

# Math 241: Laplace equation in polar coordinates; consequences and properties

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# Laplace on a disk

- Next up is to solve the Laplace equation on a disk with boundary values prescribed on the circle that bounds the disk.
- We'll use polar coordinates for this, so a typical problem might be:

$$\nabla^2 u = \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial u}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 u}{\partial \theta^2} = 0$$

on the disk of radius  $R = 3$  centered at the origin, with boundary condition

$$u(3, \theta) = \begin{cases} 1 & 0 \leq \theta \leq \pi \\ \sin^2 \theta & \pi < \theta < 2\pi \end{cases}$$

# Separation of variables

- We search for separated solutions:  $u(r, \theta) = R(r)\Theta(\theta)$ . So

$$\frac{1}{r}(rR')'\Theta + \frac{1}{r^2}R\Theta'' = 0$$

or

$$\frac{r^2 R'' + rR'}{R} = -\frac{\Theta''}{\Theta} = -\lambda.$$

- We need  $\Theta$  to be periodic with period  $2\pi$  (so that  $u$  will be well-defined as a function of  $x$  and  $y$ ) — so  $\lambda = n^2$  ( $n = 0, 1, 2, \dots$ ) and

$$\Theta = \begin{cases} a_0 & n = 0 \\ a_n \cos(n\theta) + b_n \sin(n\theta) & n = 1, 2, \dots \end{cases}$$

# Separated solutions

- The  $R$  equation becomes  $r^2 R'' + rR' + n^2 R = 0$ , for  $n = 0, 1, 2, \dots$
- This is a Cauchy-Euler equation (look in your Math 240 book) and the solution is

$$R = \begin{cases} c_1 + c_2 \ln r & n = 0 \\ c_1 r^n + c_2 r^{-n} & n = 1, 2, \dots \end{cases}$$

- Because we don't want the solution to go to infinity at the center of the disk (where  $r = 0$ ), we set  $c_2 = 0$  in both cases. So our separated solutions are:

$$u(r, \theta) = R(r)\Theta(\theta) = \begin{cases} a_0 & n = 0 \\ a_n r^n \cos(n\theta) + b_n r^n \sin(n\theta) & n = 1, 2, \dots \end{cases}$$

# Fourier series

- As usual, we'll make a series out of our separated solutions and try to match the boundary condition;

$$u(r, \theta) = a_0 + \sum_{n=1}^{\infty} a_n r^n \cos(n\theta) + b_n r^n \sin(n\theta)$$

- Each term in the series satisfies  $\nabla^2 u = 0$  and is well-defined on the disk. Now we need to match the boundary condition for  $r = 3$  (from the first slide):

$$\begin{aligned} u(3, \theta) &= a_0 + \sum_{n=1}^{\infty} 3^n (a_n \cos(n\theta) + b_n \sin(n\theta)) \\ &= \begin{cases} 1 & 0 \leq \theta \leq \pi \\ \sin^2 \theta & \pi < \theta < 2\pi \end{cases} \end{aligned}$$

# Fourier coefficients

- We calculate the Fourier coefficients as usual:

$$a_0 = \frac{1}{2\pi} \left( \int_0^\pi 1 \, d\theta + \int_\pi^{2\pi} \sin^2 \theta \, d\theta \right) = \frac{3}{4}$$

and

$$\begin{aligned} a_n &= \frac{1}{3^n \pi} \left( \int_0^\pi \cos(n\theta) \, d\theta + \int_\pi^{2\pi} \sin^2 \theta \cos(n\theta) \, d\theta \right) \\ &= \begin{cases} -\frac{1}{36} & n = 2 \\ 0 & n \neq 2, \, n > 0 \end{cases} \end{aligned}$$

and

$$\begin{aligned} b_n &= \frac{1}{3^n \pi} \left( \int_0^\pi \sin(n\theta) \, d\theta + \int_\pi^{2\pi} \sin^2 \theta \sin(n\theta) \, d\theta \right) \\ &= \begin{cases} 0 & n \text{ even} \\ \frac{1}{3^n \pi} \left( \frac{2}{n} + \frac{4}{n(n^2 - 4)} \right) & n \text{ odd} \end{cases} \end{aligned}$$

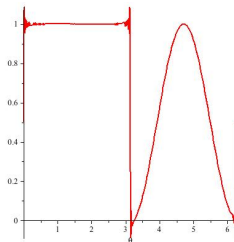
# The solution

- So the solution is

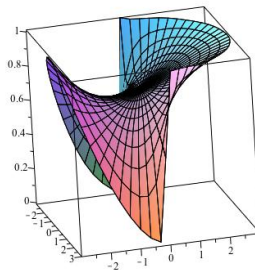
$$u(r, \theta) = \frac{3}{4} - \frac{r^2}{36} \cos(2\theta) + \sum_{n=0}^{\infty} \frac{r^{2n+1}}{3^{2n+1}\pi} \frac{2(4n^2 + 4n - 1)}{(2n + 3)(4n^2 - 1)} \sin(2n + 1)\theta.$$

# Pictures

- Boundary values, 100 terms



- Graph of  $u(r, \theta)$ , 200 terms:





