

4.2 Wavepackets

Slides: Video 4.2.6 Electron wavepackets

Text reference: Quantum Mechanics for Scientists and Engineers

Section 3.7 ("Examples of motion of wavepackets")





Wavepackets



Electron wavepackets

Quantum mechanics for scientists and engineers

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Constructing a wavepacket

We can construct a “wavepacket” by putting together a linear superposition of energy eigensolutions

For a free electron or a similar particle of mass m
the individual eigensolutions are plane waves

For propagation in the z direction, these are of the form

$$\Psi_k(z, t) \propto \exp \left\{ -i \left[\frac{E(k)}{\hbar} t - kz \right] \right\} = \exp \left\{ -i [\omega(k)t - kz] \right\}$$

for some chosen value of k , and hence of

$$\text{energy } E(k) = \frac{\hbar^2 k^2}{2m} \text{ and frequency } \omega(k) = \frac{E(k)}{\hbar}$$

Constructing a wavepacket

Since

a linear superposition of such plane wave solutions of the time-dependent Schrödinger equation is also a solution

we can have a “wavepacket”

$$\Psi_{WP}(z, t) \propto \sum_k a_k \Psi_k(z, t) = \sum_k a_k \exp\{-i[\omega(k)t - kz]\}$$

for some set of values of k in our sum

and some chosen amplitudes a_k for each such plane wave

Gaussian wavepacket

One convenient and useful set of k values and amplitudes a_k to choose is

a set of equally spaced k values

with Gaussian amplitudes or “weights” for a_k

$$\Psi_G(z, t) \propto \sum_k \exp\left[-\left(\frac{k - \bar{k}}{2\Delta k}\right)^2\right] \exp\{-i[\omega(k)t - kz]\}$$

Here \bar{k} is the center of the distribution of k values

Δk is a width parameter for the Gaussian function

Note this gives a “pulse” that is also Gaussian in space

Gaussian wavepacket

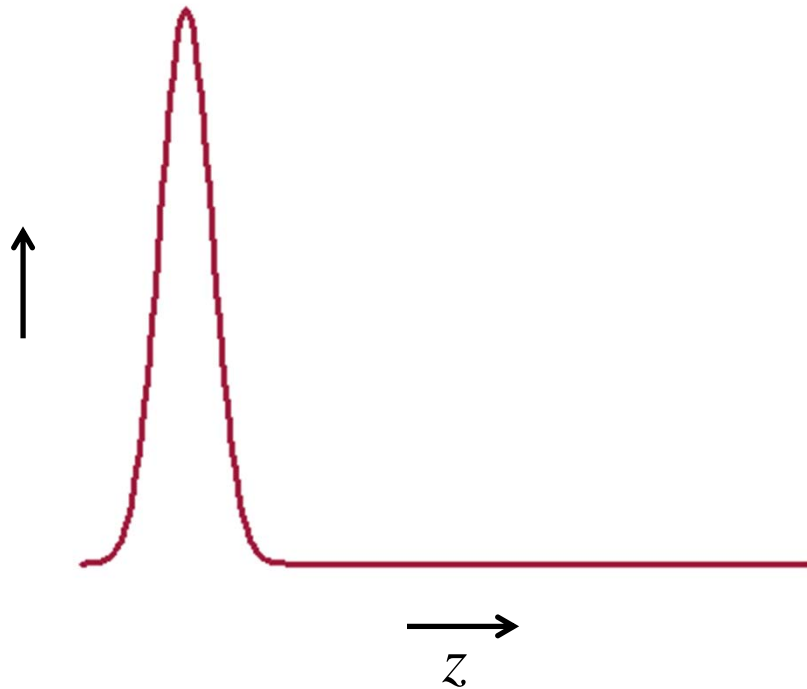
Our Gaussian wavepacket is
also Gaussian in space

As we let time evolve

simply adding up the terms
in our wavepacket sum at
each time

we can see the
wavepacket propagate
moving to the right
and getting wider

Probability density



Gaussian wavepacket

A wavepacket that increases in width as it propagates

is said to be “dispersing”

