

# Microburst Scale Size Distribution Derived with AeroCube-6

M. Shumko<sup>1</sup>, J. Sample<sup>1</sup>, A. Johnson<sup>1</sup>, D.L. Turner<sup>2</sup>, T.P. O'Brien<sup>2</sup>, B.A.  
Griffith<sup>1</sup>, J.B. Blake<sup>2</sup>, O. Agapitov<sup>3</sup>, S. G. Claudepierre<sup>2</sup>

<sup>1</sup>Department of Physics, Montana State University, Bozeman, Montana, USA

<sup>2</sup>Space Science Applications Laboratory, The Aerospace Corporation, El Segundo, California, USA

<sup>3</sup>Space Sciences Laboratory, University of California Berkeley, Berkeley, California, USA

## Key Points:

- The dual AeroCube-6 CubeSats simultaneously observed  $> 35$  keV microbursts at a variety of spatial separations ranging from 2 to  $\approx 100$  km.
- In low Earth orbit the majority of microbursts have a size on the order of a few tens of km.
- At the magnetic equator the size of most microbursts correlate to the size of correlated whistler-mode chorus waves.

---

Corresponding author: M. Shumko, [msshumko@gmail.com](mailto:msshumko@gmail.com)

## Abstract

Microbursts are an impulsive increase of electrons from the radiation belts into the atmosphere and have been directly observed in low Earth orbit and the upper atmosphere. Microbursts are believed to be generated by scattering of radiation belt electrons by whistler mode waves. Prior work has estimated that microbursts are capable of rapidly depleting the radiation belt electrons on the order of a day, hence their role to radiation belt electron losses must be considered. Radiation belt electron losses due to microbursts are not well constrained, and more work is necessary to accurately quantify their contribution as a loss process. To address this question we present a statistical study of  $> 30$  keV microburst sizes using the pair of AeroCube-6 CubeSats. The microburst size distribution in low Earth orbit and the magnetic equator was derived. In low Earth orbit, the majority of microbursts were observed while the AC6 separation was less than a few tens of km. Most observed microburst map to sizes less than 200 km at the magnetic equator which corresponds to highly correlated whistler mode chorus scale size derived in prior literature.

## 1 Plain Language Summary

<https://sharingscience.agu.org/creating-plain-language-summary/>

## 2 Introduction

Since the discovery of the Van Allen radiation belts in the 1960s by Van Allen (1959) and Vernov and Chudakov (1960), decades of research has made headway in understanding the various particle acceleration and loss mechanisms. One of the extensively studied mechanisms responsible for both acceleration and loss is wave-particle scattering between whistler-mode chorus waves and electrons (Abel & Thorne, 1998; Meredith et al., 2002; Horne & Thorne, 2003; Thorne et al., 2005; Millan & Thorne, 2007; Bortnik et al., 2008). Whistler-mode chorus waves are typically generated by a temperature anisotropy of low energy electrons up to tens of kiloelectronvolts (keV) and are typically found in the  $\sim 0 - 12$  magnetic local times (MLT) (Li, Thorne, Angelopoulos, Bortnik, et al., 2009; Li, Thorne, Angelopoulos, Bonnell, et al., 2009). Whistler-mode chorus waves interact with radiation belt electrons, and are widely believed to cause electron precipitation termed microbursts (e.g., Millan & Thorne, 2007).

Microbursts are a subsecond impulse of electrons that are observed by high altitude balloons and satellites in low Earth orbit (LEO) on the radiation belt magnetic footprints,  $\sim 4-8$  L-shell (L) (e.g., Anderson & Milton, 1964; Parks, 1967; Lorentzen, Blake, et al., 2001; Lorentzen, Looper, & Blake, 2001; O'Brien et al., 2003; Woodger et al., 2015; Crew et al., 2016; Breneman et al., 2017; Mozer et al., 2018; Greeley et al., 2019). Microburst's role as a radiation belt electron loss mechanism has been estimated to be significant, with total radiation belt electron depletion due to microbursts estimated to be on the order of a day (Lorentzen, Looper, & Blake, 2001; O'Brien et al., 2004; Thorne et al., 2005; Breneman et al., 2017). These microburst average loss estimates are not well constrained due to assumptions made regarding the microburst precipitation region. **Think about the analogy to raindrops. The temporal and spatial duty cycles are convoluted for a spacecraft... reread Paul's and Kiersten's papers.**

One of the unconstrained microburst parameters that is critical to better quantify the role of microbursts as an instantaneous loss mechanism (the number of electrons lost per microburst) is their physical size. Historically there have been various case studies that estimated microburst size. Parks (1967) found that the size of mostly low energy microbursts to be  $40 \pm 14$  km. J. Blake et al. (1996) found a microburst with a size of a few tens of km using the the Solar Anomalous and Magnetospheric Particle Explorer (SAMPEX) and concluded that typically microbursts are less than a few tens of elec-

tron gyroradii in size (order of a few km in LEO). Dietrich et al. (2010) also used SAMPEX in another case study and concluded that the observed microbursts were smaller than 4 km. More recently, Crew et al. (2016) used the Focused Investigation of Relativistic Electron Bursts: Intensity, Range, and Dynamics (FIREBIRD-II) CubeSats and found an example of a microburst larger than 11 km, and Shumko et al. (2018) also used FIREBIRD-II to identify a microburst with a size greater than  $51 \pm 1$  km. The large variance in prior results imply that there is a distribution of microburst scale sizes which this study aims to estimate.

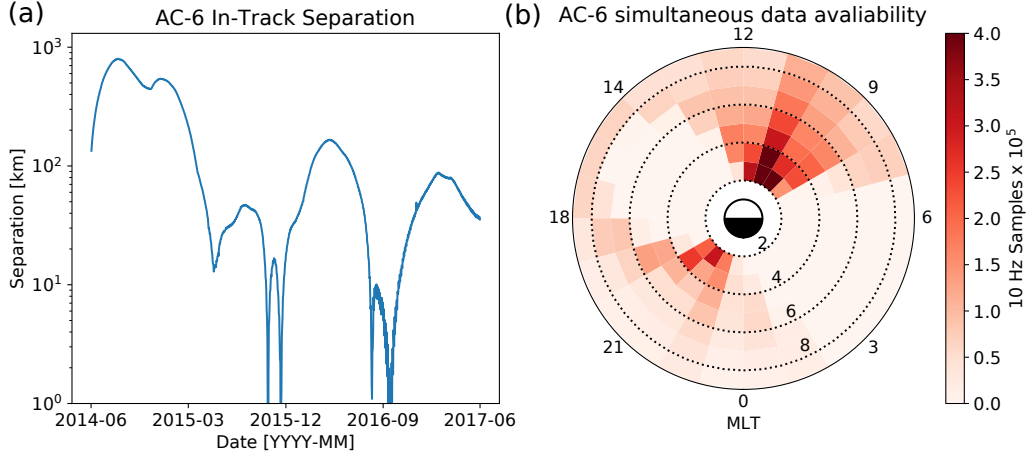
Besides addressing the instantaneous radiation belt electron losses due to individual microbursts, the microburst size distribution is useful to identify the wave mode(s) responsible for scattering microbursts. By mapping the microburst size distribution in LEO to the magnetic equator it can be compared to the wave scales derived in prior literature. This comparison can be used to identify the waves and their properties (e.g. amplitude or coherence) responsible for scattering microburst electrons.

This study addresses these two questions by estimating the microburst size distribution in LEO and the magnetic equator. The twin AeroCube-6 (AC6) CubeSats are utilized for this study because they were ideally equipped to observe microbursts simultaneously over a span of three years while their total separation varied between 2 and 800 km, mostly in latitude (in-track in orbit). This paper first describes the AC-6 mission, including their orbit and instrumentation in section 3. Section 4.1 develops the methodology used to identify microbursts observed by each spacecraft and how they were combined to make a list of the temporally coincident microbursts is described. Section 4.2 describes the methodology to estimate the microburst size distributions in LEO and the magnetic equator as a function of spacecraft separation. Then a model is developed to relate a hypothesized microburst shape, size distribution, and random microburst locations in the vicinity of AC6 to how these effects will be observed by AC6, a two-point measurement platform. Lastly, in section 6 we discuss these results and compare the microburst sizes estimated here to the size distribution of the whistler-mode chorus waves that are believed to cause microbursts.