# Generally

The solution of the Dirichlet problem  $\nabla^2 u = 0$  on a disk of radius  $R$  with boundary condition  $u = f(\theta)$  on the boundary of the disk, is

$$u(r, \theta) = a_0 + \sum_{n=1}^{\infty} a_n r^n \cos(n\theta) + b_n r^n \sin(n\theta)$$

where  $a_0 = \frac{1}{2\pi} \int_0^{2\pi} f(\theta) d\theta$ , and for  $n > 0$ ,

$$a_n = \frac{1}{\pi} \int_0^{2\pi} f(\theta) \cos(n\theta) d\theta, \quad b_n = \frac{1}{\pi} \int_0^{2\pi} f(\theta) \sin(n\theta) d\theta.$$

From

$$u(r, \theta) = a_0 + \sum_{n=1}^{\infty} a_n r^n \cos(n\theta) + b_n r^n \sin(n\theta)$$

we see that when  $r = 0$ , we get that

$$u(0, \theta) = a_0 = \frac{1}{2\pi} \int_0^{2\pi} f(\theta) d\theta.$$

$u(0, \theta)$  is the value of  $u$  at the center of the circle, and the last expression is the average value of  $f$  (that is to say, the average value of  $u$ ) on the circle of radius  $R$ .

This will be true for any disk contained in the domain where  $u$  is harmonic (i.e., where  $\nabla^2 u = 0$ ).

# Mean-value property

The reasoning on the last slide proves:

## Theorem (Mean-value property for harmonic functions)

*Let  $u$  be a harmonic function ( $\nabla^2 u = 0$  in a region that contains a disk  $D$  whose center is the point  $p$ , whose radius is  $r$  and whose circumference is the circle  $C$ ). Then*

$$u(p) = \frac{1}{2\pi} \int_0^{2\pi} u(p + re^{i\theta}) d\theta = \frac{1}{2\pi r} \int_C u ds.$$

*In other words, the value of  $u$  at the center of the disk  $D$  is equal to the average of the values of  $u$  on the circumference of the disk.*

# Maximum principle

Since the average of a set of numbers cannot exceed the maximum of the numbers, we have

## Theorem (Maximum principle for harmonic functions)

*If  $u$  is harmonic on a region  $R$ , then the maximum value of  $u$  must occur at a boundary point of  $R$ . Moreover, if the maximum point also occurs at an interior point of  $R$ , then  $u$  must be a constant function.*

Clearly, there is also a minimum principle for harmonic functions. So if we know that the maximum and minimum of a harmonic function on a domain cannot be exceeded the maximum on the boundary or be less than the minimum on the boundary.

If we multiply the equality in the mean-value property by  $r dr$  and integrate from 0 to  $R$ , we get

$$\int_0^R u(p) r dr = \frac{1}{2\pi} \int_0^R \int_0^{2\pi} u(r + e^{i\theta}) r dr d\theta$$

Since  $u(p)$  is a constant the left side equals  $u(p)R^2/2$ .

Recognizing the double integral on the right as the integral of  $u dA$  over the entire disk  $D$  of radius  $R$  centered at  $p$  gives:

## Theorem (Solid mean value property for harmonic functions)

*If  $u$  is harmonic on a disk  $D$  with center  $p$  and radius  $R$ , then the value of  $u$  at the center of the disk is equal to the average of  $u$  over the disk:*

$$u(p) = \frac{1}{\pi R^2} \iint_D u dA.$$

We can use the solid MVP to prove a remarkable fact about harmonic functions. But first we need a little geometric fact:

## Lemma

*Let  $x$  and  $y$  be points in the plane, let  $D_x(R)$  be the disk of radius  $R$  centered at  $x$ , and let  $D_y(R)$  be the disk of radius  $R$  centered at  $y$ . Then the area of the set of points contained in **one but not both** of  $D_x(R)$  or  $D_y(R)$  is less than  $2\pi R|x - y|$ .*

To see this, note that if  $q$  is in  $D_x(R)$  but not  $D_y(R)$ , then  $|q - x| < R$  and  $|q - y| > R$ , so  $R - |x - y| < |q - x| < R$  (for the left inequality, note that

$$|q - x| = |(q - y) - (x - y)| > |q - y| - |x - y| > R - |x - y|.$$

And the ring of points at distance between  $R - |x - y|$  and  $R$  from  $x$  has area less than  $2\pi R|x - y|$ .

# Liouville's theorem

Now for our remarkable fact:

## Theorem (Liouville's theorem for harmonic functions)

*Suppose  $u$  is a harmonic function defined on the entire plane. If  $u$  is bounded (that is, if there is a number  $M$  so that  $-M < u < M$  everywhere in the plane, then  $u$  is a constant function.*

To prove this, let  $x$  and  $y$  be any points in the plane. Then

$$u(x) - u(y) = \frac{1}{\pi R^2} \left( \iint_{D_x(R)} u \, dA - \iint_{D_y(R)} u \, dA \right)$$

For large  $R$ , a lot of the two integrals cancel because the disks overlap. And because  $u$  is bounded we can conclude

$$|u(x) - u(y)| \leq \frac{1}{\pi R^2} M(\text{area}(S))$$

where  $S$  is the set of points contained in one but not both of  $D_x(R)$  or  $D_y(R)$ .

# End of proof of Liouville's theorem

From the geometric lemma, we conclude that

$$|u(x) - u(y)| \leq \frac{1}{\pi R^2} (2\pi R |x - y|)$$

Since this last quantity goes to zero as  $R \rightarrow \infty$ , we conclude that  $u(x) = u(y)$ . And since  $x$  and  $y$  were arbitrary,  $u$  must be constant. This concludes the proof of Liouville's theorem.