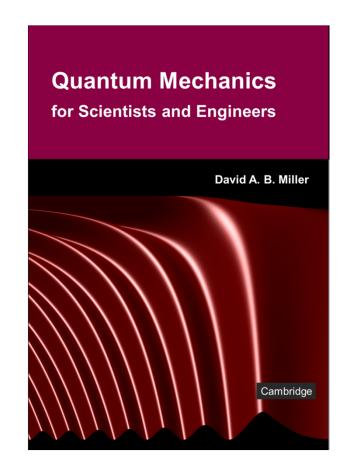
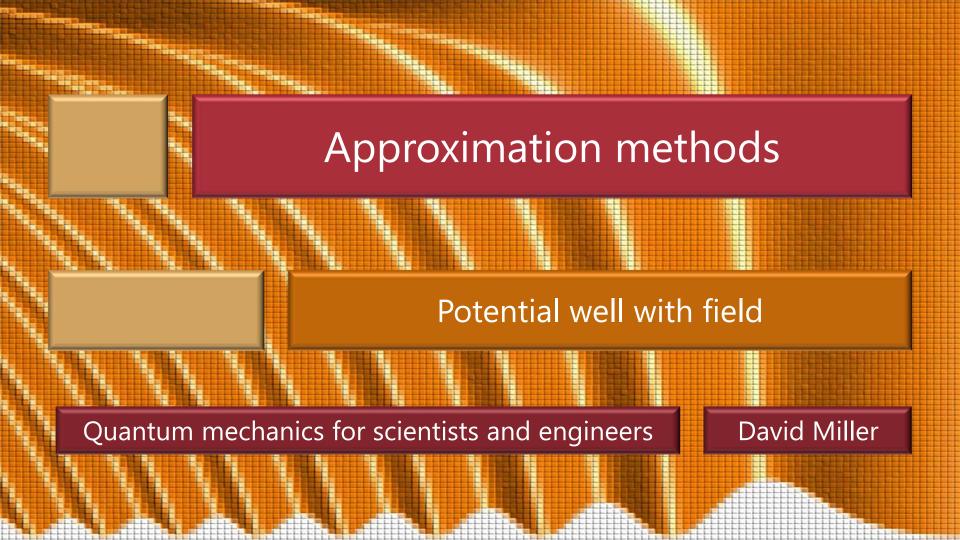
8.2 Approximation methods

Slides: Video 8.2.2 Potential well with field

Text reference: Quantum Mechanics for Scientists and Engineers

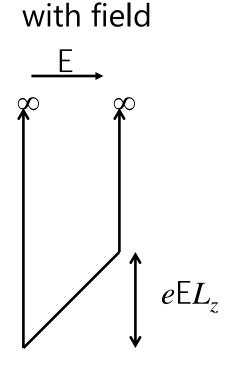
Section 6.1





Potential well with field

We are considering an electron in a potential well with infinitely high walls and with an applied electric field E without field



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The energy of an electron in an electric field E
  simply increases linearly with distance
     A positive electric field in the positive z
      direction
        pushes the electron in the negative z
         direction
          with a force of magnitude eE
So the potential energy of the electron
  increases in the positive z direction
     with the form eEz
```

We choose the potential to be zero in the middle of the well

Hence, within the well

the potential energy is

$$V(z) = e E(z - L_z / 2)$$

and the Hamiltonian becomes

$$\hat{H} = -\frac{\hbar^2}{2m} \frac{d^2}{dz^2} + e E(z - L_z/2)$$

We can usefully define dimensionless units

A convenient unit of energy is

the confinement energy of the first state of the original infinitely deep well

$$E_1^{\infty} = \frac{\hbar^2}{2m} \left(\frac{\pi}{L_z}\right)^2$$

and in those units the eigenenergy of the *n*th state will be

$$\eta_n = \frac{E_n}{E_1^{\infty}}$$

A convenient unit of field E_o gives one energy unit of potential change from one side of the well to the other

$$\mathsf{E}_o = \frac{E_1^{\infty}}{eL_z}$$

So, the (dimensionless) field will be $f = E/E_a$

A convenient distance unit is the thickness L_z so the dimensionless distance will be

$$\xi = z / L_{z}$$

From the original Hamiltonian

$$\hat{H} = -\frac{\hbar^2}{2m} \frac{d^2}{dz^2} + e E(z - L_z / 2)$$

dividing by E_1^{∞} and using dimensionless units gives

$$\hat{H} = -\frac{1}{\pi^2} \frac{d^2}{d\xi^2} + f(\xi - 1/2)$$

and a time-independent Schrödinger equation

$$\hat{H}\phi(\xi) = \eta\phi(\xi)$$

For the "unperturbed" problem without field we write the "unperturbed" Hamiltonian within the well as

$$\hat{H}_o = -\frac{1}{\pi^2} \frac{d^2}{d\xi^2}$$

The normalized solutions of the corresponding Schrödinger equation

$$\hat{H}_{\scriptscriptstyle o} \psi_{\scriptscriptstyle n} = \varepsilon_{\scriptscriptstyle n} \psi_{\scriptscriptstyle n}$$
 are then

$$\psi_n(\xi) = \sqrt{2}\sin(n\pi\xi)$$

