

# UNIFIED ANALYSIS OF MIXED METHODS FOR ELASTICITY WITH WEAKLY SYMMETRIC STRESS

JEONGHUN J. LEE

**ABSTRACT.** We propose a framework for the unified analysis of mixed methods for elasticity with weakly symmetric stress. Based on a commuting diagram in the weakly symmetric elasticity complex and extending a previous stability result, stable mixed methods are obtained by combining Stokes stable and elasticity stable finite elements. We show that the framework can be used to analyze most existing mixed methods for the elasticity problem with elementary techniques. The framework also allows us to construct new stable mixed finite elements.

## 1. INTRODUCTION

In the Hellinger–Reissner formulation of linear elasticity, for a given external body force and boundary conditions, the stress and displacement are sought as a saddle point of the Hellinger–Reissner functional. In this saddle point problem the stress tensor is directly obtained without volumetric locking in nearly incompressible materials [5]. However, the symmetry condition of the stress tensor gives a highly nontrivial obstacle in finding stable mixed finite elements for the saddle point problem. Another way to find mixed methods for the problem is to impose the symmetry condition weakly by requiring the stress to be orthogonal against a certain space of skew-symmetric tensors [1, 13]. This alternative approach turned out to be successful and various stable mixed finite elements have been developed based on this idea [3, 7, 9, 12, 15, 16, 17, 21, 24, 25]. In this paper we will call them weak symmetry elements.

There are several different ways to analyze the the stability of weak symmetry elements. In early research [3, 16] a connection between the Stokes equation and the linear elasticity equation with weak symmetry is used for the proof of stability. An analysis using mesh-dependent norms is also proposed [24, 25]. A breakthrough is made in [7] in the development of the weakly symmetric elasticity complex. In [7] a commuting diagram in the elasticity complex and a diagram chasing type argument are used to prove the stability of the Arnold–Falk–Winther family. In [9] an analysis, based on a connection with the Stokes equation, is revisited with the commuting diagram in the weakly symmetric elasticity complex. In the analysis the stability proof is reduced to proving existence of an interpolation operator satisfying some conditions and some new elements were proposed in [9] using this approach. This idea is adopted in [12, 17] to construct new elements with the aid of cleverly-designed matrix bubble functions and the results in [7].

---

1991 *Mathematics Subject Classification.* 65N30, 65N12.

*Key words and phrases.* linear elasticity, weakly symmetric stress, mixed finite elements, error analysis.

The aforementioned ways for the stability proof, although they are interesting, are not elementary. In fact, some of them need concepts which are not familiar to most numerical analysts. The goal of this paper is to present a unified framework for the analysis of weak symmetry elements with elementary techniques. In addition to providing an easier analysis of existing elements, the framework gives an insight in developing new weak symmetry elements.

The paper is organized as follows. In section 2, we summarize the notations and review the Hellinger–Reissner formulation of linear elasticity with weakly symmetric stress. In section 3, we introduce an abstract framework for the unified analysis and prove a priori error estimates. In section 4, we give examples of weak symmetry elements to which the abstract framework can be applied.

## 2. PRELIMINARIES

**2.1. Notations.** Let  $\Omega$  be a bounded domain in  $\mathbb{R}^n$  ( $n = 2, 3$ ) with a Lipschitz boundary. We use  $H^m(\Omega)$ ,  $m \geq 0$  to denote standard Sobolev spaces [14] based on the  $L^2$  norm ( $H^0(\Omega) = L^2(\Omega)$ ) and for a finite dimensional inner product space  $\mathbb{X}$ ,  $H^m(\mathbb{X})$  is the space of  $\mathbb{X}$ -valued functions such that each component is in  $H^m(\Omega)$ . The associated norm is denoted with  $\|\cdot\|_m$ . For  $p, q \in L^2(\mathbb{X})$  we will use  $(p, q)$  to denote the  $L^2$  inner product. We denote the spaces of all, symmetric, and skew-symmetric  $n \times n$  matrices by  $\mathbb{R}^{n \times n}$ ,  $\mathbb{R}_{\text{sym}}^{n \times n}$ , and  $\mathbb{R}_{\text{skw}}^{n \times n}$ , respectively.

We use grad and div to denote the standard gradient and divergence operators. However, we use curl to denote two different operators for different  $n$ , namely, if  $n = 2$ , then

$$\text{curl} : H^1(\mathbb{R}) \rightarrow L^2(\mathbb{R}^2), \quad \text{curl} \phi = (-\partial_{x_2} \phi \quad \partial_{x_1} \phi),$$

and if  $n = 3$ , then curl is the standard three dimensional curl operator.

By  $H(\text{div})$  we denote the space of square integrable  $\mathbb{R}^n$ -valued functions on  $\Omega$  such that the divergence of functions is also square integrable and the  $H(\text{div})$  norm is defined by  $\|\tau\|_{\text{div}}^2 = \|\tau\|_0^2 + \|\text{div} \tau\|_0^2$ . The  $H(\text{curl})$  space and  $\|\cdot\|_{\text{curl}}$  are defined similarly if  $n = 3$ . When we apply the operators grad, div, curl to a matrix-valued or vector-valued function the operations need to be well-defined as row-wise operators. By  $H(\text{div}; \mathbb{R}^n)$  we denote the space of functions in  $L^2(\mathbb{R}^{n \times n})$  such that each row is in  $H(\text{div})$ . The space  $H(\text{curl}; \mathbb{R}^3)$  is defined similarly for  $n = 3$ . The  $H(\text{div})$  and  $H(\text{curl})$  norms of the spaces are naturally defined.

**2.2. Hellinger–Reissner formulation of linear elasticity.** For a given displacement  $u : \Omega \rightarrow \mathbb{R}^n$ , the linear strain tensor  $\epsilon(u)$  is

$$\epsilon(u) = \frac{1}{2}(\text{grad} u + (\text{grad} u)^T),$$

where  $(\text{grad} u)^T$  is the transpose of  $\text{grad} u$ . From generalized Hooke's law the stress tensor is  $\sigma = C\epsilon(u)$  where  $C$  is the stiffness tensor such that  $C(x) : \mathbb{R}_{\text{sym}}^{n \times n} \rightarrow \mathbb{R}_{\text{sym}}^{n \times n}$  for all  $x \in \Omega$  and

$$c_0 \tau : \tau \leq C(x) \tau : \tau \leq c_1 \tau : \tau, \quad \tau \in \mathbb{R}_{\text{sym}}^{n \times n},$$

with positive constants  $c_0, c_1$  independent of  $x \in \Omega$ . For each  $x \in \Omega$ ,  $C(x)^{-1}$  is also bounded and positive definite. If an elastic medium is isotropic, then  $C^{-1} \tau$  has the

form

$$(2.1) \quad C^{-1}\tau = \frac{1}{2\mu} \left( \tau - \frac{\lambda}{2\mu + n\lambda} \operatorname{tr}(\tau) I \right),$$

where  $\mu(x), \lambda(x) > 0$  are the Lamé parameters,  $\operatorname{tr}(\tau)$  is the trace of  $\tau$ , and  $I$  is the identity matrix.