### 3 Instrumentation

The AC6 mission consists of a pair of 0.5U (10x10x5 cm) CubeSats built by The Aerospace Corporation and launched on June 19th, 2014 into a 620 x 700 km, 98° inclination orbit. The two satellites, designated as AC6-A and AC6-B, separated after launch and drifted apart. AC6 has an active attitude control system which allows them to adjust the atmospheric drag experienced by each AC6 unit by orienting their “wings” with respect to the ram direction. By changing their orientation, AC6 was able to achieve fine separation control and maintain a separation between 2-800 km. Figure 1a shows the AC6 separation for the duration of the mission. Figure 1b shows the number of quality (data flag of 0) 10 Hz samples taken simultaneously by AC6 as a function of L and MLT. Furthermore Fig. 1b shows that the AC6 orbit was roughly dawn-dusk, sun-synchronous and precessed only a few hours in MLT over a three year period.

Each AC6 unit is equipped with a three Aerospace microdosimeters (licensed to Teledyne Microelectronics, Inc). The dosimeter used for this study is dos1 and is identical on both AC6 units. Dos1 has a 35 keV electron threshold and samples at 10 Hz. More detailed technical information on AC6 is described in O’Brien et al. (2016).



**Figure 1.** AC6 mission distributions for (a) spacecraft separation and (b) number of simultaneous quality 10 Hz samples as a function of L and MLT.

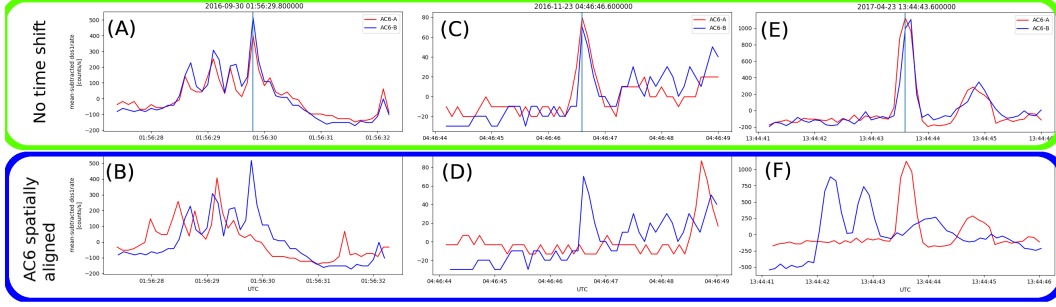
## 4 Methodology

### 4.1 Microburst Detection

The first step to find microbursts observed simultaneously by AC6 is to identify them on each spacecraft separately. Microbursts were detected with two different methods that yielded quantitatively similar results. The first method is the burst parameter (O'Brien et al., 2003). This algorithm has been successfully used in other microburst studies, mainly with the microbursts observed by SAMPEX [add citations](#). For AC6, a burst parameter threshold of 5 has good trade-off between false positive and false negative microburst detections. Another microburst detection algorithm based on wavelet spectra frequency filtering was developed and the resulting list of microbursts is similar to the list from the burst parameter.

Data cleaning to remove microburst-like transmitter noise was essential to this study. The transmitters on AC6 can cause unphysical count impulses in the dosimeters that resembles periodic trains of microbursts. One source of transmitter noise was observed at times when AC6 was in contact with the ground stations above the US for data downloads and commanding, thus the microburst detections made above the US that were mostly at low L were discarded.

Another source of noise is crosslink transmissions between AC6-A and AC6-B. These transmissions occurred when either spacecraft transitioned from the survey mode to 10 Hz mode. This noise is sometimes not caught by the data quality flag, so the following empirically-derived criteria were developed to remove those detections. The dosimeter with a 250 keV nominal electron threshold, dos2, was used because it had a similar response to noise while rarely responded to microbursts. Since the transmitter noise is very periodic with a  $\approx 0.2$  s period, cross-correlation (CC) and autocorrelation (AC) methods were applied to the dos1 and dos2 time series. Detections were removed if the following two criteria were met: either dos1 or dos2 time series had a AC peak at a 0.2 or 0.4 s lag and the dos1-dos2 CC was greater than 0.9. The AC lag criteria alone sometimes falsely removed legitimate trains of microbursts, so the second criteria insured that the detection was removed if there was an unphysically high correlation across an order of magnitude in energy.



**Figure 2.** Examples of microbursts observed simultaneously by AC6. Panels (a), (c) and (e) shows the temporally-aligned time series at spacecraft separations of 5.6 km, 16.5 km, and 68.5 km, respectively. Panels (b), (d), and (f) show the spatially aligned time series corresponding to the time series in the same column. The clear temporal correlation and lack of spatial correlation demonstrates that these events are microbursts. I know not everyone has a microscope on their computer, and I plan to clean up this plot.

The lists of microbursts observed individually by AC6 were then merged into a list of temporally correlated microbursts, i.e. microbursts that were observed simultaneously by both AC6 units, with the following procedure. The general idea is that a microburst detected by one spacecraft will cross-correlate well with the time series from the other spacecraft if it observed a similar microburst, and poorly if there was no microburst observed by the other spacecraft. Each microburst detection made by either spacecraft was cross-correlated with the time series from the other spacecraft. Windows with 1 and 1.2 s widths were used to CC the time series. Slightly different window sizes were used to account for random count variation due to Poisson noise. Microbursts detections that had a cross-correlation greater than 0.8 were considered temporally coincident. This CC threshold was chosen as it is low enough to flag user-identified temporally coincident microbursts superposed with noise, and high enough to reject most non-coincident events. Figure 2, panels (a), (c), and (e) show examples of microbursts observed by both AC6 units when they were separated by 6, 17, and 69 km, respectively.

The last CC criteria requires that the temporal CC must be greater than the spatial CC + 0.3. The spatial CC was calculated by shifting the AC6-B time series by the in-track lag to CC in the same spatial location, i.e. latitude. This criteria was applied to remove curtains, stationary structures observed by AC6 that are narrow in latitude (J. B. Blake & O'Brien, 2016) that can be confused with microbursts. Figure 2, panels (b), (d), and (f) show the shifted time series to confirm that there were no spatially correlated, non-microburst structures present. Lastly the merged microburst list was spot checked by two authors to remove poorly correlated events. After filtering out transmitter noise and applying the CC criteria 662 simultaneous microburst detections were found and used for the analysis in the following section

## 4.2 Microburst Size Distribution in LEO and Magnetic Equator

The temporally coincident microbursts, which from now on will be referred to as microbursts, are now used to estimate the fraction of microbursts observed above AC6 separation,  $s$ . When AC6 observes a microburst at  $s$ , the microburst's size must be greater than  $s$ . This fact, along with the arguments presented in Section 4 in Joy et al. (2002) who studied the most probable Jovian magnetopause and bow shock stand off distances, are used to investigate the dependence of the number of microbursts observed above  $s$ ,

as a function of  $s$ . This dependence is the microburst complementary cumulative distribution  $\bar{F}(s)$ .

The cumulative number of microbursts observed above  $s$  is the ratio of  $N(s)$ , the normalized number of microbursts observed above  $s$ , to  $N(0)$ , the total number of microbursts observed

$$\bar{F}(s) = \frac{N(s)}{N(0)}. \quad (1)$$

The  $N(s)$  normalization attempts to account for AC6's uneven sampling as a function of  $s$  and is defined by

$$N(s) = \sum_{i=s}^{\infty} n_i \frac{S_{max}}{S_i} \quad (2)$$

where  $n_i$  is the number of microbursts observed by AC6 in  $i$ th separation bin. The normalization term  $S_{max}/S_i$  is a ratio of the number of samples observed in the most sampled separation bin to the number of samples in the  $i$ th bin. This normalization factor corrects AC6's non-uniform sampling in separation by using the number of samples shown in Fig. 3c. With this normalization,  $\bar{F}(s)$  can be interpreted as the fraction of microbursts observed above  $s$  assuming AC6 sampled evenly in separation. Microburst  $\bar{F}(s)$  in LEO is shown by the black curve in Fig. 3a for  $4 < L < 8$  and split into one  $L$ -wide bins with the colored curves. The separation bin width used in Fig. 3 is 5 km. To check for bias in  $\bar{F}(s)$  due to the separation bins,  $\bar{F}(s)$  was resampled using other bin widths and offsets. Bin widths as large as 20–30 km and bin offsets did not qualitatively effect the curves in Fig. 3a.

