

# Molecular Dynamics Simulations of Spherical Colloid Size Distribution Effects

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## Introduction

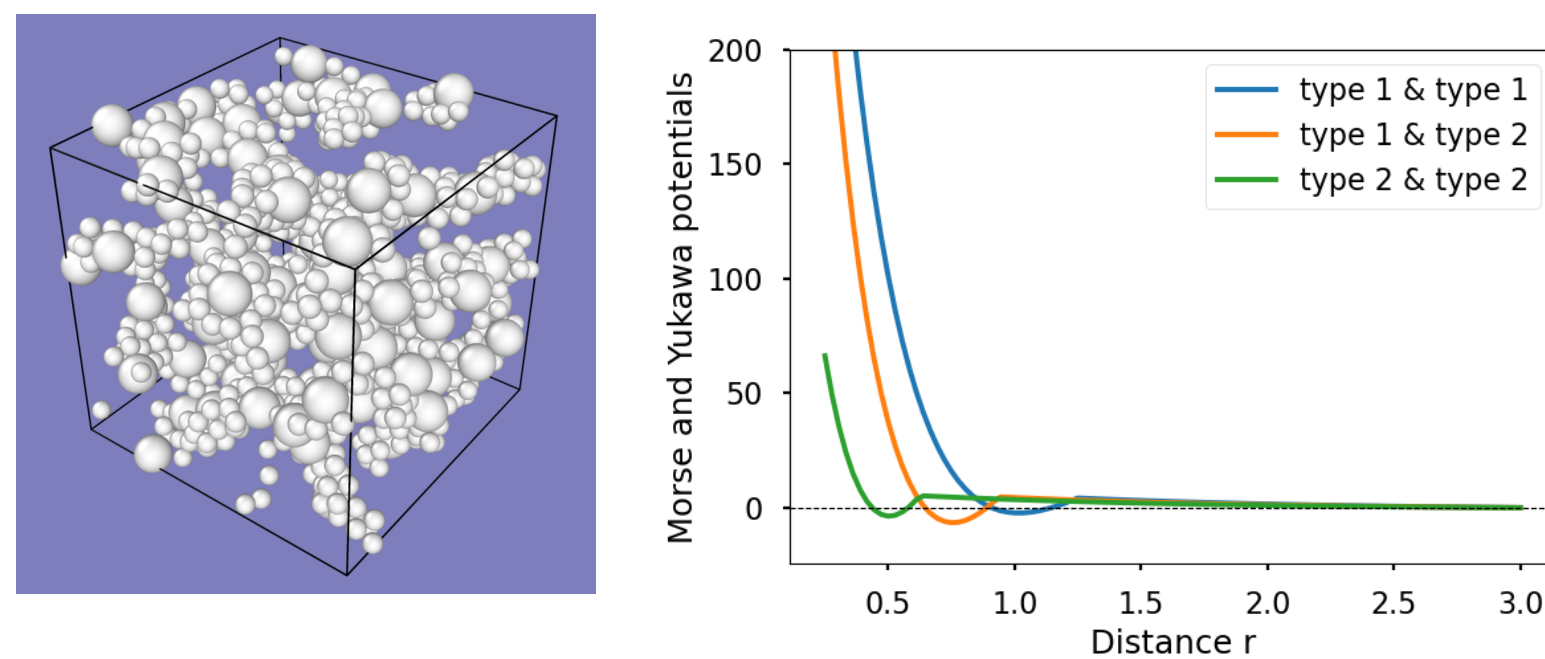
The mechanical stability of a colloidal gel depends on isostaticity networks formed during gelation [1, 2]. Molecular dynamics simulations allow a microscopic view into the percolation of these networks. This project uses the open-source simulator LAMMPS to quantify how particle contacts evolve for monodisperse and bidisperse colloidal spheres under a hybrid interaction potential.

