

Integrating Kinetic and Gravitational Time Differentials With Total Energy

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The results from the previous investigations have lead to the following causal relationship: changes in inertial time differentials (ITDs) are related to work done, either changes in specific kinetic or specific potential energy. These relationships are summarized in Equation (1).

$$\frac{dt}{dt'} = \sqrt{1 - \frac{w}{e_T}} \quad (1a)$$

$$\frac{dt}{dt'} = \sqrt{1 - \frac{\Delta e_K}{e_T}} \quad (1b)$$

$$\frac{dt}{dt'} = \sqrt{1 - \frac{\Delta e_P}{e_T}} \quad (1c)$$

Additionally, previous investigations revealed an update to the relativistic mass, kinetic energy and total energy models, as summarized in Equation (2).

$$m = \frac{m_0}{1 - v^2/c^2} = \text{Invariant} \quad (2a)$$

$$\Delta K = \frac{1}{2}mv^2 \quad (2b)$$

$$E = \frac{1}{2}mc^2 = \frac{1}{2}m_0c^2 + \frac{1}{2}mv^2 \quad (2c)$$

I now intend to relate changes in ITDs to changes in total specific energy, and then I intend to update the total energy model to include potential energy.

ITDs can be written in terms of fractions of total specific energy for both specific potential energy and specific kinetic energy, which means a relationship between ITD and total specific energy exists. Changes in specific kinetic or potential energy are related to changes in ITDs, but this is only half of the picture because we tacitly assumed all else remained equal. Now we test what if all else does not remain equal to discover a more precise cause to changes in ITDs.

In reviewing Equation (1), simple analysis reveals that transferring some amount of specific kinetic energy to some amount of specific potential energy (or vice versa) would not cause an overall change in the ITD.

For this proof, consider an object that starts with some amount of specific potential energy, who then transfers all of it to kinetic energy (no longer in a gravity potential somehow).

Proof :

Let $e_P > 0$.

$$\text{Let } \frac{1}{\gamma} = \frac{dt}{dt'} \quad (3a)$$

$$\frac{1}{\gamma_P^2} = 1 - \frac{\Delta e_P}{e_T} \quad (3b)$$

$$1 - \frac{1}{\gamma_P^2} = \frac{\Delta e_P}{e_T} \quad (3c)$$

$$\left(1 - \frac{1}{\gamma_P^2}\right) e_T = \Delta e_P = \Delta e_K \quad (3d)$$

$$\left(1 - \frac{1}{\gamma_P^2}\right) e_T = \Delta e_K \quad (3e)$$

$$1 - \frac{1}{\gamma_P^2} = \frac{\Delta e_K}{e_T} \quad (3f)$$

$$\frac{1}{\gamma_P^2} = 1 - \frac{\Delta e_K}{e_T} \quad (3g)$$

$$\frac{1}{\gamma_P^2} = \frac{1}{\gamma_K^2} \blacksquare \quad (3h)$$

Invoking the method of agreement: observing that changes in specific potential energy and changes in specific kinetic energy induced no changes in ITD, proves inductively that they are not the fundamental causes to changes in ITDs—they each play half a role.

The same change in total specific energy caused the same change in ITDs proves inductively, via method of agreement, that changes in ITD are caused by a change in total specific energy, and vice versa.

Next, I derive the change in total specific energy's relationship to change in ITD. This derivation begins by trying to solve the effective ITD for an object stationary within a gravitational field, and determining how a change in kinetic energy, as measured from the stationary position within the gravitational field, affects the ITD. This situation is shown in Figure 1.

The desire is to measure total effective ITD. The ITD from the initial inertial frame to the stationary point inside a gravitational potential can be measured. Also, the ITD between this stationary point and the moving point inside the gravity potential can be measured. These two ITDs are relatable to the total effective ITD using the chain rule, as shown in Equation (4).

