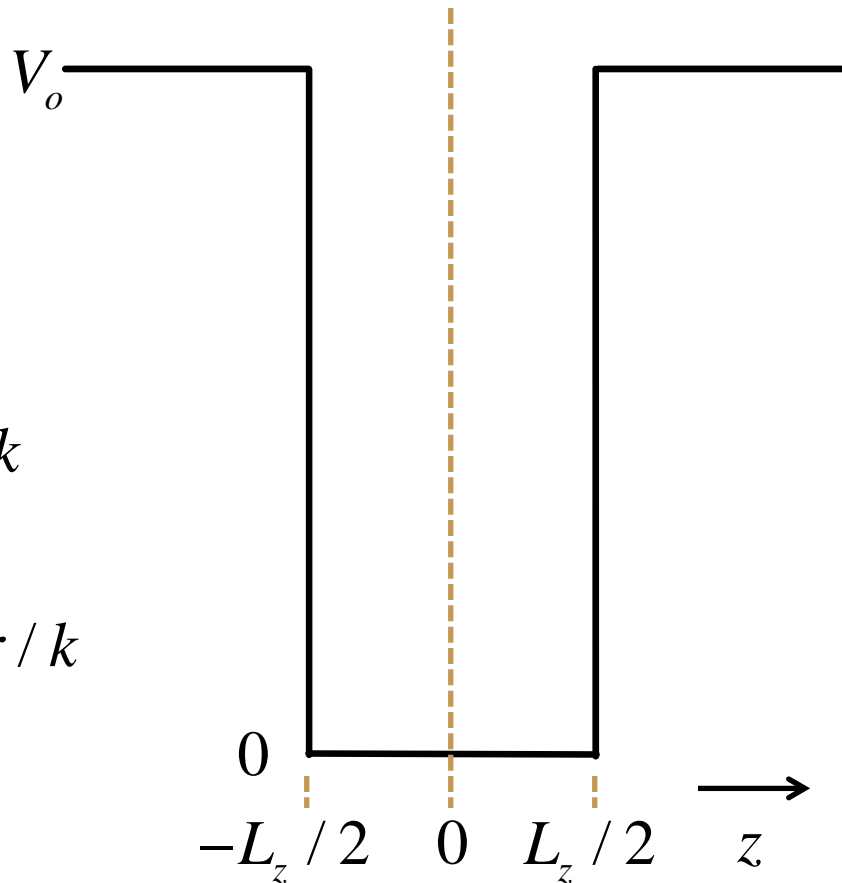
which leaves  $-\cot(kL_z / 2) = \kappa / k$

