

Extension to Quadrilateral Element of Three Field Hu-Washizu Elasticity Formulation Based on Biorthogonal Systems

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1 Linear elastic continuum problem

In this section we briefly recovery the equations governing the linear elastic problem. The equilibrium equation is:

$$-\text{div}(\boldsymbol{\sigma}) = \mathbf{f} , \quad (1)$$

while in small deformation is:

$$\mathbf{d} = \boldsymbol{\varepsilon}(\mathbf{u}) = \frac{1}{2}(\boldsymbol{\nabla} \mathbf{u} + \boldsymbol{\nabla} \mathbf{u}^T) . \quad (2)$$

In the case of linear elasticity we have:

$$\boldsymbol{\sigma} = \lambda \text{tr}(\boldsymbol{\varepsilon}) \mathbf{I} + 2\mu \boldsymbol{\varepsilon} \quad (3)$$

where μ and λ are the Lamé constant. By some algebra one obtains:

$$\boldsymbol{\sigma} = \begin{pmatrix} \lambda(\varepsilon_{11} + \varepsilon_{22}) & 0 \\ 0 & \lambda(\varepsilon_{11} + \varepsilon_{22}) \end{pmatrix} + 2\mu \begin{pmatrix} \varepsilon_{11} & \varepsilon_{12} \\ \varepsilon_{12} & \varepsilon_{22} \end{pmatrix} , \quad (4)$$

and rearranging the equation (4):

$$\boldsymbol{\sigma} = \begin{pmatrix} (\lambda + 2\mu)\varepsilon_{11} + \lambda \varepsilon_{22} & 2\mu \varepsilon_{12} \\ 2\mu \varepsilon_{12} & (\lambda + 2\mu) \varepsilon_{22} + \lambda \varepsilon_{11} \end{pmatrix} . \quad (5)$$

2 Briefly introduction to modify Hu-Washizu

We define the trial variables: $\boldsymbol{\varepsilon}(\mathbf{u})$, \mathbf{d} and $\boldsymbol{\sigma}$, while the test variables are: $\boldsymbol{\varepsilon}(\mathbf{v})$, \mathbf{e} and $\boldsymbol{\tau}$.

$$- \int_{\Omega} \text{div}(\mathbf{C} : \mathbf{d})) \cdot \mathbf{v} = \mathbf{f} \quad (6)$$

$$a((\mathbf{u}, \mathbf{d}), (\mathbf{v}, \mathbf{e})) + b((\mathbf{v}, \mathbf{e}), \boldsymbol{\sigma}) = l(\mathbf{v}) \quad (7)$$

$$b((\mathbf{u}, \mathbf{d}), \boldsymbol{\tau}) = 0 \quad (8)$$

where:

$$a((\mathbf{u}, \mathbf{d}), (\mathbf{v}, \mathbf{e})) = \int_{\Omega} \mathbf{d} : (\mathbf{C} : \mathbf{e}) \, dx + \alpha \int_{\Omega} (\boldsymbol{\varepsilon}(\mathbf{u}) - \mathbf{d}) : (\boldsymbol{\varepsilon}(\mathbf{v}) - \mathbf{e}) \, dx , \quad (9)$$

$$b((\mathbf{u}, \mathbf{d}), \boldsymbol{\tau}) = \int_{\Omega} (\boldsymbol{\varepsilon}(\mathbf{u}) - \mathbf{d}) : \boldsymbol{\tau} \, dx . \quad (10)$$

The modify weak formulation of the problem is:

$$\begin{cases} \alpha \int_{\Omega} (\varepsilon(\mathbf{u}) - \mathbf{d}) : \varepsilon(\mathbf{v}) \, dx + \int_{\Omega} \varepsilon(\mathbf{v}) : \boldsymbol{\sigma} \, dx & = \int_{\Omega} \mathbf{f} \cdot \mathbf{v} \, dx \\ \int_{\Omega} \mathbf{d} : \mathbf{C} \mathbf{e} \, dx - \alpha \int_{\Omega} (\varepsilon(\mathbf{u}) - \mathbf{d}) : \mathbf{e} \, dx - \int_{\Omega} \mathbf{e} : \boldsymbol{\sigma} \, dx & = 0 \\ \int_{\Omega} (\varepsilon(\mathbf{u}) - \mathbf{d}) : \boldsymbol{\tau} \, dx & = 0 \end{cases} \quad (11)$$

by rearranging:

$$\begin{cases} \alpha \int_{\Omega} \varepsilon(\mathbf{u}) : \varepsilon(\mathbf{v}) \, dx - \alpha \int_{\Omega} \mathbf{d} : \varepsilon(\mathbf{v}) \, dx + \int_{\Omega} \varepsilon(\mathbf{v}) : \boldsymbol{\sigma} \, dx & = \int_{\Omega} \mathbf{f} \cdot \mathbf{v} \, dx \\ -\alpha \int_{\Omega} \varepsilon(\mathbf{u}) : \mathbf{e} \, dx + \int_{\Omega} \mathbf{d} : \mathbf{C} \mathbf{e} \, dx + \alpha \int_{\Omega} \mathbf{d} : \mathbf{e} \, dx - \int_{\Omega} \mathbf{e} : \boldsymbol{\sigma} \, dx & = 0 \\ \int_{\Omega} \varepsilon(\mathbf{u}) : \boldsymbol{\tau} \, dx - \int_{\Omega} \mathbf{d} : \boldsymbol{\tau} \, dx & = 0 \end{cases} \quad (12)$$

It is possible to rewrite the system in equation (12) in matrix form in the following way:

$$\begin{bmatrix} \alpha \mathbf{A} & -\alpha \mathbf{B} & \mathbf{W} \\ -\alpha \mathbf{B}^T & \mathbf{K} + \alpha \mathbf{M} & -\mathbf{D} \\ \mathbf{W}^T & -\mathbf{D}^T & \mathbf{0} \end{bmatrix} \begin{bmatrix} \mathbf{x}_u \\ \mathbf{x}_d \\ \mathbf{x}_\sigma \end{bmatrix} = \begin{bmatrix} \mathbf{b}_f \\ \mathbf{0} \\ \mathbf{0} \end{bmatrix}, \quad (13)$$

where $\mathbf{A} = \int_{\Omega} \varepsilon(\mathbf{u}) : \varepsilon(\mathbf{v})$, $\mathbf{B} = \int_{\Omega} \mathbf{d} : \varepsilon(\mathbf{v})$, $\mathbf{W} = \int_{\Omega} \boldsymbol{\sigma} : \varepsilon(\mathbf{v})$, $\mathbf{K} = \int_{\Omega} \mathbf{C} \mathbf{e} : \mathbf{d}$, $\mathbf{M} = \int_{\Omega} \mathbf{e} : \mathbf{d}$, $\mathbf{D} = \int_{\Omega} \boldsymbol{\sigma} : \mathbf{e}$. \mathbf{D} is a diagonal matrix. Using this property it is possible condense statically \mathbf{x}_d and \mathbf{x}_σ , and we obtain the following system in the only unknown \mathbf{x}_u :

$$[\alpha \mathbf{A} - \alpha (\mathbf{B} \mathbf{D}^{-1} \mathbf{W}^{-T} + \mathbf{W} \mathbf{D}^{-1} \mathbf{B}^T) + \mathbf{W} \mathbf{D}^{-1} (\mathbf{K} + \alpha \mathbf{M}) \mathbf{D}^{-1} \mathbf{W}^T] \mathbf{x}_u = \mathbf{b}_f \quad (14)$$

