Density of States of Hodge Laplacians: for Simplicial Sparsification

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Abstract:

Keywords: density of states, simplicial complexes, sparsification, generalized effective resistance

Introduction

Well, it would be nice to have one. You know, somewhere here, maybe...

Related works.

- ♦ benson
- ♦ osting, speilman
- ♦ kyng and fast solvers
- ♦ cohen and geoemetrical solvers
- ♦ hecs precon

Main Contributions. Our main contributions are as follows:

- we show that sparsification measure (which is necessary to sample sparse spectral approximation of an up-Laplacian operator of simplicial complex) is directly related to local densities of states of a higher-order down-Laplacian;
- ♦ we develop a novel method of kernel ignoring decomposition (KID) for efficient approximation of LDoS;
- finally, we exploit existing error estimates to derive KID parameters sufficient for an arbitrary approximation quality and demonstrate the corresponding competitive algorithmic complexity for the method.

Outline. The rest of the paper is organized as follows: we provide the brief introduction of simplicial complexes and its densities of states as a fast computable proxies of spectral information, CREF. CREF states the sparsification result and demonstrates its connection to the local spectral densities with CREF describing a novel approach which efficiently computes the sparsifying probability measure. We provide numerical experiments in CREF, CREF concludes.

Preliminaries

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I. Simplicial Complexes, [Lim20]

Simplicial complexes Let $V = \{v_1, \dots v_{m_0}\}$ be the set of nodes. Each subset of nodes $\sigma = \{v_{i_1}, \dots v_{i_k}\}$ is known as **simplex** of order k-1 with all its maximal proper subsets knowns as **faces** of σ . Then a collection of simplices $\mathcal K$ is a **simplicial complex** if and only if each simplex σ enters $\mathcal K$ with all its faces (inclusion rule). As a result, each simplicial complex can be written as a collections of simplices of fixed orders, $\mathcal K = \{\mathcal V_0(\mathcal K), \mathcal V_1(\mathcal K), \mathcal V_2(\mathcal K), \dots\}$ where $\mathcal V_0(\mathcal K)$ is a set of 0-simplices (nodes), $\mathcal V_1(\mathcal K)$ is a set of edges, $\mathcal V_2(\mathcal K)$ of triangles, and so on, with $m_k = |\mathcal V_k(\mathcal K)|$ denoting the number of simplices in the corresponding set.

Boundary operators and Laplacians Due to the inclusion principle, simplices inside $\mathcal K$ form a nested structure through boundary relation; specifically, one can map a k-order simplex to a collection of (k-1)-order simplices on its border. Formally, let $\mathcal C_k$ be a linear space of all formal sums of simplices in $\mathcal V_k(\mathcal K)$; then the boundary map $\mathcal B_k$ is defined as an alternating sum:

$$B_k: \mathcal{C}_k \mapsto \mathcal{C}_{k-1}, \qquad B_k[v_1, v_2, \dots v_k] = \sum_{i=1}^k (-1)^{i-1}[v_1, \dots v_{i-1}, v_{i+1}, \dots v_k]$$
 (Eqn. 1)

such that boundary of a boundary is 0, $B_k B_{k+1} = 0$, (fundamental lemma of homology). Provided one fixes an order of simplices and their orientation (lexicographically, so the simplicity's sake), one can define the following entities describing the higher-order structure of \mathcal{K} :

Def. 1 (Homology group and higher-order Laplacian) Since $\operatorname{im} B_{k+1} \subset \ker B_k$, the quotient space $\mathcal{H}_k = \frac{\ker B_k}{\operatorname{im} B_{k+1}}$, known as k-th homology group, is correctly defined and the following isomorphisms hold

$$\mathcal{H}_k \cong \ker B_k \cap \ker B_{k+1}^{\top} \cong \ker \left(B_k^{\top} B_k + B_{k+1} B_{k+1}^{\top} \right).$$

The matrix $L_k = B_k^\top B_k + B_{k+1} B_{k+1}^\top$ is called the k-th order **Laplacian operator**; the two terms $L_k^{\downarrow} = B_k^\top B_k$ and $L_k^{\uparrow} = B_{k+1} B_{k+1}^\top$ are referred to as the **down-Laplacian** and the **up-Laplacian**, respectively.