but not  $\tan(kL_z / 2) = \kappa / k$

then the solutions  $\tan(kL_z / 2) = \kappa / k$

and  $-\cot(kL_z / 2) = \kappa / k$

are contradictory



# Particle in a finite potential well

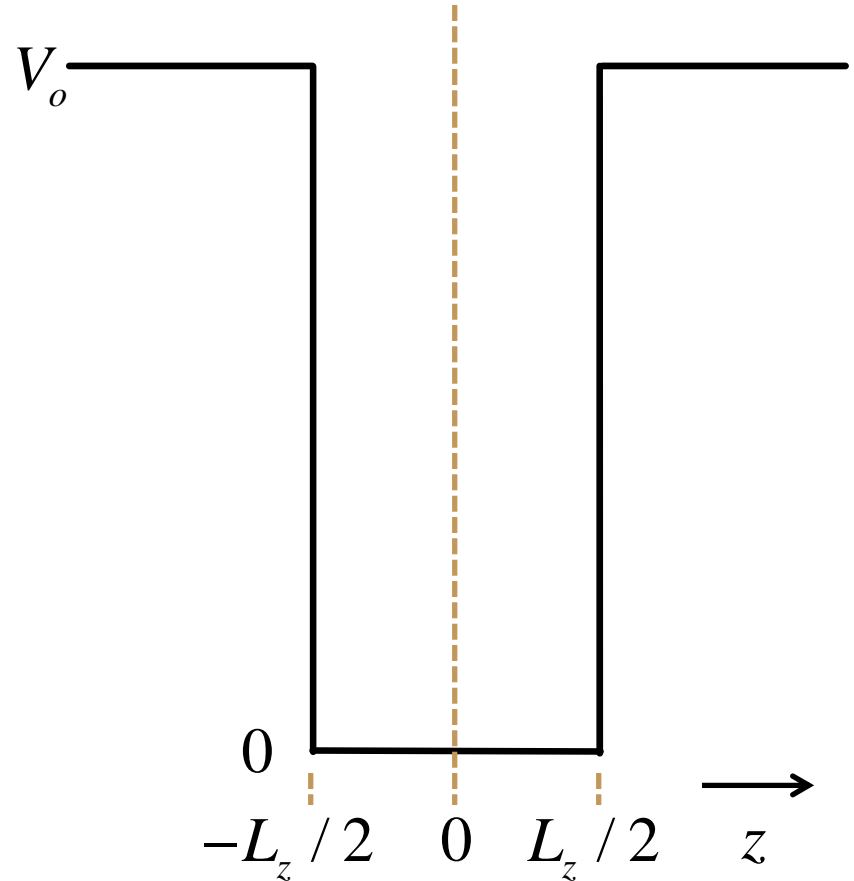
So the only possibilities are

1 -  $F = G$

and  $\tan(kL_z / 2) = \kappa / k$

2 -  $F = -G$

and  $-\cot(kL_z / 2) = \kappa / k$





# Particle in a finite potential well

1 -  $F = G$

and  $\tan(kL_z / 2) = \kappa / k$

Note from  $2AS_L = (F - G)X_L$

and  $2AC_L = -\frac{\kappa}{k}(F - G)X_L$

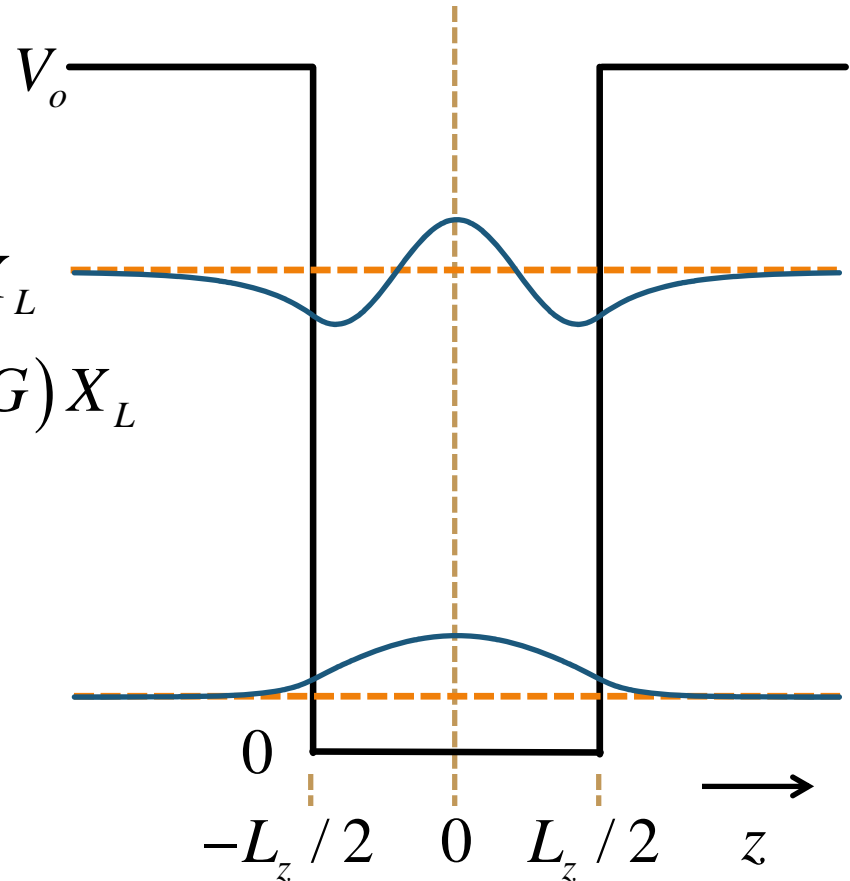
$S_L$  and  $C_L$  cannot both be 0

so  $A = 0$

Hence in the well we have

$$\psi(z) \propto \cos kz$$

which is an even function



# Particle in a finite potential well

1 -  $F = -G$

and  $-\cot(kL_z / 2) = \kappa / k$

Note from  $2BC_L = (F + G)X_L$

and  $2BS_L = \frac{\kappa}{k}(F + G)X_L$

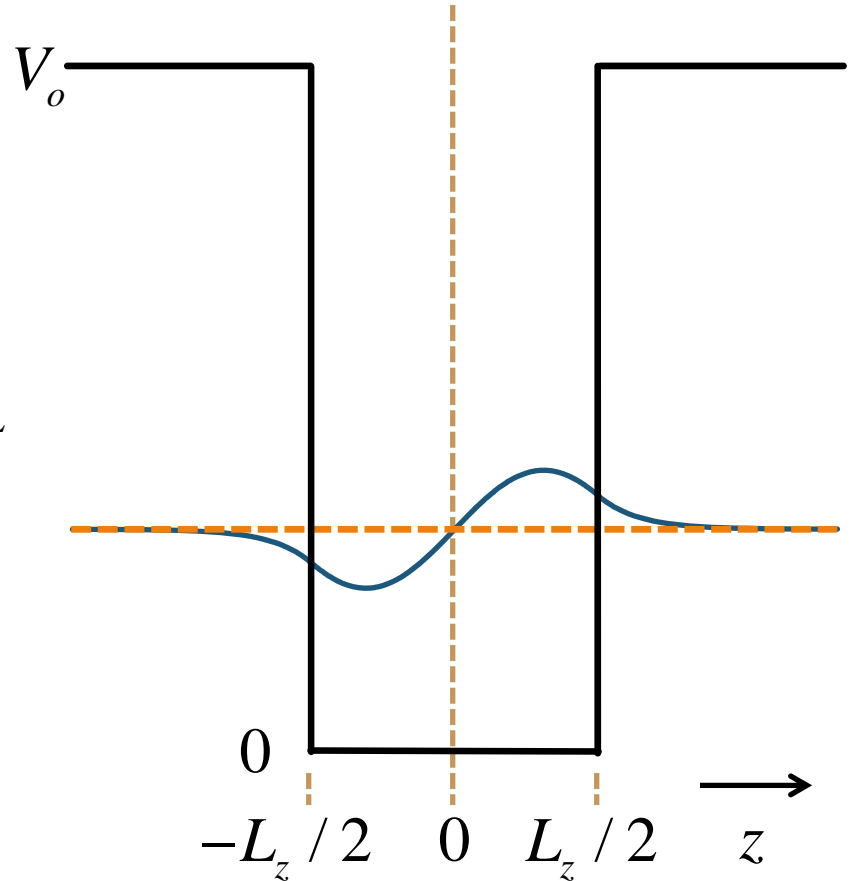
