# Constraining the dense matter equation of state using neutron star observations

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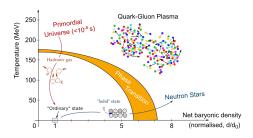
Mentors: Oleg Korobkin, Soumi De/ Rahul Somasundaram (T5/T2)

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## **Motivation**

The strong force is still poorly understood at high densities and low temperatures.



 Understanding the composition of neutron stars can thus give important information on the behavior of the strong interaction.

# Compact Object as an Astrophysical Laboratory

- Neutron stars are the dead remnants of massive stars.
- They can be up to twice as massive as our sun.
- A teaspoon of neutron star matter would weigh 4 billion tons.

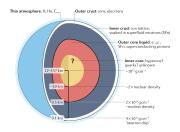


Figure: Typical Neutron Star structure, P.C: Google

The macroscopic properties of nuclear matter are described by equation of state .

## Equation of State(EoS)

Pressure and Energy density relation of nuclear matter.

- EoS 1: Meta Model+Speed of sound approach
- EoS 2: MIT Bag Model describes quark matter, strange matter on overall density scale.

# EoS 1: Meta Model up to $(1-2)n_{sat}$

 The general properties of nuclear interactions are often characterized in terms of the nuclear empirical parameters:

$$e_{sat} = E_{sat} + \frac{1}{2!}K_{sat}x^2 + \frac{1}{3!}Q_{sat}x^3 + \frac{1}{4!}Z_{sat}x^4 + \cdots$$
 (1)

$$e_{sym} = E_{sym} + L_{sym}x + \frac{1}{2!}K_{sym}x^2 + \frac{1}{3!}Q_{sym}x^3 + \frac{1}{4!}Z_{sym}x^4 + \cdots$$
 (2)

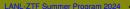
$$x = \frac{n - n_{sat}}{3n_{sat}} \tag{3}$$

Energy per nucleon in nuclear matter, defined as

$$\frac{E}{A} \approx e_{sat} + e_{sym} \tag{4}$$

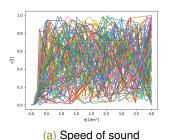
$$\epsilon = n\frac{E}{A}; \quad p = -\frac{dE}{dV} = n^2 \frac{dE/A}{dn}$$

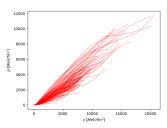
J. Margueron et al., Phys.Rev.C 97 (2018) 2, 025805



# Speed of Sound approach

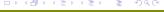
$$n_{break} < n < 25n_{sat} \quad p(n) = p(n_i) + \int_{n_i}^n c_s^2(n')\mu(n')dn'$$
$$\epsilon(n) = \epsilon(n_i) + \int_{n_i}^n \mu(n')dn'$$





(b) pressure-energy density plot

R. Somasundaram et al, Phys. Rev. C 107, 025801



## MIT Bag Model

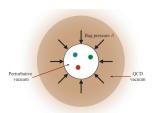


Figure: Hadron bag

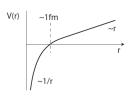


Figure: QCD Potential

# Quarks are confined inside the bag! Non-interacting.

- number density:  $n = \sum_{u,d,s} \int_0^\infty \frac{g}{2\pi^3} f(p) d^3 p$
- Energy density:  $\epsilon = \sum_{u,d,s} \int_0^\infty \frac{g}{2\pi^3} \epsilon(p) f(p) d^3 p + B$
- Pressure:  $P = \sum_{u,d,s} \int_0^\infty \frac{g}{2\pi^3} \frac{p^2}{3\epsilon} f(p) d^3 p B$
- g is the degeneracy factor=  $2 \times 3$ (spin× color)

$$f(p) = \frac{1}{e^{\frac{\epsilon - \mu}{T}} + 1} \tag{5}$$

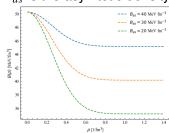
Quarks are appear to be quasi-free at short distance: K. Johnson, Acta Phys. Polon. B6

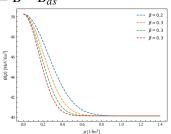
## Density dependent bag pressure

#### Density dependent Bag pressure

$$B(\rho) = B_{as} + \frac{\Delta B}{2} \left[ 1 + e^{-\beta(\rho/\rho_0)^2} \right]$$
 (6)

 $B_{as}$  is the asymtotic density,  $\Delta B = B - B_{as}$ 





## Zero Temperature Solution

We construct EOS making various approximation: at zero temperature quarks are filled up to Fermi level with Fermi energy  $E_f < \mu_f$ , distribution function approaches to unity, means all the Fermi level below Fermi energy are completely filled and abobe are empty.

