

# Matrix equilibrium states

Mark Piraino

University of Victoria

July 23, 2019

# Equilibrium/Gibbs states

Let  $\Sigma = \{0, \dots, M-1\}$  and  $\sigma : \Sigma^{\mathbb{Z}} \rightarrow \Sigma^{\mathbb{Z}}$  be the shift map.

- “Scalar” equilibrium/Gibbs states. Let  $\varphi : \Sigma^{\mathbb{Z}} \rightarrow \mathbb{R}$ .
  - Equilibrium: Maximize

$$h_{\mu}(\sigma) + \int \varphi d\mu$$

over  $\sigma$  invariant measures.

- Gibbs: There exist constants  $C > 0$  and  $P$  such that

$$C^{-1}e^{-nP+S_n\varphi(x)} \leq \mu_{\varphi}[x_0x_1\cdots x_{n-1}] \leq Ce^{-nP+S_n\varphi(x)}$$

for all  $x \in \Sigma^{\mathbb{Z}}$ .

- “Matrix” equilibrium/Gibbs states.
  - $\varphi : \Sigma^{\mathbb{Z}} \rightarrow \mathbb{R}$  is replaced by  $\mathcal{A} : \Sigma^{\mathbb{Z}} \rightarrow M_d(\mathbb{R})$ .
  - $S_n\varphi(x)$  is replaced by  $\log \|\mathcal{A}(\sigma^{n-1}x)\cdots\mathcal{A}(\sigma x)\mathcal{A}(x)\|$ .

# Equilibrium/Gibbs states

Let  $\mathcal{A} : \Sigma^{\mathbb{Z}} \rightarrow M_d(\mathbb{R})$  depend only on 1-coordinate,  
 $\mathcal{A} = (A_0, A_1, \dots, A_{M-1})$  and  $\beta \in \mathbb{R}$ .

- Equilibrium: Maximize

$$h_{\mu}(\sigma) + \underbrace{\beta \lim_{n \rightarrow \infty} \frac{1}{n} \int \log \|A_{x_{n-1}} \cdots A_{x_1} A_{x_0}\| d\mu}_{\Lambda(\mathcal{A}, \mu)}$$

over  $\sigma$  invariant measures.

- Gibbs: There exist constants  $C > 0$  and  $P$  such that

$$C^{-1} e^{-nP} \|A_{x_{n-1}} \cdots A_{x_0}\|^{\beta} \leq \mu_{\mathcal{A}, \beta}[x_0 \cdots x_{n-1}] \leq C e^{-nP} \|A_{x_{n-1}} \cdots A_{x_0}\|^{\beta}$$

for all  $x \in \Sigma^{\mathbb{Z}}$ .

**Comment:**  $P$  is the maximum of  $h_{\mu}(\sigma) + \beta \Lambda(\mathcal{A}, \mu)$ ,  $\beta \mapsto P(\beta)$  is called the (matrix) pressure function.

# Some reasons to care about MES

- They appear “in the wild”
  - Hidden Markov measures.
  - The Kusuoka measure.
- You care about their “cousin”: equilibrium states for the singular value potential.
  - Applications to self-affine fractals.
- You care about the function  $P(\beta)$ .
  -

$$\left. \frac{dP}{d\beta} \right|_{\beta=0} = \lambda$$

where  $\lambda$  is maximal Lyapunov exponent with respect to the uniform Bernoulli measure.

- Provided you take some assumptions on  $\mathcal{A}$  and  $\lambda = 0$

$$\frac{\log \|A_{x_{n-1}} \cdots A_{x_1} A_{x_0}\|}{\sqrt{n}} \xrightarrow{\text{dist}} \mathcal{N}(0, \sigma^2) \text{ and } \left. \frac{d^2 P}{d\beta^2} \right|_{\beta=0} = \sigma^2.$$

# Some questions about MES

- Existence and uniqueness:
  - Provided the collection  $\mathcal{A}$  does not preserve a proper and non-trivial subspace and  $\beta > 0$  matrix equilibrium states are unique and satisfy the Gibbs inequality. (Feng and Käenmäki '11)
- Ergodic and statistical properties:
  - Mixing? If yes can we find a mixing rate?
  - Isomorphic to a Bernoulli shift?
- Equilibrium states for (regular enough) scalar potentials are known to have very nice properties to what extent is the same true for matrix equilibrium states?

