

# Microburst Scale Size Derived from Multiple Bounces of a Microburst Simultaneously Observed with the FIREBIRD-II CubeSats

Mykhaylo Shumko<sup>1</sup>, John Sample<sup>1</sup>, Arlo Johnson<sup>1</sup>, Bern Blake<sup>2</sup>, Alex Crew<sup>3</sup>, Harlan  
Spence<sup>4</sup>, David Klumpar<sup>1</sup>, Oleksiy Agapitov<sup>5</sup>, Matthew Handley<sup>1</sup>

<sup>1</sup>Department of Physics, Montana State University, Bozeman, Montana, USA

<sup>2</sup>Space Science Applications Laboratory, The Aerospace Corporation, Los Angeles, California, USA

<sup>3</sup>The Johns Hopkins University Applied Physics Laboratory LLC, Laurel, Maryland, USA

<sup>4</sup>Institute for the Study of Earth, Oceans, and Space, University of New Hampshire, Durham, New Hampshire, USA

<sup>5</sup>Space Sciences Laboratory, UC Berkeley, Berkeley, California, USA

## Key Points:

- The lat/lon scale sizes at LEO were  $28.8 \pm 0.8$  km and  $50.8 \pm 11.4$  km, respectively.
- Deduced lower bound equatorial scale size is similar to the whistler-mode chorus source scale.

## Abstract

The FIREBIRD-II CubeSats simultaneously observed a large microburst with multiple bounces on February 2nd, 2015 during a small storm. This is the first such microburst observed by two spacecraft, and it sheds light on its spatial scale sizes and bounce periods. Its lower bound latitudinal scale size was  $28.8 \pm 0.8$  km and the longitudinal scale size was  $50.8 \pm 11.4$  km in low earth orbit. Using the Tsyganenko 1989 magnetic field model, these scale sizes were mapped to the magnetic equator to get the radial and azimuthal scale sizes of at least  $504 \pm 14$  km and  $530 \pm 119$  km, respectively. These lower bound equatorial scale sizes are similar to whistler-mode chorus wave source scale sizes, which supports the hypothesis that microbursts are a product of electron scattering by chorus waves. Lastly, the electron bounce period from the subsequent bounces was calculated and compared to analytical and numerical bounce times to validate numerous magnetic field models.

## 1 Introduction

The dynamics of radiation belt electrons are complex, and are driven by competition between source and loss processes. A few possible loss processes are radial diffusion [Shprits and Thorne, 2004], magnetopause shadowing [Ukhorskiy et al., 2006], and pitch angle diffusion [Selesnick et al., 2003; Abel and Thorne, 1998] due to plasma wave and Coulomb scattering. As described in [Millan and Thorne, 2007; Thorne, 2010] and references contained within, there are a variety of waves that cause pitch angle scattering, including electromagnetic ion cyclotron waves, plasmaspheric hiss, and whistler-mode chorus. Whistler-mode chorus predominantly occurs in the dawn sector [Li et al., 2009] where it accelerates electrons with large equatorial pitch angles and scatters electrons with small equatorial pitch angles [Horne and Thorne, 2003]. Some of these electrons may be impulsively scattered into the loss cone. This sudden increase (duration  $\approx 100$  ms) in precipitating electron flux is believed to be caused by scattering with whistler-mode chorus waves [e.g. Thorne et al., 2005; Breneman et al., 2017].

Microbursts were first termed by Anderson and Milton [1964] who used high altitude balloon observations of bremsstrahlung X-rays emitted from scattered electrons impacting the atmosphere. Since then, microbursts have routinely been observed with other balloon missions [Parks, 1967; Woodger et al., 2015; Anderson et al., 2017]. Relativistic microbursts have not been observed by high altitude balloons in the dawn sector [Millan et al., 2002], which may be caused by a pitch angle anisotropy of relativistic microbursts due to relatively weaker pitch angle scattering of relativistic electrons by chorus [Lee et al., 2012]. Microbursts have also been observed in low earth orbit (LEO) with, e.g. the SAMPEX  $> 1$  MeV channel [Nakamura et al., 1995, 2000; Blake et al., 1996; Lorentzen et al., 2001a,b; O'Brien et al., 2003, 2004; Blum et al., 2015] and FIREBIRD-II with its  $> 200$  keV energy channels [Crew et al., 2016; Breneman et al., 2017]. Microbursts and chorus waves predominantly occur in the dawn sector, are bursty [Lorentzen et al., 2001b], and observationally Breneman et al. [2017] directly linked microbursts scattered by whistler-mode chorus. Understanding microburst precipitation is important to radiation belt dynamics since they have been modeled and empirically estimated to deplete the relativistic electron population of the outer radiation belt on time scales of hours to a few days [O'Brien et al., 2004; Thorne et al., 2005; Shprits et al., 2007; Breneman et al., 2017].

An important parameter in the estimation of instantaneous radiation belt electron losses due to microbursts is their scale size. Parks [1967] used balloon measurements of bremsstrahlung X-rays to estimate the scale size of predominantly low energy microbursts as  $40 \pm 14$  km. In Blake et al. [1996] a microburst with multiple bounces observed by SAMPEX, was estimated to have a latitudinal scale size of “at least a few tens of kilometers”, and they concluded that typically microbursts are less than a few tens of electron gyroradii in size (at  $L = 5$  at LEO, the gyroradii of 1 MeV electrons is on the order of 100 m). Dietrich et al. [2010]

used SAMPEX along with ground-based very low frequency stations to conclude that microbursts have scale sizes less than 4 km.

