

Molecular Dynamics Simulations of Spherical Colloid Size Distribution Effects

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Introduction

The mechanical stability of a colloidal gel depends on isostaticity networks formed during gelation [1, 2]. Molecular dynamics simulations allow a microscopic view into the percolation of these networks. This project uses the open-source simulator LAMMPS to quantify how particle contacts evolve for monodisperse and bidisperse colloidal spheres under a hybrid interaction potential.

Methods

The Python and LAMMPS scripts used for gel initialization are adapted from a GitHub repository for shearing a polydisperse Morse gel [3]. A pairwise *yukawa/colloid* potential was added to account for long-range Coulombic repulsion.

$$V_{\text{Morse}} = D_0(e^{-2\alpha(r-r_0)} - 2e^{\alpha(r-r_0)})$$

$$V_{\text{Yukawa}} = \frac{A}{\kappa} e^{-\kappa(r-(R_i+R_j))}$$

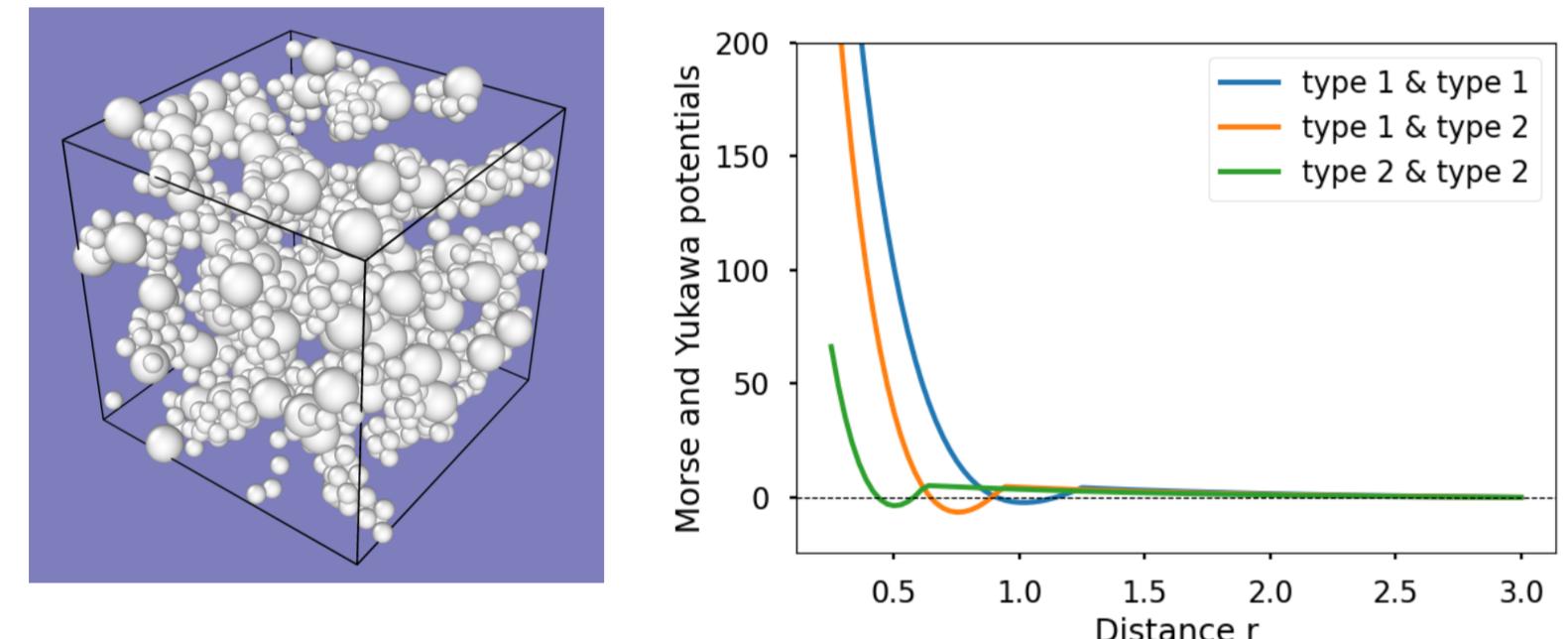


Fig. 1. (Left) Snapshot of simulation box with 1000 particles in OVITO. (Right) Sum of attractive and repulsive potentials applied for each pair of particle types.