3 Finite element discretization

We consider a quasi-uniform triangulation \mathcal{T}_h of the polygonal domain Ω consists of simply, either quadrilateral or hexahedral. We take into account of standard bilinear finite element space $K_h \subset H^1(\Omega)$ defined on the triangulation \mathcal{T}_h , where:

$$K_h := \{v \in C^0(\Omega) : v|_T \in \mathcal{Q}_1(T), T \in \mathcal{T}_h\}, \quad K_h^0 = K_h \cap H_0^1(\Omega), \quad (15)$$

and the space of bubble functions

$$B_h := \left\{ b_T \in H^1(T) : b_T|_{\partial T} = 0 \text{ and } \int_T b_T \, dx > 0, T \in \mathcal{T}_h \right\}, \quad (16)$$

and we define the spaces for strain and displacement as $\mathbf{S}_h := [K_h]^{2 \times 2}$ and $\mathbf{V}_h := [K_h^0 \oplus B_h]^2$. In the next section we discuss the different choosing of bubble functions. For the discrete stress space we use:

$$\mathbf{M}_h := \left\{ \boldsymbol{\tau}_h \in [\mathbf{M}_h]^{2 \times 2} : \int_{\Omega} \boldsymbol{\tau}_h : \mathbf{1} \, dx = 0 \right\} \subset \mathbf{S}_0, \quad (17)$$

and let $\{\phi_1, \dots, \phi_n\}$ and $\{\mu_1, \dots, \mu_n\}$ the n the basis functions for the space V_h and M_h respectively, we construct the functions μ_i using the following biorthogonality property between the space V_h and M_h :

$$\int_{\Omega} \mu_i \phi_j \, dx = c_j \delta_{ij}, \quad c_j \neq 0, \quad 1 \leq i, j \leq n, \quad (18)$$

where δ_{ij} is Kronecker symbol, and c_j is a scaling factor which can be chosen to be proportion al to the area of support of ϕ_j . The local basis function of K_h and M_h for the reference square element (see figure 1) $\hat{T} := \{(x, y) : -1 \leq x \leq 1, -1 \leq y \leq 1\}$ are:

$$\begin{aligned} \phi_1 &= \frac{1}{4}(1-x)(1-y), & \phi_2 &= \frac{1}{4}(1+x)(1-y), \\ \phi_3 &= \frac{1}{4}(1+x)(1+y), & \phi_4 &= \frac{1}{4}(1-x)(1+y). \end{aligned} \quad (19)$$

and

$$\begin{aligned} \mu_1 &= 1 - 3x - 3y + 9xy, & \mu_2 &= 1 + 3x - 3y - 9xy, \\ \mu_3 &= 1 + 3x + 3y + 9xy, & \mu_4 &= 1 - 3x + 3y + 9xy. \end{aligned} \quad (20)$$

It is important to observe that the global basis functions of the space M_h are not continuous.

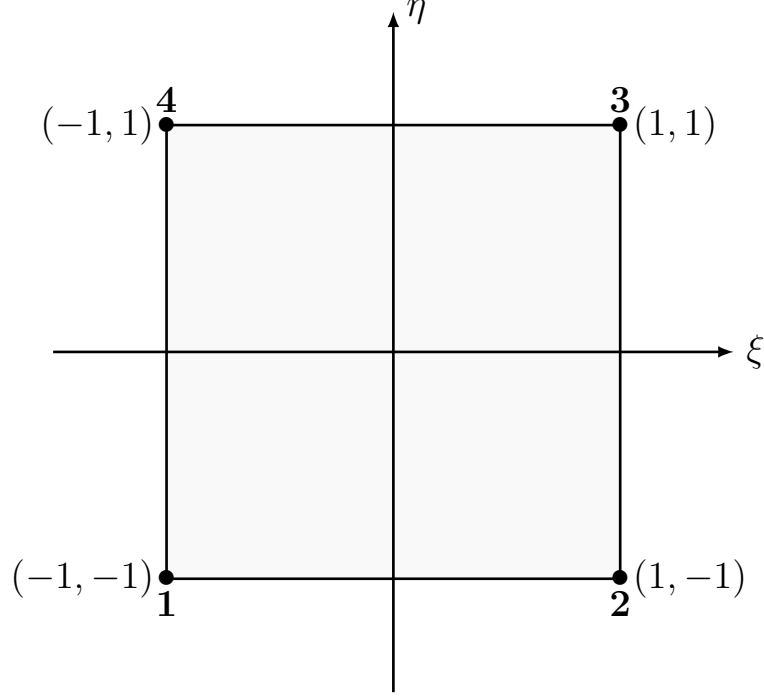


Figura 1: Reference Element

4 Bubble functions

In this section we detail the different choosing of the bubble functions. Addition of the bubble functions is essential to create a stable space. we have four types of bubbles. In the first two cases we use a modification of the standard bubble function ,that is for the reference element:

$$b_T(x, y) = (1 - x^2)(1 - y^2) , \quad (21)$$

while in the next two, we add to the standard bubble function another one.