# Constructing matrix Gibbs states

- **Sketch:** Assume  $\mathcal{A}$  is irreducible and  $\beta > 0$ . Define

$$\nu_n[x_0 \cdots x_{n-1}] = \frac{\|A_{x_{n-1}} \cdots A_{x_1} A_{x_0}\|^\beta}{\sum_{y_0 \cdots y_{n-1}} \|A_{y_{n-1}} \cdots A_{y_1} A_{y_0}\|^\beta}$$

take a weak\* limit point  $\nu$ . Satisfies Gibbs inequality but might not be shift invariant so average it

$$\frac{1}{N} \sum_{k=0}^{N-1} \sigma_*^k \nu$$

Take a weak\* limit point  $\mu$  which is the matrix Gibbs state for  $(\mathcal{A}, \beta)$ .

# Constructing matrix Gibbs states

- **Sketch:** Assume  $\mathcal{A}$  is irreducible and  $\beta > 0$ . Define

$$\nu_n[x_0 \cdots x_{n-1}] = \frac{\|A_{x_{n-1}} \cdots A_{x_1} A_{x_0}\|^\beta}{\sum_{y_0 \cdots y_{n-1}} \|A_{y_{n-1}} \cdots A_{y_1} A_{y_0}\|^\beta}$$

take a weak\* limit point  $\nu$ . Satisfies Gibbs inequality but might not be shift invariant so average it

$$\frac{1}{N} \sum_{k=0}^{N-1} \sigma_*^k \nu$$

Take a weak\* limit point  $\mu$  which is the matrix Gibbs state for  $(\mathcal{A}, \beta)$ .

- If you use this construction you have to rely exclusively on the Gibbs inequality to prove ergodic properties.

# Constructing matrix Gibbs states

- **Sketch:** Assume  $\mathcal{A}$  is irreducible and  $\beta > 0$ . Define

$$\nu_n[x_0 \cdots x_{n-1}] = \frac{\|A_{x_{n-1}} \cdots A_{x_1} A_{x_0}\|^\beta}{\sum_{y_0 \cdots y_{n-1}} \|A_{y_{n-1}} \cdots A_{y_1} A_{y_0}\|^\beta}$$

take a weak\* limit point  $\nu$ . Satisfies Gibbs inequality but might not be shift invariant so average it

$$\frac{1}{N} \sum_{k=0}^{N-1} \sigma_*^k \nu$$

Take a weak\* limit point  $\mu$  which is the matrix Gibbs state for  $(\mathcal{A}, \beta)$ .

- If you use this construction you have to rely exclusively on the Gibbs inequality to prove ergodic properties.
- We need a new construction.



# 1-Gibbs states for non-negative matrices

Set  $\mathcal{A} = (A_0, A_1)$  where  $A_0 = \begin{bmatrix} 1 & 1 \\ 0 & 1 \end{bmatrix}$ ,  $A_1 = \begin{bmatrix} 1 & 0 \\ 1 & 1 \end{bmatrix}$ .

# 1-Gibbs states for non-negative matrices

Set  $\mathcal{A} = (A_0, A_1)$  where  $A_0 = \begin{bmatrix} 1 & 1 \\ 0 & 1 \end{bmatrix}, A_1 = \begin{bmatrix} 1 & 0 \\ 1 & 1 \end{bmatrix}$ .

$$A = \begin{bmatrix} 1 & 1 \\ 0 & 1 \end{bmatrix} + \begin{bmatrix} 1 & 0 \\ 1 & 1 \end{bmatrix} = \begin{bmatrix} 2 & 1 \\ 1 & 2 \end{bmatrix} \text{ notice } A \begin{bmatrix} 1 \\ 1 \end{bmatrix} = 3 \begin{bmatrix} 1 \\ 1 \end{bmatrix}, A^T \begin{bmatrix} 1 \\ 1 \end{bmatrix} = 3 \begin{bmatrix} 1 \\ 1 \end{bmatrix}.$$

# 1-Gibbs states for non-negative matrices

Set  $\mathcal{A} = (A_0, A_1)$  where  $A_0 = \begin{bmatrix} 1 & 1 \\ 0 & 1 \end{bmatrix}, A_1 = \begin{bmatrix} 1 & 0 \\ 1 & 1 \end{bmatrix}$ .

$$A = \begin{bmatrix} 1 & 1 \\ 0 & 1 \end{bmatrix} + \begin{bmatrix} 1 & 0 \\ 1 & 1 \end{bmatrix} = \begin{bmatrix} 2 & 1 \\ 1 & 2 \end{bmatrix} \text{ notice } A \begin{bmatrix} 1 \\ 1 \end{bmatrix} = 3 \begin{bmatrix} 1 \\ 1 \end{bmatrix}, A^T \begin{bmatrix} 1 \\ 1 \end{bmatrix} = 3 \begin{bmatrix} 1 \\ 1 \end{bmatrix}.$$