Since February 1st, 2015, microbursts have been observed by FIREBIRD-II, a pair of CubeSats in LEO. Soon after launch, when the two FIREBIRD-II spacecraft were at close range, a microburst with a scale size greater than 11 km was observed [Crew *et al.*, 2016]. On the same day, a microburst with multiple bounces was simultaneously observed on both spacecraft. The microburst decay was observed over a period of a few seconds, while the spacecraft were traveling predominantly in latitude. The analysis in this paper uses FIREBIRD-II to resolve the spatial and temporal properties of the first microburst with multiple bounces observed with two spacecraft. The rest of this paper is organized as follows: in section 2, the spacecraft and the microburst observation will be introduced. In section 3, the methodology of the spacecraft time and position correction, the microburst latitudinal and longitudinal scale sizes in LEO and the magnetic equator, and electron bounce period will be explained. Lastly, in section 4, these results will be tied to the current empirical and theoretical understanding of microbursts and their connection to whistler-mode chorus scales.

## 2 Spacecraft and Observation

FIREBIRD-II is a pair of identically-instrumented 1.5U CubeSats (15 cm x 10 cm x 10 cm) that are termed Flight Unit 3 (FU3) and Flight Unit 4 (FU4) and were launched on January 31st, 2015. Their orbit has an apogee of 632 km, perigee of 433 km, and 99° inclination [Crew *et al.*, 2016]. FU3 and FU4 are flying in a string of pearls configuration with FU4 ahead, to resolve the space-time ambiguity inherent to single spacecraft missions such as SAMPEX. Each FIREBIRD-II unit has a collimated and a surface solid state detector with complementary fields of view of 54° and 90°, respectively. They are observing electron fluxes in six energy channels from ~ 230 keV to > 1 MeV. The adjustable sampling rate is 18.75 ms by default and can be as fast as 12.5 ms.

On February 2nd, 2015 at 06:12:53 UT, a microburst and subsequent bounces were observed simultaneously on both spacecraft. Figure 1 shows the High Resolution (HiRes) electron flux of the microburst, sampled at 18.75 ms. Five peaks were observed on both spacecraft. On the collimated detector, the microburst was seen up to the fourth energy channel (555 - 771 keV), while on the surface detector it was observed up to the fifth energy channel (683 - 950 keV). Only FU3 has a functioning surface detector, thus only data from the lowest four energy channels of the collimated detectors was used for this analysis. Furthermore, since FU4's 5th peak in the fourth energy channel was buried in Poisson noise, only the first four peaks were used in the spatial scale analysis.

From the black vertical black bars in Fig. 1, the first peak does not appear to be dispersed, and subsequent peaks show dispersion consistent across energy channels. This dispersion signature and amplitude decay implies that the first peak was observed soon after the electrons were scattered, followed by decaying bounces. At this time, the spacecraft was above Sweden, latitude = 63°N, longitude = 15°E, altitude = 650 km, at the eastern edge of the bounce loss cone (BLC). For this analysis, the BLC is defined as the region where locally mirroring electrons will mirror at an altitude less than 100 km in the opposite hemisphere, and could be lost due to collisions in the atmosphere [Abel and Thorne, 1998]. This unique observation of returning bounces was possible since FIREBIRD-II was at the eastern edge of the BLC; any background electron flux that was in the drift loss cone was recently lost to the South Atlantic Anomaly. In the drift loss cone, any returning bounces from microbursts will be quickly hidden in the background electron flux. For locally mirroring electrons, the mirror point in the opposite hemisphere was calculated to be 95 km using the Tsyganenko 1989 (T89) magnetic field model [Tsyganenko, 1989] with the IRBEM-Lib library. From the analysis done by Fang *et al.* [2010], the peak total ionization rate in the atmosphere for 100 keV electrons is around 80 km altitude, while 1 MeV electrons have a peak total ionization rate around 60 km altitude, so it is expected that a fraction of the microburst electrons will sur-

vive each encounter with the atmosphere. By plotting the peak flux as a function of bounce, it was found that 40 - 60 % of the microburst electrons were lost on the first bounce, similar to the 33% loss per bounce observed for a bouncing microburst observed by SAMPEX [Thorne *et al.*, 2005].

The geomagnetic conditions at the time were  $K_p = 4$ ,  $AE \approx 400$  nT and  $DST = -44$  nT, during the transition between the main and recovery phases of a geomagnetic storm. At the time of the microburst, FIREBIRD-II was at McIlwain  $L = 4.7$  and MLT = 8.3, calculated with the T89 model.

### 3 Analysis

#### 3.1 Time and position correction

At the beginning of the FIREBIRD-II mission, there were two issues with the spacecraft timing and position; their clocks were not synchronized and there was uncertainty in their separation. The approach to solve the former problem was to calculate the relative clock difference  $\delta t_c$  between the two spacecraft. The clock difference was calculated via a cross-correlation time lag analysis on trains of microbursts that hit both spacecraft simultaneously. A time lag can be calculated since the amplitudes and repetition rate of individual microbursts in these trains is unique (see Figs. S1-S4).

To correct FIREBIRD-II's relative clock difference, four time periods with coincident microbursts were hand-picked on February 2nd, 2015. The clock difference from the simultaneous microbursts was linearly fit to account for the relative clock drift ( $\approx 20$  ms/hour at this time), giving a value of  $\delta t_c = 2.28 \pm 0.12$  s at the time of the microburst analyzed here. This time shift was applied to the HiRes data in Fig. 1. The clock difference was also independently calculated using the time stamps of the FIREBIRD-II telemetry beacons during operational passes. Since they had a common time reference, the ground station computer, a time difference  $\delta t_c = 2.45^{+0.51}_{-0.98}$  s was derived.

