

# The discovery of a five-image lensed quasar at $z = 3.34$ using PanSTARRS1 and *Gaia*

Fernanda Ostrovski,<sup>1,2★</sup> Cameron A. Lemon,<sup>2,3★</sup> Matthew W. Auger,<sup>2</sup>  
Richard G. McMahon,<sup>2,3</sup> Christopher D. Fassnacht,<sup>4</sup> Geoff C.-F. Chen,<sup>4,5</sup>  
Andrew J. Connolly,<sup>6</sup> Sergey E. Koposov,<sup>2,7</sup> Estelle Pons,<sup>2,3</sup> Sophie L. Reed<sup>2,3</sup>  
and Cristian E. Rusu<sup>4,8</sup>

<sup>1</sup>Departamento de Astronomia, Instituto de Física da Universidade Federal do Rio Grande do Sul, 91501-970 Porto Alegre, Brazil

<sup>2</sup>Institute of Astronomy, University of Cambridge, Madingley Road, Cambridge CB3 0HA, UK

<sup>3</sup>Kavli Institute for Cosmology, University of Cambridge, Madingley Road, Cambridge CB3 0HA, UK

<sup>4</sup>Department of Physics, University of California Davis, 1 Shields Avenue, Davis, CA 95616, USA

<sup>5</sup>Institute of Astronomy and Astrophysics, Academia Sinica, PO Box 23-141, Taipei 10617, Taiwan

<sup>6</sup>Department of Astronomy, University of Washington, Seattle, WA 98195, USA

<sup>7</sup>Department of Physics, McWilliams Center for Cosmology, Carnegie Mellon University, 5000 Forbes Avenue, Pittsburgh, PA 15213, USA

<sup>8</sup>Subaru Telescope, National Astronomical Observatory of Japan, 650 N Aohoku Pl, Hilo, HI 96720, USA

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## ABSTRACT

We report the discovery, spectroscopic confirmation and mass modelling of the gravitationally lensed quasar system PS J0630–1201. The lens was discovered by matching a photometric quasar catalogue compiled from Pan-STARRS1 and *Wide-field Infrared Survey Explorer* photometry to the *Gaia* data release 1 catalogue, exploiting the high spatial resolution of the latter (full width at half-maximum  $\sim 0.1$  arcsec) to identify the three brightest components of the lensed quasar system. Follow-up spectroscopic observations with the William Herschel Telescope confirm the multiple objects are quasars at redshift  $z_q = 3.34$ . Further follow-up with Keck adaptive optics high-resolution imaging reveals that the system is composed of two lensing galaxies and the quasar is lensed into an  $\sim 2.8$  arcsec separation four-image cusp configuration with a fifth image clearly visible, and a  $1.0$  arcsec arc due to the lensed quasar host galaxy. The system is well modelled with two singular isothermal ellipsoids, reproducing the position of the fifth image. We discuss future prospects for measuring time delays between the images and constraining any offset between mass and light using the faintly detected *Einstein* arcs associated with the quasar host galaxy.

**Key words:** gravitational lensing: strong – methods: observational – methods: statistical – quasars: general.

## 1 INTRODUCTION

Gravitationally lensed quasars can be used as tools for a variety of astrophysical and cosmological applications (e.g. Treu 2010; Jackson 2013), including: mapping the dark matter substructure of the lensing galaxy (e.g. Metcalf & Madau 2001; Kochanek & Dalal 2004; Nierenberg et al. 2014; Vegetti et al. 2014; Birrer, Amara & Refregier 2017); determining the mass (e.g. Morgan et al. 2010) and spin (Reynolds et al. 2014) of black holes and measuring the properties of distant host galaxies (e.g. Kochanek, Keeton & McLeod 2001; Claeskens et al. 2006; Peng et al. 2006;

Ding et al. 2017). In addition, they can also be used to constrain cosmological parameters with comparable precision to baryonic acoustic oscillation methods (e.g. Suyu et al. 2010, 2013) and to probe the physical properties of quasar accretion discs through microlensing studies (e.g. Kochanek 2004; Poindexter, Morgan & Kochanek 2008; Motta et al. 2012).