The overall trend in Fig. 3a consists of a sudden cumulative probability drop off, followed by a shoulder up to  $s \approx 70$  km where  $\bar{F}(s)$  drops to nearly zero. A large negative gradient of  $\bar{F}(s)$  at a AC6 separation implies that microbursts must be smaller than that separation. To identify these separations, Fig. 3b shows the microburst probability density function (PDF), calculated by differentiating  $\bar{F}(s)$ . The microburst PDF shows a peak at  $s < 30$  km as well as a peak between 70–80 km separation. These PDF peaks are evidence of a sub 30 km microburst population and larger microbursts observed up 70–80 km separations. The shaded region around the black curves in Fig. 3a-b shows the standard error due to counting statistics. The uncertainty due to false coincidence events i.e. two unrelated microbursts lining up in time by random chance was also considered. The microburst duty cycle in a one minute window ( $\approx 1 L$ ) around each microburst was calculated. The false coincidence probability is the square of the duty cycle and was found to be less than 5% for the majority of microbursts. The false coincidence probability for each microburst was then used to randomly remove microbursts and  $\bar{F}(s)$  was recalculated in  $10^4$  trials. The spread in the  $\bar{F}(s)$  trials with microbursts randomly removed was much smaller than the uncertainty due to counting statistics alone.

To compare the microburst size to the size of their progenitor waves, the spacecraft locations during observed microbursts were mapped to the magnetic equator using the Olson-Pfitzer magnetic field model (Olson & Pfitzer, 1982) which is implemented with a Python wrapper for IRBEM-Lib (Boscher et al., 2012). As previously stated, a microburst observed in LEO has a size larger than the spacecraft separation, hence that microburst would also have a size larger than the spacecraft separation after it was mapped to the magnetic equator. Thus the procedure to estimate  $\bar{F}(s)$  is identical to the LEO size distribution but with a different normalization. The normalization factors were calculated by mapping every quality AC6 sample to the magnetic equator and binning them by equatorial separation into 100 km wide bins. Figure 4 shows the equatorial microburst size distribution in the same format as Fig. 3. The equatorial PDF trend is similar to LEO and most of the microbursts were observed when the AC6 equatorial separation was less than 300 km.

The results in Figs. 3 and 4 show the fraction of microbursts observed above a spacecraft separation and do not fully represent the microbursts size distribution due to the compounding effects from the range of microburst sizes and random locations of microbursts with respect to AC6 i.e. even if the microburst size is much larger than  $s$ , some fraction of those microbursts will graze only one AC6 unit and not be observed by the other. Thus modeling is necessary to capture the compounding influence of these statistical effects on AC6.

## 5 Modeling the Distribution of Microburst Sizes

### 5.1 Monte Carlo and Analytic Models to Estimate $\bar{F}(s)$

To account for the effects due to microbursts randomly occurring in the vicinity of AC6 with an unknown distribution of microburst sizes, Monte Carlo (MC) and analytic models were developed. These models assume a hypothesized distribution of microburst sizes expressed with a probability density function  $p(r)$ , and a microburst footprint shape to estimate  $\bar{F}(s)$ . The microburst footprint is assumed to be circular with a radius  $r$ . Various microburst size distribution hypotheses were considered: a one-size microburst population, a two-size microburst population, and continuous  $p(r)$  such as Maxwell, Weibull, and log-Normal.

The Monte Carlo model is the most intuitive and it consists of randomly scattering  $10^5$  microburst centers in a  $400 \times 400$  km grid around AC6 with random  $r$  distributed according to a hypothesized  $p(r)$ . Spacecraft A is placed at the origin, and spacecraft B is placed along the positive y-axis at distances from spacecraft A corresponding to the AC6 separation bins used in Section 4.2. Then for each spacecraft B location, the number of circular microbursts that encompass both spacecraft was counted. The modeled fraction of microbursts observed above  $s$  is then

$$\bar{F}(s) = \frac{\sum_{i>s}^{\infty} n_i}{\sum_{i>0}^{\infty} n_i}. \quad (3)$$

where as before the number of microbursts observed by both spacecraft in the  $i$ th bin is  $n_i$ .

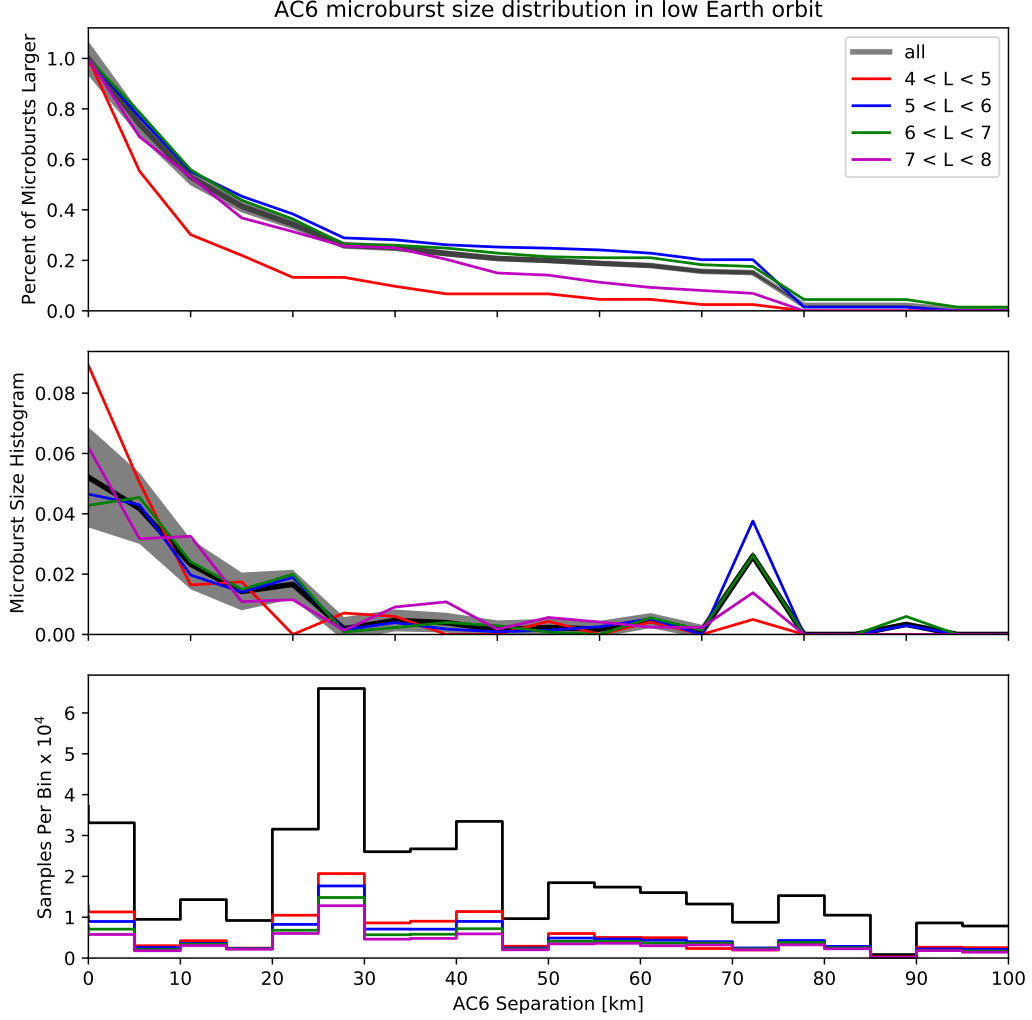
The analytic model, while identical to the MC model, highlights the geometrical concepts connecting  $p(r)$  and  $\bar{F}(s)$ . For a microburst with a diameter  $d \geq s$  there is an area between AC6 where that microburst will be observed by both spacecraft. Figure 5a-c shows the geometry with the two spacecraft indicated with black dots under varying relations between  $r$  and  $s$ . All microbursts whose center lies inside the circular area of radius  $r$  that surrounds the top spacecraft will be observed by the top spacecraft. Likewise the same argument applies to the bottom spacecraft. If it exists, the intersection of these two circular areas defines another area,  $A(r, s)$  where a microburst will be observed by both spacecraft and is given by the circle-circle intersection area equation,

$$A(r, s) = 2r^2 \cos^{-1} \left( \frac{s}{2r} \right) - \frac{s}{2} \sqrt{4r^2 - s^2}. \quad (4)$$

Example geometries where  $A(r, s) > 0$  are shown in Fig. 5b and c.