To calculate the spacecraft separation, the data was first corrected for the clock difference, and the same cross-correlation time lag analysis was applied to events that are assumed to be stationary (see Figs. S5 and S6). The time lag between spatial events,  $\delta t_d$  is the time difference between when the leading satellite observed a stationary spatial feature, to when the trailing satellite observed the same stationary spatial feature. This analysis yielded a time lag of  $\delta t_d = 2.64 \pm 0.12$  s. In the last step in the separation calculation, their velocity needed to be calculated using a Two Line Element (TLE), a data format containing the Keplerian elements for orbit propagation. With the TLE derived spacecraft velocity,  $v = 7.57$  km/s, the calculated spacecraft separation was,  $d = 19.9 \pm 0.9$  km.

An independent method to confirm the spacecraft separation was developed. The separation was calculated using TLEs. The TLE released for February 2nd was anomalous and was not used. Instead, seven TLEs released up to five days after the microburst event were backpropagated, using the SGP-4 algorithm [Hoots and Roehrich, 1980] that calculates orbital state vectors with perturbations such as Earth's atmosphere, as well as gravitational effects from the moon and sun. Then the predicted spacecraft separations at the time of the microburst event were averaged to derive a separation of  $d = 18.4 \pm 1.5$  km. These two methods give similar results, which imply that the stationary event assumption used in the cross-correlation time lag analysis is reasonable.

#### 3.2 Microburst Scale Sizes

With the spacecraft orbit, corrected separation, and timeseries from Fig. 1, a map in longitude-latitude was made and is shown in Fig. 2. The locations where FU3 saw peaks 1-5 and where FU4 saw peaks 1-4 are shown as P1-5 and P1-4, respectively. The lower bound on the latitudinal extent of the microburst was the difference in latitude between P1 on FU3

and P4 on FU4 and was found to be  $28.8 \pm 0.8$  km. The uncertainty was estimated from the spacecraft separation uncertainty calculated in section 3.1. This scale size is the largest reported by FIREBIRD-II.

To calculate the longitudinal scale size of the microburst, it is assumed that electrons observed in the last bounce by FIREBIRD-II must have begun drifting when the initial microburst was observed, to the west of the spacecraft. Following geometrical arguments, the distance that electrons drift azimuthally in a single bounce is a product of the circumference of the drift shell foot point, and the fraction of the total drift orbit traversed in a single bounce and is given by,

$$d_{az} = 2\pi(R_E + A) \cos(\lambda) \frac{t_b}{\langle T_d \rangle} \quad (1)$$

where  $R_E$  is the Earth's radius,  $A$  is the spacecraft altitude,  $\lambda$  is the magnetic latitude,  $t_b$  is the electron bounce period, and  $\langle T_d \rangle$  is the electron drift period. Parks [2003] derived  $\langle T_d \rangle$  to be,

$$\langle T_d \rangle \approx \begin{cases} 43.8/(L \cdot E) & \text{if } \alpha_0 = 90^\circ \\ 62.7/(L \cdot E) & \text{if } \alpha_0 = 0^\circ \end{cases} \quad (2)$$

where  $E$  is the electron energy in MeV,  $L$  is the  $L$  shell, and  $\alpha_0$  is the equatorial pitch angle. Electrons mirroring at FIREBIRD-II have  $\alpha_0 \approx 3.7^\circ$  and so the  $\alpha_0 = 0^\circ$  limit was used.

The microburst's longitudinal scale size is defined as the furthest distance that its electrons drifted east and were last seen. This was calculated with  $D_{az} = n d_{az}$  where  $n$  is the number of bounces observed. The stars with energy labels in Fig. 2 represent the locations of electrons with that energy when the microburst was seen at P1, and drifted eastward to be last seen at P5 for FU3 and P4 for FU4. Using this methodology, the longitudinal scale size was greater than  $38.5 \pm 8.8$  km for the 555 keV electrons and greater than  $50.8 \pm 11.4$  km for the 771 keV electrons, shown with the red dashed box in Fig. 2. The uncertainty was estimated by propagating the uncertainty in the spacecraft separation through Eq. 1.

The longitudinal and latitudinal scale sizes and their uncertainties at LEO were mapped to the magnetic equator using the T89 magnetic field model. The radial scale size (latitudinal scale mapped from LEO) is greater than  $504 \pm 14$  km and azimuthal scale size (longitudinal scale mapped from LEO) of 555 keV electrons is greater than  $451 \pm 103$  km and of 771 keV electrons is greater than  $530 \pm 119$  km.

### 3.3 Electron Bounce Period

Lastly, the observed bounce period,  $t_b$  as a function of energy was calculated. To calculate the observed  $t_b$  and uncertainties, the raw HiRes flux was baseline-subtracted and fitted. The baseline flux used in this analysis is given in *O'Brien et al.* [2004] as the flux at the 10th percentile over a specified time interval, which in this analysis was taken to be 0.5 seconds. The flux was fitted with a superposition of Gaussians for each energy channel. The flux uncertainty is from the Poisson errors of microburst and baseline fluxes summed in quadrature. Using the fit parameters, the mean  $t_b$  for the lowest four energy channels was calculated and shown in Fig. 3 with green and purple rectangles.

The bounces observed with FU3 had in-channel dispersion to earlier times in the lowest two energy channels. This hints at the underlying distribution of electron flux within those energy channels. If FIREBIRD-II had finer energy resolution, the peak flux would be towards the upper end of the first energy channel. A Gaussian fit cannot account for this in-channel dispersion, and as a first order correction, minima between peaks was used to calculate  $t_b$ , and is shown in Fig. 3 with blue rectangles. The observed energy-dependent dispersion shown in Fig. 3 is consistent with higher energy peaks returning sooner. This dispersion consistency further supports the assumption that the subsequent peaks are bounces, and not a train of microbursts scattered by bouncing chorus.