- Spherical particles with adapted potentials.** The adapted code models the colloid particles as monodisperse or bidisperse spheres. For each pairwise potential, a cutoff was defined, and parameters were made dependent on particle diameter.
- Dissipative Particle Dynamics (DPD) thermostat.** Time integration consistent with the microcanonical ensemble was performed with `dpd/tstat`, which applies a dissipative pairwise force and a random force at every timestep to capture solvent effects so that solvent molecules do not have to be included.

$$\mathbf{F}_{\text{DPD}} = (F_{\text{dissipative}} + F_{\text{random}})\hat{\mathbf{r}}_{ij}$$

Constant-energy time integration of the classical equations of motion was performed for 10^6 timesteps of 0.001. DPD causes fluctuations in total potential energy.

Simulation setup

Simulations were performed in a 3D cubic box with periodic boundary conditions. Quantities are in dimensionless $1j$ (Lennard-Jones) units. The default particle diameter is 1.0. For bidisperse gels, some particles were assigned a diameter of 0.5 with probability corresponding to a 0.5 mass fraction (converted to a number fraction). The size of the simulation box was determined by the number density, entered as a separate parameter.

Parameters were chosen arbitrarily; rescaling is necessary before comparing to real systems. See the GitHub repository linked at the right for parameter definitions.

Effect of pairwise potentials

In simple diffusion, the mean-squared displacement $\Delta x_{\text{MSD}} = \langle (x(t) - x(0))^2 \rangle$ grows linearly with time t , and $\log \Delta x_{\text{MSD}}$ grows linearly with $\log t$. A correct implementation of a pairwise potential should slow this growth. See “potential_check.”

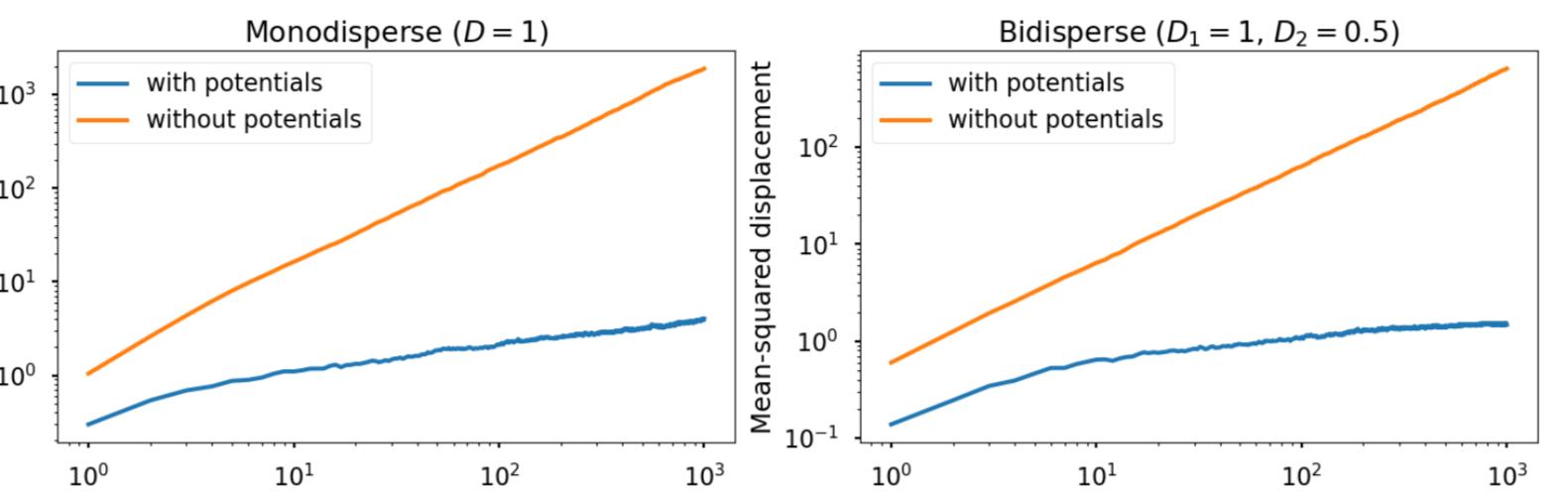


Fig. 2. Mean-squared displacement of 1000 particles with and without pairwise potentials.

The MSD plateaus on a log-log plot, consistent with the arresting effects of gelation.

Neighbor counts

The number of neighbors of each particle (coordination number) was calculated at fixed time intervals. Particles are considered neighbors when the center-to-center distance is within a specified cutoff.

