Partial Differerential Equations Thành Nam Phan Winter Semester 2021/2022

Lecture notes T_EXed by Thomas Eingartner

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Please note that I write this lecture notes for my personal use. I may write things differently than presented in the lecture. This script also contains solutions for exercises (which may be wrong). Of course, I don't push them to GitHub while the exercises still can be handed in.

Chapter 1

Introduction

A differential equation is an equation of a function and its derivatives.

Example 1.1 (Linear ODE) Let $f: \mathbb{R} \to \mathbb{R}$,

$$\begin{cases} f(t) = af(t) \text{ for all } t \geqslant 0, a \in \mathbb{R} \\ f(0) = a_0 \end{cases}$$

is a linear ODE (Ordinary differential equation). The solution is: $f(t) = a_0 e^{at}$ for all $t \ge 0$.

Example 1.2 (Non-Linear ODE) $f: \mathbb{R} \to \mathbb{R}$

$$\begin{cases} f'(t) = 1 + f^2(t) \\ f(0) = 1 \end{cases}$$

Lets consider $f(t) = \tan(t) = \frac{\sin(t)}{\cos(t)}$. Then we have

$$f'(t) = \frac{1}{\cos(t)} = 1 + \tan^2(t) = 1 + f^2(t),$$

but this solution only is good in $(-\pi, \pi)$. It's a problem to extend this to $\mathbb{R} \to \mathbb{R}$.

A PDE (Partial Differential Equation) is an equation of a function of 2 or more variables and its derivatives.

Remark 1.3 Recall for $\Omega \subseteq \mathbb{R}^d$ open and $f: \Omega \to \{\mathbb{R}, \mathbb{C}\}$ the notation of partial derivatives:

- $\partial_{x_i} f(x) = \lim_{h \to 0} \frac{f(x+he_i) f(x)}{h}$, where $e_i = (0, 0, \dots, 1, \dots, 0, 0) \in \mathbb{R}^d$
- $D^{\alpha}f(x) = \partial_{x_1}^{\alpha_1} \cdots \partial_{x_d}^{\alpha_d} f(x)$, where $\alpha \in \mathbb{N}^d$
- $Df = \nabla f = (\partial x_1, \dots, \partial_{x_d})$
- $\bullet \ \Delta f = \partial_{x_1}^2 + \dots + \partial_{x_d}^2 f$
- $D^k f = (D^{\alpha} f)_{|\alpha|=k}$, where $|\alpha| = \sum_{i=1}^d |\alpha_i|$
- $D^2 f = (\partial_{x_i} \partial_{x_j} f)_{1 \leq i, j \leq d}$

Definition 1.4 Given a function F. Then the equation of the form

$$F(D^k u(x), D^{k-1} u(x), \dots, Du(x), u(x), x) = 0$$

with the unknown function $u: \Omega \subseteq \mathbb{R}^d \longrightarrow \mathbb{R}$ is called a *PDE of order k*.

- Equations $\sum_{d} a_{\alpha}(x) D^{\alpha} u(x) = 0$, where a_{α} and u are unknown functions are called *Linear PDEs*.
- Equations $\sum_{|\alpha|=k} a_{\alpha}(x) D^{\alpha} u(x) + F(D^{k-1}u, D^{k-2}u, \dots, Du, u, x) = 0$ are called semi-linear PDEs.

Goals: For solving a PDE we want to

- Find an explizit solution! This is in many cases impossible.
- Prove a well-posted theory (existence of solutions, uniqueness of solutions, continuous dependence of solutions on the data)

We have two notations of solutions:

- 1. Classical solution: The solution is continuous differentiable (e.g. $\Delta u = f \leadsto u \in C^2$)
- 2. Weak Solutions: The solution is not smooth/continuous

Definition 1.5 (Spaces of continous and differentiable functions) Let $\Omega \subseteq \mathbb{R}^d$ be open

$$\begin{split} C(\Omega) &= \{f: \ \Omega \to \mathbb{R} \mid f \text{ continuous} \} \\ C^k(\Omega) &= \{f: \ \Omega \to \mathbb{R} \mid D^\alpha f \text{ is continuous for all } |\alpha| \ \leqslant k \} \end{split}$$

Classical solution of a PDE of order $k \rightsquigarrow C^k$ solutions!

$$L^p(\Omega) = \left\{ f: \ \Omega \to \mathbb{R} \text{ lebesgue measurable } | \int_{\Omega} |f|^p d\lambda < \infty, 1 \leqslant p < \infty \right\}$$

Sobolev Space:

$$W^{k,p}(\Omega) = \{ f \in L^p(\Omega) \mid \forall \alpha \in \mathbb{N}^n \text{ with } |\alpha| \leq k : D^{\alpha} f \in L^p(\Omega) \text{ exists} \}$$

In this course we will investigate

- Laplace / Poisson Equation: $-\Delta u = f$
- Heat Equation: $\partial_t u \Delta u = f$
- Wave Equation: $\partial_t^2 \Delta u = f$
- Schrödinger Equation: $i\partial_t u \Delta u = f$

Chapter 2

Laplace / Poisson Equation

2.1 Laplace Equation

 $-\Delta u = 0$ (Laplace) or $-\Delta u = f(x)$ (Poisson).

Definition 2.1 (Harmonic Function) Let Ω be an open set in \mathbb{R}^d . If $u \in C^2(\Omega)$ and $\Delta u = 0$ in Ω , then u is a harmonic function in Ω .

Theorem 2.2 (Gauss-Green Theorem) Let $A \subseteq \mathbb{R}^d$ open, $\vec{F} \in C^1(A, \mathbb{R}^d)$ and $K \subseteq A$ compact with C^1 boundary. Then

$$\int_{\partial K} \vec{F} \cdot \vec{\nu} \ dS(x) = \int_K \operatorname{div}(\vec{F}) \ dx$$

where ν is the outward unit normal vector field on ∂K . Thus

$$\int_{\partial V} \nabla u \cdot \vec{\nu} \ dS(x) = \int_{V} \operatorname{div}(\nabla u) \ dx = \int_{V} \Delta u(x) \ dx$$

for any $V \subseteq \Omega$ open.

Theorem 2.3 (Green's Identities) Let $A \subseteq \mathbb{R}^d$ open, $K \subseteq A$ d-dim. compactum with C^1 boundary and $f, g \in C^2(A)$

1. Green's first identity (Partial Integration):

$$\int_{K} \nabla f \cdot \nabla g \, dx = \int_{\partial K} f \frac{\partial g}{\partial \nu} \, dS - \int_{K} f \Delta g \, dx$$

where $\frac{\partial g}{\partial \nu} = \partial_{\nu} g = \nu \cdot \nabla g$

2. Green's second identity:

$$\int_{K} f \Delta g - (\Delta f) g \, dx = \int_{\partial K} \left(f \frac{\partial g}{\partial \nu} - g \frac{\partial f}{\partial \nu} \right) \, dS$$

Exercise 2.4 Let $\Omega \subseteq \mathbb{R}^d$ open, let $f:\Omega \to \mathbb{R}$ be continuous. Prove that if $\int_B f(x) \ dx = 0$, then $u \equiv 0$ in Ω .

Theorem 2.5 (Fundamential Lemma of Calculus of Variations) Let $\Omega \subseteq \mathbb{R}^d$ open, let $f \in L^1(\Omega)$. If $\int_B f(x) \ dx = 0$ for all $x \in B_r(x) \subseteq \Omega$, then f(x) = 0 a.e. (almost everywhere) $x \in \Omega$.

Remark 2.6 (Solving Laplace Equation) $-\Delta u = 0$ in \mathbb{R}^d . Consider the case when u is radial, i.e. $u(x) = v(|x|), v : \mathbb{R} \to \mathbb{R}$. Denote r = |x|, then

$$\frac{\partial r}{\partial x} = \frac{\partial}{\partial x_i} \left(\sqrt{x_1^2 + \dots + x_d^2} \right) = \frac{2x_i}{2\sqrt{x_1^2 + \dots + x_d^2}} = \frac{x_i}{r}$$

Then

$$\begin{split} \partial_{x_i} u &= \partial_{x_i} v = (\partial_r v) \frac{\partial r}{\partial x_i} = v'(r) \frac{x_i}{r} \\ \partial_{x_i}^2 u &= \partial_{x_i} \left(v(r)' \frac{x_i}{r} \right) = (\partial_{x_i} v(r)') \frac{x_i}{r} + v'(r) \partial_{x_i} \left(\frac{x_i}{r} \right) \\ &= (\partial_r v'(r)) \left(\frac{dr}{\partial_{x_i}} \right) \frac{x_i}{r} + v'(r) \left(\frac{1}{r} - \frac{x_i}{r^2} (\partial_{x_i} r) \right) = v'(r) \frac{x_i^2}{r^2} + v'r(r) \left(\frac{1}{r} - \frac{x_i^2}{r^3} \right) \end{split}$$

So we have $\Delta u = \left(\sum_{i=1}^d d_{x_i}^2\right) u = v''(r) + v'(r)(\frac{d}{r} - \frac{1}{r})$ Thus $\Delta u = v'(r) + v(r)\frac{d-1}{r}$. We consider $d \ge 2$. Laplace operator $\Delta u = 0$ now becomes $v''(r) + v'(r)\frac{d-1}{r} = 0$

$$\Rightarrow \log(v(r))' = \frac{v'(r)}{v(r)} = -\frac{d-1}{r} = -(d-1)(\log r)' \text{ (recall } \log(f)' = \frac{f'}{f})$$

$$\Rightarrow v'(r) = \frac{1}{v^{d-2} + \text{const.}}$$

$$\begin{cases} \frac{const}{r^{d-2}} + constxx + const & , d \geqslant 3 \\ const \log(r) + constxxr + const & , d = 2 \end{cases}$$

Definition 2.7 (Fundamential Solution of Laplace Equation)

$$\Phi(x) = \begin{cases} -\frac{1}{2\pi} \log(|x|), & d = 2\\ \frac{1}{(d-2)d|B_1|} \frac{1}{|x|^{d-2}}, & d \geqslant 3 \end{cases}$$

Where $|B_1|$ is the Volume of the ball $B_1(0) = B(0,1) \subseteq \mathbb{R}^d$.

