

# Homework #6: Calculus of Variations

## Solutions

March 8, 2019

### 1 Derive the Euler-Lagrange Equation from Directional Derivative

Given an energy functional  $E(f)$  of the form

$$E(f) = \int_0^1 L(f(x), \dot{f}(x), x) dx$$

we showed in class that a necessary condition for a function  $f(x)$  to be a minimizer of  $E$  is that it satisfies the Euler-Lagrange equation given by

$$L_f - \frac{d}{dx} L_{\dot{f}} = 0.$$

We derived this condition by showing that given any perturbation function  $g(x)$ , the derivative in  $\epsilon$  of  $E(f(x) + \epsilon g(x))$  vanishes at  $\epsilon=0$  if this equation is satisfied. Rederive the same Euler-Lagrange equation by showing instead that it causes the following directional derivative

$$\lim_{\epsilon \rightarrow 0} \frac{E(f + \epsilon g) - E(f)}{\epsilon}$$

to vanish for any function (“direction”)  $g$ . (Hint: use the Taylor series approximation of  $L$ .)

#### Solution:

Using the Taylor series expansion of  $L$  we obtain

$$\begin{aligned} \lim_{\epsilon \rightarrow 0} \frac{E(f + \epsilon g) - E(f)}{\epsilon} &= \lim_{\epsilon \rightarrow 0} \int_0^1 \frac{L(f(x) + \epsilon g(x), \dot{f}(x) + \epsilon \dot{g}(x), x) - L(f(x), \dot{f}(x), x)}{\epsilon} dx \\ &= \lim_{\epsilon \rightarrow 0} \int_0^1 \frac{L(f(x), \dot{f}(x), x) + \epsilon g(x) L_f(f(x), \dot{f}(x), x) + \epsilon \dot{g}(x) L_{\dot{f}}(f(x), \dot{f}(x), x) + O(\epsilon^2) - L(f(x), \dot{f}(x), x)}{\epsilon} dx \\ &= \lim_{\epsilon \rightarrow 0} \int_0^1 \frac{\epsilon g(x) L_f(f(x), \dot{f}(x), x) + \epsilon \dot{g}(x) L_{\dot{f}}(f(x), \dot{f}(x), x) + O(\epsilon^2)}{\epsilon} dx \\ &= \lim_{\epsilon \rightarrow 0} \int_0^1 (g(x) L_f(f(x), \dot{f}(x), x) + \dot{g}(x) L_{\dot{f}}(f(x), \dot{f}(x), x) + O(\epsilon)) dx \\ &= \int_0^1 (g(x) L_f(f(x), \dot{f}(x), x) + \dot{g}(x) L_{\dot{f}}(f(x), \dot{f}(x), x)) dx \end{aligned}$$

Now applying integration by parts, and assuming that the allowing perturbations  $g(x)$  vanish at  $x = 0$  and  $x = 1$  (in other words,  $g(0) = g(1) = 0$ ) we obtain the familiar right hand side

$$\int g \left( L_f - \frac{d}{dx} L_{\dot{f}} \right) dx$$

which will only vanish for all choices of  $g$  if the Euler Lagrange equation  $L_f - \frac{d}{dx} L_{\dot{f}} = 0$  is satisfied.

## 2 Finite-dimensional versus infinite-dimensional gradient

Validate the interpretation of the left hand side of the Euler-Lagrange equation as an infinite dimensional gradient by showing that the components of the “standard” (finite dimensional) gradient of

$$\hat{E}(f_0, f_1, \dots, f_N) = \sum_{n=0}^{N-1} L \left( f_n, \frac{f_{n+1} - f_n}{\Delta x}, x_n \right) \Delta x$$

converge (up to a common scalar factor of  $1/N$ ) to the values of  $L_f - \frac{\partial}{\partial x} L_{f'}$ . Note that  $\hat{E}$  is a finite approximation of  $E(f)$  from problem 1 where the variables  $f_0, f_1, \dots, f_N$  represent samples of the unknown function  $f$  at the points  $x_0, x_1, \dots, x_N$ . We will assume uniform spacing  $\Delta x = 1/N$  between samples, so that  $x_0 = 0, x_1 = \Delta x, x_2 = 2\Delta x, \dots$ . As we increase the number  $N$  of samples, we better and better approximate the function  $f$ , but the dimension of the resulting gradient vector for  $\hat{E}$  becomes larger and larger. Thus, as  $N \rightarrow \infty$ , the gradient becomes infinite dimensional.

