Partial Differerential Equations Thành Nam Phan Winter Semester 2021/2022

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Please note that I write this lecture notes for my personal use. I may write things differently than presented in the lecture. This script also contains some of my personal solutions for exercises (which may be wrong).

Chapter 1

Introduction

A differential equation is an equation of a function and its derivatives.

Example 1.1 (Linear ODE) Let $f: \mathbb{R} \to \mathbb{R}$,

$$\begin{cases} f(t) = af(t) \text{ for all } t \geqslant 0, a \in \mathbb{R} \\ f(0) = a_0 \end{cases}$$

is a linear ODE (Ordinary differential equation). The solution is: $f(t) = a_0 e^{at}$ for all $t \ge 0$.

Example 1.2 (Non-Linear ODE) $f : \mathbb{R} \to \mathbb{R}$

$$\begin{cases} f'(t) = 1 + f^2(t) \\ f(0) = 1 \end{cases}$$

Lets consider $f(t) = \tan(t) = \frac{\sin(t)}{\cos(t)}$. Then we have

$$f'(t) = \frac{1}{\cos(t)} = 1 + \tan^2(t) = 1 + f^2(t),$$

but this solution only is good in $(-\pi, \pi)$. It's a problem to extend this to $\mathbb{R} \to \mathbb{R}$.

A PDE (Partial Differential Equation) is an equation of a function of 2 or more variables and its derivatives.

Remark 1.3 Recall for $\Omega \subseteq \mathbb{R}^d$ open and $f: \Omega \to \{\mathbb{R}, \mathbb{C}\}$ the notation of partial derivatives:

- $\partial_{x_i} f(x) = \lim_{h \to 0} \frac{f(x+he_i) f(x)}{h}$, where $e_i = (0, 0, \dots, 1, \dots, 0, 0) \in \mathbb{R}^d$
- $D^{\alpha}f(x) = \partial_{x_1}^{\alpha_1} \cdots \partial_{x_d}^{\alpha_d} f(x)$, where $\alpha \in \mathbb{N}^d$
- $Df = \nabla f = (\partial x_1, \dots, \partial_{x_d})$
- $\bullet \ \Delta f = \partial_{x_1}^2 + \dots + \partial_{x_d}^2 f$
- $D^k f = (D^{\alpha} f)_{|\alpha|=k}$, where $|\alpha| = \sum_{i=1}^d |\alpha_i|$
- $D^2 f = (\partial_{x_i} \partial_{x_j} f)_{1 \leq i, j \leq d}$

Definition 1.4 Given a function F. Then the equation of the form

$$F(D^k u(x), D^{k-1} u(x), \dots, Du(x), u(x), x) = 0$$

with the unknown function $u: \Omega \subseteq \mathbb{R}^d \longrightarrow \mathbb{R}$ is called a *PDE of order k*.

- Equations $\sum_{d} a_{\alpha}(x) D^{\alpha} u(x) = 0$, where a_{α} and u are unknown functions are called *Linear PDEs*.
- Equations $\sum_{|\alpha|=k} a_{\alpha}(x) D^{\alpha} u(x) + F(D^{k-1}u, D^{k-2}u, \dots, Du, u, x) = 0$ are called semi-linear PDEs.

Goals: For solving a PDE we want to

- Find an explizit solution! This is in many cases impossible.
- Prove a well-posted theory (existence of solutions, uniqueness of solutions, continuous dependence of solutions on the data)

We have two notations of solutions:

- 1. Classical solution: The solution is continuous differentiable (e.g. $\Delta u = f \leadsto u \in C^2$)
- 2. Weak Solutions: The solution is not smooth/continuous

Definition 1.5 (Spaces of continous and differentiable functions) Let $\Omega \subseteq \mathbb{R}^d$ be open

$$\begin{split} C(\Omega) &= \{f: \ \Omega \to \mathbb{R} \mid f \text{ continuous} \} \\ C^k(\Omega) &= \{f: \ \Omega \to \mathbb{R} \mid D^\alpha f \text{ is continuous for all } |\alpha| \ \leqslant k \} \end{split}$$

Classical solution of a PDE of order $k \rightsquigarrow C^k$ solutions!

$$L^p(\Omega) = \left\{ f: \ \Omega \to \mathbb{R} \text{ lebesgue measurable } \left| \int_{\Omega} |f|^p d\lambda < \infty, \ 1 \leqslant p < \infty \right\}$$

Sobolev Space:

$$W^{k,p}(\Omega) = \{ f \in L^p(\Omega) \mid \forall \alpha \in \mathbb{N}^n \text{ with } |\alpha| \leq k : D^{\alpha} f \in L^p(\Omega) \text{ exists} \}$$

In this course we will investigate

- Laplace / Poisson Equation: $-\Delta u = f$
- Heat Equation: $\partial_t u \Delta u = f$
- Wave Equation: $\partial_t^2 \Delta u = f$
- Schrödinger Equation: $i\partial_t u \Delta u = f$

Chapter 2

Laplace / Poisson Equation

2.1 Laplace Equation

 $-\Delta u = 0$ (Laplace) or $-\Delta u = f(x)$ (Poisson).

Definition 2.1 (Harmonic Function) Let Ω be an open set in \mathbb{R}^d . If $u \in C^2(\Omega)$ and $\Delta u = 0$ in Ω , then u is a harmonic function in Ω .

Theorem 2.2 (Gauss-Green Theorem) Let $A \subseteq \mathbb{R}^d$ open, $\vec{F} \in C^1(A, \mathbb{R}^d)$ and $K \subseteq A$ compact with C^1 boundary. Then

$$\int_{\partial K} \vec{F} \cdot \vec{\nu} \ dS(x) = \int_K \operatorname{div}(\vec{F}) \ dx$$

where ν is the outward unit normal vector field on ∂K . Thus

$$\int_{\partial V} \nabla u \cdot \vec{\nu} \ dS(x) = \int_{V} \operatorname{div}(\nabla u) \ dx = \int_{V} \Delta u(x) \ dx$$

for any $V \subseteq \Omega$ open.

Theorem 2.3 (Green's Identities) Let $A \subseteq \mathbb{R}^d$ open, $K \subseteq A$ d-dim. compactum with C^1 boundary and $f, g \in C^2(A)$

1. Green's first identity (Integration by parts):

$$\int_{K} \nabla f \cdot \nabla g \, dx = \int_{\partial K} f \frac{\partial g}{\partial \nu} \, dS - \int_{K} f \Delta g \, dx$$

where $\frac{\partial g}{\partial \nu} = \partial_{\nu} g = \nu \cdot \nabla g$

2. Green's second identity:

$$\int_{K} f \Delta g - (\Delta f) g \, dx = \int_{\partial K} \left(f \frac{\partial g}{\partial \nu} - g \frac{\partial f}{\partial \nu} \right) \, dS$$

Exercise 2.4 Let $\Omega \subseteq \mathbb{R}^d$ open, let $f:\Omega \to \mathbb{R}$ be continuous. Prove that if $\int_B f(x) \ dx = 0$, then $u \equiv 0$ in Ω .

Theorem 2.5 (Fundamential Lemma of Calculus of Variations) Let $\Omega \subseteq \mathbb{R}^d$ open, let $f \in L^1(\Omega)$. If $\int_B f(x) \ dx = 0$ for all $x \in B_r(x) \subseteq \Omega$, then f(x) = 0 a.e. (almost everywhere) $x \in \Omega$.

Remark 2.6 (Solving Laplace Equation) $-\Delta u = 0$ in \mathbb{R}^d . Consider the case when u is radial, i.e. $u(x) = v(|x|), v : \mathbb{R} \to \mathbb{R}$. Denote r = |x|, then

$$\frac{\partial r}{\partial x} = \frac{\partial}{\partial x_i} \left(\sqrt{x_1^2 + \dots + x_d^2} \right) = \frac{2x_i}{2\sqrt{x_1^2 + \dots + x_d^2}} = \frac{x_i}{r}$$

Then

$$\begin{split} \partial_{x_i} u &= \partial_{x_i} v = (\partial_r v) \frac{\partial r}{\partial x_i} = v'(r) \frac{x_i}{r} \\ \partial_{x_i}^2 u &= \partial_{x_i} \left(v(r)' \frac{x_i}{r} \right) = (\partial_{x_i} v(r)') \frac{x_i}{r} + v'(r) \partial_{x_i} \left(\frac{x_i}{r} \right) \\ &= (\partial_r v'(r)) \left(\frac{dr}{\partial_{x_i}} \right) \frac{x_i}{r} + v'(r) \left(\frac{1}{r} - \frac{x_i}{r^2} (\partial_{x_i} r) \right) = v'(r) \frac{x_i^2}{r^2} + v'r(r) \left(\frac{1}{r} - \frac{x_i^2}{r^3} \right) \end{split}$$

So we have $\Delta u = \left(\sum_{i=1}^d d_{x_i}^2\right) u = v''(r) + v'(r)(\frac{d}{r} - \frac{1}{r})$ Thus $\Delta u = v'(r) + v(r)\frac{d-1}{r}$. We consider $d \ge 2$. Laplace operator $\Delta u = 0$ now becomes $v''(r) + v'(r)\frac{d-1}{r} = 0$

$$\Rightarrow \log(v(r))' = \frac{v'(r)}{v(r)} = -\frac{d-1}{r} = -(d-1)(\log r)' \text{ (recall } \log(f)' = \frac{f'}{f})$$

$$\Rightarrow v'(r) = \frac{1}{v^{d-2} + \text{const.}}$$

$$\begin{cases} \frac{const}{r^{d-2}} + constxx + const & , d \geqslant 3 \\ const \log(r) + constxxr + const & , d = 2 \end{cases}$$

Definition 2.7 (Fundamential Solution of Laplace Equation)

$$\Phi(x) = \begin{cases} -\frac{1}{2\pi} \log(|x|), & d = 2\\ \frac{1}{(d-2)d|B_1|} \frac{1}{|x|^{d-2}}, & d \geqslant 3 \end{cases}$$

Where $|B_1|$ is the Volume of the ball $B_1(0) = B(0,1) \subseteq \mathbb{R}^d$.

Remark 2.8 $\Delta\Phi(x) = 0$ for all $x \in \mathbb{R}^d$ and $x \neq 0$.

2.2Poisson-Equation

The Poisson-Equation is $-\Delta u(x) = f(x)$ in \mathbb{R}^d . The explicit solution is given by

$$u(x) = (\Phi \star f)(x) = \int_{\mathbb{R}^d} \Phi(x - y) f(y) \ dy = \int_{\mathbb{R}^d} \Phi(y) f(x - y) \ dy$$

This can be heuristically justifyfied with

$$-\Delta(\Phi \star f) = (-\Delta\Phi) \star f = \delta_0 \star f = f$$

Theorem 2.9 Assume $f \in C_c^2(\mathbb{R}^d)$. Then $u = \Phi \star f$ satisfies that $u \in C^2(\mathbb{R}^d)$ and $-\Delta u(x) = f(x)$ for all $x \in \mathbb{R}^d$

Proof. By definition we have

$$u(x) = \int_{\mathbb{R}^d} \Phi(y) f(x - y) \, dy.$$

First we check that u is continuous: Take $x_k \to x_0$ in \mathbb{R}^d . We prove that $u(x_n) \xrightarrow{n} u_0$, i.e.

$$\lim_{n \to \infty} \int_{\mathbb{R}^d} \Phi(y) f(x_n - y) \ dy = \int_{\mathbb{R}^d} \Phi(y) f(x_0 - y) \ dy$$

This follows from the Dominated Convergence Theorem. More precisely:

$$\lim_{n \to \infty} \Phi(y) f(x_n - y) = \Phi(y) f(x_0 - y) \quad \text{for all } y \in \mathbb{R}^d \setminus \{0\}$$

and

$$|\Phi(y)f(x-y)| \leq ||f||_{L^{\infty}} \cdot \mathbb{1}(|y| \leq R) \cdot |\Phi(y)| \in L^1(\mathbb{R}^d, dy)$$

where R > 0 depends on $\{x_n\}$ and supp(f) but independent of y. Now we compute the derivatives:

$$\begin{split} \partial_{x_i} u(x) &= \partial_{x_i} \int_{\mathbb{R}^d} \Phi(y) f(x-y) \ dy = \lim_{h \to 0} \int_{\mathbb{R}^d} \Phi(y) \frac{f(x+he_i-y) - f(x-y)}{h} \ dy \\ (\text{dom. conv.}) &= \int \Phi(y) \partial_{x_i} f(x-y) \ dy \\ \Rightarrow & D^{\alpha} u(x) = \int_{\mathbb{R}^d} \Phi(y) D_x^{\alpha} f(x-y) \ dy \quad \text{for all } |\alpha| \leqslant 2 \end{split}$$

 $D^{\alpha}u(x)$ is continuous, thus $u\in C^2(\mathbb{R}^d)$. Now we check if this solves the Poisson-Equation:

$$-\Delta u(x) = \int_{\mathbb{R}^d} \Phi(y)(-\Delta_x) f(x-y) \, dy = \int_{\mathbb{R}^d} \Phi(y)(-\Delta_y) f(x-y) \, dy$$
$$= \int_{\mathbb{R}^d \setminus B(0,\epsilon)} \Phi(y)(-\Delta_x) f(x-y) \, dy + \int_{B(0,\epsilon)} \Phi(y)(-\Delta_x) f(x-y) \, dy \quad (\epsilon > 0 \text{ small})$$

Now we come to the main part. We apply integration by parts (2.3):

$$\int_{\mathbb{R}^{d} \backslash B(0,\epsilon)} \Phi(y)(-\Delta_{y}) f(x-y) \, dy$$

$$= \int_{\mathbb{R}^{d} \backslash B(0,\epsilon)} (\nabla_{y} \Phi(y)) \cdot \nabla_{y} f(x-y) \, dy - \int_{\partial B(0,\epsilon)} \Phi(y) \cdot \frac{\partial f}{\partial \vec{n}} (x-y) \, dS(y)$$

$$= \int_{\mathbb{R}^{d} \backslash B(0,\epsilon)} (-\Delta_{y} \Phi(y)) f(x-y) \, dy$$

$$+ \int_{\partial B(0,\epsilon)} \frac{\partial \Phi}{\partial \vec{n}} (y) f(x-y) \, dS(y) - \int_{\partial B(0,\epsilon)} \Phi(y) \frac{\partial f}{\partial \vec{n}} (x-y) \, dS(y)$$

We have that $\nabla_y \Phi(y) = -\frac{1}{d|B_1|} \frac{y}{|y|^d}$ and $\vec{n} = \frac{y}{|y|}$ in $\partial B(0, \epsilon)$. This leads to

$$\frac{\partial \Phi}{\partial \vec{n}} = \frac{1}{d|B_1|} \frac{1}{|y|^{d-1}} = \frac{1}{d|B_1|\epsilon^{d-1}} \quad \text{for } y \in \partial B(0, \epsilon)$$

Hence:

$$\int_{\partial B(0,\epsilon)} \frac{\partial \Phi}{\partial \vec{n}}(y) f(x-y) \ dS(y) = \frac{1}{d|B_1|\epsilon^{d-1}} \int_{\partial B(0,\epsilon)} f(x-y) \ dS(y)$$
$$= \int_{\partial B(0,\epsilon)} f(x-y) \ dS(y) = \int_{\partial B(x,\epsilon)} f(y) \ dS(y) \xrightarrow{\epsilon \to 0} f(x)$$

We have to regard the following error terms:

$$\left| \int_{B(0,\epsilon)} \Phi(y) (-\Delta_y) f(x-y) \, dy \right| \leq \int_{B(0,\epsilon)} |\Phi(y)| \underbrace{\left| -\Delta_y f(x-y) \right|}_{\leq \|\Delta f\|_{L^{\infty}} \mathbb{1}(|y| \leq R)} \, dy$$

$$\leq \|\Delta f\|_{L^{\infty}} \int_{\mathbb{R}^d} \underbrace{\left| \Phi(y) |\mathbb{1}(|y| \leq R) \right|}_{L^1(\mathbb{R}^d)} \mathbb{1}(|y| \leq \epsilon) \xrightarrow{\epsilon \to 0} 0$$

Where R > 0 depends on x and the support of f but is independent of y.

$$\bullet \left| \int_{\partial B(0,\epsilon)} \Phi(y) \frac{\partial f}{\partial \vec{n}}(x-y) \ dS(y) \right| \leq \|\nabla f\|_{L^{\infty}} \int_{\partial B(0,\epsilon)} |\Phi(y)| \ dy$$

$$\leq \begin{cases} const \cdot \epsilon |\log \epsilon| \to 0, & d = 2\\ const \cdot \epsilon \to 0, & d \geqslant 3 \end{cases}$$

Conclusion: $-\Delta u(x) = f(x)$ for all $x \in \mathbb{R}^d$ proved that $u = \Phi \star f$ and $f \in C^2_c(\mathbb{R}^d)$.

Thus, if $f \in C_c^2(\mathbb{R})$, then $u = \Phi \star f$ satisfies $u \in C^2(\mathbb{R}^2)$ and $-\Delta u(x) = f(x)$ for all $x \in \mathbb{R}^d$.

Remark 2.10 The result holds for a much bigger class of functions f. For example if $f \in C_c^1(\mathbb{R})$ we can easily extend the previous proof:

$$\partial_{x_i} u = \int_{\mathbb{R}^d} \Phi(y) \partial_{x_i} f(x - y) \, dy \in C(\mathbb{R}^d) \Rightarrow u \in C^1(\mathbb{R}^d)$$

Consequently:

$$\partial_{x_i}\partial_{x_j}u = \partial_{x_i}\int_{\mathbb{R}^d} \Phi(y)\partial_{x_j}f(x-y)\,dy = \int_{\mathbb{R}^d} \partial_{x_i}\Phi(y)\partial_{x_j}f(x-y)\,dy \in C(\mathbb{R}^d)$$

So we have $u \in C^2(\mathbb{R}^d)$. Now we can compute

$$\Delta u = \sum_{i=1}^{d} \int_{\mathbb{R}^d} \partial_{x_i} \Phi(y) \partial_{x_i} f(x-y) \, dy \stackrel{(IBP)}{=} f(x).$$

Exercise 2.11 Extend this to more general functions!

2.3 Equations in general domains

Theorem 2.12 (Mean Value Theorem for Harmonic Functions) Let $\Omega \subseteq \mathbb{R}$ be open, let $u \in C^2(\Omega)$ and $\Delta u = 0$ in Ω . Then

$$u(x) = \int_{B(x,r)} u = \int_{\partial B(x,r)} u \text{ for all } x \in \Omega, B(x,r) \subseteq \Omega$$

Proof. Consider all r > 0 s.t. $B(x, r) \subseteq \Omega$,

$$f(r) = \int_{\partial B(x,r)} u$$

We need to prove that f(r) is independent of r. When it is done, then we immediately obtain

$$f(r) = \lim_{t \to 0} f(t) = u(x)$$

as u is continuous. To prove that, consider

$$f'(r) = \frac{d}{dr} \left(\int_{\partial B(0,r)} u(x+y) \, dS(y) \right)$$

$$= \frac{d}{dr} \left(\int_{\partial B(0,1)} u(x+rz) \, dS(z) \right)$$

$$(\text{dom. convergence}) = \int_{\partial B(0,1)} \frac{d}{dr} [u(x+rz)] \, dS(z)$$

$$= \int_{\partial B(0,1)} \nabla u(x+rz) z \, dS(z)$$

$$= \int_{\partial B(x,r)} \nabla u(y) \frac{y-x}{r} \, dS(y)$$

$$= \frac{1}{|B(x,r)|_{\mathbb{R}^d}} \int_{\partial B(x,r)} \nabla \cdot u(y) \cdot \vec{n_y} \, dS(y)$$

$$(\text{Gauss-Green 2.2}) = \frac{1}{|B(x,r)|_{\mathbb{R}^d}} \int_{B(x,r)} \underbrace{(\Delta u)(y)}_{=0} \, dy = 0$$

Exercise 2.13 In 1D: $\Delta u = 0 \Leftrightarrow u'' = 0 \Leftrightarrow u(x) = ax + b$ (Linear Equation)

Remark 2.14 Recall the polar decomposition. Let $x \in \mathbb{R}^d$, x = (r, w), r = |x| > 0, $\omega \in S^{d-1}$, then

$$\int_{B(0,r)} g(y) \, dy = \int_0^r \left(\int_{B(0,r)} g(y) \, dS(y) \right) dr$$

Remark 2.15 We already proved that for u harmonic we have $u(x) = f_{\partial B(x,r)} u \, dy$. Now we have

$$\int_{B(x,r)} u(y) \, dy = \int_{B(0,r)} u(x+y) \, dy$$
(Pol. decomposition)
$$= \int_0^r \left(\int_{\partial B(0,s)} u(x+y) \, dS(y) \right) ds$$

$$= \int_0^r \left(\int_{\partial B(x,s)} u(y) \, dS(y) \right) ds$$
(Mean value property)
$$= \int_0^r \left(|\partial B(x,s)| \, u(x) \right) ds = |B(x,r)| \, u(x)$$

This implies

$$\oint_{B(x,r)} u(y) \, dy = u(x) \quad \text{for any } B(x,r) \subseteq \Omega.$$

Remark 2.16 The reverse direction is also correct, namely if $u \in C^2(\Omega)$ and

$$u(x) = \int_{B(x,r)} u(y) \, dy = \int_{\partial B(x,r)} u(y) \, dy \quad \text{for all } B(x,r) \subseteq \Omega,$$

then u is harmonic, i.e. $\Delta u = 0$ in Ω . (The proof is exactly like before)

Theorem 2.17 (Maximum Principle) Let $\Omega \subseteq \mathbb{R}^d$ be open, let $u \in C^2(\Omega) \cap C(\bar{\Omega})$, $\Delta u = 0$ in Ω . Then

- a) $\max_{x \in \bar{\Omega}} u(x) = \max_{x \in \partial \Omega} u(x)$
- b) Assume that Ω is connected. Then if there is a $x_0 \in \Omega$ s.t. $u(x_0) = \max_{x \in \bar{\Omega}} u(x)$, then $u \equiv const.$ in Ω .

Proof. Given $U \subseteq \mathbb{R}^d$ open, we can write $U = \bigcup_i U_i$, where U_i is open and connected.

- b) Assume that Ω is connected and there is a $x_0 \in \Omega$ s.t. $u(x_0) = \sup_{y \in \Omega} u(x)$. Define $U = \{x \in \Omega \mid u(x) = u(x_0)\} = u^{-1}(u(x_0))$. U is closed since u is continuous. Moreover, U is open by the mean-value theorem. I.e. for all $x \in U$ there is a r > 0 s.t. $B(x,r) \subseteq U \subseteq \Omega$. Since U is connected we get $U = \Omega$, so u is constant in Ω . On the other hand, if there is no $x_0 \in \Omega$ s.t. $u(x_0) = \sup_{x \in \Omega} w$ we have $\forall x_0 \in \Omega : u(x) < \sup_{x \in \overline{\Omega}} u(x) = \sup_{x \in \partial \Omega} u(x)$
- a) Given $\Omega \subseteq \mathbb{R}^d$ open, we can write $\Omega = \bigcup_i \Omega_i$, where Ω_i is open and connected. By b) we have

$$\sup_{x \in \bar{\Omega}_i} u(x) = \sup_{x \in \partial \Omega_i} u(x), \quad \forall i$$

So we can conclude

$$\sup_{x \in \bar{\Omega}} u(x) = \sup_{x \in \partial\Omega} u(x).$$

Definition 2.18 • If $\Omega \subseteq \mathbb{R}^d$ is open, $u \in C^2(\Omega)$, then u is called *sub-harmonic* if $\Delta u \ge 0$ in Ω .

• If $\Delta u \leq 0$, then u is called *super-harmonic*.

Exercise 2.19 (E 1.4) Let $\Omega \subseteq \mathbb{R}^d$ be open and $u \in C^2(\Omega)$ be subharmonic.

a) Prove that u satisfies the Mean Value Inequality

$$\oint_{\partial B(x,r)} u(y) \, dS(y) \geqslant \oint_{B(x,r)} u(y) \, dy \geqslant u(x)$$

for all $B(x,r) \subseteq \mathbb{R}^d$.

- b) Assume further that Ω is connected and $u \in C(\bar{\Omega})$. Prove that u satisfies the strong maximum principle, namely either
 - u is constant in Ω , or
 - $\sup_{y \in \partial \Omega} u(y) > u(x)$ for all $x \in \Omega$.

My Solution. a) Let $f(r) = \int_{\partial B(x,r)} u(y) dS(y)$, then we have

$$\partial_{r} f(r) = \partial_{r} \oint_{\partial B(x,r)} u(y) \, dS(y)$$
(Dom. Convergence)
$$= \oint_{\partial B(x,r)} \partial_{r} u(y) \, dS(y)$$

$$= \oint_{\partial B(0,1)} \partial_{r} u(x+yr) \, dS(y)$$

$$= \oint_{\partial B(0,1)} \nabla u(x+yr) \cdot y \, dS(y)$$

$$= \oint_{\partial B(x,r)} \nabla u(y) \cdot \frac{y-x}{r} \, dS(y)$$

$$= \oint_{\partial B(x,r)} \nabla u(y) \cdot \vec{n}_{y} \, dS(y)$$
(Gauss-Green)
$$= \oint_{B(x,r)} \operatorname{div}(\nabla u(y)) \, dS(y)$$

$$= \oint_{B(x,r)} \underbrace{\Delta u(y)}_{\geqslant 0} \, dS(y) \geqslant 0$$

So we can conclude that

$$\oint_{\partial B(x,r)} u(y) \, dS(y) = f(r) \geqslant \lim_{r \to 0} f(r) = u(x).$$

Now regard

$$\int_{B(x,r)} u(y) \, dy = \int_0^r \left(\int_{\partial B(x,r)} u(y) \, dS(y) \right) ds$$

$$= \int_0^r \left(|\partial B(x,r)| \int_{\partial B(x,r)} u(y) \, dS(y) \right) ds$$

$$\geqslant \int_0^r |\partial B(x,r)| \cdot u(x) \, dS(y)$$

$$= u(x) \int_0^r |\partial B(x,r)| \, dS(y) = u(x) |B(x,r)|.$$

Thus we have

$$u(x) \leqslant \int_{B(x,r)} u(y)dy.$$

Finally, lets regard

$$\int_{B(x,r)} u(y) \, dy = \int_0^r \left(|\partial B(x,s)| \oint_{\partial B(x,s)} u(y) \, dS(y) \right) \, ds$$

$$(\partial_r f(r) \geqslant 0) \qquad \leqslant \int_0^r \left(|\partial B(x,s)| \oint_{\partial B(x,r)} u(y) \, dS(y) \right) \, ds$$

$$= \oint_{\partial B(x,r)} u(y) \, dS(y) \int_0^r |\partial B(x,s)| \, ds$$

$$= \oint_{\partial B(x,r)} u(y) \, dS(y) \cdot |B(x,s)|$$

and we conclude

$$\int_{B(x,r)} u(y) \, dy \leqslant \int_{\partial B(x,r)} u(y) \, dS(y).$$

b) Let $x_0 \in \Omega$ s.t. $u(x_0) = \sup_{x \in \Omega} u(x)$. Now,

$$\sup_{x \in \Omega} u(x) = u(x_0) \leqslant \int_{\partial B(x_0, r)} u(y) \, dy$$
$$\leqslant \int_{\partial B(x_0, r)} \sup_{x \in \Omega} u(x) \, dy = \sup_{x \in \Omega} u(x)$$

Since u is continuous we get $u(y) = u(x_0)$ for all $y \in B(x_0, r)$, so u is constant.

Definition 2.20 The *Poisson Equation* for given f, g on a bounded set is:

$$\begin{cases}
-\Delta u = f, & \text{in } \Omega \\
u = g, & \text{on } \partial\Omega
\end{cases}$$

Theorem 2.21 (Uniqueness) Let $\Omega \subseteq \mathbb{R}^d$ be bounded, open and connected. Let $f \in C(\Omega), g \in C(\partial\Omega)$. Then there exists at most one solution $u \in C^2(\Omega) \cap C(\bar{\Omega})$, s.t.

$$\begin{cases} -\Delta u = f, & \text{in } \Omega \\ u = g, & \text{on } \partial \Omega \end{cases}$$

Proof. Assume that we have two solutions u_1 and u_2 . Then $u := u_1 - u_2$ is a solution to

$$\begin{cases} -\Delta u = 0, & \text{in } \Omega \\ u = 0 & \text{on } \partial \Omega \end{cases}$$

By the maximum principle, we know that u=0 in Ω . More precisely, by the maximum principle we have $\forall x\in\Omega$

$$\sup_{x \in \Omega} u(x) \leqslant \sup_{x \in \partial \Omega} u(x) = 0 \quad \Rightarrow \quad u(x) \leqslant 0$$

Since -u satisfies the same property we have $\forall x \in \Omega$:

$$\sup_{x \in \Omega} (-u(x)) \leqslant \sup_{x \in \partial \Omega} (-u(x)) = 0 \quad \Rightarrow \quad -u(x) \leqslant 0 \quad \Rightarrow \quad u(x) \geqslant 0$$

So we geht u(x) = 0 in Ω .

Exercise 2.22 (Bonus 1) Let Ω be open, connected and bounded in \mathbb{R}^d . Let $u \in C^2(\Omega) \cap C(\bar{\Omega})$ s.t.

$$\begin{cases} \Delta u = 0, & \text{in } \Omega \\ u = g, & \text{on } \partial \Omega \end{cases}$$

Prove that

a) If $g \ge 0$ on $\partial \Omega$, then $u \ge 0$ in Ω .

b) If $g \ge 0$ on $\partial \Omega$ and $g \ne 0$, then u > 0 in Ω .

Lemma 2.23 (Estimates for derivatives) If u is harmonic in $\Omega \subseteq \mathbb{R}^d$, $\alpha \in \mathbb{N}_0^d$, $|\alpha| = N$ and $B(x_0, r) \subseteq \Omega$, then

$$|D^{\alpha}u(x)| \le \frac{(c_d N)^N}{r^{d+N}} \int_{B(x,r)} |u| \, dy$$

Proof. Induction: Assume $|\alpha| = N - 1$, Take $|\alpha| = N$

$$|D^{\alpha}u(x_0)| \le \frac{|S_1|}{|B_1|\frac{r}{N}} \|D^{\beta}u\|_{L^{\infty}(B(x_0,\frac{r}{n}))}, \quad D^{\alpha}u = \partial_{x_i}(D^{\beta}u)_{|\beta|=N-1}$$

Note: $x \in B(x_0, \frac{r}{N})$, so $B(x, \frac{r(N-1)}{N}) \subseteq B(x_0, r)$. By the induction hypothesis:

$$||D^{\beta}u||_{L^{\infty}(B(x_{0},\frac{r}{N}))} \leq \frac{[c_{d}(N-1)]^{N-1}}{[r^{\frac{(N-1)}{N}}]^{d+N-1}} \int_{B(x_{0},r)} |u| \, dy$$

The conclusion is:

$$|D^{\alpha}u(x_{0})| \leq \frac{|S_{1}|}{|B_{1}|\frac{r}{N}} \frac{\left[c_{d}(N-1)\right]^{N-1}}{\left(r\frac{N-1}{N^{d}}\right)^{d+N-1}} \int_{B(x_{0},r)} |u| \, dy$$

$$= \frac{|S_{1}|}{|\beta_{1}|} \frac{c_{d}^{N-1}}{\left(\frac{r}{N}\right)^{d+N} (N-1)^{d}} \int_{B(x_{0},r)} |u| \, dy$$

$$= \frac{|S_{1}|}{|\beta_{1}|} \frac{c_{d}^{N-1}}{\left(\frac{r}{N}\right)^{d+N} N^{d}} \left(\frac{N}{N-1}\right)^{d} \int_{B(x_{0},r)} |u| \, dy$$

$$\leq \frac{2^{d}|S_{1}|}{|B_{1}|} \frac{c_{d}^{N-1} N^{N}}{r^{d+N}} \int_{B(x_{0},r)} |u| \, dy \quad \text{if } c_{d} \geq \frac{2^{d}|S_{1}|}{|B_{1}|}$$

Theorem 2.24 (Regularity) Let Ω be open in \mathbb{R}^d . Let $u \in C(\Omega)$ satisfy $u(x) = \int_{\partial B} u \, dy$ for any $x \in B(x, r) \subseteq \Omega$. Then u is a harmonic function in Ω . Moreover, $u \in C^{\infty}(\Omega)$ and u is analytic in Ω .

Exercise 2.25 (E 1.1: Proof the Gauss–Green formula) Let $f := (f_i)_1^d \in C^1(\mathbb{R}^d, \mathbb{R}^d)$. Prove that for every open ball $B(y, r) \subseteq \mathbb{R}^d$ we have

$$\int_{\partial B(y,r)} f(y) \cdot \nu_y \, dS(y) = \int_{B(y,r)} \operatorname{div} f \, dx.$$

Here ν_y is the outward unit normal vector and dS is the surface measure on the sphere.

Solution. We proof this in d=3. Let $f \in C^1(\mathbb{R}^3)$

$$\int_{B(0,1)} \partial_{x_3} f \, dx = \int_{\partial B(0,1)} f x_3 \, dS(x), \quad x = (x_1, x_2, x_3) \in \mathbb{R}^3, \vec{n} = \frac{x}{|x|} \text{ on } \partial B(0,1)$$

$$B(0,1) = \{x_1^2 + x_2^2 + x_3^2 \le 1\}$$

= $\{x_1^2 + x_2^2 \le 1 - \sqrt{1 - x_1^2 - x_2^2} \le x_3 \le \sqrt{1 - x_1^2 - x_2^2}\}$

Then:

$$\int_{B(0,1)} \partial_{x_3} f \, dx = \int_{x_1^2 + x_2^2 \le 1} \left(\int_{-\sqrt{1 - x_1^2 - x_2^2} \le x_3 \le \sqrt{1 - x_1^2 - x_2^2}} \partial_{x_3} f \, dx_3 \right) \, dx_1 \, dx_2$$

$$= \int_{x_1^2 + x_2^2 \le 1} \left[f(x_1, x_2, \sqrt{1 - x_1^2 - x_2^2}) - f(x_1, x_2, -\sqrt{1 - x_1^2 - x_2^2}) \right] \, dx_1 \, dx_2$$

Lets take polar coordinates in 2D:

$$x_1 = r \cos \phi$$
 $r > 0, \phi \in [0, 2\pi)$
 $x_2 = r \sin \phi$ $\det \frac{\partial (x_1, x_2)}{\partial (r, \phi)} = r$

$$(\star) = \int_0^1 \int_0^{2\pi} [f(r\cos\phi, r\sin\phi, r) - f(r\cos\theta, r\sin\phi, -r)] r \, dr \, d\phi$$

On the other hand:

$$\int_{\partial B(0,1)} fx_3 \, dS$$

The polar coordinates in 3D are:

$$x_1 = r \cos \phi \sin \theta$$
 $r > 0, \phi \in (0, 2\pi), \theta \in (0, \pi)$
 $x_2 = r \sin \phi \sin \theta$ $\det \frac{\partial x_1, x_2, x_3}{\partial (r, \phi, t)} = r^2 \sin \theta$
 $x_3 = \cos \theta$

Then:

$$(\star\star) = \int_0^{2\pi} \int_0^{\pi} f(\cos\phi\sin\theta, \sin\phi\sin\theta, \cos\theta) \sin\theta\cos\theta \, d\theta \, d\phi$$

$$= \int_0^{2\pi} \left(\int_0^{\frac{\pi}{2}} + \int_{\frac{\pi}{2}}^{\pi} d\theta \right) \, d\phi$$

$$(r = \sin\theta) = \int_0^{2\pi} \int_0^1 f(r\cos\phi, r\sin\phi, \sqrt{1 - r^2}) r \, dr \, d\phi$$

$$- f(r\cos\phi, r\sin\phi, -\sqrt{1 - r^2}) r \, dr \, d\phi$$

Exercise 2.26 (E 1.2) Let $u \in C(\mathbb{R}^d)$ and $\int_{B(x,r)} u \, dy = 0$ for every open ball $B(x,r) \subseteq \mathbb{R}^d$. Show that u(x) = 0 for all $x \in \mathbb{R}^d$.

My Solution. Assume there is a $x_0 \in \mathbb{R}^d$ s.t. w.l.o.g. $u(x_0) > 0$. Since u is continous there is a ball $B(x_0, r)$ s.t. $u(y) > \frac{u(x_0)}{2}$ for all $y \in B(x_0, r)$. But then we get

$$\int_{B(x_0,r)} u(y) \, dy \geqslant \int_{B(x_0,r)} \frac{u(x_0)}{2} \, dy = \frac{u(x_0)}{2} \, |B(x_0,r)| > 0.$$

Exercise 2.27 (E 1.3) Let $f \in C_c^1(\mathbb{R}^d)$ with $d \ge 2$ and $u(x) := (\Phi \star f)(x)$. Prove that $u \in C^2(\mathbb{R}^2)$ and $-\Delta u(x) = f(x)$ for all $x \in \mathbb{R}^d$ (2.9 was the same for $f \in C_1(\mathbb{R})$)

Theorem 2.28 (Liouville's Theorem) If $u \in C^2(\mathbb{R}^d)$ is harmonic and bounded, then u = const.

Proof. By the bound of the derivative 2.23 we have

$$\begin{aligned} |\partial_{x_i} u(x_0)| &\leqslant \frac{c_d}{r^{d+1}} \int_{B(x_0, r)} |u| \, dy \quad \forall x_0 \in \mathbb{R}^d \, \forall r > 0 \\ &\leqslant \|u\|_{L^{\infty}} \frac{c_d}{r^{d+1}} |B(x_0, r)| \\ &\leqslant \|u\|_{L^{\infty}} \frac{c_d}{r} \xrightarrow{r \to \infty} 0 \end{aligned}$$

Thus $\partial_{x_i} u = 0$ for all $i = 1, 2, \dots d$ and u = const. in \mathbb{R}^d

Theorem 2.29 (Uniqueness of solutions to Poisson Equation in \mathbb{R}^d) If $u \in C^2(\mathbb{R}^d)$ is a bounded function and satisfies $-\Delta u = f$ in \mathbb{R}^d where $f \in C_c^2(\mathbb{R}^d)$, then we have

$$u(x) = \Phi \star f(x) + C = \int_{\mathbb{R}^d} \Phi(x - y) f(y) \, dy + C \quad \forall x \in \mathbb{R}^d$$

where C is a constant and Φ is the fundamental solution of the Laplace equation in \mathbb{R}^d .

Proof. If we can prove that v is bounded, then v = const.. We first need to show that $\Phi \star f$ is bounded.

$$\Phi = \Phi_1 + \Phi_2 = \Phi\mathbb{1}(|x| \leqslant 1) + \Phi(|x| \geqslant 1)$$
$$\Phi \star f = \Phi_1 \star f + \Phi_2 \star f$$

We have $\Phi_1 \star f \in L^1(\mathbb{R}^d)$ and $\Phi_2 \star f$ is bounded since $\Phi \to 0$ as $|x| \to \infty$ in $d \ge 3$.

Exercise 2.30 (Hanack's inequality) Let $u \in C^2(\mathbb{R}^d)$ be harmonic and non-negative. Prove that for all open, bounded and connected $\Omega \subseteq \mathbb{R}^d$, we have

$$\sup_{x \in \Omega} u(x) \leqslant C_{\Omega} \inf_{x \in \Omega} u(x),$$

where C_{∞} is a finite constant depending only on Ω .

Proof. (Exercise) Hint: $\Omega = B(x, r)$. General case cover Ω by finitely many balls, one ball is inside Ω .

Chapter 3

Convolution, Fourier Transform and Distributions

3.1 Convolutions

Definition 3.1 (Convolution) Let $f, g : \mathbb{R}^d \to \mathbb{R}$ or \mathbb{C} .

$$(f \star g)(x) = \int_{\mathbb{R}^d} f(x - y)g(y) \, dy = \int_{\mathbb{R}^d} f(y)g(x - y) \, dy = (g \star f)(x)$$

Remark 3.2 (Properties of the Convolution) • $(f \star g)(x) = f \star (g \star h)$

•
$$\hat{f} \star g = \hat{f} \star \hat{g}$$

Theorem 3.3 (Young Inequality) If $f \in L^1(\mathbb{R}^d)$ and $g \in L^p(\mathbb{R}^d)$, where $1 \leq p \leq \infty$, then $f \star g \in L^p(\mathbb{R}^d)$ and $\|f \star g\|_{L^p} \leq \|f\|_{L^1} \|g\|_{L^p}$. More generally, if $f \in L^p(\mathbb{R}^d)$, $g \in L^q(\mathbb{R}^d)$, then $f \star g \in L^1(\mathbb{R}^d)$, $\|f \star g\|_{L^1} \leq \|f\|_{L^p} \|g\|_{L^q}$, where $1 \leq p, q, r, \leq \infty$, $\frac{1}{p} + \frac{1}{q} = 1 + \frac{1}{r}$

Proof. Let $f \in L^1, g \in L^p$. With the Hölder Inequality ??, we have:

$$||f \star g||_{L^{p}}^{p} = \int_{\mathbb{R}^{d}} |f \star g(x)|^{p} dx$$

$$\leq ||f||_{L^{1}}^{\frac{p}{q}} \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} |f(x - y)||g(y)|^{p} dy dx$$

$$= ||f||_{L^{1}}^{\frac{p}{q} + 1} ||g||_{L^{p}}^{p}$$

So we have $||f \star g||_{L^p} \leq ||f||_{L^1} ||g||_{L^p}$

Theorem 3.4 (Smoothness of the Convolution) If $f \in C_c^{\infty}(\mathbb{R}^d)$, $g \in L^p(\mathbb{R}^d)$, $1 \le p \le \infty$. Then $f \star g \in C^{\infty}(\mathbb{R})$ and

$$D^{\alpha}(f \star g) = (D^{\alpha}f) \star g$$

for all $\alpha = (\alpha_1, \ldots, \alpha_d), \alpha_i \in \{0, 1, 2, \ldots\}$

Proof. First we note that $x \mapsto (f \star g)$ is continous as $x_n \to x$ in \mathbb{R}^d since

$$(f \star g)(x_n) = \int_{\mathbb{R}^d} f(x_n - y)g(y) \, dy \xrightarrow{\text{dom. conv.}} \int_{\mathbb{R}^d} f(x - y)g(y) \, dy = (f \star g)(x)$$

We can apply Dominated convergence because

$$f(x_n - y)g(y) \to f(x - y)g(y) \quad \forall y \text{ as } f \text{ is continuous and } x_n \to x$$

and

$$|f(x_n - y) \ g(y)| \le ||f||_{L^{\infty}} |g(y)| \ \mathbb{1}(|y| \le R) \in L^1(\mathbb{R}^d).$$

Where R > 0 satisfies $B(0,R) \supseteq \operatorname{supp} f + \operatorname{sup}_n |x_n|$. Now we can compute the derivatives:

$$\partial_{x_i}(f \star g)(x) = \lim_{h \to 0} \frac{(f \star g)(x + he_i) - (f \star g)(x)}{h}$$

$$= \lim_{h \to 0} \int_{\mathbb{R}^d} \frac{f(x + he_i - y) - f(x - y)}{h} g(y) \, dy$$
(Dominated Convergence)
$$= \int_{\mathbb{R}^d} \lim_{h \to 0} \frac{f(x + he_i - y) - f(x - y)}{h} g(y) \, dy$$

$$= \int_{\mathbb{R}^d} (\partial_{x_i} f)(x - y) g(y) \, dy$$

We could apply Dominated Convergence since

$$\frac{f(x+he_i-y)-f(x-y)}{h}g(y) \xrightarrow{h\to 0} (\partial_{x_i}f)(x-y)g(y) \quad \text{as } f \in C^1$$

$$\left|\frac{f(x+he_i-y)-f(x-y)}{h}g(y)\right| \leqslant \|\partial_{x_i}f\|_{L^\infty}|g(y)| \ \mathbb{1}(|y|\leqslant R) \in L^1(\mathbb{R}^d)$$

where $B(0,R) \supseteq \operatorname{supp}(f) + B(0,|x|+1)$ and $\partial_{x_i}(f \star g) = (\partial_{x_i}f) \star g \in C(\mathbb{R}^d)$ since $\partial_{x_i}f \in C_c^{\infty}(\mathbb{R}^d)$. By induction we get $D^{\alpha}(f \star g) = (D^{\alpha}f \star g) \in C(\mathbb{R}^d)$.

Remark 3.5 Question: Is there a f s.t. $f \star g = g$ for all g. In fact there is no regular function f that solves this formally:

$$f \star g = g \Rightarrow \widehat{f \star g} = \widehat{g} \Rightarrow \widehat{f}\widehat{g} = \widehat{g} \Rightarrow \widehat{f} = 1 \Rightarrow f \text{ is not a regular function!}$$

However, if f is the Dirac-Delta Distribution, $f = \delta_0$ then $\delta_0 \star g = g$ for all g. Formally:

$$\delta_0(x) = \begin{cases} 0 & x \neq 0 \\ \infty & x = 0 \\ \int \delta_0 = 1 \end{cases}$$

In fact, if $f \in L^1(\mathbb{R}^d)$, $\int f = 1$, $f_{\epsilon}(x) = \epsilon^{-d} f(\epsilon^{-1} x)$, then $f_{\epsilon} \to \delta_0$ in an appropriate sense and $f_{\epsilon} \star g \to g$ for all g nice enough.

Theorem 3.6 (Approximation by convolution) Let $f \in L^1(\mathbb{R}^d)$, $\int f = 1$, $f_{\epsilon}(x) = \epsilon^{-d} f(\frac{x}{\epsilon})$. Then for all $g \in L^p(\mathbb{R}^d)$, where $1 \leq p < \infty$, then

$$f_{\epsilon} \star g \to g \quad \text{in } L^p(\mathbb{R}^d)$$

Proof.