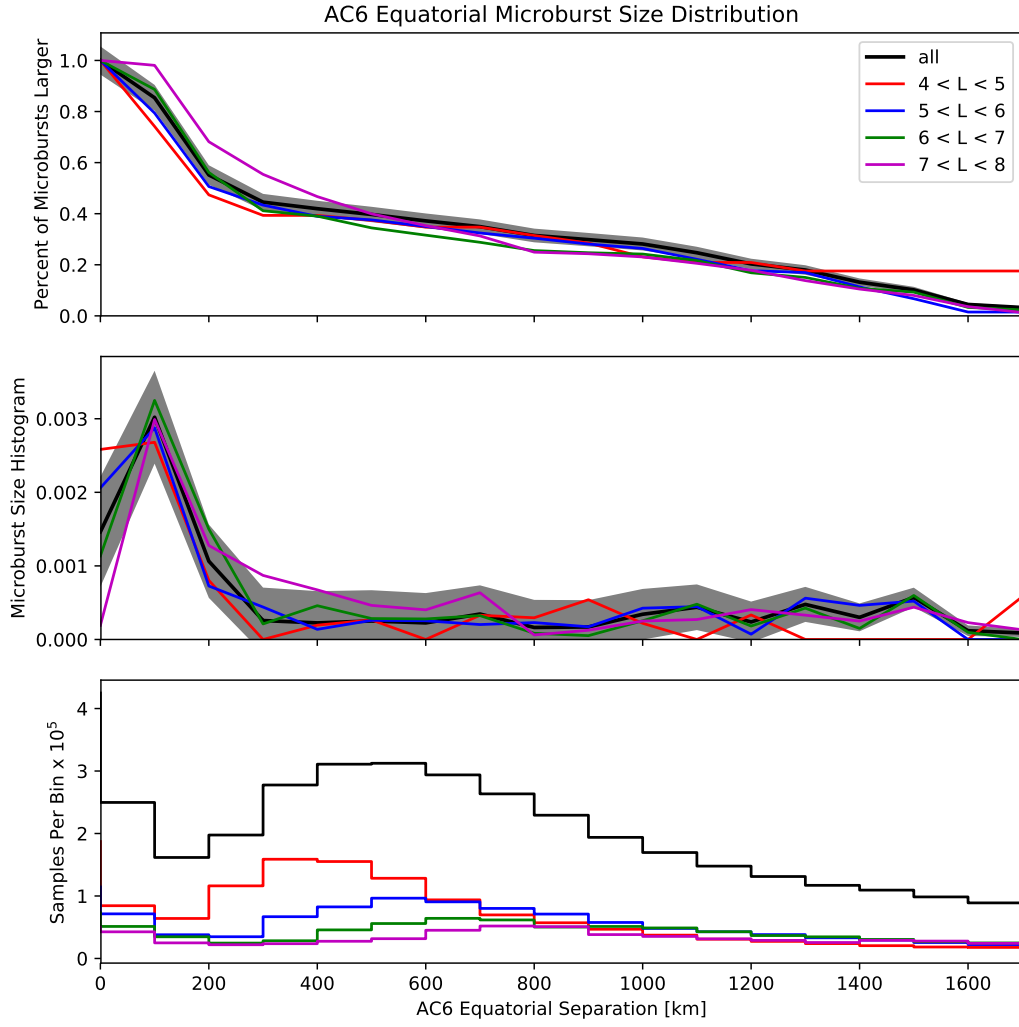
**Quarantine this section** Now we develop the analytic  $\bar{F}(s)$  given  $p(r)$ . For now we assume a one-sized  $p(r)$  i.e., a delta function in  $r$ . Then the ratio of the circle intersection areas at two distinct AC6 separations  $s_1$  and  $s_2$  is the ratio of the number of microbursts observed at those two separations





**Figure 3.** Microburst size distribution in low Earth orbit. Panel (a) shows the percent of microbursts observed above that separation after accounting for the uneven separation sampling. Panel (b) shows the microburst probability density (size histogram) as a function of separation. Lastly, panel (c) shows the number of simultaneous samples AC6 observed as a function of separation. The colored lines show the distributions binned by  $L$ , and the thick black curve for the entire radiation belt ( $4 < L < 8$ ). The gray shading around the black curve shows the uncertainty due to counting statistics.





**Figure 4.** Microburst size distribution mapped to the magnetic equator in the same format as Fig. 3.

$$\frac{n_1}{n_2} = \frac{A(r, s_1)}{A(r, s_2)}. \quad (5)$$

Similar to the MC model **Move derivation to SI**

Then the analytic  $\bar{F}(s)$  represents how quickly the fraction of microbursts observed by AC6 decreases as a function of  $s$ .

Furthermore the hypothesized  $p(r)$  is the relative occurrence of microbursts at various  $r$  and is effectively a weighting factor on  $A(r, s)$  which must be integrated out (marginalized) since AC6 is observing the cumulative effect of microbursts of all  $r$ . Lastly, a cumulative integral over a dummy variable  $s'$  is applied to the normalized areas to calculate  $\bar{F}(s)$ . With these considerations the analytic  $\bar{F}(s)$  is given by

$$\bar{F}(s) = \frac{\int_0^\infty \int_0^\infty A(r, s') p(r) dr ds'}{\int_0^\infty \int_0^\infty A(r, s') p(r) dr ds'} \quad (6)$$

where as in the Monte Carlo model the denominator normalizes  $\bar{F}(s)$  to unity. **End of quarantine**

To illustrate the effects of random microburst locations near AC6, an example analytic and Monte Carlo model of  $\bar{F}(s)$  is shown in Fig. 5d for a fixed-size  $r = 20$  km microburst population.

