notebook

October 13, 2022

```
[948]: import numpy as np
import matplotlib.pyplot as plt
import cycler
```

Define initial condition $\psi(x,0)$.

```
[949]: def f(x): return np.sin(2*np.pi*x)**6
```

Define $\partial \psi / \partial x$.

```
[950]: def f_x(c, x):
    return -12*c*np.pi*((np.sin(2*np.pi*x))**5)*np.cos(2*np.pi*x)
```

0.0.1 Problem 1

Compute solutions to the wind speed advection equation using Adams-Bashforth (3rd-order) in time and 4th-order central differencing in space.

Function F for AB3.

```
[951]: def f_x_(c, dx, dt, f_):
    return c*cdf(dt, dx, f_)
```

Define discretization schemes: - Starter: Runge-Kutta, 3rd order - Temporal: Adams-Bashforth, 3rd-order - Spatial: 4th-order compact difference

Runge-Kutta, 3rd order

```
[952]: def rk3(h, c, dt, dx, x, psi):
    # q_1 = h*f_x(c, x)
    q_1 = h*f_x(c, dx, dt, psi)
    psi_1 = psi + (1/3)*q_1
    # q_2 = h*f_x(c, psi_1) - (5/9)*q_1
    q_2 = h*f_x(c, dx, dt, psi) - (5/9)*q_1
    psi_2 = psi_1 + (15/16)*q_2
    # q_3 = h*f_x(c, psi_2) - (153/128)*q_2
    q_3 = h*f_x(c, dx, dt, psi) - (153/128)*q_2
    psi_n_1 = psi_2 + (8/15)*q_3
    return psi_n_1
```

4th-order centered difference

Adams-Bashforth, 3rd order

```
[954]: def ab3(dx, dt, c, x, psis):
    # Get previous psi values [psi(n), psi(n-1), psi(n-2)]
    psi_n, psi_n_1, psi_n_2 = psis[0], psis[1], psis[2]
    # Generate new psi value psi(n+1)
    psi_np1 = psi_n + (dt/12)*(23*f_x_(c, dx, dt, psi_n) - 16*f_x_(c, dx, dt, psi_n_1) + 5*f_x_(c, dx, dt, psi_n_2))
    return psi_np1
```

Define function to carry out the discretization and step through time.

```
[955]: def advection p1(dx, c, cfl, t max, plot=False, plot step=40, printout=False):
           # Compute timestep to meet Courant number
           dt = dx*cfl/c
           # Spatial domain
           x = np.arange(0, 1, dx)
           # Temporal domain
           t = np.arange(0, t_max, dt)
           # Initialize array for values and apply initial condition
           y = np.full((len(t)+1, len(x)), np.nan)
           y[0, :] = f(x)
           # Get exact solution
           exact = np.sin(2*np.pi*(x-c*t max))**6
           # Iterate through timesteps
           for i, t_ in enumerate(t):
               if printout and i % 10 == 0:
                   print('Step: {0} | Timestep: {1:.2f}'.format(i, t_))
               # Starter function
               if i < 2:
```

```
print('Using starter scheme...') if printout else None
          y[i+1, :] = rk3(dt, c, dt, dx, x, y[i, :])
       else:
           print('Using full scheme...') if printout else None
           # Get function values from previous timesteps
           psis = [y[i, :], y[i-1, :], y[i-2, :]]
          y[i+1, :] = ab3(dx, dt, c, x, psis)
   ''' Plotting. '''
  if plot:
      plt.rcParams["axes.prop_cycle"] = plt.cycler('color', plt.cm.viridis(np.
→linspace(0, 1, len(plot_step))))
      fig, ax = plt.subplots(figsize=(4, 3))
      for step in plot_step:
           im = ax.plot(x, y[step], marker='o', markersize=4, label='Step {0}'.

¬format(step))
      ax.set_title('Time = {0:.2f} | dx = {1:.2f} | dt = {2:.2f}'.

→format(t[step], dx, dt), fontsize=10)
      ax.set_xlim([min(x), max(x)])
      ax.set_ylim([0, 1])
      fig.tight_layout()
      fig.legend(loc='upper right', bbox_to_anchor=(1.25, 0.925),_
⇒frameon=False, fontsize=8)
   # Return the values at the maximum time
  return x, y[-1, :], exact
```