Step 1: Let $f, g \in C_c(\mathbb{R}^d)$. Then

$$(f_{\epsilon} \star g)(x) - g(x) = \int_{\mathbb{R}^{d}} f_{\epsilon}(y)g(x - y) \, dy - \int_{\mathbb{R}^{d}} f_{\epsilon}(y)g(x) \, dy$$

$$= \int_{\mathbb{R}^{d}} f_{\epsilon}(y)(g(x - y) - g(x)) \, dy$$

$$|(f_{\epsilon} \star g)(x) - g(x)| = \left| \int_{\mathbb{R}^{d}} f_{\epsilon}(y)(g(x - y) - g(x)) \, dy \right|$$

$$\leqslant \int_{\mathbb{R}^{d}} |f_{\epsilon}(y)||g(x - y) - g(x)| \, dy$$

$$\leqslant \int_{|y| \leqslant R_{\epsilon}} |f_{\epsilon}(y)||g(x - y) - g(x)| \, dy$$

$$\leqslant \int_{|y| \leqslant R_{\epsilon}} |f_{\epsilon}(y)| \, dy \left[\sup_{|z| \leqslant R} |g(x - z) - g(x)| \right] \xrightarrow{\epsilon \to 0} 0$$

We have Dominated Convergence since:

$$(f_{\epsilon} \star g)(x) - g(x) \to 0 \text{ as } \epsilon \to 0$$

and

$$|f_{\epsilon} \star g(x) - g(x)| \leqslant \|f\|_{L^{1}} \sup_{|z| \leqslant R_{\epsilon}} |g(x - z) - g(x)| \leqslant 2\|f\|_{1} \|g\|_{L^{\infty}} \mathbb{1}(|x| \leqslant R_{1}).$$

Where $B(0, R_1) \supseteq \operatorname{supp}(g) + B(0, R_{\epsilon})$, thus $f_{\epsilon} \star g \to g$ in $L^p(\mathbb{R}^d)$. To remove the technical assumptions $f, g \in C_c(\mathbb{R}^d)$, then we use a density argument. We use the fact that $C_c(\mathbb{R}^d)$ is dense in $L^p(\mathbb{R}^d)$, $1 \le p < \infty$.

Step 2: Let $g \in C_c(\mathbb{R}^d)$, $g \in L^p(\mathbb{R}^d)$. Then there is $\{g_m\} \subseteq L^p(\mathbb{R}^d)$, $g_m \to g$ in $L^p(\mathbb{R}^d)$.

$$\begin{split} \|f_{\epsilon} \star g - g\|_{L^{p}} &\leq \|f_{\epsilon} \star (g - g_{m})\|_{L^{p}} + \|f_{\epsilon} \star g_{m} - g_{m}\|_{L^{p}} + \|g_{m} - g\|_{L^{p}} \\ & (\text{Young}) &\leq \|f_{\epsilon}\|_{L^{1}} \|g - g_{m}\|_{L^{p}} + \|f_{\epsilon} \star g_{m} - g_{m}\|_{L^{p}} + \|g_{m} - g\|_{L^{p}} \\ &\leq \|f\|_{L^{1}} \|g - g_{m}\|_{L^{p}} + \|f_{\epsilon} \star g_{m} - g_{m}\|_{L^{p}} + \|g_{m} - g\|_{L^{p}} \\ &\leq (\|f\|_{L^{1}} + 1)\|g - g_{m}\|_{L^{p}} + \|f \star g_{m} - g_{m}\|_{L^{p}} \end{split}$$

So we get:

$$\limsup_{\epsilon \to 0} \|f_{\epsilon} \star g - g\|_{L^{p}} \leqslant (\|f\|_{L^{p}} + 1)\|g - g_{m}\|_{L^{p}} + \limsup_{\epsilon \to 0} \|f_{\epsilon} \star g_{m} - g_{m}\|_{L^{p}}$$

$$\underbrace{\lim\sup_{\epsilon \to 0} \|f_{\epsilon} \star g - g\|_{L^{p}}}_{\text{Oby step 1.}}$$

$$\xrightarrow{m\to\infty} 0$$

Step 3: Let $f \in L^1(\mathbb{R}^d)$ and $g \in L^p(\mathbb{R}^d)$. Take $\{f_m\} \subseteq C_c(\mathbb{R}^d)$, s.t.

$$\begin{cases} F_m \to g \in L^1(\mathbb{R}) \text{ as } m \to \infty \\ \int_{\mathbb{R}^d} F_m = 1 \text{ (it is possible since } \int_{\mathbb{R}^d}) f = 1) \end{cases}$$

Define $F_{m,\epsilon}(x) = \epsilon^{-d} F_m(\epsilon^{-1} x)$ (recall $f_{\epsilon}(x) = \epsilon^{-d} f(\epsilon^{-1} x)$). Then:

$$\begin{split} f_{\epsilon} \star g - g &= (f_{\epsilon} - F_{m,\epsilon}) \star g + F_{m,\epsilon} \star g - g \\ \Rightarrow \|f_{\epsilon} - g\|_{L^{p}} &\leq \underbrace{\|f_{\epsilon} - F_{m,\epsilon} \star g\|_{L^{p}}}_{+} + \|F_{m,\epsilon} \star g - g\|_{L^{p}} \end{split}$$

$$\underbrace{\text{Young}}_{\leqslant} \|f_{\epsilon} - F_{m,\epsilon}\|_{L^{1}} \|g\|_{L^{p}} = \|f - F_{m}\|_{L^{1}} \|g\|_{L^{p}}$$

$$\Rightarrow \limsup_{\epsilon \to 0} \|f_{\epsilon} \star g - g\|_{L^{p}} \leq \|f - F_{m}\|_{L^{1}} \|g\|_{L^{p}} = \|f - F_{m}\|_{L^{1}} \|g\|_{L^{p}}$$

Lemma 3.7 $C_c(\mathbb{R}^d)$ is dense in $L^p(\mathbb{R}^d)$, $1 \leq p < \infty$

Proof. For all $g \in L^p(\mathbb{R}^d)$ there are g_m step functions and $g_m \to m$ in $L^p(\mathbb{R}^d)$, We can assume that Ω is open and bounded and we want to approximate χ_{Ω} by $C_c(\mathbb{R}^d)$.

Lemma 3.8 (Urnson) Define

$$\Omega_{\epsilon} = \{ x \in \Omega \mid \operatorname{dist}(x, \partial \Omega) > \epsilon \}$$

Then there is a $\eta_{\epsilon} \in C_c(\mathbb{R}^d)$ s.t.

$$\begin{cases} 0 \leqslant \eta(x) \leqslant 1 & \forall x \in \mathbb{R}^d \\ \eta_{\epsilon}(x) = 1 & \text{if } x \in \Omega_{\epsilon} \\ \eta_{\epsilon}(x) = 0 & \text{if } x \notin \Omega \end{cases}$$

Lemma 3.9 (Gernal Version of Urnson) If $A, B \subseteq \mathbb{R}^d$, A closed, B closed, $A \cap B = \emptyset$. Then

$$\eta(x) = \frac{\operatorname{dist}(x, A)}{\operatorname{dist}(x, A) + \operatorname{dist}(x, B)}$$

Then $\eta \in C(\mathbb{R}^d)$, $0 \leq \eta \leq 1$ and $\eta = 0$ if $x \in B$, $\eta = 1$ if $x \in A$. App to $A = \overline{\Omega_{\epsilon}} \subset\subset \Omega$ and $B = \mathbb{R}^d \setminus \Omega$.

Theorem 3.10 (Appendix C4 in Evans) Let Ω be open in \mathbb{R}^d and for $\epsilon > 0$ define

$$\Omega_{\epsilon} = \{ x \in \Omega \mid \operatorname{dist}(x, \mathbb{R}^d \setminus \Omega) > \epsilon \}$$

Let $f \in C_c^{\infty}(\mathbb{R}^d)$, $\int_{\mathbb{R}^d} f = 1$, supp $f \subseteq B(0,1)$, $f_{\epsilon}(x) = \epsilon^{-d} f(\epsilon^{-1} x)$ supp is $B(0,\epsilon)$. Then for all $g \in L^p_{loc}(\Omega)$ (i.e. $\mathbb{1}_K g \in L^p(\Omega) \forall K$ compakt set in Ω), then:

- a) $g_{\epsilon}(x) = (f_{\epsilon} \star g)(x) = \int_{\mathbb{R}^d} f_{\epsilon}(x y)g(y) dy \int_{\Omega} f_{\epsilon}(x y)g(y) dy$ is well-defined in Ω_{ϵ} and $g_{\epsilon} \in C^{\infty}(\Omega_{\epsilon})$.
- b) $g_{\epsilon} \to g$ in $L^p_{loc}(\Omega)$ if $1 \leq p < \infty$ and $g_{\epsilon}(x) \to g(x)$ almost everywhere $x \in \Omega$.
- c) If $g \in C(\Omega)$, then $g_{\epsilon}(x) \to g(x)$ uniformly in any compact subset of Ω .

Proof. a) $D^{\alpha}(g_{\epsilon}) = (D^{\alpha}f_{\epsilon}) \star g \in C(\Omega_{\epsilon})$

b) Already proved in \mathbb{R}^d space.

Corrolary 3.11 (Lebesgue differentiation theorem) If $f \in L_{loc}^P(\mathbb{R}^d)$, then

$$\oint_{B(x,\epsilon)} |f(y) - f(x)|^p dy \to 0 \quad \text{as } \epsilon \to 0$$

Exercise 3.12 (E 2.1) Let $u \in C^2(\mathbb{R}^2)$ be convex. I.e.

$$tu(x) + u(y)(1-t) \ge u(tx + (1-t)y) \forall x, y \in \mathbb{R}^d \forall t \in [0,1]$$

a) Prove for all $x \in \mathbb{R}^d$ that H(x) = ...

Solution.

a In 1D: If u is convex $\Leftrightarrow u''(x) \ge 0$ for all $x \in \mathbb{R}$. In general: Taylor expansion for all $x, z \in \mathbb{R}^d$:

$$u(x) = u(z) + \nabla u(z)(x - y) + \int_0^1 \sum_{|\alpha| = 2} D^{\alpha} u(z + s(x - z)) \frac{(x - z)^{\alpha}}{\alpha!} ds$$

$$x = z + s(x - z), s = 1$$
 Use $z = tx + (t - 1)y \Rightarrow x - z = (1 - t)(x - y)$

$$tu(x) = tu(z) + t\nabla u(z)(1-t)(x-y) + t\int_0^1 \sum_{|\alpha|=2} D^{\alpha}u(z+s(x-z)) \frac{[(1-t)(x-y)]^{\alpha}}{\alpha!} ds$$

$$(1-t)u(y) = (1-t)u(z) + (1-t)\nabla u(z)t(y-x) + (1-t)\int_0^r \sum_{|\alpha|=2} D^{\alpha}u(z+s(y-z))\frac{[t(y-x)]^{\alpha}}{\alpha!} ds$$

$$\Rightarrow tu(x) + (1-t)u(y) = u(z) + t \int_0^1 \dots + (1-t) \int_0^1 \dots$$
$$\Rightarrow t \int_0^1 \dots + (1-t) \int_0^1 \dots \geqslant 0 \forall x, y, t, z = tx + (1-t)y$$

$$t(1-t)^2 \int_0^1 \sum_{|\alpha|=2} D^{\alpha} u(z+s(x-z)) \frac{(x-y)}{\alpha!} \, ds + (1-t)t^2 \int_0^1 \sum_{|\alpha|=2} D^{\alpha} u(z+s(y-z)) \frac{(y-z)^{\alpha}}{\alpha!} \, ds \geqslant 0$$

for all $x, y \in \mathbb{R}^d$, $t \in [0, 1]$, z = tx + (1 - t)y. Divides for t(1 - t)

$$(1-t)\int_0^1\cdots+\int_0^1\cdots\geqslant 0$$

Take $t \to 0$

$$\int_0^1 \sum_{|\alpha|=2} D^{\alpha} u(y + s(x - y)) \frac{(x - y)^{\alpha}}{\alpha!} ds \geqslant 0 \forall x, y \in \mathbb{R}^d$$

Take $y = x + a, a \in \mathbb{R}^d$

$$\int_0^1 \sum_{|\alpha|=2} D^{\alpha} u(x+a+sa) \frac{a^{\alpha}}{\alpha!} ds \geqslant 0 \forall \epsilon > 0, \forall x, a \in \mathbb{R}^d$$

Take $\epsilon \to 0$

$$\int_0^1 \sum_{|\alpha|=2} D^\alpha u(x) \frac{a^\alpha}{\alpha!} \geqslant 0 \Rightarrow \sum_{i,j=1,i\neq j} \partial_{x_i} \partial_{x_j} u(x) a_i a_j + \sum_{i=j=1}^d \partial_{x_i}^2 u(x) \frac{a_i^2}{2}$$

We get

$$\frac{1}{2}a^T H a \geqslant 0 \forall a (a_i)_{i=1}^d \in \mathbb{R}^d$$

b
$$H(x) \geqslant 0 \Rightarrow (\partial_i \partial_j u) \geqslant 0 \Rightarrow TrH(x) \geqslant 0 \Rightarrow \sum_{i=1}^d \partial_{x_i}^2 u(x) \geqslant 0 \Rightarrow \Delta u(x) \geqslant 0 \forall x \in \mathbb{R}^d$$

Exercise 3.13 (E 2.2)

Solution. Regard d=3. The function $\frac{1}{|x|}$ is harmonic in $\mathbb{R}^3\setminus\{0\}$. We prove

$$\oint_{\partial B(x,r)} \frac{dS(y)}{|y|} = \frac{1}{\max(|x|,r)}$$

If |x| > r, then $0 \notin B(x, r + \epsilon)$. Then

$$y \mapsto \frac{1}{|y|}$$

is harmonic in $B(x, r + \epsilon)$. Then by the Mean Value Property:

$$\oint_{\partial B(x,r)} \frac{dS(y)}{|y|} = \frac{1}{|x|}$$

If |x| < r: Then $\frac{1}{|y|}$ is not harmonic in B(x,r) since $0 \in B(x,r)$. Note

$$\oint_{\partial B(x,r)} \frac{dS(y)}{|y|} = \oint_{\partial B(0,r)} \frac{dS(y)}{|x-y|}$$

This function depends on x only via |x|

$$\dots = \oint_{\partial B(0,r)} \frac{dS(y)}{|Rx - Ry|}$$

for all R rotation SO(3), $dS(R_y) = dS(y)$

$$= \int_{\partial B(0,r)} \frac{dS(y)}{|Rx - y|}$$

$$= \int_{\partial B(0,r)} \frac{dS(y)}{|z - y|}$$
(Radial in z)
$$= \int_{\partial B(0,|x|)} \left(\int_{\partial B(0,|x|)} \frac{dS(y)}{|z - y|} \right) dS(z)$$
(Fubini)
$$= \int_{\partial B(0,r)} \left(\int_{\partial B(0,|x|)} \frac{dS(z)}{|z - y|} \right) dS(y)$$
(case 1 since $|y| = r > |x|$)
$$= \int_{\partial B(0,r)} \frac{1}{|y|} dS(y) = \frac{1}{r}$$

If |x| = r: Continuity: $x \mapsto f_{\partial B(0,r)} \frac{dS(y)}{|x-y|}$

Remark 3.14 For $f \in C^{|\alpha|}, g \in C^{|\beta|}$:

$$D^{\alpha+\beta}(f\star g)=(D^{\alpha}f)\star(D^{\beta}g)$$

Lemma 3.15 If $d \ge 3$ and $f : \mathbb{R}^d \to \mathbb{R}$ radial. Then:

$$\left(\frac{1}{|x|^{d-2}} \star f\right)(x) = \int_{\mathbb{R}^d} \frac{f(y)}{|x-y|^{d-2}} \, dy$$
$$= \int_{\mathbb{R}^d} \frac{f(y)}{\max(|x|^{d-2}, |y|^{d-2})} \, dy$$

Proof. (d=3) Polar coordinates:

$$\int_{\mathbb{R}^3} \frac{f(y)}{|x-y|} \, dy = \int_0^\infty \left[\int_{\partial B(0,1)} \frac{1}{|x-rw|} \, d\omega \right] f(r) \, dr$$

$$(a) = \int_0^\infty \left[\int_{\partial B(0,1)} \frac{d\omega}{\max(|x|,r)} \right] f(r) \, dr$$

$$= \int_{\mathbb{R}^3} \frac{f(y)}{\max(|x|,|y|)} \, dy$$

(b) (d=3) If f radial and non-negative

$$\int_{\mathbb{R}^3} \frac{f(y)}{|x - y|} = \int_{\mathbb{R}^3} \frac{f(y)}{|x|} \, dy = \frac{(Sf?)}{|x|}$$

Then

$$\int_{\mathbb{R}^{3}} \int_{\mathbb{R}^{3}} \frac{f_{1}(x - z_{1}) f_{2}(y - z_{2})}{|x - y|} dx dy = \int_{\mathbb{R}^{3}} \int_{\mathbb{R}^{3}} \frac{f_{1}(x) f_{2}(y)}{|x + z_{1} - y - z_{2}|} dx dy$$

$$= \int_{\mathbb{R}^{3}} \left(\int_{\mathbb{R}^{3}} f_{1}(x) dx \right) f_{2}(y) dy \leqslant \int_{\mathbb{R}^{3}} \frac{\left(\int_{\mathbb{R}^{3}} f_{1} \right)}{|y + z_{2} - z_{1}|} f_{2}(y) dy$$

$$\leqslant \frac{\left(\int_{\mathbb{R}^{3}} f_{1} \right) \left(\int_{\mathbb{R}^{3}} f_{2} \right)}{|z_{1} - z_{2}|}$$

Exercise 3.16 (Bonus 2) a) Prove that $u(x) = \frac{1}{|x|}$ is sub-harmonic in $\mathbb{R}^2 \setminus \{0\}$.

b) Prove that if $f: \mathbb{R}^2 \to \mathbb{R}$ radial, non-negative, measurable:

$$\int_{\mathbb{R}^2} \frac{f(y)}{|x - y|} \, dy \geqslant \int_{\mathbb{R}^2} \frac{f(y)}{\max(|x|, |y|)} \, dy$$

3.2 Fourier Transformation

Definition 3.17 (Fourier Transform) For $f \in L^1(\mathbb{R}^d)$ define

$$\mathcal{F}f(k) = \hat{f}(k) = \int_{\mathbb{R}^d} f(x)e^{-2\pi i k \cdot x} dx, \quad k \cdot x = \sum_{i=1}^d k_i x_i$$

Theorem 3.18 (Basic Properties) 1. If $f \in L^1(\mathbb{R}^d)$, then $\hat{f} \in L^{\infty}(\mathbb{R}^d)$ and $\|\hat{f}\|_{L^{\infty}} \leq \|f\|_{L^1}$

2. For all $f \in L^1(\mathbb{R}^d) \cap L^2(\mathbb{R}^d)$, $\|\hat{f}\|_{L^2} = \|f\|_{L^2}$. Moreover, \mathcal{F} can be extended to be a unitary transforamtion $L^2(\mathbb{R}^d) \to L^2(\mathbb{R}^d)$ s.t.

$$\|\mathcal{F}g\|_{L^2} = \|f\|_{L^2} \quad \forall f \in L^2(\mathbb{R}^d)$$

3. The inverse of F can be defined as

$$(F^{-1}f)(x) = \check{f}(x) = \int_{\mathbb{R}^d} f(x)e^{2\pi ikx} dk$$

for all $f \in L^1(\mathbb{R}^d) \cap L^2(\mathbb{R}^d)$

4.
$$\widehat{D^{\alpha}f}(k) = (2\pi i k)^{\alpha} \widehat{f}(k)$$
 as $(2\pi i k)^{\alpha} f(k) \in L^2(\mathbb{R}^d)$ $(k^{\alpha} = k_1^{\alpha_1} \cdots k_{\alpha}^{\alpha_k})$

5.
$$\widehat{f \star g}(k) = \widehat{f}(k)\widehat{g}(k)$$
 if f, g are nice enough.

Theorem 3.19 (Hausdorff-Young-Inequality) If $1 \le p \le 2$, $\frac{1}{p} + \frac{1}{p'} = 1$ and $f \in L^p(\mathbb{R}^d) \cap L^1(\mathbb{R}^n)$ then

$$\|\hat{f}\|_{L^{p'}} \leqslant \|f\|_{L^p}$$

and

$$\|\hat{f}\|_{L^p} \leqslant \|f\|_{L^p} \quad \forall f \in L^p(\mathbb{R}^d)$$

Remark 3.20 We want to apply the Fourier transform to find the solution of a PDE, e.g. the Poisson-Equation:

$$-\Delta u = f \text{ in } \mathbb{R}^d \Rightarrow |2\pi k|^2 \hat{u}(k) = \hat{f}(k) \Rightarrow \hat{u}(k) = \frac{1}{|2\pi k|^2} \hat{f}(k)$$

If we can find G s.t. $\hat{G}(k) = \frac{1}{|2\pi k|^2}$, then

$$\hat{u}(k) = \hat{G}(k)\hat{f}(k) = \widehat{G \star f}$$

$$\Rightarrow u(x) = (G \star f)(x) = \int_{\mathbb{R}^d} G(x - y)f(y) \, dy$$

In fact G is the fundamential solution of laplace quation.

Theorem 3.21 (Fourier Transform of $\frac{1}{|x|^{\alpha}}$ for $0 < \alpha < d$) We have formally

$$\widehat{\frac{c_{\alpha}}{|x|^{\alpha}}} = \frac{c_{d-\alpha}}{|k|^{d-\alpha}} \quad \forall \ 0 < \alpha < d$$

Here

$$c_{\alpha} = \pi^{-\frac{d}{2}} \Gamma\left(\frac{\alpha}{2}\right) = \pi^{-\frac{\alpha}{2}} \int_{0}^{\infty} e^{-\lambda} \lambda^{\frac{\alpha}{2} - 1} d\lambda$$

More precisely, for all $f \in C_c^{\infty}(\mathbb{R}^d)$,

$$\frac{c_{\alpha}}{|x|^{\alpha}} \star f = \left(\frac{c_{d-\alpha}}{|k|^{d-\alpha}} \hat{f}(k)\right)^{\vee}$$

Moreover if $\alpha > \frac{d}{2}$, then we also have

$$\left(\frac{c_{\alpha}}{|x|^{\alpha}} \star f\right)^{\wedge} = \frac{c_{d-\alpha}}{|f|^{d-\alpha}} \hat{f}(k)$$

Lemma 3.22 (Fourier Transform of Gaussians) In \mathbb{R}^d ,

$$\widehat{e^{-\pi|x|^2}} = e^{-\pi|k|^2}$$

More generally for all $\lambda > 0$:

$$\widehat{e^{-\pi\lambda^2|x|^2}} = \lambda^{-d}e^{-\pi\frac{|x|^2}{\lambda^2}}$$

(exercise)

Proof of Theorem. Formally:

$$\begin{split} \frac{c_{\alpha}}{|x|^{\alpha}} &= \frac{1}{|x|^{\alpha}} \pi^{-\frac{\alpha}{2}} \int_{0}^{\infty} e^{-\lambda} \lambda^{\frac{\alpha}{2}-1} d\lambda = \int_{0}^{\infty} e^{-\pi\lambda |x|^{2}} \lambda^{\frac{\alpha}{2}-1} d\lambda \\ \Rightarrow \frac{\hat{c}_{\alpha}}{|x|^{\alpha}}(k) &= \int_{0}^{\infty} \widehat{e^{-\pi\lambda |x|^{2}}}(k) \lambda^{\frac{\alpha}{2}-1} d\lambda = \int_{0}^{\infty} \lambda^{-\frac{d}{2}} e^{-\pi\frac{|k|^{2}}{\lambda}} \lambda^{\frac{\alpha}{2}-1} d\lambda \\ (\lambda \to \frac{1}{\lambda}) &= \int_{0}^{\infty} \lambda^{\frac{d}{2}e^{-\pi|k|^{2}\lambda}} \lambda^{-\frac{\alpha}{2}+1} \lambda^{-2} d\lambda \\ &= \frac{c_{d-\alpha}}{|k|^{d-\alpha}} \end{split}$$

Let $f \in C_c(\mathbb{R}^d)$. Then $\left(\frac{1}{|x|^{\alpha}} \star f\right)(x) = \int_{\mathbb{R}^d} \frac{1}{|x-y|^{\alpha}} f(y) \, dy$ is well defined as $\frac{1}{|x-y|} \in L^1_{loc}(\mathbb{R}^d, dy)$. It is bounded

$$\frac{1}{|x|^{\alpha}} \star f = \frac{1}{|x|^{\alpha}} \underbrace{\mathbb{1}(|x| \leqslant 1)}_{\in L^{\infty}(\mathbb{R}^{d})} \star \underbrace{f}_{L^{\infty}} + \underbrace{\frac{1}{|x|}\mathbb{1}(|x| > 1)}_{\in L^{\infty}} \star \underbrace{f}_{\in L^{1}} \in L^{\infty}(\mathbb{R}^{d})$$

When $|x| \to \infty$:

$$\left(\frac{1}{|x|^{\alpha}} \star f\right)(x) = \int_{\mathbb{R}^d} \frac{f(y)}{|x-y|^{\alpha}} \, dy = \int_{|y| \leqslant R} \frac{f(y)}{|x-y|^{\alpha}} \, dy \sim \frac{\int_{\mathbb{R}^d} f(y) \, dy}{|x|^{\alpha}}$$

Note that $\frac{c_{d-\alpha}}{|k|^{d-\alpha}} \underbrace{\hat{f}(k)}_{\text{bounded}} \in L^1(\mathbb{R}^d)$.

$$(...)\mathbb{1}(|k| \leq 1) + (...)\mathbb{1}(|k| > 1) \frac{1}{|k|^{d-\alpha}} |\hat{f}(k)| \, \mathbb{1}(|k| \leq 1) \leq ||f||_{L^{1}} \frac{\mathbb{1}(|k| \leq 1)}{|k|^{d-\alpha}} \in L^{1}(\mathbb{R}^{d}, dk)$$
$$\frac{1}{|k|^{d-\alpha}} |\hat{f}(k)|\mathbb{1}(k > 1) \leq |\hat{f}(k)| \in L^{2}(\mathbb{R}^{d}, dK) \text{ as } f \in L^{2}(\mathbb{R}^{d})$$

Lemma 3.23 If $f \in C_c^{\infty}(\mathbb{R}^d)$, then $\hat{f} \in L^1(\mathbb{R}^d)$

Proof. (Exercise) Hint: $|\widehat{D^{\alpha}f}| = |2\pi k|^{|\alpha|} |\widehat{f}(k)| \rightsquigarrow |\widehat{f}(k)| \leqslant \frac{1}{|k|^{|k|}}$ as $|k| \to \infty$. Compute:

$$\left(\frac{c_{d-\alpha}}{|k|^{d-\alpha}}\hat{f}(k)\right)^{\vee}(x) = \int_{\mathbb{R}^d} \frac{c_{d-\alpha}}{|k|^{d-\alpha}}\hat{f}(k)e^{2\pi ikx} dk$$

$$= \int_{\mathbb{R}^d} \left(\int_0^{\infty} e^{-\pi|k|^2\lambda} \lambda^{\frac{d-\alpha}{2}-1} d\lambda\right) \hat{f}(k)e^{2\pi ikx} dk$$

$$= \int_0^{\infty} \left(\int_{\mathbb{R}^d} e^{-\pi|k|^2\lambda} \hat{f}(k)e^{2\pi ikx} dk\right) \lambda^{\frac{d-\alpha}{2}-1} d\lambda$$

$$= \int_0^{\infty} \left(e^{-\pi k^2\lambda} \hat{f}(x)\right)^{\vee} \lambda^{\frac{d-\alpha}{2}-1} d\lambda$$

$$= \int_0^{\infty} \left(\lambda^{-\frac{d}{2}} e^{-\pi \frac{x^2}{\lambda}} (k) \hat{f}(k)\right)^{\vee} \lambda^{\frac{d-\alpha}{2}-1} d\lambda$$

$$= \int_0^{\infty} \left(\lambda^{-\frac{d}{2}} e^{-\pi \frac{x^2}{\lambda}} \lambda^{\frac{d-\alpha}{2}-1} d\lambda\right)$$

$$= \left(\int_0^{\infty} \lambda^{-\frac{d}{2}} e^{-\pi \frac{x^2}{\lambda}} \lambda^{\frac{d-\alpha}{2}-1} d\lambda\right) \star f$$

Assume $d > \alpha > \frac{d}{2}$. Then $\frac{c_{\alpha}}{|x|^{\alpha}} \star f \in L^{\infty}$ and behaves $\frac{c_{\alpha}(\int f)}{|x|^{\alpha}}$ as $|x| \to \infty$. This implies:

$$\int_{\mathbb{R}^d} \left| \ \frac{c_\alpha}{|x|^\alpha} \star f \right|^2 \leqslant c + \int_{|x| \geqslant R} \frac{c}{|x|^{2d}} \, dx < \infty$$

Thus the Fourier Transform $\frac{\widehat{c_{\alpha}}}{\|x\|^{\alpha}} \star f$ exists. Combining with

$$\frac{c_{\alpha}}{|x|^{\alpha}} \star f = \left(\frac{c_{d-\alpha}}{|f|^{d-\alpha}} \hat{f}(k)\right)^{\vee}$$

$$\Rightarrow \widehat{\frac{c_{\alpha}}{|x|^{\alpha}}} \star f = \frac{c_{d-\alpha}}{|k|^{d-\alpha}} \hat{f}(k)$$

Remark 3.24 If $d \geqslant 3$

$$\begin{split} \hat{G}(k) &= \frac{1}{|2\pi k|^2} \\ \Rightarrow G(x) &= \left(\frac{1}{|2\pi k|^2}\right)^{\vee} = \frac{1}{d(d-2(k)|x|^{d-2})} = \Phi(x) \end{split}$$

3.3 Theory of Distribution

Let $\Omega \subseteq \mathbb{R}^d$ be open.

- $D(\Omega) = C_c^{\infty}(\Omega)$ the space of test functions.
- $\phi_n \to \phi$ in $D(\Omega)$ if $\exists K \subseteq \Omega$, $\operatorname{supp}(\phi_n)$, $\operatorname{supp}(\phi) \subseteq K$ and $||D^{\alpha}(\phi_n \phi)||_{L^{\infty}} \to 0$ for all $\alpha = (\alpha_1, \dots, \alpha_d), d_i \in \{0, 1, 2, \dots\}$.

$$D'(\Omega) = \{T : D(\Omega) \to \mathbb{R} \text{ or } \mathbb{C} \text{ linear and continuous} \}$$

the space of distributions.

Motivation: $L^2(\Omega)' = L^2(\Omega), (L^p(\Omega))' = (L^q(\Omega)), \frac{1}{p} + \frac{1}{q} = 1.$

Example 3.25 ("normal functions" are distributions) If $f \in L^1_{loc}(\Omega)$, then $T = T_f$ defined by:

$$T(\phi) = \int_{\Omega} f(x)\phi(x) dx$$

is a distribution for all $\phi \in D(\Omega)$, i.e. $T \in D'(\Omega)$. Indeed, it is clear that $T(\phi)$ is well-defined for all $\phi \in D(\Omega)$ and $\phi \mapsto T(\phi)$ is linear. Let us check that $\phi \mapsto T(\phi)$ is continuous. Take $\phi_n \to \phi$ in $D(\Omega)$ and prove that $T(\phi_n) \to T(\phi)$. Since $\phi_n \to \phi$ in $D(\Omega)$, there is a compact K s.t. $\text{supp}(\phi_n)$, $\text{supp}(\phi) \subseteq K \subseteq \Omega$.

Question: Why is $f \mapsto T_f$ injective?

Lemma 3.26 (Fundamental lemma of calculus of variants) Let $\Omega \subseteq \mathbb{R}^d$ be open. If $f, g \in L^1_{loc}(\Omega)$ and $\int_{\Omega} f \phi \, dy = \int_{\Omega} g \phi \, dy$ for all $\phi \in D(\Omega)$, then f = g in $L^1_{loc}(\Omega)$

Example 3.27 (Dirac delta function) Let $\Omega \subseteq \mathbb{R}^d$ open. Define $T: D(\Omega) \to \mathbb{R}$ or \mathbb{C} by $T(\phi) = \phi(x_0)$. Let $x_0 \in \Omega$. Then $T \in D'(\Omega)$ and we denote it by δ_{x_0} . It is clear that $\phi \mapsto T(\phi) = \phi(x_0)$ is well-defined and linear for all $\phi \in D(\Omega)$. Take $\phi_n \to \phi$ in $D(\Omega)$ and prove $T(\phi_n) \to T(\phi)$, i.e. $\phi_n(x_0) \to \phi(x_0)$ (obvious.)

Example 3.28 (Principle Value) The function $f(x) = \frac{1}{x}$ is not in $L^1_{loc}(\mathbb{R})$, but we can still define

$$\int_{\mathbb{R}} f(x)\phi(x) dx = \int_{\mathbb{R}} \frac{\phi(x)}{x} dx$$

for all $\phi \in D(\mathbb{R})$ s.t. $\phi(0) = 0$. In fact,

$$\phi(x) = |\phi(x) - \phi(0)| \le (\sup |\phi'|)(x),$$

so $\frac{|\phi(x)|}{|x|} \in L^{\infty}(\mathbb{R})$ and compactly supported. So $\frac{\phi(x)}{x} \in L^{1}(\mathbb{R})$. Define $T : D(\mathbb{R}) \to \mathbb{R}$ or \mathbb{C} by

$$T(\phi) = \lim_{\epsilon \to 0} \int_{|x| > \epsilon} \frac{\phi(x)}{x} dx \quad \forall \phi \in D(\mathbb{R}) \text{ s.t. } \phi(0) = 0$$

We denote $T = p.v.(\frac{1}{x})$. We check that $T \in D'(\mathbb{R})$: For all $\epsilon > 0$ we have

$$\left|\frac{\phi(x)}{x}\right| \leqslant \frac{\|\phi\|_{L^{\infty}}}{\epsilon}$$

for all $|x| \ge \epsilon$ and ϕ is compactly supported. So we get for all $\epsilon > 0$:

$$\mathbb{1}(|x| \ge \epsilon) \frac{\phi(x)}{x} \in L^1(\mathbb{R}) \leadsto \int_{|x| \ge \epsilon} \frac{\phi(x)}{x} \, dx < \infty$$

We can write:

$$\int_{|x| \ge \epsilon} \frac{\phi(x)}{x} \, dx = \int_{|x| \ge 1} \frac{\phi(x)}{x} \, dx + \int_{\epsilon \le |x| \le 1} \frac{\phi(x)}{x} \, dx$$

The second part can be written as:

$$\int_{\epsilon \leqslant |x| \leqslant 1} \frac{\phi(x)}{x} \, dx = \int_{\epsilon}^{1} \frac{\phi(x)}{x} \, dx + \int_{-1}^{-\epsilon} \frac{\phi(x)}{x} \, dx = \int_{\epsilon}^{1} \frac{\phi(x) - \phi(-x)}{x} \, dx$$

Since $\phi \in C_c^{\infty}(\mathbb{R})$ it holds that $|\phi(x) - \phi(-x)| \leq 2\|\phi'\|_{L^{\infty}}(x)$.

$$\Rightarrow \frac{\phi(x) - \phi(-x)}{x} \in L^{\infty}(\mathbb{R}) \Rightarrow \frac{\phi(x) - \phi(-x)}{x} \in L^{1}(0, 1)$$
$$\Rightarrow \int_{0}^{1} \frac{\phi(x) - \phi(-x)}{x} dx = \lim_{\epsilon \to 0} \int_{\epsilon}^{1} \frac{\phi(x) - \phi(-x)}{x} dx$$

Remark 3.29 The function $\frac{1}{|x|^d}$ is not in $L^1_{loc}(\mathbb{R}^d)$ but $\exists T \in D'(\mathbb{R}^d)$ s.t. $T(\phi) = \int_{\mathbb{R}^d} \frac{\phi(x)}{|x|^d} dx$ for all $\phi \in C^\infty_c(\mathbb{R}^d)$ s.t. $\phi(0) = 0$

Let in the following $\mathbb{K} \in \{\mathbb{R}, \mathbb{C}\}.$

Definition 3.30 (Derivatives of distributions) Let $\Omega \subseteq \mathbb{R}^d$ and $T \in D'(\Omega)$. Define for $\alpha \in \mathbb{N}^d$:

$$D^{\alpha}T: D(\Omega) \longrightarrow \mathbb{K}$$

 $\phi \longmapsto (-1)^{|\alpha|}T(D^{\alpha}\phi)$

Motivation: $f \in C_c^{\infty}(\Omega)$

$$\int_{\Omega} (D^{\alpha} f) \phi = (-1)^{|\alpha|} \int_{\Omega} f(D^{\alpha} \phi)$$

"If the classical derivative exists, then it is the same as the distributional derivative." We write

$$(D^{\alpha}T)(\phi) = T_{D^{\alpha}f}(\phi) = (-1)^{|\alpha|}T_f(D^{\alpha}\phi).$$

Remark 3.31 For all $T \in D'(\Omega)$ it holds $D^{\alpha}T \in D'(\Omega)$ for all $\alpha \in \mathbb{N}^d$. Clearly

$$\phi \longmapsto (D^{\alpha}T)(\phi) = (-1)^{|\alpha|}T(D^{\alpha}\phi)$$

is linear. Moreover, if $\phi_n \to \phi$ in $D(\Omega)$, then $D^{\alpha}\phi_n \to D^{\alpha}\phi$ in $D(\Omega)$, so

$$(D^{\alpha}T)(\phi_n) = (-1)^{|\alpha|}T(D^{\alpha}\phi_n) \xrightarrow{n \to \infty} (-1)^{|\alpha|}T(D^{\alpha}\phi) = (D^{\alpha}T)(\phi)$$

Example 3.32 Consider $f: x \mapsto |x|$, then $f \in C(\mathbb{R})$ but $f \notin C^1(\mathbb{R})$. However,

$$f'(x) = g(x) = \begin{cases} 1 & x \ge 0 \\ -1 & x < 0 \end{cases} \in L^1_{loc}$$

Lets check f'=g, i.e. $-f(\phi')=f'(\phi)=g(\phi)$ for all $\phi\in D(\mathbb{R})$. Thus we need to prove:

$$-\int_{\mathbb{R}} f(x)\phi'(x) dx = \int_{\mathbb{R}} g(x)\phi(x) dx \quad \forall \phi \in D(\mathbb{R})$$

namely:

$$\underbrace{-\int_{\mathbb{R}} |x| \phi'(x) \, dx}_{:=(\star)} = \int_{0}^{\infty} \phi(x) \, dx - \int_{-\infty}^{0} \phi(x) \, dx$$

Now we have

$$(\star) = -\int_0^\infty x \phi'(x) \, dx + \int_{-\infty}^0 x \phi'(x) \, dx.$$

By integration by parts:

$$\int_0^\infty x \phi'(x) \, dx = \underbrace{[x \phi(x)]_0^\infty}_{-0} - \int_0^\infty \phi(x) \, dx = -\int_0^\infty \phi(x) \, dx$$

and similary:

$$\int_{-\infty}^{0} x\phi'(x) dx = -\int_{-\infty}^{0} \phi(x) dx$$

Thus f' = g in $D'(\Omega)$. We claim that $g' = 2\delta_0$ in $D'(\mathbb{R})$. In fact, for all $\phi \in D(\mathbb{R})$, then:

$$g'(\phi) = -g(\phi') = -\int_{\mathbb{R}} g\phi' \, dx = -\int_{-\infty}^{0} (-1)\phi' \, dx - \int_{0}^{\infty} (1)\phi' \, dx$$
$$= -\int_{0}^{\infty} \phi' \, dx + \int_{-\infty}^{0} \phi' \, dx = \left[\phi(0) - \underbrace{\phi(\infty)}_{=0}\right] + \left[\phi(0) - \underbrace{\phi(-\infty)}_{=0}\right] = 2\phi(0) = 2\delta_{0}(\phi)$$

So $g' = 2\delta_0$ in $D'(\mathbb{R})$.

Exercise 3.33 Prove that $(D^{\alpha}\delta_x)(\phi) = (-1)^{|\alpha|}(D^{\alpha}\phi)(x)$ for all $\phi \in D(\mathbb{R})$ for all $x \in \mathbb{R}$.

Definition 3.34 (Convergence of distributions) Let $\Omega \subseteq \mathbb{R}^d$ be open, then

$$T_n \xrightarrow{n \to \infty} T$$

in $D'(\Omega)$ if $T_n(\phi) \xrightarrow{n \to \infty} T(\phi)$ for all $\phi \in D(\Omega)$.

Exercise 3.35 Let $f \in L^1(\mathbb{R}^d)$, $\int f = 1$ For $\epsilon > 0$, define $f_{\epsilon}(x) = \epsilon^{-d} f(\epsilon^{-1} x)$. Then: $f_{\epsilon} \to \delta_0$ in $D'(\Omega)$.

Exercise 3.36 Let $\Omega \subseteq \mathbb{R}^d$ be open and $T_n \to T$ in $D'(\Omega)$. Then: $D^{\alpha}T_n \to D^{\alpha}T$ in $D'(\Omega)$ for all $\alpha = (\alpha_1, \dots, \alpha_d)$

Definition 3.37 (Convolution of distributions) Let $T \in D'(\mathbb{R})$ and $f \in L_c^{\infty}(\mathbb{R}^d)$. Define

$$(T \star f)(y) = T(f_y)$$

We write $f_y(x) = f(x - y)$ and $\tilde{f}(x) = f(-x)$.

Theorem 3.38 Let $T \in D'(\mathbb{R})$. Then for all $f \in D(\mathbb{R})$:

1. $y \mapsto T(f_y)$ is $C^{\infty}(\mathbb{R}^d)$ and

$$D_{y}^{\alpha}(T(f_{y})) = (D^{\alpha}T)(f_{y}) = (-1)^{|\alpha|}T(D^{\alpha}f_{y})$$

2. For all $g \in L^1(\mathbb{R}^d)$ and compactly supported, then

$$\int_{\mathbb{R}^d} g(y)T(f_y) \, dy = T(\underbrace{f \star g}_{\in C_c^{\infty}(\mathbb{R})})$$

Proof. 1. We prove that $y \mapsto T(f_y)$ is continuous. Take $y_n \to y$ in \mathbb{R}^d , then:

$$T(f_{y_n}) \to T(f_y)$$

since $f_{y_n} \to f_y$ in $D(\mathbb{R}^d)$. We check this: Since $f \subseteq C_c^{\infty}(\mathbb{R}^d)$, it holds that $\operatorname{supp} f \subseteq B(0,R) \subseteq \mathbb{R}^d$. Since $y_n \to y$ in \mathbb{R}^d . We have $\sup_n |y_n| < \infty$. Thus f_{y_n}, f_y are supported in $\overline{B(0,R+\sup_n |y_n|)} = K$ compact. Moreover

$$|f_{y_n}(x) - f_y(x)| = |f(x - y_n) - f(x - y)| \le ||\nabla f||_{L^{\infty}} ||y_n - y|| \to 0$$

So we get $||f_{y_n} - f_y||_{L^{\infty}} \to 0$ Similary:

$$||D^{\alpha}f_{u_{\infty}}-D^{\alpha}f_{n}||_{L^{\infty}}\to 0$$

Exercise 3.39 (E 3.1 Lebesgue Differentiation Theorem) Let $f \in L^1_{loc}(\mathbb{R}^d)$. Prove that that for almost every $x \in \mathbb{R}^d$:

$$\oint_{B(x,r)} |f(x) - f(y)| \, dy \xrightarrow{r \to 0} 0$$

Proof. Clearly the same result holds with $\mathbb{R}^d \leadsto \Omega \subseteq \mathbb{R}^d$ open. Also it suffices to consider $f \in L^1(\mathbb{R}^d)$. From the last time discussion, by a density argument there exists $r_n \to 0$ s.t.

$$\oint_{B(x,r_n)} |f(y) - f(x)| \, dy = 0$$

for a.e. $x \in \mathbb{R}^d$. We prove that for all $\epsilon > 0$, te set $A_{\epsilon} = \{x \in \mathbb{R}^d \mid \limsup_{r \to 0} f_{B(x,r)} \mid f(y) - f(x) \mid dy > \epsilon\}$ has measure 0. This will imply that

$$\bigcup_{n=1}^{\infty} A_{\frac{1}{n}} = \left\{ x \in \mathbb{R}^d \mid \limsup_{r \to 0} \int_{B(x,r)} |f(y) - f(x)| \, dy > 0 \right\}$$

has measure 0, which is what wie want to show. First, we show that $|A_{\epsilon}| = 0$: Take $\{f_n\} \subseteq C_c^{\infty}, f_n \to f \text{ in } L^1(\mathbb{R}^d)$. By the triangle inequality:

$$|f(y) - f(x)| \le |f(y) - f_n(y)| + |f_n(y) - f_n(x)| + |f_n(x) - f(x)|$$

So we get

$$\oint_{B(x,r)} |f(y) - f(x)| dy$$

$$\leq \oint_{B(x,r)} |f(y) - f_n(y)| dy + \oint_{B(x,r)} |f_n(y) - f_n(x)| + |f_n(x) - f(x)|$$

$$\Rightarrow \lim_{r \to 0} \sup \dots \leq \limsup_{r \to 0} (\dots) + 0 + |f_n(x) - f(x)|$$

Thus, for all $x \in A_{\epsilon}$, then:

$$\limsup_{r \to 0} \int_{B(x,r)} |f_n(y) - f(y)| \, dy + |f_n(x) - f(x)| > 2\epsilon$$

Observation: If $a, b \ge 0$, $a + b > 2\epsilon$ then either $a > \epsilon$ or $b > \epsilon$. Therefore $A_{\epsilon} \subseteq (S_{n,\epsilon} \bigcup \tilde{S}_{n,\epsilon})$, where

$$S_{n,\epsilon} = \{x \mid |f_n(x) - f(x)| > \epsilon\}$$

$$\tilde{S}_{n,\epsilon} = \{x \mid \limsup_{r \to 0} \int_{B(x,r)} |f_n(y) - f(y)| \, dy > \epsilon\}$$

Consequently: $|A_{\epsilon}| \leq |S_{n,\epsilon}| + |\tilde{S}_{n,\epsilon}|$ for all $n \geq 1$. By the Markov / Chebyshev inequality:

$$|S_{n,\epsilon}| \leqslant \int_{S_{n,\epsilon}} \frac{|f_n(x) - f(x)|}{\epsilon} \, dx = \int_{\mathbb{R}^d} \frac{|f_n(x) - f(x)|}{\epsilon} \, dx = \frac{\|f_n - f\|_{L^1}}{\epsilon}$$

We want to prove a simpler bound for $\tilde{S}_{n\epsilon}$. For all $x \in \tilde{S}_{n\epsilon}$:

$$\limsup_{r \to 0} \int_{B(x,r)} |f_n(x) - f(y)| \, dy > \epsilon$$

So there is a $r_x \in (0,1)$ s.t.

$$\int_{B(x,r_x)=B_x} |f_n(y) - f(y)| \, dy > \epsilon$$

Thus $\tilde{S}_{n\epsilon} \subseteq \left(\bigcup_{x \in \tilde{S}_{n,\epsilon}} B_x\right)$.