## Methods

The Python and LAMMPS scripts used for gel initialization are adapted from a GitHub repository for shearing a polydisperse Morse gel [3]. A pairwise `yukawa/colloid` potential was added to account for long-range Coulombic repulsion.

$$V_{\text{Morse}} = D_0(e^{-2\alpha(r-r_0)} - 2e^{\alpha(r-r_0)})$$

$$V_{\text{Yukawa}} = \frac{A}{\kappa} e^{-\kappa(r-(R_i+R_j))}$$



**Fig. 1.** (Left) Snapshot of simulation box with 1000 particles in OVITO. (Right) Sum of attractive and repulsive potentials applied for each pair of particle types.

- **Spherical particles with adapted potentials.** The adapted code models the colloid particles as monodisperse or bidisperse spheres. For each pairwise potential, a cutoff was defined, and parameters were made dependent on particle diameter.
- **Dissipative Particle Dynamics (DPD) thermostat.** Time integration consistent with the microcanonical ensemble was performed with `dpd/tstat`, which applies a dissipative pairwise force and a random force at every timestep to capture solvent effects so that solvent molecules do not have to be included.

$$\mathbf{F}_{\text{DPD}} = (F_{\text{dissipative}} + F_{\text{random}})\mathbf{\hat{r}}_{ij}$$

Constant-energy time integration of the classical equations of motion was performed for  $10^6$  timesteps of 0.001. DPD causes fluctuations in total potential energy.

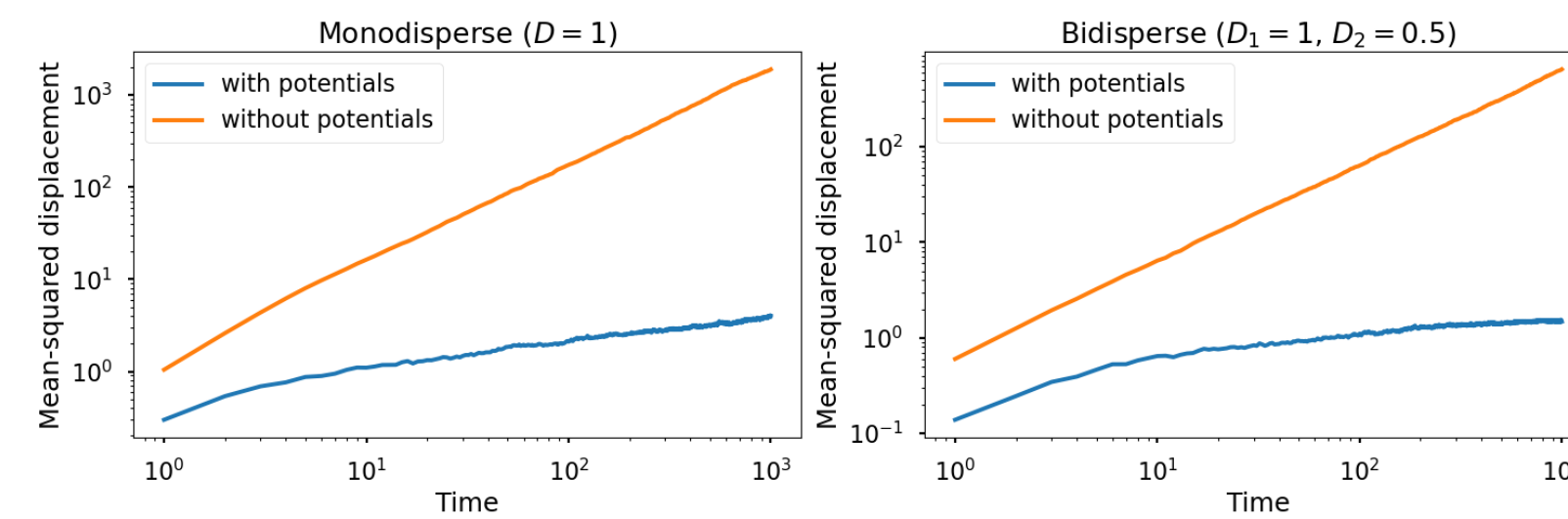
## Simulation setup

Simulations were performed in a 3D cubic box with periodic boundary conditions. Quantities are in dimensionless 1j (Lennard-Jones) units. The default particle diameter is 1.0. For bidisperse gels, some particles were assigned a diameter of 0.5 with probability corresponding to a 0.5 mass fraction (converted to a number fraction). The size of the simulation box was determined by the number density, entered as a separate parameter.