The homology group \mathcal{H}_k describes the k-th topology of the simplicial complex \mathcal{K} : $\beta_k = \dim \mathcal{H}_k = \dim \ker L_k$ coincides exactly with the number of k-dimensional holes in the complex, known as the k-th Betti number. In the case k=0, the operator $L_0=L_0^\uparrow$ is exactly the classical graph Laplacian whose kernel corresponds to the connected components of the graph, while $L_0^{\downarrow}=0$. For k=1 and k=2, the elements of $\ker L_1$ and $\ker L_2$ describe the simplex 1-dimensional holes and voids respectively, and are frequently used in the analysis of trajectory data, [SBH+20, BGL16].

redo

Weighted case and spectral structures Considering weighted generalization of simplicial complexes, one needs to preserve the fundamental lemma of homology, $B_k B_{k+1} = 0$. As a result, let us assume the collection of weight matrices $W_k \in \mathbb{R}_{m_k \times m_k}$ such that W_k is diagonal with each $(W_k)_{ii} > 0$ containing the weight of the i-the simplex in $\mathcal{V}_k(\mathcal{K})$. Then a generic weighting scheme for boundary operators can be written as:

$$B_k \mapsto W_{k-1}^{-1} B_k W_k$$
 (Eqn. 2)

Note that with such weighting scheme the dimensionality of the homology group is preserved, dim ker $L_k = \dim \ker \widehat{L}_k$, [GST23].

At each order k, weighted boundary operators B_k define the space decomposition (Hodge decomposition):

$$\mathbb{R}^{m_k} = \overbrace{\operatorname{im} B_k^{\top} \oplus \operatorname{ker} \left(B_k^{\top} B_k + B_{k+1} B_{k+1}^{\top} \right) \oplus \operatorname{im} B_{k+1}}^{\operatorname{ker} B_k}$$
 (Eqn. 3)

implying that that kernels of up- and down-Laplacians are necessarily non-trivial and high-dimensional, $\ker L_k^{\downarrow} \supseteq \operatorname{im} B_{k+1}$ and $\ker L_k^{\uparrow} \subseteq \operatorname{im} B_k^{\intercal}$. More explicit relation between these subspaces is known as a principle spectral inheritance, [GST23].

II. Networks' Density of States, [DBB19]

Spectral densities As we established above, the spectral structure of the higher-order Laplacian operators carry meaningful information of the complex's topological structure, YADA YADA

At the same time, the overall cost of extracting full spectral information from corresponding Laplacian operators L_k is computationally expensive with a complexity $\mathcal{O}(m_k^3)$. Instead, one can consider the following functional description:

the usual song

Def. 2 **(Density of States)** Fro a given symmetric matrix $A = Q\Lambda Q^{\top}$ with $Q^{\top}Q = I$ and diagonal $\Lambda = \text{diag}(\lambda_1, \dots \lambda_n)$, the **spectral density** or **density of states** (DoS)

$$\mu(\lambda \mid A) = \frac{1}{n} \sum_{i=1}^{n} \delta(\lambda - \lambda_i)$$
 (Eqn. 4)

Additionally, let \mathbf{q}_i be a corresponding unit eigenvector of A (such that $A\mathbf{q}_i = \lambda_i \mathbf{q}_i$ and $Q = (\mathbf{q}_1 \mid \mathbf{q}_2 \mid \cdots \mid \mathbf{q}_n)$); then one can define a set of local (entry-wise) densities of states (LDoS):

$$\mu_k(\lambda \mid A) = \sum_{i=1}^n \left| \mathbf{e}_k^{\mathsf{T}} \mathbf{q}_i \right|^2 \delta(\lambda - \lambda_i)$$
 (Eqn. 5)

with \mathbf{e}_k being the corresponding versor.

Here DoS function $\mu(\lambda \mid A)$ contains the overall spectrum of the operator A while the family of LDoS $\mu_k(\lambda \mid A)$ describe the contribution of the simplex $\sigma_k \in \mathcal{V}_k(\mathcal{K})$ to the spectral information.

Finally, one should note that by its definition DoS and LDoS are generalized functions, hence the quality of their computation is difficult to assess directly; instead, one considers its histogram representations:

$$h_i = \int_{x_i}^{x_i + \Delta_X} \mu(\lambda \mid A) d\lambda, \qquad h_i^{(k)} = \int_{x_i}^{x_i + \Delta_X} \mu_k(\lambda \mid A) d\lambda \tag{Eqn. 6}$$

which correspond to discretized output of the convolution of spectral densities with a smooth approximation of identity function $K_{\Delta h}$, $[\mu(\lambda \mid A) * K_{\Delta h}]$.