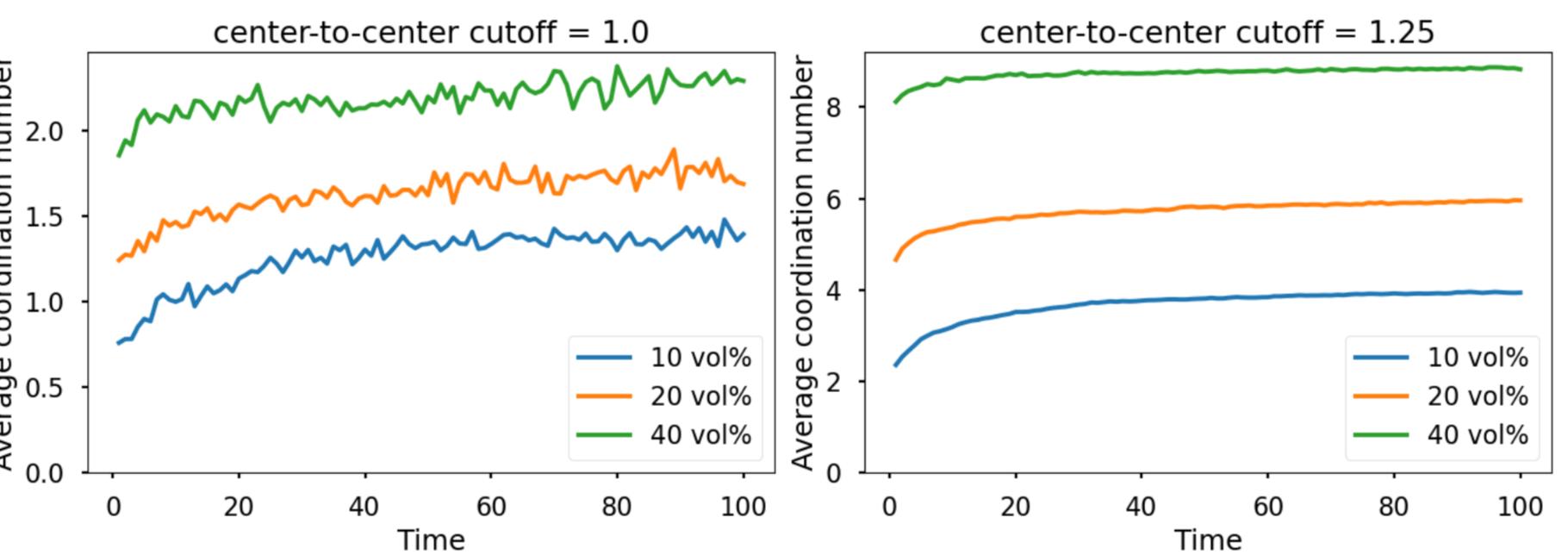


Fig. 3. Evolution of average coordination number in a monodisperse gel ($D = 1$), with two different cutoffs for counting neighbors. The simulation includes 1000 particles.