Remark 2.8 $\Delta\Phi(x) = 0$ for all $x \in \mathbb{R}^d$ and $x \neq 0$.

2.2Poisson-Equation

The Poisson-Equation is $-\Delta u(x) = f(x)$ in \mathbb{R}^d . The explicit solution is given by

$$u(x) = (\Phi \star f)(x) = \int_{\mathbb{R}^d} \Phi(x - y) f(y) \ dy = \int_{\mathbb{R}^d} \Phi(y) f(x - y) \ dy$$

This can be heuristically justifyfied with

$$-\Delta(\Phi \star f) = (-\Delta\Phi) \star f = \delta_0 \star f = f$$

Theorem 2.9 Assume $f \in C_c^2(\mathbb{R}^d)$. Then $u = \Phi \star f$ satisfies that $u \in C^2(\mathbb{R}^d)$ and $-\Delta u(x) = f(x)$ for all $x \in \mathbb{R}^d$

Proof. By definition we have

$$u(x) = \int_{\mathbb{R}^d} \Phi(y) f(x - y) \, dy.$$

First we check that u is continuous: Take $x_k \to x_0$ in \mathbb{R}^d . We prove that $u(x_n) \xrightarrow{n} u_0$, i.e.

$$\lim_{n \to \infty} \int_{\mathbb{R}^d} \Phi(y) f(x_n - y) \ dy = \int_{\mathbb{R}^d} \Phi(y) f(x_0 - y) \ dy$$

This follows from the Dominated Convergence Theorem. More precisely:

$$\lim_{n \to \infty} \Phi(y) f(x_n - y) = \Phi(y) f(x_0 - y) \quad \text{for all } y \in \mathbb{R}^d \setminus \{0\}$$

and

$$|\Phi(y)f(x-y)| \leq ||f||_{L^{\infty}} \cdot \mathbb{1}(|y| \leq R) \cdot |\Phi(y)| \in L^1(\mathbb{R}^d, dy)$$

where R > 0 depends on $\{x_n\}$ and supp(f) but independent of y. Now we compute the derivatives:

$$\begin{split} \partial_{x_i} u(x) &= \partial_{x_i} \int_{\mathbb{R}^d} \Phi(y) f(x-y) \ dy = \lim_{h \to 0} \int_{\mathbb{R}^d} \Phi(y) \frac{f(x+he_i-y) - f(x-y)}{h} \ dy \\ (\text{dom. conv.}) &= \int \Phi(y) \partial_{x_i} f(x-y) \ dy \\ \Rightarrow & D^{\alpha} u(x) = \int_{\mathbb{R}^d} \Phi(y) D_x^{\alpha} f(x-y) \ dy \quad \text{for all } |\alpha| \leqslant 2 \end{split}$$

 $D^{\alpha}u(x)$ is continuous, thus $u\in C^2(\mathbb{R}^d)$. Now we check if this solves the Poisson-Equation:

$$-\Delta u(x) = \int_{\mathbb{R}^d} \Phi(y)(-\Delta_x) f(x-y) \, dy = \int_{\mathbb{R}^d} \Phi(y)(-\Delta_y) f(x-y) \, dy$$
$$= \int_{\mathbb{R}^d \setminus B(0,\epsilon)} \Phi(y)(-\Delta_x) f(x-y) \, dy + \int_{B(0,\epsilon)} \Phi(y)(-\Delta_x) f(x-y) \, dy \quad (\epsilon > 0 \text{ small})$$

Now we come to the main part. We apply integration by parts (2.3):

$$\int_{\mathbb{R}^{d} \backslash B(0,\epsilon)} \Phi(y)(-\Delta_{y}) f(x-y) \, dy$$

$$= \int_{\mathbb{R}^{d} \backslash B(0,\epsilon)} (\nabla_{y} \Phi(y)) \cdot \nabla_{y} f(x-y) \, dy - \int_{\partial B(0,\epsilon)} \Phi(y) \cdot \frac{\partial f}{\partial \vec{n}} (x-y) \, dS(y)$$

$$= \int_{\mathbb{R}^{d} \backslash B(0,\epsilon)} (-\Delta_{y} \Phi(y)) f(x-y) \, dy$$

$$+ \int_{\partial B(0,\epsilon)} \frac{\partial \Phi}{\partial \vec{n}} (y) f(x-y) \, dS(y) - \int_{\partial B(0,\epsilon)} \Phi(y) \frac{\partial f}{\partial \vec{n}} (x-y) \, dS(y)$$

We have that $\nabla_y \Phi(y) = -\frac{1}{d|B_1|} \frac{y}{|y|^d}$ and $\vec{n} = \frac{y}{|y|}$ in $\partial B(0, \epsilon)$. This leads to

$$\frac{\partial \Phi}{\partial \vec{n}} = \frac{1}{d|B_1|} \frac{1}{|y|^{d-1}} = \frac{1}{d|B_1|\epsilon^{d-1}} \quad \text{for } y \in \partial B(0, \epsilon)$$

Hence:

$$\int_{\partial B(0,\epsilon)} \frac{\partial \Phi}{\partial \vec{n}}(y) f(x-y) \ dS(y) = \frac{1}{d|B_1|\epsilon^{d-1}} \int_{\partial B(0,\epsilon)} f(x-y) \ dS(y)$$
$$= \int_{\partial B(0,\epsilon)} f(x-y) \ dS(y) = \int_{\partial B(x,\epsilon)} f(y) \ dS(y) \xrightarrow{\epsilon \to 0} f(x)$$

We have to regard the following error terms:

$$\left| \int_{B(0,\epsilon)} \Phi(y) (-\Delta_y) f(x-y) \, dy \right| \leq \int_{B(0,\epsilon)} |\Phi(y)| \underbrace{\left| -\Delta_y f(x-y) \right|}_{\leq \|\Delta f\|_{L^{\infty}} \mathbb{1}(|y| \leq R)} \, dy$$

$$\leq \|\Delta f\|_{L^{\infty}} \int_{\mathbb{R}^d} \underbrace{\left| \Phi(y) |\mathbb{1}(|y| \leq R) \right|}_{L^1(\mathbb{R}^d)} \mathbb{1}(|y| \leq \epsilon) \xrightarrow{\epsilon \to 0} 0$$

Where R > 0 depends on x and the support of f but is independent of y.

$$\bullet \left| \int_{\partial B(0,\epsilon)} \Phi(y) \frac{\partial f}{\partial \vec{n}}(x-y) \ dS(y) \right| \leq \|\nabla f\|_{L^{\infty}} \int_{\partial B(0,\epsilon)} |\Phi(y)| \ dy$$

$$\leq \begin{cases} const \cdot \epsilon |\log \epsilon| \to 0, & d = 2\\ const \cdot \epsilon \to 0, & d \geqslant 3 \end{cases}$$

Conclusion: $-\Delta u(x) = f(x)$ for all $x \in \mathbb{R}^d$ proved that $u = \Phi \star f$ and $f \in C^2_c(\mathbb{R}^d)$.

Thus, if $f \in C_c^2(\mathbb{R})$, then $u = \Phi \star f$ satisfies $u \in C^2(\mathbb{R}^2)$ and $-\Delta u(x) = f(x)$ for all $x \in \mathbb{R}^d$.

Remark 2.10 The result holds for a much bigger class of functions f. For example if $f \in C_c^1(\mathbb{R})$ we can easily extend the previous proof:

$$\partial_{x_i} u = \int_{\mathbb{R}^d} \Phi(y) \partial_{x_i} f(x - y) \, dy \in C(\mathbb{R}^d) \Rightarrow u \in C^1(\mathbb{R}^d)$$

Consequently:

$$\partial_{x_i}\partial_{x_j}u = \partial_{x_i}\int_{\mathbb{R}^d} \Phi(y)\partial_{x_j}f(x-y)\,dy = \int_{\mathbb{R}^d} \partial_{x_i}\Phi(y)\partial_{x_j}f(x-y)\,dy \in C(\mathbb{R}^d)$$

So we have $u \in C^2(\mathbb{R}^d)$. Now we can compute

$$\Delta u = \sum_{i=1}^{d} \int_{\mathbb{R}^d} \partial_{x_i} \Phi(y) \partial_{x_i} f(x-y) \, dy \stackrel{(IBP)}{=} f(x).$$

Exercise 2.11 Extend this to more general functions!