4.1 One Bubble function (type 1)

As a first choice of bubble function we use:

$$\hat{b}_T(x, y) = c_T \cdot \phi_T(x, y) \cdot b_T(x, y) , \quad (22)$$

where c_T is a coefficient in order to obtain $\hat{b}_T(x_g, y_g) = 1$ (where \mathbf{g} is the centroid of the elements), ϕ_K is the standard bilinear basis function corresponding to the lower-left corner of the square T . In the case of reference square element we obtain:

$$\hat{b}_T(x, y) = (1 - x)(1 - y)(1 - x^2)(1 - y^2) . \quad (23)$$

4.2 One Bubble function (type 2)

The second choice of bubble function we take:

$$\hat{b}_T(x, y) = c_T \cdot (a + bx + cy) \cdot b_T(x, y) , \quad (24)$$

where $a, b, c \in \mathbb{R}$ and $a, b, c \neq 0$. For simplicity we set $a = b = c = 1$ and we obtain for the reference square:

$$\hat{b}_T(x, y) = (1 + x + y)(1 - x^2)(1 - y^2). \quad (25)$$

4.3 Two Bubble functions

Using two bubble functions, where the first is the standard bubble function and the second bubble is a modification of the standard bubble:

$$\begin{aligned} \hat{b}_{T1}(x, y) &= b_T, \\ \hat{b}_{T2}(x, y) &= c_T \cdot (ax + by) \cdot b_T, \end{aligned} \quad (26)$$

where $a, b \in \mathbb{R}$ and $a^2 + b^2 \neq 0$. For the sake of simplicity we adopt $a = b = 1$. One obtains:

$$\begin{aligned} \hat{b}_{T1}(x, y) &= (1 - x^2)(1 - y^2), \\ \hat{b}_{T2}(x, y) &= (x + y)(1 - x^2)(1 - y^2). \end{aligned} \quad (27)$$

4.4 Two Bubble functions, which one mixed

As a finally choice of bubbles we use a standard bubble function plus one mixed bubble function for the two components of displacement.

$$\begin{aligned} \hat{b}_{T1}(x, y) &= b_T, \\ \hat{b}_{T2,x}(x, y) &= (\nabla \phi_1)_x \cdot b_T, \\ \hat{b}_{T2,y}(x, y) &= (\nabla \phi_1)_y \cdot b_T, \end{aligned} \quad (28)$$

where $(\nabla \phi_1)_i$ is i -th component of the gradient of the first shape function ϕ . In this way we have as shape function for the displacement using the mixed bubble function the vector $[\hat{b}_{T2,x}(x, y), \hat{b}_{T2,y}(x, y)]$.

5 Numerical example

In this section we report some examples using the presented formulation to proven the good behaviour.

5.1 Square problem

First example is a unit square domain with homogeneous Dirichlet boundary conditions. The Lamé constant are fix to $\lambda = 123$ and $\mu = 79.3$. By imposition of the previously exact solution one obtain for the body force f

$$\begin{aligned} f_1 &= -\pi^2 \cos(\pi x) \sin(\pi y) (\lambda + \mu + 2\lambda \cos(\pi y) + 12\mu \cos(\pi y)), \\ f_2 &= -\pi^2 \sin(\pi x) (\lambda \cos(\pi y) + 3\mu \cos(\pi y) + 2\lambda (2 \cos(\pi y)^2 - 1) + 2\mu (2 \cos(\pi y)^2 - 1)) \end{aligned} \quad (29)$$

The exact solution is

$$u_1 = \cos(\pi x) \sin(2\pi y), \quad u_2 = \sin(\pi x) \cos(\pi y). \quad (30)$$

The problem is study using two type of mesh: first of all using a square mesh and before using a trapezoidal mesh. The two types of mesh are shown in figures 2(a) and 2(b). Figures 3(a), 3(b), 4(a) and 4(b) shown the error in norm L^2 in the case of regular mesh for the different types of bubble functions used and types of coefficient α . All types of element converge in a good way. In Figures 5(a), 5(b), 6(a) and 6(b) we report the previously results in the case of trapezoidal meshes.

