

Q heat, T temperature

[heat is kinetic energy transferred btw systems]
[temperature is avg kinetic energy of a system]

$$(1) \quad \Delta T \propto Q/\Delta m$$

thin cylinder with length coord z

[draw]

$$(2) \quad \Delta m \propto \Delta z$$

temperature [is a] function $T(z, t)$

[look at cross-section from z to z + Δz :]

[heat the cylinder and study T:]

[heat to cross-section prop to change in $\partial T/\partial z$ from z to z + Δz , not change in T itself]

$$(3) \quad Q \propto [(\partial T/\partial z)(z + \Delta z, t) - (\partial T/\partial z)(z, t)] \Delta t$$

altogether

$$\Delta T/\Delta t = c (1/\Delta z) [(\partial T/\partial z)(z + \Delta z, t) - (\partial T/\partial z)(z, t)]$$

taking limits,

$$\partial T/\partial t = c \partial^2 T/\partial z^2 \quad [\text{"heat eq for cylinder"}]$$

now suppose cylinder is from $z = 0$ to $z = 1$
and $T(0, 0) = T(1, 0) = 0$ [abs zero]

to solve heat eq: start with separable case

$$T(z, t) = A(z)B(t)$$

$$\text{then } B'(t)/B(t) = c A''(z)/A(z)$$

$$\begin{aligned} \text{so there is } \lambda \text{ s.t. } \quad A'' &= \lambda A \\ B' &= c\lambda B \end{aligned}$$

$$\begin{aligned} \lambda > 0: \quad A &= c_1 e^{\sqrt{\lambda}z} + c_2 e^{-\sqrt{\lambda}z} \\ \text{boundary cond's force } c_1 &= c_2 = 0 \end{aligned}$$

$$\lambda = 0: \quad [\text{what happens?}] \text{ again } A = 0$$

only $\lambda < 0$ interesting:

$$\begin{aligned} A &= c_1 \cos(\sqrt{|\lambda|}z) + c_2 \sin(\sqrt{|\lambda|}z) \\ B &= e^{c\lambda t} \end{aligned}$$

boundary cond's force $c_1 = 0$ and $\sqrt{|\lambda|} = \pi n$

Thm (Fourier) separable solns look like

$$T(z, t) = \sin(\pi n z) e^{-c(\pi n)^2 t}$$

gen'l solution is a superposition:

$$T(z, t) = \sum_{n \geq 1} a_n \sin(\pi n z) e^{-c(\pi n)^2 t}$$

[does anyone know the name for these coeffs?]

Thm (Fourier) inversion formula
for “nice” T :

$$a_n = 2 \int_0^1 T(z, 0) \sin(\pi n z) dz$$

now suppose at $t = 0$: Dirac delta at $z = 1/2$
abs zero elsewhere

$$\int_0^1 T(z, 0) \varphi(z) dz = \varphi(1/2)$$

then

$$a_n = 2 \sin(\pi n/2) = \begin{array}{ll} 0 & n \equiv 0 \pmod{4} \\ 2 & n \equiv 1 \pmod{4} \\ 0 & n \equiv 2 \pmod{4} \\ -2 & n \equiv 3 \pmod{4} \end{array}$$

$$\begin{aligned} T(z, t) &= 2 \left[\sin(\pi z) e^{-c\pi^2 t} \right. \\ &\quad - \sin(3\pi z) e^{-c(3\pi)^2 t} \\ &\quad + \sin(5\pi z) e^{-c(5\pi)^2 t} \\ &\quad \left. - \dots \right] \\ &= \sum_{n \in \mathbb{Z}} e^{2\pi i(n + 1/2)(z - 1/2)} \\ &\quad e^{-c(4\pi^2)(n + 1/2)^2 t} \end{aligned}$$

by de Moivre

[renormalize to:]

Df the Jacobi theta function

$$\Theta(z, \tau) = \sum_{n \in \mathbb{Z}} e^{2\pi i n z} e^{\pi i n^2 \tau}$$

for z, τ s.t. $\text{Im}(\tau) > 0$ [“upper half-plane”]

[$c = 1/(4\pi)$]:

$$T = e^{\pi i(z - 1/2) - t/4} \Theta(z + it/2 - 1/2, it)$$

$\Theta(z, \tau)$ highly symmetric: focus on $\theta(\tau) = \Theta(0, \tau)$

$$\theta(\tau + 1) = \theta(\tau)$$

$$\theta(-1/\tau) = \sqrt{-i\tau} \theta(\tau), \text{ where } \text{Im } \sqrt{-i\tau} > 0$$

#2 deeper: uses “Poisson summation”

also related to number theory via $q = e^{\pi i \tau}$:

$$\theta = \sum_n q^{n^2},$$

$$\theta^2 = \sum_m r_2(m) q^m$$

where $r_2(m) = |\{(a, b) \in \mathbb{Z}^2 \mid m = a^2 + b^2\}|$

Brahmagupta:

$$r_2(m), r_2(n) > 0 \text{ implies } r_2(mn) > 0$$

[why?]

$$(a^2 + b^2)(c^2 + d^2)$$

$$= (ac + bd)^2 + (ad - bc)^2$$

Fermat: studied $r_2(m)$ for prime m

$$2 = 1^2 + 1^2$$

$$3$$

$$5 = 1^2 + 2^2$$

$$7$$

$$11$$

$$13 = 2^2 + 3^2 \quad [\text{pattern?}]$$

let $d_1(m) = |\{\text{divisors} \equiv 1 \pmod{4} \text{ of } m\}|$
 $d_3(m) = |\{\text{divisors} \equiv 3 \pmod{4} \text{ of } m\}|$

$$\xi(\tau) = 1 + 4 \sum_{m \geq 1} (d_1(m) - d_3(m)) q^m$$

Thm (Jacobi) $\theta^2 = \xi$
 $r_2(m) = 4(d_1(m) - d_3(m))$

Sketch $\xi(\tau) = 1 + 4 \sum_m q^m / (1 + q^{2m})$
 $= \sum_{n \in \mathbb{Z}} 2 / (q^n + q^{-n})$

can then show that θ^2, ξ are functions f s.t.

- $f(\tau + 2) = f(\tau)$ and $f(-1/\tau) = -i\tau f(\tau)$
- f analytic and nonvanishing on $\{\text{Im}(\tau) > 0\}$
- as $\text{Im}(\tau) \rightarrow \infty$,
 $f(\tau) \rightarrow 1$ and $f(1 - 1/\tau) \sim -4i\tau e^{\pi i \tau / 2}$

the ratio $F = \xi/\theta^2$ satisfies

- $F(\tau + 2) = F(\tau)$ and $F(-1/\tau) = F(\tau)$
- F is analytic and bounded on $\{\text{Im}(\tau) > 0\}$

[F is a “holomorphic modular form of weight 0 for the theta group”]

analogue of Liouville’s thm: all such F constant

[F descends to a bdd analytic fn on $\{\text{Im}(\tau) > 0\}/\sim$ and this surface is compact]

but constant terms of θ^2 and ξ are both 1