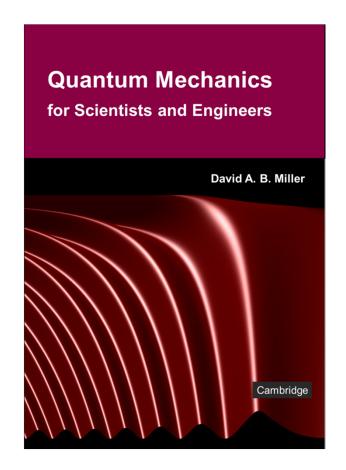
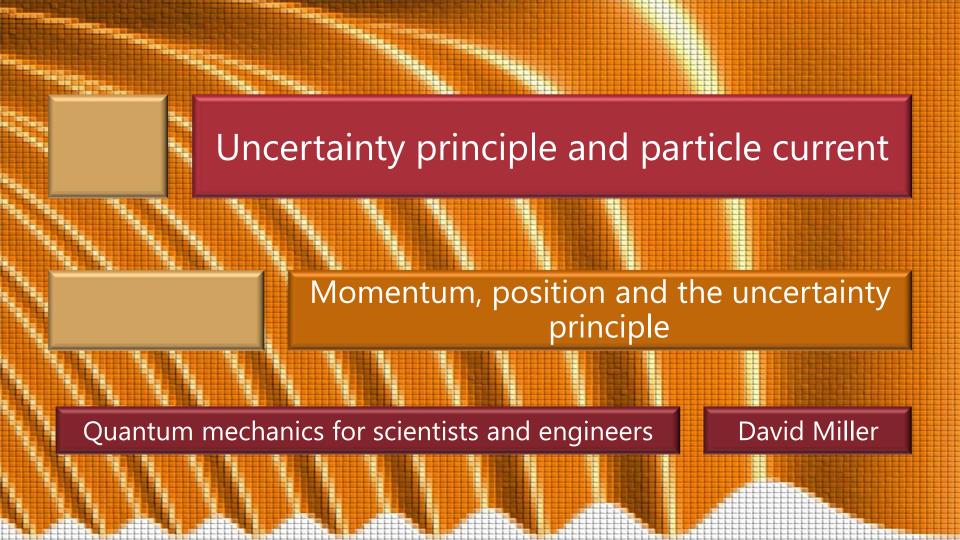
# 5.1 Uncertainty principle and particle current

Slides: Video 5.1.1 Momentum, position, and the uncertainty principle

Text reference: Quantum Mechanics for Scientists and Engineers

Sections 3.12 – 3.13





### Momentum and the momentum operator

#### For momentum

we write an operator  $\hat{p}$ 

We postulate this can be written as

$$\left[\hat{p} \equiv -i\hbar\nabla\right]$$

with

$$\nabla \equiv \mathbf{x}_o \frac{\partial}{\partial x} + \mathbf{y}_o \frac{\partial}{\partial y} + \mathbf{z}_o \frac{\partial}{\partial z}$$

where  $\mathbf{x}_{o'}$ ,  $\mathbf{y}_{o'}$  and  $\mathbf{z}_{o}$  are unit vectors in the x, y, and z directions

#### Momentum and the momentum operator

With this postulated form  $\hat{p} \equiv -i\hbar \nabla$  we find that

$$\frac{\hat{p}^2}{2m} \equiv -\frac{\hbar^2}{2m} \nabla^2$$

and we have a correspondence between the classical notion of the energy E

$$E = \frac{p^2}{2m} + V$$

and the corresponding Hamiltonian operator of the Schrödinger equation

$$\hat{H} = -\frac{\hbar^2}{2m} \nabla^2 + V = \frac{\hat{p}^2}{2m} + V$$

#### Momentum and the momentum operator

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Note that
  \hat{p} \exp(i\mathbf{k} \cdot \mathbf{r}) = -i\hbar \nabla \exp(i\mathbf{k} \cdot \mathbf{r}) = \hbar \mathbf{k} \exp(i\mathbf{k} \cdot \mathbf{r})
   This means the plane waves \exp(i\mathbf{k}\cdot\mathbf{r}) are
     the eigenfunctions of the operator \hat{p}
       with eigenvalues \hbar \mathbf{k}
We can therefore say for these eigenstates that
    the momentum is \mathbf{p} = \hbar \mathbf{k}
        Note that p is a vector, with three
         components with scalar values
            not an operator
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# Position and the position operator

