Partial Differerential Equations Thành Nam Phan Winter Semester 2021/2022

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Chapter 1

Introduction

A differential equation is an equation of a function and its derivatives.

Example 1.1 (Linear ODE) Let $f : \mathbb{R} \to \mathbb{R}$,

$$\begin{cases} f(t) = af(t) \text{ for all } t \ge 0, a \in \mathbb{R} \\ f(0) = a_0 \end{cases}$$

is a linear ODE (Ordinary differential equation). The solution is: $f(t) = a_0 e^{at}$ for all $t \ge 0$.

Example 1.2 (Non-Linear ODE) $f: \mathbb{R} \to \mathbb{R}$

$$\begin{cases} f'(t) = 1 + f^2(t) \\ f(0) = 1 \end{cases}$$

Lets consider $f(t) = \tan(t) = \frac{\sin(t)}{\cos(t)}$. Then we have

$$f'(t) = \frac{1}{\cos(t)} = 1 + \tan^2(t) = 1 + f^2(t),$$

but this solution only is good in $(-\pi, \pi)$. It's a problem to extend this to $\mathbb{R} \to \mathbb{R}$.

A PDE (Partial Differential Equation) is an equation of a function of 2 or more variables and its derivatives.

Remark 1.3 Recall for $\Omega \subseteq \mathbb{R}^d$ open and $f: \Omega \to \{\mathbb{R}, \mathbb{C}\}$ the notation of partial derivatives:

- $\partial_{x_i} f(x) = \lim_{h \to 0} \frac{f(x+he_i) f(x)}{h}$, where $e_i = (0, 0, \dots, 1, \dots, 0, 0) \in \mathbb{R}^d$
- $D^{\alpha}f(x) = \partial_{x_1}^{\alpha_1} \cdots \partial_{x_d}^{\alpha_d} f(x)$, where $\alpha \in \mathbb{N}^d$
- $Df = \nabla f = (\partial x_1, \dots, \partial_{x_d})$
- $\Delta f = \partial_{x_1}^2 + \dots + \partial_{x_d}^2 f$
- $D^k f = (D^{\alpha} f)_{|\alpha| = k}$, where $|\alpha| = \sum_{i=1}^d |\alpha_i|$
- $D^2 f = (\partial_{x_i} \partial_{x_j} f)_{1 \le i, j \le d}$

Definition 1.4 Given a function F. Then the equation of the form

$$F(D^k u(x), D^{k-1} u(x), \dots, Du(x), u(x), x) = 0$$

with the unknown function $u: \Omega \subseteq \mathbb{R}^d \longrightarrow \mathbb{R}$ is called a *PDE of order k*.

- Equations $\sum_{d} a_{\alpha}(x)D^{\alpha}u(x) = 0$, where a_{α} and u are unknown functions are called *Linear PDEs*.
- Equations $\sum_{|\alpha|=k} a_{\alpha}(x) D^{\alpha} u(x) + F(D^{k-1}u, D^{k-2}u, \dots, Du, u, x) = 0$ are called semi-linear PDEs.

Goals: For solving a PDE we want to

- Find an explizit solution! This is in many cases impossible.
- Prove a well-posted theory (existence of solutions, uniqueness of solutions, continuous dependence of solutions on the data)

We have two notations of solutions:

- 1. Classical solution: The solution is continuous differentiable (e.g. $\Delta u = f \leadsto u \in C^2$)
- 2. Weak Solutions: The solution is not smooth/continuous

Definition 1.5 (Spaces of continous and differentiable functions) Let $\Omega \subseteq \mathbb{R}^d$ be open

$$\begin{split} C(\Omega) &= \{f: \ \Omega \to \mathbb{R} \mid f \text{ continuous} \} \\ C^k(\Omega) &= \{f: \ \Omega \to \mathbb{R} \mid D^\alpha f \text{ is continuous for all } |\alpha| \ \leq k \} \end{split}$$

Classical solution of a PDE of order $k \rightsquigarrow C^k$ solutions!

$$L^p(\Omega) = \left\{ f: \ \Omega \to \mathbb{R} \text{ lebesgue measurable } | \int_{\Omega} |f|^p d\lambda < \infty, 1 \le p < \infty \right\}$$

Sobolev Space:

$$W^{k,p}(\Omega) = \{ f \in L^p(\Omega) \mid \forall \alpha \in \mathbb{N}^n \text{ with } |\alpha| \le k : D^\alpha f \in L^p(\Omega) \text{ exists} \}$$

In this course we will investigate

- Laplace / Poisson Equation: $-\Delta u = f$
- Heat Equation: $\partial_t u \Delta u = f$
- Wave Equation: $\partial_t^2 \Delta u = f$
- Schrödinger Equation: $i\partial_t u \Delta u = f$

Chapter 2

Laplace / Poisson Equation

2.1 Laplace Equation

 $-\Delta u = 0$ (Laplace) or $-\Delta u = f(x)$ (Poisson).