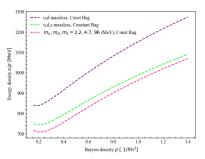
$$P = -B(\rho) + \sum_{f} \frac{1}{4\pi^2} \left[ \mu_f k_f \left( u_f^2 - \frac{5}{2} m_f^2 \right) + \frac{3}{2} m_f^4 ln \left( \frac{\mu_f + k_f}{m_f} \right) \right]$$
 (7)

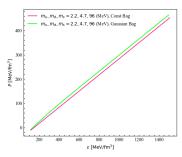
$$\epsilon = B(\rho) + \sum_{f} \frac{3}{4\pi^2} \left[ \mu_f k_f \left( u_f^2 - \frac{5}{2} m_f^2 \right) - \frac{1}{2} m_f^4 ln \left( \frac{\mu_f + k_f}{m_f} \right) \right]$$
 (8)

$$\rho = \sum_{f} \frac{k_f^3}{3\pi^2} \tag{9}$$

 $k_f$ =Fermi momentum of quark flavor f given by  $k_f = \sqrt{\mu_f^2 - m_f^2}$ ,  $\mu_f$  is the chemical potential of flavor f.

## **Quark Matter EOS**





- Minimum energy per baryon occurs correspond to the zero pressure across the baryon density ensuring thermodynamic self-consistency.
- Massive quark flavors appreciably satisfies the stability window of  $^{56}Fe = 928 \text{ MeV}(\text{E/A})$  up to several times of nuclear saturation density.

## Stellar Structure Equations

Tolman-Oppenheimer-Volkoff(TOV) Equation for compact object generally written as,

$$\frac{dm(r)}{dr} = \frac{4\pi r^2 \epsilon(r)}{c^2} \tag{10}$$

$$\frac{dP}{dr} = -\frac{G\epsilon(r)m(r)}{c^2r^2} \left[ 1 + \frac{P(r)}{\epsilon(r)} \right] \left[ 1 + \frac{4\pi r^3 P(r)}{m(r)c^2} \right] \left[ 1 - \frac{2Gm(r)}{c^2r} \right]^{-1}$$
(11)

The boundary conditions:  $P(0) = P_c$  and P(R=r)=0 with m(r)=M to ensure hydrostatic equilibrium.

We aslo worked on enthaply formulation of solving TOV: Fasten enough

## **Tidal Deformation**

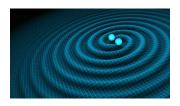


Figure: Binary neutron star inspiral.

$$Q_{ij} = -\Lambda \epsilon_{ij}$$

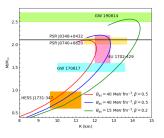
 $\epsilon_{ij}$ : External tidal field,

 $Q_{ij}$ : Quadrupole moment.

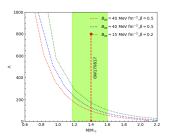
Tidal deformability:  $\Lambda = \frac{2}{3}k_2(R/M)^5$ ,



### Results: 1



(a) Mass-Radius relation

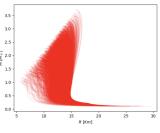


(b) Tidal deformability-Mass relation

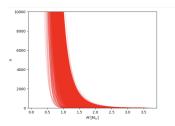
Figure: Zero Temperature Bag Model, B<sup>1/4</sup>=(130-140) MeV

### Results: 2

We have constructed 100000 EOS based on Meta Model and speed of sound approach and estimated the Mass-radius and Tidal deformabilities for each individual EOS.



(a) Mass-Radius relation

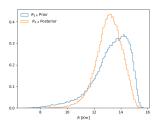


(b) Tidal deformability-Mass relation

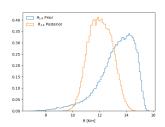
### NICER Observations

The likelihood of a certain EOS given an M-R posterior p(M,R|NICER)from a NICER measurement can be written as,

$$\mathcal{L}(EOS|NICER) = \int_{0}^{M_{TOV}} dM \, p(M, R(EOS)|NICER)$$
 (12)



(c) J0030+451 2019 M-R dataset



(d) PSR J0437-4715 2024 M-R dataset

Figure: EOS inference based on the NICER measurement of Pulsars

# **Next Step**

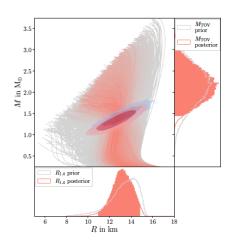


Figure: M-R posterior for PSR J0030+451

# Learning Outcome and Conclusion

- Studied Neutron star EOS formulation and hand on experience of solving stellar structure equations for structural studies of NS: i.e nuclear matter equation of state.
- With density dependent Bag model, we find equatorial radius of 11 – 13.7 km and Tidal deformability of 200 in dimensionless unit, which consistent with GW observations.
- We also able to incorporate compact object mass of mass 0.77  $M\odot$  within radius 9.7 11.2 km
- Using the PSR datasets we find the the likelihood of MM+CSM model EOS and canonical posterior distribution on an equatorial radius at mass 1.4  $M_{\odot}$ .

I would like express my sincere gratitude to my mentors: Oleg, Soumi and Rahul. Thank you all!

**Questions?**