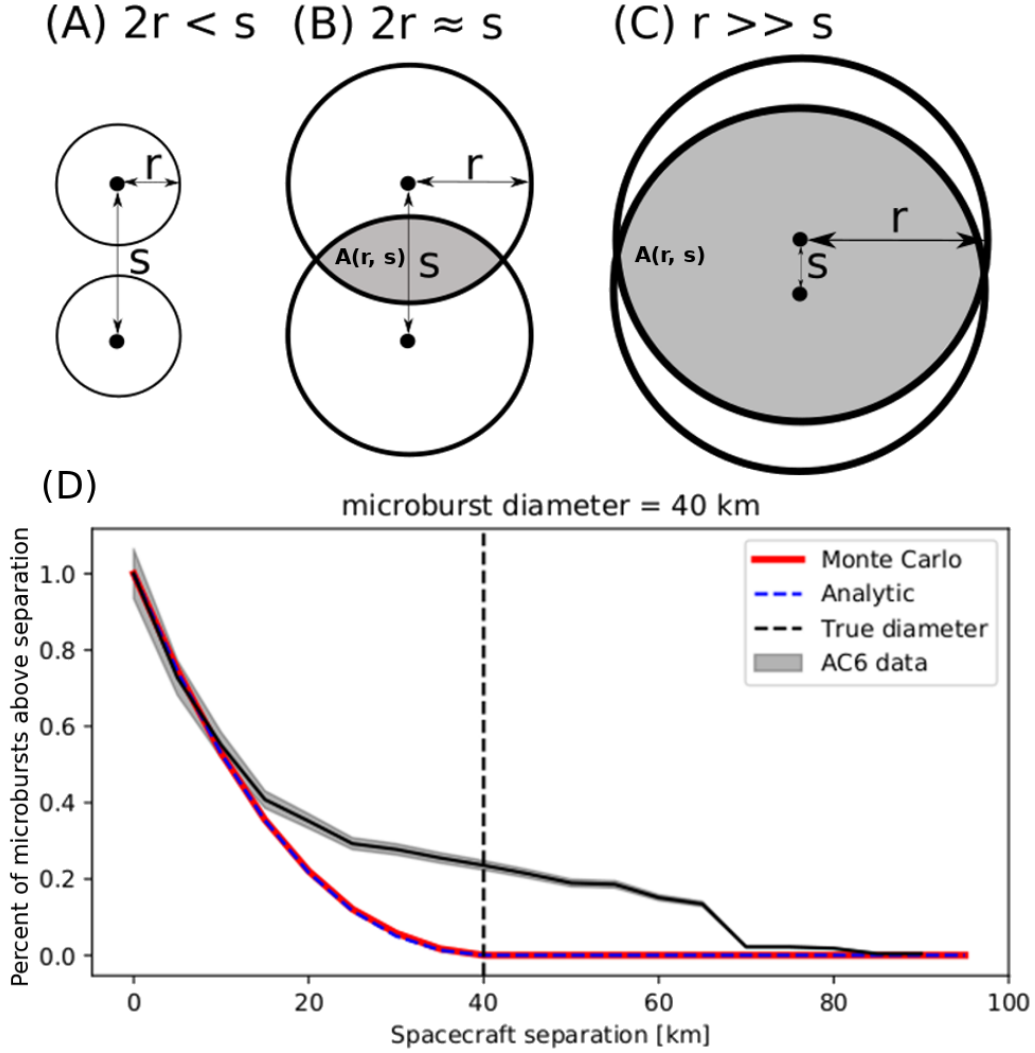
## 5.2 Estimating optimal parameters for microburst size models

**Work on this.** At this stage a few hypothesized  $p(r)$  are tested and optimal model parameters are estimated. For each hypothesized  $p(r)$ , Bayesian inference is used to estimate the optimal model parameters with a Metropolis Markov Chain Monte Carlo (MCMC) sampler (Metropolis et al., 1953). Bayes theorem is central to Bayesian inference. Given the observed  $\bar{F}(s)$  designated here with  $x$ , and  $p(r)$  model parameters with  $\theta$ , Bayes theorem can be written as

$$p(\theta|x) = \frac{p(x|\theta)p(\theta)}{p(x)} \quad (7)$$

where  $p(\theta)$  is the prior PDF for each model parameter that describes the prior level of knowledge, however weak, about the model parameters. The likelihood is  $p(x|\theta)$  that describes the probability of obtaining  $x$  given the tested model parameters. The likelihood facilitates the MC (or analytic) microburst model to map from  $\theta$  to  $\bar{F}(s)$ . The posterior distribution,  $p(\theta|x)$  is the updated  $\theta$  distribution given the observations and prior. The posterior is used here to make conclusions regarding the range of model parameters that are consistent with the observations after assuming a likelihood function and prior knowledge of the model parameters. Lastly,  $p(x)$  is the marginal likelihood that describes the probability of obtaining  $x$  after marginalizing over all prior variables. Calculation of  $p(x)$  is difficult, and often not necessary for model parameter estimation.

With all of the above terminology, the important takeaway is that the posterior distribution for each model parameter is interpreted as the range of  $\theta$  consistent with  $x$ . A 95% credible interval (CI) is reported here that is interpreted as: assuming a hypothesized  $p(\theta, r)$ , there is a 95% probability that the true  $\theta$  is inside the CI. To sample the posterior distribution, the  $\theta$  parameter space is explored with the Metropolis sampler.



**Figure 5.** Panels A-C show the varying geometries of the analytic model. The two spacecraft are shown as black dots. The enclosing black circle around each spacecraft bounds the area where a microburst will be observed by one or both AC6 units if the microburst's center lies somewhere inside the circle. Panel (A) shows the case where microburst diameter is smaller than the AC6 separation and all microbursts will be observed by either unit A or B and never simultaneously. Panel (B) shows the intermediate case where the microburst diameter is comparable to the AC6 separation and some fraction of microbursts will be observed simultaneously. The fraction of the microbursts simultaneously observed is proportional to the circle intersection area shown with grey shading. Panel (C) shows the case where the microburst diameter is much larger than the spacecraft separation and nearly all microbursts will be observed by both spacecraft. Lastly panel (D) shows  $\bar{F}(s)$  from the AC6 data with a solid black line, and a MC and analytic  $\bar{F}(s)$  curves for a single-sized microburst distribution with a 40 km diameter.

While the Metropolis sampler is explained in detail in Metropolis et al. (1953) and a good introduction given in Sambridge et al. (2006), a brief overview is warranted. The Metropolis sampler samples the posterior distribution in  $N$  trials. For each trial a random  $\theta$  is picked from the prior distribution and the posterior probability is estimated. Then the Metropolis MCMC picks a new  $\theta$  nearby in parameter space. If the posterior probability of the new  $\theta$  is higher than the current  $\theta$ , the MCMC accepts the new  $\theta$ , otherwise there is a random chance that the new  $\theta$  will be accepted or rejected. This accept/reject criteria allows the sampler to trend to more probable  $\theta$  while also exploring the neighboring regions. After the  $N$  trials, a histogram is made using the accepted  $\theta$  to produce the posterior distribution for each model parameter.

A one-size microburst population model is explored first and the microburst PDF is given by

$$p(d) = \delta(s - d) \quad (8)$$

where  $\delta$  is the Dirac Delta function and  $d$  is the diameter of all microbursts. The range of  $d$  that is consistent with the observed  $\bar{F}(s)$  is shown in Fig. 6. Assuming this model, there is a 95% probability that the microburst size is between 38 and 129 km with the optimal size of 73 km, estimated from minimizing the least squares penalty function.

A generalization of the one-size model is a two-size microburst population model that assumes the following microburst PDF

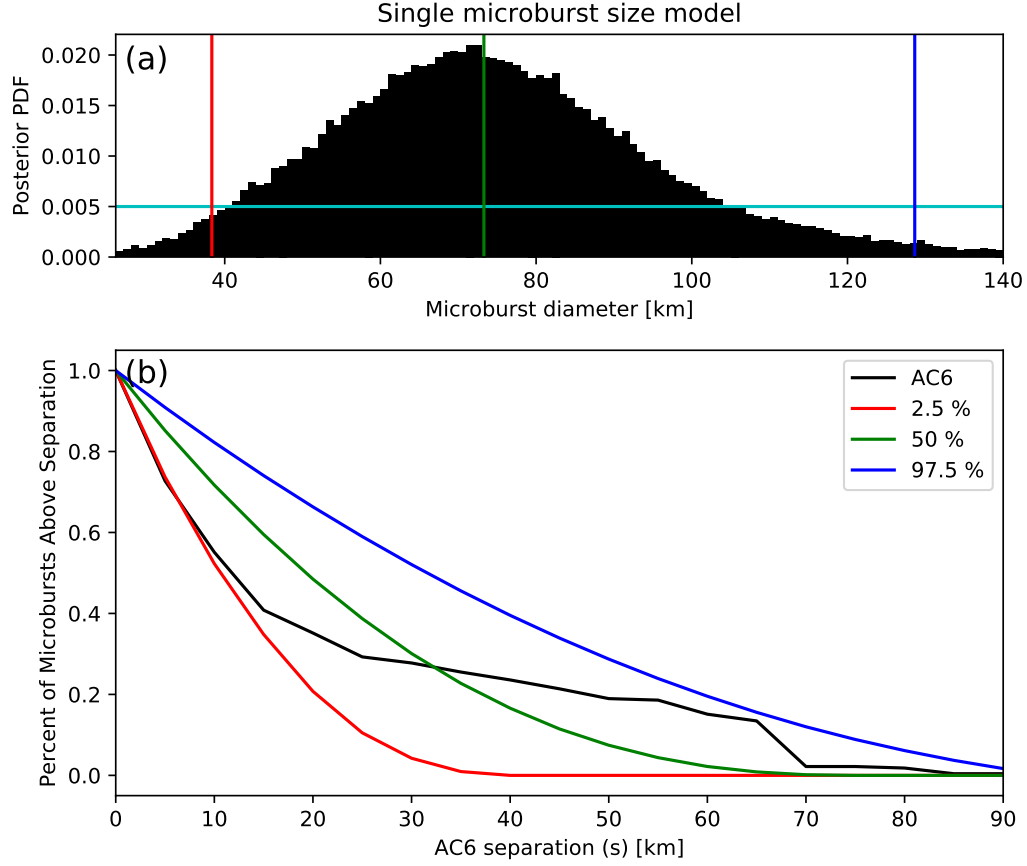
$$p(d) = a\delta(s - d_0) + (1 - a)\delta(s - d_1) \quad (9)$$