Execute and plot the results

```
[]:  # Wave speed (c)
     c = 0.1
     # Spatial increments
     dxs = [1/20, 1/40, 1/80, 1/160]
     # Strings to represent spatial increments (for plotting purposes only)
     dxs_str = ['1/20', '1/40', '1/80', '1/160']
     # Courant number and Courant number list
     cfls = [0.1, 0.2, 0.4, 0.8]
     # Maximum time
     t max = 50
     # Boolean to control prints to console
     printout = False
     # Computation mode: variable dx or CFL values
     dx_calc = True
     # Initialize array to hold each run's output array
     values = {}
```

```
# Iterate over dx values
if dx_calc:
    cfl = 0.1
    for i, dx in enumerate(dxs):
        # End step
        end_step = int(t_max/(cfl*dx/c))
        # Perform calculation, plot individual results
        arr = advection_p1(dx, c, cfl, t_max, plot=False, plot_step=[end_step],_
 →printout=printout)
        # Store array
        values['dx = \{0\}'.format(dxs_str[i])] = {'x': arr[0], 'y': arr[1], ___
 ⇔'exact': arr[2]}
# Iterate over Courant number values
else:
    dx_index = 2
    dx = dxs[dx_index]
    for i, cfl in enumerate(cfls):
        # End step
        end_step = int(t_max/(cfl*dx/c))
        # Perform calculation, plot individual results
        arr = advection_p1(dx, c, cfl, t_max, plot=False, plot_step=[0,__
 -end_step // 16, end_step // 2, 3*end_step // 4], printout=printout)
        # Store array
        values['CFL = \{0\}'.format(cfl)] = \{'x': arr[0], 'y': arr[1], 'exact': \[ \] \}
 →arr[2]}
''' Plotting.'''
# Initialize figure
fig, ax = plt.subplots(figsize=(4, 3))
# Define formatting index, color list, marker list
f_index, markers = 0, ['o', '^', '+', '2']
# Set color cycling
plt.rcParams["axes.prop_cycle"] = plt.cycler('color', plt.cm.Blues(np.
 \rightarrowlinspace(0.375, 1, len(dxs))))
# Plot the actual values
for key, value in values.items():
    im = ax.plot(value['x'], value['y'], marker=markers[f_index], markersize=5,_
 ⇒ls='--', label=key)
    if key == 'dx = 1/160' or key == 'CFL = 0.1':
        ax.plot(value['x'], value['exact'], c='r', label='exact', zorder=0)
    f index += 1
# Metadata
if dx calc:
    ax.set_title('Courant number: {0:.1f}'.format(cfl), fontsize=10)
```

```
else:
    ax.set_title('dx = {0}'.format(dxs_str[dx_index]), fontsize=10)
ax.set_xlim([min(value['x']), max(value['x'])])
ax.set_ylim([0, 1])
fig.tight_layout()
fig.legend(loc='upper right', bbox_to_anchor=(1.25, 0.925), frameon=False,_u
 ⇔fontsize=8)
# plt.savefig('figs/p1c1.png', dpi=300, bbox_inches='tight')
# Initialize figure
fig, ax = plt.subplots(figsize=(4, 3))
# Collect RMS values
rmse = [rms(value['y'], value['exact']) for _, value in values.items()]
print(rmse)
# Plot the actual values
im = ax.loglog(dxs, rmse, marker='o', color='k')
ax.set_xticks([0.005, 0.01, 0.05])
ax.set_ylim([1e-3, 1e0])
# Metadata
ax.set_ylabel('RMSE')
ax.set xlabel('dx') if dx calc else ax.set xlabel('CFL')
ax.set_title('Root mean squared error', fontsize=10)
plt.gca().invert_xaxis()
fig.tight_layout()
# plt.savefig('figs/p1c2.png', dpi=300, bbox_inches='tight')
```

0.0.2 Problem 2

Compute solutions to the wind speed advection equation using leapfrog with an Asselin filter in time and Lele's 4th-order compact differencing in space.