Ostrovski et al. (2017) found that, for the Dark Energy Survey (DES – DES Collaboration 2005; DES Collaboration et al. 2016), no simulated gravitationally lensed quasar system with image separation less than  $1.5$  arcsec is segmented into multiple catalogue sources due to limitations on survey resolution. As a result, only 23 per cent of simulated systems with pairs of quasar images are segmented into two sources. This means that to identify lensed quasars as groups of close sources of similar colours, one has to either

\* E-mail: fostrovski@ast.cam.ac.uk (FO); cl522@ast.cam.ac.uk (CAL)

employ 2D modelling techniques (Ostrovski et al., in preparation) or rely on cross-matching to higher resolution imaging surveys such as *Gaia* (full width at half-maximum  $\sim 0.1$  arcsec; Lemon et al. 2017).

*Gaia* (Gaia Collaboration et al. 2016a) is a space observatory performing a full sky survey to study the Milky Way. *Gaia* data release 1 (DR1 – Lindegren et al. 2016) contains positions for a total of  $1.1 \times 10^9$  sources across the whole sky with a point source limiting magnitude of  $G = 20.7$  (Gaia Collaboration et al. 2016b).

Here, we present the discovery of the lensed quasar PS J0630–1201 from a preliminary search for gravitationally lensed quasars from Pan-STARRS1 (PS – Chambers et al. 2016) combined with the *Wide-field Infrared Survey Explorer* (*WISE* – Wright et al. 2010) by cross-matching to *Gaia* DR1 detections. All magnitudes are quoted on the AB system. Conversions from Vega to AB for the *WISE* data are  $W1_{AB} = W1_{Vega} + 2.699$  and  $W2_{AB} = W2_{Vega} + 3.339$  (Jarrett et al. 2011).

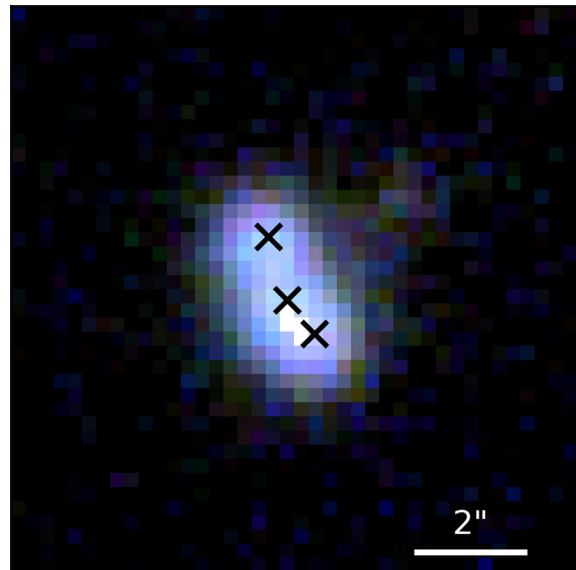
## 2 LENS DISCOVERY AND CONFIRMATION

To select gravitationally lensed quasar candidates, we first create a *morphology-independent* (i.e. not restricted to objects consistent with the point spread function – PSF) photometric quasar catalogue using Gaussian mixture models, a supervised machine learning method of generative classification, as described in Ostrovski et al. (2017). We use *WISE* and PS photometry to define  $g - i$ ,  $i - W2$ ,  $z - W1$  and  $W1 - W2$  colours and we employ three model classes (point sources, extended sources and quasars) to generate a catalogue of 378 061 quasar candidates from a parent sample of 80 028 181 objects with  $i < 20$ . To remove stellar contaminants at low galactic latitude, we apply conservative colour cuts based on the comparisons between the known quasar distribution and the distribution of our candidates as well as removal of objects with high point source probabilities. We also discard objects with a neighbouring candidate within 5 arcsec yielding a total of 296 967 objects, since the spatial resolution of PS is enough to flag these systems as potential lensed quasar candidates or binary pairs.