Kernel Polynomial Method (KPM) By its very definition, the introduction of the spectral densities $\mu(\lambda \mid L_k)$ and $\{\mu_k(\lambda \mid A)\}$ requires the complete spectral information of the original operator L_k and, hence, is not immediately

computationally more efficient. At the same time, one can use the functional nature of DoS and LDoS to obtain approximation through expansion.

Namely, let $T_m(x)$ be Chebyhev polynomials of the first kind. Specifically,

$$T_0(x) = 1$$
, $T_1(x) = x$, $T_{m+1}(x) = 2xT_m(x) - Tm - 1(x)$ (Eqn. 7)

which as an orthogonal (with respect to the weight function $w(x)=\frac{2}{(1+\delta_{0m})\pi\sqrt{1-x^2}}$) on [-1,1]. Let us transform the operator $A\to H$ such that $\sigma(H)\in[-1,1]$ (e.g. $H=\frac{2}{\lambda_{\max}}A-I$)

Sparsification of Simplicial Complexes

Th III.II.1

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(Simplicial Sparsification, [OPW22]) Let ${\mathcal K}$ be a simplicial complex restricted to its p-skeleton, $\mathcal{K}=\bigcup_{i=0}^p \mathcal{V}_i(\mathcal{K})$. Let $L_k^\uparrow(\mathcal{K})$ be its k-th up-Laplacian and let $m_k=|\mathcal{V}_k(\mathcal{K})|$. For any $\varepsilon>0$, a sparse simplicial complex ${\cal L}$ can be sampled as follows:

- (1) compute the probability measure ${f p}$ on ${\cal V}_{k+1}({\cal K})$ proportional to the generalized resistance vector $\mathbf{r} = \operatorname{diag}\left(B_{k+1}^{\top}(L_k^{\uparrow})^{\dagger}B_{k+1}\right)$, where $\operatorname{diag}(A)$ denotes the vector of the diagonal entries of A;
- (2) sample q simplices τ_i from $\mathcal{V}_{k+1}(\mathcal{K})$ according to the probability measure **p**, where q is chosen so that $q(m_k) \geq 9C^2m_k\log(m_k/\varepsilon)$, for some absolute constant C > 0;
- (3) form a sparse simplicial complex ${\cal L}$ with all the sampled simplexes of order k and all its faces with the weight $\frac{w_{k+1}(\tau_i)}{q(m_k)\mathbf{p}(\tau_i)}$; weights of repeated simplices are accumulated.

Then, with probability at least 1/2, the up-Laplacian of the sparsifier $\mathcal L$ is ε -close to the original one, i.e. it holds $L_k^{\uparrow}(\mathcal{L}) \approx L_k^{\uparrow}(\mathcal{K})$.

The bottleneck of the subsampling above is the construction of the appropriate measure \mathbf{p} or, more precisely, generalized effective resistance \mathbf{r} . Indeed, one needs a fast pseudo-inverse operator $\left(L_k^\uparrow\right)^\dagger$ in order to compute

Th III.II.2

(GER through LDoS) For a given simplicial complex *mcK* with the *k*-th order up-Laplacian $L_k^{\uparrow} = B_{k+1} W_{k+1}^2 B_{k+1}^{\dagger}$, a generalized effective resistance r can be computed through family of local densities of states $\{\mu_i(\lambda \mid L_{k+1}^{\downarrow})\}$:

$$\mathbf{r}_{i} = \int_{\mathbb{R}} (1 - \mathbb{1}_{0}(\lambda)) \mu_{i}(\lambda \mid L_{k+1}^{\downarrow}) d\lambda$$

$$(\lambda))\mu_i(\lambda \mid L_{k+1}^{\downarrow})d\lambda$$
 (Eqn. 8)

weighted egdges === BAD ?