Lemma 3.40 (Vitali Covering) If F is a collection of balls in \mathbb{R}^d with bounded radius, then there exists a sub-collection $G \subseteq F$ s.t.

• G has disjoint balls

• $\bigcup_{B \in F} B \subseteq \bigcup_{B \in G} 5B, 5B(x,r) = B(x,5r)$

Remark 3.41 The condition of the boundedness of the radius is necessary. Otherwise, consider $\{B(0,n)\}_{n=1}^{\infty}$

Here consider $F=\{B_x\}_{x\in \tilde{S}_{n\epsilon}}$. With the vitali covering leamm there is a $G\subseteq F$ s.t. G contains disjoint balls and:

$$\tilde{S}_{n,\epsilon} \subseteq \bigcup_{B \in F} B \subseteq \bigcup_{B \in G} 5B$$

So we get

$$|\tilde{S}_{n,\epsilon}| \ \leqslant | \bigcup_{B \in G} 5B| \leqslant \sum_{B \in G} |5B| = \sum_{B \in G} 5^d |B|$$

On the other hand, for all $B \in G \subseteq F$:

$$\oint_{B} |f_n(y) - f(y)| \ dy > \epsilon \Rightarrow \int_{B} |f_n - f| > \epsilon |B|$$

This implies:

$$\sup_{B \in G} \int_{B} |f_n - f| > \epsilon \sum_{B \in G} |B|$$

Since balls in G are disjoint:

$$\int_{\mathbb{R}^d} \geqslant \int_{\bigcup_{B \in G}} |f_n - f| \, dy > \epsilon \sum_{B \in G} |B| \geqslant \frac{\epsilon}{5^d} |\tilde{S}_{n,\epsilon}|$$

So

$$|\tilde{S}_{n\epsilon}| \leqslant \frac{5^d}{\epsilon} ||f_n - f||_{L^1}$$

In summary:

$$|A_{\epsilon}| \le |S_{n,\epsilon}| + |\tilde{S}_{n,\epsilon}| \le \frac{5^d + 1}{\epsilon} ||f_n - f||_{L^1} \to 0$$

as $n \to \infty$. So $|A_{\epsilon}| = 0$ for all $\epsilon > 0$

- **Remark 3.42** 1. The proof can be done by using the Besicovitch covering lemma: For all $E \subseteq \mathbb{R}^d$ s.t. E is bounded. Let F = collection of balls s.t. for all $x \in E$ there is a $B_x \in F$ s.t. x is the center of B_x . There is a sub-collection $G \subseteq F$ s.t.
 - $E \subseteq \bigcup_{B \in G} B$
 - Any point in E belongs to at most C_d balls in C_T (C_d depends only on \mathbb{R}^d), i.e.

$$\mathbb{1}_{E}(x) \leqslant \sum_{B \in G} \mathbb{1}_{B}(x) \leqslant C_{d} \mathbb{1}_{E}(x) \forall x$$

2. By a simpler argument we can prove the weak L^1 -estimate:

$$\{x \mid f^{\star}(x) > \epsilon\} \leqslant \frac{c_d}{\epsilon} ||f||_{L^1(\mathbb{R}^d)}$$

(Hardy-Littlewood maximal function)

Exercise 3.43 (E 3.2) Let $1 \leq p, q, r \leq 2$, $\frac{1}{p} + \frac{1}{q} = 1 + \frac{1}{r}$. Recall that if $f \in L^p(\mathbb{R}^d)$, $g \in L^q(\mathbb{R}^d)$, then $f \star g \in L^r(\mathbb{R}^d)$ by Young's Inequality, and its Fourier transform is well-defined by the Hausdorff-Young inequality. Prove that

$$\widehat{f \star g}(k) = \widehat{f}(k)\widehat{g}(k) \quad \forall k \in \mathbb{R}^d$$

Hint: In the lecture we already discussed the case $f, g \in C_c(\mathbb{R}^d)$.

Solution.

Step 1) $f, g \in C_c^{\infty}(\mathbb{R}^d)$ (Fubini)

Step 2) $f \in L^p, g \in L^q$, find $f_n, g_n \in C_c^{\infty}$ s.t. $f_n \to f$ in $L^p, g_n \to g$ in L^q . $\widehat{f_n \star g_n} = \widehat{f_n \hat{q}_n}$ pointwise a.e. we have

(Hausdorff-Young)
$$\begin{split} \|\widehat{f\star g} - \widehat{f_n\star g_n}\|_{L^{r'}} \\ &\leqslant \|\widehat{f\star g} - \widehat{f_n\star g_n}\|_{L^r} \\ &= \|(f-f_n)\star g_n + f_n\star (g_n-g)\|_{L^r} \\ &\leqslant \|(f-f_n)\star g_n\|_{L^r} + \|f_n\star (g_n-g)\|_{L^r} \\ &(\text{Young}) \leqslant \|f-f_n\|_{L^p}\|g_n\| + \|f_n\|_{L^p}\|g_n-g\|_{L^p} \xrightarrow{n\to\infty} 0 \end{split}$$

Moreover:

$$\|\hat{f}_{n}\hat{g}_{n} - \hat{f}\hat{g}\|_{L^{r'}} = \|(\hat{f}_{n}\hat{f})\hat{g}_{n} + \hat{f}(\hat{g}_{n} - \hat{g})\|_{L^{r'}}$$

$$(\text{H\"older}) \leq \|\hat{f}_{n} - \hat{f}\|_{L^{p'}} \|\hat{g}_{n}\|_{L^{q'}} + \|\hat{f}\|_{L^{q'}}$$

$$(\text{Hausdorff-Young (3.19)}) \leq \|f_{n} - f\|_{L^{p}} \|g_{n}\|_{L^{q}} + \|f\|_{L^{p}} \|g_{n} - g\|_{L^{p}} \xrightarrow{n \to \infty} 0$$
So $\hat{f}_{n}\hat{g}_{n} \to \hat{f}\hat{g}$ in $L^{r'}$ $\widehat{f \star g} = \hat{f}\hat{g}$ in $L^{r'}$ $\frac{1}{r'} = \frac{1}{r'} + \frac{1}{q'}$

Exercise 3.44 (E 3.3) $f \in C_c^{\infty}(\mathbb{R}^d)$. Prove $|\hat{f}(k)| \leq \frac{C_N}{(1+|k|)^N}$

Solution. Since $f \in C_c^{\infty}$ we have that $D^{\alpha} f \in C_c^{\infty}$. Recall

$$\widehat{D^{\alpha}f}(k) = (-2\pi i k)^{\alpha} \widehat{f}(k)$$

For example

$$\widehat{-\Delta f}(k) = |2\pi i k|^2 \widehat{f}(k)$$
(Induction) $\rightsquigarrow \widehat{(-\Delta)^N} f(k) = |2\pi k|^{2N} \widehat{f}(k)$

So we can conclude

$$\hat{f}(k) = \frac{\widehat{(-\Delta)^N} f(k)}{|2\pi k|^{2N}} \forall k \in \mathbb{R}^d$$

1.
$$f \in C_c^{\infty} \subseteq L^1(\mathbb{R}^d) \Rightarrow \hat{f} \in L^{\infty}$$

2.
$$(-\Delta)^N f \in C_c^{\infty} \subseteq L^1(\mathbb{R}^d) \Rightarrow \widehat{(-\Delta)^N} f \in L^{\infty}$$

Conclusion:
$$\hat{f}(k) \leqslant \begin{cases} C & \forall k \\ \frac{C_N}{|k|^{2N}} & \forall k \end{cases}$$
 So $\hat{f}(k) \leqslant \frac{C_N}{(1+|k|)^N}$

Exercise 3.45 (E 3.4)

Proof. Siehe Goodnotes

Exercise 3.46 (Bonus 3) Let $f \in L^1(\mathbb{R}^d)$ such that

$$|\hat{f}(k)| \leqslant \frac{C_N}{(1+|k|)^N}$$

for all $k \in \mathbb{R}^d$, for all $N \ge 1$. $(C_N \text{ is independent of } k)$. Prove that $f \in C^{\infty}(\mathbb{R}^d)$

$$(f \in C^{\infty})$$
 i.e. $\exists \tilde{f} \in C^{\infty}$ s.t. $f = \tilde{f}$ a.e.

Theorem 3.47 Take $T \in D'(\mathbb{R}), f \in C_c^{\infty}(\mathbb{R}^d) = D(\mathbb{R}^d), f_y(x) = f(x-y)$

- a) $y \mapsto T(f_y) \in C^{\infty}(\mathbb{R}^d)$ and $D_y^{\alpha}(T(f_y)) = (D^{\alpha}T)(f_y) = (-1)^{|\alpha|}T(D_x^{\alpha}f_y)$
- b) $\forall g \in L^1(\mathbb{R}^d)$ and compactly supported

$$\int_{\mathbb{R}^d} g(y)T(f_y) \, dy = T(\underbrace{f \star g}_{\in C_{c}^{\infty}})$$

Proof. a) $y \mapsto T(f_y)$ is continuous since $y_n \to y$ in \mathbb{R}^d , then $f_{y_n} \to f_y$ implies $T(f_{y_n}) \to T(f_y)$. Let's check that $y \mapsto T(f_y) \in C^1$:

$$\lim_{h \to 0} \frac{T(f_{y-he_i}) - T(f_y)}{h} = \lim_{h \to 0} T\left(\frac{f_{y-he_i} - f_y}{h}\right)$$

We have $\xrightarrow{f_{y-he_i}-f_y} \xrightarrow{h\to 0} (\partial_i f)_y$ in $D(\mathbb{R}^d)$

- $\exists K$ compact set such that $\operatorname{supp}(f_{y-e_i}-f_y)$, $\operatorname{supp} \partial_i f \subseteq K$ as |h| small.
- $\frac{f_{y-he_i}(x) f_y(x)}{h} (\partial_i f)_y(x)$ $= \frac{f(x-y+he_i) f(x-y)}{h} (\partial_i f)(x-y)$

$$\left| \int_0^1 \partial_i f(x - y + the_i) dt - \partial_i f(x - y) \right| \xrightarrow{h \to 0} 0 \text{ uniformly in } x$$

Similary:

$$\left| D_x^{\alpha} \left(\frac{f(x-y+he_i) - f(x-y)}{h} - (\partial_i f)(x-y) \right) \right|$$

$$= \left| \frac{D^{\alpha} f(x-y+he_i) - D^{\alpha} f(x-y)}{h} - \partial_i (D^{\alpha} f)(x-y) \right| \xrightarrow{h \to 0} 0$$

uniformly in x. Conclude:

$$\lim_{h \to 0} \frac{T(f_{y-he_i}) - T(f_y)}{h} \xrightarrow{h \to 0} T((\partial_i f)_y) \in C(\mathbb{R}^d)$$

So we geht that $y \mapsto T(f_y) \in C^1$ and $-\partial_{y_i} T(f_y) = T((\partial_i f)_y)$

By induction:

$$D_y^{\alpha}T(f_y) = (-1)^{|\alpha|}T((D^{\alpha}f)_y) = (D^{\alpha}T)(f_y) \quad \forall \alpha \in \mathbb{N}^d$$

b) Heuristic: T = T(x)

$$\int_{\mathbb{R}^d} g(y)T(f_y) \, dy = \int_{\mathbb{R}^d} g(y) \left(\int_{\mathbb{R}^d} T(x)f(x-y) \, dx \right) \, dy$$
$$= \int_{\mathbb{R}^d} T(x) \left(\int_{\mathbb{R}^d} g(y)f(x-y) \, dy \right) \, dx$$
$$= \int_{\mathbb{R}^d} T(x)(f \star g)(x) \, dx = T(f \star g)$$

Step 1: $g \in C_c^{\infty}(\mathbb{R}^d)$

(Rieman Sum)
$$\int_{\mathbb{R}^d} g(y)T(f_y) dy = \lim_{\Delta_N \to 0} \Delta_N \sum_{j=1}^N g(y_j)T(f_{y_j})$$
$$= \lim_{\Delta_N \to 0} T\left(\Delta_N \sum_{j=1}^N g(y_j)f_{y_j}\right)$$
$$= T(f \star g)$$

because

$$\lim_{\Delta_N \to 0} \Delta_N \sum_{j=1}^N g(y_j) f_{y_j}(x) \to (f \star g)(x) \text{ in } D(\mathbb{R}^d)$$

$$\lim_{\Delta_N \to 0} \Delta_N \sum_{j=1}^N g(y_j) f(x - y_j) \xrightarrow{\text{Riemann}} \int_{\mathbb{R}^d} g(y) f(x - y) \ dy = (f \star g)(x)$$

Proof of:

$$\lim_{\Delta_N \to 0} \Delta_N \sum_{j=1}^N g(y_j) f(x - y_j) \to (f \star g)(x) \text{ in } D(\mathbb{R}^d)$$

1) Since $f, g \in C_c^{\infty}$ we have $f \star g \in C_c^{\infty}$. And we have

$$x \mapsto \Delta_N \sum_{j=1}^N g(y_j) f(x - y_j) \in C^{\infty}$$

since $f \in C^{\infty}$ supported in $(\operatorname{supp} g + \operatorname{supp} f)$. So all functions are C_c^{∞} and supported in $(\operatorname{supp} g + \operatorname{supp} f)$.

2)

$$\left| \lim_{\Delta_N \to 0} \Delta_N \sum_{j=1}^N g(y_j) f(x - y_j) - \int_{\mathbb{R}^d} g(y) f(x - y) \, dy \right| \xrightarrow{\Delta_N \to 0} 0$$

uniformly in x. (Result from the Riemann-Sum)

3)

$$\left| D_x^{\alpha} (\Delta_N \sum_{j=1}^N g(y_j) f(x-y) - (f \star g)(x)) \right|$$

$$= \left| \Delta_N \sum_{j=1}^N g(y_j) D^{\alpha} f(x-y) - (D^{\alpha} f) \star g(x) \right| \xrightarrow{\Delta_N \to 0} 0$$

uniformly in x for all α .

Step 2: Take $g \in L^1(\mathbb{R}^d)$ and compactly supported. Then $\exists \{g_n\} \subseteq C_c^{\infty}(\mathbb{R}^d)$, supp $g_n \subseteq \text{supp } g + B(0,1)$ such that $g_n \to g$ in $L^1(\mathbb{R}^d)$. By Step 1:

$$\int_{\mathbb{R}^d} g_n(y) T(f_y) \, dy = T(g_n \star f)$$

Take $n \to \infty$:

$$\int_{\mathbb{R}^d} g_n(y) T(f_y) \, dy \to \int_{\mathbb{R}^d} g(y) T(f_y) \, dy$$

since $g_n \to g$ in L^1 compactly supported and $y \mapsto T(f_y) \in C^{\infty} \subseteq L^{\infty}(K)$. Moreover (exercise):

$$\underbrace{g_n \star f}_{\in C^{\infty}} \to g \star f \quad \text{in } D(\mathbb{R}^d)$$

So $T(g_n \star f) \xrightarrow{n \to \infty} T(g \star f)$. Finally we optain:

$$\int g(y)T(f_n)\,dy = T(g\star f)$$

Theorem 3.48 Let $\Omega \subseteq \mathbb{R}^d$ be open. Let $T \in D'(\Omega)$ and $f \in C_c^{\infty}(\Omega)$. Denote

$$\Omega_f = \{ y \in \mathbb{R}^d \mid \operatorname{supp} f_y = y + \operatorname{supp} f \subseteq \Omega \}$$

- a) $y \mapsto T(f_y) \in C^{\infty}(\Omega_f)$ and $D_y^{\alpha}(T(f_y)) = (D^{\alpha}T)(f_y) = (-1)^{|\alpha|}T((D^{\alpha}f)_y)$
- b) For all $g \in L^1(\Omega_q)$ compactly supported in Ω_f and it holds:

$$\int_{\Omega} g(y)T(f_y) \, dy = T(f \star g).$$

Theorem 3.49 Let $T \in D'(\Omega)$ s.t. $\nabla T = 0$ in $D'(\Omega)$. Then: T = const. in Ω .

Proof. $(\Omega = \mathbb{R}^d)$ for all $f \in C_c^{\infty}$, $y \mapsto T(f_y) \in C^{\infty}(\mathbb{R}^d)$ and $\partial_{y_i} T(f_y) = (\partial_j T)(f_y) = 0$ for all $i = 1, \ldots, d$. Then by the result of the theorem for C^{∞} functions, $y \mapsto T(f_y) = const$ independent of y. Consequently:

$$T(f_y) = T(f_0) = T(f) \quad \forall y \in \mathbb{R}^d \ \forall f \in C_c^{\infty}(\mathbb{R}^d)$$

For any $g \in C^{\infty}(\mathbb{R}^d)$:

$$\left(\int_{\mathbb{R}^d} g \, dy\right) T(f) = \int_{\mathbb{R}^d} g(y) T(f_y) \, dy = T(f \star g) = T(g \star f) = \left(\int_{\mathbb{R}^d} f \, dy\right) T(g)$$

So $\frac{T(f)}{\int_{\mathbb{R}^d} f}$ is independent of f (as soon as $\int f \neq 0$). So we get that $T(f) = const \int_{\mathbb{R}^d} f$, where const is independent of f.

Remark 3.50 If $u \in C^1(\mathbb{R}^d)$, then:

$$u(x+y) - u(x) = \int_0^1 \sum_{j=1}^d y_j (\partial_j u)(x+ty_j) dt = \int_0^1 y \nabla u(x+ty) dt$$

So we get that if $\nabla u = 0$, then u(x+y) - u(x) = 0 for all x, y, so u = const.

Theorem 3.51 (Taylor expansion for distributions) Let $T \in D'(\mathbb{R}^d)$ and $f \in C_c^{\infty}(\mathbb{R}^d)$. Then $y \mapsto T(f_u) \in C^{\infty}$ and

$$T(f_y) - T(f) = \int_0^1 \sum_{j=1}^d y_j(\partial_j T)(f_{ty}) dt.$$

In particular, if $g \in L^1_{loc}$ and $\nabla g \in L^1_{loc}$, then $\forall y \in \mathbb{R}^d$:

$$g(x+y) - g(x) = \int_0^1 g(x+ty)y \, dt$$

for a.e. $x \in \mathbb{R}^d$.

Proof. $y \mapsto T(f_y)$ is C^{∞} and $\frac{d}{dt}[T(f_{ty})] = (\nabla T)(f_{ty})y$ So we get

$$T(f_y) - T(f) = \int_0^1 \frac{d}{dt} (T(f_{ty})) dt$$
$$= \int_0^1 (\nabla T) (f_{ty}) y dt$$
$$= \int_0^1 \sum_{j=1}^d (\partial_j T) (f_{ty}) y_j dt$$

Corrolary 3.52 Let $g \in L^1_{loc}(\mathbb{R}^d)$ s.t. $\hat{\sigma}_j g \in L^1_{loc}(\mathbb{R}^d)$ for all $j = 1, 2, \dots, d$ (i.e. $g \in W^{1,1}_{loc}(\mathbb{R}^d)$). Then for all $y \in \mathbb{R}^d$:

$$g(x+y) - g(x) = \int_0^1 y \cdot \nabla g(x+ty) dt$$
$$= \int_0^1 \sum_{j=1}^d y_j \partial g(x+ty) dt$$

for a.e. x.

Proof. For all $f \in C_c^{\infty}$ we have

$$\int_{\mathbb{R}^d} f(x)[g(x+y) - g(x)] dx = \int_{\mathbb{R}^d} g(x)[f(x-y) - f(x)] dx$$

$$= g(f_y) - g(f)$$

$$= \int_0^1 \sum_{j=1}^d y_j (\partial_j g)(f_{ty}) dt$$

$$= \int_0^1 \sum_{j=1}^d y_j \int_{\mathbb{R}^d} \sum_{j=1}^d y_j \left[\int_{\mathbb{R}^d} (\partial_j g)(x) f_{ty}(x) dx \right]$$

$$= \int_0^1 \sum_{j=1}^d y_i \left[\int_{\mathbb{R}^d} (\partial_j g)(x+ty) f(x) dx \right] dt$$

$$= \int_{\mathbb{R}^d} f(x) \left[\int_0^1 \sum_{j=1}^d y_j \partial_j g(x+ty) dt \right] dx$$

For all $\phi \in C_c^{\infty}$: = g(x+y) - g(x) a.e. $x \in \mathbb{R}^d$.

Remark 3.53 If $T \in D'(\Omega)$, $\Omega \subseteq \mathbb{R}^d$ open, if $y\nabla T = 0$, then T = const.

Theorem 3.54 (Equivalence of the classical and distributional derivatives) Let $\Omega \subseteq \mathbb{R}^d$. Then the following are equivalent:

1.
$$T \in D'(\Omega)$$
 s.t. $\partial_{x_i} T = g_i \in C(\Omega)$ for all $i = 1, \ldots, d$.

2.
$$T = f \in C^1(\Omega)$$
 and $g_i = \partial_{x_i} f$

Proof.

(2) \Rightarrow (1): If $T = f \in C^1(\Omega)$, then: $\partial_{x_i} f \in C(\Omega)$.

$$\partial_{x_i} T(\phi) = -T(\partial_{x_i} \phi) = -\int_{\Omega} f(\partial_{x_i} \phi) = \int_{\Omega} (\partial_{x_i} f) \phi$$

for all $\phi \in D(\Omega)$, so $\partial_{x_i} T = \partial_{x_i} f$.

(1) \Rightarrow (2): Why is $T = f \in C^1(\Omega)$? As $\partial_{x_i} f = g_i$:

$$f(x+y) - f(x) = \int_0^1 \nabla f(x+ty)y \, dt = \int_0^1 \sum_{i=1}^d g_i(x+ty)y_i \, dt$$

So we get

$$f(y) = f(0) + \int_0^1 \sum_{i=1}^d g_i(ty)g_i dt.$$

We expect that $f \in C^1$ and $\partial_{x_i} f = g_i$. But this is not trivial to prove.

$$\frac{f(y+he_i)-f(y)}{h} = \int_0^1 \sum_{i=1}^d \left[g_i(ty+the_i)(y_i+h\delta_{ij})\right] dt$$

$$= \int_0^1 g_i(ty+the_i) dt + \int_0^1 \sum_{j\neq i} \frac{\left[g_i(ty+the_i)-g_i(ty)\right]}{h} y_i dt$$

$$\xrightarrow{h\to 0} \int_0^1 g_i(ty) dt + \text{is difficult ...}$$

Lets take $\phi \in C_c^{\infty}$, then:

$$T(\phi_y) - T(\phi) = \int_0^1 \underbrace{\nabla T}_{(g_i)_{i=1}^d} (\phi_{ty}) y \, dt$$

$$= \int_0^1 \sum_{i=1}^d \left(\int_{\Omega} g_i(x) \underbrace{\phi_{ty}}_{=\phi(x-ty)} dx \right) \, dt$$

$$= \int_{\mathbb{R}^d} \left(\sum_{i=1}^d \int_0^1 g_i(x) \phi(x-ty) y_i \, dt \right) \, dx$$

$$= \int_{\mathbb{R}^d} \left(\sum_{i=1}^d \int_0^1 g_i(x+ty) \phi(x) y_i \, dt \right) \, dx$$

$$= \int_{\mathbb{R}^d} \left(\sum_i \int_0^1 g_i(x+ty) y_i \, dt \right) \phi(x) \, dx$$

Integrating against $\psi(y)$ with $\psi \in C_c^{\infty}$:

$$\int_{\mathbb{R}^d} T(\phi_y)\psi(y) \, dy - T(\phi) \int_{\mathbb{R}^d} \psi(y) \, dy$$

$$= \int_{\mathbb{R}^d} \left(\int_{\mathbb{R}^d} \sum_i \int_0^1 g_i(x+ty) y_i \psi(y) \, dt \, dy \right) \psi(x) \, dx$$

$$\Rightarrow T(\phi \star \psi) - T(\phi) \int \psi = \dots$$

$$\Rightarrow \int_{\mathbb{R}^d} T(\psi_y) \phi(y) \, dy - T(\phi) \int \psi = \dots$$

Take $\psi \in C_c^{\infty}(\mathbb{R}^d)$ such that $\int \psi = 1$. Then:

$$T(\phi) = \int_{\mathbb{R}^d} \underbrace{T(\psi_x) - \left(\int_{\mathbb{R}^d} \sum_{i=1}^d \int_0^1 g_i(x+ty) y_i \psi(y) \, dt \, dy\right)}_{f(x)} \phi(x) \, dx$$

for all $\phi \in C_c^{\infty}$, so $T = f \in C(\Omega)$. Thus $T = f \in C(\Omega)$ and $\partial_{x_i} T = g_i \in C(\Omega)$. Then we need to prove that $f \in C^1(\Omega)$ and $\partial_{x_i} f = g_i$ (classical derivative). Since $f \in W_{loc}^{1,1}$:

$$f(x+y) - f(x) = \int_0^1 \sum_{i=1}^d g_i(x+ty)y_i dt \quad \forall x, y$$

In particular:

$$\frac{f(x+he_i) - f(x)}{h} = \int_0^1 \frac{1}{h} \sum_{i=1}^d g_i(x+the_i) h \delta_{ij} dt$$
$$= \int_0^1 g_i(x+the_i) dt \xrightarrow{h \to 0} g_i(x)$$

So we get $\partial_{x_i} f(x) = g_i(x) \in C(\Omega)$ in the classical sense. So $f \in C^1(\Omega)$.

Definition 3.55 (Sobolev Spaces) Let $\Omega \subseteq \mathbb{R}^d$ be open. We define for $1 \leq p \leq \infty$:

$$W^{1,p}(\Omega) = \{ f \in L^p(\Omega) \mid \partial_{x_i} f \in L^p(\Omega) \ \forall i = 1, \dots, d \}$$

$$W^{k,p}(\Omega) = \{ f \in L^p(\Omega) \mid D^{\alpha} f \in L^p(\Omega) \ \forall |\alpha| \le k \}$$

$$W^{k,p}_{loc}(\Omega) = \{ f \in L^p_{loc}(\Omega) \mid D^{\alpha} f \in L^p_{loc}(\Omega) \ \forall |\alpha| \le k \}$$

Theorem 3.56 (Approximation of $W^{1,p}_{loc}(\Omega)$ by $C^{\infty}(\Omega)$) Let $\Omega \subseteq \mathbb{R}^d$ be open, let $f \in W^{1,p}_{loc}(\Omega)$. Then there exists $\{f_n\} \subseteq C^{\infty}(\Omega)$ such that $f_n \to f$ in $W^{1,p}_{loc}(\Omega)$, i.e. for all $K \subseteq \Omega$ compact: $\|f_n - f\|_{L^p(K)} + \sum_{i=1}^d \|\partial_{x_i}(f_n - f)\|_{L^p(K)} \to 0$.

Proof. Case
$$\Omega = \mathbb{R}^d$$
: Take $g \in C_c^{\infty}$, $\int g = 1$, $g_{\epsilon}(x) = \epsilon^{-d}g(\epsilon^{-1}x)$. Then $g_{\epsilon} \star f \in C_c^{\infty}$. Since $f \in L_{loc}^p(\Omega)$ we have $g_{\epsilon} \star f \to f$ in L_{loc}^p as $\epsilon \to 0$. Moreover $\partial_{x_i}(g_{\epsilon} \star f) = (g_{\epsilon} \star \partial_{x_i} f) \xrightarrow{\epsilon \to 0} \partial_{x_i} f$ in L_{loc}^p . Then we can take $f_n = g_{\frac{1}{n} \star f}$.

Remark 3.57 In general, if we want to compute the distributional derivative $D^{\alpha}f$, then we can find $f_n \to f$ in $D'(\Omega)$ and compute $D^{\alpha}f_n$. Then $D^{\alpha}f_n \to D^{\alpha}f$

in $D^{\alpha}(\Omega)$. As an example we can compute $\nabla |f|$ with $f \in W^{1,p}_{loc}(\Omega)$.

$$(\nabla |f|)(x) = \begin{cases} \nabla f(x) & f(x) > 0 \\ -\nabla f(x) & f(x) < 0 \\ 0 & f(x) = 0 \end{cases}$$

Theorem 3.58 (Chain Rule) Let $G \in C^1(\mathbb{R}^d)$ with $|\nabla G|$ is bounded. Let $f = (f_i)_{i=1}^d \subseteq W^{1,p}_{loc}(\Omega)$. Then $x \mapsto G(f(x)) \in W^{1,p}_{loc}(\Omega)$ and

$$\partial_{x_i} G(f) = \sum_{k=1}^d (\partial_k G)(f) \cdot \partial_{x_i} f_k \quad \text{in } D'(\Omega)$$

Moreover, if $G(0) \in L^p(\Omega)$ (i.e. either $|\Omega| < \infty$ or G(0) = 0), then if $f = (f_i)_{i=1}^d \subseteq W^{1,p}(\Omega)$, then $G(f) \in W^{1,p}(\Omega)$.

Proof. Since $G \in C^1$ we have that G is bounded in any compact set. Moreover $\|\nabla G\|_{L^{\infty}} < \infty$ implies:

$$|G(f) - G(0)| \leq ||\nabla G||_{L^{\infty}} |f| \in L_{loc}^p$$

So $G(f) \in L^p_{loc}$. Let us compute $\partial_{x_i} G(f)$. Let $\{f^{(n)}\}_{n=1}^{\infty} \subseteq C^{\infty}$ such that $f^{(n)} \to f$ in $W^{1,p}_{loc}$, then:

$$|G(f^{(n)}) - G(f)| \le ||\nabla G||_{L^{\infty}} |f^{(n)} - f| \to 0 \text{ in } L^{p}_{loc}$$

So $G(f^{(n)}) \to G(f)$ in L^p_{loc} , thus $\partial_{x_i} G(f^{(n)}) \to \partial_{x_i} G(f)$ in $D'(\Omega)$. On the other hand, by the standard Chain-Rule for C^1 -functions:

$$\partial_{x_i} G(f^{(k)}) = \sum_{k=1}^d \underbrace{\partial_k G(f^{(k)})}_{\text{(b,d,} \to \partial_k G(f))} \underbrace{\partial_i f_k^{(n)}}_{\text{(b,d,} \to \partial_k G(f))} \to \sum_{k=1}^d \partial_k G(f) \partial_i f_k \text{ in } L^p_{loc}(\Omega)$$

Thus

$$\partial_{x_i} G(f) = \sum_{k=1}^d \underbrace{\partial_k G(f)}_{\in L^{\infty}} \underbrace{\partial_i f_k}_{\in L^p_{loc}} \in L^p_{loc} \text{ in } D'(\Omega)$$

So $G(f) \in W^{1,p}_{loc}(\Omega)$. Aussume that $G(0) \in L^p(\Omega)$ (i.e. $|\Omega| < \infty$ or G(0) = 0). If $f \in W^{1,p}(\Omega)$, then $G(f) \in W^{1,p}(\Omega)$ since

$$|G(f) - G(0)| \leq \|\nabla G\|_{L^{\infty}} |f| \in L^p \Rightarrow G(f) \in L^p$$

and

$$\partial_{x_i} G(f) = \sum_k \underbrace{\partial_k G}_{\in L^{\infty}} \underbrace{\partial_i f_k}_{\in L^p} \in L^p \Rightarrow G(f) \in W^{1,p}(\Omega)$$

Theorem 3.59 (Derivative of absolute value) Let $\Omega \subseteq \mathbb{R}^d$ be open. Let $f \in W^{1,p}(\Omega)$. Then $|f| \in W^{1,p}(\Omega)$ and if f is real-valued:

$$(\nabla |f|)(x) = \begin{cases} \nabla f(x) & f(x) > 0 \\ -\nabla f(x) & f(x) < 0 \\ 0 & f(x) = 0 \end{cases}$$

Proof. Exercise. Hint: Use the Chain-Rule for $G_{\epsilon}(x) = \sqrt{\epsilon^2 + x^2} - \epsilon \rightarrow |x|$ as $\epsilon \rightarrow 0$

3.4 Distribution vs. measures

Let μ be a Borel measure in \mathbb{R}^d s.t. $\mu(K) < \infty$ for all compact $K \subseteq \mathbb{R}^d$. Then define

$$T:\ D(\mathbb{R}^d) \longrightarrow \mathbb{C}$$

$$\phi \longmapsto \int_{\mathbb{R}^d} \phi(x)\,d\mu(x) \quad \forall \phi \in C_c^\infty$$

 \rightsquigarrow T is a distribution since if $\phi_n \to \phi$ in $D(\Omega)$, then

$$|T(\phi_n) - T(\phi)| \le \int_{\mathbb{R}^d} |\phi_n - \phi| \, d\mu(x) \le \|\phi_n - \phi\|_{L^{\infty}} \left(\int_K d\mu \right) \xrightarrow{n \to \infty} 0$$

Example 3.60 ∂_0 in $D'(\mathbb{R}^d)$ is a Borel probability measure.

Theorem 3.61 (Positive distributions are measures) Let $\Omega \subseteq \mathbb{R}^d$ be open, let $T \in D'(\Omega)$. Assume $T \geq 0$, i.e. $T(\phi) \geq 0$ for all $\phi \in D(\Omega)$ satisfying $\phi(x) \geq 0$ for all x. Then there is a Borel positive measure μ on Ω such that $\mu(K) < \infty$ for all $K \subseteq \Omega$ compact and:

$$T(\phi) = \int_{\Omega} \phi(x) \, d\mu(x) \quad \forall \phi \in D^{(\Omega)}$$

Proof. See Lieb-Loss Analysis. Sketch: If $O \subseteq \mathbb{R}^d$ is open, then

$$\mu(O) = \sup\{T(\phi) \mid \phi \in D(\Omega), 0 \le \phi \le 1, \operatorname{supp} \phi \subseteq O\}$$

For all $A \subseteq \Omega$ (not necessarily open),

$$\mu(A) = \inf \{ \mu(O) \mid O \text{ open}, A \subseteq O \}$$

The mapping $\mu: 2^{\Omega} \to [0, \infty]$ is an outer measure, i.e.

- 1. $\mu(\emptyset) = 0$
- 2. $\mu(A) \leq \mu(B)$ if $A \subseteq B$
- 3. $\mu\left(\bigcup_{i=1}^{\infty} A_i\right) \leqslant \sum_{i=1}^{\infty} \mu(A_i)$

From the outer measure we can find a σ -algebra Σ and μ is a measure on Ω s.t. E is measurable iff

$$\mu(E) = \mu(E \cap A) + \mu(E \cap A^{\complement}).$$

So all open sets are measurable, thus outer regularity (by def $\mu(A) = \inf\{\mu(O) \mid O \text{ open } \supseteq A\}$, so inner regularity $\mu(A) = \sup\{\mu(K) \mid K \text{ compact } \subseteq A\}$.

Exercise 3.62 (E 4.1) Prove that if $T_n \to T$ in $D'(\mathbb{R}^d)$, then $D^{\alpha}T_n \to D^{\alpha}T$ in $D^{\alpha}(\mathbb{R}^d)$ for all $\alpha \in \mathbb{N}^d$.

Exercise 3.63 (E 4.2)

Exercise 3.64 (E 4.3) $f \in L^1(\mathbb{R}^d)$, $\int f = 1$ $f_{\epsilon}(x) = \epsilon^{-d} f(\epsilon^{-1} x)$. Then $f_{\epsilon} \to \delta_0$ in $D'(\mathbb{R}^d)$.

Exercise 3.65 (E 4.4) Let $\{f_n\} \subseteq L^1$, supp $f \subseteq B(0,1), f_n \to f$ in L^1 . Prove for all $g \in C_c^{\infty}$ that $f_n \star g \to f \star g$ in $D(\mathbb{R}^d)$.

Solution. Since $f_n \in L^1$, supp $f \subseteq B(0,1)$ and $g \in C_c^{\infty}$ we have $f_n \star g \in C_c^{\infty}$ and

$$\operatorname{supp}(f_n \star g) \subseteq (\operatorname{supp} g) + \overline{B(0,1)} = K.$$

Since $f_n \to f$ in L^1 there is a subsequence $f_{n_k} \to f$ almost everywhere, so f supp in $\overline{B(0,1)}$, so $f \star g \in C_c^{\infty}$, supp $(f \star g) \subseteq K$. We have:

$$|f_n \star g(x) - f \star g(x)| = \left| \int_{\mathbb{R}^d} (f_n(y) - f(y))g(x - y) \, dy \right|$$

$$\leq \int_{\mathbb{R}^d} |f_n(y) - f(y)||g(x - y)| \, dy$$

$$\leq ||g||_{L^{\infty}} ||f_n - f||_{L^1} \xrightarrow{n \to \infty} 0$$

thus $||f_n \star g - f \star g||_{L^{\infty}} \to 0$. Similary:

$$||D^{\alpha}(f_n \star g) - D^{\alpha}(f \star g)||_{L^{\infty}} = ||f_n \star \underbrace{(D^{\alpha}g)}_{\in C^{\infty}} - f \star (D^{\alpha}g)||_{L^{\infty}} \xrightarrow{n \to \infty} 0$$

for all $\alpha \in \mathbb{N}^d$, so $f_n \star g \to f \star g$ in $D(\mathbb{R}^d)$.

Exercise 3.66 (E 4.5) Compute distributional derivatives f', f'' of f(x) = x|x-1|.

Solution. We prove
$$f'(x) = g(x) := \begin{cases} 2x - 1 & x > 1 \\ 1 - 2x & x < 1 \end{cases}$$
. Take $\phi \in C_c^{\infty}(\mathbb{R}^d)$.

$$-f'(\phi) = -\int_{\mathbb{R}^d} f\phi' \, dy$$

$$= -\int_{-\infty}^1 f\phi' \, dy - \int_1^{\infty} f\phi' \, dy$$

$$= [f\phi]_{-\infty}^1 - \int_{-\infty}^1 f'\phi \, dy + [f\phi]_1^{\infty} - \int_1^{\infty} f'\phi \, dy$$

$$= [f\phi]_{-\infty}^1 - \int_{-\infty}^1 g\phi \, dy + [f\phi]_1^{\infty} - \int_1^{\infty} g\phi \, dy$$

$$= f(1-)\phi(1) - f(1+)\phi(1) - \int_{\mathbb{R}^d} g\phi \, dy$$

$$= 0 - \int_{\mathbb{R}^d} g\phi \, dy$$

Now we compute f'' = g'. Take $\phi \in C_c^{\infty}(\mathbb{R}^d)$:

$$-(g')(\phi) = \int_{\mathbb{R}^d} g\phi' \, dy$$

$$= \int_{-\infty}^1 g\phi' \, dy + \int_1^{\infty} g\phi' \, dy$$

$$= [g(1-) - g(1+)]\phi(1) - \int_{-\infty}^1 g'\phi \, dy - \int_1^{\infty} g'\phi \, dy$$

$$= [g(1-) - g(1+)]\phi(1) - \int_{-\infty}^1 (-2)\phi \, dy - \int_1^{\infty} 2\phi \, dy$$

$$= -2\phi(1) + \int_{-\infty}^{\infty} [2\mathbb{1}_{(-\infty,1)}(x) - 2\mathbb{1}_{(1,\infty)}(x)]\phi(x) \, dx$$

$$= -2\delta_1(\phi) + \int_{-\infty}^{\infty} [2\mathbb{1}_{(-\infty,1)}(x) - 2\mathbb{1}_{(1,\infty)}(x)]\phi(x) \, dx$$

$$\Rightarrow g' = \underbrace{2\delta_1}_{\notin L^1_{loc}} - \underbrace{2\mathbb{1}_{(-\infty,1)} + 2\mathbb{1}_{(1,\infty)}}_{\int L^1_{loc}}$$

Chapter 4

Weak Solutions and Regularity

Definition 4.1 Consider the linear PDE:

$$\sum_{\alpha} c_{\alpha} D^{\alpha} u(x) = F(x), \quad c_{\alpha} \text{ constant}, F \text{ given}$$

A function u is called a weak solution (a distributional solution) if

$$\sum_{\alpha} c_{\alpha} D^{\alpha} u = F \quad \text{in } D'(\Omega).$$

Namely,

$$\sum_{\alpha} (-1)^{|\alpha|} c_{\alpha} \int_{\Omega} u D^{\alpha} \phi = \int_{\Omega} F \phi, \quad \forall \phi \in D(\Omega)$$

Regularity: Given some condition on the data F, what can we say about the smoothness of u? Can we say that the equation holds in the classical sense? We derived G (the solution of the Laplace Equation) before in two ways:

- 1. $\Delta G(x) = 0$ for all $x \neq 0$, assuming G(x) = G(|x|) and $d \geq 2$
- 2. $\hat{G}(k) = \frac{1}{|2\pi k|^2}$ for $d \ge 3$

Theorem 4.2 For all $d \ge 1$ we have $G \in L^1_{loc}(\mathbb{R}^d)$ and $-\Delta G = \delta_0$ in $D'(\mathbb{R}^d)$.

Proof. Take $\phi \in D(\mathbb{R}^d)$. Then:

$$(-\Delta G_y)(\phi) = G_y(-\Delta \phi) = \int_{\mathbb{R}^d} G_y(x)(-\Delta \phi)(x) dx$$
$$= \int_{\mathbb{R}^d} G(y - x)(-\Delta \phi)(x) dx$$
$$= [G \star (-\Delta \phi)](y) = (-\Delta)(G \star \phi)(y)$$

Recall for all $f \in C^2$, $-\Delta(G \star f) = f$ pointwise. So we can conclude $-\Delta G_y = \delta_y$ in $D'(\mathbb{R}^d)$.

Remark 4.3 In
$$d = 1$$
, $G(x) = -\frac{1}{2}|x|$, so $-G'(x) = \text{sgn}(x)/2$, so $-G''(x) = \delta_0$.

Remark 4.4 Formally:

$$-\Delta(G_y \star \phi) = (-\Delta G_y) \star \phi(x) = (\delta_0 \star \phi)(x) = \int \delta_0(y)\phi(x-y) \, dy = \delta_0(\phi(x-\bullet))$$

Theorem 4.5 (Poisson's equation with L^1_{loc} data) Let $f \in L^1_{loc}(\mathbb{R}^d)$ s.t. $\omega_d f \in L^1(\mathbb{R}^d)$ where

$$\omega_d(x) = \begin{cases} 1 + |x| & d = 1\\ \log(1 + |x|) & d = 2\\ \frac{1}{1 + |x|^{d-2}} & d \geqslant 3, \end{cases}$$

then $u(x)=(G\star f)(x)\in L^1_{loc}(\mathbb{R}^d).$ Moreover $-\Delta u=f$ in $D'(\mathbb{R}^d).$ In fact, $u\in W^{1,1}_{loc}(\mathbb{R}^d)$ and:

$$\partial_{x_i} u(x) = (\partial_{x_i} G) \star f(x) = \int_{\mathbb{R}^d} (\partial_{x_i} G)(x - y) f(y) dy$$

Remark 4.6 We can also replace \mathbb{R}^d by Ω and get $-\Delta u = f$ in $D'(\Omega)$.