Throughout this paper we assume the homogeneous displacement boundary condition  $u = 0$  on  $\partial\Omega$  for simplicity. For a given  $f \in L^2(\mathbb{R}^n)$ , the Hellinger–Reissner functional  $\mathcal{J} : (H(\operatorname{div}; \mathbb{R}^n) \cap L^2(\mathbb{R}_{\operatorname{sym}}^{n \times n})) \times L^2(\mathbb{R}^n) \rightarrow \mathbb{R}$  is defined by

$$(2.2) \quad \mathcal{J}(\tau, v) = \int_{\Omega} \left( \frac{1}{2} C^{-1}\tau : \tau + \operatorname{div} \tau \cdot v - f \cdot v \right) dx,$$

and it is known that  $\mathcal{J}$  has a unique critical point

$$(\sigma, u) \in H(\Omega, \operatorname{div}; \mathbb{R}_{\operatorname{sym}}^{n \times n}) \times L^2(\Omega; \mathbb{R}^n),$$

which is the solution of the elasticity problem with the boundary condition  $u = 0$ .

For the approach with weakly imposed symmetry of stress we define  $A$  as the extension of  $C^{-1}$  on  $\mathbb{R}^{n \times n}$  such that  $A$  is the identity map for skew-symmetric matrices. We define function spaces  $\Sigma$ ,  $U$ , and  $\Gamma$  by

$$\Sigma = H(\operatorname{div}; \mathbb{R}^n), \quad U = L^2(\mathbb{R}^n), \quad \Gamma = L^2(\mathbb{R}_{\operatorname{skw}}^{n \times n}),$$

and a functional  $\tilde{\mathcal{J}} : \Sigma \times U \times \Gamma \rightarrow \mathbb{R}$  by

$$(2.3) \quad \tilde{\mathcal{J}}(\tau, v, \eta) = \int_{\Omega} \left( \frac{1}{2} A\tau : \tau + \operatorname{div} \tau \cdot v + \tau : \eta - f \cdot v \right) dx.$$

The functional  $\tilde{\mathcal{J}}$  has a unique critical point  $(\sigma, u, \gamma)$  (see [6]) and the first two components coincide with the critical point of  $\mathcal{J}$  in (2.2). By variational methods, the critical point  $(\sigma, u, \gamma)$  of  $\tilde{\mathcal{J}}$  satisfies

$$(2.4) \quad (A\sigma, \tau) + (u, \operatorname{div} \tau) + (\gamma, \tau) = 0, \quad \tau \in \Sigma,$$

$$(2.5) \quad -(\operatorname{div} \sigma, v) = (f, v), \quad v \in U,$$

$$(2.6) \quad (\sigma, \eta) = 0, \quad \eta \in \Gamma.$$

The associated discrete problem with finite element spaces  $\Sigma_h \times U_h \times \Gamma_h \subset \Sigma \times U \times \Gamma$  is seeking  $(\sigma_h, u_h, \gamma_h) \in \Sigma_h \times U_h \times \Gamma_h$  such that

$$(2.7) \quad (A\sigma_h, \tau) + (u_h, \operatorname{div} \tau) + (\gamma_h, \tau) = 0, \quad \tau \in \Sigma_h,$$

$$(2.8) \quad -(\operatorname{div} \sigma_h, v) = (f, v), \quad v \in U_h,$$

$$(2.9) \quad (\sigma_h, \eta) = 0, \quad \eta \in \Gamma_h.$$

In this approach the numerical stress  $\sigma_h$  is not symmetric but is weakly symmetric due to the last equation of the above.

### 3. ABSTRACT FRAMEWORK

In this section we introduce an abstract framework for unified analysis of weak symmetry elements. This is a generalization of the approach in [20] with inspirations from [9, 18].

Throughout this paper  $c$  is a generic positive constant independent of mesh sizes. We first recall the Babuška–Brezzi stability conditions for (2.7–2.9), which are

(S1) There is  $c$  such that

$$c\|\tau\|_{\text{div}}^2 \leq (A\tau, \tau),$$

for  $\tau \in \Sigma_h$  satisfying  $(\text{div } \tau, v) + (\tau, \eta) = 0$  for all  $(v, \eta) \in U_h \times \Gamma_h$ .

(S2) There is  $c$  such that

$$\inf_{0 \neq (v, \eta) \in U_h \times \Gamma_h} \sup_{0 \neq \tau \in \Sigma_h} \frac{(\text{div } \tau, v) + (\tau, \eta)}{\|\tau\|_{\text{div}}(\|v\|_0 + \|\eta\|_0)} \geq c.$$

Now we recall a commuting diagram of the elasticity complex in [7]. Let

$$\Xi = \begin{cases} H^1(\mathbb{R}^2) & \text{if } n = 2, \\ H(\text{curl}; \mathbb{R}^3) & \text{if } n = 3. \end{cases}$$

We also define  $S$  and  $\chi$  as

$$S \begin{pmatrix} \xi_1 \\ \xi_2 \end{pmatrix} = \frac{1}{2} \begin{pmatrix} \xi_1 & \xi_2 \end{pmatrix} \quad \text{for } \xi \in \Xi, \quad \chi(r) = \begin{pmatrix} 0 & r \\ -r & 0 \end{pmatrix} \quad \text{for } r \in \mathbb{R} \quad \text{if } n = 2,$$

$$S\xi = \frac{1}{2}(\xi^T - (\text{tr } \xi)I) \quad \text{for } \xi \in \Xi, \quad \chi \begin{pmatrix} r_1 \\ r_2 \\ r_3 \end{pmatrix} = \begin{pmatrix} 0 & -r_3 & r_2 \\ r_3 & 0 & -r_1 \\ -r_2 & r_1 & 0 \end{pmatrix} \quad \text{if } n = 3.$$

Note that  $S$  and  $\chi$  are invertible algebraic operators. One can verify by a direct computation that  $S$  maps  $\Xi$  to  $H(\text{div})$  if  $n = 2$ , and to  $H(\text{div}; \mathbb{R}^3)$  if  $n = 3$ , so  $\chi \text{div } S$  maps  $\Xi$  to  $\Gamma$ . One can also verify by a direct computation that

$$\text{skw curl } \xi = \chi \text{div } S\xi, \quad \xi \in \Xi,$$

where  $\text{skw } \tau = (\tau + \tau^T)/2$  for  $\tau \in L^2(\Omega; \mathbb{R}^{n \times n})$ . For finite element spaces  $\Xi_h \subset \Xi$ ,  $\Gamma_h \subset \Gamma$ , and the  $L^2$  projection  $Q_h$  into  $\Gamma_h$ , it holds that

$$Q_h \text{skw curl } \xi = Q_h \chi \text{div } S\xi, \quad \xi \in \Xi_h,$$

which implies that the triangle in the following diagram commutes. Note that the

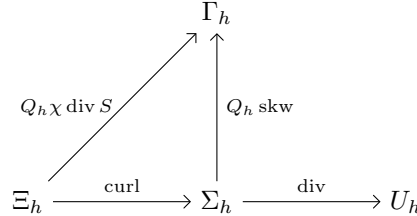


FIGURE 1. A finite element version of the commuting diagram in the weakly symmetric elasticity complex

bottom row of this diagram is not necessarily an exact sequence.

**Definition 3.1.** A triple of finite elements  $(\Sigma_h, U_h, R_h)$  is elasticity stable if  $\Sigma_h \subset \Sigma$ ,  $U_h \subset U$ ,  $R_h \subset \Gamma$  and the following hold:

(A1)  $\text{div } \Sigma_h = U_h$

(A2) There exists  $c$  such that for any  $(v, \rho) \in U_h \times R_h$ , there exists  $\tau \in \Sigma_h$  satisfying

$$\text{div } \tau = v, \quad (\tau, \rho') = (\rho, \rho') \quad \forall \rho' \in R_h, \quad \|\tau\|_{\text{div}} \leq c(\|v\|_0 + \|\rho\|_0).$$

It is not difficult to see that **(A1)** implies **(S1)** because  $A$  is positive definite, and **(A2)** implies **(S2)**. However, an elasticity stable triple  $(\Sigma_h, U_h, R_h)$  may not be an appropriate mixed finite element for (2.7–2.9) in the standard context. For instance, if  $\Sigma_h = BDM_1(\mathbb{R}^2)$ ,  $U_h = \mathcal{P}_0^d(\mathbb{R}^2)$ , then  $(\Sigma_h, U_h, 0)$  is elasticity stable. Thus an elasticity stable triple  $(\Sigma_h, U_h, \Gamma_h)$  should have a reasonable order of approximation of  $(\Sigma, U, \Gamma)$  to be an appropriate mixed finite element for (2.7–2.9).