Definition 2.1 (Harmonic Function) Let Ω be an open set in \mathbb{R}^d . If $u \in C^2(\Omega)$ and $\Delta u = 0$ in Ω , then u is a harmonic function in Ω .

Theorem 2.2 (Gauss-Green Theorem) Let $A \subseteq \mathbb{R}^d$ open, $\vec{F} \in C^1(A, \mathbb{R}^d)$ and $K \subseteq A$ compact with C^1 boundary. Then

$$\int_{\partial K} \vec{F} \cdot \vec{\nu} \ dS(x) = \int_K \operatorname{div}(\vec{F}) \ dx$$

where ν is the outward unit normal vector field on ∂K . Thus

$$\int_{\partial V} \nabla u \cdot \vec{\nu} \ dS(x) = \int_{V} \operatorname{div}(\nabla u) \ dx = \int_{V} \Delta u(x) \ dx$$

for any $V \subseteq \Omega$ open.

Theorem 2.3 (Green's Identities) Let $A \subseteq \mathbb{R}^d$ open, $K \subseteq A$ d-dim. compactum with C^1 boundary and $f, g \in C^2(A)$

1. Green's first identity (Partial Integration):

$$\int_{K} \nabla f \cdot \nabla g \, dx = \int_{\partial K} f \frac{\partial g}{\partial \nu} \, dS - \int_{K} f \Delta g \, dx$$

where $\frac{\partial g}{\partial \nu} = \partial_{\nu} g = \nu \cdot \nabla g$

2. Green's second identity:

$$\int_{K} f \Delta g - (\Delta f) g \, dx = \int_{\partial K} \left(f \frac{\partial g}{\partial \nu} - g \frac{\partial f}{\partial \nu} \right) \, dS$$

Exercise 2.4 Let $\Omega \subseteq \mathbb{R}^d$ open, let $f:\Omega \to \mathbb{R}$ be continuous. Prove that if $\int_B f(x) \ dx = 0$, then $u \equiv 0$ in Ω .

Theorem 2.5 (Fundamential Lemma of Calculus of Variations) Let $\Omega \subseteq \mathbb{R}^d$ open, let $f \in L^1(\Omega)$. If $\int_B f(x) \ dx = 0$ for all $x \in B_r(x) \subseteq \Omega$, then f(x) = 0 a.e. (almost everywhere) $x \in \Omega$.

Remark 2.6 (Solving Laplace Equation) $-\Delta u = 0$ in \mathbb{R}^d . Consider the case when u is radial, i.e. $u(x) = v(|x|), v : \mathbb{R} \to \mathbb{R}$. Denote r = |x|, then

$$\frac{\partial r}{\partial x} = \frac{\partial}{\partial x_i} \left(\sqrt{x_1^2 + \dots + x_d^2} \right) = \frac{2x_i}{2\sqrt{x_1^2 + \dots + x_d^2}} = \frac{x_i}{r}$$

Then

$$\begin{split} \partial_{x_i} u &= \partial_{x_i} v = (\partial_r v) \frac{\partial r}{\partial x_i} = v'(r) \frac{x_i}{r} \\ \partial_{x_i}^2 u &= \partial_{x_i} \left(v(r)' \frac{x_i}{r} \right) = (\partial_{x_i} v(r)') \frac{x_i}{r} + v'(r) \partial_{x_i} \left(\frac{x_i}{r} \right) \\ &= (\partial_r v'(r)) \left(\frac{dr}{\partial_{x_i}} \right) \frac{x_i}{r} + v'(r) \left(\frac{1}{r} - \frac{x_i}{r^2} (\partial_{x_i} r) \right) = v'(r) \frac{x_i^2}{r^2} + v'r(r) \left(\frac{1}{r} - \frac{x_i^2}{r^3} \right) \end{split}$$

