

## 5.1 Uncertainty principle and particle current

Slides: Video 5.1.1 Momentum, position, and the uncertainty principle

Text reference: Quantum Mechanics for Scientists and Engineers

Sections 3.12 – 3.13





# Uncertainty principle and particle current



Momentum, position and the uncertainty principle

Quantum mechanics for scientists and engineers

David Miller

# Momentum and the momentum operator

For momentum

we write an operator  $\hat{p}$

We postulate this can be written as

$$\hat{p} \equiv -i\hbar\nabla$$

with

$$\nabla \equiv \mathbf{x}_o \frac{\partial}{\partial x} + \mathbf{y}_o \frac{\partial}{\partial y} + \mathbf{z}_o \frac{\partial}{\partial z}$$

where  $\mathbf{x}_o$ ,  $\mathbf{y}_o$ , and  $\mathbf{z}_o$  are unit vectors  
in the  $x$ ,  $y$ , and  $z$  directions

# Momentum and the momentum operator

With this postulated form  $\hat{p} \equiv -i\hbar\nabla$  we find that

$$\frac{\hat{p}^2}{2m} \equiv -\frac{\hbar^2}{2m} \nabla^2$$

and we have a correspondence between the classical notion of the energy  $E$

$$E = \frac{p^2}{2m} + V$$

and the corresponding Hamiltonian operator of the Schrödinger equation

$$\hat{H} = -\frac{\hbar^2}{2m} \nabla^2 + V = \frac{\hat{p}^2}{2m} + V$$

# Momentum and the momentum operator

Note that

$$\hat{p} \exp(i\mathbf{k} \cdot \mathbf{r}) = -i\hbar \nabla \exp(i\mathbf{k} \cdot \mathbf{r}) = \hbar\mathbf{k} \exp(i\mathbf{k} \cdot \mathbf{r})$$

This means the plane waves  $\exp(i\mathbf{k} \cdot \mathbf{r})$  are the eigenfunctions of the operator  $\hat{p}$

with eigenvalues  $\hbar\mathbf{k}$

We can therefore say for these eigenstates that

the momentum is  $\mathbf{p} = \hbar\mathbf{k}$

Note that  $\mathbf{p}$  is a vector, with three components with scalar values  
not an operator

# Position and the position operator

For the position operator

the postulated operator is almost trivial  
when we are working with functions of  
position

It is simply the position vector,  $\mathbf{r}$ , itself

At least when we are working in a  
representation that is in terms of position

we therefore typically do not write  $\hat{\mathbf{r}}$

though rigorously we should

The operator for the  $z$ -component of position  
would, for example, also simply be  $z$  itself

# The uncertainty principle

Here we illustrate the position-momentum uncertainty principle by example

We have looked at a Gaussian wavepacket before

We could write this as a sum over waves of different  $k$ -values, with Gaussian weights

or we could take the limit of that process by using an integration

$$\Psi_G(z, t) \propto \int_k \exp \left[ - \left( \frac{k - \bar{k}}{2\Delta k} \right)^2 \right] \exp \left\{ -i \left[ \omega(k)t - kz \right] \right\} dk$$

# The uncertainty principle

We could rewrite

$$\Psi_G(z, t) \propto \int_k \exp\left[-\left(\frac{k - \bar{k}}{2\Delta k}\right)^2\right] \exp\left\{-i[\omega(k)t - kz]\right\} dk$$

at time  $t = 0$  as

$$\Psi(z, 0) = \int_k \Psi_k(k) \exp(ikz) dk$$

where

$$\Psi_k(k) \propto \exp\left[-\left(\frac{k - \bar{k}}{2\Delta k}\right)^2\right]$$



# The uncertainty principle

In

$$\Psi_k(k) \propto \exp\left[-\left(\frac{k - \bar{k}}{2\Delta k}\right)^2\right]$$

$\Psi_k(k)$  is the representation of the wavefunction in  $k$  space

$|\Psi_k(k)|^2$  is the probability  $P_k$

strictly, the probability density

that if we measured the momentum of the particle

actually the  $z$  component of momentum

it would be found to have value  $\hbar k$

# The uncertainty principle

With

$$\Psi_k(k) \propto \exp\left[-\left(\frac{k - \bar{k}}{2\Delta k}\right)^2\right]$$

then this probability (density) of finding a value  $\hbar k$  for the momentum would be

$$P_k = |\Psi_k(k)|^2 \propto \exp\left[-\frac{(k - \bar{k})^2}{2(\Delta k)^2}\right]$$

This Gaussian corresponds to the statistical Gaussian probability distribution  
with standard deviation  $\Delta k$

# The uncertainty principle

Note also that  $\Psi(z, 0) = \int_k \Psi_k(k) \exp(ikz) dk$

is the Fourier transform of  $\Psi_k(k)$

and, as is well known

the Fourier transform of a Gaussian is  
a Gaussian

specifically here

$$\Psi(z, 0) \propto \exp\left[-(\Delta k)^2 z^2\right]$$

# The uncertainty principle

If we want to rewrite

$$|\Psi(z, 0)|^2 \propto \exp\left[-2(\Delta k)^2 z^2\right]$$

in the standard form

$$|\Psi(z, 0)|^2 \propto \exp\left[-\frac{z^2}{2(\Delta z)^2}\right]$$

where the parameter  $\Delta z$

would now be the standard deviation  
in the probability distribution for  $z$

then  $\Delta k \Delta z = 1/2$

# The uncertainty principle

From  $\Delta k \Delta z = 1/2$

if we now multiply by  $\hbar$  to get the standard deviation  
we would measure in momentum

we have

$$\Delta p \Delta z = \frac{\hbar}{2}$$

which is the relation between the standard  
deviations we would see in

measurements of position and  
measurements of momentum

# The uncertainty principle

This relation

$$\Delta p \Delta z = \frac{\hbar}{2}$$

is as good as we can get for a Gaussian

For example

a Gaussian pulse will broaden in space as it propagates

even though the range of  $k$  values remains the same

# The uncertainty principle

It also turns out that the Gaussian shape  
is the one with the minimum possible  
product of  $\Delta p$  and  $\Delta z$

So quite generally

$$\Delta p \Delta z \geq \frac{\hbar}{2}$$

which is the uncertainty principle  
for position and momentum in  
one direction

# The uncertainty principle in Fourier analysis

Uncertainty principles are well known in Fourier analysis

One cannot simultaneously have both  
a well defined frequency and  
a well defined time

If a signal is a short pulse

it is necessarily made up out of a range of  
frequencies

$$\Delta\omega\Delta t \geq \frac{1}{2}$$

The shorter the pulse is  
the larger the range of frequencies