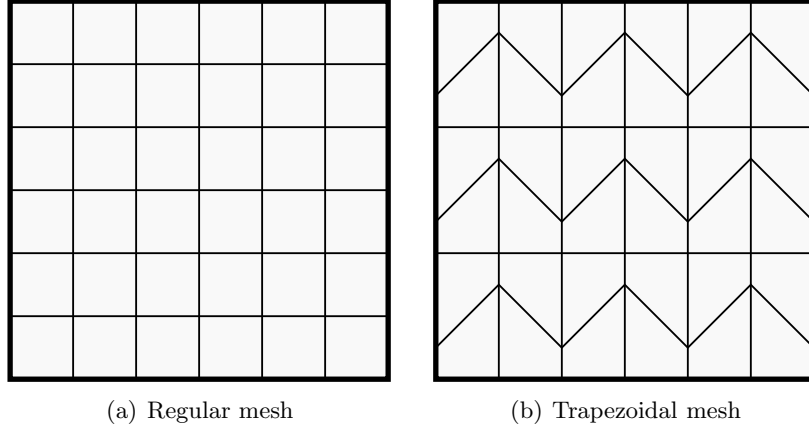


Figura 2: Square Problem

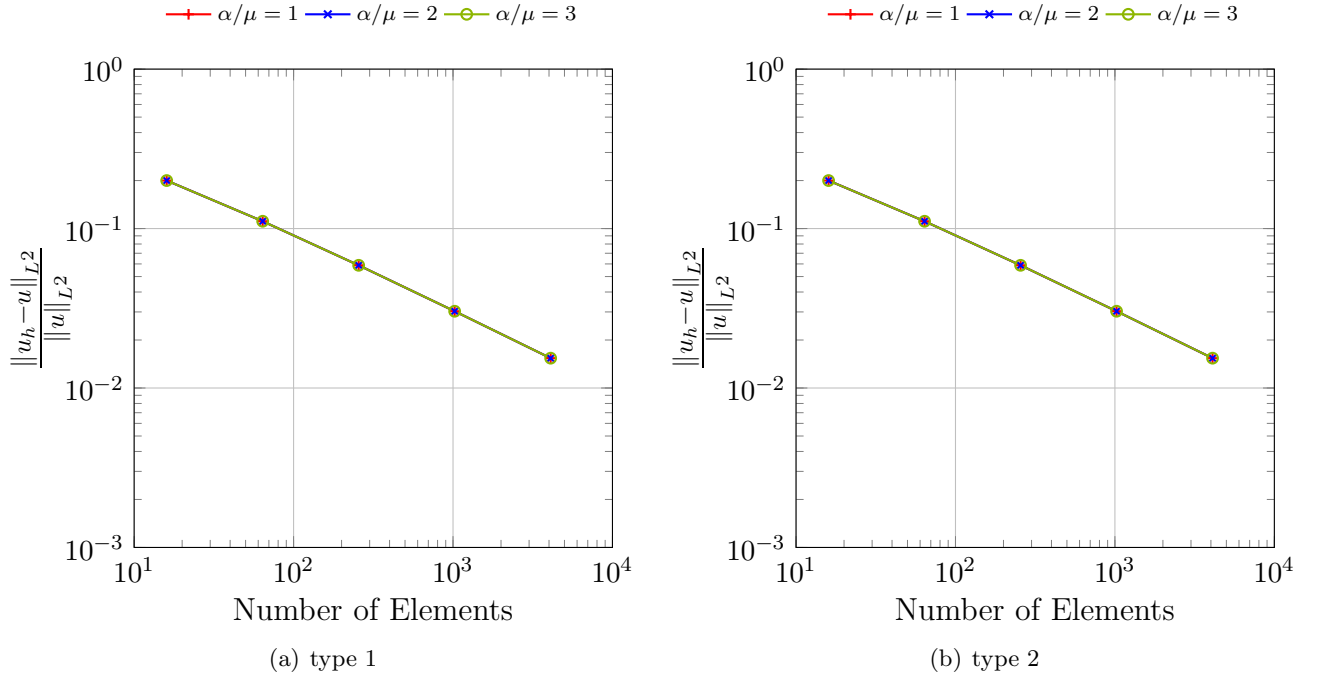


Figure 3: The relative error vs. the number of elements measured relative to the L^2 norm (Case one bubble function and regular mesh)

5.2 Cantilever beam problem

Now we consider the beam with length $L = 10$ and height $l = 2$ as we shown in figure (). The Young modulus is set equal to $E = 1500$ and the Poisson $\nu = 0.4999$ and subjected to a distributed load as in figure 7 with $f = 300$. The exact solution is:

$$\begin{aligned}
 u(x, y) &= \frac{2f}{El}(1 - \nu^2)x \left(\frac{l}{2} - y \right), \\
 v(x, y) &= \frac{f}{El} \left[x^2 + \frac{\nu}{1 - \nu} (y^2 - ly) \right].
 \end{aligned} \tag{31}$$

We use to model the beam two types of mesh: regular and trapezoidal as in the previously example (see figures 2(a) and 2(b)). we shown in figures 10(a), 10(b), 11(a) and 11(b) the L^2 -norm error for different

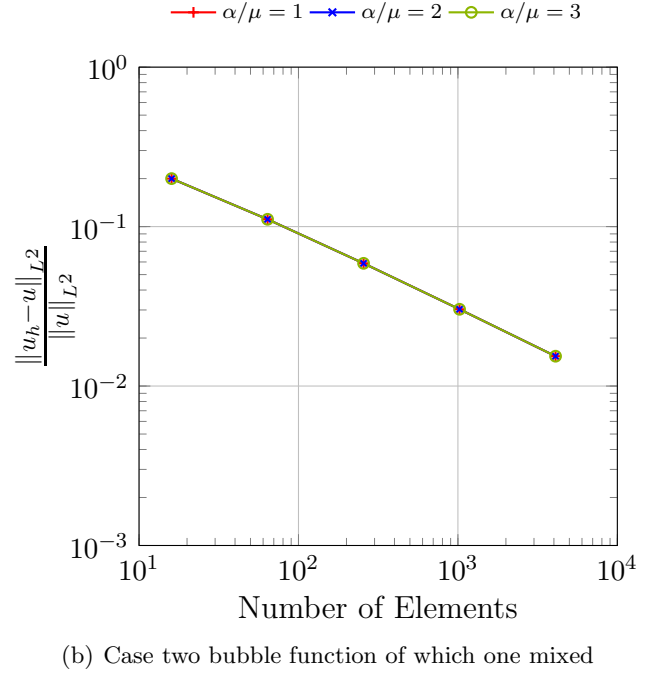
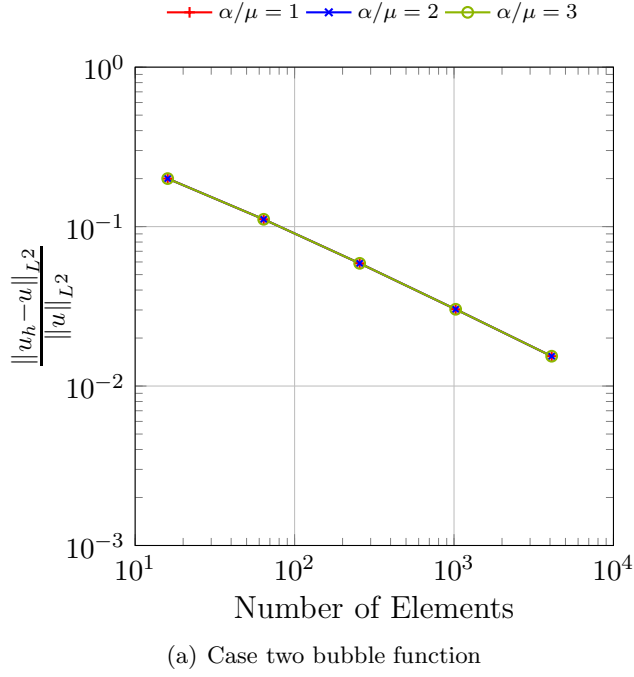


Figure 4: The relative error versus the number of elements measured relative to the L^2 norm (regular mesh)