So we have $\Delta u = \left(\sum_{i=1}^d d_{x_i}^2\right) u = v''(r) + v'(r) \left(\frac{d}{r} - \frac{1}{r}\right)$ Thus $\Delta u = v'(r) + v(r) \frac{d-1}{r}$. We consider $d \geq 2$. Laplace operator $\Delta u = 0$ now becomes $v''(r) + v'(r) \frac{d-1}{r} = 0$

becomes
$$v'(r) + v'(r) - \frac{v'(r)}{r} = 0$$

$$\Rightarrow \log(v(r))' = \frac{v'(r)}{v(r)} = -\frac{d-1}{r} = -(d-1)(\log r)' \text{ (recall } log(f)' = \frac{f'}{f})$$

$$\Rightarrow v'(r) = \frac{1}{v^{d-2} + \text{const.}}$$

$$\begin{cases} \frac{const}{r^{d-2}} + constxx + const & , d \ge 3 \\ const \log(r) + constxxr + const & , d = 2 \end{cases}$$

$$\begin{cases} \frac{const}{r^{d-2}} + constxx + const & , d \ge 3\\ const \log(r) + constxxr + const & , d = 2 \end{cases}$$

Definition 2.7 (Fundamential Solution of Laplace Equation)

$$\Phi(x) = \begin{cases} -\frac{1}{2\pi} \log(|x|), & d = 2\\ \frac{1}{(d-2)d|B_1|} \frac{1}{|x|^{d-2}}, & d \ge 3 \end{cases}$$

Where $|B_1|$ is the Volume of the ball $B_1(0) = B(0,1) \subseteq \mathbb{R}^d$.

Remark 2.8 $\Delta\Phi(x) = 0$ for all $x \in \mathbb{R}^d$ and $x \neq 0$.

2.2Poisson-Equation

The Poisson-Equation is $-\Delta u(x) = f(x)$ in \mathbb{R}^d . The explicit solution is given by

$$u(x) = (\Phi \star f)(x) = \int_{\mathbb{R}^d} \Phi(x - y) f(y) \ dy = \int_{\mathbb{R}^d} \Phi(y) f(x - y) \ dy$$

This can be heuristically justifyfied with

$$-\Delta(\Phi \star f) = (-\Delta\Phi) \star f = \delta_0 \star f = f$$

Theorem 2.9 Assume $f \in C_c^2(\mathbb{R}^d)$. Then $u = \Phi \star f$ satisfies that $u \in C^2(\mathbb{R}^d)$ and $-\Delta u(x) = f(x)$ for all $x \in \mathbb{R}^d$

Proof. By definition we have

$$u(x) = \int_{\mathbb{R}^d} \Phi(y) f(x - y) \, dy.$$

First we check that u is continuous: Take $x_k \to x_0$ in \mathbb{R}^d . We prove that $u(x_n) \xrightarrow{n} u_0$, i.e.

$$\lim_{n \to \infty} \int_{\mathbb{R}^d} \Phi(y) f(x_n - y) \ dy = \int_{\mathbb{R}^d} \Phi(y) f(x_0 - y) \ dy$$

This follows from the Dominated Convergence Theorem. More precisely:

$$\lim_{n \to \infty} \Phi(y) f(x_n - y) = \Phi(y) f(x_0 - y) \quad \text{ for all } y \in \mathbb{R}^d \setminus \{0\}$$

and

$$|\Phi(y)f(x-y)| \leq ||f||_{L^{\infty}} \cdot \mathbb{1}(|y| \leq R) \cdot |\Phi(y)| \in L^1(\mathbb{R}^d, dy)$$

where R > 0 depends on $\{x_n\}$ and supp(f) but independent of y. Now we compute the derivatives:

$$\begin{split} \partial_{x_i} u(x) &= \partial_{x_i} \int_{\mathbb{R}^d} \Phi(y) f(x-y) \ dy = \lim_{h \to 0} \int_{\mathbb{R}^d} \Phi(y) \frac{f(x+he_i-y) - f(x-y)}{h} \ dy \\ (\text{dom. conv.}) &= \int \Phi(y) \partial_{x_i} f(x-y) \ dy \\ \Rightarrow & D^{\alpha} u(x) = \int_{\mathbb{R}^d} \Phi(y) D_x^{\alpha} f(x-y) \ dy \quad \text{for all } |\alpha| \le 2 \end{split}$$