We then exploit the excellent spatial resolution and all-sky coverage of the *Gaia* DR1 catalogue to identify lensed quasar candidates by resolving the photometric quasars into multiple components. *Gaia* DR1 is known to be enormously incomplete for close-separation objects but nevertheless can identify multiple components of lensed quasars (see Lemon et al. 2017 for details). We cross-match our photometric quasar candidates with the *Gaia* DR1 catalogue using a 3 arcsec search radius and find that 1401 of the quasar candidates have two *Gaia* objects associated with their PS position, whilst 26 of the candidates have three *Gaia* associations. Visual inspection of these 26 objects revealed one lens candidate, PS J0630–1201, shown in Fig. 1 with *Gaia* positions overlaid on the PS cut-out, to be of interest. Its PS catalogue position, and PS and *WISE* photometry are listed in Table 1. The other objects with triple *Gaia* matches were ruled out as common contaminants, mainly single quasars with other objects nearby (that, despite being resolved in PS images, were not quasar candidates in our catalogue), as well as apparently faint interacting starbursting galaxies and likely duplicate entries in the *Gaia* catalogue (objects with separations of  $< 0.1$  arcsec).

### 2.1 Spectroscopic follow-up at the WHT

Spectroscopic follow-up observations to confirm that the multiple components are multiply imaged quasars were performed with the



**Figure 1.** PS J0630–1201 as a  $g, r, i$  PS colour composite. The single band stacked images were obtained from the PS cut-out server (<http://ps1images.stsci.edu/cgi-bin/ps1cutouts>). The black crosses mark the positions of overlapping objects from the *Gaia* catalogue. North is up and east is to the left.

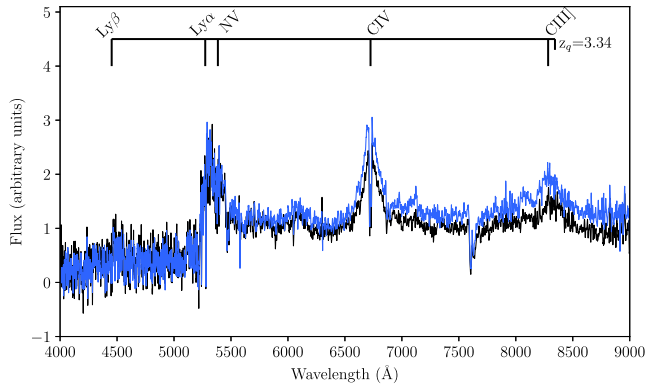
**Table 1.** PS J0630–1201 PanSTARRS and *WISE* selection photometry.

RA	06 <sup>h</sup> 30 <sup>m</sup> 09 <sup>s</sup> .11
Dec.	−12 <sup>d</sup> 01 <sup>m</sup> 20 <sup>s</sup> .0
$g$	$19.291 \pm 0.002$
$r$	$18.566 \pm 0.002$
$i$	$18.321 \pm 0.001$
$z$	$17.949 \pm 0.001$
$Y$	$17.873 \pm 0.001$
$W1$	$16.886 \pm 0.027$
$W2$	$16.643 \pm 0.030$

dual-arm Intermediate dispersion Spectrograph and Imaging System (ISIS) on the William Herschel Telescope (WHT) on the night of 2017 April 01 UT. The R300B grating was employed on the blue arm, providing wavelength coverage from  $\sim 3200$  to  $\sim 5400$  Å with  $\sim 4$  Å resolution, and the R158R grating on the red arm resulted in coverage from  $\sim 5300$  to  $\sim 10\,200$  Å with a resolution of  $\sim 7.7$  Å; the 5300 Å dichroic was used to split the beam to the blue and red arms. Two exposures of 600 s each were obtained with a slit position angle (PA) of  $22^\circ.5$ , i.e. along the direction of the three brightest images, and two clearly separated traces were visible on the red arm, with the two southern-most images not deblended, whilst no clear separation was apparent in the blue data. Extractions of these traces are shown in Fig. 2, revealing very similar spectra indicating a lensed quasar at  $z_q = 3.34$ . No features of a lower redshift lensing galaxy are visible in the spectra.