Proof of Theorem 4.5. First we check that $u \in L^1_{loc}$. Take any Ball $B(0,R) \subseteq \mathbb{R}^d$, prove $\int_{\mathbb{R}} |u| dy < \infty$. We have

$$\begin{split} \int_{B} |u| \, dy &= \int_{B} \left| \int_{\mathbb{R}^{d}} G(x - y) f(y) \, dy \right| \, dx \\ &\leq \int_{B} \int_{\mathbb{R}^{d}} |G(x - y)| |f(y)| \, dy \, dx \\ &= \int_{\mathbb{R}^{d}} \left(\int_{B} |G(x - y)| \, dx \right) |f(y)| \, dy \end{split}$$

If $y \notin B = B(0, R)$, then by Newtons's theorem (Mean-value theorem):

$$\int_{B(0,R)} |G(x-y)| \, dx = |B(0,R)||G(y)| \le C|B|\omega_d(y)$$

If $y \in B$, then $|y| \le R$, so $|x - y| \le 2R$ if $x \in B$.

$$\int_{B(0,R)} \left| G(x-y) \right| dx \leqslant \int_{|x-y| \leqslant 2R} \left| G(x-y) \right| dx = \int_{|z| \leqslant 2R} \left| G(z) \right| dz \leqslant c_R$$

as $G \in L^1_{loc}$. Thus

$$\int_{B} |u| \, dy \leqslant c_{B} \int_{|y| \geqslant R} \omega_{d}(y) |f(y)| \, dy + c_{B} \int_{|y| \leqslant R} |f(y)| \, dy < \infty$$

Let us prove $-\Delta u = f$ in $D'(\mathbb{R}^d)$. Take $\phi \in D(\mathbb{R}^d)$. Then:

$$(-\Delta u)(\phi) = u(-\Delta \phi)$$

$$= \int_{\mathbb{R}^d} u(x)(-\Delta \phi)(x) dx$$

$$= \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} G(x-y)f(y)(-\Delta \phi)(x) dx dy$$

$$= \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} G(y-x)f(y)(-\Delta \phi)(x) dx dy$$

$$= \int_{\mathbb{R}^d} [G \star (-\Delta \phi)](y)f(y) dy$$

$$= \int_{\mathbb{R}^d} -\Delta (G \star \phi)(y)f(y) dy$$

$$= \int_{\mathbb{R}^d} \phi(y)f(y) dy$$

So $-\Delta u = f$ in $D'(\mathbb{R}^d)$. We check that $\partial_i G \star f \in L^1_{loc}(\mathbb{R}^d)$. Note that

$$|\hat{\sigma}_i G(x)| \le c \frac{1}{|x|^{d-1}} \in L^1_{loc}(\mathbb{R}^d)$$

and

$$\int_{B(0,R)} |\partial_i G(x-y)| dx \leq \begin{cases} C_r \omega_d(y) & |y| \geq R \\ C_r & |y| \leq R \end{cases}$$

So $\int_{B(0,R)} |(\partial_i G \star f)|(y) < \infty$ for all R > 0. For all $\phi \in D(\mathbb{R}^d)$:

$$-(\partial_{i}u)(\phi) = u(\partial_{i}\phi) = \int_{\mathbb{R}^{d}} u(x)\partial_{i}\phi(x) dx$$

$$= \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} G(x-y)f(y)\partial_{i}\phi(x) dx dy$$

$$= \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} G(y-x)f(y)\partial_{i}\phi(x) dx dy$$

$$= \int_{\mathbb{R}^{d}} (G \star \partial_{i}^{y}\phi)(y)f(y) dy$$

$$= \int_{\mathbb{R}^{d}} (\partial_{i}^{y}G \star \phi)(y)f(y) dy$$

$$= \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} \partial_{i}^{y}G(y-x)f(y)\phi(x) dx dy$$

$$= \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} -(\partial_{i}G)(x-y)f(y)\phi(x) dx dy$$

$$= -\int_{\mathbb{R}^{d}} (\partial_{i}G \star f)(x)\phi(x) dx$$

So $\partial_i u = \partial_i G \star f \in L^1_{loc}(\mathbb{R}^d)$. Thus $u \in L^1_{loc}$, $\partial_i u \in L^1_{loc}$ for all i. So $u \in W^{1,1}_{loc}(\mathbb{R}^d)$. \blacksquare Regularity: We consider the Laplace Equation $\Delta u = 0$ in \mathbb{R}^d .

Lemma 4.7 (Weyl) If $\Omega \subseteq \mathbb{R}^d$ open and $T \in D'(\Omega)$ s.t. $\Delta T = 0$ in $D'(\Omega)$, then: $T = f \in C^{\infty}(\Omega)$ and f is a harmonic function.

Proof. $(\Omega = \mathbb{R}^d)$. Take $\phi \in C_c^{\infty}$, then $y \mapsto T(\phi_y) = T(\phi(-y))$ is C^{∞} and $\Delta_y T(\phi_y) = T((\Delta\phi)_y) = (\Delta T)(\phi_y) = 0$. Take $g \in C_c^{\infty}$, g is radial. Then:

$$\int_{\mathbb{R}^d} T(\phi_y) g(y) \, dy \stackrel{\text{(exercise)}}{=} \int_{\mathbb{R}^d} T(\phi) g(y) \, dy = T(\phi) \left(\int_{\mathbb{R}^d} g \, dy \right)$$

Exercise 4.8 Let $f \in C^{\infty}(\mathbb{R}^d)$ be a harmonic function and $g \in C_c^{\infty}$, g is radial. Then:

$$\int_{\mathbb{R}^d} f(x)g(x) \, dx = f(0) \left(\int_{\mathbb{R}^d} g(x) \, dx \right)$$

On the other hand:

$$\int_{\mathbb{R}^d} T(\phi_y) g(y) \, dy = T(\phi \star g) = T(g \star \phi) = \int_{\mathbb{R}^d} T(g_y) \phi(y) \, dy$$

Take $\int_{\mathbb{R}^d} g \, dy = 1$, then:

$$T(\phi) = \int_{\mathbb{R}^d} T(g_y)\phi(y) \ dy$$

For all $\phi \in C_c^{\infty}$. Then $T = T(g_y) \in C^{\infty}$

Now lets regard the Poisson Equation $-\Delta u = f$ in $D'(\mathbb{R}^d)$.

Remark 4.9 Any solution has the form $u = G \star g + h$ where $\Delta h = 0$ in $D'(\mathbb{R}^d)$. By Weyls Lemma (4.7), $h \in C^{\infty}$, then we only need to consider the regularity of $G \star f$.

Remark 4.10 The regularity is a local question, namely if we write

$$f = f_1 + f_2 = f\phi + f(1 - \phi),$$

where $\phi = 1$ in a ball B and $\phi \in C_c^{\infty}$.

Then $G \star f = G \star f_1 + G \star f_2$. Here $f_2 = f(1 - \phi) = 0$ in B. With Weyls Lemma (4.7), $G \star f_2 \in C^{\infty}$.

Theorem 4.11 (Low Regularity of Poisson Equation) Lef $f \in L^p(\mathbb{R}^d)$ and compactly supported. Then

- a) If $p \ge 1$, then
 - $G \star f \in C^1(\mathbb{R}^d)$ if d = 1.
 - $G \star f \in L^q_{loc}(\mathbb{R}^d)$ for any $q < \infty$ if d = 2.
 - $G \star f \in L^q_{loc}(\mathbb{R}^d)$ for $q < \frac{d}{d-2}$ if $d \ge 3$.
- b) If $\frac{d}{2} , then <math>G \star f \in C^{0,\alpha}_{loc}(\mathbb{R}^d)$ for all $0 < \alpha < 2 \frac{d}{p}$, i.e.

$$|(G \star f)(x) - (G \star f)(y)| \le C_k |x - y|^{\alpha} \quad \forall x, y \in K$$

with K compact in \mathbb{R}^d .

c) If p>d, then $G\star f\in C^{1,\alpha}_{loc}(\mathbb{R}^d)$ for all $0<\alpha<1-\frac{d}{p}.$

where G is den fundamental solution of the laplace equation.

Example 4.12 Let r = |x|

$$u(x) = \omega(r) = \log(|\log(r)|)$$

if $0 < r < \frac{1}{2}$, so u is well-defined in $B = B(0, \frac{1}{2})$. We conclude:

$$-\Delta_{\mathbb{R}^3} u(x) = -\omega''(r) - \frac{2\omega'(r)}{r} = f(x) \in L^{\frac{3}{2}(B)}$$

But the Theorem (b) tells us that if $f \in L^{\frac{3}{2}}$ then u is continuous but $u \notin C(B)$.

Proof of theorem 4.11. a) (p = 1) Why is $G \star f \in L^q_{loc}$? Recall from the proof of Youngs inequality:

$$\begin{split} |(G\star f)(x)| &= \bigg|\int_{\mathbb{R}^d} G(x-y)f(y)\,dy\bigg| \\ \text{(H\"{o}lder)} &= \left(\int_{\mathbb{R}^d} |G(x-y)|^q |f(y)|\,yd\right)^{\frac{1}{q}} \left(\int_{\mathbb{R}^d} |f(y)|\,dy\right)^{\frac{1}{q'}} \end{split}$$

Where $\frac{1}{q} + \frac{1}{q'} = 1$. Then:

$$|(G \star f)(x)|^q \leqslant C \int_{\mathbb{R}^d} |G(x-y)|^q |f(y)| \, dy$$

For any Ball $B = B(0, R) \subseteq \mathbb{R}^d$:

$$\int_{B} |G \star f(x)|^{q} dx \leq C \int_{B} \left(\int_{\mathbb{R}^{d}} |G(x-y)|^{q} |f(y)| dy \right) dx$$
$$= C \int_{\mathbb{R}^{d}} \left(\int_{B} |G(x-y)|^{q} dx \right) |f(y)| dy$$

 $G(x) \sim \frac{1}{|x|^{d-2}} \rightsquigarrow |G|^q = \frac{1}{|x|^{(d-2)q}} \in L^1_{loc}(\mathbb{R}^d)$ if $(d-2)q < 2 \Leftrightarrow q < \frac{d}{d-2}$. Here, $y \in \operatorname{supp} f$, so $|y| \leqslant R_1$, then $|x-y| \leqslant R + R$ if $|x| \leqslant R$. With $y \in \operatorname{supp} f$, this implies:

$$\int_{B(0,R)} |G(x-y)|^q \, dx \le \int_{|z| \le R+R_1} |G(z)|^q \, dz < \infty$$

b)

$$(G \star f)(x) - (G \star f)(y) = \int_{\mathbb{R}^d} (G(x-z) - G(y-z))f(z) dz$$

So

$$|G\star f(x)-(G\star f)(y)|\leqslant C\int_{\mathbb{R}^d}\left|\frac{1}{|x-z|^{d-2}}-\frac{1}{|y-z|^{d-2}}\right||f(z)|\,dz$$

for all $x, y \in \mathbb{R}^d$:

$$\begin{split} \left| \frac{1}{|x|^{d-2}} - \frac{1}{|y|^{d-2}} \right| &= \left| \left(\frac{1}{|x|} - \frac{1}{|y|} \right) \left(\frac{1}{|x|^{d-3}} + \dots + \frac{1}{|y|^{d-3}} \right) \right| \\ &\leqslant C \frac{||x| - |y||}{|x||y|} \max \left(\frac{1}{|x|^{d-3}}, \frac{1}{|y|^{d-3}} \right) \\ &= C \frac{|x - y|}{|x||y|} \max \left(\frac{1}{|x|^{d-3}}, \frac{1}{|y|^{d-3}} \right) \\ &\leqslant C \max(|x|, |y|)^{1-\alpha} \frac{|x - y|^{\alpha}}{|x||y|} \max \left(\frac{1}{|x|^{d-3}}, \frac{1}{|y|^{d-3}} \right) \end{split}$$

as

$$||x| - |y|| \le \min(|x - y|, \max(|x|, |y|)) \le |x - y|^{\alpha} \max(|x|, |y|)^{1 - \alpha}$$

Thus, for all $x, y \in \mathbb{R}^d$:

$$\left| \frac{1}{|x|^{d-2}} - \frac{1}{|y|^{d-2}} \right| \leqslant C|x - y|^{\alpha} \frac{\max(|x|, |y|)^{1-\alpha}}{|x||y|} \max\left(\frac{1}{|x|^{d-3}}, \frac{1}{|y|^{d-3}}\right)$$
$$\leqslant C|x - y|^{\alpha} \max\left(\frac{1}{|x|^{d-2+\alpha}}, \frac{1}{|y|^{d-2+\alpha}}\right)$$

So we get

$$\left| \frac{1}{|x-y|^{d-2}} - \frac{1}{|y-z|^{d-2}} \right| \le C|x-y|^{\alpha} \max\left(\frac{1}{|x-z|^{d-2+\alpha}}, \frac{1}{|y-z|^{d-2+\alpha}} \right)$$

Therefore:

$$\begin{split} |G \star f(x) - G \star f(y)| \\ &\leqslant C \int_{\mathbb{R}^d} |x - y|^{\alpha} \max \left(\frac{1}{|x - z|^{d - 2 + \alpha}}, \frac{1}{|y - z|^{d - 2 + \alpha}} \right) |f(z)| \, dz \\ &\leqslant C |x - y|^{\alpha} \left(\sup_{\xi \in \mathbb{R}^d} \int_{\mathbb{R}^d} \frac{1}{|\xi - z|^{d - 2 + \alpha}} |f(z)| \, dz \right) \end{split}$$

Claim: If $f \in L^p(\mathbb{R}^d)$ is compactly supported, $d \ge p > \frac{d}{2}$, then:

$$\sup_{\xi \in \mathbb{R}^d} \int_{\mathbb{R}^d} \frac{1}{|\xi - z|^{d-2+d}} |f(z)| \ dz < \infty$$

for all $0 < \alpha < 2 - \frac{d}{p}$. Assume supp $f \subseteq \overline{B(0, R_1)}$. Consider 2 cases:

• If $|\xi| > 2R_1$, then: $|\xi - z| \ge R_1$ for all $z \in B(0, R_1)$. Hence:

$$\int_{\mathbb{R}^d} \frac{1}{|\xi - z|^{d-2+\alpha}} |f(z)| \, dz \leqslant \frac{1}{R_1^{d-2+\alpha}} ||f||_{L^1} < \infty$$

• If $|\xi| \leq 2R_1$, then: $|\xi - z| \leq 3R_1$ for all $z \in B(0, R_1)$:

$$\begin{split} \int_{\mathbb{R}^{d}} \frac{1}{|\xi - z|^{d - 2 + \alpha}} |f(z)| \, dz & \leq \int_{|\xi - z|} \frac{1}{|\xi - z|^{d - 2 + \alpha}} |f(z)| \, dz \\ \text{(H\"{o}lder)}, \left(\frac{1}{p} + \frac{1}{q} = 1\right) & \leq \left(\int_{\mathbb{R}^{d}} |f(z)|^{p} \, dz\right)^{\frac{1}{p}} \\ & \cdot \left(\int_{|\xi - z|} \frac{1}{|\xi - z|^{(d - 2 + \alpha)q}}\right)^{\frac{1}{q}} \\ & = \|f\|_{L^{p}} \left(\int_{|z| \leqslant 3R_{1}} \frac{1}{|z|^{(d - 2 + \alpha)q}} \, dz\right)^{\frac{1}{q}} < \infty \end{split}$$

c) $(d \ge 3)$ We already know:

$$\partial_i(G \star f) = (\partial_i G \star f) \in L^1_{loc}(\mathbb{R}^d)$$

as $\omega_d f \in L^1(\mathbb{R}^d)$. We claim that $\partial_i G \star f \in C^{0,\alpha}(\mathbb{R}^d)$. So $G \star f \in C^{1,\alpha}(\mathbb{R}^d)$ by the equivalence between the classical and the distributional derivatives. Exercise. Hint:

$$|\partial_i G \star f(x) - \partial_i G \star f(y)| \le \int_{\mathbb{R}^d} |\partial_i G(x - z) - \partial_i G(y - z)| |f(z)| dz,$$

$$\partial_i G(x) = \frac{-x_i}{d|B_1||x|^d}$$
. \rightsquigarrow Need to estimate $|\partial_i G(x) - \partial_i G(y)| \leq C|x - y|^{\alpha}$.

Theorem 4.13 (High regularity for Poisson's equation) Let $f \in C^{0,\alpha}(\mathbb{R}^d)$, $0 < \alpha < 1$ be compactly supported. Then $G \star f \in C^{2,\alpha}(\mathbb{R}^d)$.

Remark 4.14 $(-\Delta u = f)$ and $f \in C(\mathbb{R}^d)$ does not imply that $u \in C^2(\mathbb{R}^d)$. (exercise)

Remark 4.15 If $f \in C^{k,\alpha}(\mathbb{R}^d)$, $k \in \{0,1,\ldots\}$, $0 < \alpha < 1$ is compactly supported, then $G \star f \in C^{k+2,\alpha}(\mathbb{R}^d)$. This more general statement is a consequence of the theorem since

$$D^{\beta}(G\star f) = G\star\underbrace{(D^{\beta}f)}_{\in C^{0,\alpha}}$$

for all $\beta = (\beta_1, \dots, \beta_d), |\beta| \leq k$.

Proof of theorem 4.13. Since $f \in L^p$ for all $p \leq \infty$ by the low regularity (4.11) we have $G \star f \in C^{1,\alpha}$ and $\partial_i(G \star f) = \partial_i G \star f$ in the classical sense. We will compute the distributional derivatives $\partial_i \partial_j (G \star f)$ and prove that they are Hölder continuous. Compute $\partial_j \partial_i (G \star f)$: For all $\phi \in C_c^{\infty}(\mathbb{R}^d)$ we have

$$\begin{aligned} -(\partial_j \partial_i G \star f)(\phi) &= (\underbrace{\partial_i (G \star f)}_{\in C})(\partial_j \phi) \\ &= \int_{\mathbb{R}^d} ((\partial_i G) \star f)(x) \partial_j \phi(x) \, dx \\ &= \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} \partial_i G(x - y) f(y) \partial_j \phi(x) \, dx \, dy \\ &= \int_{\mathbb{R}^d} f(y) \left[\int_{\mathbb{R}^d} \partial_i G(x - y) \partial \phi(x) \, x \right] \, dy \\ &\stackrel{?}{=} \int_{\mathbb{R}^d} \Box \phi(y) \, dy \end{aligned}$$

Recall: $\partial_i G(x) = \frac{-x_i}{d|B_1||x|^d}$, $\partial_i \partial_j G(x) = \frac{1}{|B_1|} \left[\frac{x_i x_j}{|x|^2} - \frac{\delta_{ij}}{d} \right] \frac{1}{|x|^d}$. We have:

$$\int_{\mathbb{R}^d} \partial_i G(x - y) \partial_j \phi(x) \, dx = \lim_{\epsilon \to 0^+} \int_{|x - y| \ge \epsilon} \partial_i G(x - y) \partial_j \phi(x) \, dx$$

By dominated convergence we have $|\partial_i G(x-y)\partial_j \phi(x)| \in L^1(dx)$. By the Gauss-Green-Theorem (2.2) for all $\epsilon > 0$:

$$\int_{|x-y| \ge \epsilon} \partial_i G(x-y) \partial_j \phi(x) dx$$

$$= \int_{\partial B(y,\epsilon)} \partial_i G(x-y) \phi(x) \omega_j dS(x) - \int_{|x-y| \ge \epsilon} \partial_j \partial_i G(x-y) \phi(x) dx$$

Where $\omega = \frac{x-y}{|x-y|}$. For the boundary term:

$$\begin{split} -\int_{\partial B(y,\epsilon)} \partial_i G(x-y) \phi(x) \omega_j \, dS(x) &= -\epsilon^{d-1} \int_{\partial B(0,1)} \partial_i G(\epsilon \omega) \phi(y+\epsilon \omega) \omega_j \, d\omega \\ (\star) &= \int_{\partial B(0,1)} \frac{1}{d|B_1|} \omega_i \omega_j \phi(y+\epsilon \omega) \, d\omega \\ &\xrightarrow{\epsilon \to 0} \int_{\partial B(0,1)} \frac{1}{d|B_1|} \; \omega_i \omega_j \phi(y) \, d\omega \\ &= \frac{1}{d} \delta_{i,j} \phi(y) \end{split}$$

$$(\star)$$
 $\partial_i G(x) = \frac{-x_i}{d|B_1||x|^d}$, so $\partial_i G(\epsilon \omega) = -\frac{-\omega_i}{d|B_1|} \frac{1}{\epsilon^{d-1}}$. for all $|\omega| = 1$.

Now we split:

$$\begin{split} &-\int_{|x-y|} \underset{\geq \epsilon}{\partial_i \partial_j G(x-y) \phi(x)} \, dx \\ &= -\int_{|x-y|} \underset{\geq 1}{\partial_i \partial_j G(x-y) \phi(x)} \, dx - \int_{1 \geqslant |x-y| \geqslant \epsilon} \partial_i \partial_j G(x-y) \phi(x) \, dx \end{split}$$

The key observation is: $\int_{\partial B(0,r)} \partial_i \partial_j G(x) dx = 0$ since

$$\partial_i \partial_j G(x) = \frac{1}{|B_1|} \left(\omega_i \omega_j - \frac{\partial_{ij}}{d} \right) \frac{1}{|x|^d},$$

 $\omega = \frac{x}{|x|}$. For example if i = 1, j = 2, r = 1:

$$\int_{\partial B(0,1)} \partial_1 \partial_2 G(x) \, dS(x) = \frac{1}{|B_1|} \int_{\partial B(0,1)} \omega_1 \omega_2 \, d\omega,$$

 $\partial B(0,1) = \{\omega \mid |\omega| = 1\}$. Consider: $\omega \mapsto R\omega, (\omega_1, \dots, \omega_d) \mapsto (-\omega_1, \omega_2, \dots, \omega_d)$. Then

$$-\int_{1\geqslant |x-y|\geqslant \epsilon} \partial_i \partial_j G(x-y)\phi(y) \, dx = 0.$$

So

$$-\int_{1\geqslant |x-y|\geqslant \epsilon} \partial_i \partial_j G(x-y)\phi(x) \, dx = -\int_{1\geqslant |x-y|\geqslant \epsilon} \partial_i \partial_j G(x-y)(\phi(x)-\phi(y)) \, dx$$

In summary:

$$\begin{split} \partial_i \partial_j (G \star f)(\phi) &= \int_{\mathbb{R}^d} f(y) \left(\int_{\mathbb{R}^d} \partial_i G(x-y) \partial_j \phi(x) \, dx \right) \, dy \\ &= \int_{\mathbb{R}^d} f(y) \frac{1}{d} \partial_{ij} \phi(y) \, dy \\ &- \int_{\mathbb{R}^d} f(y) \left(\int_{|x-y| > 1} \partial_i \partial_j G(x-y) \phi(x) \, dx \right) \\ &- \int_{\mathbb{R}^d} \left[\lim_{\epsilon \to 0} \int_{1 \geqslant |x-y| \geqslant \epsilon} \underbrace{\frac{\partial_i \partial_j G(x-y) (\phi(x) - \phi(y)) \, dx}{\sum_{|x-y|^d} |x-y| \|\nabla \phi\|_L \infty \leqslant \frac{C}{|x-y|^{d-1}} \epsilon L^1_{loc}(dx) \forall y} \right] \, dy \\ &= \int_{\mathbb{R}^d} \frac{\delta_{ij}}{d} f(x) \phi(x) \, dx - \int_{\mathbb{R}^d} \phi(x) \left(\int_{|x-y| > 1} \partial_i \partial_j G(x-y) f(y) \, dy \right) \, dx \\ &- \int_{\mathbb{R}^d} \phi(x) \left[\int_{|x-y| \leqslant 1} \partial_i \partial_j G(x-y) (f(y) - f(x)) \, dy \right] \, dx \end{split}$$

Conclusion:

$$\partial_i \partial_j (G \star f)(x) = -\frac{\delta_{ij}}{d} f(x) + \int_{|x-y|>1} \partial_i \partial_j G(x-y) f(y) \, dy$$
$$+ \int_{|x-y| \leq 1} \partial_i \partial_j G(x-y) \left(f(y) - f(x) \right) \, dy$$

The first term $f \in C^{0,\alpha}$. The second term is also at least $C^{0,\alpha}$ since $\partial_i \partial_j G(x)$ is smooth as |x| > 1. We need to prove that the thirt term

$$W_{ij}(x) = \int_{|x-y| \le 1} \partial_i \partial_j G(x-y) (f(y) - f(x)) \, dy$$

is Hölder-continuous, $|W_{ij}(x) - W_{ij(y)}| \leq C|x - y|^{\alpha}$. Recall:

$$|\partial_i \partial_j G(x-y)(f(y)-f(x))| \leqslant C \frac{1}{|x-y|^d} |x-y|^\alpha = \frac{C}{|x-y|^{d-\alpha}} \in L^1_{loc}(dy)$$

We write

$$W_{ij}(x) = \int_{|x-y| \le 1} \partial_i \partial_j G(x-y) (f(y) - f(x)) \, dy$$
$$= \int_{|z| \le 1} \partial_i \partial_j G(z) (f(x+z) - f(x)) \, dz$$

So we get:

$$W_{ij} - W_{ij}(y) = \int_{|z| \le 1} \partial_i \partial_j G(z) (f(x+z) - f(y+z) - f(x) + f(y)) dz$$

Easy thought: Use $\partial_i \partial_j G(z) | \leq \frac{C}{|z|^d}$ and

$$|f(x+z) - f(y+z) - f(x) + f(y)|$$

$$\leq \begin{cases} |f(x+z) - f(x)| + |f(y+z) - f(y)| \leq C|z|^{\alpha} \\ |f(x+z) - f(y+z)| + |f(x) - f(y)| \leq C|x-y|^{\alpha} \end{cases}$$

Thus:

$$|W_{ij}(x) - W_{ij}(y)| \le C \int_{|z| \le 1} \frac{1}{|z|^d} \min(|z|^\alpha, |x - y|^\alpha) \, dz$$

$$\le C \int_{|z| \le 1} \frac{1}{|z|^d} (|z|^\alpha)^\epsilon (|x - y|^\alpha)^{1 - \epsilon}, \quad 0 < \epsilon < 1$$

$$\le C \left(\int_{|z| \le 1} \frac{1}{|z|^{d - \alpha \epsilon}} \right) |x - y|^{\alpha (1 - \epsilon)}$$

$$\le C_\epsilon |x - y|^{\alpha (1 - \epsilon)}$$

thus it is easy to prove $|W_{ij}(x) - W_{ij}(y)| \leq C_{\alpha}|x - y|^{\alpha}$ for all $\alpha' \leq \alpha$. However, to get $\alpha' = \alpha$ we need a more precise estimate. We split:

$$W_{ij}(x) - W_{ij}(y) = \int_{|z| \le 1} \dots = \int_{|z| \le \min(4|x-y|,1)} + \int_{4|x-y| < |z| \le 1}$$

For the first domain:

$$\int_{|z| \leq 4|x-y|} |\partial_{ij}G(z)||f(x+z) - f(y+z) - f(y) + f(x)| dz$$

$$\leq C \int_{|z| \leq 4|x-y|} \frac{1}{|z|^d} |z|^\alpha dz = const \cdot |x-y|^\alpha$$

For the second domain:

$$\int_{4|x-y|<|z|\leq 1} \partial_{ij} G(z) (f(x+z) - f(y+z) + f(y)f(x)) dz$$

$$= \int_{4|x-y|$$

since $\int_{4|x-y|<|z|\leq 1} \partial_{ij} G(z) dz = 0$. Then

$$(\ldots) = \int_{4|x-y| < |z-x|} \partial_{ij} G(z-x) f(z) \, dz - \int_{4|x-y| < |z-y|} \partial_{ij} G(z-y) f(z) \, dz.$$

Denote $A = \{z \mid 4|x-y| < |z-x| \leqslant 1\}, B = \{z \mid 4|x-y| < |z-y| \leqslant 1\}.$ Consider

$$\int_{A} \partial_{ij} G(z-x) f(z) dz - \int_{B} \partial_{ij} G(z-y) f(z) dz$$

$$= \int_{A \setminus B} + \int_{B \setminus A} + \int_{A \cap B} (\partial_{ij} G(z-x) - \partial_{ij} G(z-y)) f(z) dz$$

Lets regard the intersection. We have

$$\partial_{ij}G(x) = \frac{1}{|B_1|} \frac{1}{|x|^d} (\omega_i \omega_j - \frac{1}{d} \delta_{ij})$$
$$|\partial_{ij}G(x) - \partial_{ij}G(y)| \le C|x - y| \left(\frac{1}{|x|^{d+1}} + \frac{1}{|y|^{d+1}}\right)$$

Now,

$$|\partial_{ij}G(z-x) - \partial_{ij}G(z-y)| \le C|x-y|\left(\frac{1}{|z-x|^{d+1}} + \frac{1}{|z-y|^{d+1}}\right)$$

So we have

$$\left| \int_{A \cap B} (\partial_{ij} G(z - x) - \partial_{ij} G(z - y)) f(z) dz \right|$$

$$\leq C \int_{A \cap B} |x - y| \left(\frac{1}{|z - x|^{d+1}} + \frac{1}{|z - y|^{d+1}} \right) |f(z)| dz = (\dots)$$

Now we replace f(z) by f(z) - f(x), then:

$$\left| \int_{A \cap B} (\partial_{ij} G(z - x) - \partial_{ij} G(z - y))(f(z) - f(x)) dz \right|$$

$$\leqslant C \int_{A \cap B} |x - y| \left(\frac{1}{|z - x|^{d+1}} + \frac{1}{|z - y|^{d+1}} \right) |z - x|^{\alpha} dz$$

$$= C \underbrace{\int_{A \cap B} |x - y| \frac{1}{|z - x|^{d+1-\alpha}} dz}_{(I)} + \underbrace{C \int_{A \cap B} |x - y| \frac{1}{|z - y|^{d+1}} |z - x|^{\alpha} dz}_{(II)}$$

Now,

$$\begin{split} (I) &\leqslant C|x-y| \int_{4|x-y|<|z-x|\leqslant 1} \frac{1}{|z-x|^{d+1-\alpha}} \, dz \\ &= C|x-y| \int_{4x-y<|z|\leqslant 1} \frac{1}{|z|^{d+1-\alpha}} \, dz \\ &\leqslant C|x-y| \int_{4|x-y|}^{1} \frac{1}{r^{d+1-\alpha}} r^{d-1} \, dr \\ &= C|x-y| \int_{4|x-y|}^{1} \frac{1}{r^{2-\alpha}} \, dr \\ &\leqslant C|x-y| \left[-1 + \frac{1}{(4|x-y|)^{1-\alpha}} \right] \\ &\leqslant C|x-y|^{\alpha} \end{split}$$

$$\begin{split} (II) & \leqslant C|x-y| \int_{A \cap B} \frac{1}{|z-y|^{d+1}} |z-x|^{\alpha} \, dz \\ & \leqslant C|x-y| \int_{A \cap B} \frac{1}{|z-y|^{d+1}} \left(|z-y|^{\alpha} + |x-y|^{\alpha} \right) \, dz \\ & \leqslant \underbrace{C|x-y| \int_{B} \frac{1}{|z-y|^{d+1-\alpha}} \, dz}_{\text{similar to (I)}} \, dz + C|x-y|^{1+\alpha} \int_{B} \frac{1}{|z-y|^{d+1}} \, dz \end{split}$$

and

$$C|x-y|^{1+\alpha} \int_{B} \frac{1}{|z-y|^{d+1}} \, dz \leqslant \int_{4|x-y|} \frac{1}{r^{d+1}} r^{d-1} \, dr \leqslant \frac{C}{|x-y|}$$

Consider $A \backslash B$:

$$\left| \int_{A \backslash B} \right| \leqslant C \|f\|_{L^{\infty}} \int_{A \backslash B} \frac{1}{|z - x|^d} \, dz$$

where

$$A = \{z \mid 4|x - y| < |z - x| \le 1\}$$

$$B = \{z \mid 4|x - y| < |z - y| \le 1\}$$

$$A \setminus B = \{z \in A \mid |z - y| \le 4|x - y|\} \cup \{z \in A \mid |z - y| > 1\} = E_1 \cup E_2$$

for

$$E_1 = \{ z \mid |z - y| \le 4|x - y| < |z - x| \le 1 \}$$

$$\subseteq \{ z \mid 4|x - y| \le |x - z| \le 5|x - y| \}.$$

$$|x-z| \le |x-y| + |y-z| \le 5|x-y|$$
 in E_1 . We have

$$\begin{split} \int_{E_1} \frac{1}{|z-x|^d} \, dz &\leqslant \int_{4|x-y|} \frac{1}{|z-x|^{d-\alpha}} \, dz \\ &= \int_{4|x-y|} \frac{1}{|z|^{d-\beta}} \, dz \\ &= \int_{4|x-y|} \frac{1}{|r^d} r^{d-1} \, dr \\ &= \int_{4|x-y|} \frac{1}{r^{1-\alpha}} \, dr \\ &\leqslant C|x-y|^{\alpha} \end{split}$$

Now in E_2 : $|z - x| \ge |z - y| - |y - x| \ge 1 - |y - x|$.

$$\int_{E_2} \frac{1}{|z - x|^{d - \alpha}} dz \le \int \frac{1}{|z - x|^{d - \alpha}} dz = \int_{1 - |x - y|}^{1} \frac{1}{r^{d - \alpha}} r^{d - 1} dr$$

$$\le const. \left| 1 - \frac{1}{(1 - |x - y|)^{\alpha}} \right| \le C|x - y|^{\alpha}$$

Exercise 4.16 (E 5.1) Prove that if f is a harmonic function in \mathbb{R}^d and $g \in C_c(\mathbb{R}^d)$ is radial, then

$$\int_{\mathbb{R}^d} f(x)g(x) dx = f(0) \int_{\mathbb{R}^d} g(x) dx$$

Solution. $x = r\omega, r > 0, |\omega| = 1$

$$\int_{\mathbb{R}^d} f(x)g(x) dx \stackrel{\text{(Polar)}}{=} \int_0^\infty \left(\int_{\partial B(0,1)} f(r\omega)g(r\omega) d\omega \right) dr$$

$$= \int_0^\infty \left(g_0(r) \int_{\partial B(0,1)} f(r\omega) d\omega \right) dr$$
(Mean value theorem (2.12))
$$= \int_0^\infty \left(g_0(r)f(0) \int_{\partial B(0,1)} d\omega \right) dr$$

$$= f(0) \int_0^\infty \left(\int_{\partial B(0,1)} g(r\omega) d\omega \right) dr$$

$$= f(0) \int_{\mathbb{R}^d} g(x) dx$$

Remark 4.17 Let $g \in C_c(\mathbb{R}^d)$ be radial. Why is $\int_{\mathbb{R}^3} \frac{g(x)}{|x|} dx \neq \infty$? Because $f(x) = \frac{1}{|x|}$ is harmonic in $\mathbb{R}^d \setminus \{0\}$ and sub-harmonic in \mathbb{R}^d , $-\Delta f = c\delta_0$.

Exercise 4.18 (E 5.2) Let $1 \leq p < \infty$. Let $\Omega \subseteq \mathbb{R}^d$ be open. Consider the Sobolev Space

$$W^{1,p}(\Omega) = \{ f \in L^p(\Omega) \mid \partial_{x_i} f \in L^p(\Omega), \forall i = 1, 2, \dots, d \}$$

with the norm

$$||f||_{W^{1,p}} = ||f|| + \sum_{i=1}^d ||\partial_{x_i} f||_{L^p(\Omega)}.$$

Prove that $W^{1,p}(\Omega)$ is a Banach space. Here $x = (x_i)_{i=1}^d \in \mathbb{R}^d$. Hint: You can use the fact that $L^p(\Omega)$ is a Banach Space.

Solution. $W^{1,p}(\Omega) \subseteq L^p(\Omega) \times L^p(\Omega) \cdots \times L^p(\Omega) = (L^p(\Omega))^{d+1}$. For an element $f \in W^{1,p}(\Omega)$ we can think of it as $f \mapsto (f, \partial_1 f, \partial_2 f, \dots, \partial_d f)$, so $W^{1,p}(\Omega)$ is a subspace of $(L^p(\Omega))^{d+1}$, which is a norm-space. Why is $W^{1,p}(\Omega)$ closed in $(L^p(\Omega))^{d+1}$? Take $\{f_n\}_{n=1}^{\infty} \subseteq W^{1,p}(\Omega)$ such that $f_n \to f$ in L^p an $\partial_i f_n \to g_i$ in L^p for all $i=1,\dots,d$. We prove that $(f,g_1,\dots,g_d) \in W^{1,p}(\Omega)$, i.e. $f \in W^{1,p}$ and $g_i = \partial_i f$ for all $i=1,\dots,d$. We know that $f_n \to f$ in $L^p(\Omega)$, so $f_n \to f$ in $D'(\Omega)$ and $\partial_i f_n \to \partial_i f$ in $D'(\Omega)$. On the other hand we have $partial_i f_n \to g_i$ in $L^p(\Omega)$, so $\partial_i f_n \to g_i$ in $D'(\Omega)$. So we get $\partial_i f = g_i \in L^p(\Omega)$ for all $i=1,\dots,d$ in $D'(\Omega)$. So we can conclude $f \in W^{1,p}(\Omega)$ and $\partial_i f = g_i$ for all $i=1,\dots,d$.

Exercise 4.19 (E 5.3) Let f be a real-valued function in $W^{1,p}(\mathbb{R}^d)$ for some $1 \le p < \infty$. Prove that $|f| \in W^{1,p}(\mathbb{R}^d)$ and

$$(\nabla |f|)(x) = \begin{cases} \nabla f(x) & f(x) > 0 \\ -\nabla f(x) & f(x) < 0 \\ 0 & f(x) = 0 \end{cases}$$

Solution. Consider $G_{\epsilon}(t) = \sqrt{\epsilon^2 + t^2} - \epsilon$ for $\epsilon > 0$, $t \in \mathbb{R}$. Clearly we have $G_{\epsilon}(t) \to |t|$ as $\epsilon \to 0$ and

$$G'_{\epsilon}(t) = \frac{2t}{2\sqrt{\epsilon^2 + t^2}} = \frac{t}{\sqrt{\epsilon^2 + t^2}}$$

so $|G'_{\epsilon}(t)| \leq 1$, $G_{\epsilon}(0) = 0$. By the chain rule, $G_{\epsilon}(f) \in W^{1,p}(\mathbb{R}^d)$ and

$$(\partial_i G_{\epsilon}(f))(x) = G'_{\epsilon}(f)\partial_i f(x) = \frac{f(x)}{\sqrt{\epsilon^2 + f^2(x)}}\partial_i f(x) \in L^p(\mathbb{R}^d)$$

for all $i=1,\ldots,d$. Note then when $\epsilon\to 0$ that $G_\epsilon(f)(x)\to |f(x)|$ pointwise, so $G_\epsilon(f)\to |f|$ in $L^p(\mathbb{R}^d)$. $|G_\epsilon(f)(x)-G_\epsilon(0)|\leqslant |f(x)|\in L^p(\mathbb{R}^d)$ by dominated convergence.

$$\partial_i G_{\epsilon}(f)(x) = \frac{f(x)}{\sqrt{\epsilon^2 + f^2(x)}} \partial_i f(x) \xrightarrow{\epsilon \to 0} g_i(x) := \begin{cases} \partial f_i(x) & f(x) > 0 \\ -\partial_i f(x) & f(x) < 0 \\ 0 & f(x) = 0 \end{cases}$$
$$|\partial_i G_{\epsilon}(f)(x)| \leqslant \left| \frac{f(x)}{\sqrt{\epsilon^2 + f^2(x)}} \right| |\partial_i f(x)| \leqslant |\partial_i f(x)| \leqslant |D_i f(x)| \leqslant$$

So we get $\partial_i G_{\epsilon}(f) \xrightarrow{\epsilon \to 0} g_i$ in $L^p(\mathbb{R}^d)$ by Dominated Convergence. So we conclude: $\partial_i(|f|) = g_i \in L^p(\mathbb{R}^d)$ for all $i = 1, \ldots, d$, so $|f| \in W^{1,p}(\mathbb{R}^d)$, $|f| \in L^p$.

Exercise 4.20 (E 5.4) Let $\Omega \subseteq \mathbb{R}^d$ be open and bounded, $f \in L^1(\Omega)$,

$$u(x) = \int_{\Omega} G(x - y) f(y) \, dy$$

Let $-\Delta u = f$ in $D'(\Omega)$, $u \in L^1_{loc}(\Omega)$, $f \in L^1_{loc}(\mathbb{R}^d)$ and $\omega_d f \in L^1(\mathbb{R}^d)$, where

$$\omega_d(x) = \begin{cases} 1 + |x| & d = 1\\ \log(1 + |x|) & d = 1\\ \frac{1}{(1+|x|)^{d-2}} & d \geqslant 3 \end{cases}$$

Prove that

$$G \star f = \int_{\mathbb{R}^d} G(x - y) f(y) \, dy \in L^1_{loc}(\mathbb{R}^d)$$

and $-\Delta(G \star f) = f$ in $D'(\mathbb{R}^d)$.

Solution. Define $\tilde{f} = \mathbb{1}_{\Omega}(x)f(x) = \begin{cases} f(x) & x \in \Omega \\ 0 & x \notin \Omega \end{cases}$. Then

$$u(x) = \int_{\Omega} G(x - y) f(y) dy = \int_{\mathbb{R}^d} G(x - y) \tilde{f}(y) dy = (G \star \tilde{f})(x)$$

We have $u \in L^1_{loc}(\mathbb{R}^d)$, so $u \in L^1(\Omega)$. Then $-\Delta u = \tilde{f}$ in $D'(\mathbb{R}^d)$, so $-\Delta u = f$ in $D'(\Omega)$. Claim: $-\Delta u = f$ in $D'(\mathbb{R}^d)$, so $-\Delta u = f$ in $D'(\Omega)$ if $\Omega \subseteq \mathbb{R}^d$, $\tilde{f}|_{\Omega} = f$. Take $\phi \in C^\infty_c(\Omega)$. We need: $(-\Delta u)(\phi) \stackrel{?}{=} \int_{\Omega} f \phi$. We have $\phi \in C^\infty_c(\Omega)$, so $\phi C^\infty_c(\mathbb{R}^d)$. This implies:

$$(-\Delta u)(\phi) = \int_{\mathbb{R}^d} \tilde{f}\phi = \int_{\substack{\Omega, \\ \text{supp } \phi \subseteq \Omega}} \tilde{f}\phi = \int_{\Omega} f\phi$$

Exercise 4.21 (E 5.5) Let $B = B\left(0, \frac{1}{2}\right) \subseteq \mathbb{R}^3$. Consider $u: B \to \mathbb{R}$, defined by $u(x) = \log |\log |x||$.

Prove that the distributional derivative $f = -\Delta u$ is a function in $L^{\frac{3}{2}}(B)$.

Solution.

$$\omega(r) = \log(-\log(r)), \quad \text{for } r \in \left(0, \frac{1}{2}\right)$$

$$\omega'(r) = \frac{1}{-\log(r)} \left(-\frac{1}{r}\right) = \frac{1}{r \log r}$$

$$\omega''(r) = -\frac{1}{(r \log(r))^2} (r \log(r))' = -\frac{\log(r) + 1}{(r \log r)^2}$$

So we have

$$-\Delta u = w''(r) = \frac{1}{(r \log r)^2} - \frac{1}{r^2 \log(r)} = f(r)$$

We show that $f \in L^{\frac{3}{2}}$:

$$\int_{B} |f(x)|^{\frac{3}{2}} dx = const \int_{0}^{\frac{1}{2}} \left| \frac{1}{r^{2}(\log r)^{2}} - \frac{1}{r^{2}\log r} \right|^{\frac{3}{2}} r^{2} dr$$

$$\tilde{\leq} \int_{0}^{\frac{1}{2}} \frac{1}{r} \left| \frac{1}{(\log(r))^{2}} - \frac{1}{(\log(r))} \right|^{\frac{3}{2}} dr$$

$$\begin{pmatrix} r = e^{-x}, \\ x \in (\log(2), \infty), \\ dr = -e^{-x} dx \end{pmatrix} \quad \tilde{\leq} \int_{\log(2)}^{\infty} e^{x} \left(\frac{1}{x^{2}} + \frac{1}{x} \right)^{\frac{3}{2}} e^{-x} dx$$

$$\tilde{\leq} \int_{\log(2)}^{\infty} \frac{1}{x^{\frac{3}{2}}} dx < \infty$$

Where $\tilde{<}$ means up to a constant. Now, $u(x) = \omega(r) = \log(-\log(r))$.

$$-\Delta u(x) = f(r) = \frac{1}{r^2(\log(r))^2} - \frac{1}{r^2\log(r)}$$

for all $x \neq 0, |x| = r < \frac{1}{2}$. Why is $-\Delta u(x) = f$ in D'(B)? Take $\phi \in C_c^{\infty}(B)$, check: $\int_B u(-\Delta \phi) = \int_B f \phi$.

$$\int_{|x|<\frac{1}{2}} u(-\Delta\phi) dx = \lim_{\epsilon \to 0^+} \int_{\epsilon < |x|<\frac{1}{2}} u(x)(-\Delta\phi)(x) dx$$

by Dominated convergence. $u \in L^1(B)$. For all $\epsilon > 0$:

$$\begin{split} \int_{\epsilon < |x| < \frac{1}{2}} u(x)(-\Delta \phi)(x) \, dx &= \int_{|x| > \epsilon} u(x)(-\Delta \phi)(x) \, dx \\ &= \int_{\partial B(0,\epsilon)} u(x) \nabla \phi(x) \frac{x}{|x|} \, dS(x) + \int_{|x| > \epsilon} \nabla u(x) \nabla \phi(x) \, dx \end{split}$$

The boundary term vanishes as $\epsilon \to 0$ since

$$\left| u(x)\nabla\phi(x)\frac{x}{|x|} \right| \leqslant \|\nabla\phi\|_{L^{\infty}}|u(x)| = C\log|\log(r)|$$

$$\left| \int_{\partial B(0,\epsilon)} u(x) \nabla \phi(x) \frac{x}{|x|} dS(x) \right| \leq C \int_{\partial B(0,\epsilon)} \log |\log(\epsilon)| dS(x)$$

$$= C \log |\log \epsilon| \underbrace{|\partial B(0,\epsilon)|}_{\sim \epsilon^2} \xrightarrow{\epsilon \to 0} 0$$

$$\int_{|x|>\epsilon} \nabla u(x) \nabla \phi(x) \, dx = \sum_{i=1}^d \int_{|x|>\epsilon} \partial_i u(x) \partial_i \phi(x) \, dx$$

$$= \sum_{i=1}^d \left(-\int_{\partial B(0,\epsilon)} \partial_i u(x) \phi(x) \frac{x_i}{|x|} \, dS(x) - \int_{|x|>\epsilon} \underbrace{\partial_i \partial_i u(x)}_{f(x)} \phi(x) \, dx \right)$$

The boundary term vanishes as $\epsilon \to 0$ as

$$\left| \int_{\partial B(0,\epsilon)} \partial u(x) \phi(x) \frac{x_i}{|x|} dS(x) \right| \leq \|\phi\|_{L^{\infty}} \int_{\partial B(0,\epsilon)} |\partial_i u(x)| dS(x)$$

$$(\star) \qquad \leq C \frac{1}{|\epsilon \log(r)|} |\partial B(0,\epsilon)| \to 0$$

as $\epsilon \to 0$. $(\star)u = u(r), u(x) = \omega(|x|), \partial_i u(x) = \omega(|x|) \frac{x_i}{|x|}, |\partial_i u(x)| \leq |\omega(|x|)| = \left|\frac{1}{r \log(r)}\right|$. Finally:

$$\int_{|x|>\epsilon} f(x)\phi(x) dx \xrightarrow{\epsilon \to 0} \int_{\mathbb{R}^d} f(x)\phi(x) dx$$

Since $f\phi \in L^1$ and Dominated Convergence.