$S_L$  and  $C_L$  cannot both be 0

so  $B = 0$

Hence in the well we have

$$\psi(z) \propto \sin kz$$

which is an odd function



# Particle in a finite potential well

Though we have found the nature of the solutions

we have not yet formally solved for the eigenenergies

$E$

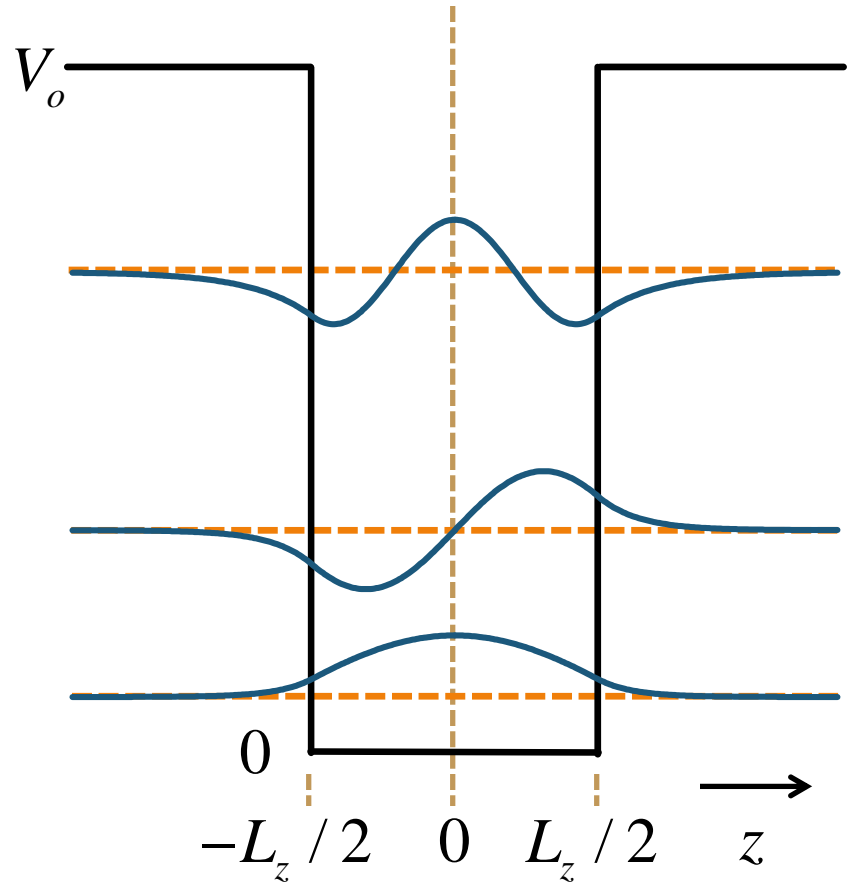
and hence for  $k$  and  $\kappa$

We do this by solving

$$\tan(kL_z / 2) = \kappa / k$$

and

$$-\cot(kL_z / 2) = \kappa / k$$



# Solving for the eigenenergies

Change to “dimensionless” units

Use the energy of the first level in the  
“infinite” potential well width  $L_z$

leading to a dimensionless  
eigenenergy

and a dimensionless barrier height

$$E_1^\infty = \frac{\hbar^2}{2m} \left( \frac{\pi}{L_z} \right)^2$$

$$\varepsilon \equiv E / E_1^\infty$$

$$v_o \equiv V_o / E_1^\infty$$

Also

$$k = \sqrt{2mE / \hbar^2} = \left( \pi / L_z \right) \sqrt{E / E_1^\infty} = \left( \pi / L_z \right) \sqrt{\varepsilon}$$