Superposed in Fig. 3, are  $t_b$  curves for various models including an analytical solution in a dipole [Schulz and Lanzerotti, 1974], and numerical models: T89, Tsyganenko 2004 (T04) [Tsyganenko and Sitnov, 2005], and Olson & Pfizter Quiet [Olson and Pfizter, 1982]. The numerical  $t_b$  curves were calculated using a Python wrapper for IRBEM-Lib. It traces the magnetic field line between mirror points, to calculate  $t_b$  assuming conservation of energy and the first adiabatic invariant for electrons mirroring at FIREBIRD-II. A discrepancy of  $\approx 20\%$  between the observed bounce periods derived from the fits and the models is seen at the lower energies, and is reduced at the higher energy channels. The low energy discrepancy is removed by using the minima to calculate  $t_b$  for the  $< 555$  keV peaks. Across all but one energy channel, the T04 model has the largest discrepancy compared to other models.

## 4 Discussion

The twin FIREBIRD-II CubeSats have enabled a direct estimate of the lower bound scale size of a large microburst. This microburst is larger than the latitudinal scale sizes of  $> 1$  MeV microbursts reported in Blake *et al.* [1996], and similar to the scale sizes of  $> 15$  keV microbursts observed with a high altitude balloon [Parks, 1967]. Furthermore, the latitudinal scale size in this study was roughly  $\sim 2.6$  times larger than other simultaneous microbursts reported in Crew *et al.* [2016] and  $\sim 10$  times larger than than reported in Dietrich *et al.* [2010]. No energy dependence on the scale size was observed.

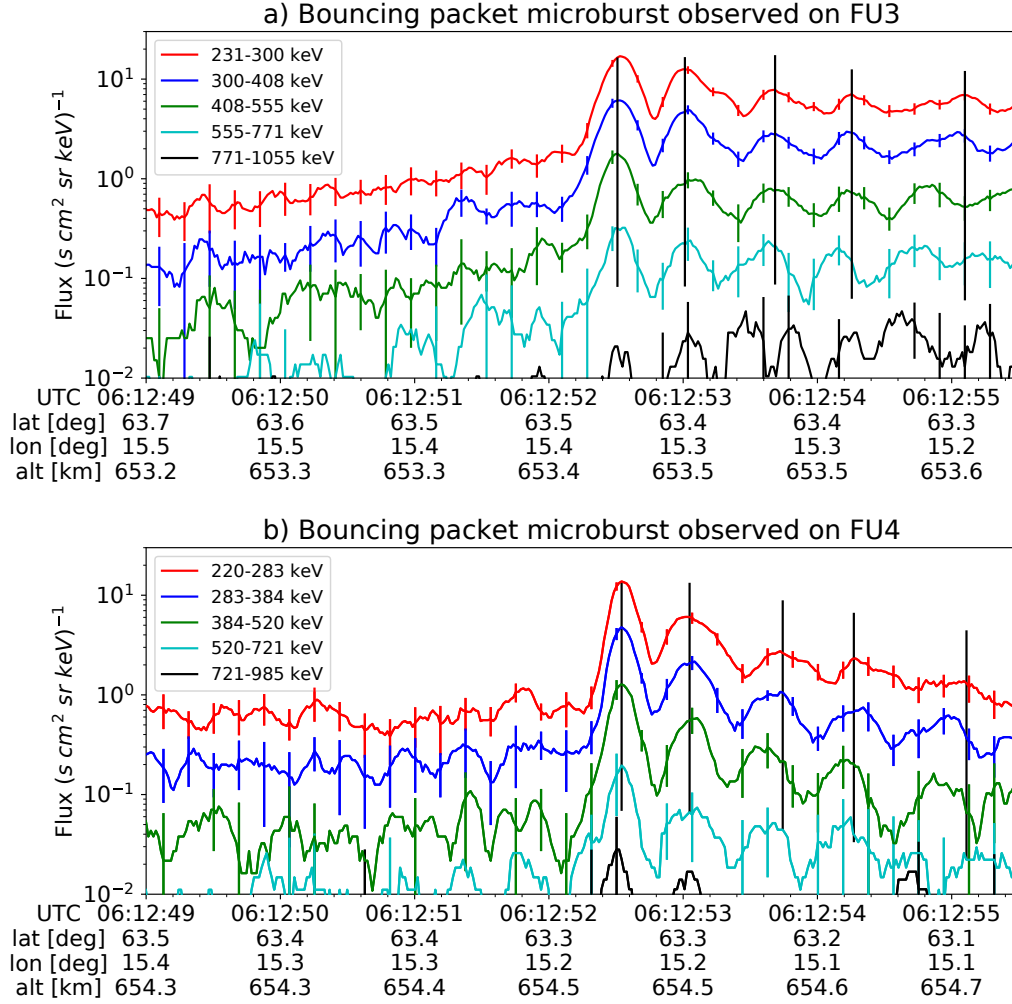
The microburst scale size obtained in Section 3.2 and scaled to the geomagnetic equator can be compared with the scales of chorus waves presumably responsible for the rapid burst electron precipitation. Early direct estimates of the chorus source scales were made by the coordinated measurement by ISEE-1, 2. The wave power correlation scale was estimated to be about several hundred kilometers across the background magnetic field [Gurnett *et al.*, 1979]. Santolik *et al.* [2003] determined the correlation lengths of chorus-type whistler waves to be around 100 km based on multipoint CLUSTER WBD measurements near the chorus source region during the magnetic storm of 18 April 2002 at L shell about 4. Agapitov *et al.* [2010, 2011, 2017] recently showed that the spatial extent of chorus source region can be larger: from 600 km in the outer radiation belt to more than 1000 km in the outer magnetosphere. The lower bound azimuthal and latitudinal scales obtained in Section 3.2 and scaled to the magnetic equator, were similar to the whistler-mode chorus source scale sizes reported in Agapitov *et al.* [2011, 2017]. No wave measurements from nearby spacecraft were available at this time to confirm the presence of chorus. Nevertheless, since AE  $\sim 400$  nT at this time, relatively strong chorus waves were statistically more likely to occur [Li *et al.*, 2009]. The microburst electrons were likely scattered by chorus, similar to the conclusions made by e.g. Lorentzen *et al.* [2001a]; O'Brien *et al.* [2003]; Breneman *et al.* [2017].