Exercise 4.22 (Bonus 5) Construct $u \in L^1(\mathbb{R}^3)$ compactly supported s.t. $-\Delta u \in L^{\frac{3}{2}}(\mathbb{R}^3)$ and u is not continuous at 0.

Hint: Related to E 5.5. $u_0(x) = \omega(r) = \log(|\log(r)|)$ if $0 < r = |x| < \frac{1}{2}$. Consider χu_0 where $\chi \in C_c^{\infty}$, $\chi = 0$ if $|x| > \frac{1}{2}$, $\chi = 1$ if $|x| < \frac{1}{4}$. You can prove that $\Delta(\chi u_0) = (\Delta \chi) u_0 + 2\nabla \chi \nabla u_0 + \chi(\underbrace{\Delta u_0}_{c,L^{\frac{3}{2}}})$ in $D'(\mathbb{R}^3)$. (almost everywhere, in distributional

sense, integration by parts)

Theorem 4.23 (Regularity on Domains) Let $\Omega \subseteq \mathbb{R}^d$ be open. Assume $u, f \in D'(\Omega)$ such that $-\Delta u = f$ in $D'(\Omega)$.

- a) If $f \in L^1_{loc}(\Omega)$, then
 - $u \in C^1(\Omega)$ if d = 1
 - $u \in L^q_{loc}(\Omega)$ for all $q < \infty$ if d = 2
 - $u \in L^q_{loc}(\Omega)$ for all $q < \frac{d}{d-2}$ if $d \geqslant 3$
- b) If $f \in L^q_{loc}(\Omega)$, $d \geqslant p < \frac{d}{2}$, then $u \in C^{0,\alpha}_{loc}(\Omega)$, where $0 < \alpha < 2 \frac{d}{p}$
- c) If $f \in L^p_{loc}(\Omega)$, p > df, then $u \in C^{1,\alpha}_{loc}(\Omega)$, where $0 \le \alpha < 1 \frac{d}{p}$
- d) If $f \in C^{0,\alpha}_{loc}(\Omega)$ for some $0 < \alpha < 1$, then $u \in C^{2,\alpha}_{loc}(\Omega)$
- e) If $f \in C^{m,\alpha}_{loc}(\Omega)$, then $u \in C^{m+2,\alpha}_{loc}(\Omega)$

Proof. Let $\mathbb{K} \in \{\mathbb{R}, \mathbb{C}\}$. Take a ball $\overline{B} \subseteq \Omega$. Define $f_B : \mathbb{R}^d \to \mathbb{K}$,

$$f_B(x) = (\mathbb{1}_B f)(x) = \begin{cases} f(x) & x \in B \\ 0 & x \notin B \end{cases}$$

Then if $f \in L^1_{loc}(\Omega)$, f_B is compactly supported. From the previous theorems: $G \star f_B \in L^1_{loc}(\mathbb{R}^d)$ and $-\Delta(G \star f_B) = f_B$ in $D'(\mathbb{R}^d)$. On the other hand, $-\Delta u = f$ in $D'(\Omega)$, so $-\Delta(u - G \star f_B) = 0$ in D'(B). Indeed, for all $\phi \in C_c^{\infty}(B)$, then:

$$(-\Delta u)(\phi) = \int_{\Omega} f\phi = \int_{B} f_{B}\phi = -\int_{\mathbb{R}^{d}} f_{B}\phi = (-\Delta)(G \star f_{B})(\phi)$$

Then $-\Delta u = -\Delta(G \star f_B)$ in D'(B). Then $u - G \star f_B$ is harmonic in B and by Weyls lemma we have $u - G \star f_B \in C^{\infty}(B)$. So the smoothness of u in B is the same to that of $G \star f$.

Exercise 4.24 (E 6.1) Show that If $\chi \in C^{\infty}(\mathbb{R}^d)$, then $f \in W^{1,p}(\mathbb{R}^d)$, $1 \leq p < \infty$, then $\chi f \in W^{1,p}_{loc}(\mathbb{R}^d)$ and

$$\partial_i(\chi f) = (\partial_i \chi) f + \chi(\partial_i f)$$
 in $D'(\mathbb{R}^d)$

Solution. $\chi f \in L^p_{loc}(\mathbb{R}^d)$ obvious. $\partial(\chi f) \in L^p_{loc}(\mathbb{R}^d)$ is nontrivial but follows from $\partial_i(\chi f) = \underbrace{(\partial_i \chi) f + \chi(\partial f)}_{\mathcal{L}^p}$ in $D'(\mathbb{R}^d)$. To compute the distributional derivative

 $\partial_i(\chi f)$, then: Take $\phi \in C_c^{\infty}(\mathbb{R}^d)$:

$$-\int_{\mathbb{R}^d} \chi f(\partial \phi) = \int_{\mathbb{R}^d} (?) \phi$$

We have

$$-\int_{\mathbb{R}^d} \chi f(\partial_i \phi) = -\int_{\mathbb{R}^d} f(\chi \partial_i \phi)$$

$$(\chi \partial_i \phi = (\partial_i \chi) \phi + \chi(\partial_i \phi)) = -\int_{\mathbb{R}^d} f(\partial_i (\chi \phi) - (\partial_i \chi) \phi)$$

$$= -\int_{\mathbb{R}^d} f \partial_i (\underbrace{\chi \phi}_{\in C_c^{\infty}}) + \int_{\mathbb{R}^d} f(\partial_i \chi) \phi$$

$$= \int_{\mathbb{R}^d} (\partial_i f) \chi \phi + \int f(\partial_i \chi) \phi$$

$$= \int_{\mathbb{R}^d} ((\partial_i f) \chi + f(\partial_i \chi)) \phi$$

So
$$\partial_i(\chi f) = (\partial_i f)\chi + f(\partial_i \chi)$$
 in $D'(\mathbb{R}^d)$.

Remark 4.25 Question: If $\chi \in C^1(\mathbb{R}^d)$, $f \in W^{1,p}(\mathbb{R}^d)$. Is this it still correct that $\partial_i(\chi f) = (\partial_i \chi)f + \chi(\partial_i f)$ in $D'(\mathbb{R}^d)$?

Proof. It suffices to show that we still can apply intergration by parts.

$$(\star) \quad -\int f\partial_i g \stackrel{?}{=} \int (\partial_i f)g$$

Approximation: (*) is correct if $g \in C_c^{\infty}$

• If $g \in C_c^1$, there is $\{g_n\} \subseteq C_c^{\infty}$ s.t. $g_n \to g$ in $W_{loc}^{1,p}$, $\frac{1}{p} + \frac{1}{q} = 1$.

$$\int (\partial_i g) f \xleftarrow{n \to \infty} - \int \underbrace{f}_{L^p} \underbrace{\partial g_n}_{\to \partial_i g \text{ in } L^q} = \int \underbrace{(\partial_i f)}_{\in L^p} \underbrace{g_n}_{\to g \text{ in } L^q} \xrightarrow{n \to \infty} \int (\partial_i f) g$$

Exercise 4.26 (E 6.2) \mathbb{R}^2 , $G(x) = -\frac{1}{2\pi} \log |x|$. Let $f \in L^p(\mathbb{R}^d)$, compactly supported. Define $u(x) = (G \star f)(x) = \int_{\mathbb{R}^2} G(x-y) f(y) \, dy$

- 1. If p = 1, then $u \in L_{loc}^q(\mathbb{R}^2)$ for all $q < \infty$.
- 2. If p > 2, then $u \in C^{1,\alpha}$ with $0 < \alpha < 1 \frac{2}{p}$.

Solution. 1. Take any ball B = B(0, R) and:

$$\int_{B} |u(x)|^{q} dx = \int_{B} \left(\int_{\mathbb{R}^{d}} |G(x-y)| |f(y)| dy \right)^{q} dx$$

$$\leqslant C \int_{B} \left(\int_{\mathbb{R}^{2}} |G(x-y)|^{q} |f(y)| dy \right) dx$$

$$= C \int_{\mathbb{R}^{2}} \left(\int_{B} |G(x-y)|^{q} dx \right) |f(y)| dy$$

Recall from the proof of Youngs inequality:

$$\begin{split} |u(x)| &= \left| \int_{\mathbb{R}^2} G(x-y) f(y) \, dy \right| \\ &\leqslant \int_{\mathbb{R}^2} |G(x-y)| |f(y)| \ dy \\ &\leqslant \left(\int_{\mathbb{R}^2} |G(x-y)|^q |f(y)| \, dy \right)^{\frac{1}{q}} \left(\int_{\mathbb{R}^2} |f(y)| \, dy \right)^{\frac{1}{q}}, \quad \frac{1}{q} + \frac{1}{q'} = 1 \end{split}$$

Assume supp $f \subseteq \overline{B(0,R)}$. Then if $y \in \text{supp } f$ and $x \in B(0,R)$, then $|x-y| \le |x| + |y| \le R + R_1$. For all $y \in \text{supp } f$:

$$\int_{B(0,R)} |G(x-y)|^q dx \le \int_{|x-y| \le R+R_1} |G(x-y)|^q dx$$

$$= \int_{|z| \le R+R_1} |G(z)|^q dz < \infty$$

as $G \in L^q_{loc}$ $(|G(z)| = \frac{1}{2\pi} |\log(z)| \leq \frac{C_{R+R_1,\epsilon}}{|z|^{\epsilon}}$ for all $|z| \leq R + R_1$, so

$$\int_{|z| \leqslant R + R_1} |G(z)|^q \leqslant C_{R + R_1, \epsilon} \int_{|z| \leqslant R + R_1} \frac{1}{|z|^{\epsilon q}} dz < \infty$$

if $\epsilon q < 2$.

2. Recall $\partial_i u \in L^1_{loc}(\mathbb{R}^2)$ and:

$$\partial_i u(x) = (\partial_i G \star f)(x) = c \int_{\mathbb{R}^2} \frac{x_i - y_i}{|x - y|^2} f(y) \, dy$$

First we show $\partial_i u \in C^{0,\alpha}$:

$$|\partial_i u(x) - \partial_i u(z)| = \left| C \int_{\mathbb{R}^2} \left(\frac{x_i - y_i}{|x - y|^2} - \frac{z_i - y_i}{|z - y|^2} \right) f(y) dy \right|$$

$$\leqslant C \int_{\mathbb{R}^2} \left| \frac{x_i y_i}{|x - y|^2} - \frac{z_i - y_i}{|z - y|^2} \right| |f(y)| dy$$

$$\stackrel{?}{\leqslant} C|x - y|^{\alpha}$$

Note that

$$\left| \frac{x_i - y_i}{|x - y|^2} - \frac{z_i - y_i}{|z - y|^2} \right| = \left| (x_i - y_i) \left(\frac{1}{|x - y|^2} - \frac{1}{|z - y|^2} \right) + \frac{x_i - z_i}{|z - y|^2} \right|$$

$$\leq |x_i - y_i| \left| \frac{1}{|x - y|^2} - \frac{1}{|z - y|^2} \right| + \frac{|x_i - z_i|}{|z - y|^2}$$

$$\leq C|z - x|^{\alpha} \left(\frac{1}{|x - y|^{1+\alpha}} + \frac{1}{|z - y|^{1+\alpha}} + \frac{|x - z|}{|z - y|^2} \right)$$

Here $|x_i - z_i| \le |x - z|$ and $|x_i - y_i| \le |x - y|$ and:

$$\begin{split} \underbrace{\left|\frac{1}{|x-y|^2} - \frac{1}{|z-y|^2}\right|}_{\text{sym } x \leftrightarrow z} &= \left|\frac{1}{|x-y|} - \frac{1}{|z-y|}\right| \, \left|\frac{1}{|x-y|} + \frac{1}{|z-y|}\right| \\ &= \frac{||z-y| - |x-y|}{|x-y||z-y|} \left|\frac{1}{|x-y|} + \frac{1}{|z-y|}\right| \\ &\leqslant |z-x|^{\alpha} \frac{\max(|z-y|, |x-y|)^{1-\alpha}}{|x-y||z-y|} \left(\frac{1}{|x-y|} + \frac{1}{|z-y|}\right) \\ &\leqslant C|z-x|^{\alpha} \left(\frac{1}{|x-y|^{2+\alpha}} + \frac{1}{|z-y|^{2+\alpha}}\right) \end{split}$$

By the symmetrie $x \leftrightarrow z$:

$$LHS \leqslant C|z - x|^{\alpha} \left(\frac{1}{|x - y|^{1 + \alpha}} + \frac{1}{|z - y|^{1 + \alpha}} \right) + \frac{|x - y|}{|x - y|^2}$$

$$\Rightarrow LHS \leqslant C \dots + |x - z| \min \left(\frac{1}{|z - y|^2}, \frac{1}{|x - y|^2} \right)$$

$$\leqslant (|x - y| + |z - y|)^{1 - \alpha}$$

$$C|z - x|^{\alpha} \left(\frac{1}{|x - y|^{1 + \alpha}} + \frac{1}{|z - y|^{1 + \alpha}} \right)$$

In summary:

$$|\partial_i u(x) - \partial_i u(z)| \le C \int_{\mathbb{R}^2} \left| \frac{x_i - y_i}{|x - y|^2} - \frac{z_i - y_i}{|z - y|^2} \right| |f(y)| \, dy$$
$$= C|x - y|^{\alpha} \int_{\mathbb{R}^2} \left(\frac{1}{|x - y|^{1 + \alpha}} + \frac{1}{|z - y|^{1 + \alpha}} \right) |f(y)| \, dy$$

Consider if $|x| > 2R_1$:

$$\int_{\mathbb{R}^2} \frac{1}{|x-y|^{1+\alpha}} |f(y)| \, dy \leqslant \int_{\mathbb{R}^2} \frac{1}{R_1^{1+\alpha}} \, \left| f(y) \right| dy \leqslant C$$

supp $f \subseteq B(0,R_1)$. If $|x| < 2R_1$, then $|x-y| \leq 3R$ if $y \in B(0,R_1)$. Hence:

$$\begin{split} & \int_{|x-y| \leqslant 3R_1} \frac{1}{|x-y|^{1+\alpha}} |f(y)| \, dy \\ & \leqslant \left(\int_{|x-y| \leqslant 3R_1} \frac{1}{|x-y|^{(1+\alpha)p'}} \right)^{\frac{1}{p'}} \left(\int |f(y)|^p \, dy \right)^{\frac{1}{p}} \\ & = \int_{|z| \leqslant 3R_1} \frac{1}{|z|^{(1+\alpha)p'}} \, dz < \infty \end{split}$$

So $\alpha < 1 - \frac{2}{p}$.

Exercise 4.27 (E 6.3) Let $f \in C^{0,\alpha}_{loc}$ and $-\Delta u = f$ in $D'(\Omega)$. Prove $u \in C^{2,\alpha}_{loc}(\Omega)$.

Solution. Take an open ball $B \subseteq \bar{B} \subseteq \Omega$. We prove $u \in C^{2,\alpha}(B)$. There is an open Ω_B s.t. $\bar{B} \subseteq \bar{\Omega}_B \subseteq \Omega$. Then there is a $\chi_B \in C_c^{\infty}(\mathbb{R}^d)$ s.t. $\chi_B(x) = 1$ if $x \in B$ and $\chi_B(x) = 0$ if $x \notin \Omega_B$. Define

$$f_B(x) = \chi_B(x)f(x) : \mathbb{R}^d \to \mathbb{R}$$

We prove that $f_B \in C^{0,\alpha}(\mathbb{R}^d)$. Since $f \in C^{0,\alpha}_{loc}(\Omega)$ we have $f \in C^{0,\alpha}(\Omega)$, so $|f(x) - f(y)| \leq C|x - y|^{\alpha}$ for all $x, y \in \Omega_B$. Then:

$$|f_B(x) - f_B(y)| = |\chi_B(x)f(x) - \chi_B(y)f(y)|$$

$$\leq |(\chi_B(x) - \chi_B(y))f(x) + \chi_B(y)(f(x) - f(y))|$$

$$\leq C|x - y|^{\alpha}||f||_{L^{\infty}} + C||\chi||_{L^{\infty}(\Omega_B)}||x - y|^{\alpha} \leq C_{\Omega_B}||x - y|^{\alpha}$$

What about other cases? If x,y are bot not in Ω_B , then $|f_B(x)-f_B(y)|=0$, then if $x\in\Omega_B$ and $y\notin\Omega_B$: $|f_B(x)-f_B(y)|=|f_B(x)|=|\chi_B(x)||f(x)|=|\chi_B(x)-\chi_B(y)||f(x)|\leqslant C|x-y|^\alpha$. Conclusion: $|f_B(x)-f_B(y)|\leqslant C|x-y|^\alpha$ for all $x,y\in\mathbb{R}^d$, i.e. $f_B\in C^{0,\alpha}(\mathbb{R}^d)$. Also f_B is compactly supported. By a theorem in the lecture: $G\star f_B\in C^{2,\alpha}(\mathbb{R}^d)$. Finally: $-\Delta u=f$ in $D'(\Omega), -\Delta(G\star f_B)=f_B$ in $D'(\mathbb{R}^d)$. So we conclude $-\Delta u=f=f_B=-\Delta(G\star f_B)$ in $D'(B), -\Delta(u-G\star f_B)=0$ in D'(B), so $u-G\star f_B\in C^{\infty}(B)$, so $u\in C^{2,\alpha}(B)$.

Exercise 4.28 (E 6.4) $u, f \in L^2(\mathbb{R}^d), -\Delta u = f$ in $D'(\mathbb{R}^d)$. Prove $u \in W^{2,2}(\mathbb{R}^d), \|u\|_{W^{2,2}(\mathbb{R}^d)} \leq C (\|u\|_{L^2} + \|f\|_{L^2}).$

$$\begin{split} W^{2,2}(\mathbb{R}^d) &= \{g \in L^2(\mathbb{R}^d) \mid D^{\alpha}g \in L^2 \text{ for all } |\alpha| \leqslant 2 \} \\ &= \{g \in L^2(\mathbb{R}^d) \mid \widehat{D^{\alpha}g}(k) = (-2\phi i k)^{\alpha} \hat{g}(k) \in L^2(\mathbb{R}^d) \text{ for all } |\alpha| \leqslant 2 \} \\ &= \{g \in L^2(\mathbb{R}^d) \mid (1 + |k|^2) \hat{g}(k) \in L^2(\mathbb{R}^d) \} \end{split}$$

 $||u||_{W^{2,2}(\mathbb{R}^d)}$ is comparable $\int_{\mathbb{R}^d} (1+|k|^2)^2 |\hat{g}(k)|^2 dk$. If $D^{\alpha}g \in L^2$, then $\widehat{D^{\alpha}g}(k) = (-2\pi i k)^{\alpha} \hat{g}(k)$. For any $\phi \in C_c^{\infty}(\mathbb{R}^d)$:

$$\begin{split} \widehat{D^{\alpha}g}(k)\widehat{\phi}(k), dk &= \int (D^{\alpha}g)\phi = (-1)^{|\alpha|} \int g(D^{\alpha}\phi) \\ &= (-1)^{|\alpha|} \int \overline{\widehat{g}}(k)\widehat{D^{\alpha}} \ \phi(k) \\ &= (-1)^{|\alpha|} \int \overline{\widehat{g}}(k)(-2\pi i k)^{\alpha} \widehat{\phi}(k) \ dk \end{split}$$

so $\hat{D}^{\alpha}g(k)=(-1)^{|k|}\hat{g}(k)\overline{(-2\pi ikx)^{\alpha}}=\hat{g}(k)(-2\pi ik)^{\alpha}.$ This implies:

$$||u||_{W^{2,2}(\mathbb{R}^d)} \leq C \int_{\mathbb{R}^d} (1+|k|^2)^2 |\hat{u}(k)|^2 dk$$

$$= C \left(||u||_{L^2}^2 + \int_{\mathbb{R}^d} |k|^4 |\hat{u}(k)|^2 dk \right)$$

$$\leq C \left(||u||_{L^2}^2 + ||f||_{L^2}^2 \right)$$

$$\leq C (||u||_{L^2}^2 + ||f||_{L^2}^2)^2$$

Remark 4.29 (Bonus 6) Let $f, g \in W^{1,2}(\mathbb{R}^d)$. Prove that $fg \in W^{1,1}(\mathbb{R}^d)$ and

$$\partial_i(fg) = (\partial_i f)g + f(\partial_i g)$$
 in $D'(\mathbb{R}^d)$

Chapter 5

Existence for Poisson's Equation on Domains

Let $\Omega \subseteq \mathbb{R}^d$ be open. Consider Poisson's equation.

$$\begin{cases} -\Delta u = f & \text{in } \Omega \\ u = g & \text{on } \partial \Omega \end{cases}$$

for given data (f,g) and u the unknown function.

- Classical solutions: $f \in C^2(\bar{\Omega}) \leadsto$ explicit representation formula.
- Weak solution: $f \in L^p(\Omega)$, $g \in L^p(\partial\Omega) \rightsquigarrow u \in W^{2,p}(\Omega)$. We are going to establish the existence by *Energy Methods*. (Calculus of variations)

Definition 5.1 (C^1 -Domains) Let $\Omega \subseteq \mathbb{R}^d$ be open. We say that Ω is of class C^1 (i.e. $\partial \Omega \in C^1$) if for all $x_0 \in \partial \Omega$ there is a bijective function $h: U \to Q$, where

- $x_0 \in U$ open in \mathbb{R}^d
- $Q = \{x = (x_1, \dots, x_d) = (x', x_d)\} \in \mathbb{R}^{d-1} \times \mathbb{R} \mid |x'| < 1, |x_d| < 1\}$
- $h \in C^1(\bar{U})$ and $h^{-1} \in C^1(\bar{Q})$ (C^1 -diffeomorphism)
- h(U) = Q

$$h(U \cap \Omega) = Q_{+} = Q \cap \mathbb{R}^{d}_{+} = \{x = (x', x_{d}) \in Q \mid x_{d} > 0\}$$

$$h(U \cap \partial\Omega) = Q_{0} = Q \cap \partial\mathbb{R}^{d}_{+} = \{x = (x', x_{d}) \in Q \mid x_{d} = 0\}$$

$$h(U \setminus \bar{\Omega}) = Q_{-} = Q \cap \mathbb{R}^{d}_{-} = \{x = (x', x_{d}) \in Q \mid x_{d} < 1\}$$

(From Brezis' book)

Remark 5.2 The set Q can be replaced by a ball, i.e. Ω is of C^1 if for all $x_0 \in \partial \Omega$ there is a function $U \to B(0,1) \subseteq \mathbb{R}^d$.

- $x_0 \in U$ with $U \subseteq \mathbb{R}^d$ open.
- $h \in C^1(\bar{U}), h^{-1} \in C^1(\overline{B(0,1)})$
- $h(U \cap \Omega) = B(0,1) \cap \mathbb{R}^d_+, h(U \cap \partial\Omega) = B(0,1) \cap \mathbb{R}^d$.

Remark 5.3 (An equivalent definition form Evan's book App. C) Let $\Omega \subseteq \mathbb{R}^d$ be open. Then Ω is C^1 if for all $x_0 \in \partial \Omega$ there is a r > 0 and a C^1 -function $\gamma : \mathbb{R}^{d-1} \to \mathbb{R}$ s.t. (upon relabeling and reorienting the axes if necessary) such that:

$$\Omega \cap B(x_0, r) = \{x = (x', x_d) \in B(x_0, r) \mid x_d < \gamma(x_0)\}\$$

Proof of the equivalence of the two definitions.

Def. 2 \Rightarrow Def. 1: In fact, given $x_0 \in \partial \Omega$ and γ we can define

$$h(x', x_d) = (x', x_d - \gamma(x')) \in C^1(\mathbb{R}^d, \mathbb{R}^d)$$
$$h^{-1}(x', x_d) = (x', x_d + \gamma(x')) \in C^1(\mathbb{R}^d, \mathbb{R}^d)$$

Def. 1 \Rightarrow Def. 2: We need the inverse function theorem and the implicit function theorem. Let $x_0 \in \partial \Omega$, let $h: U \to B(0,1)$ as in Def. 1. Denote $h = (h_1,h_2,\ldots,h_d)$. Since h is invertible near x_0 , by the inverse function theorem we have for the Jacobi matrix $Jh(x_0) = (\partial_j h_i(x_0))_{1 \leqslant i,j \leqslant d}$ is invertible. So we have $\nabla h_d(x_0) = (\partial_j h_d(x_0))_{1 \leqslant j \leqslant d} \neq \vec{0}^{\mathbb{R}^d}$, so there is a $j \in \{1,2,\ldots,d\}$ s.t. $\partial_j h_d(x_0) \neq 0$. By relabeling and reorienting the axes, we can assume that $\partial_d h_d(x_0) > 0$. By continuity there is a r > 0 such that $\partial_d h_d(x) > 0$ for all $x \in B(x_0,r)$. Define $\gamma: \mathbb{R}^{d-1} \to \mathbb{R}$ s.t. in $B(x_0,r)$:

$$x = (x', x_d) \in \partial \Omega \iff h_d(x', x_d) = 0 \iff x_d = \gamma(x'),$$

 $h_d: \mathbb{R}^d \to \mathbb{R}$. This gives a solution γ if $\partial_d h_d > 0$ in $B(x_0, r)$. (For implicit function theorem, $\partial_d h_d(x_0) \neq 0$) Question: Why in $B(x_0, r)$?

$$x = (x', x_d) \in \Omega \iff x_d > \gamma(x')$$

Since $\partial_d h_d(x) > 0$ for all $x \in B(x_0, r)$ we have that $x_d \mapsto h_d(x', x_d)$ is strictly increasing, hence

$$x = (x', x_d) \in \Omega$$

$$\iff h(x', x_d) \in \mathbb{R}^d_+$$

$$\iff h_d(x', x_d) > 0 = h_d(x', \gamma(x'))$$

$$\iff x_d > \gamma(x')$$

Theorem 5.4 (Gauss-Green formula / Integration by parts) Let $\Omega \subseteq \mathbb{R}^d$ be open and bounded. Then

1. For all $u, v \in C^1(\bar{\Omega})$:

$$\int_{\Omega} (\partial_i u)v = -\int_{\Omega} u(\partial_i v) + \int_{\partial\Omega} uvn_i dS,$$

where $\vec{n} = (n_i)_{i=1}^d$ is the outwarded unit normal vector.

2. For all $u, v \in C^2(\bar{\Omega})$:

$$\int_{\Omega} u(-\Delta v) = \int_{\Omega} \nabla u \nabla v - \int_{\partial \Omega} u \frac{\partial v}{\partial \vec{n}} dS$$

where $\frac{\partial v}{\partial \vec{n}} = \nabla v \vec{n} = \sum_{i=1}^{d} \partial_i v n_i$.

Classical solutions via Green's function:

$$\begin{cases} -\Delta u = f & \text{in } \Omega \\ u = g & \text{on } \partial \Omega \end{cases}$$

Let $\Omega \subseteq \mathbb{R}^d$ be open, bounded, $\partial \Omega \in C^1$. Assume there exists a $u \in C^2(\bar{\Omega})$, $f \in C(\bar{\Omega})$, $g \in C(\partial \Omega)$. Let G be the fundamential solution of the Laplace Equation in \mathbb{R}^d . We use integration by parts in $\Omega \setminus B(x, \epsilon)$:

$$\begin{split} & \int_{\Omega \backslash B(x,\epsilon)} u(y)(-\Delta G)(y-x) \, dy \\ & = \int_{\Omega \backslash B(x,\epsilon)} \nabla u(y) \nabla G(y-x) \, dy - \int_{\partial \Omega \cup \partial B(x,\epsilon)} u(y) \frac{\partial G}{\partial \vec{n}}(y-x) \, dS(y) \\ & \int_{\Omega \backslash B(x,\epsilon)} G(y-x)(-\Delta u)(y) \, dy \\ & = \int_{\Omega \backslash B(x,\epsilon)} \nabla G(y-x) \nabla u(y) \, dy - \int_{\partial \Omega \cup \partial B(x,\epsilon)} G(y-x) \frac{\partial u}{\partial \vec{n}}(y) \, dS(y) \end{split}$$

This implies:

$$\begin{split} &\int_{\Omega \backslash B(x,\epsilon)} \left[u(y) (-\Delta G(y-x)) - G(y-x) (-\Delta u)(y) \right] \, dy \\ &= -\int_{\partial \Omega \cup \partial B(x,\epsilon)} \left[u(y) \frac{\partial G}{\partial \vec{n}} (y-x) - G(y,x) \frac{\partial u}{\partial \vec{n}} (y) \right] \, dS(y) \end{split}$$

for all $x \in \Omega, x \in B(x, \epsilon) \subseteq \Omega$. When $\epsilon \to 0$, then the left hand side converges to $-\int_{\Omega} G(y-x)f(y)\,dy$ and the right hand side (for $d \geqslant 2$) we have $\partial_j G(y) = \frac{-y_j}{d|B_1||y|^d}$, so

$$\frac{\partial G}{\partial \vec{n}} = \nabla G \vec{n} = \nabla G(y) \left(\frac{-y}{|y|} \right) = \sum_{i=1}^{d} \frac{-y_i}{d|B_1||y|^d} \frac{-y_j}{|y|} = \frac{1}{d|B_1||y|^{d-1}} \operatorname{on} \partial B(0, \epsilon)$$

so we have

$$\frac{\partial G}{\partial \vec{n}}(y-x) = \frac{1}{d|B_1|\epsilon^{d-1}}$$

on $\partial B(x,\epsilon)$. Hence

$$\int_{\partial B(x,\epsilon)} u(y) \frac{\partial G}{\partial \vec{n}}(y-x) dS(y) = \frac{1}{d|B_1|\epsilon^{d-1}} \int_{\partial B(x,\epsilon)} u(y) dS(y)$$
$$= \int_{\partial B(x,\epsilon)} u(y) dS(y) \xrightarrow{\epsilon \to 0} u(x)$$

On the other hand:

$$\left| \int_{\partial B(x,\epsilon)} G(y-x) \frac{\partial u(y)}{\partial \vec{n}} \, dS(y) \right| \leq C \epsilon^{d-1} \sup_{|z|=\epsilon} |G(z)| \xrightarrow{\epsilon \to 0} 0$$

since $|G(z)| \le \frac{C}{|z|^{d-2}}$ if $d \ge 3$, $|G(z)| \le C|\log(z)|$ if d=2 and $|G(z)| \le C|z|$ if d=1. In summary:

$$-\int_{\Omega} G(y-x)f(y) dy = -\int_{\partial\Omega} \left[u(y) \frac{\partial G}{\partial \vec{n}}(y-x) - G(y-x) \frac{\partial u}{\partial \vec{n}}(y) \right] dS(y) - u(x)$$

$$\Leftrightarrow u(x) = \int_{\Omega} G(y-x)f(y) dy + \int_{\partial\Omega} \left[G(y-x) \frac{\partial u}{\partial \vec{n}}(y) - g(y) \frac{\partial G}{\partial \vec{n}}(y-x) \right] dS(y)$$

Problem: We don't know anything about $\frac{\partial u}{\partial \vec{n}}$ on $\partial \Omega$. Trick: We can resolve that by using the *corrector* function: $\Phi_x = \Phi_x(y)$ which solves:

$$\begin{cases} -\Delta \Phi_x = 0 & \text{in } \Omega \\ \Phi_x(y) = G(y - x) & \text{on } \partial \Omega \end{cases}$$

We assume that Φ_x exists.

Definition 5.5 (Green's function) $\tilde{G}(x-y) = G(y-x) - \Phi_x(y)$ for all $x, y \in \Omega$, $x \neq y$.

Exercise 5.6 (E 7.1) Let $\Omega \subseteq \mathbb{R}^d$ be open and bounded with C^1 boundary. For $x \in \Omega$, assume there exist $\Phi_x(y)$, $y \in \bar{\Omega}$, s.t.

$$\begin{cases} \Delta_y \Phi_x(y) = 0\\ \Phi_x(y) = G(y - x) \end{cases},$$

 $G(z)=rac{1}{d(d-2)|B_1||z|^{d-2}}, d\geqslant 3.$ Prove that $\Phi_x(y)=\Phi_y(x)$ for all $x,y\in\Omega$. Then $\tilde{G}(x,y)=G(y-x)-\Phi_x(y)$ is symmetric, i.e. $\tilde{G}(x,y)=\tilde{G}(y,x)$.

Solution. Assume $x \neq y$. Define

$$f(z) = \tilde{G}(x, z) = G(z - x) - \Phi_x(z)$$

$$g(z) = \tilde{G}(y, z) = G(z - y) - \Phi_y(z)$$

Integration by parts:

$$\begin{split} \int_{\Omega \setminus (B(x,\epsilon) \cup B(y,\epsilon))} (f\Delta g - g\Delta f) &= \int_{\partial \Omega \cup \partial B(x,\epsilon) \cup \partial B(y,\epsilon)} \left(f \frac{\partial g}{\partial \vec{n_z}} - g \frac{\partial f}{\partial \vec{n_z}} \right) dS(z) \\ &= \int_{\partial B(x,\epsilon) \cup \partial B(y,\epsilon)} \left(f \frac{\partial g}{\partial \vec{n_z}} - g \frac{\partial f}{\partial \vec{n_z}} \right) dS(z) \end{split}$$

Consider $f \frac{\partial g}{\partial n_z^2}$ on $\partial B(x, \epsilon)$. Since g is only singular at y, so $\left| \frac{\partial g}{\partial \bar{n}} \right| \leqslant C$ on $\partial B(x, \epsilon)$. This implies:

$$\begin{split} \int_{\partial B(x,\epsilon)} \left| f \frac{\partial g}{\partial \vec{n_z}} \right| \, dS(z) &\leqslant C \int_{\partial B(x,\epsilon)} |f| \, dS(z) \\ &\leqslant C \int_{\partial B(x,\epsilon)} \left(\frac{1}{|x-z|^{d-2}} + \|\Phi_x\|_{L^{\infty}(\Omega)} \right) \, dS(z) \\ &\leqslant C \epsilon^{d-1} \left(\frac{1}{\epsilon^{d-2}} + 1 \right) \leqslant C \epsilon \xrightarrow{\epsilon \to 0} 0 \end{split}$$

Consider $f \frac{\partial g}{\partial \vec{n_z}}$ on $\partial B(y, \epsilon)$. Decompose $\frac{\partial g}{\partial \vec{n}} = \left[\nabla_z G(z-y) - \nabla_z \Phi_y(z)\right] \frac{(z-y)}{|z-y|}$. Since $\Phi_y(z)$ is harmonic in Ω , we have that

$$\int_{\partial B(y,\epsilon)} \left| f \nabla_z \Phi_y(z) \frac{-(z-y)}{|z-y|} \right| \le C \int_{\partial B(y,\epsilon)} |f| \le C \epsilon^{d-1} \xrightarrow{\epsilon \to 0} 0$$

Thus the main contribution from $f \frac{\partial g}{\partial \vec{n}}$ is

$$\begin{split} &\int_{\partial B(y,\epsilon)} f(z) \nabla_z G(z-y) \frac{-(z-y)}{|z-y|} \, dS(z) \\ &= \int_{\partial B(y,\epsilon)} f(z) \frac{-(z-y)}{d|B_1||z-y|^d} \frac{-(z-y)}{|z-y|} \, dS(z) \\ &= \frac{1}{d|B_1|\epsilon^{d-1}} \int_{\partial B(y,\epsilon)} f(z) \, dS(z) \\ &= \oint_{\partial B(y,\epsilon)} f(z) \, dS(z) = f(y) \end{split}$$

In summary:

$$\int_{\partial B(x,\epsilon)\cup\partial B(y,\epsilon)} f \frac{\partial g}{\partial \vec{n_z}} dS(z) \xrightarrow{\epsilon\to 0} f(y)$$

Similary:

$$\int_{\partial B(x,\epsilon)\cup\partial B(y,\epsilon)} g \frac{\partial f}{\partial \vec{n_z}} dS(z) \xrightarrow{\epsilon \to 0} g(x)$$

So we have that f(y) = g(x), so

$$f(y) = G(y - x) - \Phi_x(y)$$

$$g(x) = G(x - y) - \Phi_y(x).$$

So $\Phi_x(y) = \Phi_y(x)$ for all $x \neq y \in \Omega$. This implies $\Phi_x(y) = \Phi_y(x)$ for all $x, y \in \Omega$.

Theorem 5.7 Let $\Omega \subseteq \mathbb{R}^d$ be open, bounded and C^1 . If $u \in C^2(\Omega)$ solves

$$\begin{cases} -\Delta u = f & \text{in } \Omega \\ u = g & \text{on } \partial\Omega \end{cases},$$

then

$$u(x) = -\int_{\partial\Omega} g(y) \frac{\partial G}{\partial \vec{n_y}}(x, y) dS(y) + \int_{\Omega} \tilde{G}(x, y) dy$$

Proof. We need to prove:

$$\int_{\Omega} \Phi_x(y) f(y) \, dy + \int_{\partial \Omega} \left(-g(y) \frac{\partial \Phi_x(y)}{\partial \vec{n}_y} + G(y - x) \frac{\partial u}{\partial \vec{n}}(y) \right) = 0$$

By integration by parts:

$$\int_{\Omega} \Phi_{x}(y)f(y) \, dy = \int_{\Omega} \Phi_{x}(y)(-\Delta u(y)) \, dy$$

$$= \int_{\Omega} \left[\Phi_{x}(y)(-\Delta u(y)) + (\Delta \Phi_{x}(y))u(y) \right] \, dy$$

$$(\Delta \Phi_{x}(y) = 0) = \int_{\partial \Omega} \left(-\Phi_{x}(y) \frac{\partial u}{\partial \vec{n}} + \frac{\partial \Phi_{x}(y)}{\partial \vec{n}} \underbrace{u(y)}_{g(y)} \right) dS(y)$$

How can we compute $\Phi_x(y)$? It is not easy for general domains. But let us prove on two cases:

- $\Omega = \mathbb{R}^d_+$ (half-space)
- $\Omega = B(0, r)$ (a ball)

5.1 Green's function on the upper half plane

We use the following notation:

$$\mathbb{R}_{+}^{d} = \{ x = (x_{1}, x_{2}, \dots, x_{d}) = (x', x_{d}) \in \mathbb{R}^{d-1} \times \mathbb{R} \mid x_{d} > 0 \}$$
$$\partial \mathbb{R}_{+}^{d} = \{ x = (x', x_{d}) \mid x_{d} = 0 \} = \mathbb{R}^{d-1} \times \{0\}$$

For all $x \in \mathbb{R}^d$ we want to find the correction function $\Phi_x(y)$ with $y \in \overline{\mathbb{R}^d_+}$ s.t.

$$\begin{cases} +\Delta_y \Phi_x(y) = 0 & \text{in } \mathbb{R}^d_+ \\ \Phi_x(y) = G(y - x) & \text{in } \partial \mathbb{R}^d_+ \end{cases}$$

Definition 5.8 (Reflection for \mathbb{R}^d_+) For all $x = (x', x_d) \in \mathbb{R}^d$, $\tilde{x} = (x', -x_d) \in \mathbb{R}^d$, (if $x \in \mathbb{R}^d_+ \Rightarrow \tilde{x}\mathbb{R}^d_-$)

Claim: $\Delta_y \Phi_x(y) = G(y - \tilde{x})$ is a corrector function.

- $\Delta_y \Phi_x(y) = \Delta_y G(y \tilde{x}) = 0$ for all $y \in \mathbb{R}^d_+$ for all $x \in \mathbb{R}^d_+$ (as $\tilde{x} \in \mathbb{R}^d_- = \mathbb{R}^d \setminus \overline{\mathbb{R}^d_+}$)
- $\Phi_x(y) = G(y \tilde{x}) = G(y x)$ on $y \in \partial \mathbb{R}^d_+$. In fact, $y \in \partial \mathbb{R}^d_+$, so $y_d = 0$, so

$$G(y - \tilde{x}) = G_0(|y - \tilde{x}|) = G_0\left(\sqrt{\sum_{i=1}^{d-1} |x_i - y_i|^2 + |x_d|^2}\right) = G_0(|y - x|)$$

Consider f = 0 and

$$\begin{cases} -\Delta = 0 & \text{in } \mathbb{R}^d_+ \\ u = g & \text{on } \partial \mathbb{R}^d_+ \end{cases}$$

Then we expect

$$u(x) = -\int_{\partial\Omega} g(y) \frac{\partial \tilde{G}}{\partial \vec{n_y}}(x, y) dS(y)$$

We compute

$$\frac{\partial \tilde{G}}{\partial \vec{n_y}}(x-y) = \sum_{j=1}^d \frac{\partial \tilde{G}}{\partial y_j}(x,y)\vec{n_j} = -\frac{\partial \tilde{G}}{\partial y_d}(x,y) = \frac{\partial}{\partial y_d}(G(y-\tilde{x}) - G(y-x)) = \dots$$

because $\tilde{G}(x,y) = G(y-x) - \Phi_x(y) = G(y-x) - G(y-\tilde{x})$.

$$\dots = \frac{1}{d|B_1|} \left[\frac{-(y_d - \tilde{x}_d)}{|y - \tilde{x}|^d} - \frac{-(y_d - x_d)}{|y - x|^d} \right]$$
$$(y \in \partial \mathbb{R}^d_+) = \frac{1}{d|B_1|} \left[\frac{\tilde{x}_d}{|y - x|} - \frac{x_d}{|y - x|^d} \right] = \frac{-2x_d}{d|B_1||y - x|^d}$$

We expect

$$u(x) = -\int_{\partial \mathbb{R}^d} g(y) \frac{\partial \tilde{G}}{\partial \vec{n_y}}(x, y) \, dS(y) = \int_{\partial \mathbb{R}^d} g(y) \frac{2x_d}{d|B_1||y - x|^d} \, dS(y)$$

Theorem 5.9 Assume $g \in C(\mathbb{R}^{d-1}) \cap L^{\infty}(\mathbb{R}^{d-1})$ Then

$$u(x) = \int_{\partial \mathbb{R}^d_+} g(y) K(x, y) \, dS(y)$$

and

$$K(x,y) = \frac{2x_d}{d|B_1||y-x|^d}$$
 for all $x \in \mathbb{R}^d_+$.

satisfies that $u \in C^{\infty}(\mathbb{R}^d_+) \cap L^{\infty}(\mathbb{R}^d_+)$ and

$$\begin{cases} \Delta u = 0 & \text{in } \mathbb{R}^d_+ \\ \lim_{\substack{x \to 0 \\ x \in \mathbb{R}^d_+}} u(x) = g(x_0) & \forall x_0 \in \partial \mathbb{R}^d_+ \end{cases}$$

Proof. Claim: For all $y \in \partial \mathbb{R}^d_+$, $x \mapsto K(x,y)$ is harmonic in \mathbb{R}^d_+ (i.e. $\Delta_x K(x,y) = 0$ in \mathbb{R}^d_+)

• Argument from Evans:

$$K(x,y) = -\frac{\partial}{\partial y_d}, \ \tilde{G}(y-x) = -\frac{\partial}{\partial y_d}(G(y-x) - G(y-\tilde{x}))$$

We know that for all $x \in \mathbb{R}^d_+$, $y \mapsto \tilde{G}(y,x)$ is haromnic in $\mathbb{R}^d_+ \setminus \{x\}$. By symmetry we have $\tilde{G}(y,x) = \tilde{G}(x,y)$ for all $x,y \in \mathbb{R}^d_+$. So for all $y \in \mathbb{R}^d_+$, $x \mapsto \tilde{G}(y,x)$ is harmonic in $\mathbb{R}^d_+ \setminus \{y\}$. Then for all $y \in \mathbb{R}^d_+$: $-\frac{\partial}{\partial y_d} \tilde{G}(y,x) = K(x,y)$ is harmonic $x \in \mathbb{R}^d_+ \setminus \{y\}$. By a limit argument, for all $y \in \partial \mathbb{R}^d_+$, $x \mapsto K(x,y)$ is harmonic for all $x \in \mathbb{R}^d_+$.