where the diameters of the two microburst populations is given by  $d_0$  and  $d_1$  and  $a$  is the parameter that quantifies the relative fractions of the two populations. The result of this model is shown in Fig. 7. The fit is slightly better than the one-size model, although that is to be expected assuming two more free model parameters. A majority, 98 %, of microbursts, have a size between 12 and 47 km with a rare population with a size between 76 and 234 km. The set of parameters that minimize least squares is 99.5 % of microbursts are small with a size of 21 km and the remaining 0.5 % of microbursts have a 140 km size.

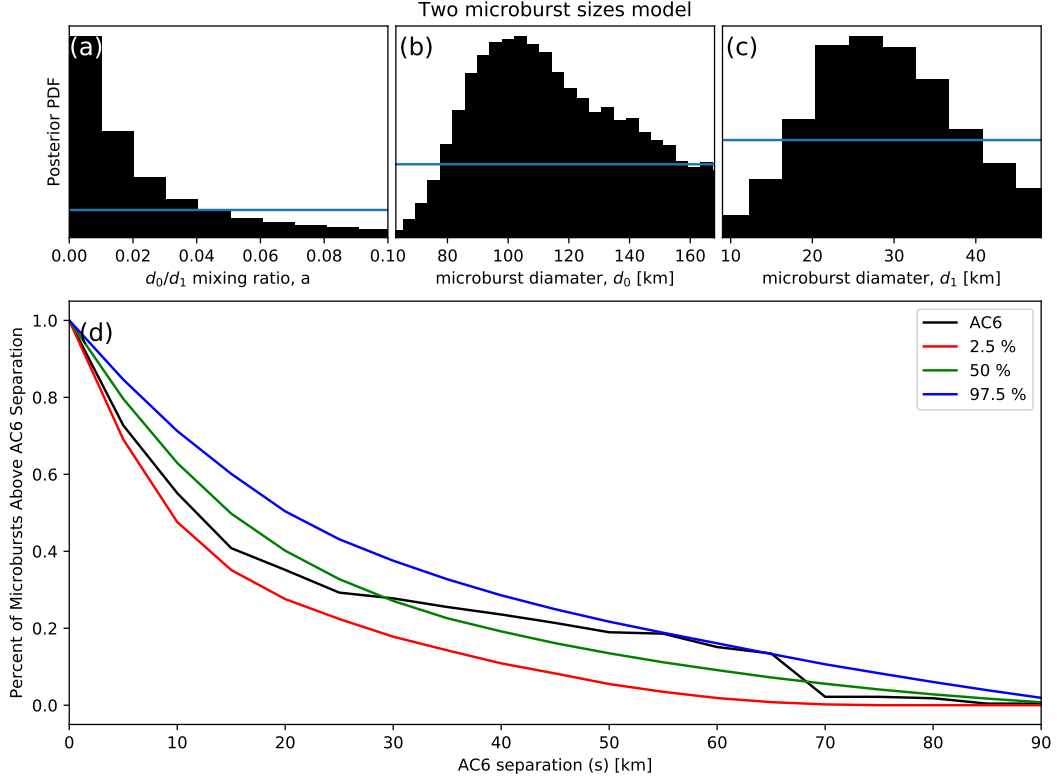
## 6 Discussion

The LEO microburst  $\bar{F}(s)$  estimated in section 4.2 shows that a majority of coincident microbursts were observed by AC6 when they were separated by less than a few tens of km. This conclusion is consistent with prior literature and most similar to Parks (1967) who reported that  $> 15$  keV microbursts are  $40 \pm 14$  km and the bouncing packet example shown in J. Blake et al. (1996) with a size of “at least a few tens of kilometers”. The AC6 microburst size distribution is much larger than the sizes reported in Dietrich et al. (2010) who used very low (VLF) frequency transmission paths and SAMPEX to conclude that microbursts must be smaller than 4 km from a small number of microbursts observed during one SAMPEX pass. Dietrich et al. (2010) arrived at their conclusion by looking for temporal coincidence of microbursts and FAST events, subsecond VLF transmission perturbations, but the connection between FAST events and microbursts is not well understood. Lastly, our new results are consistent with FIREBIRD-II observations of a  $> 11$  km microburst reported by Crew et al. (2016). The small number of microbursts observed by AC6 up to  $s \approx 70$  km are consistent with the  $> 51$  km bouncing packet microburst reported in Shumko et al. (2018).

The microburst PDF shown in Fig. 3b suggests that the microburst size distribution is bimodal. The quality of the AC6 data is insufficient to definitively conclude that there are two distinct microburst populations. This has been suggested before by J. Blake et al. (1996) who noted that the  $> 150$  keV and  $> 1$  MeV microbursts are not always well correlated e.g. Fig. 10 in J. Blake et al. (1996). The different microburst popula-



**Figure 6.** Range of plausible microburst sizes assuming all microbursts are one fixed size. Panel (a) shows the posterior probability density function of microburst diameters in black. The red, green, and blue vertical lines at 38, 73, and 129 km represent the 2.5, 50, and 97.5 posterior percentiles, respectively. A uniform prior between 0 and 200 km was assumed for this MCMC run and is shown in cyan. Panel (b) shows the percent of microbursts observed above an AC6 separation for  $4 < L < 8$  in black. The 2.5, 50 and 97.5 size percentiles were estimated from the posterior and plotted in red, green, and blue curves, respectively.



**Figure 7.** Plausible microburst percent curves assuming microburst size distribution is bi-modal consisting of two sizes  $d_0$  and  $d_1$  with a mixing term that quantifies the relative occurrence of the  $d_0$  to  $d_1$  microburst populations. Panel (a) shows the posterior distribution for the microburst population mixing term,  $a$  with a median value of 0.02. The  $a$  prior was uniform between 0 and 0.2. Panel (b) shows the posterior distribution for  $d_0$ , the larger microburst population estimated with a uniform prior between 50 and 200 km and the posterior median diameter of 122 km. Panel (c) shows the posterior distribution for  $d_1$ , the smaller microburst population, estimated using a uniform prior between 0 and 50 km with a median diameter of 28 km. Panel (d) is similar to Fig. 6b and shows the AC6 microburst fraction for  $4 < L < 8$  in black. A set of 1000 random parameter triples ( $a$ ,  $d_0$ , and  $d_1$ ) were drawn from the posterior and used to generate a family of  $\bar{F}(s)$  curves. At each  $s$  the range of consistent  $\bar{F}(s)$  were quantified by the 2.5, 50 and 97.5 percentiles and shown with the red, green, and blue curves, respectively.

tion hypothesis can be better tested with an AC6-like mission with better energy resolution and homogeneous MLT coverage.