Using the fit parameters from section 3.3, the exponential E-folding energy,  $E_0$  was calculated to be  $E_0 \sim 100$  keV. This is similar to the results in Lee *et al.* [2005] who used STSAT-1 and Datta *et al.* [1997] who used a sounding rocket. It is soft for a typical microburst observed with FIREBIRD-II. There was no statistically significant change in  $E_0$  for subsequent bounces.

## 5 Conclusions

This was a first observation of a large microburst with multiple bounces made possible by the twin FIREBIRD-II CubeSats. Its lower bound LEO latitudinal and longitudinal scale sizes of  $28.8 \pm 0.8$  km and  $50.8 \pm 11.4$  km make it one of the largest observed. No energy dependence on the scale size was observed. Furthermore, its lower bound LEO scale sizes mapped to the magnetic equator were  $504 \pm 14$  km radially and  $530 \pm 119$  km azimuthally. The bounce periods calculated with the Gaussian fit and minima methods are in good agreement with the numerically estimated bounce periods, confirming that the energy-dependent dispersion was due to bouncing. The similarity of the microburst and chorus source region scale



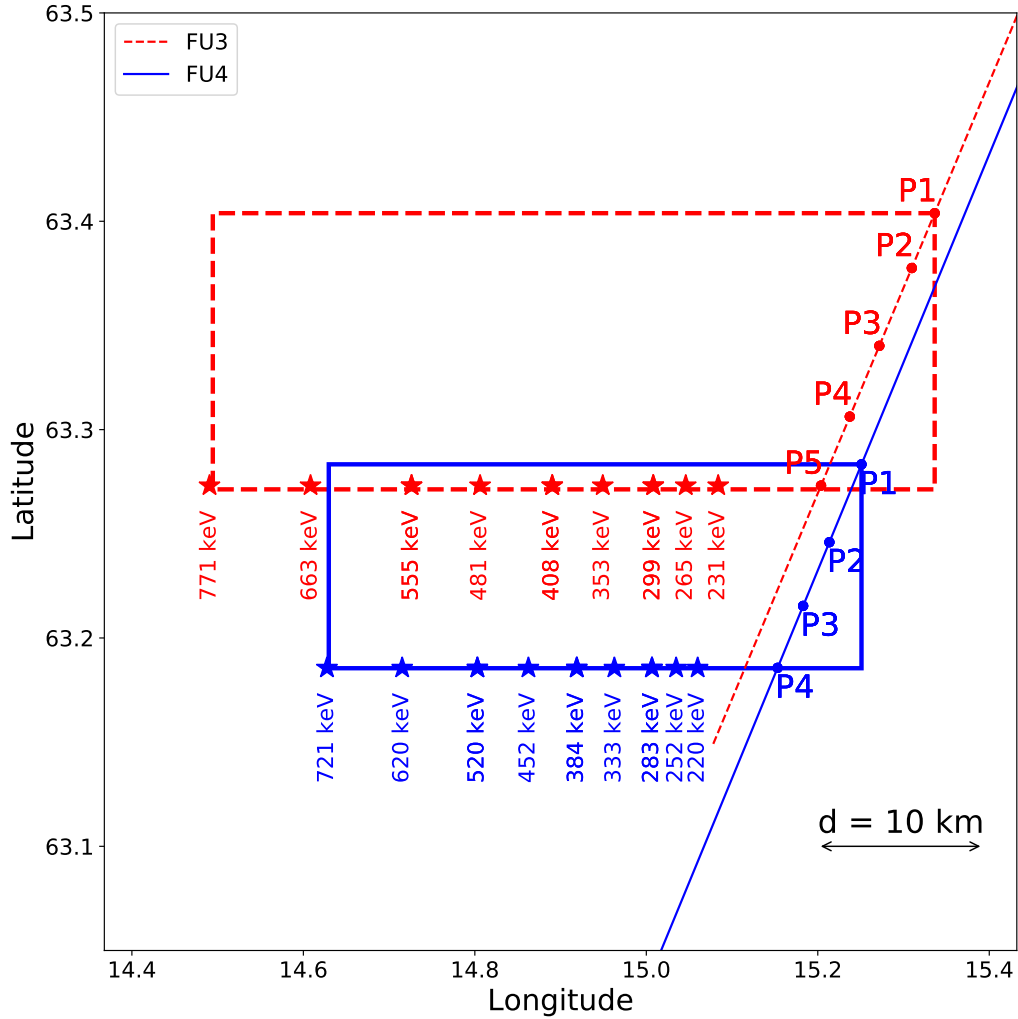


**Figure 1.** HiRes data of the microburst observed at February 2nd, 2015 at 06:12:53 UT, smoothed with a 150 ms rolling average. The subsequent bounces showed some energy dispersion. As discussed in section 3 a time correction of -2.28 s was applied to FU3. While the flux from five energy channels is shown, only channels with reasonable counting statistics were used for the spatial scale analysis. Vertical colored bars show the  $\sqrt{N}$  error every 10th data point and vertical black bars are lined up with the peaks in the 220-283 keV energy channel to help identify dispersion.

sizes, as well as magnetospheric location and condition, support the idea that microbursts are produced by electrons scattered by whistler-mode chorus waves.