$$\kappa = \sqrt{2m(V_o - E) / \hbar^2} = \left( \pi / L_z \right) \sqrt{(V_o - E) / E_1^\infty} = \left( \pi / L_z \right) \sqrt{v_o - \varepsilon}$$

# Solving for the eigenenergies

Consequently  $\frac{\kappa}{k} = \sqrt{\frac{V_o - E}{E}} = \sqrt{\frac{v_o - \varepsilon}{\varepsilon}}$

$$\frac{kL_z}{2} = \frac{\pi}{2} \sqrt{\frac{E}{E_1^\infty}} = \frac{\pi}{2} \sqrt{\varepsilon} \quad \text{and} \quad \frac{\kappa L_z}{2} = \frac{\pi}{2} \sqrt{\frac{V_o - E}{E_1^\infty}} = \frac{\pi}{2} \sqrt{v_o - \varepsilon}$$

So  $\tan(kL_z / 2) = \kappa / k$  becomes  $\tan\left[(\pi / 2)\sqrt{\varepsilon}\right] = \sqrt{(v_o - \varepsilon) / \varepsilon}$

or  $\sqrt{\varepsilon} \tan\left[(\pi / 2)\sqrt{\varepsilon}\right] = \sqrt{(v_o - \varepsilon)}$

and  $-\cot(kL_z / 2) = \kappa / k$  becomes  $-\cot\left[(\pi / 2)\sqrt{\varepsilon}\right] = \sqrt{(v_o - \varepsilon) / \varepsilon}$