Parameters were chosen arbitrarily; rescaling is necessary before comparing to real systems. See the GitHub repository linked at the right for parameter definitions.

## Effect of pairwise potentials

In simple diffusion, the mean-squared displacement  $\Delta x_{\text{MSD}} = \langle (x(t) - x(0))^2 \rangle$  grows linearly with time  $t$ , and  $\log \Delta x_{\text{MSD}}$  grows linearly with  $\log t$ . A correct implementation of a pairwise potential should slow this growth. See “potential\_check.”

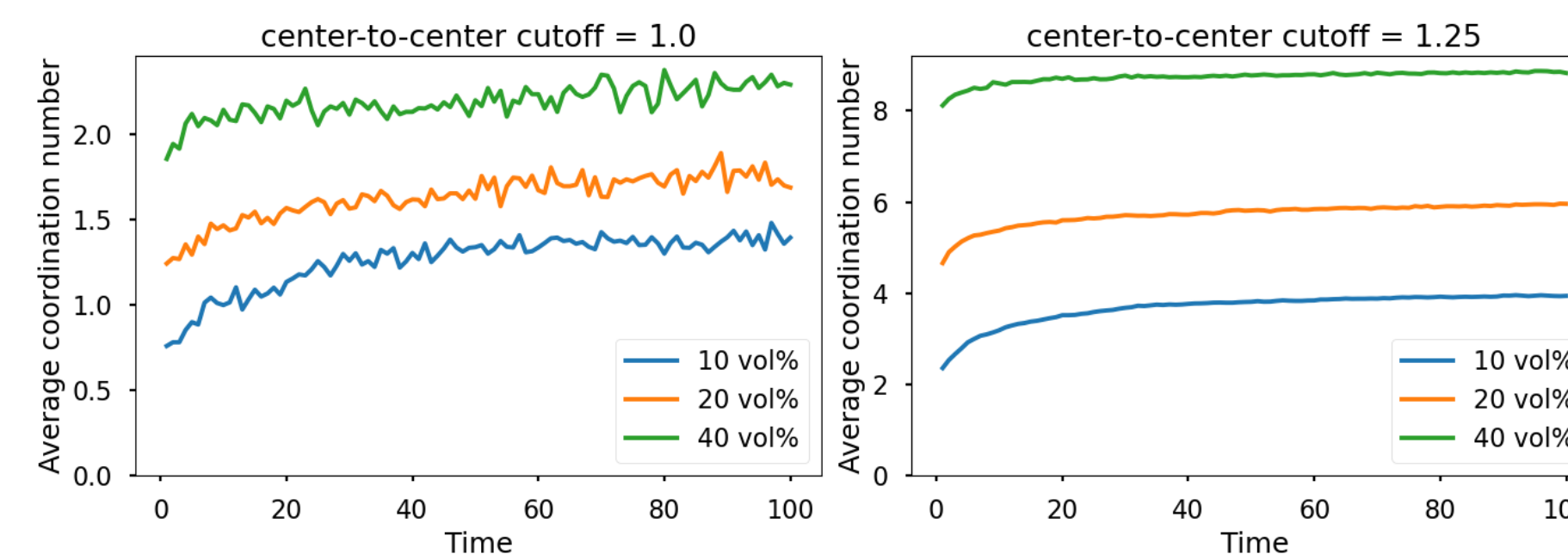