• A direct proof:

$$K(x,y) = \frac{2x_d}{d|B_1|} \frac{1}{|x-y|^d}$$

for all $x \in \mathbb{R}^d_+$, $y \in \partial \mathbb{R}^d_+$. For $i \neq d$, $x = (x_1, \dots, x_d)$,

$$\begin{split} \partial_{x_i} K(x,y) &= \frac{2x_d}{d|B_1|} \frac{(-d)}{|x-y|^{d+1}} \frac{x_i - y_i}{|x-y|} = \frac{-2x_d}{|B_1|} \frac{x_i - y_i}{|x-y|^{d+2}} \\ \partial_{x_i}^2 K(x,y) &= -\frac{2x_d}{|B_1|} \left[\frac{1}{|x-y|^{d+1}} - \frac{(d+2)}{|x-y|^{d+3}} (x_i - y_i) \frac{(x_i - y_i)}{|x-y|} \right] \\ &= -\frac{2x_d}{|B_1|} \left[\frac{1}{|x-y|^{d+1}} - \frac{(d+2)}{|x-y|^{d+4}} (x_i - y_i)^2 \right] \end{split}$$

Moreover:

$$\begin{split} \partial_{x_d} K(x,y) &= \frac{2}{d|B_1|} \frac{1}{|x-y|^d} + \frac{2x_d}{d|B_1|} (-d) \frac{(x_d - y_d)}{|x-y|^{d+2}} \\ (y_d = 0) &= \frac{2}{d|B_1|} \frac{1}{|x-y|^d} + \frac{2x_d^2}{|B_1||x-y|^{d+2}} \\ \partial_{x_d}^2 K(x,y) &= \frac{-2}{|B_1|} \frac{(x_d - y_d)}{|x-y|^{d+2}} + \frac{4x_d}{|B_1||x-y|^{d+2}} - \frac{2(d+2)|B_1|}{x} \frac{(x_d - y_d)}{|x-y|^{d+4}} \end{split}$$

Then:

$$\begin{split} \Delta_x K(x,y) &= \sum_{i=1}^{d-1} \ \partial_{x_i}^2 K(x,y) + \partial_{x_i}^2 K(x,y) \\ &= -\frac{2x_d}{|B_1|} \left[\frac{d-1}{|x-y|^{d+2}} - (d+2) \sum_{i=1}^{d-1} \frac{(x_i - y_i)^2}{|x-y|^{d+4}} \right. \\ &+ \frac{1+2}{|x-y|^{d+2}} - \frac{(d+2)x_d(x_d - y_d)}{|x-y|^{d+4}} \right] \\ &= -\frac{2x_d}{|B_1|} \left[\frac{d+2}{|x-y|^{d+2}} - (d+2) \frac{1}{|x-y|^{d+4}} \left(\sum_{i=1}^{d} |x_i - y_i|^2 \right) \right] = 0 \end{split}$$

for all $x \in \mathbb{R}^d_+$, $y \in \partial \mathbb{R}^d_+$. Claim (exercise) for all $x \in \mathbb{R}^d_+$,

$$\int_{\partial \mathbb{R}^d_{\perp}} K(x, y) \, dy = 1$$

Consider

$$u(x) = \int_{\partial \mathbb{R}^d_{\perp}} K(x, y) g(y) \, dy, \quad x \in \mathbb{R}^d_{+}$$

Since $g \in L^{\infty}(\mathbb{R}^{d-1}) = L^{\infty}(\partial \mathbb{R}^d_+)$ and $K(x,y) \ge 0$, hence

$$|u(x)| \le \left(\int_{\partial \mathbb{R}^d_+} K(x,y) \, dy\right) \|g\|_{L^{\infty}}$$

Thus $||u||_{L^{\infty}} \leq ||g||_{L^{\infty}}$. Moreover

$$D_x^{\alpha} u(x) = \int_{\partial \mathbb{R}^d_{\perp}} D_x^{\alpha} K(x, y) g(y) \, dy$$

bounded, so $u \in C^{\infty}(\mathbb{R}^d_+)$, $x \mapsto K(x,y)$ is smooth as $x \neq y$.

$$\Delta_x u(x) = \int_{\partial \mathbb{R}^d_+} \underbrace{\Delta_x K(x, y)}_{=0} g(y) \, dy = 0$$

So u is harmonic in \mathbb{R}^d_+ . ($\Rightarrow u \in C^{\infty}$ by Weyl's lemma). Take $x_0 \in \partial \mathbb{R}^d_+$ and $x \in \mathbb{R}^d_+$. Then:

$$|u(x) - g(x_0)| = \left| \int_{\partial \mathbb{R}^d_+} K(x, y) (g(y) - g(x_0)) \, dy \right|$$

$$\leq \int_{\partial \mathbb{R}^d_+} K(x, y) |g(y) - g(x_0)| \, dy$$

$$= \underbrace{\int_{|y - x_0| \leq L|x - x_0|}}_{(I)} + \underbrace{\int_{|y - x_0| > L|x - x_0|}}_{(II)}$$

$$(I) = \int_{|y-x_0| \leqslant L|x-x_0|} K(x,y)|g(y) - g(x_0)| \, dy$$
$$= \sup_{|y-x_0| \leqslant L|x-x_0|} |g(y) - g(x_0)| \xrightarrow{x \to x_0} 0 \quad \forall L > 0$$

(II): If
$$|y - x_0| > L|x - x_0|$$
, then $|y - x| > \frac{1}{2}|y - x_0| > \frac{L}{2}|x - x_0|$ if $L \ge 2$.

$$\int_{|y-x_0|>|L|x-x_0|} K(x,y)|g(y)-g(x_0)|\,dy \leqslant C \int_{y\in\partial\mathbb{R}_+^d} \frac{x_d}{|x_0-y|}\,dy$$

$$Cx_d \int_{\substack{z\in\mathbb{R}^{d-1}\\|z|>L|x-x_0|}} \frac{1}{|z|^d}\,dz = const.\frac{x_d}{L|x-x_0|} \leqslant \frac{const.}{L} \xrightarrow{L\to\infty} 0$$

$$x_d = |x_d - (x_0)_d| \le |x - x_0|$$

5.2 Green's function for a ball

Let B = B(0,1). For all $x \in B$, for all $y \in \bar{B}$ we want to find the corrector function $\Phi_x(y)$ s.t.

$$\begin{cases} \Delta_y \Phi_x(y) = 0 & \text{in } B \\ \Phi_x(y) = G(y - x) & \text{on } \partial B \end{cases}$$

where for $d \ge 3$: $G(z) = \frac{1}{d(d-2)|B_1||z|^{d-2}}$.

Definition 5.10 (Reflection / Duality through the sphere ∂B) For all $x \in \mathbb{R}^d \setminus \{0\}$ we define $\tilde{x} = \frac{x}{|x|^2}$. Clearly we have for all $x \in B$ that if |x| < 1, then $|\tilde{x}| = \left|\frac{x}{|x|^2}\right| = \frac{1}{|x|} > 1$, so $\tilde{x} \notin \bar{B}$

Lemma 5.11 For $d \ge 3$ the function $\Phi_x(y) = G(|x|(y-\tilde{x}))$ is a corrector function.

Proof.

$$\Phi_x(y) = \frac{1}{d(d-2)|B_1||x|^{d-2}|y-\tilde{x}|^{d-2}}$$

for all $x \in B, x \neq 0$, for all $y \in \overline{B}$. Then clearly $y \mapsto \Phi_x(y)$ is harmonic in B (Since $\frac{1}{|z|^{d-2}}$ is harmonic in $\mathbb{R}\setminus 0$). Let's check the boundary: Let $y \in \partial B$, i.e. |y|=1. Then

$$\begin{aligned} ||x|(y-\tilde{x})| &= |x| \left| y - \frac{x}{|x|^2} \right| \\ &= |x| \sqrt{|y|^2 - 2\frac{xy}{|x|^2} + \left| \frac{x}{|x|^2} \right|^2} \\ &= \sqrt{|x|^2 |y|^2 - 2xy + 1} \\ (|y| &= 1) &= \sqrt{|x|^2 - 2xy + |y|^2} = |x - y| \end{aligned}$$

Thus $\Phi_x(y) = G(|x||y - \tilde{x}|) = G(y - x)$ for all $0 \neq x \in B$, for all $y \in \partial B$. Let's compute the Poisson kernel: If want to solve

$$\begin{cases} -\Delta u = 0 & \text{in } B \\ u = g & \text{on } \partial B \end{cases}$$

then

$$u(x) = -\int_{\partial B} \frac{\partial \tilde{G}}{\partial \vec{n}_y}(x, y) g(y) dS(y).$$

$$\tilde{G}(x,y) = G(y-x) - \Phi_x(y) = G(y-x) - G(|x|(y-\tilde{x})) \text{ for all } x \in B \setminus \{0\}, \ y \in \bar{B}.$$

$$\frac{\partial \tilde{G}}{\partial \vec{n}_y} = \sum_{i=1}^d \partial_{y_i} \tilde{G}y_i$$

Here

$$\begin{split} \partial_{y_i} \tilde{G} &= \partial_{y_i} G(y-x) - \partial_{y_i} [G(|x|(y-\tilde{x}))] \\ &= \frac{-(y_i-x_i)}{d|B_1||y-x|^d} + \frac{y_i-\tilde{x}_i}{d|B_1||x|^{d-2}|y-\tilde{x}|^d} \\ \Rightarrow \frac{\partial \tilde{G}}{\partial \vec{n}_y} &= \sum_{i=1}^d [\dots] y_i \\ &= \frac{-y(y-x))}{d|B_1||y-x|^d} + \frac{y(y-\tilde{x})}{d|B_1||x|^{d-2}|y-\tilde{x}|^d} \\ &= \frac{1}{d|B_1||y-x|^d} (-y(y-x) + y(y-\tilde{x})|x|^2) \\ &= \frac{1}{d|B_1||y-x|^d} [-|y|^2 + xy + |y|^2|x|^2 - xy] \\ &= \frac{-1+|x|^2}{d|B_1||y-x|^d} \end{split}$$

as $y \in \partial B$.

Theorem 5.12 (Poisson Formula for a Ball) Let $B=B(0,1), g\in C(\partial B)$. Define for all $x\in B$:

$$u(x) = \int_{\partial B} K(x, y)g(y) dS(y),$$

 $K(x,y) = -\frac{\partial \tilde{G}}{\partial \vec{n}_y}(x,y) = \frac{1-|x|^2}{d|B_1||y-x|^d} \text{ for all } x \in B, \text{ for all } y \in \partial B. \text{ Then } u \in C^\infty(B),$ $\Delta u = 0 \text{ and for all } x_0 \in \partial B \text{ we have } \lim_{x \in B} x_0 \ u(x) = g(x_0). \text{ This holds for all } d \geqslant 2.$

Proof. We need to check:

- 1. For all $y \in \partial B$, $x \mapsto K(x, y)$ is harmonic in B.
- 2. $\int_{\partial B} K(x,y) dS(y) = 1$ for all $x \in B$ (exercise)

Now for all $x \in B$, for all $y \in \partial B$:

$$K(x,y) = \frac{1 - |x|^2}{d|B_1||y - x|^d}$$

$$\partial_{x_i} K(x,y) = \frac{-2x_i}{d|B_1|} \frac{1}{|x - y|^d} - \frac{1 - |x|^2}{|B_1|} \frac{x_i - y_i}{|x - y|^{d+2}}$$

$$\partial_{x_i}^2 K(x,y) = -\frac{2}{d|B_1|} \frac{1}{|x - y|^d} + \frac{2x_i}{|B_1|} \frac{x_i - y_i}{|x - y|^{d+2}} + \frac{2x_i}{|B_1|} \frac{x_i - y_i}{|x - y|^{d+2}}$$

$$-\frac{1 - |x|^2}{|B_1|} \frac{1}{|x - y|^{d+2}} + \frac{1 + |x|^2}{|B_1|} (d+2) \frac{(x_i - y_i)^2}{|x - y|^{d+4}}$$

$$\Delta_x K = \sum_{i=1}^d \partial_{x_i}^2 K = -\frac{2}{|B_1|} \frac{1}{|x - y|^d} + \frac{4x(x - y)}{|B_1||x - y|^{d+2}}$$

$$-\frac{d(1 - |x|^2)}{|B_1|} \frac{1}{|x - y|^{d+2}} + (d+2) \frac{1 - |x|}{|B_1|} \frac{1}{|x - y|^{d+2}}$$

$$= \frac{2}{|B_1||x - y|^{d+2}} [-|x|^2 + 2xy - |y|^2 + 2|x|^2 - 2xy + 1 - |x|^2]$$

$$= \frac{2}{|B_1||x - y|^{d+2}} [-|x|^2 + 2xy - |y|^2 + 2|x|^2 - 2xy + 1 - |x|^2]$$

 $1-|y|^2=0$ as $y\in\partial B$. Thus $\Delta_x K(x,y)=0$, for all $x\in B$, for all $y\in\partial B$.

$$|u(x)| = \left| \int_{\partial B} K(x, y) g(y) \, dS(y) \right| \le ||g||_{L^{\infty}(\partial B)}$$

 $\int_{\partial B} K(x,y), dS(y) = ||g||_{L^{\infty}},$

$$\Delta_x u(x) = \int_{\partial B} \underbrace{\Delta_x K(x, y)}_{=0} g(y) \, dS(y) = 0$$

Take $x \in B$, $x \to x_0 \in \partial B$.

$$|u(x) - g(x_0)| = \left| \int_{\partial B} K(x, y) (g(y) - g(x_0)) \, dS(y) \right|$$

$$\leq \int_{A_1} + \int_{A_2} K(x, y) |g(y) - g(x_0)| \, dS(y),$$

where

$$A_1 = \{ y \in \partial B \mid |y - x_0| \leq |x - x_0|^{\alpha} \}$$

$$A_2 = \{ y \in \partial B \mid |y - x_0| > |x - x_0|^2 \}$$

On A_1 we have:

$$\int_{A_1} \dots \leqslant \sup_{\substack{|z-x_0| \leqslant |x-x_0|^{\alpha} \\ z \in \partial B}} \int_{\partial B} K(x,y) \, dS(y) \xrightarrow{x \to x_0} 0$$

since $G \in C(\partial B)$. On A_2 :

$$|y - x_0| > |x - x_0|^{\alpha}$$

$$\Rightarrow |y - x| \ge |y - x_0| - |x - x_0| \ge |x - x_0|^{\alpha} - |x - x_0| \ge \frac{1}{2}x - x_0^{\alpha}$$

if $\alpha < 1$ and $|x - x_0|$ small. So we get

$$K(x,y) = \frac{1 - |x|^2}{d|B_1||x - y|^d} \le C \frac{1 - |x|^2}{|x - x_0|^{d\alpha}} \le C|x - x_0|^{1 - d\alpha}$$

Thus

$$\int_{A_2} K(x,y) |g(y) - g(x_0)| \, dS(y) \leqslant C \|g\|_{L^{\infty}} |x - x_0|^{1 - d\alpha} \xrightarrow{x \to x_0} 0$$

if
$$1 - d\alpha > 0 \Leftrightarrow \alpha < \frac{1}{d}$$
.

Exercise 5.13 (E 7.2) Define $\mathbb{R}^d_+ = \{(x', x_d) \in \mathbb{R}^{d-1} \times \mathbb{R} \mid x_d > 0\}$. Let $K(x, y) = \frac{2x_d}{d|B_1||x-y|^d}$ for all $x \in \mathbb{R}^d_+, y \in \partial \mathbb{R}^d_+ = \{(y', 0) \mid y' \in \mathbb{R}^{d-1}\} \simeq \mathbb{R}^{d-1}$. Prove

$$\int_{\partial \mathbb{R}^d_+} K(x, y) \, dS(y) = 1 \quad \forall x \in \mathbb{R}^d_+$$

Solution. Denote $x = (x', x_d), y = (y', 0), x', y' \in \mathbb{R}^{d-1}, x_d > 0.$

$$\int_{\partial \mathbb{R}^d_+} K(x,y) \, dS(y) = \int_{\mathbb{R}^{d-1}} \frac{2x_d}{d|B_1| \left(|x'-y'|^2 + x_d^2\right)^{\frac{d}{2}}} \, dy' = \dots$$

as
$$|x - y| = |(x' - y', x_d)| = \sqrt{|x' - y'|^2 + x_d^2}$$
.

$$(y' - x' \mapsto y') \qquad \dots = \int_{\mathbb{R}^{d-1}} \frac{2x_d}{d|B_1| (|y'|^2 + x_d^2)^{\frac{d}{2}}} dy'$$

$$(y' = x_d z) \qquad = \int_{\mathbb{R}^{d-1}} \frac{2x_d}{d|B_1| (x_d^2(|z|^2 + 1))^{\frac{d}{2}}} (x_d^{d-1}) dz$$

$$= \int_{\mathbb{R}^{d-1}} \frac{2}{d|B_1| (|z|^2 + 1)^{\frac{d}{2}}} dz$$

$$= \int_0^\infty \frac{2\omega_{d-1}}{d|B_1|} \frac{1}{(r^2 + 1)^{\frac{d}{2}}} r^{d-2} dr$$

$$= \frac{2\omega_{d-1}}{\omega_d} \int_0^\infty \frac{1}{(r^2 + 1)^{\frac{d}{2}}} r^{d-2} dr$$

Set d = 2: $\omega_1 = 1, |\omega_2| = 2\pi$

$$\frac{2}{\pi} \int_0^\infty \frac{1}{r^2 + 1} dr = \frac{2}{\pi} \int_0^{\frac{\pi}{2}} \frac{1}{(\tan t)^2 + 1} [(\tan t)^2 + 1] dt = 1$$

we we set $r = \tan t, t \in (0, \frac{\pi}{2}), \frac{dr}{dt} = (\tan t)' = 1 + (\tan t)^2$

For d = 3:

$$\frac{2\cdot 2\pi}{4\pi} \int_0^\infty \frac{1}{(r^2+1)^{\frac{3}{2}}} r \, dr = \int_0^\infty \frac{d}{dr} \left[\frac{-1}{(r^2+1)^{\frac{1}{2}}} \right] dr = \frac{-1}{(r^2+1)^{\frac{1}{2}}} \bigg]_0^\infty = 1$$

Exercise 5.14 (7.3) Let $g \in C(\partial \mathbb{R}^d_+) \cap L^{\infty}(\partial \mathbb{R}^d_+)$ $(\partial \mathbb{R}^d_+ \simeq \mathbb{R}^{d-1})$.

$$u(x) = \int_{\partial \mathbb{R}^d_+} K(x, y) g(y) dS(y)$$
 $K(x, y) = \frac{2x_d}{d|B_1||x - y|^d}, x \in \mathbb{R}^d_+$

Prove that if g(y) = |y|, if $|y| \le 1$, then $|\nabla u|$ is unbounded in $B(0,r) \cap \mathbb{R}^d_+$ for all r > 0.

Solution.

$$\begin{split} \partial_{x_d} u(x) &= \int_{\partial \mathbb{R}^d_+} \partial x_d K(x,y) g(y) \, dy \quad \forall x \in \mathbb{R}^d_+ \\ &= \frac{2}{d|B_1|} \int_{\partial \mathbb{R}^d_+} \left[\frac{1}{|x-y|^d} - \frac{dx_d^2}{|x-y|^{d+2}} \right] g(y) \, dy \\ &= \frac{2}{d|B_1|} \int_{\partial \mathbb{R}^d_+} \frac{1}{|x-y|^{d+2}} [|x-y|^2 - dx_d^2] g(y) \, dy \\ &= \frac{2}{d|B_1|} \int_{\partial \mathbb{R}^d_+} \frac{1}{(|x'-y'| + x_d^2|^{\frac{d+2}{2}}} \left[|y'|^2 - (d-1)x_d^2 \right] g(y) \, dy \end{split}$$

Assume that $\partial_d u$ is bounded in $B(0,r) \cap \mathbb{R}^d_+$ Then:

$$|u(0, x_d) - \underbrace{u(0, 0)}_{q(0)=0}| \le C|x_d|$$

if x_d small. Consider:

$$\begin{split} \limsup_{x_d \to 0^+} \frac{u(0, x_d)}{x_d} &= \limsup_{x_d \to 0^+} c \int_{\mathbb{R}^{d-1}} \frac{1}{(|y'|^2 + x_d^2)^{\frac{d}{2}}} g(y) \, dy' \\ &\geqslant \int_{\mathbb{R}^{d-1}} \frac{1}{|y'|^d} g(y) \, dy = \int_{|y'| \leqslant 1} + \int_{|y'| > 1} \\ &\quad to \int_{\mathbb{R}^{d-1}} \frac{1}{|y'|^{d-1}} \, dy' = \infty \end{split}$$

Exercise 5.15 (Bonus 7) Recall the Poisson kernel on a ball $B(0,r) \subseteq \mathbb{R}^d$:

$$K(x,y) = \frac{r^2 - |x|^2}{d|B_1|r} \frac{1}{|x - y|^d}$$

for all $x \in B(0,r)$, $y \in \partial B(0,r)$. Prove:

$$\int_{\partial B(0,r)} K(x,y) \, dS(y) = 1$$

for all $x \in B(0,r)$. (It suffices if you can prove d=2 and d=3)

5.3 Energy Method

Consider $u \in C^2(\Omega)$ for $\Omega \subseteq \mathbb{R}^d$ open, bounded and with C^1 boundary and

$$\begin{cases} -\Delta u = f & \text{in } \Omega \\ u = g & \text{on } \partial \Omega. \end{cases}$$

Take $\phi \in C_c^{\infty}(\Omega)$, then by integration by parts:

$$0 = \int_{\Omega} (-\Delta u - f)\phi = \int_{\Omega} \nabla u \nabla \phi - \int_{\Omega} f \phi$$

Key observation: This is the derivative of the energy functional

$$E(u) = \frac{1}{2} \int_{\Omega} |\nabla u|^2 - \int_{\Omega} f u$$

If u is a minimizer of E, then it solves the equation $-\Delta u = f$ in Ω . The boundary condition u = g does not appear on E, but this is encoded in the set of *admissible* functions. (The set of candidates of solutions). For the classical solutions, we have

Theorem 5.16 (Dirichlet's principle) Let $\Omega \subseteq \mathbb{R}^d$ be open, bounded with C^1 -boundary. Let $f \in C(\bar{\Omega})$ and $g \in C(\partial B)$. Then the following statements are equivalent:

1.
$$u \in C^2(\bar{\Omega})$$
 solves
$$\begin{cases} -\Delta u = f & \text{in } \Omega \\ u = g & \text{on } \partial\Omega \end{cases}$$

2. u is a minimizer of the variational problem $E = \inf_{v \in A} E(v)$, where

$$E(v) = \frac{1}{2} \int_{\Omega} |\nabla v|^2 - \int_{\Omega} f v,$$

$$A = \{ v \in C^2(\bar{\Omega}) \mid v = g \text{ on } \partial\Omega \}.$$

Moreover there is at most a solution / minimizer (uniqueness).

Proof. The result holds even for complex-valued functions. Let us write the proof for real-valued functions.

1. \Rightarrow 2.: Let $u \in C^2(\bar{\Omega})$ be a solution of $\begin{cases} -\Delta u = f & \text{in } \Omega \\ u = g & \text{on } \partial \Omega \end{cases}$. Then we prove $E(u) \leqslant E(v)$ for all $v \in A$. If $v \in A$, then u - v = 0 on $\partial \Omega$. Using this and $-\Delta u = f$ in Ω , we have:

$$\begin{aligned} 0 &= \int_{\Omega} (-\Delta u - f) \cdot (u - v) \, dy \\ (\text{Part. Int.}) &= \int_{\Omega} \nabla u (\nabla u - \nabla v) \, dy - \int_{\Omega} f(u - v) \, dy \\ &= \left[\frac{1}{2} \int_{\Omega} |\nabla u|^2 \, dy - \int_{\Omega} fu \, dy \right] - \left[\frac{1}{2} \int_{\Omega} |\nabla v|^2 \, dy - \int_{\Omega} fv \, dy \right] \\ &+ \frac{1}{2} \int_{\Omega} |\nabla u|^2 + \frac{1}{2} \int_{\Omega} \\ &= E(u) - E(v) + \frac{1}{2} \underbrace{\int_{\Omega} |\nabla u - \nabla v|^2}_{\geqslant 0} \end{aligned}$$

 $E(u) \leq E(v)$, so u is a minimizer of $\inf_{v \in A} E(v)$. Moreover u is the unique minimizer on A. Since E(u) = E(v) we have $\int_{\Omega} |\nabla (u - v)|^2 = 0$, so u - v = const., so u - v = 0 in $\bar{\Omega}$.

2. \Rightarrow 1.: Assume that u is a minimizer of $\inf_{v \in A} E(v)$. Then $E(u) \leqslant E(v)$ for all $v \in A$. Take $\phi \in C_c^{\infty}(\Omega)$, then $u + t\phi \in A$ for all $t \in \mathbb{R}$.

$$\begin{split} &\Rightarrow E(u) \leqslant E(u+t\phi) \text{ for all } t \in \mathbb{R} \\ &\Rightarrow t \mapsto E(u+t\phi) \text{ has a minimizer at } t=0 \\ &\Rightarrow 0 = \frac{d}{dt} E(u+t\phi)|_{t=0} \\ &= \frac{d}{dt} \left(\frac{1}{2} \int_{\Omega} |\nabla u + t \nabla \phi|^2 - \int_{\Omega} f(u+t\phi) \right) \bigg|_{t=0} \\ &= \frac{d}{dt} \left(\frac{1}{2} \int_{\Omega} |\nabla u|^2 + t^2 |\nabla \phi|^2 + 2t \nabla u \nabla \phi - \int_{\Omega} f(u+t\phi) \right) \bigg|_{t=0} \\ &\int_{\Omega} \nabla u \nabla \phi - \int_{\Omega} f \phi = \int_{\Omega} (-\Delta u - f) \phi \\ &\text{ for all } \phi \in C_c^{\infty}(\Omega). \text{ So } -\Delta u - f = 0 \text{ in } \Omega \text{ and } u = g \text{ since } u \in A. \end{split}$$

Direct method of calculus of variations. Think $f : \mathbb{R} \to \mathbb{R}$, $f \in C(\mathbb{R})$, $f(x) \to \infty$ as $|x| \to \infty$. There is a $x_0 \in \mathbb{R}$ s.t. $f(x_0) = \inf_{x \in \mathbb{R}} f(x)$.

Step 1: $E = \inf_{x \in \mathbb{R}} f(x) > -\infty$

Step 2: Take a minizing sequence $\{x_n\} \subseteq \mathbb{R}$, $f(x_n) \to E$. Up to a subsequence $x_n \to x_0$ in \mathbb{R} (compactness)

Step 3: Lower semicontinuity $E = \liminf_{n \to \infty} f(x_n) \ge f(x_0)$

If we apply the direct method to $\inf_{v \in A} E(v)$,

$$E(v) = \frac{1}{2} \int_{\Omega} |\nabla v|^2 - \int_{\Omega} f v,$$

 $A = \{ v \in C^2(\bar{\Omega}), v = g \text{ on } \partial\Omega \}$

Step 1: Easy $E = \int_{v \in A} E(v) > -\infty$

Step 2: There is a minimizing sequence $\{v_n\} \subseteq A$ s.t. $E(v_n) \to E$. We don't know if there is a subsequence of $\{v_n\}$ that converges to $u \in A$. The lack of compactness is a serious problem! We need to find the rigt set A! Consider again

$$\begin{cases} -\Delta u = f & \text{in } \Omega \\ u = g & \text{on } \partial \Omega \end{cases}$$

Consider the simple case g=0. $\Delta u-f$ in $\Omega\Leftrightarrow \nabla u\nabla\phi...$ The right set A should be $A=\{v\mid \int_{\Omega}|\nabla v|^2<\infty, v=0 \text{ on }\partial\Omega\}.$ Rigorously we take $W_0^{1,2}(\Omega)=\overline{C_c^\infty(\Omega)}W^{1,2}(\Omega)$ (Notation: $H_0^1=W_0^{1,2},H^1=W^{1,2}$) Recall that $W^{1,p}$ is a banach space with norm $\|f\|_{W^{1,p}(\Omega)}=\|f\|_{L^p(\Omega)}+\|\nabla f\|_{L^p(\Omega)}.$ We know that $C_c^\infty(\Omega)$ is dense in $W_{loc}^{1,p}(\Omega)$, i.e. for all $u\in W_{loc}^{1,p}(\Omega)$ there is $\|u_n\|\subseteq C_c^\infty$ s.t. $u_n\to u$ in $W^{1,p}(K)$ for all $K\subseteq\Omega$ compact. However in general $C_c^\infty(\Omega)$ is not dense in $W^{1,p}(\Omega)$, i.e. $W_0^{1,p}(\Omega)=\overline{C_c^\infty(\Omega)}W^{1,p}(\Omega)\subsetneq W^{1,p}(\Omega).$ Clearly $W_0^{1,p}$ is a closed subspace of $W^{1,p}(\Omega)\to W_0^{1,p}(\Omega)$ is a Banach space with $\|\cdot\|_{W^{1,p}(\Omega)}.$ Why does $W_0^{1,p}(\Omega)$ encode the 0-boundary condition? Note that by definition for all $u\in W_0^{1,p}(\Omega)$ there is a sequence $\{u_n\}\subseteq C_c^\infty(\Omega), u_n\to u$ in $W^{1,p}(\Omega)$ up to a subsequence $u_n(x)\to u(x)$ for almost every $x\in\Omega.$ Note $u_n|_{\partial\Omega}=0\to u|_{\partial\Omega}=0$ since $\partial\Omega$ must be of 0-measure.

Theorem 5.17 (Characterization for $W_0^{1,p}$) Let Ω be open, bounded with C^1 -boundary. Let $u \in W^{1,p}(\Omega) \cap C(\bar{\Omega})$. Then the following statements are equivalent:

- a) u = 0 on $\partial \Omega$
- b) $u \in W_0^{1,p}(\Omega)$

(Later we will remove the condition $C(\bar{\Omega})$ by introducing the Trace operator.)

Remark 5.18 If d=1, it holds that $W^{1,p} \subseteq C(\bar{\Omega})$. Then the theorem gives a full characterization for $W_0^{1,p}$, but if $d \ge 2$, then in general $W^{1,p} \nsubseteq C(\Omega)$. (later)

Proof of theorem 5.17.

 $a) \Rightarrow b$:

Lemma 5.19 If $u \in W^{1,p}(\Omega)$ and supp $u \subseteq \Omega$, then $u \in W_0^{1,p}(\Omega)$.

Proof. Since $K := \operatorname{supp} u$ is a compact subset in Ω , we can find a function $\chi \in C_c^{\infty}(\Omega)$, $\chi = 1$ on K. Moreover since $u \in W^{1,p}(\Omega)$, there is a sequence $\{u_n\} \subseteq C_c^{\infty}(\Omega)$ s.t. $u_n \to u$ in $W_{loc}^{1,p}(\Omega)$. We claim that $\chi u_n \to \chi u$ in $W_{loc}^{1,p}(\Omega)$. (exercise, $\nabla(\chi u) = \nabla \chi u + \chi \nabla u$). This implies $\chi u_n \to u$ in $W^{1,p}(\operatorname{supp} \chi)$, thus $\chi u_n \to u$ in $W^{1,p}(\Omega)$, so $u \in W_0^{1,p}(\Omega)$.

Assume $u \in W^{1,p}(\Omega) \cap C(\bar{\Omega})$ and u = 0 on $\partial\Omega$. Take $G \in C^1(\mathbb{R})$ s.t. $|G(t)| \leq t$ for all t, G(t) = t if $t \geq 2$ and G(t) = 0 if $t \leq 1$. Then let

$$u_n(x) := \frac{1}{n} G(nu(x)) \in W^{1,p}(\Omega)$$

$$\stackrel{\text{(Chain-rule)}}{\Rightarrow} \quad \nabla u_n(x) = \frac{1}{n} G'(nu(x)) n \nabla u(x) = G'(nu(x)) \ \nabla u(x)$$

Moreover, u_n is compactly supported in Ω , so $u_n \in W_0^{1,p}(\Omega)$ by the lemma and $u_n \to u$ in $W^{1,p}(\Omega)$, so $u \in W_0^{1,p}(\Omega)$ since $W_0^{1,p}$ is a closed space. Recall that $u \in C(\bar{\Omega})$ and u = 0 on $\partial \Omega$. Thus for all $\epsilon > 0$ there is a compact $K_\epsilon \subseteq \Omega$ s.t. $\sup_{x \in \Omega \setminus K_\epsilon} |u(x)| \le \epsilon$. For any given $n \in \mathbb{N}$, $u_n(x) \neq 0$, so $G(nu(x)) \neq 0$. This implies n|u(x)| > 1, hence $|u(x)| > \frac{1}{n}$. Thus $u_n(x) = 0$ for all x such that $|u(x)| \le \frac{1}{n}$, so $\sup u_n \subseteq K_{\frac{1}{n}}$ compact in Ω . Next, let us check $u_n \to u$ in $W^{1,p}(\Omega)$.

$$\int_{\Omega} |u_n(x) - u(x)|^p \, dx \to 0$$

since $u_n(x) = \frac{1}{n}G(nu(x)) \xrightarrow{n \to \infty} u(x)$ for all $x \in \Omega$ and $|u_n(x)| \leq \frac{1}{n}|G(nu(x))| \leq \frac{1}{n}|nu(x)| \leq |u(x)| \in L^p(\Omega)$.

$$\int_{\Omega} |\nabla u_n(x) - \nabla u(x)|^p dx = \int_{\Omega} |G'(nu(x)) - 1|^p |\nabla u(x)|^p dx \to 0$$

as $|G'(v(x)) - 1| \to 0$ for all x s.t. $u(x) \neq 0$ and $\nabla u(x) = 0$ on $\{x \mid u(x) = 0\}$. (exercise)

(b) \Rightarrow (a): Let $u \in W^{1,p}(\Omega) \cap C(\bar{\Omega})$ and $u \in W_0^{1,p}(\Omega)$. Then we prove u = 0 on $\partial\Omega$. Lets regard the case $\Omega = Q_+ = \{(x',x_d) \mid \mathbb{R}^{d-1} \times \mathbb{R} \mid |x'| < 1, 0 < x_d < 1\}$. We prove that if $u \in W_0^{1,p}(Q_+) \cap C(Q_+)$, then u = 0 on $Q_0 = \{(x',0) \mid x' \in \mathbb{R}^{d-1}, |x'| < 1\}$. Since $u \in W_0^{1,p}(Q_+)$ there is $\{u_n\} \subseteq C_c^{\infty}(Q_+)$ s.t. $u_n \to u$ in $W^{1,p}(Q_+)$ for all $x = (x',x_d) \in Q_+$, then:

$$u_n(x', x_d) = \underbrace{u_n(x', 0)}_{=0} + \int_0^{x_d} \partial_d u_n(x', t) dt$$

Hence

$$|u_n(x',x_d)| \le \int_0^{x_d} |\partial_d u_n(x',t)| dt$$

This implies:

$$\int_{0 < x_d < \epsilon} \int_{|x'| \le 1} |u_n(x', x_d)| dx' dx_d$$

$$\leq \int_{0 < x_d < \epsilon} \int_{|x'| < 1} \left(\int_0^{x_d} |\partial_d u_n(x', t)| dt \right) dx' dx_d$$

$$\leq \epsilon \int_{|x'| < 1} \int_0^{\epsilon} |\partial_d u_n(x', t)| dx' dt$$

$$\Rightarrow \frac{1}{\epsilon} \int_0^{\epsilon} \int_{|x'| \le 1} |u_n(x', x_d)| \ dx' dx_d \le \int_0^{\epsilon} \int_{|x'| < 1} |\partial_d u_n(x', x_d)| \ dx' dx_d$$

for all $n \in \mathbb{N}$, $\epsilon > 0$. Take now $n \to \infty$, use $u_n \to u$ in $W^{1,p}(\Omega)$. Then:

$$\frac{1}{\epsilon} \int_0^{\epsilon} \int_{|x'| \leq 1} |u(x', x_d)| \, dx' \, dx_d \leq \int_0^{\epsilon} \int_{|x'| < 1} |\partial_x u_n(x', x_d)| \, dx' \, dy$$

for all $\epsilon > 0$. Take $\epsilon \to 0$:

$$\int_{|x'| \le 1} |u(x', 0)| \, dx' \le 0$$

here we use $u \in C(\bar{\Omega})$ for the left side and Dominated Convergence for the right side. Thus u(x',0)=0 for all $|x'|\leqslant 1$, i.e. u=0 on $\partial\Omega$. Let's regard the general case: Let Ω be open, bounded and with C^1 -boundary. Lets define local charts By definition for all $x\in\partial\Omega$, there is a U_x open, such there is a bijective map $h:U_x\to Q$, and h,h^{-1} are C^1 . Then clearly $\partial\Omega\subseteq\bigcup_{x\in\partial\Omega}U_x$. Since $\partial\Omega$ is compact, there is a finite subcover $\{U_i\}_{i=1}^N$ s.t. $\partial\Omega\subseteq\bigcup_{i=1}^NU_i$. We can find U_0 open s.t. $\bar{U}_0\subseteq\Omega$ and $\Omega\subseteq\bigcup_{i=0}^NU_i$.

Lemma 5.20 There is a sequence $\{x_i\}_{i=0}^N \subseteq C^{\infty}(\mathbb{R}^d)$ s.t.

- 1. $\chi_i \ge 0$, $\sum_{i=0}^{N} \chi_i = 1$ in \mathbb{R}^d ($\{\chi_i\}$ is a partition of unity)
- 2. For all i = 1, ..., N, supp χ_i is in U_i , i.e. $\chi_i \in C_c^{\infty}(U_i)$.
- 3. i = 0, supp $\chi_0 \subseteq \mathbb{R}^d \setminus \partial \Omega$ and $\chi_0 \setminus \Omega \in C_c^{\infty}(\Omega)$. (exercise)

Given $u \in W_0^{1,p}(\Omega) \cap C(\bar{\Omega})$. Then $u = \sum_{i=0}^N \chi_i u$, where $\chi_i \geqslant 0$, $\chi_0 \in C_c^{\infty}(\Omega)$, $\chi_i \in C_c^{\infty}(U_i)$. Since $\chi_0 u$ is supported in a copact set inside Ω , $\chi_0 u = 0$ on $\partial \Omega$. It remains to show that for all $i = 1, \dots, N$, $\chi_i u = 0$ on $U_i \cap \partial \Omega$. Then $\chi_i u(h^{-1}x) \in W_0^{1,p}(Q) \cap C(\bar{\Omega})$. This implies $\chi_i u(h^{-1}x) = 0$ on Q_0 , so $\chi_i u(x) = 0$ on $U_i \cap \partial \Omega$. Why $W_0^{1,p}(U_i \cap \Omega) \to W_0^{1,p}(Q_+)$. If $v \in W_0^{1,p}(U_i \cap \Omega)$, then $v_n \to v, v_n \in C_c^{\infty}$. $v_n \circ h^{-1} \to v \circ h^{-1} \Rightarrow v \circ h^{-1} \in W_0^{1,p}(Q_+)$

Exercise 5.21 (E 8.1) Let $u \in W^{1,1}_{loc}(\mathbb{R}^d)$. Let $B = u^{-1}(\{0\})$. Prove that $\nabla u(x) = 0$ for a.e. $x \in B$.

Solution. We have already seen that if $f,g\in W^{1,1}_{loc}(\mathbb{R}^d)$, then $\max(f,g)\in W^{1,1}_{loc}$. This implies that if $u=u^+-u^-\in W^{1,1}_{loc}$, then $u^+,v^+\in W^{1,1}_{loc}$ since $u^+=\max(u,0)$ and $u^-=\max(-u,0)$. We have that $\nabla u=\nabla u^+-\nabla u^-$. Claim:

$$\nabla u^{+} = \begin{cases} 0 & u(x) \le 0 \\ \nabla u & u(x) > 0 \end{cases} \quad \nabla u^{-} = \begin{cases} 0 & u(x) \ge 0 \\ \nabla u & u(x) < 0 \end{cases}$$

$$\int_{\mathbb{R}^d} (\partial_i u^+) \phi = -\int_{\mathbb{R}^d} u^+ \partial_i \phi = -\int_{\{u(x) \le 0\}} 0 \partial_i \phi - \int_{\{u(x) > 0\}} u \partial_i \phi$$
$$= \int_{\{u(x) \le 0\}} 0 \phi + \int_{\{u(x) > 0\}} \partial_i u \phi$$

Alternative way: We showed for $f \in W^{1,p}(\mathbb{R}^d)$, that

$$\nabla |f|(x) = \begin{cases} (\nabla f)(x) & f(x) > 0 \\ -(\nabla f)(x) & f(x) < 0 \\ 0 & f(x) = 0 \end{cases}$$

 $u_+ = \frac{1}{2}(u+|u|)$. Hence $\nabla u_+ = \frac{1}{2}(\nabla u + \nabla |u|)$. Remark: If $A \subseteq \mathbb{R}$ has measure zero, then $\nabla u 1_{\{u(x) \in A\}} = 0$ a.e. (Th. 6.19 Lieb-Loss Analysis)

Exercise 5.22 (E 8.2) Let $\Omega, U \subseteq \mathbb{R}^d$ be open, $U \cap \Omega \neq \emptyset$, $u \in W_0^{1,p}(\Omega)$, $1 \leq p < \infty$, $\chi \in C_c^{\infty}(U)$. Prove: $\chi u \in W_0^{1,p}(\Omega \cap U)$ Hint: Recall $W_0^{1,p}(\Omega) = \overline{C_c^{\infty}(\Omega)}^{\|\cdot\|_{W^{1,p}}}$

Solution. By definition there is a sequence $(u_n)_{n\in\mathbb{N}}\subseteq C_c^\infty(\Omega)$ s.t. $u_n\xrightarrow[n\to\infty]{\|\cdot\|_{W^{1,p}}}u$, i.e.

$$||u_n - u||_p + ||\nabla u_n - \nabla u||_p \xrightarrow{n \to \infty} 0$$

Define $f_n: \mathbb{R}^d \to \mathbb{C}$, $f_n(x) := u_n(x)\chi(x)$. Note $f_n \in C_c^{\infty}(\Omega \cap U)$ for all $n \in \mathbb{N}$. Claim: $(f_n)_{n \in \mathbb{N}}$ is Cauchy with respect to $\|\cdot\|_{W^{1,p}}$. Proof:

$$||f_n - f_m||_p = ||\chi(u_n - u_m)||_p \le ||\chi||_{\infty} \underbrace{||u_n - u_m||_p}_{n,m\to\infty} \xrightarrow{n,m\to\infty} 0$$

$$\nabla f_n = \nabla(\chi u_n) = (\nabla \chi)u_n + \chi \nabla u_n$$

$$\|\nabla f_n - \nabla f_m\|_p \le \|\nabla \chi(u_n - u_m)\|_p + \|\chi(\nabla u_n - \nabla u_m)\|_p$$

$$\le \|\nabla \chi\|_{\infty} \underbrace{\|u_n - u_m\|_p}_{n, m \to \infty} + \underbrace{\|\chi\|}_{<\infty} \underbrace{\|\nabla u_n - \nabla u_m\|_p}_{n, m \to \infty} \xrightarrow{n, m \to \infty} 0$$

Thus, there is a $f \in W_0^{1,p}(\Omega \cap U)$ s.t. $||f_n - f||_{W^{1,p}} \xrightarrow{n \to \infty} 0$. We know:

$$||f_n - \chi u||_{L^p} = ||\chi u_n - \chi u||_p$$

$$\leq ||\chi||_{\infty} \underbrace{||u_n - u||_p}_{\to 0} \xrightarrow{n \to \infty} 0$$

Since limits in L^p are unique, we get $\chi u = f \in W_0^{1,p}(\Omega \cup U)$.

Exercise 5.23 (E 8.3) Let $\Omega, U \subseteq \mathbb{R}^d$ open and bounded, $h: \bar{U} \to \bar{\Omega}$ C^1 -diffeomorphisms, $u \in W_0^{1,p}(\Omega), \ 1 \leqslant p < \infty$. Prove $(x \mapsto u(h(x)) \in W_0^{1,p}(U)$.

Solution. Since $u \in W_0^{1,p}(\Omega)$ there is a sequence $(u_n)_{n \in \mathbb{N}} \subseteq C_c^{\infty}(\Omega)$ s.t.

$$\|u - u_n\|_p + \|\nabla u - \nabla u_n\|_p \xrightarrow{n \to \infty} 0$$

Define for all $n \in \mathbb{N}$ $f_n : U \to \mathbb{C}$, $f_n(x) = u_n(h(x))$. Note $f_n \in C_c^1(U)$. Claim 1: $(f_n)_{n \in \mathbb{N}}$ is Cauchy wrt. $\|\cdot\|_{W^{1,p}}$.

$$||f_n - f_m||_p^p = \int_U |u_n(h(x)) - u_m(h(x))|^p dx$$

$$= \int_\Omega |u_n(y) - u_m(y)|^p dy \underbrace{\det(Dh^{-1})(y)}_{\leqslant C < \infty} \xrightarrow{n, m \to \infty} 0$$

$$(\nabla f_n)(x) = \nabla (u_n(h(x))) = (\nabla u_n)(h(x))(Dh)(x)$$

$$\begin{split} \|\nabla f_n - \nabla f_m\|_p^p &= \int_U \left| \left[(\nabla u_n)(h(x)) - (\nabla u_m)(h(x)) \right] \underbrace{(Dh)(x)}_{bdd.} \right|^p \, dx \\ &\leqslant C \int_U \left| (\nabla u_n)(h(x)) - (\nabla u_m)(h(x)) \right|^p \, dx \\ &= C \int_\Omega \left| (\nabla u_n)(y) - (\nabla u_m)(y) \right|^p \underbrace{\left| \det Dh^{-1}(y) \right|}_{\leqslant \tilde{C}} \, dx \xrightarrow{n,m \to 0} 0 \end{split}$$

Claim 2: $||f_n - u \circ h||_p \xrightarrow{n \to \infty} 0$.

$$||f_n - u \circ h||_p = \int_U |u_n(h(x)) - u(h(x))|^p dx$$

$$= \int_\Omega |u_n(y) - u(y)|^p \underbrace{\det Dh^{-1}(y)}_{\leqslant C} dy \xrightarrow{n \to \infty} 0$$

Conclusion: Since $(f_n)_{n\in\mathbb{N}}\subseteq C^1_c(U)$ is Cauchy with respect to $\|\cdot\|_{W^{1,p}}$, there is a $f\in W^{1,p}_0(U)$ s.t. $f_n\xrightarrow[\|\cdot\|_{W^{1,p}}]{}f$. Since limits in L^p are unique by claim 2 we get $u\circ h=f\in W^{1,p}_0(U)$.

Exercise 5.24 (E 8.4) Let $\Gamma \subseteq \mathbb{R}^d$ be compact, $\{U_i\}_{i=1}^N$ open s.t. $\Gamma \subseteq \bigcup_{i=1}^N U_i$. Prove: There exists $\{\chi_i\}_{i=0}^N \subseteq C^{\infty}(\mathbb{R}^d)$ s.t.