 $D^{\alpha}u(x)$ is continuous, thus $u\in C^2(\mathbb{R}^d)$. Now we check if this solves the Poisson-Equation:

$$-\Delta u(x) = \int_{\mathbb{R}^d} \Phi(y)(-\Delta_x) f(x-y) \, dy = \int_{\mathbb{R}^d} \Phi(y)(-\Delta_y) f(x-y) \, dy$$
$$= \int_{\mathbb{R}^d \setminus B(0,\epsilon)} \Phi(y)(-\Delta_x) f(x-y) \, dy + \int_{B(0,\epsilon)} \Phi(y)(-\Delta_x) f(x-y) \, dy \quad (\epsilon > 0 \text{ small})$$

Now we come to the main part. We apply integration by parts (2.3):

$$\int_{\mathbb{R}^{d} \backslash B(0,\epsilon)} \Phi(y)(-\Delta_{y}) f(x-y) \, dy$$

$$= \int_{\mathbb{R}^{d} \backslash B(0,\epsilon)} (\nabla_{y} \Phi(y)) \cdot \nabla_{y} f(x-y) \, dy - \int_{\partial B(0,\epsilon)} \Phi(y) \cdot \frac{\partial f}{\partial \vec{n}} (x-y) \, dS(y)$$

$$= \int_{\mathbb{R}^{d} \backslash B(0,\epsilon)} \underbrace{(-\Delta_{y} \Phi(y))}_{=0} f(x-y) \, dy$$

$$+ \int_{\partial B(0,\epsilon)} \frac{\partial \Phi}{\partial \vec{n}} (y) f(x-y) \, dS(y) - \int_{\partial B(0,\epsilon)} \Phi(y) \frac{\partial f}{\partial \vec{n}} (x-y) \, dS(y)$$

We have that $\nabla_y \Phi(y) = -\frac{1}{d|B_1|} \frac{y}{|y|^d}$ and $\vec{n} = \frac{y}{|y|}$ in $\partial B(0, \epsilon)$. This leads to

$$\frac{\partial \Phi}{\partial \vec{n}} = \frac{1}{d|B_1|} \frac{1}{|y|^{d-1}} = \frac{1}{d|B_1|\epsilon^{d-1}} \quad \text{ for } y \in \partial B(0, \epsilon)$$

Hence:

$$\int_{\partial B(0,\epsilon)} \frac{\partial \Phi}{\partial \vec{n}}(y) f(x-y) \ dS(y) = \frac{1}{d|B_1|\epsilon^{d-1}} \int_{\partial B(0,\epsilon)} f(x-y) \ dS(y)$$
$$= \oint_{\partial B(0,\epsilon)} f(x-y) \ dS(y) = \oint_{\partial B(x,\epsilon)} f(y) \ dS(y) \xrightarrow{\epsilon \to 0} f(x)$$

We have to regard the following error terms:

$$\left| \int_{B(0,\epsilon)} \Phi(y) (-\Delta_y) f(x-y) \ dy \right| \leq \int_{B(0,\epsilon)} |\Phi(y)| \underbrace{\left| -\Delta_y f(x-y) \right|}_{\leq \|\Delta f\|_{L^{\infty}} \mathbb{1}(|y| \leq R)} dy$$

$$\leq \|\Delta f\|_{L^{\infty}} \int_{\mathbb{R}^d} \underbrace{\left| \Phi(y) |\mathbb{1}(|y| \leq R)}_{L^1(\mathbb{R}^d)} \mathbb{1}(|y| \leq \epsilon) \xrightarrow{\epsilon \to 0} 0$$

Where R > 0 depends on x and the support of f but is independent of y.

$$\left| \int_{\partial B(0,\epsilon)} \Phi(y) \frac{\partial f}{\partial \vec{n}}(x - y) \ dS(y) \right| \leq \|\nabla f\|_{L^{\infty}} \int_{\partial B(0,\epsilon)} |\Phi(y)| \ dy$$

$$\leq \begin{cases} const \cdot \epsilon |\log \epsilon| \to 0, & d = 2\\ const \cdot \epsilon \to 0, & d \geq 3 \end{cases}$$

Conclusion: $-\Delta u(x) = f(x)$ for all $x \in \mathbb{R}^d$ proved that $u = \Phi \star f$ and $f \in C_c^2(\mathbb{R}^d)$.

Thus, if $f \in C_c^2(\mathbb{R})$, then $u = \Phi \star f$ satisfies $u \in C^2(\mathbb{R}^2)$ and $-\Delta u(x) = f(x)$ for all $x \in \mathbb{R}^d$.