Runge-Kutta, 3rd order

```
[956]: def rk3_2(h, c, dt, dx, x, psi):
    # q_1 = h*f_x(c, x)
    q_1 = h*f_x_2(c, dx, dt, psi)
    psi_1 = psi + (1/3)*q_1
    # q_2 = h*f_x(c, psi_1) - (5/9)*q_1
    q_2 = h*f_x_2(c, dx, dt, psi) - (5/9)*q_1
    psi_2 = psi_1 + (15/16)*q_2
    # q_3 = h*f_x(c, psi_2) - (153/128)*q_2
    q_3 = h*f_x_2(c, dx, dt, psi) - (153/128)*q_2
    psi_n_1 = psi_2 + (8/15)*q_3
    return psi_n_1
```

4th-order Lele compact difference

```
[957]: def cdf_compact(dt, dx, f_):
           \# f_{-} is an array of function values at a given time
           # dt is the timestep
           # Create upper, central, and bottom diagonals
           a = np.full(len(f_), 5/24)
           b = np.full(len(f_), 14/24)
           c = np.full(len(f_), 5/24)
           # Create RHS
           rhs = np.zeros((len(f), ))
           for i in range(0, len(f_)):
               # Using modulo for indices to handle boundary conditions, since BCs are
        \rightarrowperiodic
               rhs[i] = (11*(f_[(i+1) \% len(f_))] - f_[(i-1) \% len(f_)])/2 + (f_[(i+2)_{\bot})
        4\% len(f_)] - f_[(i-2) \% len(f_)])/4)/(12*dx)
           f_prime = cyc_tridiag(len(f_), a, b, c, rhs)
           return f_prime
```

Leapfrog with Asselin filter

```
[958]: def leapfrog(dx, dt, c, x, gamma, psis):
    # Get previous psi values [psi(n), psi(n-1)]
    psi_n, psi_n_1, psi_n_2 = psis
    # Asselin filter - previous step
    psi_n_1_ass = psi_n_1 + gamma*(psi_n - 2*psi_n_1 + psi_n_2)
    # Compute values at next time step
    psi_np1 = psi_n_1_ass + 2*dt*f_x_2(c, dx, dt, psi_n)
    # Compute filtered values at current time step
    psi_n_ass = psi_n + gamma*(psi_n_1_ass - 2*psi_n + psi_np1)
    return psi_n_ass, psi_np1
```

Define $\partial \psi / \partial x$.

```
[959]: def f_x_2(c, dx, dt, f_):
    return c*cdf_compact(dt, dx, f_)
```

Define function to carry out the discretization and step through time

```
[960]: def advection_p2(dx, c, cfl, t_max, plot=False, plot_step=40, printout=False):
    # Compute timestep to meet Courant number
    dt = dx*cfl/c
    # Spatial domain
    x = np.arange(0, 1, dx)
    # Temporal domain
```

```
t = np.arange(0, t_max, dt)
  # Asselin filter value
  gamma = 0.1
  # Initialize array for values and apply initial condition
  y = np.full((len(t)+1, len(x)), np.nan)
  y[0, :] = f(x)
  # Get exact solution
  exact = np.sin(2*np.pi*(x-c*t_max))**6
  # Iterate through timesteps
  for i, t_ in enumerate(t):
      if printout and i % 10 == 0:
          print('Step: {0} | Timestep: {1:.2f}'.format(i, t_))
      # Starter function
      if i < 3:
          print('Using starter scheme...') if printout else None
          y[i+1, :] = rk3_2(dt, c, dt, dx, x, y[i, :])
      else:
          print('Using full scheme...') if printout else None
           # Get function values from previous timesteps
          psis = [y[i, :], y[i-1, :], y[i-2, :]]
          y[i, :], y[i+1, :] = leapfrog(dx, dt, c, x, gamma, psis)
   ''' Plotting. '''
  if plot:
      plt.rcParams["axes.prop_cycle"] = plt.cycler('color', plt.cm.viridis(np.
→linspace(0, 1, len(plot_step))))
      fig, ax = plt.subplots(figsize=(4, 3))
      for step in plot step:
          im = ax.plot(x, y[step], marker='o', markersize=4, label='Step {0}'.

¬format(step))
      ax.set_title('Time = {0:.2f} | dx = {1:.2f} | dt = {2:.2f}'.
→format(t[step], dx, dt), fontsize=10)
      ax.set xlim([min(x), max(x)])
      ax.set_ylim([0, 1])
      fig.tight layout()
      fig.legend(loc='upper right', bbox_to_anchor=(1.25, 0.925),
⇒frameon=False, fontsize=8)
      plt.show()
  # Return the values at the maximum time
  return x, y[-1, :], exact
```