**Definition 3.2.** Suppose that  $\Xi_h \subset \Xi$ ,  $R_h \subset \Gamma$  are finite element spaces and  $Q_h : \Gamma \rightarrow R_h$  is the  $L^2$  projection. The pair  $(\Xi_h, R_h)$  is Stokes stable if

**(B)** There exists  $c$  such that for any  $\rho \in R_h$  there is  $\xi \in \Xi_h$  satisfying

$$(Q_h \chi \operatorname{div} S\xi, \rho) \geq c \|\rho\|_0^2, \quad \|\operatorname{curl} \xi\|_0 \leq c \|\rho\|_0.$$

The condition **(B)** implies that  $Q_h \chi \operatorname{div} S : \Xi_h \rightarrow \Gamma_h$  is surjective. Furthermore, for any  $\rho \in R_h$ , there exists  $\xi \in \Xi_h$  such that  $Q_h \chi \operatorname{div} S\xi = \rho$  and  $\|\operatorname{curl} \xi\|_0 \leq c \|\rho\|_0$ .

**Theorem 3.3.** Let  $\Xi_h \subset \Xi$ ,  $\Sigma_h \subset \Sigma$ ,  $U_h \subset U$ ,  $\Gamma_h \subset \Gamma$  be four finite element spaces. For a subspace  $\Gamma_h^0$  of  $\Gamma_h$  its orthogonal complement is denoted by  $\Gamma_h^1$ . Suppose that  $(\Sigma_h, U_h, \Gamma_h^0)$  is elasticity stable and  $(\Xi_h, \Gamma_h^1)$  is Stokes stable with  $\chi \operatorname{div} S\Xi_h \perp \Gamma_h^0$ . Then  $(\Sigma_h, U_h, \Gamma_h)$  is elasticity stable.

*Proof.* Since we assume that  $(\Sigma_h, U_h, \Gamma_h^0)$  is elasticity stable, we only need to check **(A2)** to show that  $(\Sigma_h, U_h, \Gamma_h)$  is elasticity stable. Let  $v \in U_h$ ,  $\eta = \eta_0 + \eta_1 \in \Gamma_h^0 \oplus \Gamma_h^1$  be given. By **(A2)** for  $(\Sigma_h, U_h, \Gamma_h^0)$  there is  $\tau_0 \in \Sigma_h$  such that

$$\|\tau_0\|_{\operatorname{div}} \leq c(\|v\|_0 + \|\eta\|_0), \quad \operatorname{div} \tau_0 = v, \quad (\tau_0, \eta'_0) = (\eta_0, \eta'_0), \quad \forall \eta'_0 \in \Gamma_h^0.$$

Let  $Q_h^1$  be the  $L^2$  projection into  $\Gamma_h^1$ . Since  $(\Xi_h, \Gamma_h^1)$  is Stokes stable, there exists  $\xi \in \Xi_h$  such that  $Q_h^1 \chi \operatorname{div} S\xi = \eta_1 - Q_h^1 \tau_0$  and  $\|\operatorname{curl} \xi\|_0 \leq c\|\eta_1 - Q_h^1 \tau_0\|_0$ . We take  $\tau = \tau_0 + \operatorname{curl} \xi$  and check that the conditions in **(A2)** holds. Since  $\operatorname{div} \tau_0 = v$ ,

$$\operatorname{div} \tau = \operatorname{div} \tau_0 + \operatorname{div} \operatorname{curl} \xi = \operatorname{div} \tau_0 = v.$$

For  $\eta' = \eta'_0 + \eta'_1 \in \Gamma_h^0 \oplus \Gamma_h^1$

$$\begin{aligned} (\tau, \eta') &= (\tau_0 + \operatorname{skw} \operatorname{curl} \xi, \eta') \\ &= (\tau_0 + \chi \operatorname{div} S\xi, \eta') \\ &= (\tau_0, \eta'_0) + (\tau_0 + \chi \operatorname{div} S\xi, \eta'_1) \quad (\because \chi \operatorname{div} S\xi \perp \eta'_0) \\ &= (\tau_0, \eta'_0) + (Q_h^1(\tau_0 + \chi \operatorname{div} S\xi), \eta'_1) \\ &= (\eta_0, \eta'_0) + (\eta_1, \eta'_1) \quad (\because Q_h^1 \chi \operatorname{div} S\xi = \eta_1 - Q_h^1 \tau_0) \\ &= (\eta, \eta'). \end{aligned}$$

Note that  $\|\eta_1\|_0 \leq \|\eta\|_0$ . By the triangle inequality,  $\|\tau_0\|_{\operatorname{div}} \leq c(\|v\|_0 + \|\eta\|_0)$  and  $\|\operatorname{curl} \xi\|_0 \leq c\|\eta_1 - Q_h^1 \tau_0\|_0 \leq c(\|\eta\|_0 + \|\tau_0\|_0)$ ,

$$\|\tau\|_{\operatorname{div}} \leq \|\operatorname{curl} \xi\|_0 + \|\tau_0\|_{\operatorname{div}} \leq c(\|v\|_0 + \|\eta\|_0).$$

Thus **(A2)** holds and  $(\Sigma_h, U_h, \Gamma_h)$  is elasticity stable.  $\square$

Now we show an improved a priori error estimate.

**Theorem 3.4.** Suppose that  $(\Sigma_h, U_h, \Gamma_h)$  is elasticity stable and there exists an interpolation operator  $\Pi_h : H^1(\Omega, \mathbb{R}^{n \times n}) \rightarrow \Sigma_h$  such that

$$\operatorname{div} \Pi_h \tau = P_h \operatorname{div} \tau, \quad \tau \in H^1(\Omega; \mathbb{R}^{n \times n}),$$

where  $P_h$  is the  $L^2$  projection into  $U_h$ . Let  $(\sigma, u, \gamma)$  and  $(\sigma_h, u_h, \gamma_h)$  be the solutions of (2.4–2.6) and (2.7–2.9). Then the following inequality

$$(3.1) \quad \|\sigma - \sigma_h\|_0 + \|P_h u - u_h\|_0 + \|\gamma - \gamma_h\|_0 \leq c(\|\sigma - \Pi_h \sigma\|_0 + \|\gamma - Q_h \gamma\|_0),$$

holds with  $Q_h$  the  $L^2$  projection into  $\Gamma_h$ .