2.3 Equations in general domains

Theorem 2.12 (Mean Value Theorem for Harmonic Functions) Let $\Omega \subseteq \mathbb{R}$ be open, let $u \in C^2(\Omega)$ and $\Delta u = 0$ in Ω . Then

$$u(x) = \int_{B(x,r)} u = \int_{\partial B(x,r)} u \text{ for all } x \in \Omega, B(x,r) \subseteq \Omega$$

Proof. Consider all r > 0 s.t. $B(x, r) \subseteq \Omega$,

$$f(r) = \int_{\partial B(x,r)} u$$

We need to prove that f(r) is independent of r. When it is done, then we immediately obtain

$$f(r) = \lim_{t \to 0} f(t) = u(x)$$

as u is continuous. To prove that, consider

$$f'(r) = \frac{d}{dr} \left(\int_{\partial B(0,r)} u(x+y) \, dS(y) \right)$$

$$= \frac{d}{dr} \left(\int_{\partial B(0,1)} u(x+rz) \, dS(z) \right)$$

$$(\text{dom. convergence}) = \int_{\partial B(0,1)} \frac{d}{dr} [u(x+rz)] \, dS(z)$$

$$= \int_{\partial B(0,1)} \nabla u(x+rz) z \, dS(z)$$

$$= \int_{\partial B(x,r)} \nabla u(y) \frac{y-x}{r} \, dS(y)$$

$$= \frac{1}{|B(x,r)|_{\mathbb{R}^d}} \int_{\partial B(x,r)} \nabla \cdot u(y) \cdot \vec{n_y} \, dS(y)$$

$$(\text{Gauss-Green 2.2}) = \frac{1}{|B(x,r)|_{\mathbb{R}^d}} \int_{B(x,r)} \underbrace{(\Delta u)(y)}_{=0} \, dy = 0$$

Exercise 2.13 In 1D: $\Delta u = 0 \Leftrightarrow u'' = 0 \Leftrightarrow u(x) = ax + b$ (Linear Equation)

Remark 2.14 Recall the polar decomposition. Let $x \in \mathbb{R}^d$, x = (r, w), r = |x| > 0, $\omega \in S^{d-1}$, then

$$\int_{B(0,r)} g(y) \, dy = \int_0^r \left(\int_{B(0,r)} g(y) \, dS(y) \right) dr$$

Remark 2.15 We already proved that for u harmonic we have $u(x) = f_{\partial B(x,r)} u \, dy$. Now we have

$$\int_{B(x,r)} u(y) \, dy = \int_{B(0,r)} u(x+y) \, dy$$
(Pol. decomposition)
$$= \int_0^r \left(\int_{\partial B(0,s)} u(x+y) \, dS(y) \right) ds$$

$$= \int_0^r \left(\int_{\partial B(x,s)} u(y) \, dS(y) \right) ds$$
(Mean value property)
$$= \int_0^r \left(|\partial B(x,s)| \, u(x) \right) ds = |B(x,r)| \, u(x)$$

This implies

$$\oint_{B(x,r)} u(y) \, dy = u(x) \quad \text{for any } B(x,r) \subseteq \Omega.$$

Remark 2.16 The reverse direction is also correct, namely if $u \in C^2(\Omega)$ and

$$u(x) = \int_{B(x,r)} u(y) \, dy = \int_{\partial B(x,r)} u(y) \, dy \quad \text{for all } B(x,r) \subseteq \Omega,$$

then u is harmonic, i.e. $\Delta u = 0$ in Ω . (The proof is exactly like before)

Theorem 2.17 (Maximum Principle) Let $\Omega \subseteq \mathbb{R}^d$ be open, let $u \in C^2(\Omega) \cap C(\bar{\Omega})$, $\Delta u = 0$ in Ω . Then

- a) $\sup_{x \in \bar{\Omega}} u(x) = \sup_{x \in \partial \Omega} u(x)$
- b) Assume that Ω is connected. Then if there is a $x_0 \in \Omega$ s.t. $u(x_0) = \sup_{x \in \overline{\Omega}} u(x)$, then $u \equiv const.$ in Ω .

Proof. Given $U \subseteq \mathbb{R}^d$ open, we can write $U = \bigcup_i U_i$, where U_i is open and connected.

- b) Assume that Ω is connected and there is a $x_0 \in \Omega$ s.t. $u(x_0) = \sup_{y \in \Omega} u(x)$. Define $U = \{x \in \Omega \mid u(x) = u(x_0)\} = u^{-1}(u(x_0))$. U is closed since u is continuous. Moreover, U is open by the mean-value theorem. I.e. for all $x \in U$ there is a r > 0 s.t. $B(x,r) \subseteq U \subseteq \Omega$. Since U is connected we get $U = \Omega$, so u is constant in Ω . On the other hand, if there is no $x_0 \in \Omega$ s.t. $u(x_0) = \sup_{x \in \Omega} w$ we have $\forall x_0 \in \Omega : u(x) < \sup_{x \in \overline{\Omega}} u(x) = \sup_{x \in \partial \Omega} u(x)$
- a) Given $\Omega \subseteq \mathbb{R}^d$ open, we can write $\Omega = \bigcup_i \Omega_i$, where Ω_i is open and connected. By b) we have

$$\sup_{x \in \bar{\Omega}_i} u(x) = \sup_{x \in \partial \Omega_i} u(x), \quad \forall i$$

So we can conclude

$$\sup_{x \in \bar{\Omega}} u(x) = \sup_{x \in \partial\Omega} u(x).$$

Definition 2.18 • If $\Omega \subseteq \mathbb{R}^d$ is open, $u \in C^2(\Omega)$, then u is called *sub-harmonic* if $\Delta u \ge 0$ in Ω .

• If $\Delta u \leq 0$, then u is called *super-harmonic*.

Exercise 2.19 (E 1.4) Let $\Omega \subseteq \mathbb{R}^d$ be open and $u \in C^2(\Omega)$ be subharmonic.

a) Prove that u satisfies the Mean Value Inequality

$$\int_{\partial B(x,r)} u(y) \, dS(y) \geqslant \int_{B(x,r)} u(y) \, dy \geqslant u(x)$$

for all $B(x,r) \subseteq \mathbb{R}^d$.

- b) Assume further that Ω is connected and $u \in C(\bar{\Omega})$. Prove that u satisfies the strong maximum principle, namely either
 - u is constant in Ω , or
 - $\sup_{y \in \partial \Omega} u(y) > u(x)$ for all $x \in \Omega$.

My Solution. a) Let $f(r) = \int_{\partial B(x,r)} u(y) dS(y)$, then we have

$$\partial_{r} f(r) = \partial_{r} \oint_{\partial B(x,r)} u(y) \, dS(y)$$
(Dom. Convergence)
$$= \oint_{\partial B(x,r)} \partial_{r} u(y) \, dS(y)$$

$$= \oint_{\partial B(0,1)} \partial_{r} u(x+yr) \, dS(y)$$

$$= \oint_{\partial B(0,1)} \nabla u(x+yr) \cdot y \, dS(y)$$

$$= \oint_{\partial B(x,r)} \nabla u(y) \cdot \frac{y-x}{r} \, dS(y)$$

$$= \oint_{\partial B(x,r)} \nabla u(y) \cdot \vec{n}_{y} \, dS(y)$$
(Gauss-Green)
$$= \oint_{B(x,r)} \operatorname{div}(\nabla u(y)) \, dS(y)$$

$$= \oint_{B(x,r)} \Delta u(y) \, dS(y) \geqslant 0$$

So we can conclude that

$$\oint_{\partial B(x,r)} u(y) \, dS(y) = f(r) \geqslant \lim_{r \to 0} f(r) = u(x).$$

Now regard

$$\int_{B(x,r)} u(y) \, dy = \int_0^r \left(\int_{\partial B(x,r)} u(y) \, dS(y) \right) ds$$

$$= \int_0^r \left(|\partial B(x,r)| \int_{\partial B(x,r)} u(y) \, dS(y) \right) ds$$

$$\geqslant \int_0^r |\partial B(x,r)| \cdot u(x) \, dS(y)$$

$$= u(x) \int_0^r |\partial B(x,r)| \, dS(y) = u(x) |B(x,r)|.$$

Thus we have

$$u(x) \leqslant \int_{B(x,r)} u(y)dy.$$

Finally, lets regard

$$\int_{B(x,r)} u(y) \, dy = \int_0^r \left(|\partial B(x,s)| \oint_{\partial B(x,s)} u(y) \, dS(y) \right) \, ds$$

$$(\partial_r f(r) \geqslant 0) \qquad \leqslant \int_0^r \left(|\partial B(x,s)| \oint_{\partial B(x,r)} u(y) \, dS(y) \right) \, ds$$

$$= \oint_{\partial B(x,r)} u(y) \, dS(y) \int_0^r |\partial B(x,s)| \, ds$$

$$= \oint_{\partial B(x,r)} u(y) \, dS(y) \cdot |B(x,s)|$$

and we conclude

$$\int_{B(x,r)} u(y) \, dy \leqslant \int_{\partial B(x,r)} u(y) \, dS(y).$$

b) Let $x_0 \in \Omega$ s.t. $u(x_0) = \sup_{x \in \Omega} u(x)$. Now,

$$\sup_{x \in \Omega} u(x) = u(x_0) \leqslant \int_{\partial B(x_0, r)} u(y) \, dy$$
$$\leqslant \int_{\partial B(x_0, r)} \sup_{x \in \Omega} u(x) \, dy = \sup_{x \in \Omega} u(x)$$

Since u is continuous we get $u(y) = u(x_0)$ for all $y \in B(x_0, r)$, so u is constant.