**Fig. 2.** Mean-squared displacement of 1000 particles with and without pairwise potentials.

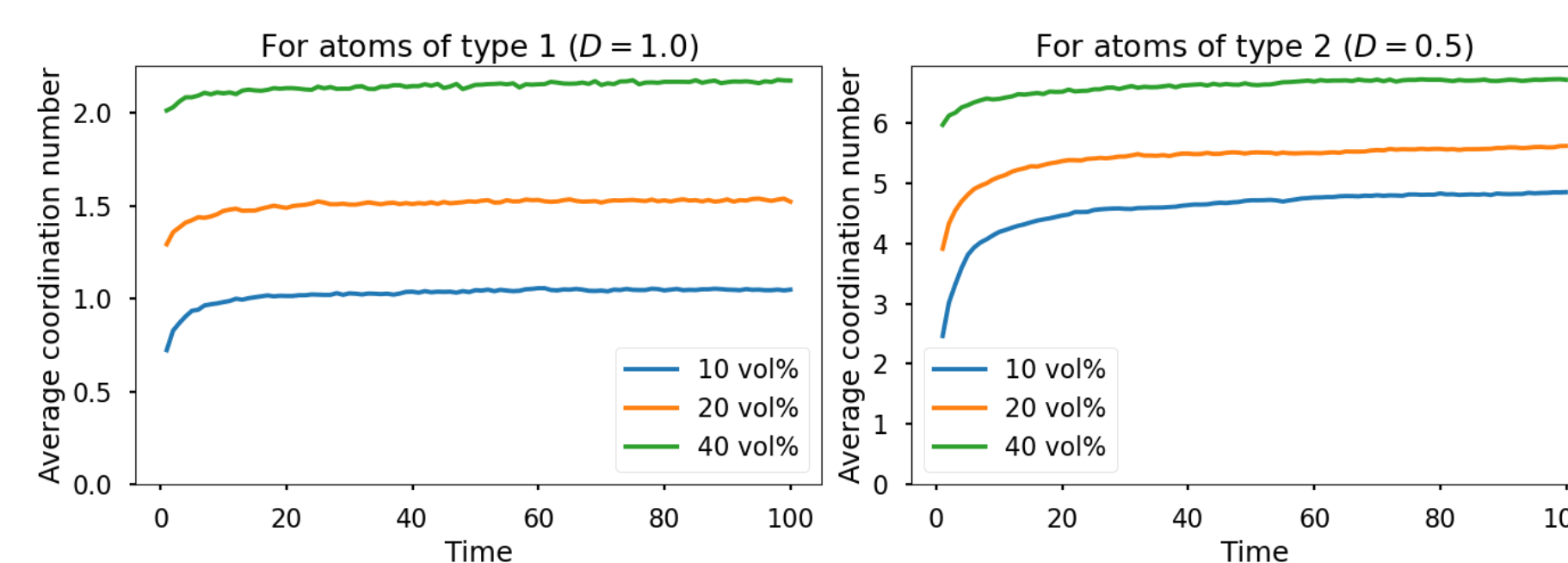
The MSD plateaus on a log-log plot, consistent with the arresting effects of gelation.

## Neighbor counts

The number of neighbors of each particle (coordination number) was calculated at fixed time intervals. Particles are considered neighbors when the center-to-center distance is within a specified cutoff.



**Fig. 3.** Evolution of average coordination number in a monodisperse gel ( $D = 1$ ), with two different cutoffs for counting neighbors. The simulation includes 1000 particles.