*Proof.* The proof is same as that of Theorem 4.1 in [20] but we include the details here to be self-contained. The difference of (2.4–2.6) and (2.7–2.9) gives

$$(3.2) \quad (A(\sigma - \sigma_h), \tau) + (u - u_h, \operatorname{div} \tau) + (\gamma - \gamma_h, \tau) = 0, \quad \tau \in \Sigma_h,$$

$$(3.3) \quad (\operatorname{div}(\sigma - \sigma_h), v) = 0, \quad v \in U_h,$$

$$(3.4) \quad (\sigma - \sigma_h, \eta) = 0, \quad \eta \in \Gamma_h.$$

Let  $\Sigma_{h,0} = \{\tau \in \Sigma_h : \operatorname{div} \tau = 0\}$  and consider an auxiliary problem of seeking  $(\sigma'_h, \gamma'_h) \in \Sigma_{h,0} \times \Gamma_h$  such that

$$(3.5) \quad (A\sigma'_h, \tau) + (\gamma'_h, \tau) + (\sigma'_h, \eta) = F(\tau) + G(\eta), \quad (\tau, \eta) \in \Sigma_{h,0} \times \Gamma_h,$$

with a bounded linear functional  $(F, G)$  on  $\Sigma_{h,0} \times \Gamma_h$ . As a special case of **(A2)**, for  $v = 0$  and any given  $\eta \in \Gamma_h$  there exists  $\tau \in \Sigma_{h,0}$  such that  $(\tau, \eta') = (\eta, \eta')$  for all  $\eta' \in \Gamma_h$  and  $\|\tau\|_0 \leq c\|\eta\|_0$ . From this observation and **(A1)**,  $\Sigma_{h,0} \times \Gamma_h$  is a stable mixed finite element for the problem (3.5) with the  $L^2$  norms. By restricting  $\tau \in \Sigma_{h,0}$ , the sum of (3.2) and (3.4) is

$$(A(\sigma - \sigma_h), \tau) + (\gamma - \gamma_h, \tau) + (\sigma - \sigma_h, \eta) = 0,$$

which is equivalent to

$$(3.6) \quad (A(\sigma_h - \Pi_h \sigma), \tau) + (\gamma_h - Q_h \gamma, \tau) + (\sigma_h - \Pi_h \sigma, \eta) \\ = (A(\sigma - \Pi_h \sigma), \tau) + (\gamma - Q_h \gamma, \tau) + (\sigma - \Pi_h \sigma, \eta).$$

Note that  $\sigma_h - \Pi_h \sigma \in \Sigma_{h,0}$  because  $\operatorname{div} \sigma_h = P_h \operatorname{div} \sigma = \operatorname{div} \Pi_h \sigma$  by (3.3) and  $\operatorname{div} \Sigma_h = U_h$ . By the Babuška–Brezzi stability of (3.5), there exists  $(\tau, \eta) \in \Sigma_{h,0} \times \Gamma_h$  such that  $\|\tau\|_0 + \|\eta\|_0 \leq c$  and

$$\|\sigma_h - \Pi_h \sigma\|_0 + \|\gamma_h - Q_h \gamma\|_0 \leq (A(\sigma_h - \Pi_h \sigma), \tau) + (\gamma_h - Q_h \gamma, \tau) + (\sigma_h - \Pi_h \sigma, \eta).$$

Combining this, (3.6), and the Cauchy–Schwarz inequality with  $\|\tau\|_0 + \|\eta\|_0 \leq c$ ,

$$\|\sigma_h - \Pi_h \sigma\|_0 + \|\gamma_h - Q_h \gamma\|_0 \leq (A(\sigma - \Pi_h \sigma), \tau) + (\gamma - Q_h \gamma, \tau) + (\sigma - \Pi_h \sigma, \eta) \\ \leq c(\|\sigma - \Pi_h \sigma\|_0 + \|\gamma - Q_h \gamma\|_0).$$

By the triangle inequality and the above one,

$$\|\sigma - \sigma_h\|_0 + \|\gamma - \gamma_h\|_0 \leq \|\sigma - \Pi_h \sigma\|_0 + \|\Pi_h \sigma - \sigma_h\|_0 + \|\gamma - Q_h \gamma\|_0 + \|Q_h \gamma - \gamma_h\|_0 \\ \leq c(\|\sigma - \Pi_h \sigma\|_0 + \|\gamma - Q_h \gamma\|_0),$$

so (3.1) for  $\|\sigma - \sigma_h\|_0$  and  $\|\gamma - \gamma_h\|_0$  is proved. To estimate  $\|u_h - P_h u\|_0$ , observe that (3.2) gives

$$(3.7) \quad (A(\sigma - \sigma_h), \tau) + (P_h u - u_h, \operatorname{div} \tau) + (\gamma - \gamma_h, \tau) = 0, \quad \tau \in \Sigma_h,$$

because  $\operatorname{div} \tau \in U_h$  is orthogonal to  $u - P_h u$ . By **(A2)** there is  $\tau$  in (3.7) such that  $\operatorname{div} \tau = P_h u - u_h$  and  $\|\tau\|_{\operatorname{div}} \leq c\|P_h u - u_h\|_0$ . Then we have

$$\|P_h u - u_h\|_0^2 = -(A(\sigma - \sigma_h), \tau) - (\gamma - \gamma_h, \tau) \\ \leq c(\|\sigma - \sigma_h\|_0 + \|\gamma - \gamma_h\|_0)\|P_h u - u_h\|_0.$$

Combining the result with the estimates of  $\|\sigma - \sigma_h\|_0$  and  $\|\gamma - \gamma_h\|_0$ , we have

$$\|P_h u - u_h\|_0 \leq c(\|\sigma - \Pi_h \sigma\|_0 + \|\gamma - Q_h \gamma\|_0),$$

as desired.  $\square$

If  $\Sigma_h$  and  $\Gamma_h$  provide higher order approximations than that of  $U_h$ , then (3.1) implies that  $\|P_h u - u_h\|_0$  is superconvergent. A local post-processing can be used to get a new numerical solution  $u_h^*$  such that the convergence rate of  $\|u - u_h^*\|_0$  is as good as that of  $\|\sigma - \sigma_h\|_0 + \|\gamma - \gamma_h\|_0$ . A higher order superconvergence of  $\|P_h u - u_h\|_0$  can be obtained by the Aubin–Nitsche duality argument and the elliptic regularity of  $\Omega$  when  $f \in U_h$ . In this case, a higher order local post-processing can be used to obtain  $u_h^{**}$  which is a higher order approximation of  $u$  in  $L^2(\mathbb{R}^n)$ . A careful discussion can be found in [20] for second order rectangular elements. It is straightforward to generalize the argument in [20] to higher order elements.

#### 4. EXAMPLES

In this section we show examples which can be analyzed by the abstract framework.

By  $\mathcal{T}_h$  we denote a shape-regular mesh of  $\Omega$  and  $h$  is the maximum diameter of the elements in  $\mathcal{T}_h$ . By  $\mathcal{P}_k(D)$  and  $\mathcal{P}_k(D; \mathbb{X})$ , we denote the spaces of  $\mathbb{R}$  and  $\mathbb{X}$ -valued polynomials of degree  $\leq k$  on  $D \subset \Omega$ . For a rectangle  $D$ ,  $\mathcal{Q}_k(D)$  is the space of polynomials of degree at most  $k$  in each variable  $x_i$ ,  $1 \leq i \leq n$ . Now we define