Definition 2.20 The *Poisson Equation* for given f, g on a bounded set is:

$$\begin{cases}
-\Delta u = f, & \text{in } \Omega \\
u = g, & \text{on } \partial\Omega
\end{cases}$$

Theorem 2.21 (Uniqueness) Let $\Omega \subseteq \mathbb{R}^d$ be bounded, open and connected. Let $f \in C(\Omega), g \in C(\partial\Omega)$. Then there exists at most one solution $u \in C^2(\Omega) \cap C(\bar{\Omega})$, s.t.

$$\begin{cases} -\Delta u = f, & \text{in } \Omega \\ u = g, & \text{on } \partial \Omega \end{cases}$$

Proof. Assume that we have two solutions u_1 and u_2 . Then $u := u_1 - u_2$ is a solution to

$$\begin{cases} -\Delta u = 0, & \text{in } \Omega \\ u = 0 & \text{on } \partial \Omega \end{cases}$$

By the maximum principle, we know that u=0 in Ω . More precisely, by the maximum principle we have $\forall x\in\Omega$

$$\sup_{x \in \Omega} u(x) \leqslant \sup_{x \in \partial \Omega} u(x) = 0 \quad \Rightarrow \quad u(x) \leqslant 0$$

Since -u satisfies the same property we have $\forall x \in \Omega$:

$$\sup_{x \in \Omega} (-u(x)) \leqslant \sup_{x \in \partial \Omega} (-u(x)) = 0 \quad \Rightarrow \quad -u(x) \leqslant 0 \quad \Rightarrow \quad u(x) \geqslant 0$$

So we geht u(x) = 0 in Ω .

Exercise 2.22 (Bonus 1) Let Ω be open, connected and bounded in \mathbb{R}^d . Let $u \in C^2(\Omega) \cap C(\bar{\Omega})$ s.t.

$$\begin{cases} \Delta u = 0, & \text{in } \Omega \\ u = g, & \text{on } \partial \Omega \end{cases}$$

Prove that

a) If $g \ge 0$ on $\partial \Omega$, then $u \ge 0$ in Ω .

b) If $g \ge 0$ on $\partial \Omega$ and $g \ne 0$, then u > 0 in Ω .

My Solution. a) We have that $\Delta(-u) = 0$, so -u is harmonic in Ω . Since Ω is open and bounded we can apply the Maximum Principle (2.17) and get that

$$\sup_{x\in\bar{\Omega}}-u(x)\leqslant \sup_{x\in\partial\Omega}-g(x)\leqslant 0.$$

This implies $\inf_{x \in \Omega} u(x) \ge 0$, so $u \ge 0$ in Ω .

b) We prove this by contraposition. Assume there is a $x_0 \in \Omega$ s.t. $u(x_0) = 0$. Since we have $u \ge 0$ on Ω by a), it follows that

$$0 = -u(x_0) = \sup_{x \in \Omega} -u(x) \leqslant \sup_{x \in \partial\Omega} -g(x) \leqslant 0,$$

so -u attains its maximum on Ω . Hence -u=0=u is constant by the strong maximum principle because Ω is connected, in fact $0=u|_{\partial\Omega}=g$.

Lemma 2.23 (Estimates for derivatives) If u is harmonic in $\Omega \subseteq \mathbb{R}^d$, $\alpha \in \mathbb{N}_0^d$, $|\alpha| = N$ and $B(x_0, r) \subseteq \Omega$, then

$$|D^{\alpha}u(x)| \leqslant \frac{(c_d N)^N}{r^{d+N}} \int_{B(x,r)} |u| \, dy$$

Proof. Induction: Assume $|\alpha| = N - 1$, Take $|\alpha| = N$

$$|D^{\alpha}u(x_0)| \le \frac{|S_1|}{|B_1|\frac{r}{N}} \|D^{\beta}u\|_{L^{\infty}(B(x_0,\frac{r}{n}))}, \quad D^{\alpha}u = \partial_{x_i}(D^{\beta}u)_{|\beta|=N-1}$$

Note: $x \in B(x_0, \frac{r}{N})$, so $B(x, \frac{r(N-1)}{N}) \subseteq B(x_0, r)$. By the induction hypothesis:

$$||D^{\beta}u||_{L^{\infty}(B(x_{0},\frac{r}{N}))} \leq \frac{[c_{d}(N-1)]^{N-1}}{[r^{\frac{(N-1)}{N}}]^{d+N-1}} \int_{B(x_{0},r)} |u| \, dy$$

The conclusion is:

$$\begin{split} |D^{\alpha}u(x_{0})| &\leqslant \frac{|S_{1}|}{|B_{1}|\frac{r}{N}} \frac{[c_{d}(N-1)]^{N-1}}{(r\frac{N-1}{N^{d}})^{d+N-1}} \int_{B(x_{0},r)} |u| \, dy \\ &= \frac{|S_{1}|}{|\beta_{1}|} \frac{c_{d}^{N-1}}{\left(\frac{r}{N}\right)^{d+N}} \frac{1}{(N-1)^{d}} \int_{B(x_{0},r)} |u| \, dy \\ &= \frac{|S_{1}|}{|\beta_{1}|} \frac{c_{d}^{N-1}}{\left(\frac{r}{N}\right)^{d+N}} \frac{1}{N^{d}} \left(\frac{N}{N-1}\right)^{d} \int_{B(x_{0},r)} |u| \, dy \\ &\leqslant \frac{2^{d}|S_{1}|}{|B_{1}|} \frac{c_{d}^{N-1}N^{N}}{r^{d+N}} \int_{B(x_{0},r)} |u| \, dy \quad \text{if } c_{d} \geqslant \frac{2^{d}|S_{1}|}{|B_{1}|} \end{split}$$

Theorem 2.24 (Regularity) Let Ω be open in \mathbb{R}^d . Let $u \in C(\Omega)$ satisfy $u(x) = \int_{\partial B} u \, dy$ for any $x \in B(x, r) \subseteq \Omega$. Then u is a harmonic function in Ω . Moreover, $u \in C^{\infty}(\Omega)$ and u is analytic in Ω .

Exercise 2.25 (E 1.1: Proof the Gauss–Green formula) Let $f := (f_i)_1^d \in C^1(\mathbb{R}^d, \mathbb{R}^d)$. Prove that for every open ball $B(y, r) \subseteq \mathbb{R}^d$ we have

$$\int_{\partial B(y,r)} f(y) \cdot \nu_y \, dS(y) = \int_{B(y,r)} \operatorname{div} f \, dx.$$

Here ν_y is the outward unit normal vector and dS is the surface measure on the sphere.

Solution. We proof this in d=3. Let $f \in C^1(\mathbb{R}^3)$

$$\int_{B(0,1)} \partial_{x_3} f \, dx = \int_{\partial B(0,1)} f x_3 \, dS(x), \quad x = (x_1, x_2, x_3) \in \mathbb{R}^3, \vec{n} = \frac{x}{|x|} \text{ on } \partial B(0,1)$$

$$B(0,1) = \{x_1^2 + x_2^2 + x_3^2 \leqslant 1\}$$

$$= \{x_1^2 + x_2^2 \leqslant 1 - \sqrt{1 - x_1^2 - x_2^2} \leqslant x_3 \leqslant \sqrt{1 - x_1^2 - x_2^2}\}$$

Then:

$$\begin{split} \int_{B(0,1)} \partial_{x_3} f \, dx &= \int_{x_1^2 + x_2^2 \leqslant 1} \left(\int_{-\sqrt{1 - x_1^2 - x_2^2} \leqslant x_3 \leqslant \sqrt{1 - x_1^2 - x_2^2}} \partial_{x_3} f \, dx_3 \right) \, dx_1 \, dx_2 \\ &= \int_{x_1^2 + x_2^2 \leqslant 1} \left[f(x_1, x_2, \sqrt{1 - x_1^2 - x_2^2}) \right. \\ &\left. - f(x_1, x_2, -\sqrt{1 - x_1^2 - x_2^2}) \right] \, dx_1 \, dx_2 \end{split}$$