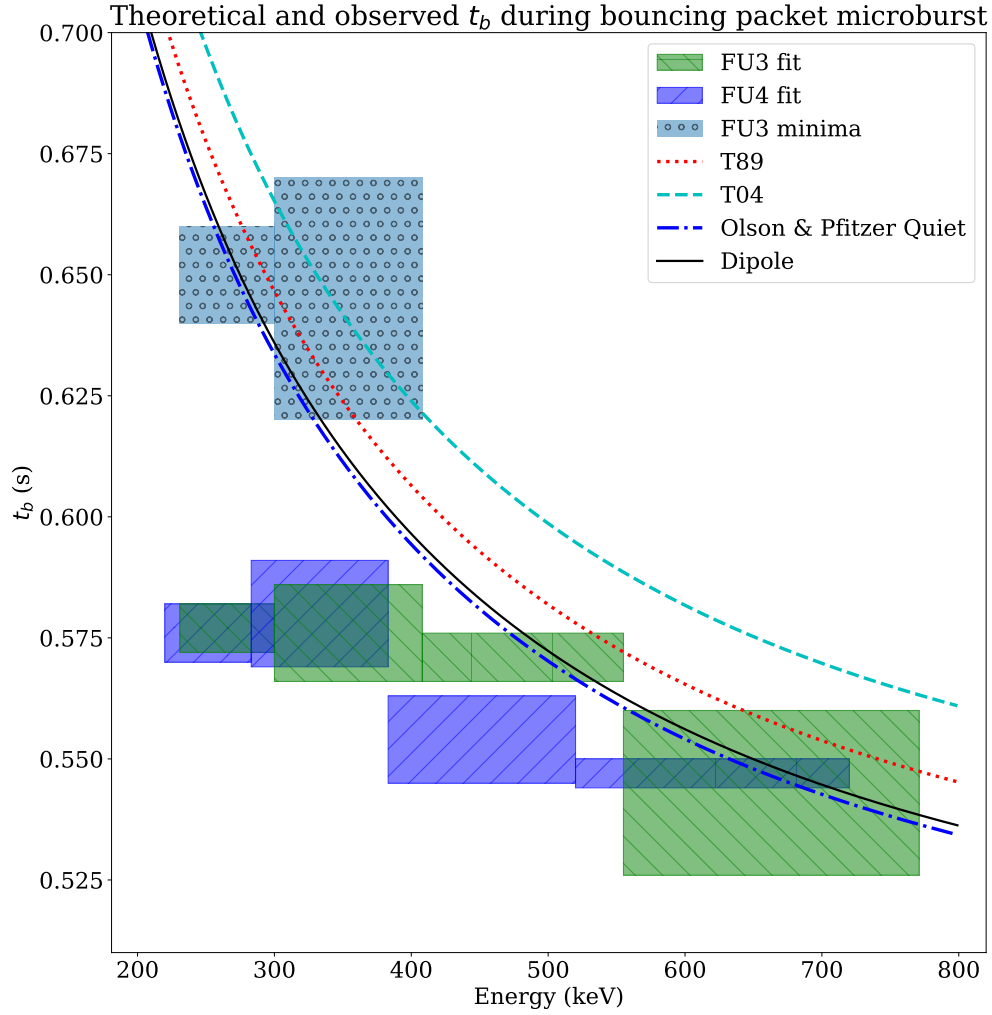
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**Figure 2.** The topology of the FIREBIRD-II orbit and the multiple bounces of the microburst projected onto latitude and longitude with axis scaled to equal distance. Attributes relating to FU3 shown in red dashed lines, and FU4 with blue solid lines. The spacecraft path is shown with the diagonal lines, starting at the upper right corner. The labels P1-4 for FU4 and P1-5 for FU3 indicate where the spacecraft were when the  $N^{th}$  peak was seen in the lowest energy channel in the HiRes data. The stars with the accompanying energy labels represent the locations of the electrons with that energy that started at time of P1, and were seen at the last peak on each spacecraft. The rectangles represent the lower bound of the microburst scale size, assuming that the majority of the electrons were in the upper boundary of energy channel 4.





**Figure 3.** Observed and theoretical  $t_b$  for electrons of energies from 200 to 770 keV. The solid black line is  $t_b$  in a dipole magnetic field, derived in *Schulz and Lanzerotti* [1974]. The red and cyan dashed lines are the  $t_b$  derived using the T89, and T04 magnetic field models with IRBEM. Lastly, the blue dashed curve is the  $t_b$  derived using the Olson & Pfitzer Quiet model. The green and purple rectangles represent the observed  $t_b$  for FU3 and FU4 using a Gaussian fit, respectively. The blue rectangles represent the observed  $t_b$  calculated with the minima between the bounces. The width of the boxes represent the width of those energy channels, and the height represents the uncertainty from the fit.