$$\begin{aligned} \mathcal{P}_k^d(\mathbb{X}) &= \{p \in L^2(\mathbb{X}) \mid p|_T \in \mathcal{P}_k(T; \mathbb{X}), \quad T \in \mathcal{T}_h\}, & k \geq 0, \\ \mathcal{P}_k^c(\mathbb{X}) &= \{p \in H^1(\mathbb{X}) \mid p|_T \in \mathcal{P}_k(T; \mathbb{X}), \quad T \in \mathcal{T}_h\}, & k \geq 1, \\ \mathcal{Q}_k^d(\mathbb{X}) &= \{p \in L^2(\mathbb{X}) \mid p|_T \in \mathcal{Q}_k(T; \mathbb{X}), \quad T \in \mathcal{T}_h\}, & k \geq 0, \\ \mathcal{Q}_k^c(\mathbb{X}) &= \{p \in H^1(\mathbb{X}) \mid p|_T \in \mathcal{Q}_k(T; \mathbb{X}), \quad T \in \mathcal{T}_h\}, & k \geq 1, \\ RTN_k &= \{p \in H(\text{div}) \mid p|_T \in \mathcal{P}_{k-1}(T; \mathbb{R}^n) + \mathbf{x}\mathcal{P}_{k-1}(T)\}, & k \geq 1, \\ BDM_k &= \{p \in H(\text{div}) \mid p|_T \in \mathcal{P}_k(T; \mathbb{R}^n)\}, & k \geq 1, \end{aligned}$$

where  $\mathbf{x}$  is the vector function  $(x_1, \dots, x_n)$ . Note that the lowest order RTN element is denoted by  $RTN_1$  in this paper, which is different from [11]. The rectangular RTN and BDM elements [11] are denoted by  $rRTN_k$  and  $rBDM_k$  with  $k \geq 1$ . We also define  $RTN_k(\mathbb{R}^n)$  and  $BDM_k(\mathbb{R}^n)$  as the subspaces of  $H(\text{div}; \mathbb{R}^n)$  such that each row of an element in those spaces is in  $RTN_k$  and  $BDM_k$ , respectively. Throughout this section  $\Sigma_h$ ,  $U_h$ ,  $\Gamma_h$ ,  $\Xi_h$  are the finite element spaces in Figure 1.

##### 4.1. Elements with continuous $\Gamma_h$ .

4.1.1. *PEERS*. Let  $b_T$  be the standard cubic bubble function on a triangle  $T \in \mathcal{T}_h$  and

$$B = \{\xi \mid \xi|_T = p b_T, p \in \mathbb{R}^2, T \in \mathcal{T}_h\}.$$

The PEERS [3] is

$$\Sigma_h = RTN_1(\mathbb{R}^2) + \text{curl } B, \quad U_h = \mathcal{P}_0^d(\mathbb{R}^2), \quad \Gamma_h = \mathcal{P}_1^c(\mathbb{R}^{2 \times 2}_{\text{skw}}).$$

Let  $\Gamma_h^0 = 0$  and  $\Xi_h = \mathcal{P}_1^c(\mathbb{R}^2) + B$ . It is not difficult to see that  $(\Sigma_h, U_h, 0)$  is elasticity stable from the stability of  $(RTN_1, \mathcal{P}_0)$  for the mixed Poisson equation. Moreover, the stability of the MINI element for the Stokes equation [4] implies

that  $(\Xi_h, \Gamma_h)$  is Stokes stable because  $S$  and  $\chi$  are algebraic isomorphisms. Thus  $(\Sigma_h, U_h, \Gamma_h)$  is elasticity stable.

4.1.2. *Taylor–Hood based elements.* Let

$$\Sigma_h = BDM_k(\mathbb{R}^n), \quad U_h = \mathcal{P}_{k-1}^d(\mathbb{R}^n), \quad \Gamma_h = \mathcal{P}_k^c(\mathbb{R}_{\text{skw}}^{n \times n}),$$

for  $k \geq 1$  and take  $\Gamma_h^0 = 0$ ,

$$\Xi_h = \begin{cases} \mathcal{P}_{k+1}^c(\mathbb{R}^2), & \text{if } n = 2, \\ \mathcal{P}_{k+1}^c(\mathbb{R}^{3 \times 3}), & \text{if } n = 3. \end{cases}$$

By definition,  $(\Sigma_h, U_h, 0)$  is elasticity stable. Moreover, the stability of Taylor–Hood elements for the Stokes equation yields that  $(\Xi_h, \Gamma_h)$  is Stokes stable. Thus  $(\Sigma_h, U_h, \Gamma_h)$  is elasticity stable.

In the two dimensional case these elements were noticed in [15]. In the three dimensional case it seems that the same elements have not appeared in the literature but similar elements were proposed in [9] with slightly larger space for  $\Sigma_h$  using the Raviart–Thomas–Nédélec spaces.

4.1.3. *Two dimensional rectangular element.* Let  $\Omega$  be a bounded two dimensional domain with a rectangular mesh  $\mathcal{T}_h$ . Here we propose a rectangular version of the Taylor–Hood based elements in two dimensions. Let  $S_2$  be the serendipity element with 8 local degrees of freedom [10] and set

$$\Sigma_h = rBDM_1(\mathbb{R}^2), \quad U_h = \mathcal{P}_0^d(\mathbb{R}^2), \quad \Gamma_h = Q_1^c(\mathbb{R}_{\text{skw}}^{2 \times 2}),$$

and let  $\Xi_h = S_2(\mathbb{R}^2)$ ,  $\Gamma_h^0 = 0$ . Since  $(rBDM_1, \mathcal{P}_0^d)$  is stable for mixed Poisson equation with  $\text{div } rBDM_1 = \mathcal{P}_0^d$ ,  $(\Sigma_h, U_h, 0)$  is elasticity stable. It is also known that  $(\Xi_h, \Gamma_h)$  is Stokes stable [23], so  $(\Sigma_h, U_h, \Gamma_h)$  is elasticity stable.

4.2. **Elements with discontinuous  $\Gamma_h$ .** To apply the framework for elements with discontinuous  $\Gamma_h$  we need preliminary results.

**Lemma 4.1.** *Suppose that  $\Sigma_h = BDM_1(\mathbb{R}^n)$ ,  $U_h = 0$ , and  $\Gamma_h = \mathcal{P}_0^d(\mathbb{R}_{\text{skw}}^{n \times n})$ . Then  $(\Sigma_h, U_h, \Gamma_h)$  is elasticity stable.*

A proof of this lemma can be found in [7, 9, 18]. For a simple proof we refer to Proposition 2.10 in [18]. From this lemma the following can be easily obtained.

**Corollary 4.2.** *The two triples*

$$(BDM_k(\mathbb{R}^n), \mathcal{P}_{k-1}^d(\mathbb{R}^n), \mathcal{P}_0^d(\mathbb{R}_{\text{skw}}^{n \times n})) \quad \text{and} \quad (RTN_{k+1}(\mathbb{R}^n), \mathcal{P}_k^d(\mathbb{R}^n), \mathcal{P}_0^d(\mathbb{R}_{\text{skw}}^{n \times n}))$$

*are elasticity stable for  $k \geq 1$ .*

Let  $b_T$  be the standard cubic/quartic bubble function on a triangle/tetrahedron  $T \in \mathcal{T}_h$ . In the two and three dimensional cases  $B_k \subset \Xi$ ,  $k \geq 1$  is defined by

$$B_k = \begin{cases} \{\eta \in L^2(\Omega; \mathbb{R}^2) \mid \eta|_T \in b_T \mathcal{P}_{k-1}(T; \mathbb{R}^2)\}, & \text{if } n = 2, \\ \{\eta \in L^2(\Omega; \mathbb{R}^{3 \times 3}) \mid \eta|_T \in b_T \mathcal{P}_{k-1}(T; \mathbb{R}^{3 \times 3})\}, & \text{if } n = 3. \end{cases}$$