### 2.2 AO follow-up with Keck

After the spectroscopic confirmation of multiple quasars, the system was observed on 2017 April 11 UT with the NIRC2 camera mounted on Keck II, using Keck’s adaptive optics (AO) system. The NIRC2 narrow camera was used, giving a  $10$  arcsec  $\times$   $10$  arcsec field of view



**Figure 2.** WHT/ISIS spectra of component C (black) and the blended components A and B (blue). On top, we have marked the location of the source quasar emission lines at  $z_q = 3.34$ .

and 10 mas pixels, and four 180 s exposures were obtained with the  $K'$  filter. These data clearly resolve the three quasars observed with WHT spectroscopy (A, B and C in Fig. 3), and also reveal two additional point-like objects (D and E) and two extended objects (G1 and G2). Note that most of the structure around the bright images is an artefact (‘waffling’) due to AO correction problems with the low-bandwidth wavefront sensor.

The PSF of image C appears to have a structure extending down from the core of the PSF that is not seen in images A or B but could be consistent with a lensed arc. We therefore produced a pixellated model of the PSF around the images ABC to remove it and increase the dynamic range of the image. In the first iteration of this procedure, an arc between images B and C was clearly visible, and we therefore re-fitted for the PSF excluding pixels that contain arc flux. The residuals of this fit are shown in the middle and right-hand panels of Fig. 3.

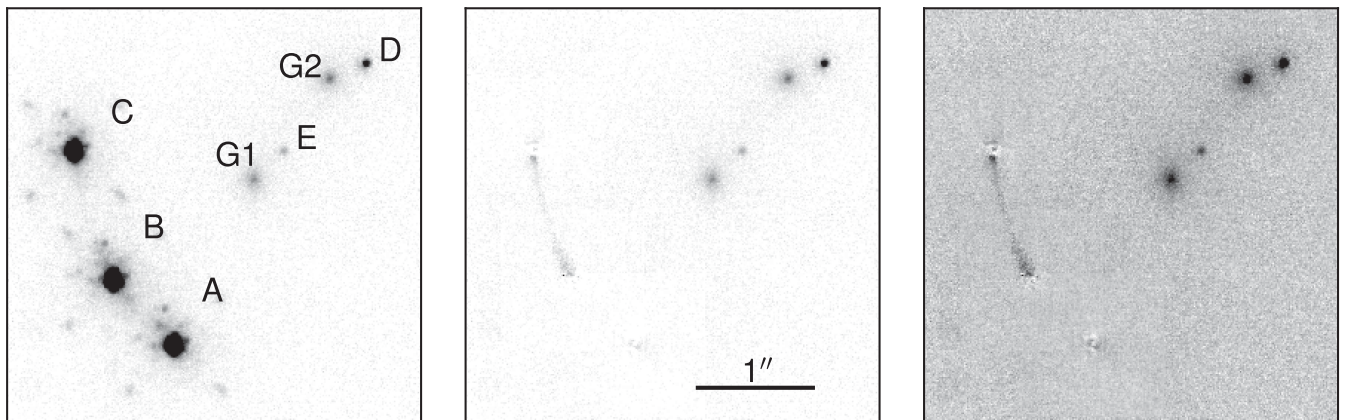
### 3 LENS MASS MODELLING

The AO imaging data show that PS J0630–1201 is a lensed quasar in a ‘cusp’ configuration, with two lensing galaxies. The presence of a fifth point source, E in Fig. 3, is intriguing as it is located approximately where the fifth (demagnified) image would be expected to appear. To understand the nature of E as a possible fifth quasar

image, we initially create a mass model using the positions of the brighter quasar images (A–D) and the two lensing galaxies, and use this model to predict the location of any additional images. We first determine the position of each point source by modelling them with Gaussian or Moffat profiles, and we fit Sérsic profiles to G1 and G2. We also perform moment-based centroiding in a range of apertures, and use the spread of all of our measurements to estimate the uncertainties on the positions; these are typically  $\sim 1$  mas for the point sources and  $\sim 10$ – $20$  mas for the galaxies. However, we also find that the relative positions of the quasar images change from exposure to exposure, presumably due to atmospheric fluctuations, and we therefore impose a 5 mas minimum uncertainty on each position. We simultaneously fit the PS imaging data (*grizY*) whilst fitting the Keck data ( $K'$ ), and our photometric and astrometric results are given in Table 2.