It gets wider because

the change in ω with k

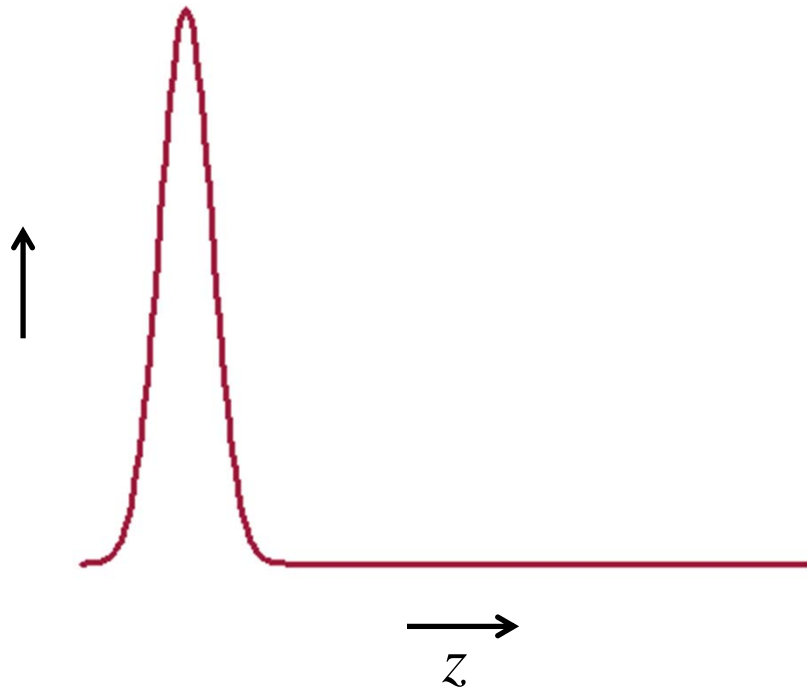
is not even linear

(here it is quadratic)

an effect known as

group velocity
dispersion

Probability density



Wavepacket at a barrier

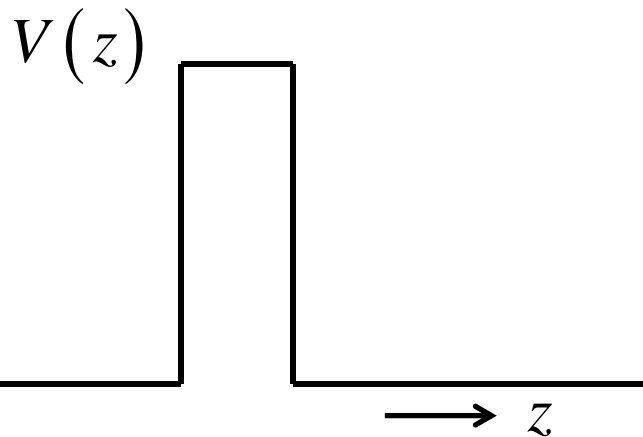
Suppose we want to understand
a wavepacket hitting a barrier
from the left

We proceed in the same way

but now we use

a superposition of the
energy eigenfunctions $\Psi_{Bk}(z, t)$
of the Schrödinger
equation

with a barrier potential $V(z)$



Eigensolutions with a barrier

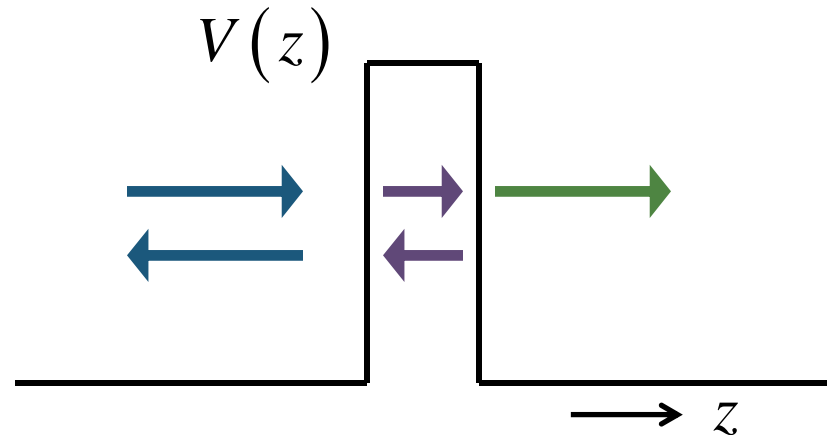
The energy eigensolutions $\Psi_{Bk}(z, t)$