**Solution:** Taking the partial derivative of  $\hat{E}$  with respect to one of its arguments  $f_i$  yields the finite dimensional gradient component as follows

$$\begin{aligned} \hat{E}_{f_i} &= \Delta x \frac{\partial}{\partial f_i} \left( \dots + L \left( f_{i-2}, \frac{f_{i-1} - f_{i-2}}{\Delta x}, x_{i-2} \right) + L \left( f_{i-1}, \frac{f_i - f_{i-1}}{\Delta x}, x_{i-1} \right) + L \left( f_i, \frac{f_{i+1} - f_i}{\Delta x}, x_i \right) + L \left( f_{i+1}, \frac{f_{i+2} - f_{i+1}}{\Delta x}, x_{i+1} \right) + \dots \right) \\ &= \frac{1}{N} \left( \frac{\partial}{\partial f_i} L \left( f_{i-1}, \frac{f_i - f_{i-1}}{\Delta x}, x_{i-1} \right) + \frac{\partial}{\partial f_i} L \left( f_i, \frac{f_{i+1} - f_i}{\Delta x}, x_i \right) \right) \\ &= \frac{1}{N} \left( \frac{1}{\Delta x} L_{f'} \left( f_{i-1}, \frac{f_i - f_{i-1}}{\Delta x}, x_{i-1} \right) - \frac{1}{\Delta x} L_{f'} \left( f_i, \frac{f_{i+1} - f_i}{\Delta x}, x_i \right) + L_f \left( f_i, \frac{f_{i+1} - f_i}{\Delta x}, x_i \right) \right) \\ &= \frac{1}{N} \left( L_f \left( f_i, \frac{f_{i+1} - f_i}{\Delta x}, x_i \right) - \frac{L_{f'} \left( f_i, \frac{f_{i+1} - f_i}{\Delta x}, x_i \right) - L_{f'} \left( f_{i-1}, \frac{f_i - f_{i-1}}{\Delta x}, x_{i-1} \right)}{\Delta x} \right) \\ &= \frac{1}{N} \left( L_f \left( f(x_i), \frac{f(x_i + \Delta x) - f(x_i)}{\Delta x}, x_i \right) - \frac{L_{f'} \left( f(x_i), \frac{f(x_i + \Delta x) - f(x_i)}{\Delta x}, x_i \right) - L_{f'} \left( f(x_i - \Delta x), \frac{f(x_i) - f(x_i - \Delta x)}{\Delta x}, x_i - \Delta x \right)}{\Delta x} \right) \end{aligned}$$

If we now ignore the common scalar factor of  $\frac{1}{N}$  we obtain that the rescaled (but same direction) gradient of  $\hat{E}$  has components

$$L_f \left( f(x_i), \frac{f(x_i + \Delta x) - f(x_i)}{\Delta x}, x_i \right) - \frac{L_{f'} \left( f(x_i), \frac{f(x_i + \Delta x) - f(x_i)}{\Delta x}, x_i \right) - L_{f'} \left( f(x_i - \Delta x), \frac{f(x_i) - f(x_i - \Delta x)}{\Delta x}, x_i - \Delta x \right)}{\Delta x}$$

If we let  $\Delta x \rightarrow 0$  (which is the same as letting  $N \rightarrow \infty$ ) we get

$$\begin{aligned} \lim_{\Delta x \rightarrow 0} & L_f \left( f(x_i), \frac{f(x_i + \Delta x) - f(x_i)}{\Delta x}, x_i \right) - \frac{L_{f'} \left( f(x_i), \frac{f(x_i + \Delta x) - f(x_i)}{\Delta x}, x_i \right) - L_{f'} \left( f(x_i - \Delta x), \frac{f(x_i) - f(x_i - \Delta x)}{\Delta x}, x_i - \Delta x \right)}{\Delta x} \\ &= L_f \left( f(x_i), \lim_{\Delta x \rightarrow 0} \frac{f(x_i + \Delta x) - f(x_i)}{\Delta x}, x_i \right) - \frac{\partial}{\partial x} L_{f'} \left( f(x_i), \lim_{\Delta x \rightarrow 0} \frac{f(x_i + \Delta x) - f(x_i)}{\Delta x}, x_i \right) \\ &= L_f \left( f(x_i), f'(x_i), x_i \right) - \frac{\partial}{\partial x} L_{f'} \left( f(x_i), f'(x_i), x_i \right) \end{aligned}$$