Lets take polar coordinates in 2D:

$$x_1 = r \cos \phi$$
 $r > 0, \phi \in [0, 2\pi)$
 $x_2 = r \sin \phi$ $\det \frac{\partial(x_1, x_2)}{\partial(r, \phi)} = r$

$$(\star) = \int_0^1 \int_0^{2\pi} [f(r\cos\phi, r\sin\phi, r) - f(r\cos\theta, r\sin\phi, -r)] r \, dr \, d\phi$$

On the other hand:

$$\int_{\partial B(0,1)} fx_3 \, dS$$

The polar coordinates in 3D are:

$$x_1 = r \cos \phi \sin \theta$$
 $r > 0, \phi \in (0, 2\pi), \theta \in (0, \pi)$
 $x_2 = r \sin \phi \sin \theta$
$$\det \frac{\partial x_1, x_2, x_3}{\partial (r, \phi, t)} = r^2 \sin \theta$$

$$x_3 = \cos \theta$$

Then:

$$(\star \star) = \int_0^{2\pi} \int_0^{\pi} f(\cos \phi \sin \theta, \sin \phi \sin \theta, \cos \theta) \sin \theta \cos \theta \, d\theta \, d\phi$$

$$= \int_0^{2\pi} \left(\int_0^{\frac{\pi}{2}} + \int_{\frac{\pi}{2}}^{\pi} d\theta \right) \, d\phi$$

$$(r = \sin \theta) = \int_0^{2\pi} \int_0^1 f(r \cos \phi, r \sin \phi, \sqrt{1 - r^2}) r \, dr \, d\phi$$

$$- f(r \cos \phi, r \sin \phi, -\sqrt{1 - r^2}) r \, dr \, d\phi$$

Exercise 2.26 (E 1.2) Let $u \in C(\mathbb{R}^d)$ and $\int_{B(x,r)} u \, dy = 0$ for every open ball $B(x,r) \subseteq \mathbb{R}^d$. Show that u(x) = 0 for all $x \in \mathbb{R}^d$.

My Solution. Assume there is a $x_0 \in \mathbb{R}^d$ s.t. w.l.o.g. $u(x_0) > 0$. Since u is continous there is a ball $B(x_0, r)$ s.t. $u(y) > \frac{u(x_0)}{2}$ for all $y \in B(x_0, r)$. But then we get

$$\int_{B(x_0,r)} u(y) \, dy \geqslant \int_{B(x_0,r)} \frac{u(x_0)}{2} \, dy = \frac{u(x_0)}{2} \, |B(x_0,r)| > 0.$$

Exercise 2.27 (E 1.3) Let $f \in C_c^1(\mathbb{R}^d)$ with $d \ge 2$ and $u(x) := (\Phi \star f)(x)$. Prove that $u \in C^2(\mathbb{R}^2)$ and $-\Delta u(x) = f(x)$ for all $x \in \mathbb{R}^d$ (2.9 was the same for $f \in C_1(\mathbb{R})$)

Theorem 2.28 (Liouville's Theorem) If $u \in C^2(\mathbb{R}^d)$ is harmonic and bounded, then u = const.

Proof. By the bound of the derivative 2.23 we have

$$\begin{aligned} |\partial_{x_i} u(x_0)| &\leq \frac{c_d}{r^{d+1}} \int_{B(x_0, r)} |u| \, dy \quad \forall x_0 \in \mathbb{R}^d \, \forall r > 0 \\ &\leq \|u\|_{L^{\infty}} \frac{c_d}{r^{d+1}} |B(x_0, r)| \\ &\leq \|u\|_{L^{\infty}} \frac{c_d}{r} \xrightarrow{r \to \infty} 0 \end{aligned}$$

Thus $\partial_{x_i} u = 0$ for all $i = 1, 2, \dots d$ and u = const. in \mathbb{R}^d

Theorem 2.29 (Uniqueness of solutions to Poisson Equation in \mathbb{R}^d) If $u \in C^2(\mathbb{R}^d)$ is a bounded function and satisfies $-\Delta u = f$ in \mathbb{R}^d where $f \in C_c^2(\mathbb{R}^d)$, then we have

$$u(x) = \Phi \star f(x) + C = \int_{\mathbb{R}^d} \Phi(x - y) f(y) \, dy + C \quad \forall x \in \mathbb{R}^d$$

where C is a constant and Φ is the fundamental solution of the Laplace equation in \mathbb{R}^d .

Proof. If we can prove that v is bounded, then v = const.. We first need to show that $\Phi \star f$ is bounded.

$$\Phi = \Phi_1 + \Phi_2 = \Phi\mathbb{1}(|x| \leqslant 1) + \Phi(|x| \geqslant 1)$$
$$\Phi \star f = \Phi_1 \star f + \Phi_2 \star f$$

We have $\Phi_1 \star f \in L^1(\mathbb{R}^d)$ and $\Phi_2 \star f$ is bounded since $\Phi \to 0$ as $|x| \to \infty$ in $d \ge 3$.

Exercise 2.30 (Hanack's inequality) Let $u \in C^2(\mathbb{R}^d)$ be harmonic and non-negative. Prove that for all open, bounded and connected $\Omega \subseteq \mathbb{R}^d$, we have

$$\sup_{x \in \Omega} u(x) \leqslant C_{\Omega} \inf_{x \in \Omega} u(x),$$

where C_{∞} is a finite constant depending only on Ω .

Proof. (Exercise) Hint: $\Omega = B(x, r)$. General case cover Ω by finitely many balls, one ball is inside Ω .

Chapter 3

Convolution, Fourier Transform and Distributions

Definition 3.1 (Convolution) Let $f, g : \mathbb{R}^d \to \mathbb{R}$ or \mathbb{C} .

$$(f \star g)(x) = \int_{\mathbb{R}^d} f(x - y)g(y) \, dy = \int_{\mathbb{R}^d} f(y)g(x - y) \, dy = (g \star f)(x)$$

Remark 3.2 (Properties of the Convolution) • $(f \star g)(x) = f \star (g \star h)$

$$\bullet \ \hat{f \star g} = \hat{f} \star \hat{g}$$

Theorem 3.3 (Young Inequality) If $f \in L^1(\mathbb{R}^d)$ and $g \in L^p(\mathbb{R}^d)$, where $1 \leq p \leq \infty$, then $f \star g \in L^p(\mathbb{R}^d)$ and $\|f \star g\|_{L^p} \leq \|f\|_{L^1} \|g\|_{L^p}$. More generally, if $f \in L^p(\mathbb{R}^d)$, $g \in L^q(\mathbb{R}^d)$, then $f \star g \in L^1(\mathbb{R}^d)$, $\|f \star g\|_{L^1} \leq \|f\|_{L^p} \|g\|_{L^q}$, where $1 \leq p, q, r, \leq \infty$, $\frac{1}{p} + \frac{1}{q} = 1 + \frac{1}{r}$

Proof. Let $f \in L^1, g \in L^p$. With the Hölder Inequality ??, we have:

$$||f \star g||_{L^{p}}^{p} = \int_{\mathbb{R}^{d}} |f \star g(x)|^{p} dx$$

$$\leq ||f||_{L^{1}}^{\frac{p}{q}} \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} |f(x - y)||g(y)|^{p} dy dx$$

$$= ||f||_{L^{1}}^{\frac{p}{q} + 1} ||g||_{L^{p}}^{p}$$

So we have $||f \star g||_{L^p} \leq ||f||_{L^1} ||g||_{L^p}$

Theorem 3.4 (Smoothness of the Convolution) If $f \in C_c^{\infty}(\mathbb{R}^d)$, $g \in L^p(\mathbb{R}^d)$, $1 \le p \le \infty$. Then $f \star g \in C^{\infty}(\mathbb{R})$ and

$$D^{\alpha}(f \star q) = (D^{\alpha}f) \star q$$

for all $\alpha = (\alpha_1, \dots, \alpha_d), \alpha_i \in \{0, 1, 2, \dots\}$

Proof. First we note that $x \mapsto (f \star g)$ is continous as $x_n \to x$ in \mathbb{R}^d since

$$(f \star g)(x_n) = \int_{\mathbb{R}^d} f(x_n - y)g(y) \, dy \xrightarrow{\text{dom. conv.}} \int_{\mathbb{R}^d} f(x - y)g(y) \, dy = (f \star g)(x)$$

We can apply Dominated convergence because

$$f(x_n - y)g(y) \to f(x - y)g(y) \quad \forall y \text{ as } f \text{ is continuous and } x_n \to x$$

and

$$|f(x_n - y) \ g(y)| \le ||f||_{L^{\infty}} |g(y)| \ \mathbb{1}(|y| \le R) \in L^1(\mathbb{R}^d).$$

Where R > 0 satisfies $B(0,R) \supseteq \operatorname{supp} f + \operatorname{sup}_n |x_n|$. Now we can compute the derivatives:

$$\partial_{x_i}(f \star g)(x) = \lim_{h \to 0} \frac{(f \star g)(x + he_i) - (f \star g)(x)}{h}$$

$$= \lim_{h \to 0} \int_{\mathbb{R}^d} \frac{f(x + he_i - y) - f(x - y)}{h} g(y) \, dy$$
(Dominated Convergence)
$$= \int_{\mathbb{R}^d} \lim_{h \to 0} \frac{f(x + he_i - y) - f(x - y)}{h} g(y) \, dy$$

$$= \int_{\mathbb{R}^d} (\partial_{x_i} f)(x - y) g(y) \, dy$$

We could apply Dominated Convergence since

$$\frac{f(x+he_i-y)-f(x-y)}{h}g(y) \xrightarrow{h\to 0} (\partial_{x_i}f)(x-y)g(y) \quad \text{as } f\in C^1$$

$$\left|\frac{f(x+he_i-y)-f(x-y)}{h}g(y)\right| \leqslant \|\partial_{x_i}f\|_{L^\infty}|g(y)| \ \mathbb{1}(|y|\leqslant R) \in L^1(\mathbb{R}^d)$$

where $B(0,R) \supseteq \operatorname{supp}(f) + B(0,|x|+1)$ and $\partial_{x_i}(f \star g) = (\partial_{x_i}f) \star g \in C(\mathbb{R}^d)$ since $\partial_{x_i}f \in C_c^{\infty}(\mathbb{R}^d)$. By induction we get $D^{\alpha}(f \star g) = (D^{\alpha}f \star g) \in C(\mathbb{R}^d)$.