Execute and plot the runs

```
[]: # Wave speed (c)
     c = 0.1
     # Spatial increments
     dxs = [1/20, 1/40, 1/80, 1/160]
     # Strings to represent spatial increments (for plotting purposes only)
     dxs_str = ['1/20', '1/40', '1/80', '1/160']
     # Courant number and Courant number list
     cfls = [0.1, 0.2, 0.4, 0.8]
     # Maximum time
     t max = 50
     # Boolean to control prints to console
     printout = False
     # Computation mode: variable dx or CFL values
     dx_calc = False
     # Initialize array to hold each run's output array
     values = {}
     # Iterate over dx values
     if dx_calc:
        cfl = 0.1
        for i, dx in enumerate(dxs):
             # End step
             end_step = int(t_max/(cfl*dx/c))
             # Perform calculation, plot individual results
            arr = advection_p2(dx, c, cfl, t_max, plot=False, plot_step=[0],__
      →printout=printout)
             # Store array
            values['dx = \{0\}'.format(dxs_str[i])] = {'x': arr[0], 'y': arr[1], ___
      # Iterate over Courant number values
     else:
        dx_index = 2
        dx = dxs[dx index]
        for i, cfl in enumerate(cfls):
             # End step
             end_step = int(t_max/(cfl*dx/c))
             # Perform calculation, plot individual results
            arr = advection_p2(dx, c, cfl, t_max, plot=False, plot_step=[end_step],__
      →printout=printout)
             # Store array
            values['CFL = {0}'.format(cfl)] = {'x': arr[0], 'y': arr[1], 'exact':
      →arr[2]}
```

```
''' Plotting. '''
# Initialize figure
fig, ax = plt.subplots(figsize=(4, 3))
# Define formatting index, color list, marker list
f_index, markers = 0, ['o', '^', '+', '2']
# Set color cycling
plt.rcParams["axes.prop_cycle"] = plt.cycler('color', plt.cm.Blues(np.
 →linspace(0.375, 1, len(dxs))))
# Plot the actual values
for key, value in values.items():
    im = ax.plot(value['x'], value['y'], marker=markers[f_index], markersize=5,__
 ⇔ls='--', label=key)
   if key == 'dx = 1/160' or key == 'CFL = 0.1':
        ax.plot(value['x'], value['exact'], c='r', label='exact', zorder=0)
   f index += 1
# Metadata
if dx_calc:
   ax.set_title('Courant number: {0:.1f}'.format(cfl), fontsize=10)
else:
   ax.set_title('dx = {0}'.format(dxs_str[dx_index]), fontsize=10)
ax.set_xlim([min(value['x']), max(value['x'])])
ax.set_ylim([0, 1])
fig.tight_layout()
fig.legend(loc='upper right', bbox_to_anchor=(1.25, 0.925), frameon=False,_u
⇔fontsize=8)
# plt.savefig('figs/p2c1.png', dpi=300, bbox_inches='tight')
# Initialize figure
fig, ax = plt.subplots(figsize=(4, 3))
# Collect RMS values
rmse = [rms(value['y'], value['exact']) for _, value in values.items()]
print(rmse)
# Plot the actual values
im = ax.loglog(dxs, rmse, marker='o', color='k')
ax.set_xticks([0.005, 0.01, 0.05])
ax.set_ylim([1e-3, 1e0])
# Metadata
ax.set ylabel('RMSE')
ax.set_xlabel('dx') if dx_calc else ax.set_xlabel('CFL')
ax.set_title('Root mean squared error', fontsize=10)
plt.gca().invert_xaxis()
fig.tight_layout()
# plt.savefig('figs/p2c2.png', dpi=300, bbox_inches='tight')\
```

0.0.3 Problem 3

(a) Derive coefficient values that give a 3rd-order accurate scheme in space and time.

(a) (i). Derive a 3rd-order accurate scheme for space.

Test the coefficients to ensure 3rd-order accuracy.

```
[806]: a, b, c, d, e = -1/6, 1, -1/2, -1/3, 1

#

print('Oth-order: {0:.3f}'.format(a + b + c + d + 0*e))
# '

print('1st-order: {0:.3f}'.format(-2*a - b + 0*c + d + e))
# ''

print('2nd-order: {0:.3f}'.format(2*a + (1/2)*b + 0*c + (1/2)*d + 0*e))
# '''

print('3rd-order: {0:.3f}'.format((-4/3)*a - (1/6)*b + 0*c + (1/6)*d + 0*e))
# '''

print('4th-order: {0:.3f}'.format((16/24)*a + (1/24)*b + 0*c + (1/24)*d + 0*e))
```