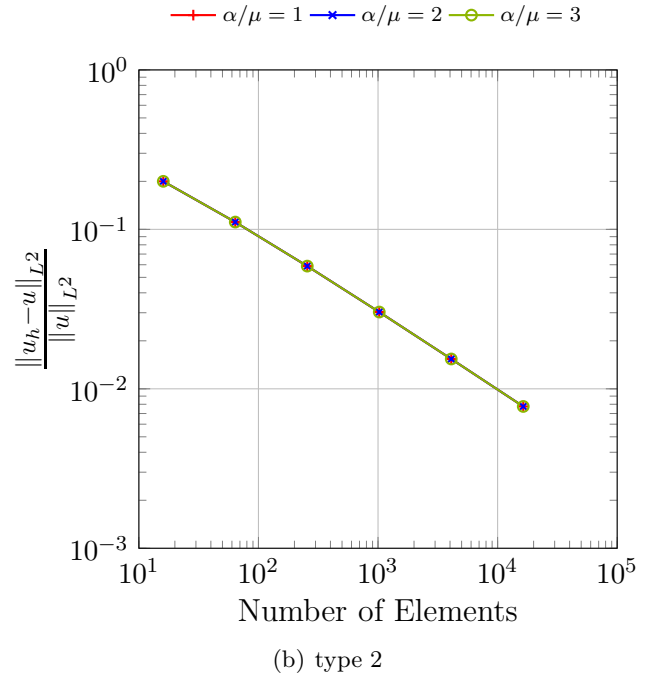
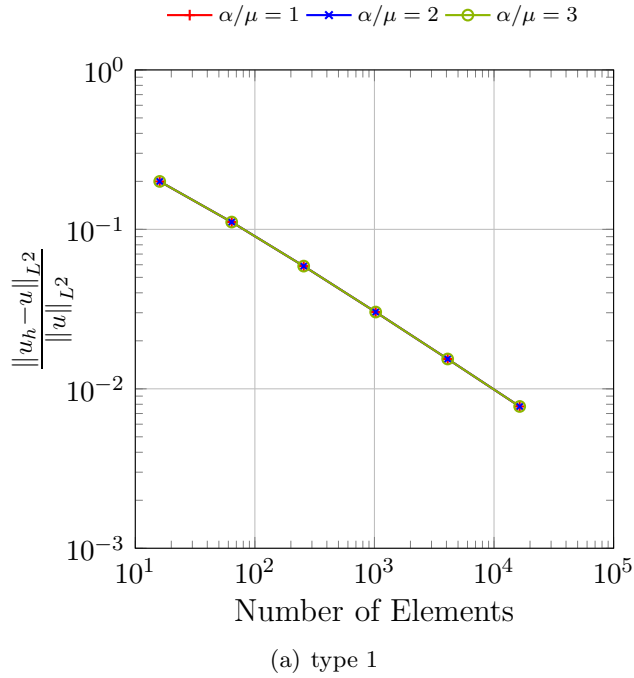


Figure 5: The relative error vs. the number of elements measured relative to the L^2 norm (Case one bubble function and Trapezoidal mesh)

types of bubble functions used in the case of $\alpha/\mu := 1, 2, 3$, while in figures 10(a), 10(b), 11(a) and 11(b) the same plots using trapezoidal meshes. In the all cases the elements distorted have a good behaviour respect to the regular mesh.

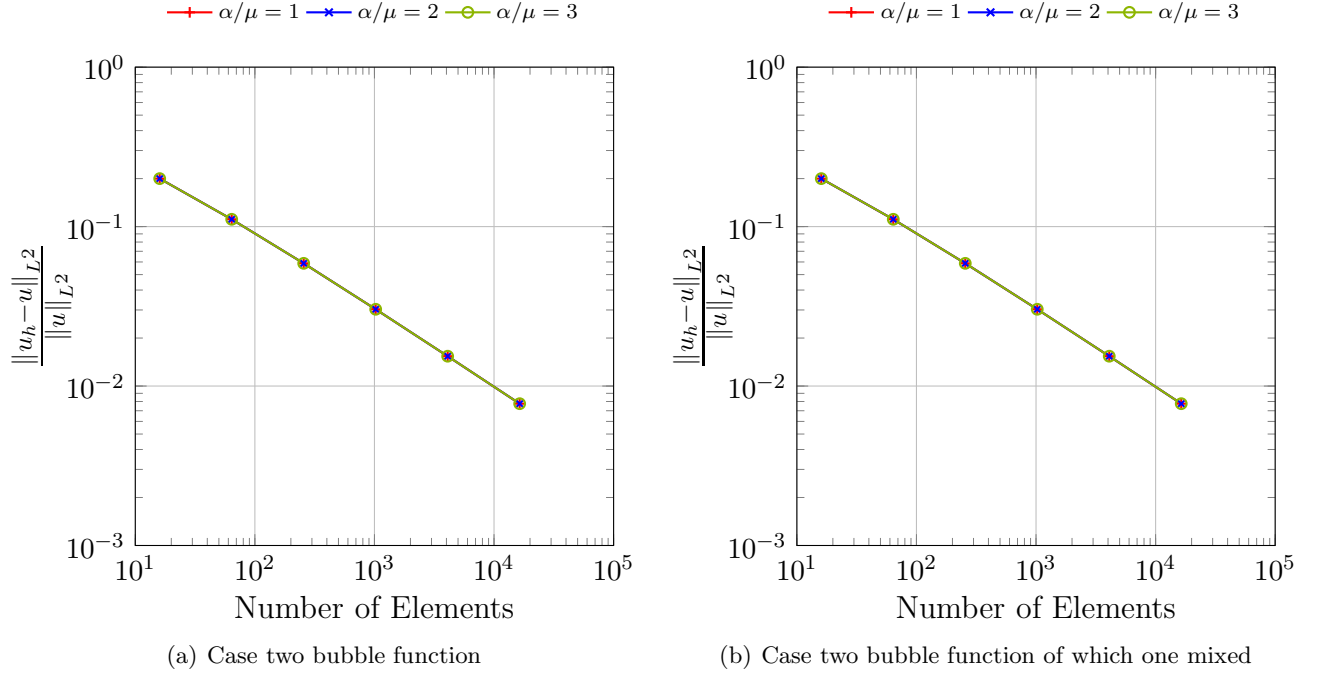


Figure 6: The relative error vs. the number of elements measured relative to the L^2 norm (Case one bubble function and Trapezoidal mesh)

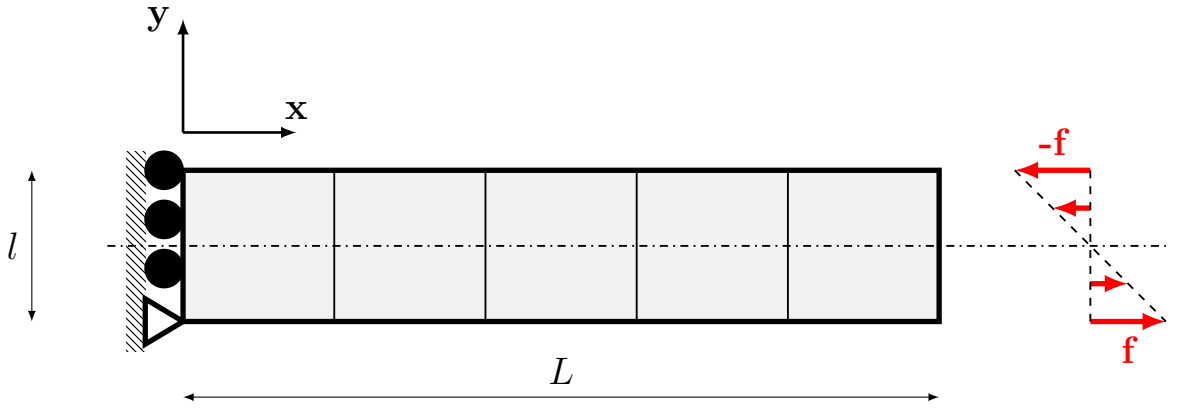


Figure 7: Beam cantilever geometry

5.3 Cook's membrane

The final example is the Cook's membrane. That is a typical benchmark and consist of a beam with vertex: $(0,0)$, $(48,44)$, $(48,60)$ and $(0,44)$. The left vertical edge is clamped and the right vertical edge subjected to the vertical distributed forces with resultant $F = 100$ as it shown in figure 12. The material properties are taken to be $E = 250$ and $\nu = 0.4999$, so that a nearly incompressible response is obtained. We report in figures 13(a), 13(b), 13(c) and 13(d) the vertical displacement of the point A versus the number of element per side for different choosing of the parameter $\alpha = \{1, \mu, 2\mu, 3\mu\}$. All elements return different behaviour using different coefficients α . In the case of $\alpha = 1$, figure 13(a), the obtained results completely not converge to the reference solution.