Remark 2.10 The result holds for a much bigger class of functions f. For example if $f \in C_c^1(\mathbb{R})$ we can easily extend the previous proof:

$$\partial_{x_i} u = \int_{\mathbb{R}^d} \Phi(y) \partial_{x_i} f(x - y) \, dy \in C(\mathbb{R}^d) \Rightarrow u \in C^1(\mathbb{R}^d)$$

Consequently:

$$\partial_{x_i}\partial_{x_j}u = \partial_{x_i}\int_{\mathbb{R}^d} \Phi(y)\partial_{x_j}f(x-y)\,dy = \int_{\mathbb{R}^d} \partial_{x_i}\Phi(y)\partial_{x_j}f(x-y)\,dy \in C(\mathbb{R}^d)$$

So we have $u \in C^2(\mathbb{R}^d)$. Now we can compute

$$\Delta u = \sum_{i=1}^{d} \int_{\mathbb{R}^{d}} \partial_{x_{i}} \Phi(y) \partial_{x_{i}} f(x-y) \, dy \stackrel{(IBP)}{=} f(x).$$

Exercise 2.11 Extend this to more general functions!

2.3 Equations in general domains

Theorem 2.12 (Mean Value Theorem for Harmonic Functions) Let $\Omega \subseteq \mathbb{R}$ be open, let $u \in C^2(\Omega)$ and $\Delta u = 0$ in Ω . Then

$$u(x) = \int_{B(x,r)} u = \int_{\partial B(x,r)} u$$
 for all $x \in \Omega, B(x,r) \subseteq \Omega$

Exercise 2.13 In 1D: $\Delta u = 0 \Leftrightarrow u'' = 0 \Leftrightarrow u(x) = ax + b$ (Linear Equation)

Proof. (Of theorem)2.12 Consider all r > 0 s.t. $B(x,r) \subseteq \Omega$,

$$f(r) = \int_{\partial B(x,r)} u$$

We need to prove that f(r) is independent of r. When it is done, then we immediately obtain

$$f(r) = \lim_{t \to 0} f(t) = u(x)$$

as u is continuous. To prove that, consider

$$f'(r) = \frac{d}{dr} \left(\int_{\partial B(0,r)} u(x+y) \, dS(y) \right)$$

$$= \frac{d}{dr} \left(\int_{\partial B(0,1)} u(x+rz) \, dS(z) \right)$$

$$(\text{dom. convergence}) = \int_{\partial B(0,1)} \frac{d}{dr} [u(x+rz)] \, dS(z)$$

$$= \int_{\partial B(0,1)} \nabla u(x+rz) z \, dS(z)$$

$$= \int_{\partial B(x,r)} \nabla u(y) \frac{y-x}{r} \, dS(y)$$

$$= \frac{1}{|B(x,r)|_{\mathbb{R}^d}} \int_{\partial B(x,r)} \nabla \cdot u(y) \cdot \vec{n_y} \, dS(y)$$

$$(\text{Gauss-Green 2.2}) = \frac{1}{|B(x,r)|_{\mathbb{R}^d}} \int_{B(x,r)} \underbrace{(\Delta u)(y)}_{=0} \, dy = 0$$

Remark 2.14 Recall the polar decomposition. Let $x \in \mathbb{R}^d$, x = (r, w), r = |x| > 0, $\omega \in S^{d-1}$, then

$$\int_{B(0,r)} g(y) \, dy = \int_0^r \left(\int_{B(0,r)} g(y) \, dS(y) \right) dr$$

Remark 2.15 We already proved that for u harmonic we have $u(x) = f_{\partial B(x,r)} u \, dy$. Now we have

$$\int_{B(x,r)} u(y) \, dy = \int_{B(0,r)} u(x+y) \, dy$$

$$(\text{Pol. decomposition}) = \int_0^r \left(\int_{\partial B(0,s)} u(x+y) \, dS(y) \right) ds$$

$$= \int_0^r \left(\int_{\partial B(x,s)} u(y) \, dS(y) \right) ds$$

$$(\text{Mean value property}) = \int_0^r \left(|\partial B(x,s)| \, u(x) \right) ds = |B(x,r)| \, u(x)$$

This implies

$$\oint_{B(x,r)} u(y) \, dy = u(x) \quad \text{for any } B(x,r) \subseteq \Omega.$$

Remark 2.16 The reverse direction is also correct, namely if $u \in C^2(\Omega)$ and

$$u(x) = \int_{B(x,r)} u(y) \, dy = \int_{\partial B(x,r)} u(y) \, dy \quad \text{for all } B(x,r) \subseteq \Omega,$$

then u is harmonic, i.e. $\Delta u = 0$ in Ω . (The proof is exactly like before)

Exercise 2.17 (E 1.1) Proof the Gauss–Green formula: Let $f := (f_i)_1^d \in C^1(\mathbb{R}^d, \mathbb{R}^d)$. Prove that for every open ball $B(y, r) \subseteq \mathbb{R}^d$ we have

$$\int_{\partial B(y,r)} f(y) \cdot \nu_y \, dS(y) = \int_{B(y,r)} \operatorname{div} f \, dx.$$

Here ν_y is the outward unit normal vector and dS is the surface measure on the sphere.