Oth-order: 0.000 1st-order: 0.000 2nd-order: 0.000 3rd-order: 0.000 4th-order: -0.083

- (a) (ii) Do this for the spatial representation of time of the $\partial^2 \psi / \partial x^2$ term.
 - Trial 1: do regular Taylor table
 - Trial 2: shift by a derivative up (instead of starting at ψ , start at $\partial \psi$).

```
[967]: from sympy import *

# Initialize symbolic variable
h = Symbol('h')
# Define the matrix system
```

 $\left\{ \left(\frac{2.22044604925031 \cdot 10^{-16}}{h^2}, \ -\frac{1.0}{h^2}, \ \frac{2.0}{h^2}, \ -\frac{1.0}{h^2} \right) \right\}$

Test the coefficients to make sure they're 3rd-order accurate

```
[865]: a, b, c, d, e = 0, -1, 2, -1, 1

#

print('Oth-order: {0:.3f}'.format(a + b + c + d + 0*e))
# '

print('1st-order: {0:.3f}'.format(-2*a - b + 0*c + d + 0*e))
# ''

print('2nd-order: {0:.3f}'.format(2*a + (1/2)*b + 0*c + (1/2)*d + e))
# '''

print('3rd-order: {0:.3f}'.format((-4/3)*a - (1/6)*b + 0*c + (1/6)*d + 0*e))
# '''

print('4th-order: {0:.3f}'.format((16/24)*a + (1/24)*b + 0*c + (1/24)*d + 0*e))
```

Oth-order: 0.000 1st-order: 0.000 2nd-order: 0.000 3rd-order: 0.000 4th-order: -0.083

(b) (iii) Do this for the spatial representation of time of the $\partial^3 \psi / \partial x^3$ term.

```
[884]: from sympy import *

# Initialize symbolic variable
h = Symbol('h')
# Define the matrix system
```

```
In the second seco
```

[884]: $\left\{ \left(\frac{1.0}{h^3}, -\frac{3.0}{h^3}, \frac{3.0}{h^3}, -\frac{1.0}{h^3} \right) \right\}$

Test the coefficients to make sure they're 3rd-order accurate

```
[887]: a, b, c, d, e = 1, -3, 3, -1, 1

#

print('Oth-order: {0:.3f}'.format(a + b + c + d + 0*e))
# '

print('1st-order: {0:.3f}'.format(-2*a - b + 0*c + d + 0*e))
# ''

print('2nd-order: {0:.3f}'.format(2*a + (1/2)*b + 0*c + (1/2)*d + 0*e))
# '''

print('3rd-order: {0:.3f}'.format((-4/3)*a - (1/6)*b + 0*c + (1/6)*d + e))
# '''

print('4th-order: {0:.3f}'.format((16/24)*a + (1/24)*b + 0*c + (1/24)*d + 0*e))
```

Oth-order: 0.000 1st-order: 0.000 2nd-order: 0.000 3rd-order: 0.000 4th-order: 0.500

0.0.4 Auxiliary functions

Root mean square (RMS) calculation

[947]:

Cyclic tridiagonal solver

```
[946]: def cyc_tridiag(jmx, a, b, c, f):
           ''' Written by D. Durran, translated by S. Ditkovsky. '''
           # jmx = dimension of all arrays
           \# a = sub (lower) diagonal
           # b = center diagonal
           # c = super (upper) diagonal
           # f = right hand side
           fmx = f[-1]
           # Create work arrays
           q = np.empty(jmx)
           s = np.empty(jmx)
           #forward elimination sweep
           q[0] = -c[0]/b[0]
           f[0] = f[0]/b[0]
           s[0] = -a[0]/b[0]
           for j in range(jmx-1):
               p = 1./(b[j+1] + a[j+1]*q[j])
               q[j+1] = -c[j+1]*p
               f[j+1] = (f[j+1] - a[j+1]*f[j])*p
               s[j+1] = -a[j+1]*s[j]*p
           #Backward pass
           q[-1] = 0.0
           s[-1] = 1.0
           for j in reversed(range(jmx-1)):
               s[j] = s[j] + q[j]*s[j+1]
               q[j] = f[j] + q[j]*q[j+1]
           #final pass
           f[-1] = (fmx-c[-1]*q[0] - a[-1]*q[-2])/(c[-1]*s[0] + a[-1]*s[-2] + b[-1])
           for j in range(jmx-1):
               f[j] = f[-1]*s[j] + q[j]
```

return f