The model results from section 5 emphasize that care must be taken to compare the  $\bar{F}(s)$  curves observed by AC6 and the true microburst size distribution due to the compounding effect of an unknown microburst size distribution, unknown microburst shape, and random microburst locations near AC6. By naively assuming there is only one microburst size, the results in Fig. 6 suggest that there is a 95% probability that the microburst size is somewhere between 38 and 129 km, a relatively wide range of values. The two-size model microburst from Fig. 7 has a smaller variance around the AC6  $\bar{F}(s)$ , which is expected with the addition of two more free parameters. The two size model is interpreted as 98% of microbursts sizes are between 12 and 47 km and larger microbursts are very uncommon. A variety of continuous  $p(r)$  such as the Maxwellian, Weibull and log-normal were also tested with similar fits and are not shown. While the two-size model quantitatively fits the observations the best out of all  $p(r)$  tested, nature surely does not only have two microburst sizes. Rather the most realistic PDF hypothesis would be bimodal and continuous. Due to lack of prior observations and theoretical predictions, it is difficult to identify a more appropriate  $p(r)$  hypothesis at this time.

The equatorial microburst  $\bar{F}(s)$  estimated in section 4.2 can be compared to prior multi-point measurements of chorus source sizes made near the magnetic equator. The International Sun-Earth Explorers (ISEE 1 and 2) were used by Gurnett et al. (1979) to make one of the first direct chorus source scale measurements. Gurnett et al. (1979) estimated that the wave power correlation scale was on the order of a few hundred km across the background magnetic field. Using the Cluster Wide Band Data measurements Santolik et al. (2003) found the correlation scale of whistler mode chorus waves to be around 100 km near the source region at  $L \approx 4$  and midnight MLT sector. Furthermore, Turner et al. (2017) used the four satellites comprising the Magnetospheric Multiscale Mission and found that rising tone whistler mode chorus elements were phase coherent up to 70 km at  $L \approx 8$ . Lastly, Agapitov et al. (2010, 2011, 2017, 2018) used multiple sets of spacecraft missions with wave measurements near the chorus source region to statistically show that the extent of chorus source region can extend from 600 km in the outer radiation belt to greater than 1,000 km in the outer magnetosphere.

More generally, the microburst probability density from Fig. 4b shows that the majority of microbursts were observed when the equatorial AC6 separation was less than 200 km. This result show that the waves responsible for scattering microburst electrons must have correlated properties on those scales. The wave properties necessary for scattering microburst electrons e.g. coherence, polarization, wave normal angle, can be identified by studying the waves properties that are only observed by multiple equatorial spacecraft at small separations. These properties can then aid wave-particle scattering model development by constraining the wave properties and scattering modes responsible for scattering microburst electrons which in turn can be used to predict the distribution of microburst sizes in LEO.

## 7 Conclusions

In conclusion, the twin AC6 CubeSats enabled the detailed statistical study of microburst sizes from a two point measurement platform. Roughly 60% of the  $> 30$  keV microbursts were simultaneously observed while AC6 was separated by less than 20 km and the rest were observed up to  $\approx 80$  km separation. The microburst cumulative distribution function is essential to model and quantify the relationship between the number of microbursts observed as a function of separation to a hypothesized microburst size distributions. The AC6 microburst data, together with modeling, has hinted at the existence of a bimodal microburst size PDF with the majority of microbursts smaller than 40 km and a rare microburst population with a size around 100 km. The bimodal size



hypothesis may be more comprehensively addressed from LEO spacecraft with more microburst observations for spacecraft separations up to 100 km, homogeneous MLT coverage, and differential energy channels. Moreover, to disentangle the compounding effect that affects two-point microburst measurements, a X-ray imager on a high altitude balloon can observe the atmospheric microburst footprint and determine the microburst size, shape, and any spatial correlations without ambiguity.

When mapped to the magnetic equator, most microbursts were observed while the mapped AC6 separation was less than 200 km. This correlates well with the sizes of highly correlated chorus waves and it suggests that the wave properties crucial for scattering microbursts must be correlated on relatively small regions. By studying the wave properties that are correlated on a few hundred km scales, the dominant wave scattering modes may be identified.

## Acknowledgments

This work was made possible with the help from the many engineers and scientists at the Aerospace Corporation who successfully designed, built, and operated AC6. M. Shumko was supported by NASA Headquarters under the NASA Earth and Space Science Fellowship Program - Grant 80NSSC18K1204. **Other Aerospace and MSU funding sources...** The AC6 data is available at <http://rbspgway.jhuapl.edu/ac6> and the IRBEM-Lib version used for this analysis can be downloaded from <https://sourceforge.net/p/irbem/code/616/tree/>.

## References

- Abel, B., & Thorne, R. M. (1998). Electron scattering loss in earth's inner magnetosphere: 1. dominant physical processes. *Journal of Geophysical Research: Space Physics*, 103(A2), 2385–2396.
- Agapitov, O., Blum, L. W., Mozer, F. S., Bonnell, J. W., & Wygant, J. (2017). Chorus whistler wave source scales as determined from multipoint van allen probe measurements. *Geophysical Research Letters*, n/a–n/a. Retrieved from <http://dx.doi.org/10.1002/2017GL072701> (2017GL072701) doi: 10.1002/2017GL072701
- Agapitov, O., Krasnoselskikh, V., Dudok de Wit, T., Khotyaintsev, Y., Pickett, J. S., Santolik, O., & Rolland, G. (2011). Multispacecraft observations of chorus emissions as a tool for the plasma density fluctuations' remote sensing. *Journal of Geophysical Research: Space Physics*, 116(A9), n/a–n/a. Retrieved from <http://dx.doi.org/10.1029/2011JA016540> (A09222) doi: 10.1029/2011JA016540
- Agapitov, O., Krasnoselskikh, V., Zaliznyak, Y., Angelopoulos, V., Le Contel, O., & Rolland, G. (2010). Chorus source region localization in the earth's outer magnetosphere using themis measurements. *Annales Geophysicae*, 28(6), 1377–1386. Retrieved from <http://www.ann-geophys.net/28/1377/2010/> doi: 10.5194/angeo-28-1377-2010
- Agapitov, O., Mourenas, D., Artemyev, A., Mozer, F., Bonnell, J., Angelopoulos, V., ... Krasnoselskikh, V. (2018). Spatial extent and temporal correlation of chorus and hiss: Statistical results from multipoint themis observations. *Journal of Geophysical Research: Space Physics*, 123(10), 8317–8330.
- Anderson, K. A., & Milton, D. W. (1964). Balloon observations of X rays in the auroral zone: 3. High time resolution studies. *Journal of Geophysical Research*, 69(21), 4457–4479. Retrieved from <http://dx.doi.org/10.1029/JZ069i021p04457> doi: 10.1029/JZ069i021p04457
- Blake, J., Looper, M., Baker, D., Nakamura, R., Klecker, B., & Hovestadt, D. (1996). New high temporal and spatial resolution measurements by sampex of the precipitation of relativistic electrons. *Advances in Space Research*, 18(8), 171 - 186. Retrieved from <http://www.sciencedirect.com/>