My Solution (wrong). Let $\tilde{1} = (1, ..., 1) \in \mathbb{R}^d$.

$$\int_{B(y,r)} \operatorname{div}(f) \, dx = \int_{B(y,r)} \nabla f \cdot \tilde{1} \, dx$$

$$= \int_{B(y,r)} \nabla f \cdot \nabla x \, dx$$
(Partial Integration 2.3)
$$= \int_{\partial B(y,r)} f \cdot \frac{\partial x}{\partial \nu_x} \, dS(x) - \int_{B(y,r)} f \cdot \underbrace{\Delta x}_{=0} \, dx$$

$$= \int_{\partial B(y,r)} f \cdot \nu_x \cdot \underbrace{\nabla x}_{=\tilde{1}} \, dS(x)$$

$$= \int_{\partial B(y,r)} f \cdot \nu_x \, dS(x)$$

Theorem 2.18 (Maximum Principle) Let $\Omega \subseteq \mathbb{R}^d$ be open, let $u \in C^2(\Omega) \cap C(\overline{\Omega})$, $\Delta u = 0$ in Ω . Then

- 1. $\sup_{x \in \bar{\Omega}} u(x) = \sup_{x \in \partial \Omega} u(x)$
- 2. Assume that Ω is connected. Then if there is a $x_0 \in \Omega$ s.t. $u(x_0) = \sup_{x \in \overline{\Omega}} u(x)$, then $u \equiv const.$ in Ω .

Proof-Idea. Assume there exists $x_0 \in \Omega$ s.t. $u(x_0) = \sup_{x \in \overline{\Omega}} u(x)$. We have that $B(x_0, r) \subseteq \Omega$, so by the mean valum theorem we have

$$u(x_0) = \int_{B(x_0,r)} \underbrace{u(x)}_{\leq u(x_0)} \leq \int_{B(x_0,r)} u(x_0) \, dx = u(x_0)$$

So we get $u(x) = u(x_0)$ for all $x \in B(x_0, r)$.

Proof. Given $U \subseteq \mathbb{R}^d$ open, we can write $U = \bigcup_i U_i$, where U_i is open and connected.

- b) Assume that Ω is connected and there is a $x_0 \in \Omega$ s.t. $u(x_0) = \sup_{y \in \Omega} u(x)$. Define $U = \{x \in \Omega \mid u(x) = u(x_0)\} = u^{-1}(u(x_0))$. U is closed since u is continuous. Moreover, U is open by the mean-value theorem. I.e. for all $x \in U$ there is a r > 0 s.t. $B(x,r) \subseteq U \subseteq \Omega$. Since U is connected we get $U = \Omega$, so u is constant in Ω . On the other hand, if there is no $x_0 \in \Omega$ s.t. $u(x_0) = \sup_{x \in \Omega} u$ we have $\forall x_0 \in \Omega : u(x) < \sup_{x \in \overline{\Omega}} u(x) = \sup_{x \in \partial \Omega} u(x)$
- a) Given $\Omega \subseteq \mathbb{R}^d$ open, we can write $\Omega = \bigcup_i \Omega_i$, where Ω_i is open and connected. By b) we have

$$\sup_{x \in \bar{\Omega}_i} u(x) = \sup_{x \in \partial \Omega_i} u(x), \quad \forall i$$

$$\Rightarrow \quad \sup_{x \in \bar{\Omega}} u(x) = \sup_{x \in \partial \Omega} u(x)$$

Definition 2.19 • If $\Omega \subseteq \mathbb{R}^d$ is open, $u \in C^2(\Omega)$, then u is sub-harmonic if $\Delta u \geq 0$ in Ω .

• If $\Delta u \leq u$, then u is called *super-harmonic*.

Theorem 2.20 Let Ω be open in \mathbb{R}^d , $u \in C^2(\Omega)$, $\Delta u \geq 0$ in Ω .