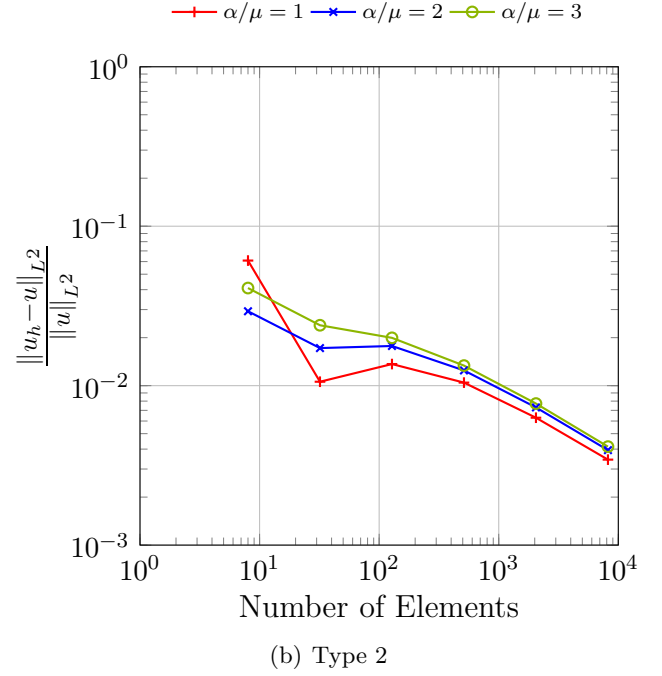
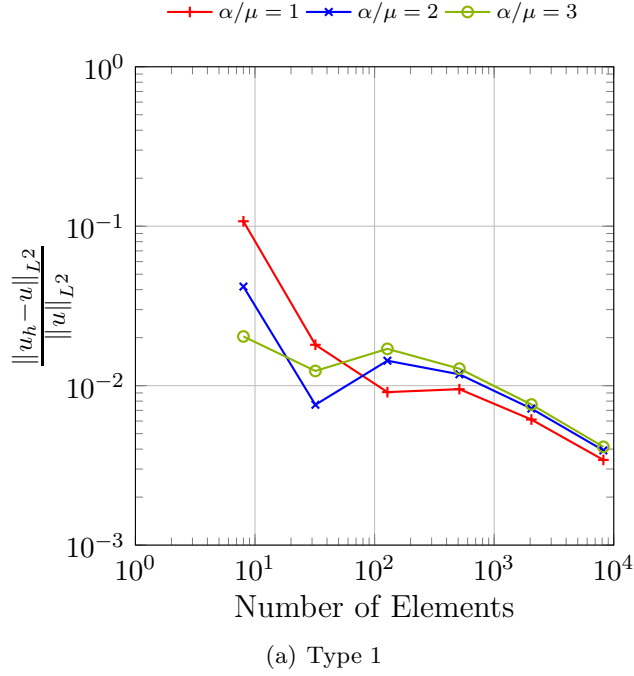


Figure 8: Beam Cantilever: the relative error vs. the number of elements measured relative to the L^2 norm (regular mesh)

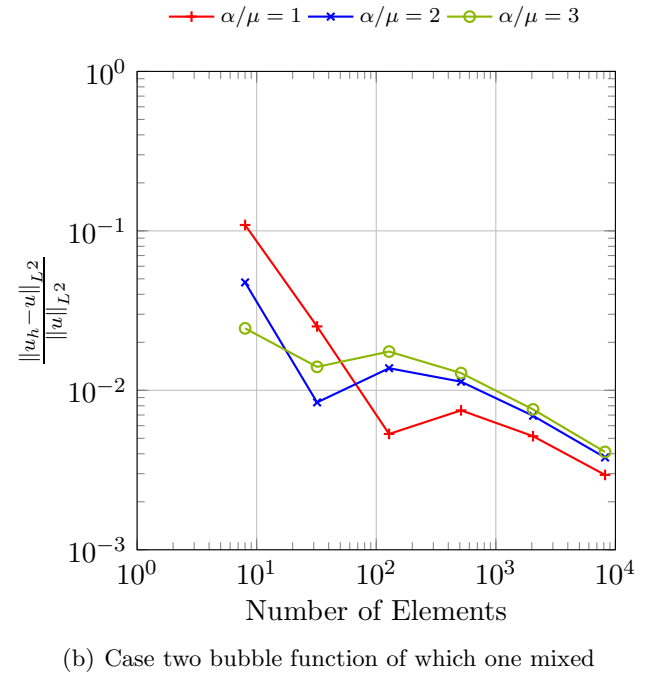
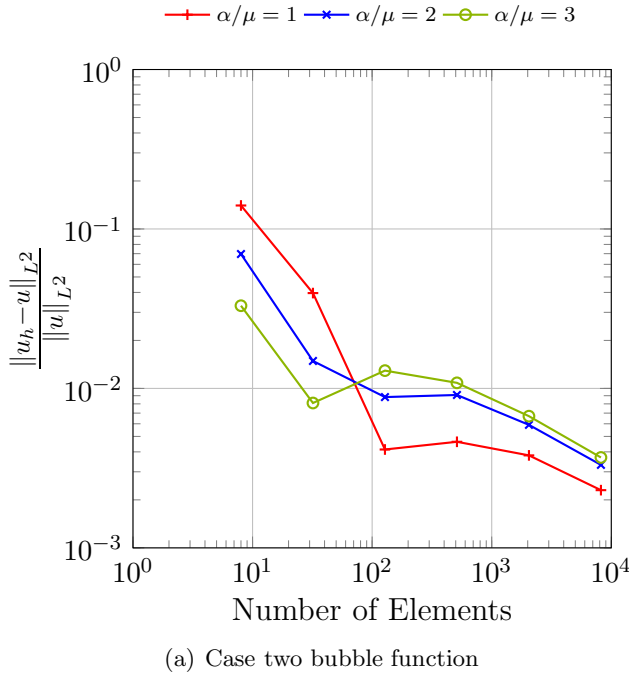