When  $n = 3$  a matrix bubble function  $\mathbf{b}_T$  on each  $T \in \mathcal{T}_h$  is defined by

$$\mathbf{b}_T = \sum_{i=0}^4 \lambda_i \lambda_{i+1} \lambda_{i+2} (\text{grad } \lambda_{i+3})^T (\text{grad } \lambda_{i+3}),$$



where  $\lambda_i$ ,  $i = 0, 1, 2, 3$  are the barycentric coordinates on  $T$ ,  $\text{grad } \lambda_i$  is a row vector, and the index  $i$  is counted modulo 4. One can see that  $\mathbf{b}_T$  is symmetric positive definite and the cross product of each row of  $\mathbf{b}_T$  and the unit normal vector  $n_e$  on an edge/face  $e \subset \partial T$  vanishes. By the integration by parts,

$$(4.1) \quad (\text{curl}(\mathbf{b}_T \text{curl } \eta_1), \eta_2) = (\mathbf{b}_T \text{curl } \eta_1, \text{curl } \eta_2), \quad \eta_1, \eta_2 \in \mathcal{P}_k(T; \mathbb{R}_{\text{skw}}^{n \times n}),$$

so the above relation gives an inner product on

$$\hat{\mathcal{P}}_k^d(\mathbb{R}_{\text{skw}}^{n \times n}) = \{\tau \in \mathcal{P}_k^d(\mathbb{R}_{\text{skw}}^{n \times n}) \mid \tau \perp \mathcal{P}_0^d(\mathbb{R}_{\text{skw}}^{n \times n})\}.$$

Moreover, the norm given by this inner product with weight  $\mathbf{b}_T$  is equivalent to the standard  $L^2$  norm on  $T$  up to constants independent of the diameter of  $T$ . For  $k \geq 1$  let

$$(4.2) \quad B(\eta) = \begin{cases} h_T^{-2} b_T \text{rot } \eta & \text{for } \eta \in \hat{\mathcal{P}}_k(T; \mathbb{R}_{\text{skw}}^{2 \times 2}) \quad \text{if } n = 2, \\ h_T^{-2} \mathbf{b}_T \text{curl } \eta & \text{for } \eta \in \hat{\mathcal{P}}_k(T; \mathbb{R}_{\text{skw}}^{3 \times 3}) \quad \text{if } n = 3, \end{cases}$$

and define  $\hat{B}_k$  as

$$\hat{B}_k = \{\xi \in \Xi : \xi|_T = B(\eta) \text{ for some } \eta \in \hat{\mathcal{P}}_k(T; \mathbb{R}_{\text{skw}}^{n \times n})\}.$$

**Lemma 4.3.** *For  $k \geq 1$  the pairs  $(\hat{B}_k, \hat{\mathcal{P}}_k^d(\mathbb{R}_{\text{skw}}^{n \times n}))$  and  $(B_k, \hat{\mathcal{P}}_k^d(\mathbb{R}_{\text{skw}}^{n \times n}))$  are Stokes stable.*

*Proof.* We prove only the three dimensional case for both pairs because the two dimensional case is similar.

In the case of  $(\hat{B}_k, \hat{\mathcal{P}}_k^d(\mathbb{R}_{\text{skw}}^{n \times n}))$  pair, for a given  $\eta \in \hat{\mathcal{P}}_k(\mathbb{R}_{\text{skw}}^{n \times n})$ , take  $\xi \in \hat{B}_k$  such that  $\xi|_T = h_T^{-2} \mathbf{b}_T \text{curl } \eta|_T$  for  $T \in \mathcal{T}_h$ . Then

$$(Q_h \chi \text{div } S\xi, \eta)_T = (Q_h \text{skw curl } \xi, \eta)_T = (\text{curl } \xi, \eta)_T = h_T^{-2} (\mathbf{b}_T \text{curl } \eta, \text{curl } \eta)_T.$$

By the standard scaling argument

$$\|\eta\|_{0,T} \sim h_T \|\text{curl } \eta\|_{0,T}, \quad \|\text{curl } \xi\|_{0,T} \sim h_T^{-1} \|\text{curl } \eta\|_{0,T} \sim \|\eta\|_{0,T},$$

so  $\|\text{curl } \xi\|_0 \leq c \|\eta\|_0$  and  $(Q_h \chi \text{div } S\xi, \eta) \geq c \|\eta\|_0^2$ .

With the  $(B_k, \hat{\mathcal{P}}_k^d(\mathbb{R}_{\text{skw}}^{n \times n}))$  pair and a given  $\eta \in \hat{\mathcal{P}}_k(\mathbb{R}_{\text{skw}}^{n \times n})$ , one can take  $\xi \in B_k$  such that  $\xi|_T = h_T^{-2} b_T \text{curl } \eta|_T$  for  $T \in \mathcal{T}_h$ . The rest of the argument is then similar to the first case, so we omit the details.  $\square$

We are now ready to present examples with discontinuous  $\Gamma_h$ .

**4.2.1. The Cockburn–Gopalakrishnan–Guzmán (CGG) and Arnold–Falk–Winther (AFW) elements.** The CGG elements [12] are

$$\Sigma_h = RTN_k(\mathbb{R}^n) + \text{curl } \hat{B}_{k-1}, \quad U_h = \mathcal{P}_{k-1}^d(\mathbb{R}^n), \quad \Gamma_h = \mathcal{P}_{k-1}^d(\mathbb{R}_{\text{skw}}^{n \times n}),$$

for  $k \geq 2$ . To apply the presented framework let  $\Gamma_h^0 = \mathcal{P}_0^d(\mathbb{R}_{\text{skw}}^{n \times n})$  and  $\Xi_h = \hat{B}_{k-1}$ . Then  $(\hat{B}_k, \hat{\mathcal{P}}_k^d(\mathbb{R}_{\text{skw}}^{n \times n}))$  is Stokes stable by Lemma 4.3 and  $(\Sigma_h, U_h, \Gamma_h^0)$  is elasticity stable by Corollary 4.2. Moreover,  $Q_h \chi \text{div } S\xi = Q_h \text{skw curl } \xi$  for  $\xi \in \Xi_h$  is orthogonal to  $\Gamma_h^0$  by the definition of  $\hat{B}_k$  and (4.1), so  $(\Sigma_h, U_h, \Gamma_h)$  is elasticity stable.

The AFW elements [7] are

$$\Sigma_h = BDM_k(\mathbb{R}^n), \quad U_h = \mathcal{P}_{k-1}^d(\mathbb{R}^n), \quad \Gamma_h = \mathcal{P}_{k-1}^d(\mathbb{R}_{\text{skw}}^{n \times n}),$$

for  $k \geq 1$ . The stability of these elements for  $k = 1$  follows as a corollary of Lemma 4.1. For  $k \geq 2$  it follows from the stability of CGG elements because  $\text{curl } \hat{B}_{k-1} \subset \mathcal{P}_k^d(\mathbb{R}^{n \times n}) \cap H(\text{div}; \mathbb{R}^n)$  and then  $RTN_k(\mathbb{R}^n) + \text{curl } \hat{B}_{k-1} \subset BDM_k(\mathbb{R}^n)$ .