Remark 3.5 Question: Is there a f s.t. $f \star g = g$ for all g. In fact there is no regular function f that solves this formally:

$$f \star g = g \Rightarrow \widehat{f \star g} = \widehat{g} \Rightarrow \widehat{f}\widehat{g} = \widehat{g} \Rightarrow \widehat{f} = 1 \Rightarrow f \text{ is not a regular function!}$$

However, if f is the Dirac-Delta Distribution, $f = \delta_0$ then $\delta_0 \star g = g$ for all g. Formally:

$$\delta_0(x) = \begin{cases} 0 & x \neq 0 \\ \infty & x = 0 \\ \int \delta_0 = 1 \end{cases}$$

In fact, if $f \in L^1(\mathbb{R}^d)$, $\int f = 1$, $f_{\epsilon}(x) = \epsilon^{-d} f(\epsilon^{-1} x)$, then $f_{\epsilon} \to \delta_0$ in an appropriate sense and $f_{\epsilon} \star g \to g$ for all g nice enough.

Theorem 3.6 (Approximation by convolution) Let $f \in L^1(\mathbb{R}^d)$, $\int f = 1$, $f_{\epsilon}(x) = \epsilon^{-d} f(\frac{x}{\epsilon})$. Then for all $g \in L^p(\mathbb{R}^d)$, where $1 \leq p < \infty$, then

$$f_{\epsilon} \star g \to g \quad \text{in } L^p(\mathbb{R}^d)$$

Proof.

Step 1: Let $f, g \in C_c(\mathbb{R}^d)$. Then

$$(f_{\epsilon} \star g)(x) - g(x) = \int_{\mathbb{R}^{d}} f_{\epsilon}(y)g(x - y) \, dy - \int_{\mathbb{R}^{d}} f_{\epsilon}(y)g(x) \, dy$$

$$= \int_{\mathbb{R}^{d}} f_{\epsilon}(y)(g(x - y) - g(x)) \, dy$$

$$|(f_{\epsilon} \star g)(x) - g(x)| = \left| \int_{\mathbb{R}^{d}} f_{\epsilon}(y)(g(x - y) - g(x)) \, dy \right|$$

$$\leqslant \int_{\mathbb{R}^{d}} |f_{\epsilon}(y)||g(x - y) - g(x)| \, dy$$

$$\leqslant \int_{|y| \leqslant R_{\epsilon}} |f_{\epsilon}(y)||g(x - y) - g(x)| \, dy$$

$$\leqslant \underbrace{\int_{|y| \leqslant R_{\epsilon}} |f_{\epsilon}(y)| \, dy}_{|z| \leqslant R} \left[\sup_{|z| \leqslant R} |g(x - z) - g(x)| \right] \xrightarrow{\epsilon \to 0} 0$$

We have Dominated Convergence since:

$$(f_{\epsilon} \star g)(x) - g(x) \to 0 \text{ as } \epsilon \to 0$$

and

$$|f_{\epsilon} \star g(x) - g(x)| \leqslant \|f\|_{L^{1}} \sup_{|z| \leqslant R_{\epsilon}} |g(x - z) - g(x)| \leqslant 2\|f\|_{1} \|g\|_{L^{\infty}} \mathbb{1}(|x| \leqslant R_{1}).$$

Where $B(0, R_1) \supseteq \operatorname{supp}(g) + B(0, R_{\epsilon})$, thus $f_{\epsilon} \star g \to g$ in $L^p(\mathbb{R}^d)$. To remove the technical assumptions $f, g \in C_c(\mathbb{R}^d)$, then we use a density argument. We use the fact that $C_c(\mathbb{R}^d)$ is dense in $L^p(\mathbb{R}^d)$, $1 \le p < \infty$.

Step 2: Let $g \in C_c(\mathbb{R}^d)$, $g \in L^p(\mathbb{R}^d)$. Then there is $\{g_m\} \subseteq L^p(\mathbb{R}^d)$, $g_m \to g$ in $L^p(\mathbb{R}^d)$.

$$\begin{split} \|f_{\epsilon} \star g - g\|_{L^{p}} &\leq \|f_{\epsilon} \star (g - g_{m})\|_{L^{p}} + \|f_{\epsilon} \star g_{m} - g_{m}\|_{L^{p}} + \|g_{m} - g\|_{L^{p}} \\ & (\text{Young}) &\leq \|f_{\epsilon}\|_{L^{1}} \|g - g_{m}\|_{L^{p}} + \|f_{\epsilon} \star g_{m} - g_{m}\|_{L^{p}} + \|g_{m} - g\|_{L^{p}} \\ &\leq \|f\|_{L^{1}} \|g - g_{m}\|_{L^{p}} + \|f_{\epsilon} \star g_{m} - g_{m}\|_{L^{p}} + \|g_{m} - g\|_{L^{p}} \\ &\leq (\|f\|_{L^{1}} + 1)\|g - g_{m}\|_{L^{p}} + \|f \star g_{m} - g_{m}\|_{L^{p}} \end{split}$$

So we get:

$$\limsup_{\epsilon \to 0} \|f_{\epsilon} \star g - g\|_{L^{p}} \leqslant (\|f\|_{L^{p}} + 1)\|g - g_{m}\|_{L^{p}} + \limsup_{\epsilon \to 0} \|f_{\epsilon} \star g_{m} - g_{m}\|_{L^{p}}$$

$$\underbrace{\lim\sup_{\epsilon \to 0} \|f_{\epsilon} \star g - g\|_{L^{p}}}_{\text{Oby step 1.}}$$

$$\xrightarrow{m\to\infty} 0$$

Step 3: Let $f \in L^1(\mathbb{R}^d)$ and $g \in L^p(\mathbb{R}^d)$. Take $\{f_m\} \subseteq C_c(\mathbb{R}^d)$, s.t.

$$\begin{cases} F_m \to ginL^1(\mathbb{R}) \text{ as } m \to \infty \\ \int_{\mathbb{R}^d} F_m = 1 \text{(it is possible since } \int_{\mathbb{R}^d} f = 1 \text{)} \end{cases}$$

Define
$$F_{m,\epsilon}(x) = \epsilon^{-d} F_m(\epsilon^{-1} x)$$
 (recall $f_{\epsilon}(x) = \epsilon^{-d} f(\epsilon^{-1} x)$). Then:

$$f_{\epsilon} \star g - g = (f_{\epsilon} - F_{m,\epsilon}) \star g + F_{m,\epsilon} \star g - g$$

$$\Rightarrow \|f_{\epsilon} - g\|_{L^{p}} \leq \underbrace{\|f_{\epsilon} - F_{m,\epsilon} \star g\|_{L^{p}}}_{+} + \|F_{m,\epsilon} \star g - g\|_{L^{p}}$$

$$\underbrace{\text{Young}}_{\leqslant} \|f_{\epsilon} - F_{m,\epsilon}\|_{L^{1}} \|g\|_{L^{p}} = \|f - F_{m}\|_{L^{1}} \|g\|_{L^{p}}$$

$$\Rightarrow \limsup_{\epsilon \to 0} \|f_{\epsilon} \star g - g\|_{L^{p}} \leq \|f - F_{m}\|_{L^{1}} \|g\|_{L^{p}} = \|f - F_{m}\|_{L^{1}} \|g\|_{L^{p}}$$

Lemma 3.7 $C_c(\mathbb{R}^d)$ is dense in $L^p(\mathbb{R}^d)$, $1 \leq p < \infty$

Proof. For all $g \in L^p(\mathbb{R}^d)$ there are g_m step functions and $g_m \to m$ in $L^p(\mathbb{R}^d)$, We can assume that Ω is open and bounded and we want to approximate χ_{Ω} by $C_c(\mathbb{R}^d)$.