## 3 Calculate the gradient descent diffusion equation

Given an initial noisy 2D image  $I$  at time  $t = 0$ , derive the image diffusion equation

$$I_t = ?$$

which decreases the following energy (which is very large for a noisy image) as fast as possible. In other words, write down its gradient descent flow.

$$E = \frac{1}{2} \iint e^{(I_x^2 + I_y^2)} dx dy$$

**Solution:**

Notice that this has the form of the Perona-Malik style energy  $\int_{\Omega} c(\|\nabla I\|) d\Omega$  where  $c(z) = \frac{1}{2}e^{z^2}$ , and therefore  $\dot{c}(z) = ze^{z^2}$ . As such we may apply the general gradient descent Perona-Malik diffusion formula derived in class as follows.

$$\begin{aligned}
I_t &= \nabla \cdot \left( \frac{\dot{c}(\|\nabla I\|)}{\|\nabla I\|} \nabla I \right) = \nabla \cdot \left( \frac{\|\nabla I\| e^{\|\nabla I\|^2}}{\|\nabla I\|} \nabla I \right) = \nabla \cdot \left( e^{\|\nabla I\|^2} \nabla I \right) \\
&= e^{\|\nabla I\|^2} \nabla \cdot \nabla I + \nabla \left( e^{\|\nabla I\|^2} \right) \cdot \nabla I = e^{\|\nabla I\|^2} \Delta I + \left( e^{\|\nabla I\|^2} \nabla \|\nabla I\|^2 \right) \cdot \nabla I \\
&= e^{\|\nabla I\|^2} \left( \Delta I + \nabla(I_x^2 + I_y^2) \cdot \nabla I \right) \\
&= e^{\|\nabla I\|^2} \left( \Delta I + 2 \begin{bmatrix} I_x I_{xx} + I_y I_{xy} \\ I_y I_{xy} + I_y I_{yy} \end{bmatrix} \cdot \begin{bmatrix} I_x \\ I_y \end{bmatrix} \right) \\
I_t &= e^{\|\nabla I\|^2} (I_{xx}(1 + 2I_x^2) + 4I_x I_y I_{xy} + I_{yy}(1 + 2I_y^2))
\end{aligned}$$

Taking the direct approach instead, we note that the Lagrangian  $L(I, I_x, I_y, x, y) = \frac{1}{2}e^{(I_x^2 + I_y^2)}$  depends only on  $I_x$  and  $I_y$  and as such, the general gradient descent formula taking right-hand side of the Euler-Lagrange equation can be written as follows.

$$\begin{aligned}
I_t &= \frac{\partial}{\partial x} L_{I_x} + \frac{\partial}{\partial y} L_{I_y} \\
&= \frac{\partial}{\partial x} \left( I_x e^{(I_x^2 + I_y^2)} \right) + \frac{\partial}{\partial y} \left( I_y e^{(I_x^2 + I_y^2)} \right) \\
&= \left( I_{xx} e^{(I_x^2 + I_y^2)} + 2I_x (I_x I_{xx} + I_y I_{xy}) e^{(I_x^2 + I_y^2)} \right) + \left( I_{yy} e^{(I_x^2 + I_y^2)} + 2I_y (I_x I_{xy} + I_y I_{yy}) e^{(I_x^2 + I_y^2)} \right) \\
&= e^{(I_x^2 + I_y^2)} ((1 + 2I_x^2) I_{xx} + 4I_x I_y I_{xy} + (1 + 2I_y^2) I_{yy}) \\
&= e^{\|\nabla I\|^2} ((1 + 2I_x^2) I_{xx} + 4I_x I_y I_{xy} + (1 + 2I_y^2) I_{yy})
\end{aligned}$$