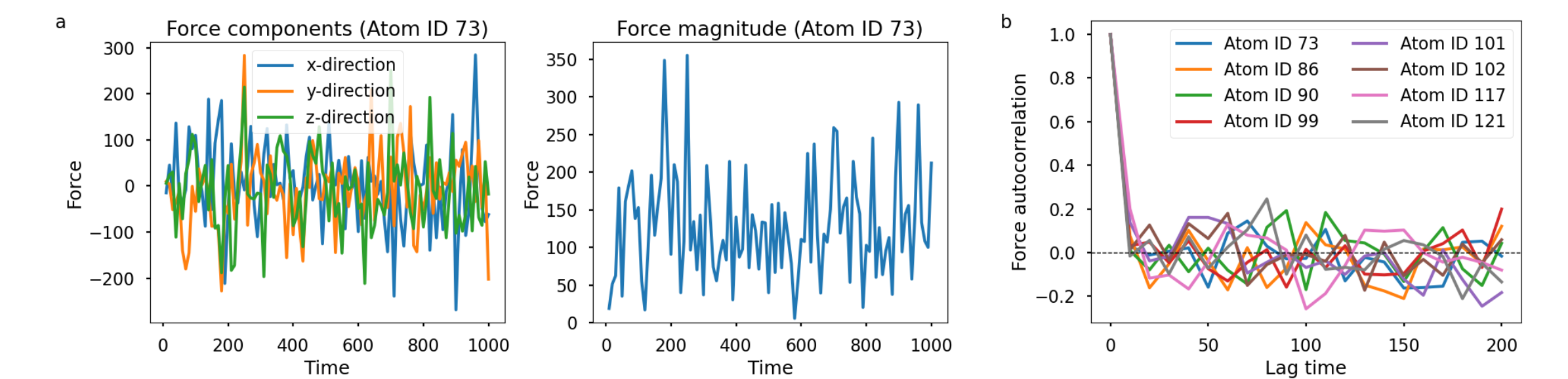
**Fig. 4.** Evolution of average coordination number in a bidisperse gel, with the  $j$  cutoff defined as  $(D_i + D_j)/2 + D_i/4$  for particle  $i$ , so that particles whose surfaces are within  $0.25D_i$  of particle  $i$ 's surface are considered neighbors. The simulation includes 1000 particles (500 of each size).

- At the same total volume fraction, atoms with  $D = 1$  approach a much lower coordination number when smaller atoms are present (also observed for a surface-to-surface cutoff of 0; see “volumefrac\_neighbors” folder).
- At a  $0.25D$  surface-to-surface neighbor cutoff, monodisperse gels with  $D = 1$  approach similar neighbor counts as particles with  $D = 0.5$  in the bidisperse gels.

## Atom-level forces

Force over time was extracted for individual atoms using “forces.py.” The normalized autocorrelation function  $R$  (defined below) was computed in “plotforces.py.”

$$R(F, t_{\text{lag}}) = \frac{\langle F(t)F(t_{\text{lag}}) \rangle}{\langle F(t)^2 \rangle}$$



**Fig. 5.** (a) Force components and magnitude over time for an arbitrarily selected type 1 ( $D = 1$ ) atom in a bidisperse gel with 10 vol% colloid, in the absence of `dpd/tstat`. (b) Autocorrelation function of the force magnitude for a selection of  $D = 1$  atoms.

Even when random energy dissipation is excluded, memory of the current force (measured by autocorrelation) decays quickly at the level of an individual atom.

## Conclusions

In a simple model with interacting spheres, post-gelation coordination number depends on particle size distribution at constant colloid volume fraction.

## Future work

Validate the DPD thermostat by including a specific solvent explicitly, and determine the appropriate parameters to enable comparison to experiment.

Vary the temperature and look for changes in the force memory of atoms.

## GitHub repository

This work was conducted as a final project for PHYS 25000 (Autumn 2025). The simulation code can be found at <https://github.com/msong2026/lammps-gels>. Sample LAMMPS scripts by Christopher Ness [4] were also consulted.

## Acknowledgments

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## References

- [1] Hideyo Tsurusawa et al. “Direct link between mechanical stability in gels and percolation of isostatic particles”. In: *Science Advances* 5.5 (2019), eaav6090. DOI: 10.1126/sciadv.aav6090.
- [2] Lilian C. Hsiao et al. “Role of isostaticity and load-bearing microstructure in the elasticity of yielded colloidal gels”. In: *Proceedings of the National Academy of Sciences* 109.40 (2012), pp. 16029–16034. DOI: 10.1073/pnas.1206742109.
- [3] Peter Crowther. *Bulk creep Molecular Dynamics simulations*. <https://github.com/merrygoat/creep-sheared-gels>. 2020.
- [4] Christopher Ness. “Simulating dense, rate-independent suspension rheology using LAMMPS”. In: *Computational Particle Mechanics* 10.6 (2023), pp. 2031–2037. DOI: 10.1007/s40571-023-00605-x.