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For the position operator
   the postulated operator is almost trivial
    when we are working with functions of
    position
     It is simply the position vector, \mathbf{r}, itself
At least when we are working in a
 representation that is in terms of position
  we therefore typically do not write \hat{\mathbf{r}}
     though rigorously we should
The operator for the z-component of position
 would, for example, also simply be z itself
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Here we illustrate the position-momentum uncertainty principle by example We have looked at a Gaussian wavepacket before We could write this as a sum over waves of different k-values, with Gaussian weights or we could take the limit of that process by using an integration

$$\Psi_G(z,t) \propto \int_k \exp \left[ -\left(\frac{k-\overline{k}}{2\Delta k}\right)^2 \right] \exp \left\{ -i\left[\omega(k)t - kz\right] \right\} dk$$

We could rewrite

$$\Psi_{G}(z,t) \propto \int_{k} \exp\left[-\left(\frac{k-\overline{k}}{2\Delta k}\right)^{2}\right] \exp\left\{-i\left[\omega(k)t-kz\right]\right\} dk$$

at time t = 0 as

$$\Psi(z,0) = \int_{k} \Psi_{k}(k) \exp(ikz) dk$$
 where

$$\Psi_k(k) \propto \exp \left[ -\left(\frac{k-\overline{k}}{2\Delta k}\right)^2 \right]$$

In

$$\Psi_k(k) \propto \exp\left[-\left(\frac{k-\overline{k}}{2\Delta k}\right)^2\right]$$

 $\Psi_k(k)$  is the representation of the wavefunction in k space  $\left|\Psi_k(k)\right|^2$  is the probability  $P_k$  strictly, the probability density that if we measured the momentum of the particle actually the z component of momentum it would be found to have value  $\hbar k$ 

With 
$$\Psi_k(k) \propto \exp \left[ -\left(\frac{k-\overline{k}}{2\Delta k}\right)^2 \right]$$

then this probability (density) of finding a value  $\hbar k$  for the momentum would be

$$P_k = \left|\Psi_k(k)\right|^2 \propto \exp\left[-\frac{\left(k - \overline{k}\right)^2}{2\left(\Delta k\right)^2}\right]$$

This Gaussian corresponds to the statistical Gaussian probability distribution with standard deviation  $\Delta k$ 

Note also that 
$$\Psi(z,0) = \int_k \Psi_k(k) \exp(ikz) dk$$
  
is the Fourier transform of  $\Psi_k(k)$   
and, as is well known  
the Fourier transform of a Gaussian is  
a Gaussian  
specifically here  
 $\Psi(z,0) \propto \exp\left[-\left(\Delta k\right)^2 z^2\right]$ 

If we want to rewrite

$$\left|\Psi(z,0)\right|^2 \propto \exp\left[-2\left(\Delta k\right)^2z^2\right]$$
 in the standard form

$$\left|\Psi(z,0)\right|^2 \propto \exp\left[-\frac{z^2}{2(\Delta z)^2}\right]$$

where the parameter  $\Delta z$ 

would now be the standard deviation in the probability distribution for z then  $\Delta k \Delta z = 1/2$ 

From 
$$\Delta k \Delta z = 1/2$$

if we now multiply by  $\hbar$  to get the standard deviation we would measure in momentum

we have

$$\Delta p \Delta z = \frac{\hbar}{2}$$

which is the relation between the standard deviations we would see in measurements of position and measurements of momentum

This relation

$$\Delta p \Delta z = \frac{R}{2}$$

is as good as we can get for a Gaussian

For example

a Gaussian pulse will broaden in space as it propagates

even though the range of k values remains the same

It also turns out that the Gaussian shape is the one with the minimum possible product of  $\Delta p$  and  $\Delta z$ . So quite generally

$$\Delta p \Delta z \ge \frac{\hbar}{2}$$

which is the uncertainty principle for position and momentum in one direction

#### The uncertainty principle in Fourier analysis

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Uncertainty principles are well known in Fourier analysis
   One cannot simultaneously have both
     a well defined frequency and
     a well defined time
If a signal is a short pulse
   it is necessarily made up out of a range of
    frequencies
                        \Delta \omega \Delta t \geq \frac{1}{2}
     The shorter the pulse is
        the larger the range of frequencies
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