1. We have the mean-value inequality

$$u(x) = \int_{B(x,r)} u \le \int_{\partial B(x,r)} u \quad \text{for all } x \in B(x,r) \subseteq \Omega$$

- 2. Assume that Ω is connected and bounded. Then either
 - u is a constant in Ω
 - $u(x) < \sup_{y \in \partial \Omega} u(y)$ for all $x \in \partial \Omega$

The Poisson Equation for given f, g on a bounded set is:

$$\begin{cases} -\Delta f = f, & \text{in } \Omega \\ u = g, & \text{on } \partial \Omega \end{cases}$$

Theorem 2.21 (Uniqueness) Let $\Omega \subseteq \mathbb{R}^d$ be bounded, open and connected. Let $f \in C(\Omega), g \in C(\partial\Omega)$. Then there exists at most one solution $u \in C^2(\Omega) \cap C(\bar{\Omega})$, s.t.

$$\begin{cases} -\Delta f = f, & \text{in } \Omega \\ u = g, & \text{on } \partial \Omega \end{cases}$$

Proof. Assume that we have two solutions u_1 and u_2 . Then $u := u_1 - u_2$ is a solution to By the maximum principle, we know that u = 0 in Ω . More precisely, by the maximum principle we have

$$\sup_{x\in\Omega}u(x)\leq \sup_{x\in\partial\Omega}u(x)=0\Rightarrow u(x)\leq 0 \forall x\in\Omega$$

Since -u satisfies the same property:

$$\sup_{x \in \Omega} (-u(x)) \le \sup_{x \in \partial \Omega} (-u(x)) = 0$$

$$\Rightarrow -u(x) \le 0 \text{ for all } x \in \Omega$$

$$\Rightarrow u(x) \ge 0 \text{ for all } x \in \Omega$$

So we geht u(x) = 0 in Ω .

Exercise 2.22 Let Ω be open, connected and bounded in \mathbb{R}^d . Let $u \in C^2(\Omega) \cap C(\bar{\Omega})$ s.t.

$$\begin{cases} \Delta u = 0, & \text{in } \Omega \\ u = g, & \text{on } \delta \Omega \end{cases}$$

Proof that

1. If $g \geq 0$ on $\partial \Omega$, then $u \geq 0$ in Ω .

2. If $g \geq 0$ on $\partial\Omega$ and $g \neq 0$, then u > 0 in Ω .

Lemma 2.23 (Estimates for derivatives) If u is harmonic in $\Omega \subseteq \mathbb{R}^d$ and $B(x_0, r) \subseteq \Omega$, then

$$|D^{\alpha}u(x_0)| \le \frac{(c_d N)^N}{r^{d+n}} \int_{B(x_0,r)} |u|$$

Theorem 2.24 (Regularity) Let Ω be open in \mathbb{R}^d . Let $u \in C(\Omega)$ satisfy $u(x) = \int_{\partial B} u$ for any $x \in B(x,r) \subseteq \Omega$. Then u is a harmonic function in Ω , namely $u \in C^2(\Omega)$ and $\Delta u = 0$ in Ω . Moreover, $u \in C^{\infty}(\Omega)$ and u is analytic in Ω .

Proof. $\int \eta = 1$. Define $\eta_{\epsilon}(x) = \epsilon^{-d} \eta(\epsilon^{-1} x)$ for all $\epsilon > 0$. Then

$$\int_{\mathbb{R}^d} \eta_{\epsilon} = \int_{\mathbb{R}^d} \eta = 1$$

We prove $u_{\epsilon} := \eta_{\epsilon} \star u = u$ for all $\epsilon > 0$. By definition:

$$u_{\epsilon}(x) = \int_{\mathbb{R}^d} \eta_{\epsilon}(x - y)u(y) \, dy = \int_0^{\infty} \left[\int_{\partial B(x, r)} \eta_{\epsilon}(x - y)u(y) \, dS(y) \right] dr$$
$$= \int_0^{\infty} \left[\int_{\partial B} n_{\epsilon}(r) \, dS(y) \right] dr = \int_0^{\infty} \left[\eta_{\epsilon}(r) |\partial B(x, r)| u(x) \right] dr$$
$$= u(x) \int_0^{\infty} \left[\eta_{\epsilon}(r) |\partial B(0, r)| \right] dr = u(x) \int_{\mathbb{R}^d} \eta_{\epsilon}(y) \, dy = u(x)$$

On the other hand, $u_{\epsilon} = \eta_{\epsilon} \star u$ is $C^{\infty}(\mathbb{R}^d)$. In fact $D^{\alpha}(\eta_{\epsilon} \star u) = (D^{\alpha}\eta_{\epsilon}) \star u$ is continuous for any α (Exercise). Then $u \in C^{\infty}(\mathbb{R}^d)$, so u is harmonic in \mathbb{R}^d , i.e. $\Delta u = 0$ in \mathbb{R}^d .