- 1. $\chi_i \geqslant 0$ for all $i, \sum_{i=0}^N \chi_i = 1$
- 2. $\operatorname{supp}(\chi_i) \subseteq U_i \text{ for all } i \in \{1, \dots, N\}$
- 3. $\operatorname{supp}(\chi_0) \subseteq \mathbb{R}^d \backslash \Gamma$

Solution. WLOG assume that $U_i \neq \emptyset$ for all i. If $\Gamma \neq 0$, then $\chi_0 = 1$ does the job. Now suppose $\Gamma \neq \emptyset$. Let $\psi \in C_c^{\infty}(B_1(0)), \psi \geqslant 0, \int \psi = 1, \psi|_{B_{\frac{1}{2}}(0)} > 0$ and for $\epsilon > 0$ let $\psi_{\epsilon}(x) = \frac{1}{\epsilon^d} \psi\left(\frac{x}{\epsilon}\right)$, so $\int \psi_{\epsilon} = 1$. Define

$$\tilde{d} := \sup\{\tilde{\tilde{d}} > 0 \mid \forall x \in \Gamma \exists i \in \{1, \dots, N\} \text{ s.t. } \operatorname{dist}(x, U_i^c) \geqslant \tilde{\tilde{d}}\}$$

Claim 1: $\tilde{d} > 0$ Suppose this was not true. Then there is a sequence $(x_n)_{n \in \mathbb{N}} \subseteq \Gamma$ s.t. for all $i \in \{1, ..., N\}$,

$$\operatorname{dist}(x_n, U_i^c) < \frac{1}{n}$$

Since Γ is compact, there is a subsequence, which we call x_n again, s.t. $x_n \xrightarrow{n \to \infty} \bar{x}$ for asome Γ . By $\Gamma \subseteq \bigcup_{i=1}^N U_i$ there is a $\epsilon_{\bar{x}} > 0$ and $i \in \{1, \dots, N\}$ s.t. $B_{\epsilon_{\bar{x}}}(\bar{x}) \subseteq U_i \notin \mathbb{R}$. Define $d := \min\{\tilde{d}, 1\} > 0$. For all $\epsilon > 0$, for all $A \subseteq \mathbb{R}^d$: $(A)_{\epsilon} := \{x \in A \mid \operatorname{dist}(x, A^c) \ge \epsilon\}$. for every $i \in \{1, \dots, N\}$ define $\phi_i : U_i \to [0, \infty)$ by

$$\phi_i(x) := \mathbb{1}_{(U_i \cap B_R(0))_{\frac{d}{4}}} \star \phi_{\frac{d}{4}}$$

Note $\phi_i \in C_c^{\infty}(U_i)$ and $(U_i \cap B_R(0))_{\frac{d}{4}} \subseteq (\operatorname{supp}(\phi_i))^0$. Define $\phi_0 : \mathbb{R}^d \setminus \Gamma \to [0, \infty)$ by $\phi_0(x) = \mathbb{1}_{(\mathbb{R}^1 \setminus \Gamma)_{\frac{d}{4}}} \star \psi_{\frac{d}{4}}$. Again, $\phi_0 \in C^{\infty}(\mathbb{R}^d \setminus \Gamma)$, $\operatorname{supp}(\phi_0))^0 \supseteq (\mathbb{R}^d \setminus \Gamma)_{\frac{d}{4}}$, $\operatorname{supp}(\phi_0) \subseteq \mathbb{R}^d \setminus \Gamma$. Claim 2: For all $x \in \mathbb{R}^d$ there is a $i \in \{0, 1, \dots, N\} : \phi_i(x) > 0$. Proof: By construction, we know for $i \in \{1, \dots, N\}$ that ϕ_i is > 0 on $(U_i \cap B_R(0))_{\frac{d}{4}}$. Moreover $\phi_0 > 0$ on $(\mathbb{R}^d \setminus \Gamma)_{\frac{d}{4}}$. thus, we are done if we can show that $\bigcup_{i=1}^N (U_i \cap B_R(0))_{\frac{d}{4}} = 0$.

 $(\mathbb{R}^d \setminus \Gamma)_{\frac{d}{4}} = \mathbb{R}^d$. Suppose there is a $x \in \mathbb{R}^d \setminus A$. Then $\operatorname{dist}(x, \Gamma) < \frac{d}{4}$. Since $\Gamma \subseteq B_{\frac{R}{2}}(0)$ and R > 2 and $d \leq 1$.

$$|x-0| \le \operatorname{dist}(x,\Gamma) + \frac{R}{2} < \frac{d}{4} + \frac{R}{2} = R - \frac{d}{4} - \frac{R}{2} + \frac{d}{2} < R - \frac{d}{4} - \frac{2}{2} + \frac{1}{2} < R - \frac{d}{4}$$

Thus $x \in (B_R(c))_{\frac{d}{4}}$. Thus, we are done if we can show that $x \in (U_i)_{\frac{d}{4}}$ for some $i \in \{1, ..., N\}$. Since $\operatorname{dist}(x, \Gamma) < \frac{d}{4}$, there is a $y \in \Gamma$ s.t. $|x - y| < \frac{d}{4}$. By definition of \tilde{d} there is a $i \in \{1, ..., N\}$ s.t. $\operatorname{dist}(y, U_i^c) \geqslant \tilde{d} \geqslant d$, i.e. for all $z \in U_i^c$ we have $|y - z| \geqslant d$. We get

$$|x-z| \geqslant |\underbrace{|x-y|}_{\leq \frac{d}{4}} - \underbrace{|y-z|}_{\geqslant d}| \geqslant \frac{3d}{4} < \frac{d}{4}$$

This implies $\operatorname{dist}(x,U_i^c) > \frac{d}{4}$, so $x \in (U_i)_{\frac{d}{4}} \notin$. Define for all $i \in \{0,\ldots,N\} : \chi_i : \mathbb{R}^d \to [0,\infty)$ by

$$\chi_i(x) = \frac{\phi_i(x)}{\sum_{j=0}^N \phi_j(x)}$$

 χ_i is well-defined by Claim 2 and $\chi_i \in C^{\infty}(\mathbb{R}^d)$. Also note that $\sum \chi_i = 1$, $\chi_i \geq 0$, which implies 1. Furthermore, since $\operatorname{supp}(\phi_i) \subseteq U_i$, we have $\operatorname{supp}(\chi_i) \subseteq U_i$ for all $i \in \{1, \ldots, N\}$, which implies 2. Finally, since $\operatorname{supp}(\phi_0) \subseteq \mathbb{R}^d \setminus \Gamma$, we get $\operatorname{supp}(\chi_0) \subseteq \mathbb{R}^d \setminus \Gamma$. This implies 3.

5.4 Variational problem for weak solutions

$$\begin{cases} -\Delta u = f & \text{in } \Omega \\ u = 0 & \text{on } \partial \Omega \end{cases}$$

("formally") for all $\phi \in C_c^{\infty}(\Omega)$, then

$$\int_{\Omega} \nabla u \nabla \phi = \int_{\Omega} f \phi$$

if $\nabla u \in L^2$, $f \in L^2$. By a density argument:

$$\int_{\Omega} \nabla u \nabla \phi = \int_{\Omega} f \phi$$

for all $\phi \in \overline{C_c^{\infty}(\Omega)}^{H^1(\Omega)} = H_0^1(\Omega)$.

Theorem 5.25 (Poincare inequality) There is a C > 0 s.t.

$$C \int_{\Omega} |\nabla v|^2 \geqslant \int_{\Omega} |v|^2$$

for all $v \in H_0^1(\Omega)$.

Remark 5.26 $H^1(\Omega)$ with $\|v\|_{H^1(\Omega)} = (\|v\|_{L^2}^2 + \|\nabla v\|_{L^2}^2)^{\frac{1}{2}}$ is a Hilbert-Space. This implies that $H^1_0(\Omega) \subseteq H^1(\Omega)$ is also a Hilbert space. By the Poincaré inequality (5.25) we have for all $v \in H^1_0(\Omega)$:

$$||v||_{H^1(\Omega)} \geqslant ||\nabla v||_{L^2} \geqslant \frac{1}{2c} ||v||_{L^2} + \frac{1}{2} ||\nabla v||_{L^2} \geqslant \frac{1}{C^1} ||v||_{H^1(\Omega)}$$

We can think of $H_0^1(\Omega)$ as a Hilbert space with $||v||_{H_0^1(\Omega)} := ||\nabla v||_{L^2(\Omega)}$.

Proof. (Of the Poincaré inequality (5.25)) We need to prove:

$$\exists C > 0: \quad C \int_{\Omega} |\nabla v|^2 \geqslant \int_{\Omega} |v|^2 \quad \forall v \in H_0^1(\Omega)$$

$$\Leftrightarrow \quad \exists C > 0: \quad C \int_{\Omega} |\nabla v|^2 \geqslant \int_{\Omega} |v|^2 \quad \forall v \in C_c^{\infty}(\Omega)$$

Assume by contradiction that this does not hold, i.e. there is no C>0 s.t. the statement holds. Thus there is a sequence $\{v_n\}\subseteq C_c^\infty(\Omega)$ s.t.

$$\int_{\Omega} |v_n|^2 = 1, \quad \int_{\Omega} |\nabla v_n|^2 \xrightarrow{n \to \infty} 0$$

Since $v_n \in C_c^2(\Omega)$ we can extend v_n by 0 outside Ω , so $v_n \in C_c^{\infty}(\mathbb{R}^d)$. Then:

$$\int_{\mathbb{R}^d} |v_n|^2 = 1, \quad \int_{\mathbb{R}^d} |\nabla v_n|^2 \to 0, \quad \text{supp } v_n \subseteq \Omega$$

By the Fourier transform:

$$\int_{\mathbb{R}^d} |\hat{v}_n(k)|^2 \, dk = 1, \quad \int_{\mathbb{R}^d} |2\pi k|^2 |\hat{v}_n(k)|^2 \, dk \to 0, \quad \text{supp } v_n \subseteq \Omega$$

We prove that

$$\int_{\mathbb{R}^d} |\hat{v}_n(k)|^2 \, dk \to 0$$

We write

$$\int_{\mathbb{R}^d} |\hat{v}_n(k)| \, dk = \int_{|k| \leqslant \epsilon} + \int_{|k| > \epsilon}$$

First, for all $\epsilon > 0$:

$$\int_{|k|>\epsilon} |\hat{v}_n(k)|^2 \leqslant \int_{\mathbb{R}^d} \frac{|k|^2}{\epsilon^2} |\hat{v}_n(k)|^2 dk \xrightarrow{n\to\infty} 0$$

Second:

$$\int_{|k| \leqslant \epsilon} |\hat{v}_n(k)|^2 dk \leqslant \left(\int_{|k| \leqslant \epsilon} 1 dk \right)^{\frac{1}{q}} \left(\int_{|k| \leqslant \epsilon} |\hat{v}_n(k)|^{2p} dk \right)^{\frac{1}{p}}, \quad 1 < p, q < \infty$$

$$\leqslant C \epsilon^{\frac{d}{q}} \|\hat{v}_n\|_{L^{2p}}^2, \quad \frac{1}{p} + \frac{1}{q} = 1 \text{ and } 1 \leqslant r \leqslant 2$$

Moreover, since Ω is bounded,

$$\|v_n\|_{L^r} \leqslant \left(\int_{\Omega} |v_n|^r\right)^{\frac{1}{r}} \leqslant \|1_{\Omega}\|_{L^s} \|v_n\|_{L^2}^{1-\theta} \leqslant C_{\Omega} \quad \forall 1 \leqslant r \leqslant 2.$$

Thus we can take r < 1 but close to 1. Then p is sufficiently large, so q is close to 1. Then

$$\int_{|k| \leq \epsilon} |\hat{v}_n(k)|^2 \leq C \epsilon^{\frac{d}{q}} \|\hat{v}_n\|_{L^{2p}}^2 \leq C \epsilon^{\frac{d}{q}} \|v_n\|_{L^r}^2 \leq C \epsilon^{\frac{d}{q}}$$

Conclusion:

$$\int_{\mathbb{R}^d} |\hat{v}_n(k)|^2 = \int_{|k| \leqslant \epsilon} + \int_{|k| > \epsilon} \leqslant C\epsilon^{\frac{d}{q}} + \int_{|k| > \epsilon} \xrightarrow{n \to \infty} C\epsilon^{\frac{d}{q}} \xrightarrow{\epsilon \to 0} 0$$

which contradicts to the assumtion $\|\hat{v}\|_{L^2} = \|v\|_{L^2} = 1$.

Exercise 5.27 Let $\Omega \subseteq \mathbb{R}^d$ be open, bounded with C^1 -boundary. Let $u \in W^{1,p}(\Omega)$, for some $1 \leq p < \infty$. Then the following is equivalent:

a)
$$u \in W_0^{1,p}(\Omega)$$

b)
$$\tilde{u}(x) = \begin{cases} u(x) & x \in \Omega \\ 0 & x \in \mathbb{R}^d \setminus \Omega \end{cases} \in W^{1,p}()$$

Theorem 5.28 (Dirichlet, Riemann, Poincare, Hilbert) Let $\Omega \subseteq \mathbb{R}^d$ be open, bounded with C^1 -boundary. Let $f \in L^2(\Omega)$. Then there exists a unique solution $u \in H_0^1(\Omega)$ of the variational problem

$$\int_{\Omega} \nabla u \nabla \phi = \int_{\Omega} f \phi$$

for all $\phi \in H_0^1(\Omega)$. $(\Rightarrow -\Delta = f \text{ in } D'(\Omega))$. Moreover, u is the unique minimizer of

$$\inf_{v \in H_0^1(\Omega)} \left(\frac{1}{2} \int_{\Omega} |\nabla v|^2 - \int_{\Omega} fv \right)$$

Proof. Let us prove that there is a solution $u \in H_0^1(\Omega)$ for $\inf_{v \in H_0^1(\Omega)} E(v)$, $E(v) = \frac{1}{2} \int_{\Omega} |\nabla v|^2 - \int_{\Omega} fv$.

Step 1: We prove $E > -\infty$. Take $v \in H_0^1(\Omega)$. By the Poincaré and Hölder inequalities:

$$\begin{split} E(v) &= \frac{1}{2} \int_{\Omega} |\nabla v|^2 - \int_{\Omega} f v \\ &\geqslant \frac{1}{2C} \|v\|_{L^2(\Omega)} - \|f\|_{L^2(\Omega)} \|v\|_{L^2(\Omega)} \\ &\geqslant \frac{1}{2C} \|v\|_{L^2(\Omega)}^2 - \left(\frac{1}{4C} \|v\|_{L^2(\Omega)}^2 + C \|f\|_{L^2(\Omega)}^2\right) \\ &\geqslant -C \|f\|_{L^2(\Omega)}^2 > -\infty \end{split}$$

We can also bound:

$$\begin{split} E(v) &= \frac{1}{2} \int_{\Omega} |\nabla v|^2 - \int_{\Omega} f v \\ &\geqslant \frac{1}{4} \int_{\Omega} |\nabla v|^2 - \frac{1}{4C} \int_{\Omega} |v|^2 - \|f\|_{L^2} \|v\|_{L^2} \\ &\geqslant \frac{1}{4} \int_{\Omega} |\nabla v|^2 - C \|f\|_{L^2}^2 \end{split}$$

Step 2: We can take a minimizing sequence $\{v_n\}\subseteq H_0^1(\Omega)$ s.t. $E(v_n)\xrightarrow{n\to\infty} E$. Then:

$$\frac{1}{4} \int_{\Omega} |\nabla v_n|^2 \leqslant E(v_n) + C \|f\|_{L^2}^2 \longrightarrow const.$$

So $|\nabla v_n|$ is bounded in $L^2(\Omega)$. We know that $H_0^1(\Omega)$ is a Hilbert space with norm $||v||_{H_0^1(\Omega)} = ||\nabla v||_{L^2(\Omega)}$ (and the norm is equivalent to the H^1 -norm). Thus $\{v_n\}$ is bounded in $H_0^1(\Omega)$.

Remark 5.29 (Reminder from functional analysis) Let H be a Hilbert space. We say that $v_n \to v$ if $||v_n - v|| \to 0$ and $v_n \to v$ weakly in H if $\langle v_n, \phi \rangle \to \langle v, \phi \rangle$ for all $\phi \in H$.

Theorem 5.30 (Banach-Alaoglu) If H is a Hilbert space and $\{v_n\}$ is a bounded sequence, then there is a subsequence $\{v_{n_k}\}$ s.t. $v_{n_k} \to v$ weakly in H

Remark 5.31 $-v_n \to v \text{ in } H \text{ iff } f(v_n) \to f(v) \text{ for all } f \in H^* = \mathcal{L}(H, \mathbb{R}).$

- If $v_n \to v$ in H, then: $\liminf_{n\to\infty} ||v_n|| \ge ||v||$ (Fatous Lemma)

In fact, for all $\phi \in H \langle v_n, \phi \rangle \to \langle v, \phi \rangle$ and $|\langle v_n, \phi \rangle| \leq ||v_n|| ||\phi||$. This implies

$$\frac{|\langle v, \phi \rangle|}{\|\phi\|} \le \liminf_{n \to \infty} \|v_n\|.$$

So we get

$$||v|| = \sup_{\phi \neq 0} \frac{\langle v, \phi \rangle|}{||\phi||} \le \liminf_{n \to \infty} ||v_n||$$

By the Banach-Alaoglu theorem, up to a subsequence, $v_n \to u$ weakly in $H_0^1(\Omega)$. We prove that u is aminimizer for \mathcal{E}

$$E \longleftarrow \mathcal{E}(v_n) = \frac{1}{2} \int |\nabla v_n|^2 - \int f v_n$$

- Since $v_n \to u$ in $H_0^1(\Omega)$ we have that

$$\liminf_{n \to \infty} \|v_n\|_{H_0^1(\Omega)}^2 \geqslant \|u\|_{H_0^1(\Omega)}^2$$

So we have

$$\liminf_{n \to \infty} \int_{\Omega} |\nabla v_n|^2 \geqslant \int_{\Omega} |\nabla u|^2.$$

– Consider the functional $\mathcal{L}: \phi \in H_0^1(\Omega) \to \int_{\Omega} f \phi$. We claim that \mathcal{L} is continuous. In fact:

$$|\mathcal{L}| = \left| \int_{\Omega} f \phi \right| \le \|f\|_{L^{2}} \|\phi\|_{L^{2}} \le C \|f\|_{L^{2}} \|\nabla f\|_{L^{2}} = C \|f\|_{L^{2}} \|\phi\|_{H_{0}^{1}(\Omega)}$$

Thus from $v_n \to v$ in $H_0^1(\Omega)$ we get $\mathcal{L}(v_n) \to \mathcal{L}(u)$, thus $\int_{\Omega} f v_n \to \int_{\Omega} f u$.

Conclusion: $E = \liminf \mathcal{E}(v_n) \geqslant \mathcal{E}(u)$, so u is a minimizer for \mathcal{E} .

Step 3: Uniqueness. If E has 2 minimizers u_1, u_2 we can prove that $u_1 = u_2$. This is because of the convexity:

$$0 \geqslant \frac{\mathcal{E}(u_1) + \mathcal{E}(u_2)}{2} - \mathcal{E}\left(\frac{u_1 + u_2}{2}\right)$$

$$= \frac{1}{8} \left[2 \int_{\Omega} |\nabla u_1|^2 + 2 \int_{\Omega} |\nabla u_2|^2 - \int_{\Omega} |\nabla (u_1 + u_2)|^2 \right]$$

$$= \frac{1}{8} \int_{\Omega} |\nabla u_1 - \nabla u_2|^2 \geqslant 0$$

This implies that $\nabla(u_1 - u_2) = 0$, so $u_1 - u_2 = const = c_0$. Since $u_1, u_2 \in H_0^1(\Omega)$, we have that $u_1 - u_2 \in H_0^1(\Omega)$ and $c_0 \in C(\overline{\Omega})$. Hence $c_0 = 0$ on $\partial\Omega$, so $c_0 = 0$.

Remark 5.32 We can also prove directly that there is a unique $u \in H_0^1(\Omega)$ s.t.

$$\int_{\Omega} \nabla u \nabla \phi = \int_{\Omega} f \phi \quad \forall \phi \in H_0^1(\Omega)$$

by Riesz theorem. So we get $\langle u, \phi \rangle_{H_0^1(\Omega)} = \mathcal{L}(\phi)$.

Recall the corrector function for the unit ball:

$$\phi_x(y) = G(|x||y - \tilde{x}|), \quad \tilde{x} = \frac{x}{|x|^2}$$

This is ok if $x \neq 0$. When $x \rightarrow 0$:

$$G(|x|(y-\tilde{x})) = G(\underbrace{|x|y-\frac{x}{|x|}}_{|\cdot|\to 1})G(z), \quad |z|=1$$

is well-defined as G is radial.

Question: If $u \in H^1(\Omega)$, then how can we define $u|_{\partial\Omega}$?

5.5 Theory of Trace

Theorem 5.33 (Trace Operator) Let $\Omega \subseteq \mathbb{R}^d$ be open, bounded with C^1 boundary. Then there is a unique linear bounded operator $T: H^1(\Omega) \to L^2(\partial\Omega)$ such that

- If $u \in H^1(\Omega) \cap C(\bar{\Omega})$, then $Tu = u|_{\partial\Omega}$ in the usual restriction sense.
- There is a C > 0 s.t. $||Tu||_{L^2(\partial\Omega)} \leq C||u||_{H^1(\Omega)}$ for all $u \in H^1(\Omega)$

Theorem 5.34 If $u \in H^1(\Omega)$, then $u \in H^1_0(\Omega)$ is equivalent to Tu = 0 in $L^2(\partial\Omega)$. $(H^1_0(\Omega) = T^{-1}(\{0\}))$. The we can discuss about

$$\begin{cases} -\Delta u = f & \text{in } \Omega \\ u|_{\partial\Omega} = g & \text{on } \partial\Omega \end{cases}$$

Lemma 5.35 (Trace inequality on \mathbb{R}^d_+) if $u \in C_c^{\infty}(\mathbb{R}^d)$, then:

 $\|u|_{\partial\mathbb{R}^d_+}\|_{L^2(\partial\mathbb{R}^d_+)}\leqslant C\|u\|_{H^1(\mathbb{R}^d)}\quad\text{with }C>0\text{ independent of }u.$

Proof. $x = (x', x_d) \in \mathbb{R}^{d-1} \times \mathbb{R}$.

$$|u(x',0)|^{2} = -\int_{0}^{\infty} \partial_{d}(|u(x',x_{d})|^{2}) dx_{d}$$

$$= -\int_{0}^{\infty} 2\partial_{d}u(x',x_{d})u(x',x_{d}) dx_{d}$$

$$\leq \int_{0}^{\infty} [|\partial_{d}u(x',x_{d})|^{2} + |u(x',x_{d})|^{2}] dx_{d}$$

This implies:

$$\int_{\mathbb{R}^{d-1}} |u(x',0)|^2 dx' \le \int_{\mathbb{R}^{d-1}} \left(\int_0^\infty [\dots] dx_d \right) dx'$$

$$= \int_{\mathbb{R}^d_+} \left[|\partial_d u|^2 + |u|^2 \right] = ||u||_{H^1(\mathbb{R}^d_+)}^2$$

Corrolary 5.36 If $u \in H^1(Q)$ and u is compactly supported, then:

$$||u||_{L^2(Q_0)} \le ||u||_{H^1(Q_+)}$$

Here

$$\begin{split} Q &= \{x = (x', x_d) \in \mathbb{R}^{d-1} \times \mathbb{R} \mid |x'| < 1, |x_d| < 1\} \\ Q_+ &= \{x = (x', x_d) \in Q \mid x_d > 0\} \\ Q_0 &= \{x = (x', x_d) \in Q \mid x_d = 0\}. \end{split}$$

Proof. We extend u by 0 outside of Q, so $u \in H^1(\mathbb{R}^d)$.

Theorem 5.37 (Extension) If $\Omega \subseteq \mathbb{R}^d$ is open, bounded with C^1 -boundary, then there is a bounded linear operator $B: H^1(\Omega) \to H^1(\mathbb{R}^d)$ s.t.

- $Bu|_{\Omega} = u$ for all $u \in H^1(\Omega)$
- $||Bu||_{H^1(\mathbb{R}^d)} \le C||u||_{H^1(\Omega)}$ and $||Bu||_{L^2(\mathbb{R}^d)} \le C||u||_{L^2(\Omega)}$.

Proof of Theorem 5.33. Since $\partial\Omega$ is C^1 there are open sets $\{U_i\}_{i=1}^N\subseteq\mathbb{R}^d$ such that $\partial\Omega\subseteq\bigcup_{i=1}^NU_i$ and for all i there is a C^1 -diffeomorphism $h_i:U_i\to Q$ s.t. $h_i(U_i)=Q$, $h_i(U_i\cap\Omega)=Q_+$, $h_i(U_i\cap\partial\Omega)=Q_+$. Then there exists a partition of unity $\{\theta_i\}_{i=0}^N\subseteq C^\infty(\mathbb{R}^d)$ s.t.

- 1. $\sum_{i=0}^{N} \theta_i = 1$ for all $x \in \mathbb{R}^d$
- 2. For all i = 1, ..., N: $\theta_i \in C_c^{\infty}(U_i)$
- 3. supp $\theta_0 \subseteq \mathbb{R}^d \setminus \partial \Omega$ (in particular $\theta_0|_{\Omega} \in C_c^{\infty}(\Omega)$)

Then given $u \in H^1(\Omega)$, we can write $u = \sum_{i=0}^N u_i$, where $u_i = \theta_i u$. By the extension theorem (5.37), $u \longrightarrow$ extended to $Bu \in H^1(\mathbb{R}^d)$, thus

$$Bu = \sum_{i=0}^{N} \theta_i(Bu) = \sum_{i=0}^{N} v_i, \quad v_i = \theta_i(Bu)$$

Then $v_i \in H^1(\mathbb{R}^d)$ and v_i is compactly supported in U_i for all $i=1,2,\ldots,N$ and $\operatorname{supp} v_0 \subseteq \mathbb{R}^d \backslash \partial \Omega, \ v_i \in H^1(\mathbb{R}^d)$ and compactly supported inside U_i . This implies $\tilde{v}_i(y) = v_i(h_i^{-1}(y)) \in H^1(Q)$ and compactly supported inside $Q, y \in Q$. Thus $\|\tilde{v}_i\|_{L^2(Q_0)} \leqslant C \|\tilde{v}_i\|_{H^1(Q_+)}$. So we have $\|v_i\|_{L^2(\partial \Omega)} \leqslant C \|\tilde{v}_i\|_{L^2(Q_0)} \leqslant C' \|\tilde{v}\|_{H^1(Q_+)} \leqslant C'' \|v_i\|_{H^1(U_i \cap \Omega)}$. Thus:

$$||u||_{L^{2}(\partial\Omega)} = \left|\left|\sum_{i=1}^{N} v_{i}\right|\right|_{L^{2}(\partial\Omega)} \leqslant \sum_{i=1}^{N} ||v_{i}||_{L^{2}(\partial\Omega)} \leqslant \sum_{i=1}^{N} C'' ||v_{i}||_{H^{1}(U_{i}\cap\Omega)}$$
$$= C'' \sum_{i=1}^{N} ||\theta_{i}u||_{H^{1}(\Omega)} \leqslant C'' \sum_{i=1}^{N} C||u||_{H^{1}(\Omega)}$$

This proof works for $u \in C(\bar{\Omega})$. This implies

$$||u||_{L^2(\partial\Omega)} \leq C||u||_{H^1(\Omega)}$$
 for all $u \in H^1(\Omega) \cap C(\bar{\Omega})$.

This allows us to define

$$T: H^1(\Omega) \longrightarrow L^2(\partial\Omega)$$

 $u \longmapsto u|_{\partial\Omega}$

by continuity. I.e. for all $u \in H^1(\Omega)$ there is $\{u_n\} \subseteq H^1(\Omega) \cap C(\bar{\Omega})$ s.t. $u_n \to u$ in H^1_0 . Then $Tu_n \to Tu$ in $L^2(\partial\Omega)$.

Lemma 5.38 (Extension for Q) Let $u \in H^1(Q_+)$. Then we define $Bu: Q \to \mathbb{R}$ by

$$Bu(x) = \begin{cases} u(x) & x \in Q_{+} \\ -u(x', -x_{d}) & x \in Q_{-} \end{cases},$$

 $x=(x,x_d).$ Then $Bu\in H^1(Q)$ and $Bu|_{Q^+}=u,\ \|Bu\|_{L^2(Q)}^2=2\|u\|_{L^2(Q_+)}^2,\ \|\nabla(Bu)\|_{L^2(Q)}^2=\|\nabla u\|_{L^2(Q_+)}^2$

Proof. It is obvious $Bu|_{Q^+} = u$ and

$$\int_{Q} |Bu|^{2} = \int_{Q_{+}} |Bu|^{2} \int_{Q_{-}} |Bu|^{2}$$

$$= \int_{Q} |u|^{2} + \int_{Q_{-} = \{(x, -x_{d}) | (x, x_{d}) \in Q_{+} \}} |u(x, -x_{d})|^{2}$$

$$= 2 \int_{Q_{+}} |u|^{2}$$

We prove:

$$\nabla(Bu)(x) = \begin{cases} \nabla u(x) & u \in Q_+ \\ \nabla u(x', -x_d) & u \in Q_- \end{cases}$$

First, $\partial_d Bu(x) = \partial_d u(x', -x_d)$ if $x \in Q_-$. Take $\phi \in C_c^{\infty}(Q)$, then:

$$\int_{Q} (Bu(x))(\partial_{d}\phi)(x) dx = \int_{Q_{+}} u\partial_{d}\phi + \int_{Q_{-}} -u(x', -x_{d})\partial_{d}[\phi(x', x_{d})] dx$$

$$(x \to -x_{d}) = \int_{Q_{+}} u\partial_{d}\phi + \int_{Q_{+}} [u(x', x_{d})(\partial_{d}\phi)(x', -x_{d})] dx$$

$$\stackrel{(\phi \notin C_{c}^{\infty}(Q_{+}))}{\approx} \int_{Q_{+}} (\partial_{d}u)\phi(x) + \int_{Q_{+}} (\partial_{d}u(x', x_{d}))\phi(x', -x_{d}) dx$$

$$= -\int_{Q_{+}} (\partial_{d}u)\phi(x) + \int_{Q_{-}} \partial_{d}u(x', -x_{d})\phi(x', x_{d}) dx$$

$$= -\int_{Q} f\phi, \quad \text{where } f(x) = \begin{cases} \partial_{d}u & x \in Q_{+} \\ -\partial_{d}u(x', -x_{d}) & x \in Q_{-} \end{cases}$$

We prove $\int_{Q_+} u \partial_d \tilde{\phi} = -\int_{Q_+} (\partial_d u) \tilde{\phi}$ where $\tilde{\phi}(x, x_d) = \phi(x, x_d) - \phi(x, -x_d)$, $\tilde{\phi} \notin C_c^{\infty}(Q_+)$. Define $\eta_{\epsilon} = 0$ when $|x_d| \leqslant \epsilon$, $\eta_{\epsilon} = 1$ if $|x_d| \geqslant 2\epsilon$, $\eta_{\epsilon} \in C^{\infty}$, $\eta_{\epsilon}(x', x_d) = \eta_0(x', \frac{x_d}{\epsilon})$, $\eta_0 = \begin{cases} 1 & |x_d| \geqslant 2 \\ 0 & |x_d| \geqslant 1 \end{cases}$. We have

$$\int_{Q_+} u \partial_d (\eta_{\epsilon} \tilde{\phi}) = -\int_{Q_+} \partial_d u (\eta_{\epsilon} \tilde{\phi})$$

We take $\epsilon \to 0$,

$$\int_{Q_+} (\partial_d u)(\eta_{\epsilon} \tilde{\phi}) \to \int_{Q_+} (\partial_d u) \tilde{\phi}$$

by dominated convergence

$$\begin{split} \int_{Q_{+}} u \partial_{d} (\eta_{\epsilon} \tilde{\phi}) &= \int_{Q_{+}} u (\partial_{d} \eta_{\epsilon}) \tilde{\phi} + \int_{Q_{+}} u \eta_{\epsilon} \partial_{d} \tilde{\phi} \\ \int_{Q_{+}} u \eta_{\epsilon} \partial_{d} \tilde{\phi} &\xrightarrow{\epsilon \to 0} u \partial_{d} \tilde{\phi} \end{split}$$

by dominated convergence.

$$\left| \int_{Q_{+}} u(\partial_{d}\eta_{\epsilon}) \tilde{\phi} \right| = \left| \int_{Q} u \frac{1}{\epsilon} (\partial_{d}\eta_{0}) \left(x, \frac{x_{d}}{\epsilon} \right) \tilde{\phi} \right|$$

$$\left(\begin{vmatrix} \tilde{\phi}(x', x_{d}) | \\ = |\phi(x, x_{d}) - \phi(x, x_{d}) | \\ \leqslant \|\partial_{d}\phi\|_{L^{\infty}} |x_{d}| \end{vmatrix} \right) \leqslant \frac{1}{\epsilon} \|\partial_{d}\eta_{0}\|_{L^{\infty}} \int_{Q_{+} \cap \{x_{d} \leqslant 2\epsilon\}} |u| \underbrace{\tilde{\phi}}_{\leqslant C|x_{d}|\leqslant C\epsilon}$$

$$\left(\text{Dominated cv } u \in L^{1}(Q_{+}) \right) \leqslant C \int_{Q_{+} \cap \{0 \leqslant x_{d} \leqslant 2\epsilon\}} |u| \xrightarrow{\epsilon \to 0} 0$$

where $u \in L^2(Q_+)$ because $u \in H^1(Q_+)$.

Exercise 5.39 (E. 9.1) Let Ω be open, bounded with C^1 -boundary. Let $u \in H_0^1(\Omega)$, $f \in L^2(\Omega)$. Show that the following statements are equivalent:

- 1) $-\Delta u = f$ in $D'(\Omega)$
- 2) $\int \nabla u \nabla \phi = \int f \phi$ for all $\phi \in H_0^1$
- 3) $E = \inf_{v} \left(\frac{1}{2} \int_{\Omega} |\nabla v|^2 \int_{\Omega} fv \right)$

Solution.

1) \Rightarrow 2) From $-\Delta u = f$ in $D'(\Omega)$ we get that

$$\int_{\Omega} u(-\Delta\phi) = \int_{\Omega} f\phi$$

for all $\phi \in C_c^{\infty}(\Omega)$. Claim: If $u \in H_0^1, \phi \in C_c^{\infty}$, then

$$\int_{\Omega} (-\Delta \phi) = \int_{\Omega} \nabla u \nabla \phi$$

Density argument: $u \in H_0^1 = \overline{C_c^{\infty}(\Omega)}^{\|\cdot\|_H}$, so there is a sequence $\{u_n\} \subseteq C_c^{\infty}(\Omega)$ s.t. $u_n \to u$ in $H^1(\Omega)$. Since $u_n, \phi \in C_c^{\infty}(\Omega)$, then by the integration by parts:

$$\int_{\Omega} u_n(-\Delta\phi) = \int_{\Omega} (\nabla u_n) \nabla \phi \forall n$$

Take $n \to \infty$, then,

$$\int_{\Omega} u(-\Delta\phi) = \int_{\Omega} (\nabla u) \nabla \phi$$

as $u_n \to u$ and $\nabla u_n \to \nabla u$ in L^2 . Claim: If $\int_{\Omega} \nabla u \nabla \phi = \int_{\Omega} f \phi$ for all $\phi \in C_c^{\infty}(\Omega)$, then

$$\int_{\Omega} \nabla u \nabla \phi = \int_{\Omega} f \phi$$

for all $\phi \in H_0^1$. (Given $\nabla u, f \in L^2$). With density argument: For all $\phi \in H_0^1$ there is a sequence $\{\phi_n\} \subseteq C_c^\infty(\Omega)$ s.t. $\phi_n \to \phi$ in H^1 . Then:

$$\int_{\Omega} \nabla u \nabla \phi_n = \int_{\Omega} f \phi_n$$

for all n. Take $n \to \infty$:

$$\int \nabla u \nabla \phi = \int f \phi$$

as $\phi_n \to \phi$, $\nabla \phi_n \to \nabla \phi$ in L^2 .

2) \Rightarrow 3) We show $E(u) \leqslant E(v)$ for all $v \in H_0^1$, i.e.

$$\frac{1}{2} \int_{\Omega} |\nabla u|^2 - \int_{\Omega} f u \leqslant \frac{1}{2} \int_{\Omega} |\nabla v|^2 - \int_{\Omega} f v$$

for all $v \in H_0^1$. Write v = u + w, then:

$$E(v) = \frac{1}{2} \int |\nabla v|^2 - \int fv$$

$$= \frac{1}{2} \int_{\Omega} |\nabla (u+w)|^2 - \int_{\Omega} f(u+w)$$

$$= \frac{1}{2} \int_{\Omega} \left[|\nabla u|^2 + |\nabla w|^2 + 2\nabla u \nabla w \right] \int_{\Omega} (fu+fw)$$

$$= E(u) + \frac{1}{2} \int_{\Omega} |\nabla w|^2 + \left(\underbrace{\nabla u \nabla w - \int fw}_{=0} \right)$$

as $w = v - u \in H_0^1$ (by (2))

$$3) \Rightarrow 1)$$

$$E(u) \leqslant E(u + t\phi)$$

for all $\phi \in H_0^1$ (or C_c^{∞}) for all $t \in \mathbb{R}$. This implies:

$$\frac{d}{dt}E(u+t\phi)|_{t=0} = 0$$

Here

$$E(u+t\phi) = \frac{1}{2} \int \underbrace{|\nabla(u+t\phi)|^2}_{|\nabla u|^2 + t^2 |\nabla \phi|^2 + 2t\nabla u \nabla \phi} - \int f(u+t\phi)$$
$$= E(u) + t \left[\int_{\Omega} \nabla u \nabla \phi \int_{\Omega} f \phi \right] + t^2 \int_{\Omega} |\nabla \phi|^2$$

This implies

$$\frac{d}{dt}E(u+t\phi)|_{t=0} = \int_{\Omega} \nabla u \nabla \phi - \int_{\Omega} f \phi$$

Conclude:

$$\int_{\Omega} \nabla u \nabla \phi = \int_{\Omega} f \phi$$

for all $\phi \in H_0^1$ or C_c^{∞} So we get

$$\int_{\Omega} u(-\Delta\phi) = \int_{\Omega} f\phi$$

for all $\phi \in C_c^{\infty}$. so we can conclude:

$$-\Delta u = f$$

in
$$D'(\Omega) \Rightarrow 1$$

Exercise 5.40 (E 9.2)

$$Q = \{(x', x_d) \mid |x'| < 1, |x_d| < 1\}$$

Given $u \in H^1(Q_+)$, define $Bu: Q \to \mathbb{R}$ as

$$Bu(x) = \begin{cases} u(x) & x \in Q_+ \\ u(\tilde{x}) & x \in Q_- \end{cases},$$

 $x = (x', x_d) \Leftrightarrow \tilde{x} = (x', -x_d), x \in Q_- \Leftrightarrow \tilde{x} \in Q_+$. In the lectures:

$$\partial_d(Bu)(x) = \begin{cases} \partial_d u(x) & x \in Q_+ \\ -(\partial_d u)(\tilde{x}) & x \in Q_- \end{cases}$$

This implies $\partial_d(Bu) \in L^2(Q)$.

1. For all i = 1, ..., d - 1, then:

$$\partial_i(Bu)(x) = \begin{cases} \partial_i u(x) & x \in Q_+ \\ \partial_i u(\tilde{x}) & x \in Q_- \end{cases}$$

2. Example $u \in H^2(Q_+)$ but $Bu \notin H^2(Q)$.

Solution. 1. For all $\phi \in C_c^{\infty}(Q)$:

$$\int_{O} Bu(x)\partial_{i}\phi(x) = \int_{O_{+}} u(x)\partial_{i}\phi(x) + \int_{O_{-}} u(\tilde{x})\partial_{i}\phi(x)$$

Write $\vec{n} = (n_1, \dots, n_d)$. Here:

$$\int_{Q_{+}} u(x)\partial_{i}\phi(x) dx = \int_{Q_{+}} -\partial_{i}u(x)\phi(x) dx + \int_{\partial Q_{+}} u(x)\phi(x)n_{i} dS$$

$$\int_{Q_{-}} u(x', -x_{d})\partial_{i}\phi(x', x_{d}) dx' dx_{d} = -\int_{Q_{+}} u(x', x_{d})\partial\phi(x', -x_{d}) dx' dx_{d}$$

$$= \int_{Q_{+}} \partial_{i}u(x)\phi(\tilde{x}) - \int_{\partial Q_{+}} u\phi n_{i} dS$$

$$= \int_{Q_{-}} -\partial_{i}u(\tilde{x})\phi(x) - \int_{\partial Q_{+}} u\phi n_{i} dS$$

with $d(-x_d) = d(x_d)$. Conclude:

$$\begin{split} \int_Q (Bu)(x)\partial_i\phi(x)\,dx &= \int_{Q_+} (-\partial_i u)(x)\phi(x) + \int_{Q_-} (-\partial_i u)(\tilde x)\phi(x) \\ &= \int_Q -h(x)\phi(x)\,dx, \quad h(x) = \begin{cases} \partial_i u(x), & x\in Q_+\\ \partial_i u(\tilde x), & x\in Q_- \end{cases} \end{split}$$

for all $\phi \in C_c^{\infty}(Q)$, so $\partial_i(Bu) \in L^2$ for all $i = 1, 2, \dots, d-1$. Thus $Bu \in H^1(Q)$.

2. 1D: Take $Q_+(0,1), Q_- = (-1,0), Q_0 = \{0\}, Q = (-1,1), \ u(x) = x \text{ in } Q_+ = (0,1), \ Bu(x) = u(x) = -x \text{ if } x \in Q_- = (-1,0), \text{ i.e. } Bu(x) = |x| \text{ if } x \in Q = (-1,1).$ We know

$$(Bu)'(x) = \begin{cases} 1 & x \in (0,1) \\ -1 & x \in (-1,0) \end{cases} \in L^2(-1,1)$$

i.e. $Bu \in H^1(Q)$.

$$(Bu)''(x) = 2\delta_0(x)$$

in D'(Q) but $\notin L^2(-1,1)$, i.e. $Bu \notin H^2(Q)$. Question: Given $u \in H^2(Q_+)$, can we find an extension $Bu \in H^2(Q)$ Yes! E.g. u(x) = x in (0,1), so Bu(x) = x in (-1,1). In general: $u \in H^2(Q) \leadsto \tilde{u} \in H^2(Q)$ but $\nabla u = 0$ on ∂Q_+ .

Exercise 5.41 (Bonus 8) Assume $u \in H^2(Q_+)$ and $\begin{cases} u = 0 \\ \nabla u = 0 \end{cases}$ on ∂Q_+ Prove that $Bu \in H^2(Q)$. (Reflection extension) (Ok in 1D)

Remark 5.42 If $u \in H^2(Q_+)$, then $\nabla u \in H^1(Q_+)$, so $\nabla u|_{\partial Q_+}$ by trace theory. In general: $\Omega \subseteq \mathbb{R}^d$, C^2 -boundary condition, then the same result holds.

Remark 5.43 In 1D: $\begin{cases} u \in H^2(0,1) \\ u(0) = 0 \\ u'(0) = 0 \end{cases}, \ u|_{Q_0} \in L^2(Q_0), \ \text{1D: } Q_0 = \{0\}. \ \text{In general: }$

If $u \in H^1(0,1)$, then u(0) is determined by trace theory. If $u \in H^2(0,1)$, u'(0) is determined. Sobolev:

$$H^1(0,1) \subseteq C([0,1])$$

 $H^2(0,1) \subseteq C^1([0,1])$

Lemma 5.44 (Poincare inequality) Let Ω be open, bounded connected with C^1 -boundary. Then for all $g \in L^2(\partial \Omega)$ s.t. $g \neq constant$ there is a C > 0 s.t.

$$||u||_{L^2(\Omega)} \leqslant C||\nabla u||_{L^2(\Omega)}$$

for all $u \in M$, where

$$M = \{ v \in H^1(\Omega) \mid v|_{\partial\Omega} = g \}.$$

Proof. We assume that the statement does not hold true. Then there is a sequence $\{u_n\} \subseteq H^1(\Omega), \ u_n|_{\partial\Omega} = g$ s.t.

$$\|\nabla u_n\|_{L^2(\Omega)} \to 0, \quad \|u_n\|_{L^2(\Omega)} = 1.$$

Since $\{u_n\}$ is bounded in $H^1(\Omega)$, by the Banach-Alaoglu theorem (5.30), up to a subsequence

$$u_n \to u_0$$
 weakly in $H^1(\Omega)$

Since $\nabla u_n \to 0$ strongly in L^2 and $\nabla u_n \to \nabla u_0$ weakly in L^2 , we have $\nabla u_0 = 0$, so $u_0|_{\partial\Omega} = const$. (here we need Ω to be connected), so $u_0|_{\partial\Omega} = const$. On the other hand, note that M is convex and closed in $H^1(\Omega)$ since the trace operator $T: H^1(\Omega) \to L^2(\partial\Omega)$ is continuous. Therefore, M is also weakly closed in $H^1(\Omega)$ by the Hahn-Banach theorem. Thus from $\{u_n\} \subseteq M$, $u_n \to u_0$ weakly in $H^1(\Omega)$ we get that $u_0 \in M$, so $u_0|_{\partial\Omega} = g$. We get a contradiction since $g \neq const$

Theorem 5.45 (Solution for Poisson Equation with inhomogeneous boundary condition) Let Ω be open, bounded with C^1 -boundary. Let $f \in L^2(\Omega)$, $g \in L^2(\partial\Omega)$. There there is a unique $u \in H^1(\Omega)$ s.t.

$$\begin{cases} -\Delta u = f & \text{in } D'(\Omega) \\ u|_{\partial\Omega} = g & \text{on } \partial\Omega \end{cases}$$

Here $u|_{\partial\Omega}=T(u)\in L^2(\partial\Omega)$ is defined by the trace operator. Moreover if Ω is connected and $g\neq constant$, then u is the unique minimizer for the variational problem

$$E = \inf_{v \in M} \frac{1}{2} \int_{\Omega} |\nabla v|^2 - \int_{\Omega} f v,$$

where $M = \{v \in H^1(\Omega), v|_{\partial\Omega} = g \text{ on } \partial\Omega\}$

Proof. First let us assume that Ω is connected and $g \neq const.$

Step 1: We prove that $E = \int_{v \in M} E(v)$ has a minimizer. By Poincares Inequality (5.44), for all $v \in M$:

$$\begin{split} E(v) &= \frac{1}{2} \int_{\Omega} |\nabla v|^2 - \int_{\Omega} fv \\ \text{(H\"older)} & \geqslant \frac{1}{2} \ \|\nabla v\|_{L^2(\Omega)}^2 - \|f\|_{L^2(\Omega)} \|v\|_{L^2(\Omega)} \\ \text{(Poincar\'e 5.44)} & \geqslant \frac{1}{2} \|\nabla v\|_{L^2(\Omega)}^2 - C \|f\|_{L^2(\Omega)} \|\nabla v\|_{L^2(\Omega)} \\ & \geqslant \frac{1}{4} \|\nabla v\|_{L^2(\Omega)} - C \|f\|_{L^2(\Omega)} \end{split}$$

Thus $E = \inf_{v \in M} E(v) > -\infty$. Moreover, taking a minimizing sequence $\{v_n\} \subseteq M$, $E(v_n) \to E$, we find that $\|\nabla v_n\|_{L^2(\Omega)}$ is bounded, and hence $\|v_n\|_{H^1(\Omega)}$ is bounded (by Poincaré inequality) again. By Banach-Alaoglu (5.30), up to a subsequence we have $v_n \to u$ weakly in $H^1(\Omega)$. Hence

$$\begin{cases} \limsup_{n \to \infty} \int_{\Omega} |\nabla v_n|^2 \geqslant \int_{\Omega} |\nabla u|^2 & \text{as } \nabla v_n \to \nabla u \text{ in } L^2 \\ \int_{\Omega} v_n f \to \int_{\Omega} u f & \text{as } v_n \to u \text{ in } L^2 \end{cases}$$

Note that $\{v_n\} \subseteq M$, $v_n \to u$ in $H^1(\Omega)$ and M is weakly closed in $H^1(\Omega)$ (as argued in the proof of Poincare inequality), therefore $u \in M$. This means that u is a minimizer for $E = \inf_{v \in M} E(v)$.