Figure 9: Beam Cantilever: the relative error vs. the number of elements measured relative to the L^2 norm (regular mesh)

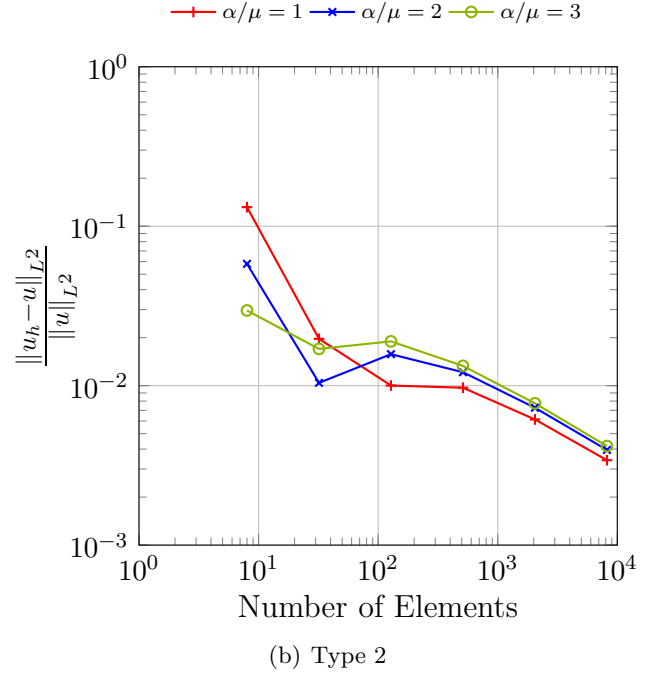
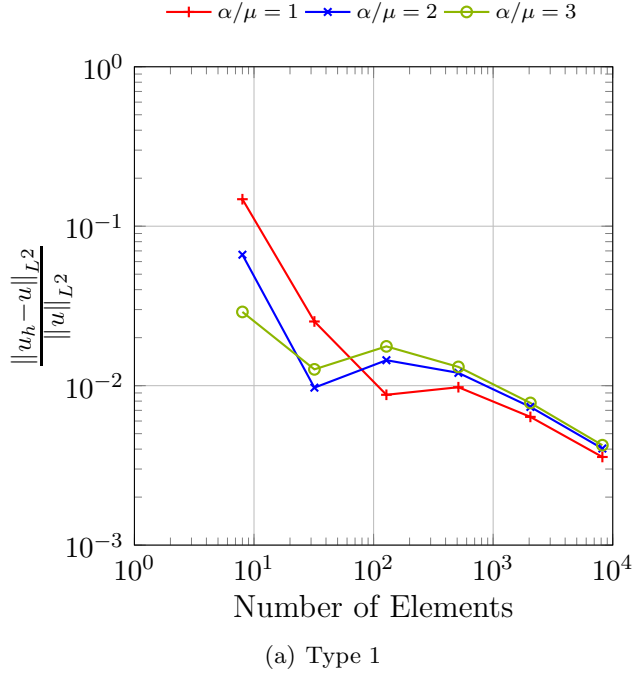


Figure 10: Beam Cantilever: the relative error vs. the number of elements measured relative to the L^2 norm (trapezoidal mesh)

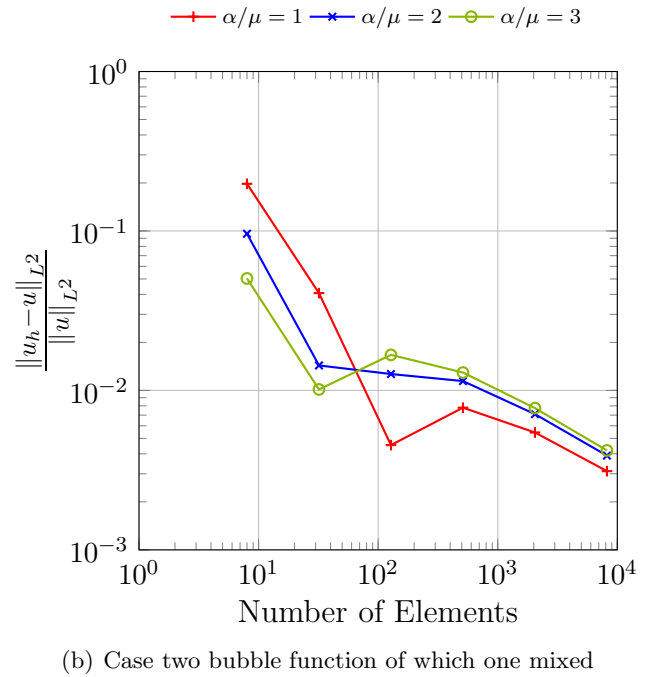
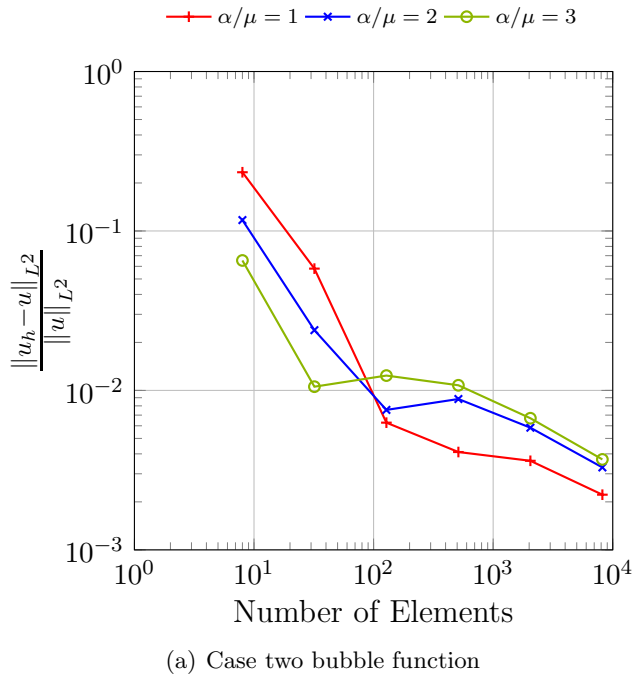


Figure 11: Beam Cantilever: the relative error vs. the number of elements measured relative to the L^2 norm (trapezoidal mesh)