we need to assume somewhere $W_k = I$

for simplicity and

do not forget about

Proof Let $B_{k+1}W_{k+1} = USV^{\top}$ where S is diagonal and invertible and both U and V are orthogonal (so it is a truncated SVD decomposition of $B_{k+1}W_{k+1}$ matrix with eliminated obsolete kernel). Then:

$$(L_k^{\uparrow})^{\dagger} = \left(B_{k+1}W_{k+1}^2B_{k+1}^{\top}\right)^{\dagger} = \left(US^2U^{\top}\right)^{\dagger} = US^{-2}U^{\top}$$
 (Eqn. 9)

$$\mathbf{r} = \operatorname{diag}\left(W_{k+1}B_{k+1}^{\top}(L_{k}^{\uparrow})^{\dagger}B_{k+1}W_{k+1}\right) =$$

$$= \operatorname{diag}\left(VSU^{\top}US^{-2}U^{\top}USV^{\top}\right) = \operatorname{diag}\left(VV^{\top}\right)$$
(Eqn. 10)

As a result, $\mathbf{r}_i = \|V_i.\|^2 = \sum_j |v_{ij}|^2$, so the *i*-th entry of the resistance is defined by the sum of square of *i*-th components of eigenvectors \mathbf{v}_j of $L_{k+1}^{\downarrow} = W_{k+1}B_{k+1}^{\intercal}B_{k+1}W_{k+1}$ operator where $\mathbf{v}_j \perp \ker L_{k+1}^{\downarrow}$. Note that

$$\mu_{i}(\lambda \mid L_{k+1}^{\downarrow}) = \sum_{j=1}^{m_{k+1}} \left| \mathbf{e}_{i}^{\top} \mathbf{q}_{j} \right|^{2} \delta\left(\lambda - \lambda_{j}\right) = \sum_{j=1}^{m_{k+1}} \left| q_{ij} \right|^{2} \delta\left(\lambda - \lambda_{j}\right) \tag{Eqn. 11}$$

SO

$$\mathbf{r}_{i} = \|V_{i}.\|^{2} = \sum_{j} |v_{ij}|^{2} = \int_{\mathbb{R}\backslash\{0\}} \sum_{j=1}^{m_{k+1}} |q_{ij}|^{2} \delta\left(\lambda - \lambda_{j}\right) d\lambda =$$

$$= \int_{\mathbb{R}\backslash\{0\}} \mu_{i}(\lambda \mid L_{k+1}^{\downarrow}) d\lambda = \int_{\mathbb{R}} (1 - \mathbb{1}_{0}(\lambda)) \mu_{i}(\lambda \mid L_{k+1}^{\downarrow}) d\lambda$$
(Eqn. 12)

Rem II.1 (Sensitivity of the Sparsification vis-a-vis sampling measure p) Compare ε with $\frac{1}{m_2}$: if below, you are fiiiiiiine INSERT FIGURE HERE

IV KID for LDoS

Note that up- and down-Laplacians L_k^{\uparrow} and L_k^{\downarrow} by their definition have non-trivial kernels of high dimensionality. Indeed, recalling the Hodge decomposition :

ref up

$$\mathbb{R}^{m_k} = \overbrace{\operatorname{im} B_k^{\top} \oplus \ker \left(B_k^{\top} B_k + B_{k+1} B_{k+1}^{\top} \right) \oplus \operatorname{im} B_{k+1}}^{\ker B_k}$$
 (Eqn. 13)

the subspaces $\ker B_k = \ker L_k^{\downarrow}$ and $\ker B_{k+1}^{\top} = \ker L_k^{\uparrow}$ include at least $\operatorname{im} B_{k+1}$ and $\operatorname{im} B_{k+1}^{\top}$. As a result, the zero eigenvalue in the corresponding DoS and LDoS for both operators exhibits a dominating pick severely affecting the quality of KPM-approximation.

do we need an illustration

(Filtration of the kernels) As given by the Hodge decomposition above, the most part of the peak in L_k^{\downarrow} is explained by the elements of im B_{k+1} . Then one can avoid the dominating kernel pick in the corresponding LDoS family $\{\mu_k(\lambda \mid L_k^{\downarrow})$ by filtering im B_{k+1} out provided one can obtain an orthonormal basis of this subspace. As a result, the filtration of the kernel becomes the question of the efficient range finder of a sparse operator im B_{k+1} . Although we do not claim that computationally efficient range finder for such case does not exist, the majority of the methods at the current moment run into the bottleneck of the QR-decomposition which is far more computationally complex then one can allow in comparison with

the straightforward computation of GER for L_{ν}^{\uparrow} .