- science/article/pii/0273117795009698 doi: [http://dx.doi.org/10.1016/0273-1177\(95\)00969-8](http://dx.doi.org/10.1016/0273-1177(95)00969-8)
- Blake, J. B., & O'Brien, T. P. (2016). Observations of small-scale latitudinal structure in energetic electron precipitation. *Journal of Geophysical Research: Space Physics*, 121(4), 3031–3035. Retrieved from <http://dx.doi.org/10.1002/2015JA021815> (2015JA021815) doi: 10.1002/2015JA021815
- Bortnik, J., Thorne, R., & Inan, U. S. (2008). Nonlinear interaction of energetic electrons with large amplitude chorus. *Geophysical Research Letters*, 35(21).
- Boscher, D., Bourdarie, S., O'Brien, P., Guild, T., & Shumko, M. (2012). *Irbem-lib library*.
- Breneman, A., Crew, A., Sample, J., Klumpar, D., Johnson, A., Agapitov, O., ... others (2017). Observations directly linking relativistic electron microbursts to whistler mode chorus: Van allen probes and FIREBIRD II. *Geophysical Research Letters*.
- Crew, A. B., Spence, H. E., Blake, J. B., Klumpar, D. M., Larsen, B. A., O'Brien, T. P., ... Widholm, M. (2016). First multipoint in situ observations of electron microbursts: Initial results from the NSF FIREBIRD II mission. *Journal of Geophysical Research: Space Physics*, 121(6), 5272–5283. Retrieved from <http://dx.doi.org/10.1002/2016JA022485> (2016JA022485) doi: 10.1002/2016JA022485
- Dietrich, S., Rodger, C. J., Clilverd, M. A., Bortnik, J., & Raita, T. (2010). Relativistic microburst storm characteristics: Combined satellite and ground-based observations. *Journal of Geophysical Research: Space Physics*, 115(A12).
- Greeley, A., Kanekal, S., Baker, D., Klecker, B., & Schiller, Q. (2019). Quantifying the contribution of microbursts to global electron loss in the radiation belts. *Journal of Geophysical Research: Space Physics*.
- Gurnett, D., Anderson, R., Scarf, F., Fredricks, R., & Smith, E. (1979). Initial results from the isee-1 and-2 plasma wave investigation. *Space Science Reviews*, 23(1), 103–122.
- Horne, R. B., & Thorne, R. M. (2003). Relativistic electron acceleration and precipitation during resonant interactions with whistler-mode chorus. *Geophysical Research Letters*, 30(10). Retrieved from <http://dx.doi.org/10.1029/2003GL016973> (1527) doi: 10.1029/2003GL016973
- Joy, S., Kivelson, M., Walker, R., Khurana, K., Russell, C., & Ogino, T. (2002). Probabilistic models of the jovian magnetopause and bow shock locations. *Journal of Geophysical Research: Space Physics*, 107(A10), SMP–17.
- Li, W., Thorne, R., Angelopoulos, V., Bonnell, J., McFadden, J., Carlson, C., ... Auster, H. (2009). Evaluation of whistler-mode chorus intensification on the nightside during an injection event observed on the THEMIS spacecraft. *Journal of Geophysical Research: Space Physics*, 114(A1).
- Li, W., Thorne, R. M., Angelopoulos, V., Bortnik, J., Cully, C. M., Ni, B., ... Magnes, W. (2009). Global distribution of whistler-mode chorus waves observed on the THEMIS spacecraft. *Geophysical Research Letters*, 36(9). Retrieved from <http://dx.doi.org/10.1029/2009GL037595> (L09104) doi: 10.1029/2009GL037595
- Lorentzen, K. R., Blake, J. B., Inan, U. S., & Bortnik, J. (2001). Observations of relativistic electron microbursts in association with VLF chorus. *Journal of Geophysical Research: Space Physics*, 106(A4), 6017–6027. Retrieved from <http://dx.doi.org/10.1029/2000JA003018> doi: 10.1029/2000JA003018
- Lorentzen, K. R., Looper, M. D., & Blake, J. B. (2001). Relativistic electron microbursts during the GEM storms. *Geophysical Research Letters*, 28(13), 2573–2576. Retrieved from <http://dx.doi.org/10.1029/2001GL012926> doi: 10.1029/2001GL012926
- Meredith, N., Horne, R., Summers, D., Thorne, R., Iles, R., Heynderickx, D., & Anderson, R. (2002). Evidence for acceleration of outer zone electrons to rela-

- tivistic energies by whistler mode chorus. In *Annales geophysicae* (Vol. 20, pp. 967–979).
- Metropolis, N., Rosenbluth, A. W., Rosenbluth, M. N., Teller, A. H., & Teller, E. (1953). Equation of state calculations by fast computing machines. *The journal of chemical physics*, 21(6), 1087–1092.
- Millan, R., & Thorne, R. (2007). Review of radiation belt relativistic electron losses. *Journal of Atmospheric and Solar-Terrestrial Physics*, 69(3), 362–377. Retrieved from <http://www.sciencedirect.com/science/article/pii/S1364682606002768> doi: <http://dx.doi.org/10.1016/j.jastp.2006.06.019>
- Mozer, F. S., Agapitov, O. V., Blake, J. B., & Vasko, I. Y. (2018). Simultaneous observations of lower band chorus emissions at the equator and microburst precipitating electrons in the ionosphere. *Geophysical Research Letters*. Retrieved from <http://dx.doi.org/10.1002/2017GL076120> doi: 10.1002/2017GL076120
- O’Brien, T. P., Blake, J. B., & W., G. J. (2016, May). *Aerocube-6 dosimeter data readme* (Tech. Rep. No. TOR-2016-01155). The Aerospace Corporation.
- O’Brien, T. P., Looper, M. D., & Blake, J. B. (2004). Quantification of relativistic electron microburst losses during the GEM storms. *Geophysical Research Letters*, 31(4). Retrieved from <http://dx.doi.org/10.1029/2003GL018621> (L04802) doi: 10.1029/2003GL018621
- O’Brien, T. P., Lorentzen, K. R., Mann, I. R., Meredith, N. P., Blake, J. B., Fennell, J. F., ... Anderson, R. R. (2003). Energization of relativistic electrons in the presence of ULF power and MeV microbursts: Evidence for dual ULF and VLF acceleration. *Journal of Geophysical Research: Space Physics*, 108(A8). Retrieved from <http://dx.doi.org/10.1029/2002JA009784> doi: 10.1029/2002JA009784
- Olson, W. P., & Pfitzer, K. A. (1982). A dynamic model of the magnetospheric magnetic and electric fields for July 29, 1977. *Journal of Geophysical Research: Space Physics*, 87(A8), 5943–5948. Retrieved from <http://dx.doi.org/10.1029/JA087iA08p05943> doi: 10.1029/JA087iA08p05943
- Parks, G. K. (1967). Spatial characteristics of auroral-zone X-ray microbursts. *Journal of Geophysical Research*, 72(1), 215–226.
- Sambridge, M., Gallagher, K., Jackson, A., & Rickwood, P. (2006). Trans-dimensional inverse problems, model comparison and the evidence. *Geophysical Journal International*, 167(2), 528–542.
- Santolik, O., Gurnett, D., Pickett, J., Parrot, M., & Cornilleau-Wehrlin, N. (2003). Spatio-temporal structure of storm-time chorus. *Journal of Geophysical Research: Space Physics*, 108(A7).
- Shumko, M., Sample, J., Johnson, A., Blake, B., Crew, A., Spence, H., ... Handley, M. (2018). Microburst scale size derived from multiple bounces of a microburst simultaneously observed with the firebird-ii cubesats. *Geophysical Research Letters*, 45(17), 8811–8818. Retrieved from <https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2018GL078925> doi: 10.1029/2018GL078925
- Thorne, R. M., O’Brien, T. P., Shprits, Y. Y., Summers, D., & Horne, R. B. (2005). Timescale for MeV electron microburst loss during geomagnetic storms. *Journal of Geophysical Research: Space Physics*, 110(A9). Retrieved from <http://dx.doi.org/10.1029/2004JA010882> (A09202) doi: 10.1029/2004JA010882
- Turner, D., Lee, J., Claudepierre, S., Fennell, J., Blake, J., Jaynes, A., ... others (2017). Examining coherency scales, substructure, and propagation of whistler mode chorus elements with magnetospheric multiscale (mms). *Journal of Geophysical Research: Space Physics*, 122(11).
- Van Allen, J. A. (1959). The geomagnetically trapped corpuscular radiation. *Journal of Geophysical Research*, 64(11), 1683–1689. Retrieved from <http://dx.doi.org/10.1029/JZ064i011p01683> doi: 10.1029/JZ064i011p01683

- 551 Vernov, S., & Chudakov, A. (1960). Investigation of radiation in outer space. In *In-*  
552 *ternational cosmic ray conference* (Vol. 3, p. 19).
- 553 Woodger, L., Halford, A., Millan, R., McCarthy, M., Smith, D., Bowers, G., . . .  
554 Liang, X. (2015). A summary of the BARREL campaigns: Technique for  
555 studying electron precipitation. *Journal of Geophysical Research: Space*  
556 *Physics*, 120(6), 4922–4935.