4.2.2. *Gopalakrishnan–Guzmán (GG) and Stenberg elements.* The GG elements [17] are

$$\Sigma_h = BDM_k(\mathbb{R}^n) + \text{curl } \hat{B}_k, \quad U_h = \mathcal{P}_{k-1}^d(\mathbb{R}^n), \quad \Gamma_h = \mathcal{P}_k^d(\mathbb{R}_{\text{skw}}^{n \times n}),$$

for  $k \geq 1$ . Let  $\Gamma_h^0 = \mathcal{P}_0^d(\mathbb{R}_{\text{skw}}^{n \times n})$  and  $\Xi_h = \hat{B}_k$ . Then  $(\Xi_h, \Gamma_h^1)$  is Stokes stable by Lemma 4.3 and  $(\Sigma_h, U_h, \Gamma_h^0)$  is elasticity stable by Corollary 4.2. Moreover,  $Q_h \chi \text{div } S \Xi_h$  is orthogonal to  $\Gamma_h^0$ , so  $(\Sigma_h, U_h, \Gamma_h)$  is elasticity stable. By checking the degree of polynomials one can see that  $\text{curl } \hat{B}_{k-1} \subset BDM_k(\mathbb{R}^n)$  holds, so only a small part of  $\text{curl } \hat{B}_k$  is necessary for  $\Sigma_h$ . See [17] for degrees of freedom of  $\Sigma_h$  for implementation.

The Stenberg elements [25] are

$$\Sigma_h = BDM_k(\mathbb{R}^n) + \text{curl } B_k, \quad U_h = \mathcal{P}_{k-1}^d(\mathbb{R}^n), \quad \Gamma_h = \mathcal{P}_k^d(\mathbb{R}_{\text{skw}}^{n \times n}),$$

for  $k \geq 1$ . The stability can be proved in a way similar to the GG elements by Lemma 4.3 and Corollary 4.2. As noticed in [18], the  $k = 1$  case, which is not included in [25], also gives a stable mixed method.

4.2.3. *Elements with barycentric subdivision grids.* Let  $\mathcal{M}_h$  be a shape-regular mesh of  $\Omega$  and  $\mathcal{T}_h$  be the mesh obtained by dividing each element in  $\mathcal{M}_h$  into  $n + 1$  subelements by connecting the vertices of the element to its barycenter. Define

$$\Sigma_h = BDM_k(\mathbb{R}^n), \quad U_h = \mathcal{P}_{k-1}^d(\mathbb{R}^n), \quad \Gamma_h = \mathcal{P}_k^d(\mathbb{R}_{\text{skw}}^{n \times n}),$$

for  $k \geq 1$ . As is noticed in [17], this triple is elasticity stable when  $n = 2$  and when  $n = 3$  with  $k \geq 2$  due to the stable finite elements for the Stokes equation in [22, 26]. Thus we only consider  $n = 3$ ,  $k = 1$  case which is not known previously. For  $M \in \mathcal{M}_h$  we define

$$\mathcal{P}_{2,0}^c(M) = \{\xi \in H^1(M; \mathbb{R}^{3 \times 3}) : \xi|_T \in \mathcal{P}_2(T; \mathbb{R}^{3 \times 3}), \xi|_{\partial M} = 0 \text{ for } T \subset M, T \in \mathcal{T}_h\},$$

and

$$\Xi_h = \{\xi \in \Xi : \xi|_M \in \mathcal{P}_{2,0}^c(M) \text{ for } M \in \mathcal{M}_h\},$$

$$\Gamma_h^0 = \{\eta \in \Gamma : \eta|_M = \mathcal{P}_0^d(M; \mathbb{R}_{\text{skw}}^{n \times n}) \text{ for } M \in \mathcal{M}_h\}.$$

Then  $(\Xi_h, \Gamma_h^1)$  is Stokes stable by Lemma 2 in [26]. Furthermore,  $\text{curl } \Xi_h$  is orthogonal to  $\Gamma_h^0$  due to the integration by parts because  $\xi \in \Xi_h$  is a bubble-like function on each macroelement  $M \in \mathcal{M}_h$ . Since  $(\Sigma_h, U_h, \Gamma_h^0)$  is elasticity stable, so is  $(\Sigma_h, U_h, \Gamma_h)$ .

4.2.4. *Rectangular elements with discontinuous  $\Gamma_h$ .* There are not much literature on rectangular/quadrilateral weak symmetry elements with discontinuous  $\Gamma_h$  [8, 20, 19, 2]. Unfortunately the properties which are crucial in our framework do not hold on quadrilateral meshes. For instance, **(A1)** fails in general for quadrilateral  $H(\text{div})$  and  $L^2$  element pairs, and there is no interpolation operator  $\Pi_h$  as in Theorem 3.4. Nonetheless, a low order element and a family of elements on quadrilateral meshes, say the Arnold–Awanou–Qiu elements, have been developed in [2].

TABLE 1. Convergence rates of errors and orders of approximation of finite elements for some weak symmetry elements.

elements	convergence rate [order of approximation]			order	$n$
	$\ \sigma - \sigma_h\ _0$	$\ P_h u - u_h\ _0$	$\ \gamma - \gamma_h\ _0$		
PEERS (4.1.1)	1 [1]	1 [1]	1 [1]	1	2
THB (4.1.2)	$k$ [ $k$ ]	$k$ [ $k-1$ ]	$k$ [ $k$ ]	$k \geq 2$	2/3
2D Rect. (4.1.3)	2 [2]	2 [1]	2 [2]	2	2
AFW (4.2.1)	$k$ [ $k+1$ ]	$k$ [ $k$ ]	$k$ [ $k$ ]	$k \geq 1$	2/3
CGG (4.2.1)	$k$ [ $k$ ]	$k$ [ $k$ ]	$k$ [ $k$ ]	$k \geq 2$	2/3
Stenberg (4.2.2)	$k$ [ $k$ ]	$k$ [ $k-1$ ]	$k$ [ $k$ ]	$k \geq 2$	2/3
GG (4.2.2)	$k$ [ $k$ ]	$k$ [ $k-1$ ]	$k$ [ $k$ ]	$k \geq 2$	2/3
Barycentric (4.2.3)	$k$ [ $k$ ]	$k$ [ $k-1$ ]	$k$ [ $k$ ]	$k \geq 2$	2/3
Awanou low (4.2.4)	1 [2]	1 [1]	1 [1]	1	2/3
Awanou high (4.2.4)	$k$ [ $k+1$ ]	$k$ [ $k+1$ ]	$k$ [ $k$ ]	$k \geq 2$	2
rAAQ (4.2.4)	$k$ [ $k$ ]	$k$ [ $k$ ]	$k$ [ $k$ ]	$k \geq 2$	2
rGG (4.2.4)	$k$ [ $k$ ]	$k$ [ $k-1$ ]	$k$ [ $k$ ]	$k \geq 1$	2/3

THB = Taylor–Hood based elements, 2D Rect. = 2D rectangular element

AFW = Arnold–Falk–Winther, CGG = Cockburn–Gopalakrishnan–Guzmán

GG = Gopalakrishnan–Guzmán, Barycentric = the elements on barycentric meshes

rAAQ = Arnold–Awanou–Qiu family on rectangular meshes, rGG = rectangular GG

It is known that the triple  $(rBDM_1(\mathbb{R}^n), \mathcal{P}_0^d(\mathbb{R}^n), \mathcal{P}_0^d(\mathbb{R}_{\text{skw}}^{n \times n}))$  is elasticity stable [8]. However, these elements are not readily extended to higher orders as in the triangular AFW elements. The higher order elements in [8] are  $(rRTN_{k+2}, \mathcal{Q}_{k+1}^d, \mathcal{Q}_k^d)$  with  $k \geq 1$ , which are not rectangular analogues of the AFW elements. The Arnold–Awanou–Qiu family [2] on rectangular meshes is  $(rRTN_{k+1}, \mathcal{Q}_k^d, \mathcal{P}_k^d)$  with  $k \geq 1$ . As discussed in [2], this family can be analyzed with the framework using Stokes stable pairs  $(\mathcal{Q}_{k+1}^c, \mathcal{P}_k^d)$ ,  $k \geq 1$  and elasticity stable pairs  $(rRTN_{k+1}, \mathcal{Q}_k^d, 0)$ ,  $k \geq 1$ . Existence of similar higher order elements in three dimensions is not clear.