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Rem II.2

Theorem III.II.2 implies that one indeed does not need the full profile $\mu_k(\lambda \mid L_{k+1}^{\downarrow})$, but LDoS ignoring the kernel, $(1-\mathbb{1}_0)\mu_k(\lambda \mid L_{k+1}^{\downarrow})$. As a result, one can attempt to construct an approximation different from Chebyshev KPM that avoids the kernel of the operator in its entirety.

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Conclusion

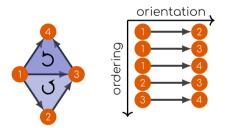
Motive filtration is still a thing

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III. Appendix

III.I Simplicial complexes



$$B_{2}[1,2,3] = \underbrace{(+1)[1,2]}_{\text{order}} + (-1)[1,3] + (+1)[2,3]$$

$$B_{2}[1,3,4] = \underbrace{(+1)[1,3]}_{\text{1st in order}} + (-1)[1,4] + (+1)[3,4]$$

Figure III.1: Example of the simplicial complex with ordering and orientation: nodes from $\mathcal{V}_0(\mathcal{K})$ in orange, triangles from $\mathcal{V}_2(\mathcal{K})$ in blue. Orientation of edges and triangles is shown by arrows; the action of \mathcal{B}_2 operator is given for both triangles.

III.II Proof of Theorem 1

Proof

Let $B_{k+1}W_{k+1} = USV^{\top}$ where S is diagonal and invertible and both U and V are orthogonal (so it is a truncated SVD decomposition of $B_{k+1}W_{k+1}$ matrix with eliminated obsolete kernel). Then:

$$(L_k^{\uparrow})^{\dagger} = \left(B_{k+1}W_{k+1}^2B_{k+1}^{\top}\right)^{\dagger} = \left(US^2U^{\top}\right)^{\dagger} = US^{-2}U^{\top}$$
 (Eqn. 14)

$$\mathbf{r} = \operatorname{diag}\left(W_{k+1}B_{k+1}^{\top}(L_{k}^{\uparrow})^{\dagger}B_{k+1}W_{k+1}\right) =$$

$$= \operatorname{diag}\left(VSU^{\top}US^{-2}U^{\top}USV^{\top}\right) = \operatorname{diag}\left(VV^{\top}\right)$$
(Eqn. 15)

As a result, $\mathbf{r}_i = \|V_i\|^2 = \sum_j |v_{ij}|^2$, so the *i*-th entry of the resistance is defined by the sum of square of *i*-th components of eigenvectors \mathbf{v}_j of $L_{k+1}^{\downarrow} = W_{k+1}B_{k+1}^{\top}B_{k+1}W_{k+1}$ operator where $\mathbf{v}_j \perp \ker L_{k+1}^{\downarrow}$. Note that

$$\mu_i(\lambda \mid L_{k+1}^{\downarrow}) = \sum_{j=1}^{m_{k+1}} \left| \mathbf{e}_i^{\top} \mathbf{q}_j \right|^2 \delta\left(\lambda - \lambda_j\right) = \sum_{j=1}^{m_{k+1}} \left| q_{ij} \right|^2 \delta\left(\lambda - \lambda_j\right) \tag{Eqn. 16}$$

$$\mu_{i}(\lambda \mid L_{k+1}^{\downarrow}) = \sum_{j=1}^{m_{k+1}} \left| \mathbf{e}_{i}^{\mathsf{T}} \mathbf{q}_{j} \right|^{2} \delta \left(\lambda - \lambda_{j}\right) = \sum_{j=1}^{m_{k+1}} \left| q_{ij} \right|^{2} \delta \left(\lambda - \lambda_{j}\right)$$

$$\mathbf{r}_{i} = \|V_{i}\|^{2} = \sum_{j} |v_{ij}|^{2} = \int_{\mathbb{R} \setminus \{0\}} \sum_{j=1}^{m_{k+1}} \left| q_{ij} \right|^{2} \delta \left(\lambda - \lambda_{j}\right) d\lambda =$$

$$= \int_{\mathbb{R} \setminus \{0\}} \mu_{i}(\lambda \mid L_{k+1}^{\downarrow}) d\lambda = \int_{\mathbb{R}} (1 - \mathbb{1}_{0}(\lambda)) \mu_{i}(\lambda \mid L_{k+1}^{\downarrow}) d\lambda$$
(Eqn. 17)