We use the positions of the galaxies and point sources to constrain a lensing mass model. The two galaxies are initially modelled as singular isothermal spheres and we use the positions of the four brighter images to infer their *Einstein* radii using the LENS-MODEL software (Keeton 2001). The best-fitting lens model predicts a fifth image near the location of image E with a flux approximately half of image D, comparable to the observed flux ratio. We consequently use the position of E to constrain a more realistic lens model, allowing the two galaxies to have some ellipticity, and we include an external shear. This model has five free parameters for each galaxy (two position parameters, an *Einstein* radius and an ellipticity with its orientation) and two additional parameters each for the position of the source and the external shear, i.e. there are 14 total parameters. We likewise have 14 constraints from the observed positions of the five quasar images and the two lensing galaxies.

Using image plane modelling, we sample the parameter space using EMCEE (Foreman-Mackey et al. 2013) and, as expected, find a best-fitting model with  $\chi^2 \sim 0$ . Both mass components are inferred to be coincident with the light, with *Einstein* radii for G1 and G2 of  $1.01 \text{ arcsec} \pm 0.01$  and  $0.58 \text{ arcsec} \pm 0.01$ , respectively. Both masses are also mildly flattened, but we find that the light and mass *orientations*, given in Table 3, are misaligned, significantly so for G2. We also note that the shear is well constrained ( $\gamma = 0.14 \pm 0.01$  with PA  $-76^\circ$ ) and is particularly large, though there is no nearby galaxy along the shear direction. The total magnification for the system is  $\sim 53$ .



**Figure 3.** Left: NIRC2  $K'$  AO imaging of PS J0630–1201. The four images ABCD are in a canonical ‘cusp’ configuration, but note the presence of two galaxies, G1 and G2, as well as the additional point source E. The additional structure around the images A, B and C is due to poor wavefront correction. Middle: the same as the left-hand panel, but with a model for the AO PSF subtracted from the three brightest images, revealing the presence of a faint arc. Right: same as the middle but with enhanced contrast to better show the lensed quasar host galaxy.



**Table 2.** Relative astrometry and photometry of lens components. *grizYK'* of D and E are combined with G2 and G1, respectively, due to blending.

	$G^a$	$g$	$r$	$i$	$z$	$Y$	$K'$	$\Delta\alpha$	$\Delta\delta$
A	$19.95 \pm 0.01$	$20.65 \pm 0.05$	$19.92 \pm 0.04$	$19.65 \pm 0.04$	$19.27 \pm 0.01$	$19.22 \pm 0.05$	$18.65 \pm 0.01$	$0.674 \pm 0.005$	$-1.421 \pm 0.005$
B	$19.76 \pm 0.01$	$20.54 \pm 0.05$	$19.73 \pm 0.05$	$19.80 \pm 0.03$	$19.18 \pm 0.02$	$19.23 \pm 0.04$	$18.49 \pm 0.01$	$1.192 \pm 0.005$	$-0.865 \pm 0.005$
C	$19.61 \pm 0.01$	$20.53 \pm 0.05$	$19.70 \pm 0.05$	$19.66 \pm 0.04$	$19.24 \pm 0.02$	$19.15 \pm 0.04$	$18.66 \pm 0.01$	$1.530 \pm 0.005$	$0.248 \pm 0.005$
D(+G2)	–	$23.03 \pm 0.16$	$21.80 \pm 0.05$	$21.59 \pm 0.04$	$21.32 \pm 0.04$	$20.85 \pm 0.03$	$21.30 \pm 0.01$	$-0.966 \pm 0.005$	$0.994 \pm 0.005$
E(+G1)	–	$23.33 \pm 0.34$	$23.16 \pm 0.14$	$21.93 \pm 0.09$	$21.83 \pm 0.09$	$21.60 \pm 0.09$	$22.76 \pm 0.01$	$-0.257 \pm 0.005$	$0.241 \pm 0.005$
G1	–	–	–	–	–	–	$20.86 \pm 0.01$	$0.000 \pm 0.013$	$0.000 \pm 0.010$
G2	–	–	–	–	–	–	$20.93 \pm 0.01$	$-0.650 \pm 0.017$	$0.862 \pm 0.005$

<sup>a</sup>Gaia DR1 catalogue  $G$ -band photometry for components A,B and C.**Table 3.** Light and mass shape parameters (position angle, PA, and axis ratio,  $q$ ) for the lens galaxies.

