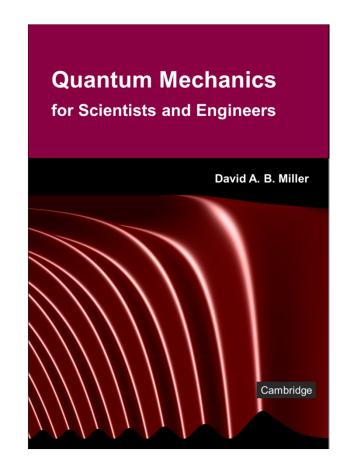
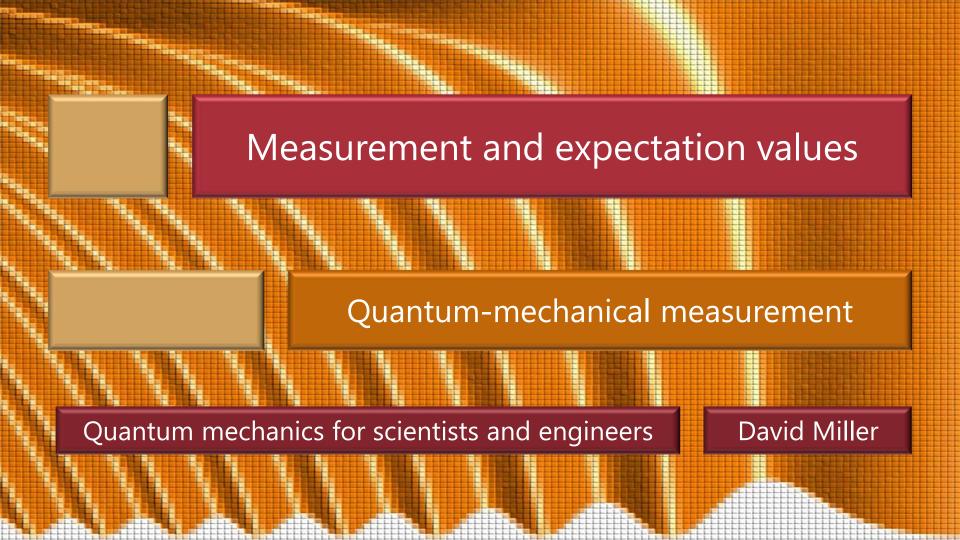
# 4.3 Measurement and expectation values

Slides: Video 4.3.1 Quantummechanical measurement

Text reference: Quantum Mechanics for Scientists and Engineers

Section 3.8





#### Probabilities and expansion coefficients

Suppose we take some (normalized) quantum mechanical wave function  $\Psi(\mathbf{r},t)$ and expand it in some complete orthonormal set of spatial functions  $\psi_n(\mathbf{r})$ At least if we allow the expansion coefficients  $c_n$  to vary in time we know we can always do this

$$\Psi(\mathbf{r},t) = \sum c_n(t)\psi_n(\mathbf{r})$$

#### Probabilities and expansion coefficients

Then the fact that  $\Psi(\mathbf{r},t)$  is normalized means we know the answer for the normalization integral

$$\int_{-\infty}^{\infty} \left| \Psi(\mathbf{r}, t) \right|^{2} d^{3}\mathbf{r} = \int_{-\infty}^{\infty} \left[ \sum_{n} c_{n}^{*}(t) \psi_{n}^{*}(\mathbf{r}) \right] \times \left[ \sum_{m} c_{m}(t) \psi_{m}(\mathbf{r}) \right] d^{3}\mathbf{r} = 1$$

Because of the orthogonality of the basis functions only terms with n = m survive the integration

Because of the orthonormality of the basis functions the result from any such term will simply be  $\left|c_n(t)\right|^2$ 

Hence we have 
$$\sum |c_n|^2 = 1$$

#### Measurement postulate

On measurement of a state the system collapses into the nth eigenstate of the quantity being measured with probability  $P_n = \left| c_n \right|^2$ 

In the expansion of the state in the eigenfunctions of the quantity being measured  $c_n$  is the expansion coefficient of the nth eigenfunction

## Expectation value of the energy

Suppose do an experiment to measure the energy E of some quantum mechanical system

We could repeat the experiment many times and get a statistical distribution of results

Given the probabilities  $P_n$  of getting a specific energy eigenstate, with energy  $E_n$ 

we would get an average answer

$$\langle E \rangle = \sum_{n} E_{n} P_{n} = \sum_{n} E_{n} \left| c_{n} \right|^{2}$$

where we call this average  $\langle E \rangle$  the "**expectation value**"