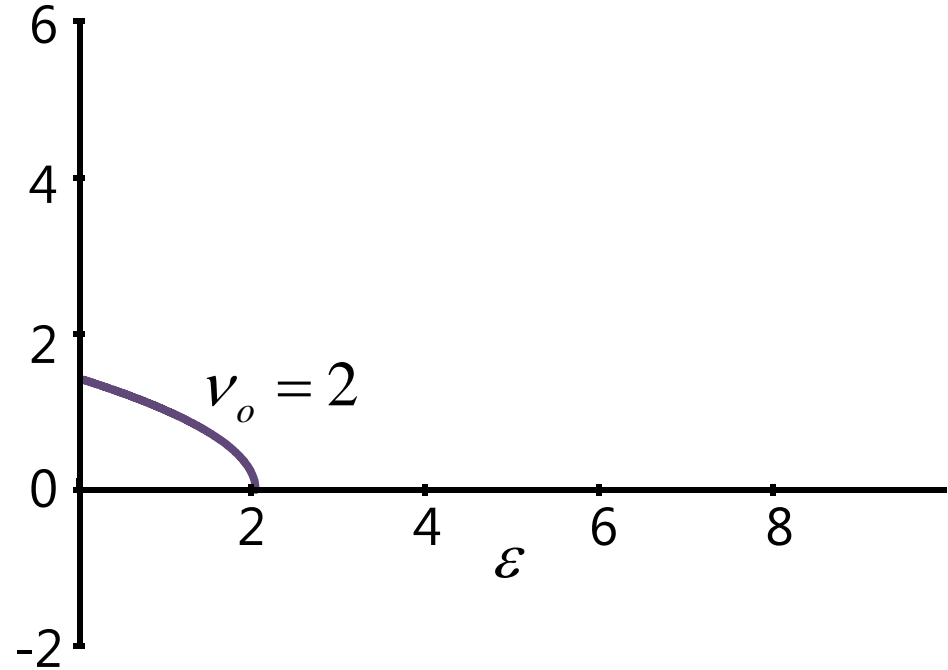
or  $-\sqrt{\varepsilon} \cot\left[(\pi / 2)\sqrt{\varepsilon}\right] = \sqrt{(v_o - \varepsilon)}$