Lemma 3.8 (Urnson) Define

$$\Omega_{\epsilon} = \{ x \in \Omega \mid \operatorname{dist}(x, \partial \Omega) > \epsilon \}$$

Then there is a $\eta_{\epsilon} \in C_c(\mathbb{R}^d)$ s.t.

$$\begin{cases} 0 \leqslant \eta(x) \leqslant 1 & \forall x \in \mathbb{R}^d \\ \eta_{\epsilon}(x) = 1 & \text{if } x \in \Omega_{\epsilon} \\ \eta_{\epsilon}(x) = 0 & \text{if } x \notin \Omega \end{cases}$$

Lemma 3.9 (Gernal Version of Urnson) If $A, B \subseteq \mathbb{R}^d$, A closed, B closed, $A \cap B = \emptyset$. Then

$$\eta(x) = \frac{\operatorname{dist}(x, A)}{\operatorname{dist}(x, A) + \operatorname{dist}(x, B)}$$

Then $\eta \in C(\mathbb{R}^d)$, $0 \leq \eta \leq 1$ and $\eta = 0$ if $x \in B$, $\eta = 1$ if $x \in A$. App to $A = \overline{\Omega_{\epsilon}} \subset\subset \Omega$ and $B = \mathbb{R}^d \setminus \Omega$.

Theorem 3.10 (Appendix C4 in Evans) Let Ω be open in \mathbb{R}^d and for $\epsilon > 0$ define

$$\Omega_{\epsilon} = \{ x \in \Omega \mid \operatorname{dist}(x, \mathbb{R}^d \setminus \Omega) > \epsilon \}$$

Let $f \in C_c^{\infty}(\mathbb{R}^d)$, $\int_{\mathbb{R}^d} f = 1$, supp $f \subseteq B(0,1)$, $f_{\epsilon}(x) = \epsilon^{-d} f(\epsilon^{-1} x)$ supp is $B(0,\epsilon)$. Then for all $g \in L^p_{loc}(\Omega)$ (i.e. $\mathbb{1}_K g \in L^p(\Omega) \forall K$ compakt set in Ω), then:

- a) $g_{\epsilon}(x) = (f_{\epsilon} \star g)(x) = \int_{\mathbb{R}^d} f_{\epsilon}(x y)g(y) dy \int_{\Omega} f_{\epsilon}(x y)g(y) dy$ is well-defined in Ω_{ϵ} and $g_{\epsilon} \in C^{\infty}(\Omega_{\epsilon})$.
- b) $g_{\epsilon} \to g$ in $L^p_{loc}(\Omega)$ if $1 \leq p < \infty$ and $g_{\epsilon}(x) \to g(x)$ almost everywhere $x \in \Omega$.
- c) If $g \in C(\Omega)$, then $g_{\epsilon}(x) \to g(x)$ uniformly in any compact subset of Ω .

Proof. a) $D^{\alpha}(g_{\epsilon}) = (D^{\alpha}f_{\epsilon}) \star g \in C(\Omega_{\epsilon})$

b) Already proved in \mathbb{R}^d space.

Corrolary 3.11 (Lebesgue differentiation theorem) If $f \in L_{loc}^P(\mathbb{R}^d)$, then

$$\oint_{B(x,\epsilon)} |f(y) - f(x)|^p dy \to 0 \quad \text{as } \epsilon \to 0$$

Exercise 3.12 (E 2.1) Let $u \in C^2(\mathbb{R}^2)$ be convex. I.e.

$$tu(x) + u(y)(1-t) \ge u(tx + (1-t)y) \forall x, y \in \mathbb{R}^d \forall t \in [0,1]$$

a) Prove for all $x \in \mathbb{R}^d$ that H(x) = ...

Solution.

a In 1D: If u is convex $\Leftrightarrow u''(x) \ge 0$ for all $x \in \mathbb{R}$. In general: Taylor expansion for all $x, z \in \mathbb{R}^d$:

$$u(x) = u(z) + \nabla u(z)(x - y) + \int_0^1 \sum_{|\alpha| = 2} D^{\alpha} u(z + s(x - z)) \frac{(x - z)^{\alpha}}{\alpha!} ds$$

$$x = z + s(x - z), s = 1$$
 Use $z = tx + (t - 1)y \Rightarrow x - z = (1 - t)(x - y)$

$$tu(x) = tu(z) + t\nabla u(z)(1-t)(x-y) + t\int_0^1 \sum_{|\alpha|=2} D^{\alpha}u(z+s(x-z)) \frac{[(1-t)(x-y)]^{\alpha}}{\alpha!} ds$$

$$(1-t)u(y) = (1-t)u(z) + (1-t)\nabla u(z)t(y-x) + (1-t)\int_0^r \sum_{|\alpha|=2} D^{\alpha}u(z+s(y-z))\frac{[t(y-x)]^{\alpha}}{\alpha!} ds$$

$$\Rightarrow tu(x) + (1-t)u(y) = u(z) + t \int_0^1 \dots + (1-t) \int_0^1 \dots$$
$$\Rightarrow t \int_0^1 \dots + (1-t) \int_0^1 \dots \geqslant 0 \forall x, y, t, z = tx + (1-t)y$$

$$t(1-t)^2 \int_0^1 \sum_{|\alpha|=2} D^{\alpha} u(z+s(x-z)) \frac{(x-y)}{\alpha!} \, ds + (1-t)t^2 \int_0^1 \sum_{|\alpha|=2} D^{\alpha} u(z+s(y-z)) \frac{(y-z)^{\alpha}}{\alpha!} \, ds \geqslant 0$$

for all $x, y \in \mathbb{R}^d$, $t \in [0, 1]$, z = tx + (1 - t)y. Divides for t(1 - t)

$$(1-t)\int_0^1\cdots+\int_0^1\cdots\geqslant 0$$

Take $t \to 0$

$$\int_0^1 \sum_{|\alpha|=2} D^{\alpha} u(y + s(x - y)) \frac{(x - y)^{\alpha}}{\alpha!} ds \geqslant 0 \forall x, y \in \mathbb{R}^d$$

Take $y = x + a, a \in \mathbb{R}^d$

$$\int_0^1 \sum_{|\alpha|=2} D^{\alpha} u(x+a+sa) \frac{a^{\alpha}}{\alpha!} ds \geqslant 0 \forall \epsilon > 0, \forall x, a \in \mathbb{R}^d$$

Take $\epsilon \to 0$

$$\int_0^1 \sum_{|\alpha|=2} D^\alpha u(x) \frac{a^\alpha}{\alpha!} \geqslant 0 \Rightarrow \sum_{i,j=1,i\neq j} \partial_{x_i} \partial_{x_j} u(x) a_i a_j + \sum_{i=j=1}^d \partial_{x_i}^2 u(x) \frac{a_i^2}{2}$$

We get

$$\frac{1}{2}a^T H a \geqslant 0 \forall a (a_i)_{i=1}^d \in \mathbb{R}^d$$

b
$$H(x) \geqslant 0 \Rightarrow (\partial_i \partial_j u) \geqslant 0 \Rightarrow TrH(x) \geqslant 0 \Rightarrow \sum_{i=1}^d \partial_{x_i}^2 u(x) \geqslant 0 \Rightarrow \Delta u(x) \geqslant 0 \forall x \in \mathbb{R}^d$$

Exercise 3.13 (E 2.2)

Solution. Regard d=3. De function $\frac{1}{|x|}$ is harmonic in $\mathbb{R}^3\setminus\{0\}$. We prove

$$\oint_{\partial B(x,r)} \frac{dS(y)}{|y|} = \frac{1}{\max(|x|,r)}$$

If |x| > r, then $0 \notin B(x, r + \epsilon)$. Then

$$y \mapsto \frac{1}{|y|}$$

is harmonic in $B(x, r + \epsilon)$. Then by the Mean Value Property:

$$\oint_{\partial B(x,r)} \frac{dS(y)}{|y|} = \frac{1}{|x|}$$

If |x| < r: Then $\frac{1}{|y|}$ is not harmonic in B(x,r) since $0 \in B(x,r)$. Note

$$f_{\partial B(x,r)} \frac{dS(y)}{|y|} = f_{\partial B(0,r)} \frac{dS(y)}{|x-y|}$$

This function depens on x only via |x|.

$$\dots = \int_{\partial B(0,r)} \frac{dS(y)}{|Rx - Ry|}$$

for all R rotation SO(3), $dS(R_y) = dS(y)$

$$= \int_{\partial B(0,r)} \frac{dS(y)}{|Rx - y|}$$

$$= \int_{\partial B(0,r)} \frac{dS(y)}{|z - y|}$$

$$= \int_{\partial B(0,|x|)} \left(\int_{\partial B(0,r)} \frac{dS(y)}{|z - y|} \right) dS(y)$$
(Fubini)
$$= \int_{\partial B(0,r)} \left(\int_{\partial B(0,|x|)} \frac{dS(z)}{|z - y|} \right) dS(y)$$
(case 1 since $|y| = r > |x|$)
$$= \int_{\partial B(0,r)} \left(\frac{1}{|y|} \right) dS(y) = \frac{1}{r}$$