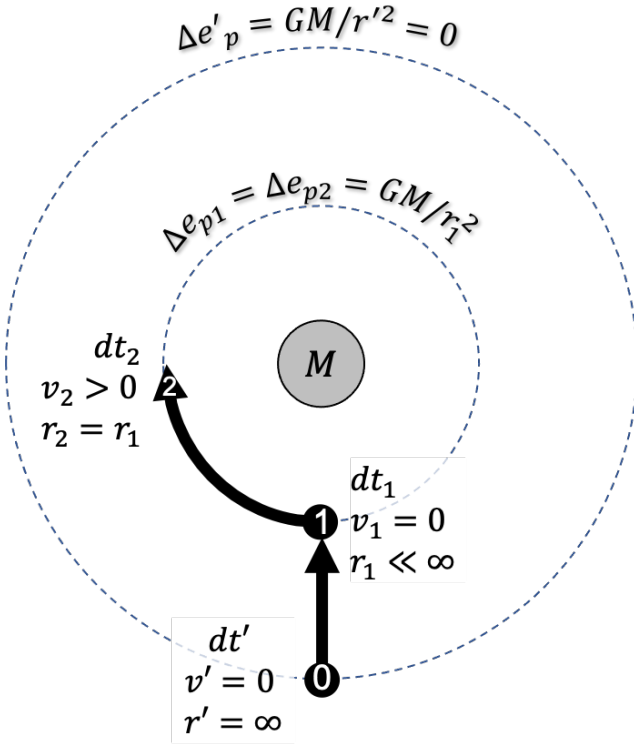


Figure 1. Total time differential effects example.

Lastly, I can now use these tools to update the total energy model in Equation (2) to include an external potential energy term as shown in Equation (6), where v' is velocity measured from an initial inertial frame infinitely far from a gravity potential (as opposed to from within the gravitational potential).

$$\frac{1}{\gamma_T} = \sqrt{1 - \frac{\Delta e_T}{e_T}} \quad (6a)$$

$$\frac{1}{\gamma_T^2} = 1 - \frac{\Delta e_T}{e_T} \quad (6b)$$

$$e_T = \frac{1}{\gamma_T^2} e_T - \Delta e_T \quad (6c)$$

$$E_T = \frac{1}{\gamma_T^2} E_T + \Delta E_T \quad (6d)$$

$$\frac{1}{2} mc^2 = \frac{1}{\gamma_T^2} \frac{1}{2} mc^2 + \Delta E_P + \Delta E_K \quad (6e)$$

$$\frac{1}{2} mc^2 = \frac{1}{2} m_0 c^2 + \Delta E_P + \frac{1}{2} m v'^2 \blacksquare \quad (6f)$$

$$\frac{dt_1}{dt'} = \sqrt{1 - \frac{\Delta e_{P1}}{e_T}} \quad (4a)$$

$$\frac{dt_2}{dt_1} = \sqrt{1 - \frac{\Delta e_{K2}}{e_T}} \quad (4b)$$

$$\frac{dt_2}{dt'} = \frac{dt_1}{dt'} \frac{dt_2}{dt_1} = \sqrt{1 - \frac{\Delta e_{P1}}{e_T}} \sqrt{1 - \frac{\Delta e_{K2}}{e_T}} \quad (4c)$$

This can be generalized to any condition where an object inside a gravity potential gains some kinetic energy as measured inside that potential, as shown in Equation (5):

$$\text{Let : } \frac{1}{\gamma_P} = \frac{dt_P}{dt'} = \sqrt{1 - \frac{\Delta e_P}{e_T}} \quad (5a)$$

$$\text{Let : } \frac{1}{\gamma_K} = \frac{dt_K}{dt_P} = \sqrt{1 - \frac{\Delta e_{K/P}}{e_T}} \quad (5b)$$

$$\frac{1}{\gamma_T} = \sqrt{1 - \frac{\Delta e_P + \Delta e_{K/P} - \Delta e_K \frac{\Delta e_P}{e_T}}{e_T}} \quad (5c)$$

$$\frac{1}{\gamma_T} = \sqrt{1 - \frac{\Delta e_P + (1 - \frac{\Delta e_P}{e_T}) \Delta e_{K/P}}{e_T}} \quad (5d)$$

$$\frac{1}{\gamma_T} = \sqrt{1 - \frac{\Delta e_P + \frac{1}{\gamma_P^2} \Delta e_{K/P}}{e_T}} \quad (5e)$$

$$\frac{1}{\gamma_T} = \sqrt{1 - \frac{\Delta e_P + \Delta e_K}{e_T}} = \sqrt{1 - \frac{\Delta e_T}{e_T}} \blacksquare \quad (5f)$$