for a particle incident from
the left are

forward and backward
propagating waves on the
left

“forward” and “backward”
exponentials inside the
barrier

forward propagating
waves on the right

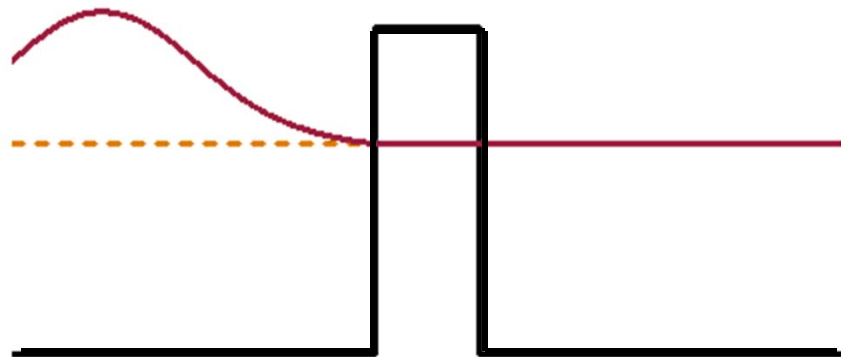


Wavepacket at a barrier

Constructing a Gaussian-weighted linear superposition of solutions

with equally-spaced k -values
on the left

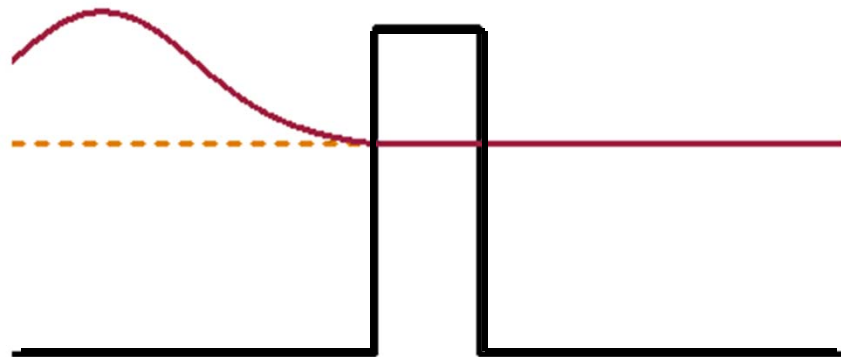
here centered round the k -value corresponding to the dashed orange line
gives an approximately Gaussian “pulse” on the left to start with



Wavepacket at a barrier

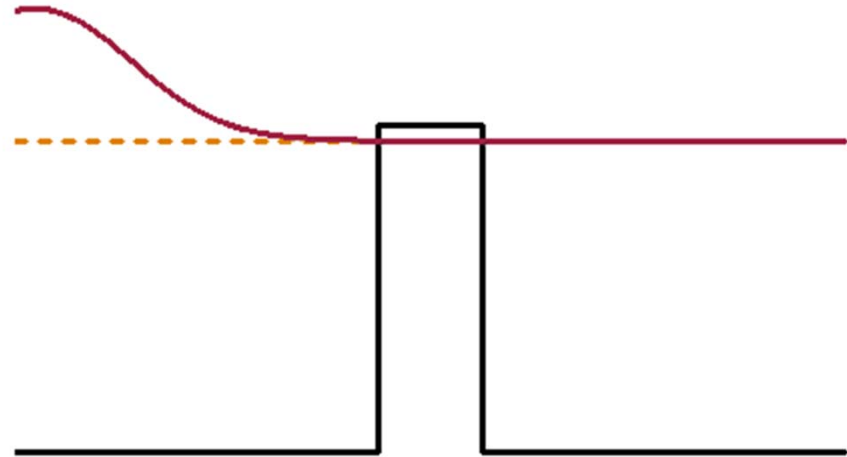
Now as time evolves, the
“pulse” moves to the right
where it partially bounces off
and partially transmits

Note that there is a significant
probability of
finding the particle in the
barrier
while the pulse is “hitting” it



Wavepacket at a barrier

With the particle incident at a
higher average energy
the transmitted pulse is
stronger
and the reflected pulse is
weaker



Wavepacket at a barrier

Even with energies above the barrier

there is still significant
reflection of the pulse

wave are generally reflected
off of any changes in the
potential