If |x| = r: Continuiti: $x \mapsto f_{\partial B(0,r)} \frac{dS(y)}{|x-y|}$

Remark 3.14 For $f \in C^{|\alpha|}, g \in C^{|\beta|}$:

$$D^{\alpha+\beta}(f\star g)=(D^{\alpha}f)\star(D^{\beta}g)$$

Lemma 3.15 If $d \ge 3$ and $f : \mathbb{R}^d \to \mathbb{R}$ radial. Then:

$$\left(\frac{1}{|x|^{d-2}} \star f\right)(x) = \int_{\mathbb{R}^d} \frac{f(y)}{|x-y|^{d-2}} \, dy$$
$$= \int_{\mathbb{R}^d} \frac{f(y)}{\max(|x|^{d-2}, |y|^{d-2})} \, dy$$

Proof. (d=3) Polar coordinates:

$$\int_{\mathbb{R}^3} \frac{f(y)}{|x-y|} \, dy = \int_0^\infty \left[\int_{\partial B(0,1)} \frac{1}{|x-rw|} \, d\omega \right] f(r) \, dr$$

$$(a) = \int_0^\infty \left[\int_{\partial B(0,1)} \frac{d\omega}{\max(|x|,r)} \right] f(r) \, dr$$

$$= \int_{\mathbb{R}^3} \frac{f(y)}{\max(|x|,|y|)} \, dy$$

(b) (d=3) If f radial and non-negative

$$\int_{\mathbb{R}^3} \frac{f(y)}{|x - y|} = \int_{\mathbb{R}^3} \frac{f(y)}{|x|} \, dy = \frac{(Sf?)}{|x|}$$

Then

$$\int_{\mathbb{R}^{3}} \int_{\mathbb{R}^{3}} \frac{f_{1}(x-z_{1})f_{2}(y-z_{2})}{|x-y|} dx dy = \int_{\mathbb{R}^{3}} \int_{\mathbb{R}^{3}} \frac{f_{1}(x)f_{2}(y)}{|x+z_{1}-y-z_{2}|} dx dy$$

$$= \int_{\mathbb{R}^{3}} \left(\int_{\mathbb{R}^{3}} f_{1}(x) dx \right) f_{2}(y) dy \leqslant \int_{\mathbb{R}^{3}} \frac{\left(\int_{\mathbb{R}^{3}} f_{1} \right)}{|y+z_{2}-z_{1}|} f_{2}(y) dy$$

$$\leqslant \frac{\left(\int_{\mathbb{R}^{3}} f_{1} \right) \left(\int_{\mathbb{R}^{3}} f_{2} \right)}{|z_{1}-z_{2}|}$$

Exercise 3.16 (Bonus 2) Consider $u(x) = \frac{1}{|x|}$ in \mathbb{R}^d .

- 1. Prove that u is sub-harmonic in $\mathbb{R}^2 \setminus \{0\}$.
- 2. Prove that if $f: \mathbb{R}^2 \to \mathbb{R}$ radial, non-negative, measurable:

$$\int_{\mathbb{R}^2} \frac{f(y)}{|x - y|} \, dy \geqslant \int_{\mathbb{R}^2} \frac{f(y)}{\max(|x|, |y|)} \, dy$$

Definition 3.17 (Fourier Transform)

Theorem 3.18 (Basic Properties) 1. For all $f \in L^1(\mathbb{R}^d) \cap L^2(\mathbb{R}^d)$, $\|\hat{f}\|_{L^2} = \|f\|_{L^2}$. Moreover, F can be extended to be a unitary transforamtion $L^2(\mathbb{R}^d) \to L^2(\mathbb{R}^d)$ s.t.

- 2. $\widehat{D^{\alpha}f}(k)=(2\pi ik)^{\alpha}\widehat{f}(k)$ as laj as $(2\pi ik)^kf(k)\in L^2(\mathbb{R}^d)$ $(k^{\alpha}=k_1^{alpha_1}\cdots k_{\alpha}^{\alpha_k})$
- 3. (Formel) $\widehat{f+g}(k) = \widehat{f}(k)\widehat{g}(k)$ if f,g are nice enough.

Theorem 3.19 (Hausdorff-Young-Inequality) If $1 \le p \le 2$ and $\frac{1}{p} + \frac{1}{p'} = 1$, then

$$F: L^P(\mathbb{R}^d) \to L^p(\mathbb{R}^d)$$

and

$$\|\hat{f}\|_{L^p} \leqslant \|f\|_{L^p} \quad \forall f \in L^p(\mathbb{R}^d)$$

Remark 3.20 If $1 \le p \le 2$ and $f \in L^p(\mathbb{R}^d)$ we can write $f = f_1 + f_2$ when $f_1 \in L^1$, $f_2 \in L^2$, e.g.

$$f = \underbrace{f\mathbb{1}(|f|\geqslant 1)}_{f_1} + \underbrace{f\mathbb{1}(|f|<1)}_{f_2}$$

$$\int_{\mathbb{R}^d} |f_2|^2 \, dy = \int_{\mathbb{R}^d} |f|^2 \mathbb{1}(|f| < 1) \leqslant \int_{\mathbb{R}^d} |f|^p \, dy < \infty$$
$$\int_{\mathbb{R}^d} |f_1| = \int_{\mathbb{R}^d} |f| \mathbb{1}(|f| \geqslant 1) \leqslant \int_{\mathbb{R}^d} |f|^p < \infty$$

thus we can define $\hat{f} = \hat{f}_1 + \hat{f}_2$ well defined in $L^{\infty}(\mathbb{R}^d) + L^2(\mathbb{R}^d)$ what it in unclear why $\hat{f} \in L^q$? Proof of HY inequality. Weee nedd Riez-Theorem interpolation theorem. If $1 \leq p_0, p_1, q_0, q_1 \leq \infty$, and $\Omega \subseteq \mathbb{R}^d$ open and is a linear operator and

 $T: L^{p_0} \to L^{q_0} \text{ and } \|T\|_{L^{p_1} \to L^{q_1}} \leqslant 1, \quad \text{ for } i = 0, 1. \text{ Then } T: L^{p_0} \to L^{q_0} \text{ and } \|T\|_{L^{p_1} \to L^{q_1}} \leqslant 1, \quad \text{ for any } 0 < \theta < 1$

Consider the Foureir Transform:

$$F: L^1 + L^1 \to L^2 + L^{\infty}$$

and

$$||F||_{L^{1}\to L^{\infty}} \leqslant 1 \text{ as } ||\hat{f}||_{L^{\infty}} \leqslant ||f||_{L^{1}} \forall f p \in L^{1}$$

$$||F||_{L^{2}\to L^{2}} = 1 \text{ as } ||\hat{f}||_{L^{2}} = ||f||_{L^{2}} \forall f \in L^{2}$$

$$\Rightarrow ||F||_{L^{p_{0}}\to L^{p_{t}}} \leqslant 1 \forall \theta \in (0, 1)$$

$$p_0 = 1, p_1 = 2, q_0 = \infty, q_1 = 2$$

Exercise 3.21 Hint: first $f, g \in C_c(\mathbb{R}^d)$ then use a density arugment)

We want to deque the Fourier transform for singular function, e.g. $\frac{1}{|x|}$. Poisson:

$$-\Delta u = f \text{ in } \mathbb{R}^d$$

$$\Rightarrow |2\pi k|^2 \hat{u}(k) = \hat{f}(k)$$

$$\Rightarrow \hat{u}(k) = \frac{1}{|2\pi k|^2} \hat{f}(k)$$

If we can find G s.t. $\hat{G}(k) = \frac{1}{|2\pi k|^2}$, then

$$\hat{u}(k) = \hat{G}(k)\hat{f}(k) = \widehat{G \star f}$$

$$\Rightarrow u(x) = (G \star g)(x) = \int_{\mathbb{R}^d} G(x - y)f(y) \, dy$$

In fact G is the fundamential solution of laplace quation.

Theorem 3.22 (Fourier Transform of $\frac{1}{|x|^d}$ for $0 < \alpha < d$) . We have permally

$$\widehat{\frac{c_{\alpha}}{|x|^{\alpha}}} = \frac{c_{d-\alpha}}{|k|^{d-\alpha}} \forall 0 < \alpha < d$$

Here

$$c_{\alpha} = \pi^{-\frac{d}{2}} \Gamma\left(\frac{\alpha}{2}\right) = \pi^{-\frac{\alpha}{2}} \int_{0}^{\infty} e^{-\lambda} \lambda^{\frac{\alpha}{2} - 1} d\lambda$$

More precisely, for all $f \in C_c^{\infty}(\mathbb{R}^d)$,

$$\frac{c_{\alpha}}{|x|^{\alpha}} \star f = \left(\frac{c_{d-\alpha}}{|k|^{d+\alpha}} (k)f\right)$$

Moreover if $\alpha < \frac{d}{2}$, then we also have

$$\left(\frac{c_{\alpha}}{|x|^{\alpha}j} \star f\right)^{\wedge} = \frac{c_{d-\alpha}}{|f|^{d-\alpha}} \hat{f}(k)$$