## References

- Abel, B., and R. M. Thorne (1998), Electron scattering loss in earth's inner magnetosphere: 1. dominant physical processes, *Journal of Geophysical Research: Space Physics*, 103(A2), 2385–2396.
- Agapitov, O., V. Krasnoselskikh, Y. Zaliznyak, V. Angelopoulos, O. Le Contel, and G. Rolland (2010), Chorus source region localization in the earth's outer magnetosphere using themis measurements, *Annales Geophysicae*, 28(6), 1377–1386, doi:10.5194/angeo-28-1377-2010.
- Agapitov, O., V. Krasnoselskikh, T. Dudok de Wit, Y. Khotyaintsev, J. S. Pickett, O. Santolik, and G. Rolland (2011), Multispacecraft observations of chorus emissions as a tool for the plasma density fluctuations' remote sensing, *Journal of Geophysical Research: Space Physics*, 116(A9), n/a–n/a, doi:10.1029/2011JA016540, a09222.
- Agapitov, O., L. W. Blum, F. S. Mozer, J. W. Bonnell, and J. Wygant (2017), Chorus whistler wave source scales as determined from multipoint van allen probe measurements, *Geophysical Research Letters*, pp. n/a–n/a, doi:10.1002/2017GL072701, 2017GL072701.
- Anderson, B., S. Shekhar, R. Millan, A. Crew, H. Spence, D. Klumpar, J. Blake, T. O'Brien, and D. Turner (2017), Spatial scale and duration of one microburst region on 13 august 2015, *Journal of Geophysical Research: Space Physics*.
- Anderson, K. A., and D. W. Milton (1964), Balloon observations of x rays in the auroral zone: 3. high time resolution studies, *Journal of Geophysical Research*, 69(21), 4457–4479, doi:10.1029/JZ069i021p04457.
- Blake, J., M. Looper, D. Baker, R. Nakamura, B. Klecker, and D. Hovestadt (1996), New high temporal and spatial resolution measurements by sampex of the precipitation of relativistic electrons, *Advances in Space Research*, 18(8), 171 – 186, doi: [http://dx.doi.org/10.1016/0273-1177\(95\)00969-8](http://dx.doi.org/10.1016/0273-1177(95)00969-8).
- Blum, L., X. Li, and M. Denton (2015), Rapid mev electron precipitation as observed by sampex/hilt during high-speed stream-driven storms, *Journal of Geophysical Research: Space Physics*, 120(5), 3783–3794, doi:10.1002/2014JA020633, 2014JA020633.
- Breneman, A., A. Crew, J. Sample, D. Klumpar, A. Johnson, O. Agapitov, M. Shumko, D. Turner, O. Santolik, J. Wygant, et al. (2017), Observations directly linking relativistic electron microbursts to whistler mode chorus: Van allen probes and firebird ii, *Geophysical Research Letters*.
- Crew, A. B., H. E. Spence, J. B. Blake, D. M. Klumpar, B. A. Larsen, T. P. O'Brien, S. Driscoll, M. Handley, J. Legere, S. Longworth, K. Mashburn, E. Mosleh, N. Ryhajlo, S. Smith, L. Springer, and M. Widholm (2016), First multipoint in situ observations of electron microbursts: Initial results from the nsf firebird ii mission, *Journal of Geophysical Research: Space Physics*, 121(6), 5272–5283, doi:10.1002/2016JA022485, 2016JA022485.
- Datta, S., R. Skoug, M. McCarthy, and G. Parks (1997), Modeling of microburst electron precipitation using pitch angle diffusion theory, *Journal of Geophysical Research: Space Physics*, 102(A8), 17,325–17,333.
- Dietrich, S., C. J. Rodger, M. A. Clilverd, J. Bortnik, and T. Raita (2010), Relativistic microburst storm characteristics: Combined satellite and ground-based observations, *Journal of Geophysical Research: Space Physics*, 115(A12).
- Fang, X., C. E. Randall, D. Lummerzheim, W. Wang, G. Lu, S. C. Solomon, and R. A. Frahm (2010), Parameterization of monoenergetic electron impact ionization, *Geophysical Research Letters*, 37(22).
- Gurnett, D., R. Anderson, F. Scarf, R. Fredricks, and E. Smith (1979), Initial results from the isee-1 and-2 plasma wave investigation, *Space Science Reviews*, 23(1), 103–122.
- Hoots, F. R., and R. L. Roehrich (1980), Models for propagation of norad element sets, *Tech. Rep. 3*, Spacetrack.
- Horne, R. B., and R. M. Thorne (2003), Relativistic electron acceleration and precipitation during resonant interactions with whistler-mode chorus, *Geophysical Research Letters*,