Consider now the general case where $\Omega \subseteq \mathbb{R}^d$ is open. Take $\epsilon > 0$ small and define $\Omega_{\epsilon} = \{x \in \Omega \mid \mathrm{dist}(x, \partial\Omega) > \epsilon\}$. Define

$$u_{\epsilon}(x) = \int_{\mathbb{R}^d} \eta_{\epsilon}(x - y)u(y) \ dy$$
 for all $x \in \Omega_{\epsilon}$

Recall that $\eta_{\epsilon}(y) = 0$ if $|y| \ge \epsilon$, then:

$$u_{\epsilon}(x) = \int_{B(x,\epsilon)} \eta_{\epsilon}(x-y)u(y) dy$$

is well-defined since $B(x,\epsilon)\subseteq \Omega$ for all $x\in \Omega_{\epsilon}$. Then by the same computation using the polar-decomposition, we find that $u_{\epsilon}(x)=u(x)$ for all $x\in \Omega$. Note that $u_{\epsilon}\in C^{\infty}(\Omega_{\epsilon})$, so $\Delta u=0$. Taking $\epsilon\to 0$, we get $u\in C^{\infty}(\Omega)$. Then we conclude that u is harmonic (We need to reverse the proof of the mean-value theorem). Analytically, we need to prove that for all $x_0\in \Omega$, there is a r>0 s.t. $B(x_0,r)\subseteq \Omega$ and

$$u(x) = u(x_0) + \sum_{\alpha \neq 0} c_{\alpha}(x - x_0)^{\alpha}$$
 for all $x \in B(x_0, r)$

Here $\alpha = (\alpha_1, \dots, \alpha_d), \alpha_i \in \{0, 1, 2, \dots\}$ and $y^{\alpha} = y_1^{\alpha_1} y_2^{\alpha_2} \dots y_d^{\alpha_d}$. We want to prove that the series converges uniformly in $B(x_0, r)$. Recall the Taylor expansion:

$$u(x) = u(x_0) + \sum_{0 < |\alpha| < N} D^{\alpha} u(x_0) \frac{(x - x_0)^{\alpha}}{\alpha!} + R_N(x)$$

where $|\alpha| = \alpha_1 + \alpha_2 + \cdots + \alpha_d$, $\alpha! = \alpha_1! \cdots \alpha_d!$ and

$$R_N(x) = \sum_{|\alpha|=N} \int_0^1 D^{\alpha} u(x_0 + t(x - x_0)) \frac{(x - x_0)^{\alpha}}{\alpha!} dt$$

We have: Take $x_0 \in \Omega$, take r > 0 small and $L = L_d > 0$ large, s.t. $B(x_0, (L+1)r) \subseteq \Omega$. Then for all $x \in B(x_0, r)$ we have $B(x, L_r) \subseteq B(x_0, (L+1)r) \subseteq \Omega$. With lemma 2.23 we get

$$|D^{\alpha}u(x)| \le \frac{(c_d N)^N}{(Lr)^{d+N}} \int_{B(x,Lr)} |u|$$

$$|R_N(x)| \leq \sum_{|\alpha|=N} ||D^{\alpha}u||_{L^{\infty}(B(x_0,r))} \frac{r^N}{\alpha!} \quad \text{for all } x \in B(x_0,r)$$

$$\leq \sum_{|\alpha|=N} \frac{(c_d N)^N}{(Lr)^{d+N}} \frac{r^N}{\alpha!} \int_{B(x_0,(L+1)r)} |u|$$

$$= \sum_{|\alpha|=N} \left(\frac{c_d N}{L}\right)^N \frac{1}{\alpha!} \underbrace{\frac{1}{(Lr)^d} \int_{B(x_0,(L+1)r)} |u|}_{X_0}$$

Multinomial theorem:

$$d^{N} = (1+1+\cdots+1)^{N} = \sum_{|\alpha|=N} \frac{N!}{\alpha!}$$

Conclusion:

$$|R_N(x)| \le \left(\frac{dC_dN}{L}\right)^N \frac{1}{N!}M$$

We need $N^N \ll N!$ (Stirly formula)