As the last example, we propose a family of new rectangular elements, say, rectangular GG elements. Following the construction of the GG elements we define  $\hat{B}^r(\eta)$  for  $\eta \in \hat{\mathcal{P}}_k^d(\mathbb{R}_{\text{skw}}^{n \times n})$  as in (4.2) but with a standard rectangular quartic bubble function  $b_T^r$  in two dimensions, and a rectangular matrix bubble function  $\mathbf{b}_T^r$  in three dimensions. We refer to [20] for precise definitions of  $b_T^r$ ,  $\mathbf{b}_T^r$ , and a proof of (4.1) with  $b_T^r$  and  $\mathbf{b}_T^r$ . The space  $\hat{B}_k^r$  is defined correspondingly. The rectangular GG elements are defined by

$$\Sigma_h = rBDM_k(\mathbb{R}^n) + \text{curl}(\hat{B}_k^r), \quad U_h = \mathcal{P}_{k-1}^d(\mathbb{R}^n), \quad \Gamma_h = \mathcal{P}_k^d(\mathbb{R}_{\text{skw}}^{n \times n}),$$

for  $k \geq 1$ . The stability of these elements can be proved with an analysis similar to the GG elements by taking  $\Gamma_h^0 = \mathcal{P}_0^d(\mathbb{R}_{\text{skw}}^{n \times n})$  and  $\Xi_h = \hat{B}_k^r$ . By counting degrees of polynomials one can see that  $\text{curl}(\hat{B}_{k-2}^r)$  is included in  $rBDM_k(\mathbb{R}^n)$  since  $rBDM_k$  contains all polynomials of degree  $k$  (cf. (3.29) and (3.30) in [11]), so the number of degrees of freedom of  $\Sigma_h$  can be reduced because only a part of  $\text{curl}(\hat{B}_k^r)$  is necessary for  $\Sigma_h$ .

## 5. CONCLUSION

We presented a framework for the analysis of mixed methods for elasticity with weakly symmetric stress. The framework enables us to analyze many existing weak symmetry elements in a unified way and it is also useful in finding new stable mixed finite elements. An extension of the framework to quadrilateral or hexahedral elements seems to be an interesting nontrivial question.

## REFERENCES

- [1] M. Amara and J. M. Thomas. Equilibrium finite elements for the linear elastic problem. *Numer. Math.*, 33(4):367–383, 1979.
- [2] Douglas N. Arnold, Gerard Awanou, and Weifeng Qiu. Mixed finite elements for elasticity on quadrilateral meshes. <http://arxiv.org/abs/1306.6821>, preprint.
- [3] Douglas N. Arnold, Franco Brezzi, and Jr. Jim Douglas. PEERS: a new mixed finite element for plane elasticity. *Japan J. Appl. Math.*, 1:347–367, 1984.
- [4] Douglas N. Arnold, Franco Brezzi, and M. Fortin. A stable finite element for the Stokes equations. *Calcolo*, 21(4):337–344 (1985), 1984.
- [5] Douglas N. Arnold, Jr. Jim Douglas, and Chaitan P. Gupta. A family of higher order mixed finite element methods for plane elasticity. *Numer. Math.*, 45(1):1–22, 1984.
- [6] Douglas N. Arnold, Richard S. Falk, and Ragnar Winther. Finite element exterior calculus, homological techniques, and applications. *Acta Numer.*, 15:1–155, 2006.
- [7] Douglas N. Arnold, Richard S. Falk, and Ragnar Winther. Mixed finite element methods for linear elasticity with weakly imposed symmetry. *Math. Comp.*, 76(260):1699–1723 (electronic), 2007.
- [8] Gerard Awanou. Rectangular mixed elements for elasticity with weakly imposed symmetry condition. *Adv. Comput. Math.*, 38(2):351–367, 2013.
- [9] Daniele Boffi, Franco Brezzi, and Michel Fortin. Reduced symmetry elements in linear elasticity. *Commun. Pure Appl. Anal.*, 8(1):95–121, 2009.
- [10] Susanne C. Brenner and L. Ridgway Scott. *The Mathematical Theory of Finite Element Methods*. Springer, Third edition, 2008.
- [11] F. Brezzi and M. Fortin. *Mixed and Hybrid Finite Element Methods*, volume 15 of *Springer Series in computational Mathematics*. Springer, 1992.
- [12] Bernardo Cockburn, Jayadeep Gopalakrishnan, and Johnny Guzmán. A new elasticity element made for enforcing weak stress symmetry. *Math. Comp.*, 79(271):1331–1349, 2010.
- [13] B. M. Fraeijs de Veubeke. Stress function approach. In *Proceedings of the World Congress on Finite Element Methods in Structural Mechanics*, volume 5, pages J.1 – J.51. 1975.
- [14] Lawrence C. Evans. *Partial differential equations*, volume 19 of *Graduate Studies in Mathematics*. American Mathematical Society, Providence, RI, 1998.
- [15] Richard S. Falk. Finite elements for linear elasticity. In D. Boffi and L. Gastaldi, editors, *Mixed Finite Elements: Compatibility Conditions*, volume 1939. Springer, 2008.
- [16] Mohamed Farhloul and Michel Fortin. Dual hybrid methods for the elasticity and the Stokes problems: a unified approach. *Numer. Math.*, 76(4):419–440, 1997.
- [17] J. Gopalakrishnan and J. Guzmán. A second elasticity element using the matrix bubble. *IMA J. Numer. Anal.*, 32(1):352–372, 2012.
- [18] J. Guzmán. A unified analysis of several mixed methods for elasticity with weak stress symmetry. *J. Sci. Comput.*, 44(2):156–169, 2010.
- [19] Mika Juntunen and Jeonghun Lee. A mesh dependent norm analysis of low order mixed finite element for elasticity with weakly symmetric stress. *Mathematical Models and Methods in Applied Sciences*, pages 1–15, 2013. To appear, DOI:10.1142/S0218202514500171.
- [20] Mika Juntunen and Jeonghun Lee. Optimal second order rectangular elasticity elements with weakly symmetric stress. *BIT Numerical Mathematics*, pages 1–21, 2013. To appear, DOI:10.1007/s10543-013-0460-2.
- [21] Mary E. Morley. A family of mixed finite elements for linear elasticity. *Numer. Math.*, 55(6):633–666, 1989.
- [22] Jinshui Qin. *On the convergence of some low order mixed finite elements for incompressible fluids*. ProQuest LLC, Ann Arbor, MI, 1994. Thesis (Ph.D.)—The Pennsylvania State University.

- [23] Rolf Stenberg. Analysis of mixed finite elements methods for the Stokes problem: a unified approach. *Math. Comp.*, 42(165):9–23, 1984.
- [24] Rolf Stenberg. On the construction of optimal mixed finite element methods for the linear elasticity problem. *Numer. Math.*, 48(4):447–462, 1986.
- [25] Rolf Stenberg. A family of mixed finite elements for the elasticity problem. *Numer. Math.*, 53(5):513–538, 1988.
- [26] Shangyou Zhang. A new family of stable mixed finite elements for the 3D Stokes equations. *Math. Comp.*, 74(250):543–554, 2005.

AALTO UNIVERSITY, DEPARTMENT OF MATHEMATICS AND SYSTEMS ANALYSIS, P.O. Box 11100,  
00076 AALTO, FINLAND, JEONGHUN.LEE@AALTO.FI