The formula

$$\mu_{\mathcal{A},1}[x_0 x_1 \cdots x_{n-1}] = 3^{-n} \begin{bmatrix} 1 & 1 \end{bmatrix} A_{x_{n-1}} \cdots A_{x_1} A_{x_0} \begin{bmatrix} 1/2 \\ 1/2 \end{bmatrix}$$

extends to a measure on  $\{0,1\}^{\mathbb{Z}}$  by Kolmogorov consistency.

# 1-Gibbs states for non-negative matrices

Set  $\mathcal{A} = (A_0, A_1)$  where  $A_0 = \begin{bmatrix} 1 & 1 \\ 0 & 1 \end{bmatrix}, A_1 = \begin{bmatrix} 1 & 0 \\ 1 & 1 \end{bmatrix}$ .

$$A = \begin{bmatrix} 1 & 1 \\ 0 & 1 \end{bmatrix} + \begin{bmatrix} 1 & 0 \\ 1 & 1 \end{bmatrix} = \begin{bmatrix} 2 & 1 \\ 1 & 2 \end{bmatrix} \text{ notice } A \begin{bmatrix} 1 \\ 1 \end{bmatrix} = 3 \begin{bmatrix} 1 \\ 1 \end{bmatrix}, A^T \begin{bmatrix} 1 \\ 1 \end{bmatrix} = 3 \begin{bmatrix} 1 \\ 1 \end{bmatrix}.$$

The formula

$$\mu_{\mathcal{A},1}[x_0 x_1 \cdots x_{n-1}] = 3^{-n} \begin{bmatrix} 1 & 1 \end{bmatrix} A_{x_{n-1}} \cdots A_{x_1} A_{x_0} \begin{bmatrix} 1/2 \\ 1/2 \end{bmatrix}$$

extends to a measure on  $\{0,1\}^{\mathbb{Z}}$  by Kolmogorov consistency. This is the 1-Gibbs state for  $\mathcal{A}$  as

$$\begin{bmatrix} 1 & 1 \end{bmatrix} A_{x_{n-1}} \cdots A_{x_1} A_{x_0} \begin{bmatrix} 1/2 \\ 1/2 \end{bmatrix} \approx \|A_{x_{n-1}} \cdots A_{x_1} A_{x_0}\|_1$$

moreover we can see that  $P = \log 3$ .

# 1-Gibbs states for non-negative matrices

We can deduce mixing properties from spectral properties of  $A = A_0 + A_1$ .

$$\begin{aligned}
 & \mu_{\mathcal{A},1}([i_0 \cdots i_{m-1}] \cap \sigma^{-n}[j_0 \cdots j_{m'-1}]) \\
 &= 3^{-(m+m'+n-m)} \sum_{k_0 \cdots k_{n-m-1}} \left\langle \underbrace{A_{j_{m'}} \cdots A_{j_0}}_{A_J} \underbrace{A_{k_{n-m-1}} \cdots A_{k_0}}_{A_K} \underbrace{A_{i_m} \cdots A_{i_0}}_{A_I} [1/2], [1] \right\rangle \\
 &= 3^{-(m+m')} \left\langle A_J 3^{-(n-m)} A^{n-m} A_I [1/2], [1] \right\rangle \text{ because } \sum_{|K|=n-m} A_K = A^{n-m} \\
 &\xrightarrow{n \rightarrow \infty} 3^{-m} \langle A_I [1/2], [1] \rangle 3^{-m'} \langle A_J [1/2], [1] \rangle \text{ because } A \text{ is positive} \\
 &= \mu_{\mathcal{A},1}([i_0 \cdots i_{m-1}]) \mu_{\mathcal{A},1}([j_0 \cdots j_{m'-1}])
 \end{aligned}$$

# 1-Gibbs states for non-negative matrices

We can deduce mixing properties from spectral properties of  $A = A_0 + A_1$ .