#### Energy expectation value example

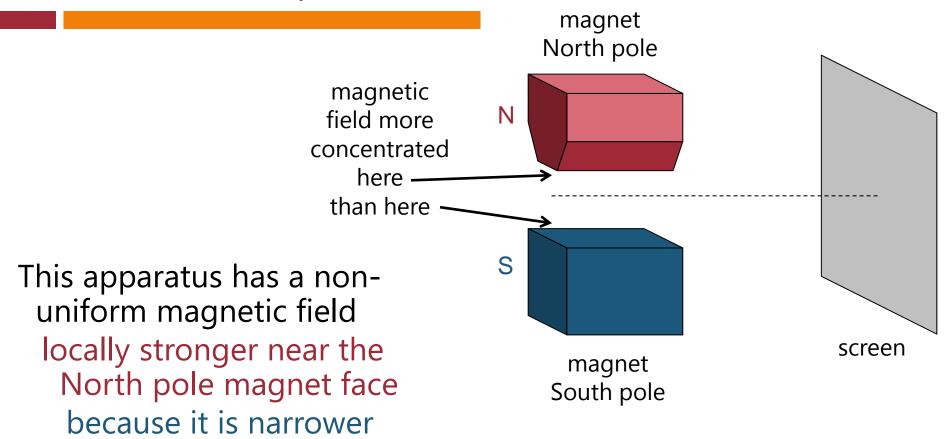
For example, for the coherent state discussed above with parameter N, we have

$$\langle E \rangle = \sum_{n=0}^{\infty} E_n \frac{N^n \exp(-N)}{n!} = \hbar \omega \left[ \sum_{n=0}^{\infty} n \frac{N^n \exp(-N)}{n!} \right] + \frac{1}{2} \hbar \omega = \left( N + \frac{1}{2} \right) \hbar \omega$$

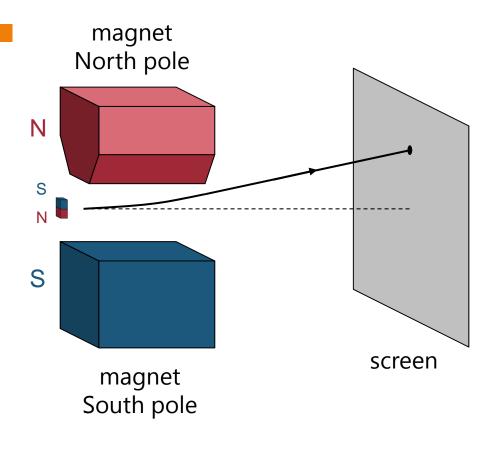
where we use the result that the average in a Poisson statistical distribution is just the parameter *N* 

Note that N does not have to be an integer This is an expectation value, not an eigenvalue

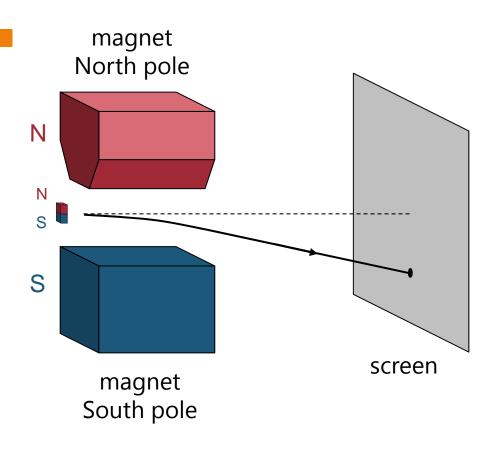
We can have states with any expectation value we want



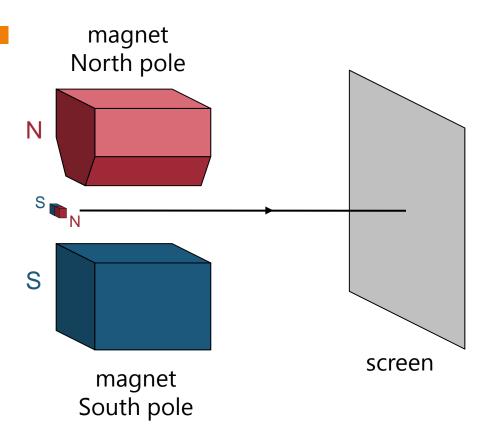
Imagine firing some small magnets initially along the dashed line Because the magnetic field is non-uniform stronger near the North pole than near the South pole a vertical magnet will be deflected up



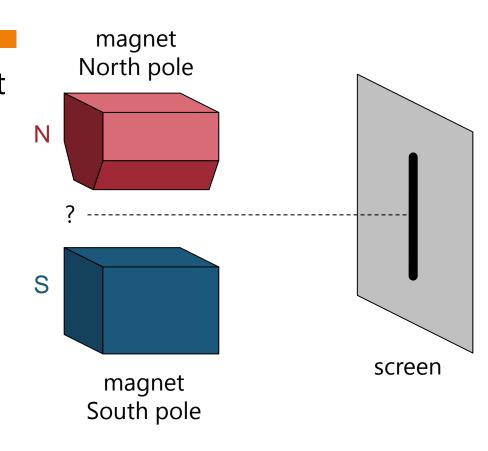
Imagine firing some small magnets initially along the dashed line Because the magnetic field is non-uniform stronger near the North pole than near the South pole a vertical magnet will be deflected up or down



A horizontally-oriented magnet will not be deflected



A horizontally-oriented magnet will not be deflected and magnets of other orientations should be deflected by intermediate amounts After "firing" many randomly oriented magnets we should end up with a line on the screen



### Electrons and the Stern-Gerlach experiment

```
Electrons have a quantum mechanical property
 called spin
  It gives them a "magnetic moment"
    just like a small magnet
What will happen if we fire electrons
  with no particular "orientation" of their spin
     into the Stern-Gerlach apparatus?
       We might expect the "line" on the screen
(Note: the actual experiment used silver atoms,
 which behave the same as electrons in this case)
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With electrons we get two dots! "Explanation" We are measuring the vertical component of the spin There are two eigenstates of this component up and down so we have collapse to the eigenstates