	PA (light)	$q$ (light)	PA (mass)	$q$ (mass)
G1	$-9 \pm 4$	$0.81 \pm 0.03$	$-21 \pm 6$	$0.81 \pm 0.02$
G2	$-69 \pm 4$	$0.79 \pm 0.03$	$-2 \pm 7$	$0.83 \pm 0.04$

**Table 4.** Predicted time delays between the image pairs in PS J0630–1201. All values are in days. The values in the top right (bottom left) are for a lens redshift of 1 (0.5). Light arrives in the images in the following order: CABED.

	A	B	C	D	E
A	–	0.9	1.8	243	208
B	0.4	–	2.7	242	207
C	0.7	1.1	–	245	210
D	97.8	97.4	98.5	–	35
E	83.7	83.3	84.4	14.1	–

## 4 DISCUSSION AND CONCLUSIONS

We have presented spectroscopic and imaging data that confirm that PS J0630–1201 is a quasar at  $z_q = 3.34$  lensed into five images by two lensing galaxies. We are able to fit the positions of the five images well with a two-SIE lens model and recover flux ratios to within 30 per cent with discrepancies likely caused by microlensing and/or differential extinction and reddening, as evidenced by the strongly varying flux ratio between images B and C from optical to near-infrared wavelengths. However, we find that for both lenses the ellipticity of the mass is not consistent with the ellipticity of the galaxy light (Table 3) and the inferred shear is quite large. This could be the result of an additional mass component, e.g. a dark matter halo that is not coincident with either galaxy (e.g. Shu et al. 2016). In that case, the weak demagnification of the fifth image might indicate that the dark matter halo is not cuspy (e.g. Collett et al. 2017), although constraints from the quasar image positions alone are not sufficient to test this. Deeper imaging of the arc of the lensed host galaxy and observations at radio wavelengths, where extinction and the effects of microlensing are no longer important, will help to constrain a more complex model for the mass distribution.

The relatively bright fifth image of PS J0630–1201 also presents the possibility of obtaining four new time delay measurements, for a total of 10 time delays. Based upon our current best lens model, these delays should range between 1 and 245 d assuming that the lens redshifts are  $z \sim 1$  (Table 4). Because of the overall compactness of the system and the presence of the two lensing galaxies, it would be difficult to obtain time delays from the fifth image with conventional seeing-limited monitoring programmes. However, if such a campaign observed a sudden brightening (or

dimming) event in one of the brighter images, dedicated monitoring with a high-resolution facility (e.g. Robo-AO; Baranec et al. 2014) could yield an observation of the delayed brightening of the fifth image.

Several other lensed quasars have been observed to have bright additional images, and these systems typically also have multiple lensing galaxies (with the exception being the radio-loud double PMN J1632–0033, which has an observed highly de-magnified central image; Winn, Rusin & Kochanek 2004). The cluster lens SDSS J1004+4112 (Inada et al. 2005) is a quad with an observed ‘central’ image, but the complexity of the mass distribution makes it very difficult to estimate time delays (Fohlmeister et al. 2007). PMN J0134–0931 is a radio-loud quasar that is being lensed into five images by two galaxies (Winn et al. 2003). However, the image separations are very small so measuring time delays – expected to range from an hour to two weeks (Keeton & Winn 2003) – will be extremely difficult. B1359+154 (Rusin et al. 2001) and SDSS J2222+2745 (Dahle et al. 2013) are six-image lens systems, and in both cases the lens is a compact group of three galaxies. These lenses have intrinsically more complex mass distributions, but time delays may still be informative for cosmography. However, most of the independent time delays are expected to be less than a day for B1359+154 (Rusin et al. 2001) and quite long (400–700 d) for SDSS J2222+2745 (Sharon et al. 2017). PS J0630–1201 therefore appears to be the most promising lens for measuring additional independent time delays.

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