$$\begin{aligned}
 & \mu_{\mathcal{A},1}([i_0 \cdots i_{m-1}] \cap \sigma^{-n}[j_0 \cdots j_{m'-1}]) \\
 &= 3^{-(m+m'+n-m)} \sum_{k_0 \cdots k_{n-m-1}} \left\langle \underbrace{A_{j_{m'}} \cdots A_{j_0}}_{A_J} \underbrace{A_{k_{n-m-1}} \cdots A_{k_0}}_{A_K} \underbrace{A_{i_m} \cdots A_{i_0}}_{A_I} [1/2], [1] \right\rangle \\
 &= 3^{-(m+m')} \left\langle A_J 3^{-(n-m)} A^{n-m} A_I [1/2], [1] \right\rangle \text{ because } \sum_{|K|=n-m} A_K = A^{n-m} \\
 &\xrightarrow{n \rightarrow \infty} 3^{-m} \langle A_I [1/2], [1] \rangle 3^{-m'} \langle A_J [1/2], [1] \rangle \text{ because } A \text{ is positive} \\
 &= \mu_{\mathcal{A},1}([i_0 \cdots i_{m-1}]) \mu_{\mathcal{A},1}([j_0 \cdots j_{m'-1}])
 \end{aligned}$$

**Question:** How do we deal with  $\beta \neq 1$ ? Matrices which are not non-negative?

# Matrix Gibbs states: Defining the operators

Start with  $\mathcal{A} = (A_0, A_1, \dots, A_{M-1})$  and  $\beta$ .

# Matrix Gibbs states: Defining the operators

Start with  $\mathcal{A} = (A_0, A_1, \dots, A_{M-1})$  and  $\beta$ .

**Trick:** Consider the action of the matrices on norm one vectors weighted by  $\|A_i u\|^\beta$ .



# Matrix Gibbs states: Defining the operators

Start with  $\mathcal{A} = (A_0, A_1, \dots, A_{M-1})$  and  $\beta$ .

**Trick:** Consider the action of the matrices on norm one vectors weighted by  $\|A_i u\|^\beta$ . Define

$$L_i : C(\mathbb{RP}^{d-1}) \rightarrow C(\mathbb{RP}^{d-1}) \text{ by } L_i f(\bar{u}) = \left\| A_i \frac{u}{\|u\|} \right\|^\beta f(\overline{A_i u}).$$

# Matrix Gibbs states: Defining the operators

Start with  $\mathcal{A} = (A_0, A_1, \dots, A_{M-1})$  and  $\beta$ .

**Trick:** Consider the action of the matrices on norm one vectors weighted by  $\|A_i u\|^\beta$ . Define

$$L_i : C(\mathbb{RP}^{d-1}) \rightarrow C(\mathbb{RP}^{d-1}) \text{ by } L_i f(\bar{u}) = \left\| A_i \frac{u}{\|u\|} \right\|^\beta f(\overline{A_i u}).$$

Transfer operator:

$$L_{\mathcal{A},\beta} = \sum_i L_i \text{ that is } L_{\mathcal{A},\beta} f(\bar{u}) = \sum_i \left\| A_i \frac{u}{\|u\|} \right\|^\beta f(\overline{A_i u}).$$

# Matrix Gibbs states: Defining the operators

Start with  $\mathcal{A} = (A_0, A_1, \dots, A_{M-1})$  and  $\beta$ .

**Trick:** Consider the action of the matrices on norm one vectors weighted by  $\|A_i u\|^\beta$ . Define

$$L_i : C(\mathbb{RP}^{d-1}) \rightarrow C(\mathbb{RP}^{d-1}) \text{ by } L_i f(\bar{u}) = \left\| A_i \frac{u}{\|u\|} \right\|^\beta f(\overline{A_i u}).$$

Transfer operator:

$$L_{\mathcal{A},\beta} = \sum_i L_i \text{ that is } L_{\mathcal{A},\beta} f(\bar{u}) = \sum_i \left\| A_i \frac{u}{\|u\|} \right\|^\beta f(\overline{A_i u}).$$

Take (if they exist)  $h > 0$   $\nu$  with  $L_{\mathcal{A},\beta} h = \rho(L_{\mathcal{A},\beta}) h$  and  $L_{\mathcal{A},\beta}^* \nu = \rho(L_{\mathcal{A},\beta}) \nu$ .

$$\mu_{\mathcal{A},\beta}[x_0 x_1 \dots x_{n-1}] = \rho(L_{\mathcal{A},\beta})^{-n} \int_{\mathbb{RP}^{d-1}} L_{x_0} L_{x_1} \dots L_{x_{n-1}} h d\nu.$$

$$\mu_{\mathcal{A},\beta}[x_0x_1\cdots x_{n-1}] = \rho(L_{\mathcal{A},\beta})^{-n} \int_{\mathbb{RP}^{d-1}} L_{x_0}L_{x_1}\cdots L_{x_{n-1}} h d\nu.$$