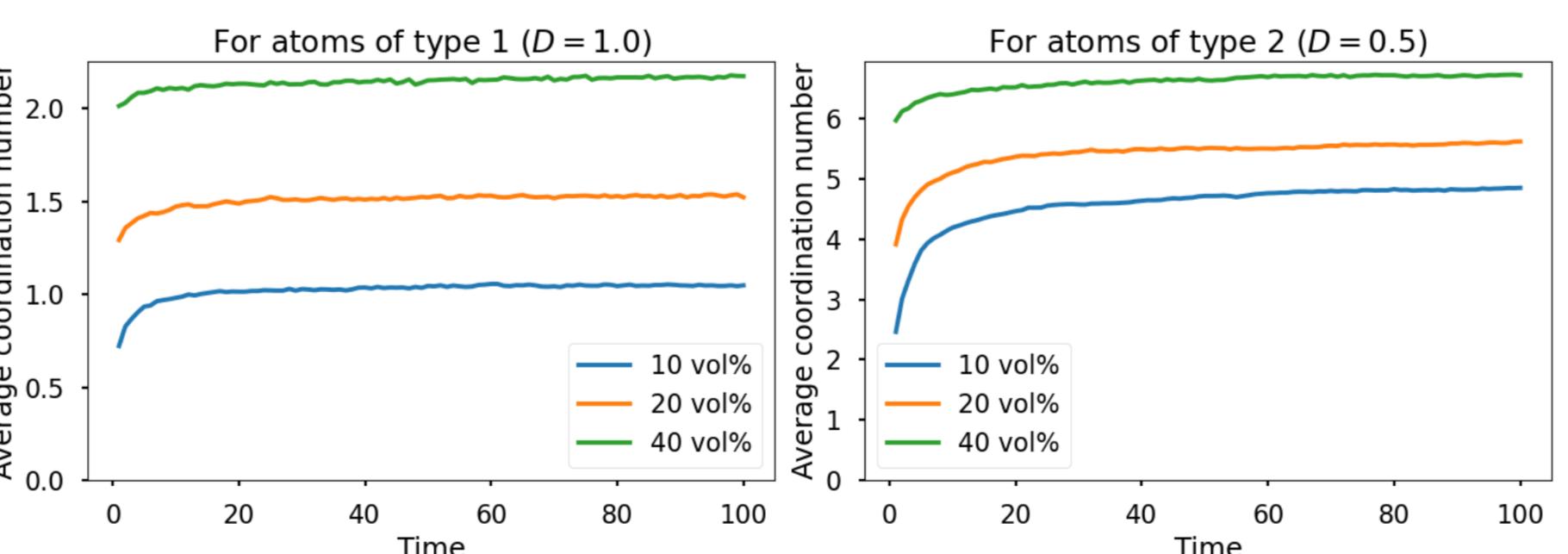


Fig. 4. Evolution of average coordination number in a bidisperse gel, with the j cutoff defined as $(D_i + D_j)/2 + D_i/4$ for particle i , so that particles whose surfaces are within $0.25D_i$ of particle i 's surface are considered neighbors. The simulation includes 1000 particles (500 of each size).

- At the same total volume fraction, atoms with $D = 1$ approach a much lower coordination number when smaller atoms are present (also observed for a surface-to-surface cutoff of 0; see “volumefrac_neighbors” folder).
- At a $0.25D$ surface-to-surface neighbor cutoff, monodisperse gels with $D = 1$ approach similar neighbor counts as particles with $D = 0.5$ in the bidisperse gels.

Atom-level forces

Force over time was extracted for individual atoms using “forces.py.” The normalized autocorrelation function R (defined below) was computed in “plotforces.py.”

$$R(F, t_{\text{lag}}) = \frac{\langle F(t)F(t_{\text{lag}}) \rangle}{\langle F(t)^2 \rangle}$$

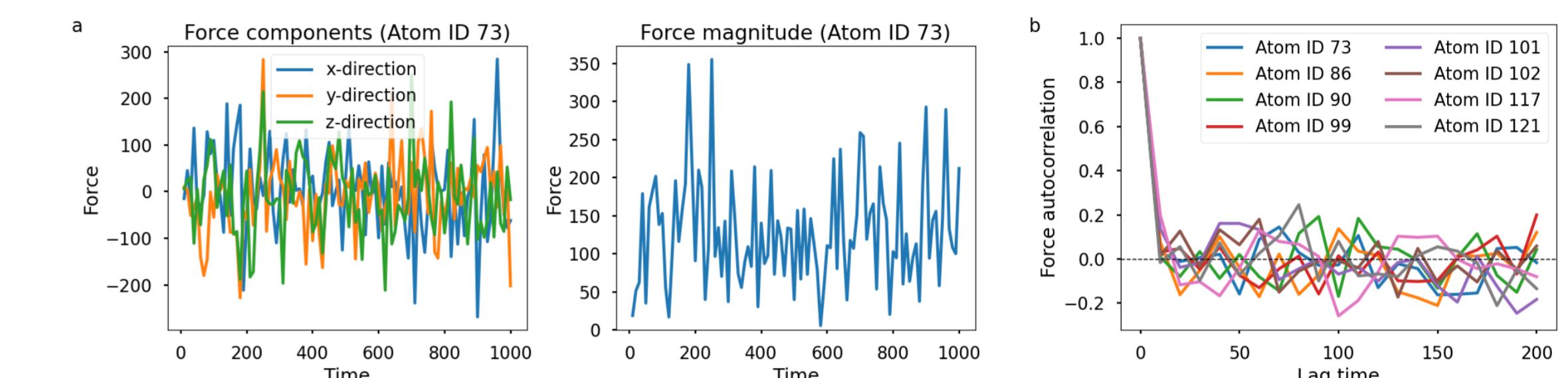


Fig. 5. (a) Force components and magnitude over time for an arbitrarily selected type 1 ($D = 1$) atom in a bidisperse gel with 10 vol% colloid, in the absence of `dpd/tstat`. (b) Autocorrelation function of the force magnitude for a selection of $D = 1$ atoms.

Even when random energy dissipation is excluded, memory of the current force (measured by autocorrelation) decays quickly at the level of an individual atom.

Conclusions

In a simple model with interacting spheres, post-gelation coordination number depends on particle size distribution at constant colloid volume fraction.

Future work

Validate the DPD thermostat by including a specific solvent explicitly, and determine the appropriate parameters to enable comparison to experiment.

Vary the temperature and look for changes in the force memory of atoms.

GitHub repository

This work was conducted as a final project for PHYS 25000 (Autumn 2025). The simulation code can be found at <https://github.com/msong2026/lammps-gels>. Sample LAMMPS scripts by Christopher Ness [4] were also consulted.

Acknowledgments

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References

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