# Graphical solution

Choose a specific well  
depth  $v_o$

and plot the curve

$$\sqrt{(v_o - \varepsilon)}$$

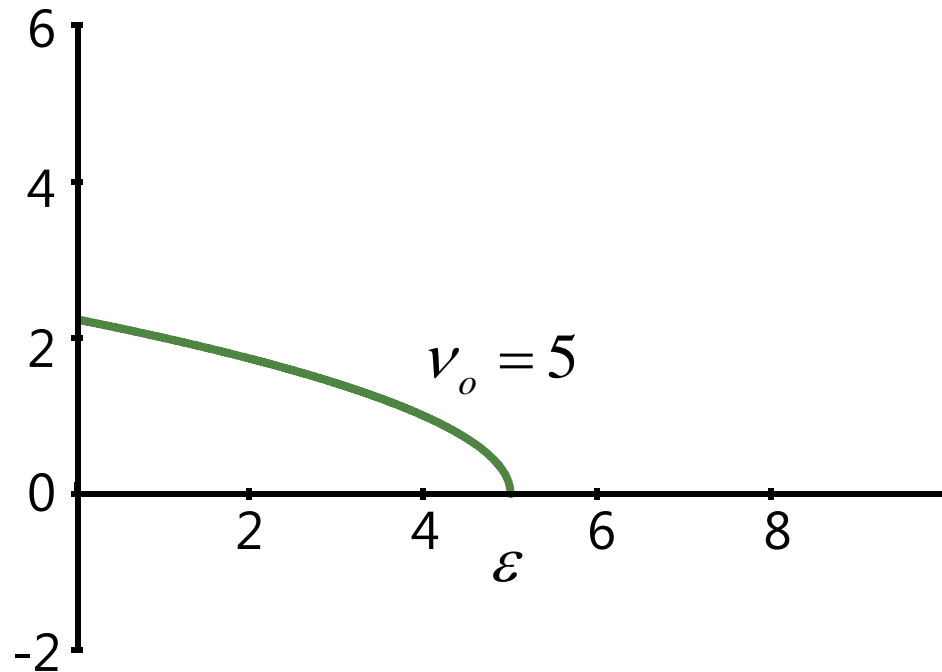


# Graphical solution

Choose a specific well  
depth  $v_o$

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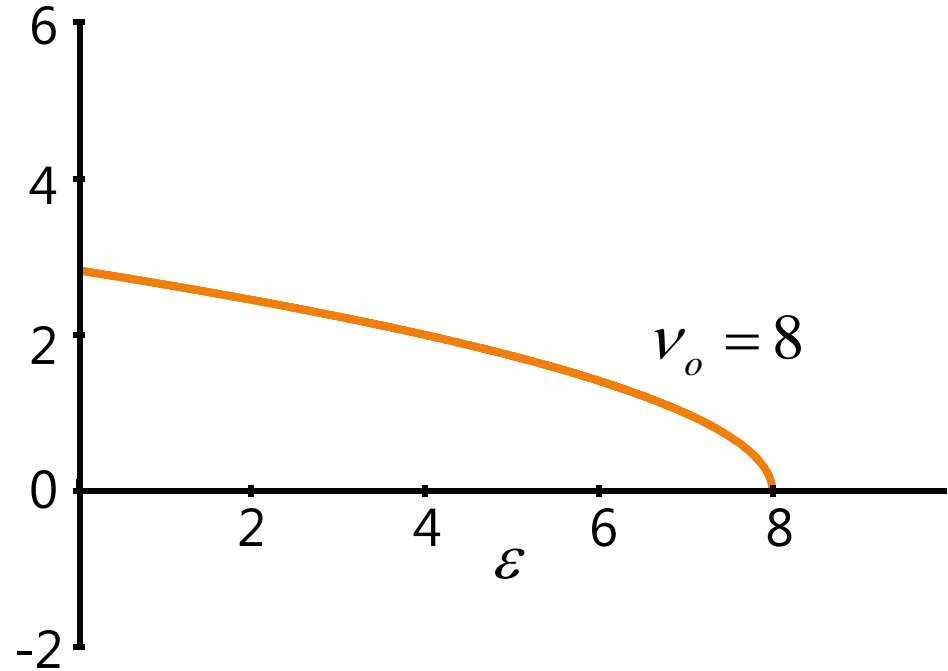