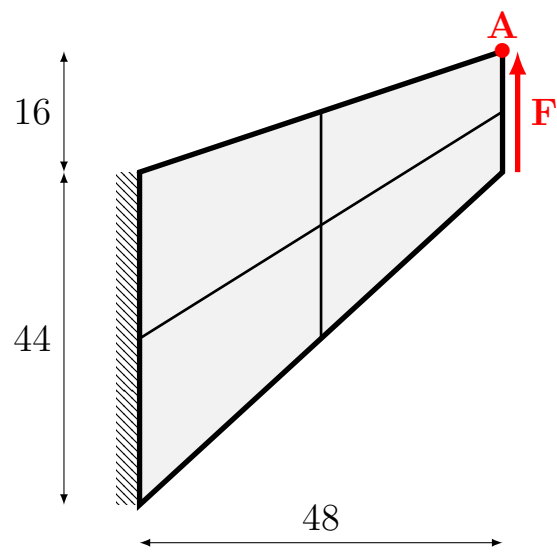


Figura 12: Cook's Membrane geometry

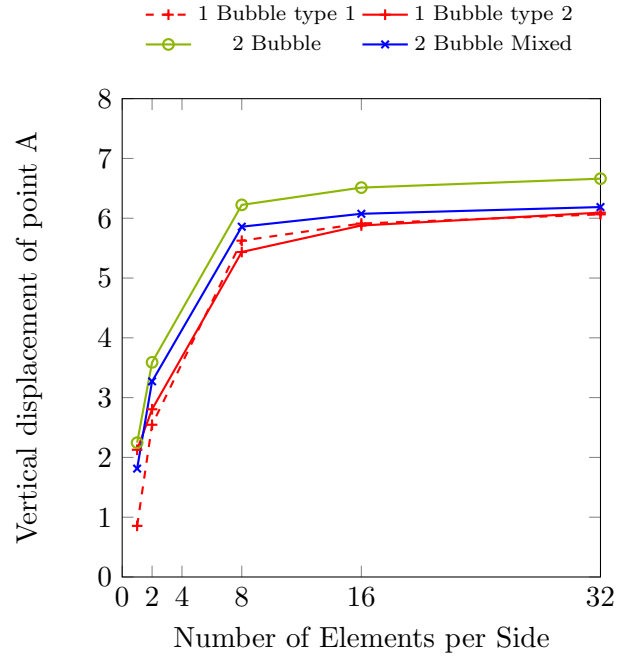
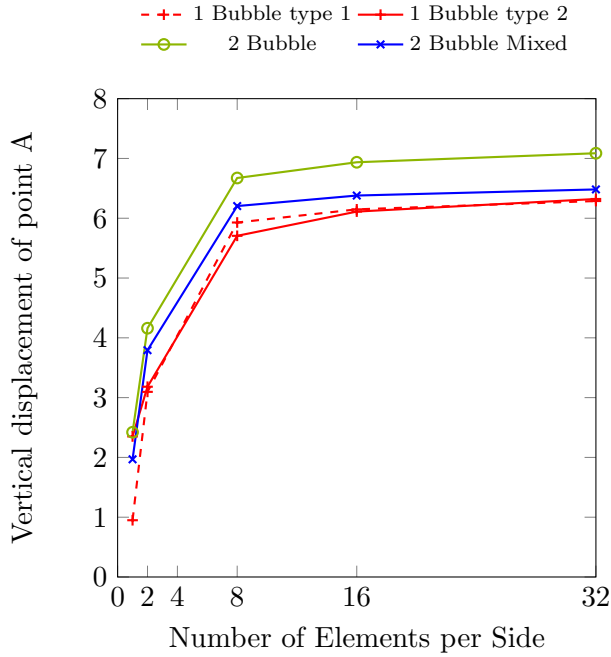
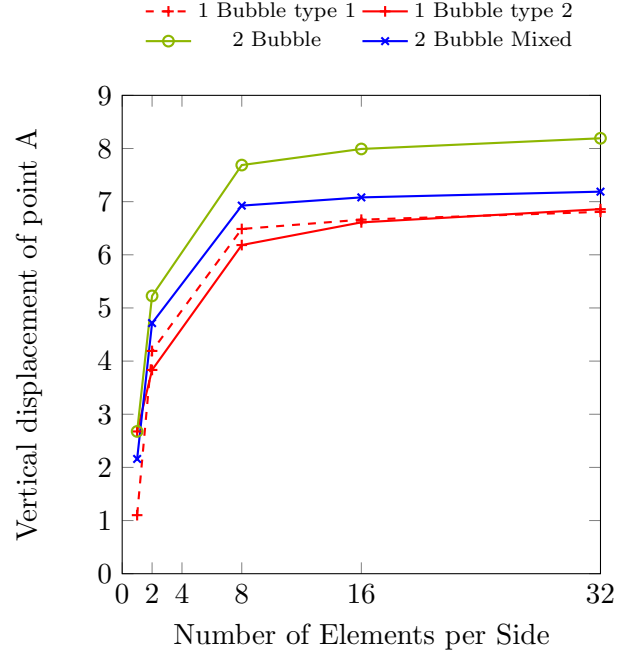
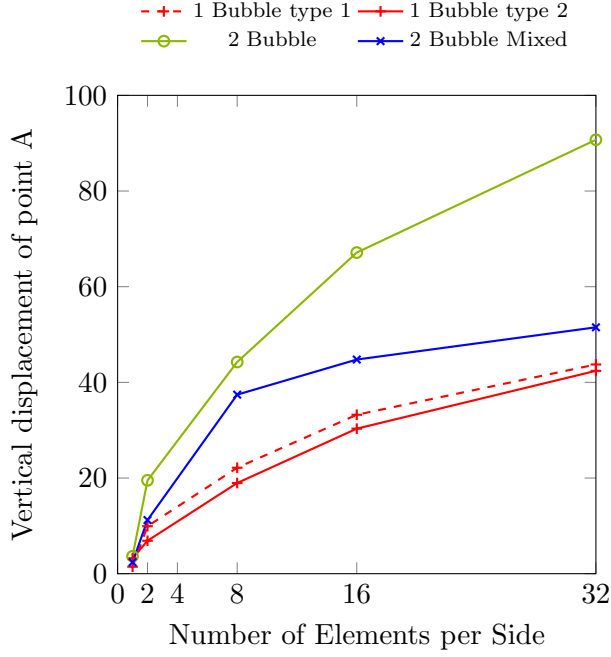


Figure 13: Vertical Displacement of point A vs. the number of elements per side