Provided the collection  $\mathcal{A}$  is proximal (read contracting) and strongly irreducible (read topologically mixing) then  $L_{\mathcal{A},\beta}$  has a spectral gap (Guivarc'h and Le Page '04). This implies  $h$  and  $\nu$  exist and

$$\rho(L_{\mathcal{A},\beta})^{-n} L_{\mathcal{A},\beta}^n f \xrightarrow{n \rightarrow \infty} \langle f, \nu \rangle h.$$

$$\mu_{\mathcal{A},\beta}[x_0x_1\cdots x_{n-1}] = \rho(L_{\mathcal{A},\beta})^{-n} \int_{\mathbb{RP}^{d-1}} L_{x_0}L_{x_1}\cdots L_{x_{n-1}} h d\nu.$$

Provided the collection  $\mathcal{A}$  is proximal (read contracting) and strongly irreducible (read topologically mixing) then  $L_{\mathcal{A},\beta}$  has a spectral gap (Guivarc'h and Le Page '04). This implies  $h$  and  $\nu$  exist and

$$\rho(L_{\mathcal{A},\beta})^{-n} L_{\mathcal{A},\beta}^n f \xrightarrow{n \rightarrow \infty} \langle f, \nu \rangle h.$$

This makes the analog of the computation

$$\begin{aligned} \mu_{\mathcal{A},1}([I] \cap \sigma^{-n}[J]) &= 3^{-|I|-|J|} \left\langle A_J 3^{-n+|J|} A^{n-|J|} A_I [1/2], [1] \right\rangle \\ &\xrightarrow{n \rightarrow \infty} 3^{-|I|} \langle A_I [1/2], [1] \rangle 3^{-|J|} \langle A_J [1/2], [1] \rangle \\ &= \mu_{\mathcal{A},1}([I]) \mu_{\mathcal{A},1}([J]). \end{aligned}$$

work which yields the following result...

# Result

## Theorem (P.)

*Suppose that  $\mathcal{A} = (A_0, \dots, A_{M-1})$  is a collection of real invertible  $d \times d$  matrices which is proximal and strongly irreducible. Then for any  $\beta \geq 0$  there exists a unique Gibbs state for  $(\mathcal{A}, \beta)$ ,  $\mu_{\mathcal{A}, \beta}$ . Moreover*

- ①  $\mu_{\mathcal{A}, \beta}$  is weak Bernoulli.
- ②  $\mu_{\mathcal{A}, \beta}$  has exponential decay of correlations for Hölder continuous functions. That is, for a fixed  $\theta \in (0, 1)$  there are constants  $D$  and  $\gamma \in (0, 1)$  such that

$$\left| \int f \cdot g \circ \sigma^n d\mu_{\mathcal{A}, \beta} - \int f d\mu_{\mathcal{A}, \beta} \int g d\mu_{\mathcal{A}, \beta} \right| \leq D \|f\|_{\theta} \|g\|_{\theta} \gamma^n$$

for all  $f, g \in \mathcal{H}_{\theta}$ ,  $n \geq 0$ .

# Where do we go from here?

- Irreducible collections:
  - Can fail to be mixing.
  - Isomorphic to a Bernoulli shift cross a finite rotation?

# Where do we go from here?

- Irreducible collections:
  - Can fail to be mixing.
  - Isomorphic to a Bernoulli shift cross a finite rotation?
- What is the connection between  $L_{\mathcal{A},\beta} : C(\mathbb{RP}^{d-1}) \rightarrow C(\mathbb{RP}^{d-1})$  and the transfer operator  $\mathcal{L} : L^1(\mu_{\mathcal{A},\beta}) \rightarrow L^1(\mu_{\mathcal{A},\beta})$ ?



# Where do we go from here?

- Irreducible collections:
  - Can fail to be mixing.
  - Isomorphic to a Bernoulli shift cross a finite rotation?
- What is the connection between  $L_{\mathcal{A},\beta} : C(\mathbb{RP}^{d-1}) \rightarrow C(\mathbb{RP}^{d-1})$  and the transfer operator  $\mathcal{L} : L^1(\mu_{\mathcal{A},\beta}) \rightarrow L^1(\mu_{\mathcal{A},\beta})$ ?
- Lots other questions:
  - $\beta < 0$ ?
  - Singular value potential?
  - $\mathcal{A} : \Sigma^{\mathbb{Z}} \rightarrow M_d(\mathbb{R})$  Hölder?
  - etc...