- 30(10), n/a–n/a, doi:10.1029/2003GL016973, 1527.
- Lee, J.-J., G. K. Parks, K. W. Min, H. J. Kim, J. Park, J. Hwang, M. P. McCarthy, E. Lee, K. S. Ryu, J. T. Lim, E. S. Sim, H. W. Lee, K. I. Kang, and H. Y. Park (2005), Energy spectra of 170–360 keV electron microbursts measured by the Korean STSAT-1, *Geophysical Research Letters*, 32(13), doi:10.1029/2005GL022996, 113106.
- Lee, J. J., G. K. Parks, E. Lee, B. T. Tsurutani, J. Hwang, K. S. Cho, K.-H. Kim, Y. D. Park, K. W. Min, and M. P. McCarthy (2012), Anisotropic pitch angle distribution of 100 keV microburst electrons in the loss cone: measurements from STSAT-1, *Annales Geophysicae*, 30(11), 1567–1573, doi:10.5194/angeo-30-1567-2012.
- Li, W., R. M. Thorne, V. Angelopoulos, J. Bortnik, C. M. Cully, B. Ni, O. LeContel, A. Roux, U. Auster, and W. Magnes (2009), Global distribution of whistler-mode chorus waves observed on the THEMIS spacecraft, *Geophysical Research Letters*, 36(9), n/a–n/a, doi:10.1029/2009GL037595, 109104.
- Lorentzen, K. R., J. B. Blake, U. S. Inan, and J. Bortnik (2001a), Observations of relativistic electron microbursts in association with VLF chorus, *Journal of Geophysical Research: Space Physics*, 106(A4), 6017–6027, doi:10.1029/2000JA003018.
- Lorentzen, K. R., M. D. Looper, and J. B. Blake (2001b), Relativistic electron microbursts during the geomagnetic storms, *Geophysical Research Letters*, 28(13), 2573–2576, doi:10.1029/2001GL012926.
- Millan, R., and R. Thorne (2007), Review of radiation belt relativistic electron losses, *Journal of Atmospheric and Solar-Terrestrial Physics*, 69(3), 362–377, doi: <http://dx.doi.org/10.1016/j.jastp.2006.06.019>, Global Aspects of Magnetosphere-Ionosphere Coupling.
- Millan, R. M., R. Lin, D. Smith, K. Lorentzen, and M. McCarthy (2002), X-ray observations of MeV electron precipitation with a balloon-borne germanium spectrometer, *Geophysical Research Letters*, 29(24).
- Nakamura, R., D. N. Baker, J. B. Blake, S. Kanekal, B. Klecker, and D. Hovestadt (1995), Relativistic electron precipitation enhancements near the outer edge of the radiation belt, *Geophysical Research Letters*, 22(9), 1129–1132, doi:10.1029/95GL00378.
- Nakamura, R., M. Isawa, Y. Kamide, D. Baker, J. Blake, and M. Looper (2000), Observations of relativistic electron microbursts in association with VLF chorus, *J. Geophys. Res.*, 105, 15,875–15,885.
- O'Brien, T. P., K. R. Lorentzen, I. R. Mann, N. P. Meredith, J. B. Blake, J. F. Fennell, M. D. Looper, D. K. Milling, and R. R. Anderson (2003), Energization of relativistic electrons in the presence of ULF power and MeV microbursts: Evidence for dual ULF and VLF acceleration, *Journal of Geophysical Research: Space Physics*, 108(A8), n/a–n/a, doi:10.1029/2002JA009784, 1329.
- O'Brien, T. P., M. D. Looper, and J. B. Blake (2004), Quantification of relativistic electron microburst losses during the geomagnetic storms, *Geophysical Research Letters*, 31(4), n/a–n/a, doi:10.1029/2003GL018621, 104802.
- Olson, W. P., and K. A. Pfister (1982), A dynamic model of the magnetospheric magnetic and electric fields for July 29, 1977, *Journal of Geophysical Research: Space Physics*, 87(A8), 5943–5948, doi:10.1029/JA087iA08p05943.
- Parks, G. (2003), *Physics Of Space Plasmas: An Introduction, Second Edition*, Westview Press.
- Parks, G. K. (1967), Spatial characteristics of auroral-zone X-ray microbursts, *Journal of Geophysical Research*, 72(1), 215–226.
- Santolik, O., D. Gurnett, J. Pickett, M. Parrot, and N. Cornilleau-Wehrlin (2003), Spatio-temporal structure of storm-time chorus, *Journal of Geophysical Research: Space Physics*, 108(A7).
- Schulz, M., and L. J. Lanzerotti (1974), *Particle Diffusion in the Radiation Belts*, Springer.
- Selesnick, R. S., J. B. Blake, and R. A. Mewaldt (2003), Atmospheric losses of radiation belt electrons, *Journal of Geophysical Research: Space Physics*, 108(A12), doi:10.1029/2003JA010160, 1468.

- Shprits, Y. Y., and R. M. Thorne (2004), Time dependent radial diffusion modeling of relativistic electrons with realistic loss rates, *Geophysical Research Letters*, *31*(8), n/a–n/a, doi:10.1029/2004GL019591, 108805.
- Shprits, Y. Y., N. P. Meredith, and R. M. Thorne (2007), Parameterization of radiation belt electron loss timescales due to interactions with chorus waves, *Geophysical Research Letters*, *34*(11), n/a–n/a, doi:10.1029/2006GL029050, 111110.
- Thorne, R. M. (2010), Radiation belt dynamics: The importance of wave-particle interactions, *Geophysical Research Letters*, *37*(22), doi:10.1029/2010GL044990, 122107.
- Thorne, R. M., T. P. O'Brien, Y. Y. Shprits, D. Summers, and R. B. Horne (2005), Timescale for mev electron microburst loss during geomagnetic storms, *Journal of Geophysical Research: Space Physics*, *110*(A9), n/a–n/a, doi:10.1029/2004JA010882, a09202.
- Tsyganenko, N. (1989), A solution of the chapman-ferraro problem for an ellipsoidal magnetopause, *Planetary and Space Science*, *37*(9), 1037 – 1046, doi: [http://dx.doi.org/10.1016/0032-0633\(89\)90076-7](http://dx.doi.org/10.1016/0032-0633(89)90076-7).
- Tsyganenko, N. A., and M. I. Sitnov (2005), Modeling the dynamics of the inner magnetosphere during strong geomagnetic storms, *Journal of Geophysical Research: Space Physics*, *110*(A3), n/a–n/a, doi:10.1029/2004JA010798, a03208.
- Ukhorskiy, A. Y., B. J. Anderson, P. C. Brandt, and N. A. Tsyganenko (2006), Storm time evolution of the outer radiation belt: Transport and losses, *Journal of Geophysical Research: Space Physics*, *111*(A11), n/a–n/a, doi:10.1029/2006JA011690, a11S03.
- Woodger, L., A. Halford, R. Millan, M. McCarthy, D. Smith, G. Bowers, J. Sample, B. Anderson, and X. Liang (2015), A summary of the barrel campaigns: Technique for studying electron precipitation, *Journal of Geophysical Research: Space Physics*, *120*(6), 4922–4935.