# Graphical solution

Choose a specific well  
depth  $v_o$

and plot the curve

$$\sqrt{(v_o - \varepsilon)}$$





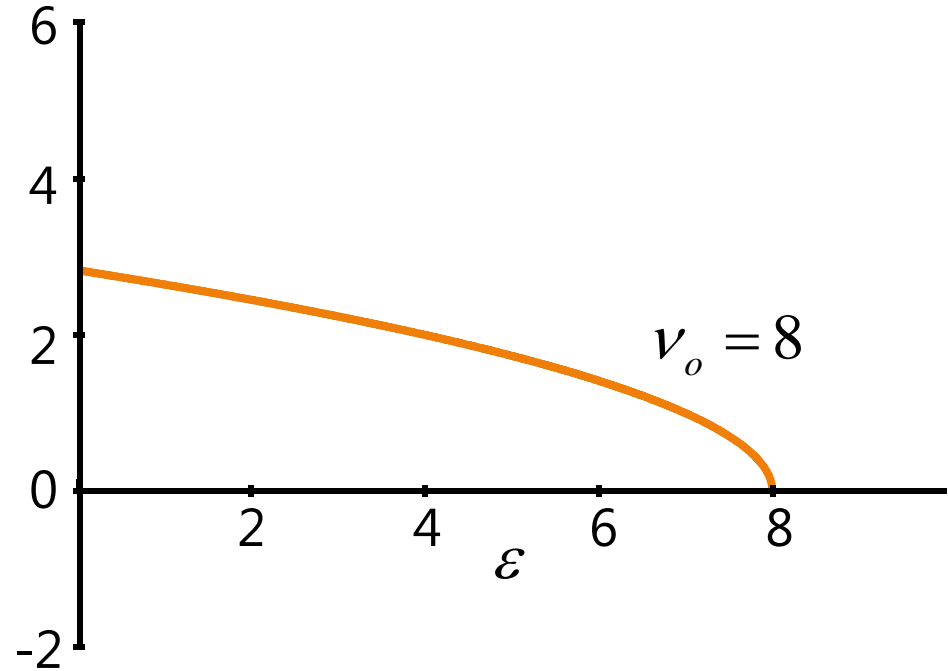
# Graphical solution

Choose a specific well  
depth  $v_o$

and plot the curve

$$\sqrt{(v_o - \varepsilon)}$$

Now add the curves



# Graphical solution

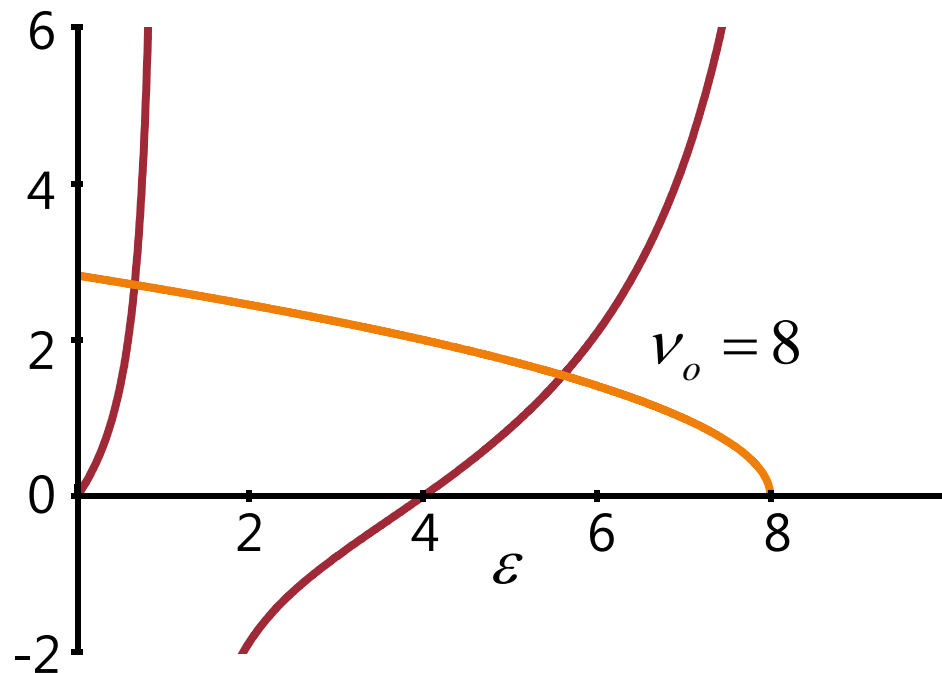
Choose a specific well  
depth  $v_o$

and plot the curve

$$\sqrt{(v_o - \varepsilon)}$$

Now add the curves

$$\sqrt{\varepsilon} \tan\left(\frac{\pi}{2}\sqrt{\varepsilon}\right)$$



# Graphical solution

Choose a specific well  
depth  $v_o$

and plot the curve

$$\sqrt{(v_o - \varepsilon)}$$

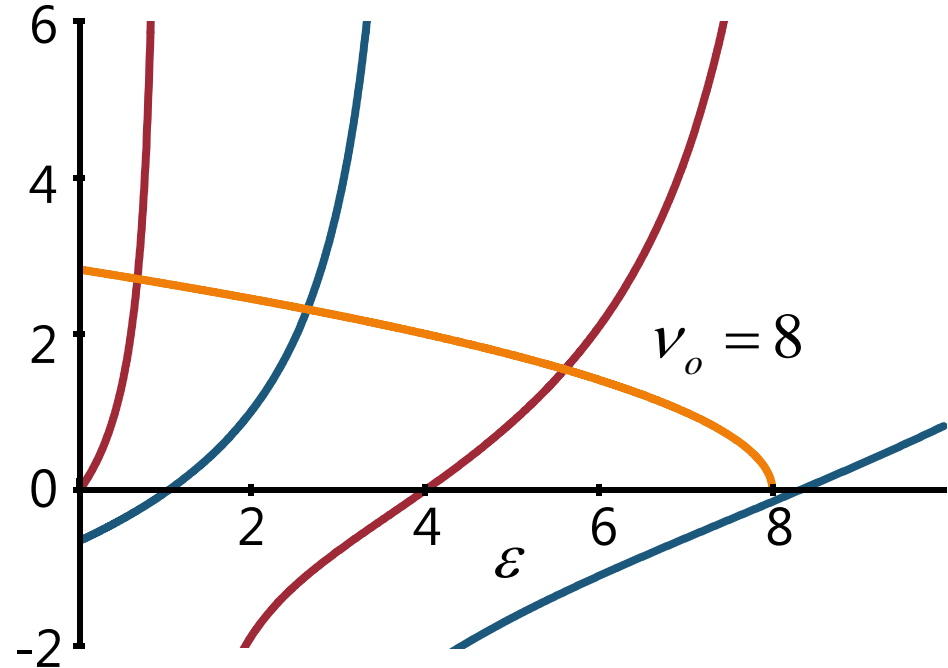
Now add the curves

$$\sqrt{\varepsilon} \tan\left(\frac{\pi}{2}\sqrt{\varepsilon}\right)$$

—

$$-\sqrt{\varepsilon} \cot\left(\frac{\pi}{2}\sqrt{\varepsilon}\right)$$

—



# Graphical solution

For a specific  $\nu_o$   
the solutions are the values  
of  $\varepsilon$  at the intersections of

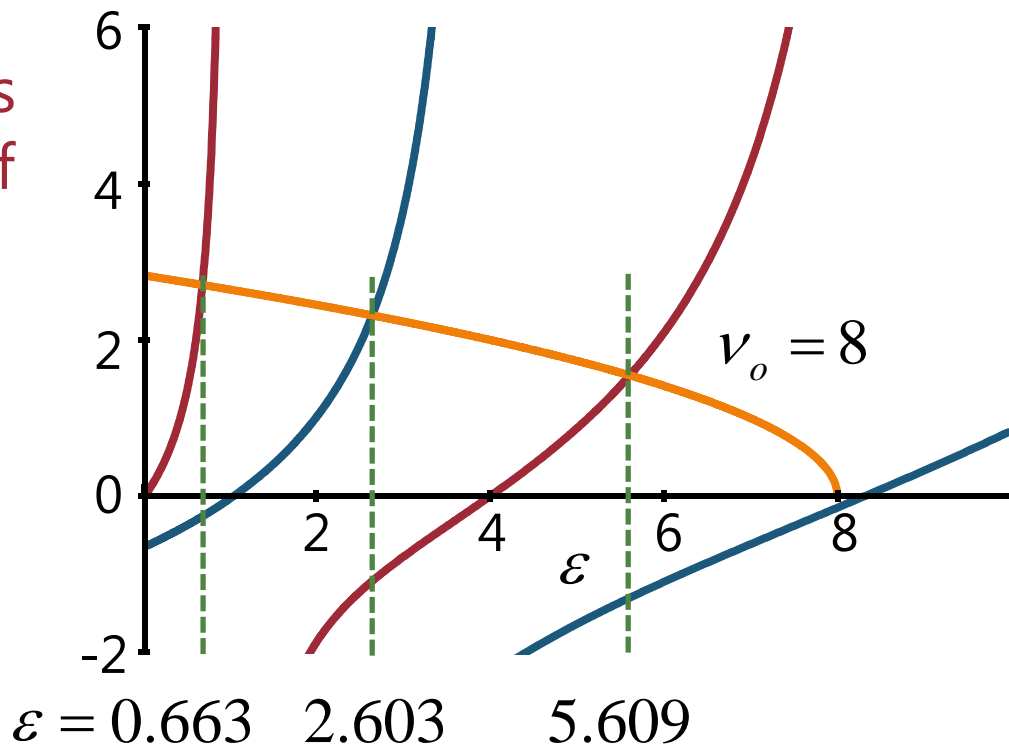
$$\sqrt{(\nu_o - \varepsilon)}$$

and

$$\sqrt{\varepsilon} \tan\left(\frac{\pi}{2}\sqrt{\varepsilon}\right)$$

or

$$-\sqrt{\varepsilon} \cot\left(\frac{\pi}{2}\sqrt{\varepsilon}\right)$$



# Solutions

These are the solutions for a well depth  $V_o$  of  $8E_1^\infty$

Note that

they are all

lower energies

than the corresponding solutions for the infinitely deep well of the same width