Step 2: Now we prove that if u is a minimizer for E, then $-\Delta u = f$ in $D'(\Omega)$. In fact, for all $\phi \in C_c^{\infty}(\Omega)$ we have

$$E(u) \le E(u + t\phi) \quad \forall t \in \mathbb{R}$$

because $u + t\phi \in M$. So we get that

$$0 = \frac{d}{dt}E(u+t\phi)|_{t=0} = \int_{\Omega} \nabla u \nabla \phi - \int_{\Omega} f \phi$$

Thus

$$\int_{\Omega} u(-\Delta \phi) = \int_{\Omega} \nabla u \nabla \phi, = \int_{\Omega} f \phi \quad \forall \phi \in C_c^{\infty}(\Omega).$$

So $-\Delta u = f$ in $D'(\Omega)$.

Step 3: We prove that Poissons equation has at most one solution. Assume that u_1 , u_2 are 2 solutions. Then $u = u_1 - u_2$ solves

$$\begin{cases} -\Delta u = 0 \text{ in } \Omega \\ u|_{\partial\Omega} = 0 \text{ on } \Omega \end{cases}$$

so u = 0.

Step 4: If $g = c_0$ is a constant, then Poissons equation can be rewritten with $\tilde{u} = u - c_0$:

$$\begin{cases} -\Delta u = f & \text{in } \Omega \\ u|_{\partial\Omega} = c_0 & \text{on } \Omega \end{cases} \Leftrightarrow \begin{cases} -\Delta \tilde{u} = f & \text{in } \Omega \\ \tilde{u} = 0 & \text{on } \Omega \end{cases}$$

If Ω is not connected, then by considering connected components of Ω we can prove that Poisson's equation always has a unique solution (for all $f \in L^2(\Omega), g \in L^2(\partial\Omega)$).

5.6 Final Remarks

We can describe $H_0^1(\Omega)$ as the kernel of the trace operator $T: H^1(\Omega) \to L^2(\partial\Omega)$

Theorem 5.46 Let $\Omega \subseteq \mathbb{R}^d$ be open, bounded with C^1 -boundary. Then:

$$H_0^1(\Omega) = \{ u \in H^1(\Omega) \mid T(u) = 0 \text{ on } \partial \Omega \}$$

Recall that if $u \in H^1(\Omega) \cap C(\overline{\Omega})$, then $T(u) = u|_{\partial\Omega}$ is the usual restriction. In this case we recover a result proved before.

Proof.

Recall that the varionational characterization of the Poisson equation

$$\begin{cases} -\Delta u = f & \text{in } \Omega \\ u = g & \text{on } \partial \Omega \end{cases}$$

is

$$\int_{\Omega} \nabla u \nabla \phi = \int_{\Omega} f \phi \quad \forall \phi \in M$$

where $M = \{v \in H^1(\Omega) \mid v = g \text{ on } \partial\Omega\}$ In fact, if $u \in H^2(\Omega)$ and

$$\int_{\Omega} \nabla u \nabla \phi = \int_{\Omega} f \phi \quad \forall \phi \in H^1(\Omega)$$

Then u satisfies the Neumann condition:

$$\frac{\partial u}{\partial n} = \nabla u \cdot \vec{n} = 0 \text{ on } \partial \Omega$$

(justification ...)

For the exercises of sheet 10: Let $\Omega = (a,b) \subseteq \mathbb{R}$ be an open bounded interval. For every $u \in H^1(\Omega)$ the values u(a) and u(b) are determined uniquely by trace theory, or by Sobolev's embedding theorem. Recall: If $u \in H^1((a,b)) \leadsto \partial \Omega = \{a,b\}$ counting measure iff $g \in L^2(\partial \Omega)$ i.e. g(a) = g(b) are well-defined.

Exercise 5.47 (E 10.1) a) Prove $H^1(\mathbb{R}) \subseteq (C(\mathbb{R}) \cap L^{\infty}(\mathbb{R}))$

Hint: You can use Fourier Transform

b)
$$H^1(\Omega) \subseteq C(\Omega)$$

Solution. a) Let $u \in H^1(\mathbb{R})$. Then $u, u' \in L^2(\mathbb{R}) \Leftrightarrow \hat{u}(k)(1 + |2\pi k|) \in L^2(\mathbb{R})$. Thus:

$$u(x) = \int_{\mathbb{R}} \hat{u}(k)e^{2\pi ikx} dk \in C(\mathbb{R}) \cap L^{\infty}(\mathbb{R})$$

if $\hat{u} \in L^1(\mathbb{R})$. So we have to show $\hat{u} \in L^1(\mathbb{R})$.

$$\begin{split} \int_{\mathbb{R}} |\hat{u}(k)| &= \int_{\mathbb{R}} \frac{|g(k)|}{1 + |2\pi k|} \\ &\leqslant \left(\int_{\mathbb{R}} |g(k)|^2 \, dk \right) \left[\int_{\mathbb{R}} \left(\frac{1}{1 + |2\pi k|} \right)^2 \, dk \right]^{\frac{1}{2}} < \infty \end{split}$$

b) Given $u \in H^1(\Omega)$, then there is an extension $\tilde{u} \in H^1(\mathbb{R})$. By a) $\tilde{u} \in C(\mathbb{R})$, so $u = \tilde{u}|_{\tilde{\Omega}} \in C(\bar{\Omega})$. Remak: We have $\|u\|_{L^{\infty}(\Omega)} \leq C\|u\|_{H^1(\Omega)}$, where $\Omega = (a,b)$ or \mathbb{R} (but only in 1D)

Recall: If $\Omega \subseteq \mathbb{R}^d (d \ge 1)$ open, bounded with C^1 -boundary. Then

$$||u||_{L^2(\Omega)} \leq C||\nabla u||_{L^2(\Omega)} \quad \forall u \in H_0^1(\Omega)$$

Actually the same bound holds if $u \in H^1(\Omega)$ and $u|_{\Gamma} = 0$ for an open subset $\Gamma \subseteq \partial \Omega$. In 1D we have:

Exercise 5.48 (E 10.2 (Poincare inequality)) Let $u \in H^1(\Omega)$, u(a) = 0. Prove that there exists a constant C > 0 such that

$$||u||_{L^2(\Omega)} \leq C||u'||_{L^2(\Omega)}$$

Solution. Let $u \in C^1(\bar{\Omega})$ and u(a) = 0. Then:

$$u(x) = u(a) + \int_0^x u'(t) dt \quad \forall x \in (a, b)$$

$$\Rightarrow |u(x)| \le \int_a^x |u'(t)| dt \le \int_a^b |u'(t)| dt = ||u'||_{L^1(\Omega)} \le C||u'||_{L^2(\Omega)}$$

as Ω is bounded. This implies:

$$\frac{1}{C} \|u\|_{L^{2}(\Omega)} \le \|u\|_{L^{\infty}(\Omega)} \le C \|u'\|_{L^{2}(\Omega)}$$

To extend this for $u \in H^1(\Omega)$, we can use a density argument. More precisely, for all $u \in H^1(\Omega)$ there is a sequence $\{u_n\} \subseteq C^1(\bar{\Omega})$ s.t $u_n \to u$ in $H^1(\Omega)$. Then:

$$||u||_{L^{2}(\Omega)} = \lim_{n \to \infty} ||u_{n}||_{L^{2}(\Omega)} \le C \lim_{n \to \infty} ||u'_{n}||_{L^{2}(\Omega)} = C ||u'||_{L^{2}(\Omega)}$$

Recall: For all $f \in W_{loc}^{1,1}(O)$ with O in \mathbb{R}^d we have

$$f(x) - f(y) = \int_0^1 \nabla f(y + t(x - y))(x - y) dt$$

if $x, y \in O$, $y + t(x - y) \in O$ for all $t \in [0, 1]$. For 1D: If $u \in H^1(a, b)$:

$$u(x) - u(y) = \int_{y}^{x} u'(t) dt \quad \forall x, y \in (a, b)$$

Exercise 5.49 (E 10.3 (Poincare inequality)) Let $u \in H^2(\Omega)$ and $f \in L^2(\Omega)$. Prove that the following statements are equivalent:

a) u solves the equation:

$$\begin{cases} -u'' = f & \text{in } D'(\Omega) \\ u'(0) = u'(1) = 0 \end{cases}$$

b)

$$\int_{\Omega} u'\phi' = \int_{\Omega} f\phi$$

for all $\phi \in H^1(\Omega)$.

Here $u \in H^2(\Omega) \Rightarrow u' \in H^1(\Omega) \Rightarrow u'(0), u'(1)$ determined uniquely by trace theorem / Sobolev inequality $H^1(\Omega) \subseteq C(\bar{\Omega})$

Solution.

b) \Rightarrow a) For all $\phi \in C_c^{\infty}(\Omega)$:

$$\int_{\Omega} f\phi = \int_{\Omega} u'\phi' = -\int_{\Omega} u\phi''$$

This implies -u''=f in $D'(\Omega)$ a.e. Thus for all $\phi\in H^1(\Omega)$:

$$\int_{\Omega} f\phi = \int_{\Omega} -u''\phi = \int_{\Omega} u'\phi' - [u'\phi]_a^b$$

By b) we conclude $0 = [u'\phi]_a^b = u'(b)\phi(b) - u'(a)\phi(a)$ for all $\phi \in H^1(\Omega)$. We can choose $\phi \in H^1(\Omega)$ s.t. $\phi(a) = 0$, $\phi(b) = 1$. This implies $\phi'(b) = 0$. Similarly, we can chose $\phi \in H^1(\Omega)$ s.t. $\phi(a) = 1$, $\phi(b) = 0$. This implies u'(a) = 0.

a) \Rightarrow b) From a) and Integration by parts:

$$\int_{\Omega} f\phi = \int_{\Omega} -u''\phi = \int_{\Omega} u'\phi' - \underbrace{[u'\phi]_a^b}_{=0 \text{ as } u'(a)=u'(b)=0}$$

This implies:

$$\int_{\Omega} f\phi = \int_{\Omega} u'\phi' \quad \forall \phi \in H^1(\Omega)$$

Exercise 5.50 (E 10.4 (Robin boundary condition)) Let $f \in L^2(\Omega)$.

a) Prove that there exists a unique $u \in M := \{\phi \in H^1(\Omega), u(a) = 0\}$ s.t.

$$\int_{\Omega} u'\phi' = \int_{\Omega} f\phi \quad \forall \phi \in M$$

b) Prove that the above function u is the unique solution to the equation

$$\begin{cases} -u'' = f & \text{in } D'(\Omega) \\ u(a) = 0 & u'(b) = 0 \end{cases}$$

Solution. a) By 10.2 we have

$$\|\phi\|_{L^2(\Omega)} \leqslant C \|\phi'\|_{L^2(\Omega)} \quad \forall \phi \in M$$

Thus: $(M, \|\phi\|_M := \|\phi'\|_{L^2(\Omega)})$ is a Hilbert space. More precisely, we know $(M, \|\cdot\|_M)$ is a closed subspace of $H^1 \leadsto$ a Hilbert space. And $\|\cdot\|_M$ is comparable to $\|\cdot\|_{H^1}$. By Riesz representation theorem there is a unique $u \in M$ s.t. $\langle \phi, u \rangle_M = F(\phi)$ for all $\phi \in M$. We use this for

$$F(\phi) = \int_{\Omega} f\phi \quad \forall \phi \in M$$

Here $|F(\phi)| \leq ||f||_{L^2} ||\phi||_{L^2}$.

b) Let $u \in M$ be the solution in (a) i.e.

$$\int_{\Omega} f\phi = \int_{\Omega} u'\phi' \quad \forall \phi \in M$$

Then we prove that u solves

$$\begin{cases} -u'' = f & \text{in } D'(\Omega) \\ u(a) = u'(b) = 0 \end{cases}$$

Since $u \in M$ we have $u \in H^1(\Omega)$ and u(a) = 0. From

$$\int_{\Omega} f\phi = \int_{\Omega} u'\phi' \quad \forall \phi \in M$$

we get for all $\phi \in C_c^{\infty}(\Omega)$:

$$\int_{\Omega} f\phi = \int_{\Omega} u'\phi' = \int_{\Omega} -u\phi''$$

So we get -u'' = f in $D'(\Omega)$. Since $f \in L^2(\Omega) \Rightarrow u'' \in L^2(\Omega) \Rightarrow u \in H^2(\Omega)$ $\Rightarrow u' \in H^1(\Omega) \Rightarrow u'(b)$ is uniquely determined. For all $\phi \in M$:

$$\int_{\Omega} f\phi = \int_{\Omega} -u''\phi = \int_{\Omega} u'\phi' - \left(u'(b)\phi(b) - u'(a)\phi(a)\right) \quad \text{as } \phi \in M$$

and $\int_{\Omega} f \phi = \int_{\Omega} u' \phi'$. This implies:

$$u'(b)\phi(b) = 0 \quad \forall \phi \in M$$

Take $\phi(x) = \frac{x-a}{b-a} \in M$, $\phi(b) = 1$. Uniqueness of the solution: Take u s.t.

$$\begin{cases} -u'' = f & \text{in } D'(\Omega) \\ u(a) = u'(b) = 0 \end{cases}$$

This implies $u \in H^2(\Omega)$. By integration by parts: For all $\phi \in H^1(\Omega)$, $\phi(a) = 0$.

$$\int_{\Omega} f \phi = \int_{\Omega} -u'' \phi = \int u' \phi' \quad \forall \phi \in M$$

Thus $u \in M$ and

$$\int_{\Omega} f \phi = \int_{\Omega} u' \phi' \quad \forall \phi \in M.$$

Exercise 5.51 (Bonus 9) Prove that the solution u in Problem E 10.4 is the unique minimizer for the minimization problem:

$$E = \inf_{v \in M} \left(\int_{\Omega} |v'|^2 - \int_{\Omega} fv \right)$$

Chapter 6

Heat Equation

6.1 Fundamental Solution

$$\begin{cases} \partial_t u = \Delta u & (x,t) \in \mathbb{R}^d \times (0,\infty) \\ u = g & (x,t) \in \mathbb{R}^d \times \{0\} \end{cases}$$

The fundamential solution is:

$$\Phi(x,t) = \frac{1}{(4\pi t)^{\frac{d}{2}}} e^{-\frac{|x|^2}{4t}}, \quad x \in \mathbb{R}^d, t > 0$$

We have:

$$\begin{cases} \partial_t \Phi = \Delta \Phi & (x,t) \in \mathbb{R}^d \times (0,\infty) \\ \int_{\mathbb{R}^d} \Phi(x,t) \, dx = 1 & \forall t > 0 \\ \lim_{t \to 0} \Phi(x,t) = \delta_0(x) & \text{in } D'(\mathbb{R}^d) \end{cases}$$

Theorem 6.1 If $g \in C(\mathbb{R}^d) \cap L^{\infty}(\mathbb{R}^d)$, then

$$u(x,t) := \int_{\mathbb{R}^d} \Phi(x-y,t)g(y) \, dy$$

satisfies

- (i) $u \in C^{\infty}(\mathbb{R}^d \times (0, \infty))$
- (ii) $\partial_t u = \Delta u$ for all $(x,t) \in \mathbb{R}^d \times (0,\infty)$
- (iii) $\lim_{t\to 0} u(x,t) = g(x)$ for all $x \in \mathbb{R}^d$

Notation 6.2 For functions of (x,t) we introduce the following notation for different regularity in x and t.

$$f \in C_1^2 \Leftrightarrow f, D_x f, D_x^2 f, \partial_t f \in C$$

Theorem 6.3 (Nonhomogeneous problem) Let $f \in C_1^2(\mathbb{R}^d, [0, \infty))$ be compactly supported. Define

$$u(x,t) = \int_0^t \int_{\mathbb{R}^d} \Phi(x-y,t-s) f(y,s) \, dy \, ds$$

Then

(i)
$$u \in C_1^2(\mathbb{R}^d \times (0, \infty))$$

(ii)
$$\partial_t u = \Delta u + f$$
 for all $x \in \mathbb{R}^d, t > 0$

(iii)
$$\lim_{t\to 0} u(x,t) = 0$$
 for all $x \in \mathbb{R}^d$.

Proof. We write

$$u(x,t) = \int_0^t \int_{\mathbb{R}^d} \Phi(y,s) f(x-y,t-s) \, dy \, ds$$

With the Leibniz integral rule we get

$$\partial_t u(x,t) = \int_0^t \int_{\mathbb{R}^d} \Phi(y,s) \partial_t f(x-y,t-s) \, dy \, ds + \int_{\mathbb{R}^d} \Phi(y,s) f(x-y,0) \, dy$$

and

$$\partial_{ij}u(x,t) = \int_0^t \int_{\mathbb{R}^d} \Phi(y,s) \partial_{ij} f(x-y,t-s) \, dy.$$

This shows that $\partial_t u$, $\partial_{ij} u$ are in $C(\mathbb{R}^d \times (0, \infty))$. Next we calculate:

$$\partial_t u - \Delta u = \int_0^t \int_{\mathbb{R}^d} \Phi(y, s) (\partial_t - \Delta_x) f(x - y, t - s) \, dy \, ds + \int_{\mathbb{R}^d} \Phi(y, s) f(x - y, 0) \, dy$$

$$= \underbrace{\int_\epsilon^t \int_{\mathbb{R}^d} \Phi(y, s) (\partial_t - \Delta_x) f(x - y, t - s) \, dy \, ds}_{=:I_\epsilon}$$

$$+ \underbrace{\int_0^\epsilon \int_{\mathbb{R}^d} \Phi(y, s) (\partial_t - \Delta_x) f(x - y, t - s) \, dy \, ds}_{J_\epsilon}$$

$$+ \underbrace{\int_{\mathbb{R}^d} \Phi(y, s) f(x - y, 0) \, dy}_{K}$$

Then

$$|J_{\epsilon}| \leq \|(\partial_{t} - \Delta_{x})f\|_{L^{\infty}} \int_{0}^{\epsilon} \int_{\mathbb{R}^{d}} \Phi(y, s) \, dy \, ds \leq C\epsilon \xrightarrow{\epsilon \to 0} 0$$

$$I_{\epsilon} = \int_{\epsilon}^{t} \int_{\mathbb{R}^{d}} \Phi(y, s)(-\partial_{s} - \Delta_{y})f(x - y, t - s) \, dy \, ds$$

$$(Green (2.3)) = \int_{\epsilon}^{t} \int_{\mathbb{R}^{d}} \underbrace{(\partial_{s} - \Delta_{y})\Phi(y, s)}_{=0} f(x - y, t - s) \, dy \, ds$$

$$- \left[\int_{\mathbb{R}^{d}} \Phi(y, s)f(x - y, t - s) \right]_{s = \epsilon}^{s = t}$$

This implies:

$$I_{\epsilon} + K = \int_{\mathbb{R}^d} \Phi(y, \epsilon) f(x - y, t - \epsilon) \, dy$$

$$\xrightarrow{\epsilon \to 0} \int_{\mathbb{R}^d} \delta_0(y) f(x - y, t) \, dy = f(x, t)$$

Thus

$$\partial_t u - \Delta u = f(x,t) \quad \forall (x,t) \in \mathbb{R}^d \times (0,\infty)$$

Finally:

$$||u(\cdot,t)||_{L^{\infty}} \leqslant ||f||_{L^{\infty}} \int_0^t \int_{\mathbb{R}^d} \Phi(y,s) \, dy \, ds = ||f||_{L^{\infty}} t \xrightarrow{t \to 0} 0$$

Exercise 6.4 If f, g are given as above, then

$$u(x,t) = \int_{\mathbb{R}^d} \Phi(x-y,t)g(y) \, gy + \int_0^t \int_{\mathbb{R}^d} \Phi(x-y,t-s)f(y,s) \, ds$$

solves

$$\begin{cases} \partial_t - \Delta u = f \\ u(\cdot, t) = g \end{cases}$$

Remark 6.5 (Duhamel formula) Consider the ODE $\partial_t w(t) = Aw(t)$ for all $A \in \mathbb{R}$. Then the solution is

$$w(t) = e^{tA}w(0).$$

More generally: If $\partial_t w(t) = Aw(t) + f(t)$, then

$$\begin{split} \partial_t(e^{-tA}w(t)) &= e^{-tA}(\partial_t w(t) - Aw(t)) = e^{-tA}f(t) = e^{-tA}f(t) \\ \Rightarrow & e^{-tA}w(t) = w(0) + \int_0^t e^{-sA}f(s)\,ds \\ \Rightarrow & w(t) = e^{tA}w(0) + \int_0^t e^{(t-s)A}f(s)\,ds \end{split}$$

More generally, if A is an operator (independent of time) then:

$$\partial_t w(t) = Aw(t) + f(t)$$

$$\Rightarrow \qquad w(t) = e^{tA}w(0) + \int_0^t e^{(t-s)A}f(s) \, ds$$

Application: If $A = \Delta$, then the operator $e^{t\Delta}$ has kernel

$$e^{t\Delta}(x,y) = \Phi(x-y,t) = \frac{1}{(4\pi t)^{\frac{d}{2}}} e^{-\frac{|x-y|^2}{4t}}.$$

This is called the *heat kernel*.

Theorem 6.6 (L^2 -data) For every $g \in L^2(\mathbb{R}^d)$, define

$$u(t,x) = \int_{\mathbb{R}^d} \Phi(x-y,t)g(y) \, dy$$

Then $u \in C^{\infty}(\mathbb{R}^d \times (0, \infty))$ and it solves the heat equation

$$\begin{cases} \partial_t u = \Delta_x u & \mathbb{R}^d \times (0, \infty) \\ \lim_{t \to 0} u(\cdot, t) = g & \text{in } L^2(\mathbb{R}^d) \end{cases}$$

Proof. Recall the heuristic computation from the heat equation using the Fourier transform

$$\begin{array}{ll} \partial_t u(x,t) = \Delta_x u(x,t) \\ \Leftrightarrow & \partial_t \hat{u}(k,t) = -|2\pi k|^2 \hat{u}(k,t) \\ \Leftrightarrow & \partial_t (e^{t|2\pi k|^2} \hat{u}(k,t)) = 0 \\ \Leftrightarrow & e^{t|2\pi k|^2} \hat{u}(k,t) = \hat{u}(k,0) = \hat{g}(k) \\ \Leftrightarrow & \hat{u}(k,t) = e^{-t|2\pi k|^2} \hat{g}(k) = \hat{\Phi}(k,t) \hat{g}(k) = \widehat{\Phi \star g} \\ \Leftrightarrow & u(x,t) = \Phi \star g = \int_{\mathbb{R}^d} \Phi(x-y,t) g(y) \, dy \end{array}$$

Here we only need the direction \Leftarrow which is rigorous if $g \in L^2(\mathbb{R}^d)$. From the Fourier transform, it is also easy to check that $u(\cdot,t) \to g$ in L^2 as $t \to 0$ (exercise). To see the smoothness, note that for all t > 0, and for all $m \in \mathbb{N}$:

$$(1 + |2\pi k|^m)\hat{u}(k,t) = \underbrace{(1 + |2\pi k|^m)e^{-t|2\pi k|^2}}_{\in L^\infty} \underbrace{\hat{g}(k)}_{\in L^2} \in L^2$$

This implies $u(\cdot,t) \in H^m(\mathbb{R}^d)$ for all $m \ge 1$, so $u(\cdot,t) \in C^\infty(\mathbb{R}^d)$ by Sobolev embedding (see below). This argument can also be used to show that $u \in C^\infty(\mathbb{R}^d \times (0,\infty))$ (exercise)

Theorem 6.7 (Sobolev embedding) If $m > \frac{d}{2}$, then $H^m(\mathbb{R}^d) \subseteq (C(\mathbb{R}^d) \cap L^{\infty}(\mathbb{R}^d))$.

Proof. We write for all $u \in H^m(\mathbb{R}^d)$:

$$\hat{u}(k) = \underbrace{\hat{u}(k)(1 + |2\pi k|^m)}_{\in L^2 \text{ as } u \in H^m} \underbrace{\frac{1}{1 + |2\pi k|^m}}_{\in L^2 \text{ as } m > \frac{d}{2}}$$

This implies $\hat{u}(k) \in L^1(\mathbb{R}^d)$ and finally $u = (\hat{u})^{\vee} \in (C(\mathbb{R}^d) \cap L^{\infty}(\mathbb{R}^d))$.

Exercise 6.8 (E 11.1) Let $g \in L^2(\mathbb{R}^d)$,

$$u(x,t) = \int_{\mathbb{R}^d} \Phi(x-y,t)g(y) dy,$$
 $\Phi(x,t) = \frac{1}{(4\pi t)^{\frac{d}{2}}} e^{-\frac{|x|^2}{4t}}$

be the fundamential solution of the heat equation

$$\begin{cases} \partial_t u - \Delta_x u = 0 & \forall (x, t) \in \mathbb{R}^d \times (0, \infty) \\ u(x, t) \to g(x) & \text{as } t \to 0. \end{cases}$$

Prove that

- a) $u \in C^{\infty}(\mathbb{R}^d \times (0, \infty))$.
- b) $||u(\cdot,t) g||_{L^2(\mathbb{R}^d)} \xrightarrow{t \to 0^+} 0$
- c) If $g \in H^1(\mathbb{R}^d)$, then $||u(\cdot,t) g||_{L^2(\mathbb{R}^d)} \le C\sqrt{t}$ as $t \to 0^+$.

Solution. a) We prove for all t > 0:

$$u(x,t)\in \bigcap_{m\geqslant 1}H^m(\mathbb{R}^d)\subseteq C^\infty(\mathbb{R}^d)$$

We use the Fourier transform:

$$\hat{\Phi}(k,t) = e^{-t|2\pi k|^2}$$

Recall $\widehat{e^{-\pi|x|^2}} = e^{-\pi|k|^2}$. From this we get $\widehat{e^{-\pi\lambda|x|^2}} = \lambda^{-\frac{d}{2}} e^{-\frac{\pi|k|^2}{\lambda}}$. Then:

$$\widehat{e^{-\frac{|x|^2}{4t}}} = \widehat{e^{-\pi \frac{1}{4\pi t}|x|^2}} = \left(\frac{1}{4\pi t}\right)^{-\frac{d}{2}} e^{-\pi|k|^2 4\pi t} = (4\pi t)^{\frac{d}{2}} e^{-t|2\pi k|^2}$$

Hence:

$$\hat{u}(k,t) = \hat{\Phi}(k,t)\hat{g}(k) = e^{-t|2\pi k|^2}\hat{g}(k) \in L^1(\mathbb{R}^d, dk) \quad \forall t > 0$$

This implies:

$$u(x,t) = \int_{\mathbb{R}^d} e^{-t|2\pi k|^2} \hat{g}(k) e^{2\pi i k x} dk \quad \forall (x,t) \in \mathbb{R}^d \times (0,\infty)$$

Consequently:

$$D_x^{\alpha}u(x,t) = \int_{\mathbb{R}^d} \underbrace{e^{-t|2\pi k|^2} \hat{g}(k)(2\pi i k)^{\alpha}}_{L^1(\mathbb{R}^d,dk)} e^{2\pi i k x} \, dk \in C(\mathbb{R}^d,(0,\infty))$$

$$D_t^{\alpha} u(x,t) = \int_{\mathbb{R}^d} (-|2\pi k|^2)^{\alpha} e^{-t|2\pi k|^2} \hat{g}(k) e^{2\pi i k x} \, dk \in C(\mathbb{R}^d, (0, \infty))$$

Also:

$$\begin{split} \partial_t u - \Delta_x u \\ &= \int_{\mathbb{R}^d} -|2\pi k|^2 e^{-t|2\pi k|^2} \hat{g}(k) e^{2\pi i k x} \, dk + \int_{\mathbb{R}^d} e^{-t|2\pi k|^2} \hat{g}(k) |2\pi i k|^2 e^{2\pi i k x} \, dk \\ &= 0 \end{split}$$

b) Finally:

$$\int_{\mathbb{R}^d} |u(x,t) - g(x)|^2 = \int_{\mathbb{R}^d} |\hat{u}(k,t) - \hat{g}(k)|^2 dk$$

$$= \int_{\mathbb{R}^d} \underbrace{|e^{-t|2\pi k|^2} - 1|^2}_{\in [0,1]} \underbrace{|\hat{g}(k)|^2}_{\in L^1(\mathbb{R}^d)} dk \xrightarrow{t \to 0} 0$$

by dominated convergence. Now,

$$\int_{\mathbb{R}^d} |u(x,t)|^2 dx = \int_{\mathbb{R}^d} |\hat{u}(k,t)|^2 dk$$

$$= \int_{\mathbb{R}^d} \underbrace{e^{-2t|2\pi k|^2}}_{\in [0,1] \text{ and } \xrightarrow{t \to 0} 0} |\hat{g}(k)|^2 dk \xrightarrow{t \to \infty} 0$$

c) Assume $g \in H^1(\mathbb{R}^d) \Leftrightarrow \int_{\mathbb{R}^d} (1 + |2\pi k|^2) |\hat{g}(k)|^2 dk < \infty$. We claim for all $s \ge 0$ that $|1 - e^{-s}| \le \min(1, Cs) \le C\sqrt{s}$: We have that $s \mapsto \left|\frac{1 - e^{-s}}{s}\right|$ is bounded and continuous in [0, 1] as $\left|\frac{1 - e^{-s}}{s}\right| \to 1$, so $\frac{1 - e^{-s}}{s} \le C$ for all $s \in [0, 1]$.

$$\int_{\mathbb{R}^d} |u(x,t) - g(x)|^2 dx = \int_{\mathbb{R}^d} \underbrace{\left| 1 - e^{-t|2\pi k|^2} \right|^2}_{\leqslant C(t|2\pi k|^2)} |\hat{g}(k)|^2 dk$$

$$\leqslant C \int_{\mathbb{R}^d} t|2\pi k|^2 |\hat{g}(k)|^2 dk$$

$$\leqslant Ct \|g\|_{H^1}^2 \quad \forall t > 0$$

Step 1: Spectral problem:

$$\begin{cases} -\Delta u_n = \lambda_n u_n & \text{in } \Omega \\ u_n|_{\partial\Omega} = 0 \end{cases}$$

Lemma 6.9 There is a $\lambda_n > 0$, $\lambda_n \xrightarrow{n \to \infty} \infty$ and an orthonormal family $\{u_n\} \subseteq L^2(\Omega)$ s.t. $u_n \in H^1_0(\Omega) \cap C^\infty(\Omega)$ solving this eigenvalue equation.

Step 2:

$$\begin{cases} \partial_t - \Delta_x u = 0 \\ u(x,0) = g(x) \end{cases} \Rightarrow \begin{cases} \partial_t \langle u_n, u \rangle_{L^2(\Omega)} = \langle u_n, \Delta_x u \rangle = \langle \Delta_x u_n, u \rangle = -\lambda_n \langle u_n, u \rangle \\ \langle u_n, u \rangle_{t=0} = \langle u_n, g \rangle \end{cases}$$

$$\Rightarrow \qquad \langle u_n, u \rangle = e^{-t\lambda_n} \langle u_n, g \rangle \qquad \forall t > 0, \forall n = 1, 2, \dots$$

$$\Rightarrow \qquad u = \sum_{n=0}^{\infty} \langle \rangle = -\sum_{n=0}^{\infty} e^{-t\lambda_n} \langle \rangle u$$

Example 6.10 $\Omega = (0, 1)$,

$$\begin{cases}
-u_n'' = \lambda_n u_n & \text{in } (0, 1) \\
u(0) = u(1) = 0
\end{cases}$$

has solution

$$\begin{cases} u_n(x) = \sqrt{2}\sin(\pi nx) & n = 1, 2, \dots \\ \lambda_n = (\pi_n)^2 \end{cases}$$

has a solution:

$$u(x,t) = \sum_{n=1}^{\infty} e^{-t\lambda_n} \underbrace{\langle u_n, g \rangle}_{g_n} u_n(x) = \sum_{n=1}^{\infty} e^{-t\pi^2 n^2} g_n \sin(\pi n x),$$
$$\int_0^1 \sin(n\pi x)^2 dx = \frac{1}{2} \quad \forall n > 1$$
$$g_n = \sqrt{2} \langle u_n, g \rangle = 2 \int_0^1 \sin(\pi n x) g(x) dx$$

Exercise 6.11 (E 11.2) Consider the heat equation in a bounded domain

$$\begin{cases} \partial_t u(x,t) = \Delta_x u(x,t) & \forall x \in \Omega, t > 0 \\ u(x,t) = 0 & \forall x \in \partial\Omega, t > 0 \\ u(x,0) = g(x) & \forall x \in \Omega \end{cases}$$

Let us focus on the simlest case $\Omega=(0,1).$ Prove that for every $g\in C^1_c(0,1),$ the function

$$u(x,t) = \sum_{n=1}^{\infty} g_n e^{-t\pi^2 n^2} \sin(n\pi x),$$
 $g_n = 2 \int_0^1 g(y) \sin(n\pi y) dy$

is a classical solution to the above heat equation.

Solution. Direct proof of heat equation. $g \in C^1_c(0,1) \subseteq H^1_0(0,1), \Rightarrow \sum_n \pi^2 n^2 |g_n|^2 = c\|g'\|^2_{L^2(0,1)} < \infty$, so $\sum_n |g_n| < \infty$.

$$u(x,0) = \underbrace{\sum_{n=1}^{\infty} g_n \sin(\pi nx)}_{\in C[0,1]} = g(x) \quad \forall x \in [0,1]$$

From $u(x,t) = \sum_{n=1}^{\infty} e^{-tn^2\pi^2} g_n \sin(\pi nx)$ we get

$$\begin{cases} \partial_t u(x,t) = \sum_{n=1}^{\infty} (-n^2 \pi^2) e^{-t\pi^2 n^2} g_n \sin(\pi n x) & \forall t > 0, \forall x \in (0,1) \\ \Delta_x u(x,t) = \sum_{n=1}^{\infty} e^{-t\pi^2 n^2} g_n [-(\pi n)^2] \sin(\pi n x) & \forall t > 0, \forall x \in (0,1) \end{cases}$$

So
$$\partial_t u - \Delta_x u = 0$$
 for all $t > 0, x \in (0, 1)$

Exercise 6.12 (E 11.3) Let $g(t) = e^{-\frac{1}{t^2}}$ and denote $g^{(n)}(t)$ the *n*-th derivative of g. Define

$$u(x,t) = \sum_{n=0}^{\infty} \frac{g^{(n)}(t)}{(2n)!} x^{2n}, \quad \forall x \in \mathbb{R}, t > 0$$

Prove that u is a classical solution to the heat equation

$$\begin{cases} \partial_t u(x,t) = \Delta_x u(x,t) & \forall x \in \mathbb{R}, t > 0 \\ \lim_{t \to 0} u(x,t) = 0 & \forall x \in \mathbb{R} \end{cases}$$

Solution. Formally:

$$\begin{cases} \partial_t u = \sum_{n=0}^{\infty} \frac{g^{(n+1)}(t)}{(2n)!} x^{2n} \\ -\Delta_x u = \sum_{n=1}^{\infty} \frac{g^{(n)}}{(2n)!} (2n)(2n-1) x^{2n-2} = \sum_{n=1}^{\infty} \frac{g^{(n)}(t)}{(2n-2)!} x^{2n-2} = \sum_{m=0}^{\infty} \frac{g^{(m+1)(t)}}{(2m)!} x^{2m} \end{cases}$$

This implies $\partial_t u = \Delta_x u$ (if the series are convergent) $(x,t) \in B \times \left[\epsilon, \frac{1}{\epsilon}\right]$ for $B \subset \mathbb{R}$ bounded, $\epsilon > 0$. Also

$$g(t) = e^{-\frac{1}{t^2}} \xrightarrow{t \to 0^+} e^{-\infty} = 0$$

$$g'(t) = e^{-\frac{1}{t^2}} \left(\frac{2}{t^3}\right) \xrightarrow{t \to 0^+} 0$$

$$g''(t) = e^{\frac{1}{t^2}} \left(-\frac{3!}{t^4} + \frac{2}{t^3}\right) \xrightarrow{t \to 0^+} 0$$

$$g'''(t) = e^{-\frac{1}{t^2}} \left(\frac{4!}{t^5} - \frac{3!}{t^4} + \frac{2}{t^3}\right)$$

Let's proof the convergence of the series:

$$u(x,t) = \sum_{n=0}^{\infty} \frac{g^{(n)}(t)}{(2n)!} x^{2n}$$

converges absolutely for $|x| \leq C, t \in \left[\epsilon, \frac{1}{\epsilon}\right], \epsilon > 0$. By induction,

$$g^{(n)}(t) = e^{-\frac{1}{t^2}} \underbrace{\left(\frac{(n+1)!}{t^{n+2}} - \frac{n!}{t^{n+1}} + \frac{(n+1)!}{t^n} - \dots\right)}_{\text{pol in } (\frac{1}{t}), \text{ all cos bounded by } (n+1)} (-1)^{n-1}$$

This implies

$$|g^{(n)}(t)| \le e^{-\frac{1}{t^2}} [(n+2)!] \left(\frac{1}{t^{n+2}} + 1\right), \quad \frac{1}{t^s} \le \left(\frac{1}{t^{n+2}} + 1\right) \forall s = 0, 1, \dots, n+2$$

Thus

$$\sum_{n\geqslant 0} \left| \frac{g^{(n)}}{(2n)!} x^{2n} \right| \leqslant \sum_{n\geqslant 0} e^{-\frac{1}{t^2}} \frac{(n+2)!}{(2n)!} \left(\frac{1}{t^{n+2}} + 1 \right) x^{2n}$$

(1)

$$\sum_{n\geqslant 0} \frac{(n+2)!}{(2n)!} x^{2n} = \sum_{n\geqslant 0} \frac{1}{(n+3)(n+4)\cdots(2n)}$$

$$\leqslant \sum_{n\geqslant 0} \frac{1}{n^{n-2}} x^{2n}$$

$$\leqslant \sum_{n\geqslant M} + \sum_{n\geqslant M} \frac{1}{M^{n-2}} x^{2n}$$

$$M^2 \sum_{n} \left(\frac{x^2}{M}\right)^n$$

$$\leqslant m^2 \frac{1}{1 - \left(\frac{x^2}{M}\right)}$$

(2)
$$t \in \left[\epsilon, \frac{1}{\epsilon}\right]$$
, so $\frac{1}{t} \leqslant \frac{1}{\epsilon}$, so $\frac{1}{t^{n+2}} \leqslant \frac{1}{\epsilon^{n+2}} \longrightarrow \sum_{n \geqslant 0} \frac{(n+2)!}{(2n)!} \frac{1}{t^{n+2}} x^{2n} \leqslant \sum_{n \geqslant 0} \frac{1}{n^{n-2}} \frac{1}{\epsilon^{n-2}} x^{2n}$

Remark 6.13 $|u(x,t)| \leq \exp\left(\frac{cx^2}{t}\right) \leadsto$ unphysical solution. Violates $|u(x,t)| \leq Ce^{C|x|^2}$ for all $\forall (x,t) \in \mathbb{R} \times [0,T]$

Exercise 6.14 (Bonus 10) Consider

$$u(x,t) = \int_{\mathbb{R}^d} \Phi(x-y,t)g(y) \, dy$$

where $\Phi(x,t) = \frac{1}{(4\pi t)^{\frac{d}{2}}} e^{-\frac{x^2}{4t}}$. Assume $g \in C_c^{\infty}(\mathbb{R}^d)$. Prove or disprove that

$$||u(\cdot,t) - g||_{L^2(\mathbb{R}^d)} \leqslant C_n t^n$$

as $t \to 0^+$ for all $n = 1, 2, \dots$

6.2 Maximum Principle

Recall the Poisson equation $-\Delta u \leq 0$ in $\Omega \subseteq \mathbb{R}^d$ open, bounded. Then

$$\sup_{\bar{\Omega}} u(x) = \sup_{\partial \Omega} u(x).$$

Theorem 6.15 (Maximum principle for bounded sets) Let $\Omega \subseteq \mathbb{R}^d$ be open and bounded. Let T > 0 and define

$$\Omega_T = \Omega \times (0, T),$$

$$\partial^* \Omega_T = (\bar{\Omega} \times \{0\}) \cup (\partial \Omega \times [0, T])$$

If $u \in C_1^2(\Omega_T) \cap C(\bar{\Omega}_T)$ solves $\partial_t u - \Delta_x u \leq 0$ in Ω_T , then

$$\max_{\overline{\Omega_T}} u = \max_{\partial^{\star} \Omega_T} u.$$

Proof. We will use Hopf's argument which is simpler that the mean-value theorem (there exists a mean-value theorem for heat equation, but it is complicated and we will not discuss it). Firstly, to illustrate the principle, we proof the maximum principle for the Poisson Equation: Asumme $u \in C^2(\Omega) \cap C(\bar{\Omega})$

Step 1) Assume $\Delta u > 0$ in Ω . Since $\bar{\Omega}$ is compact, there is a $x_0 \in \bar{\Omega}$ s.t. $u(x_0) = \max_{x \in \bar{\Omega}} u(x)$. We prove that $x_0 \in \partial \Omega$. In fact, if $x_0 \in \Omega$, then since x_0 is a (local) maximizer of u in Ω , we have $\Delta u(x_0) \leq 0$, which contradicts to the assumption that $\Delta u > 0$ in Ω . Thus $x_0 \in \partial \Omega$, and hence

$$\max_{x \in \bar{\Omega}} u(x) = u(x_0) \leqslant \max_{x \in \partial \Omega} u(x).$$

Step 2) Now assume $\Delta u \ge 0$ in Ω . Define

$$u_{\epsilon}(x) = u(x) + \epsilon |x|^2, \quad \epsilon > 0.$$

Then, $\Delta u_{\epsilon} > 0$ in Ω , hence by Step 1 and

$$u \le u_{\epsilon} \le u + \epsilon \sup_{x \in \bar{\Omega}} |x|^2$$

we have

$$\begin{split} \max_{x \in \bar{\Omega}} u(x) &\leqslant \max_{x \in \bar{\Omega}} u_{\epsilon}(x) \leqslant \max_{x \in \partial \Omega} u_{\epsilon}(x) \\ &\leqslant \max_{x \in \partial \Omega} u(x) + \epsilon \left(\sup_{x \in \bar{\Omega}} |x|^2 \right) \xrightarrow{\epsilon \to 0} \max_{x \in \partial \Omega} u(x) \end{split}$$

Proof for the heat equation:

Step 1) Assume $u \in C_1^2(\Omega \times (0,T]) \cap C(\bar{\Omega} \times [0,T])$ and

$$\partial_t u - \Delta_x u < 0$$

in $\Omega \times (0,T]$. Then, because of compactness, there is $(x_0,t_0) \in \overline{\Omega} \times [0,T]$ s.t.

$$u(x_0, t_0) = \max_{(x,t)\in\bar{\Omega}\times[0,T]} u(x,t).$$

We prove that $(x_0, t_0) \in \partial^* \Omega_T$. Assume by contradiction that $(x_0, t_0) \notin \partial^* \Omega_T$, then $x_0 \in \Omega$ and $t_0 \in (0, T]$. Since $x \mapsto u(x, t_0)$ has a (local) maximizer $x_0 \in \Omega$ we have that $\Delta_x u(x_0, t_0) \leq 0$. Since $t \mapsto u(x_0, t)$ has a (local) maximizer $t_0 \in (0, T]$ we have that $\partial_t u(x_0, t_0) \geq 0$. This implies:

$$(\partial_t u - \Delta_x u)(x_0, t_0) \geqslant 0$$

which is a contradiction to the assumption. Thus $(x_0, t_0) \in \partial^* \Omega_T$, i.e. $\max_{\bar{\Omega}_T} u = \max_{\partial^* \Omega_T} u$.

Step 2) Assume $u \in C_1^2(\Omega \times (0,T)) \cap C(\bar{\Omega} \times [0,T])$ and

$$\partial_t u - \Delta_x u \leqslant 0 \quad \text{in } \Omega \times (0, T).$$

Let $\tilde{T} \in (0,T)$ and for $\epsilon > 0$:

$$u_{\epsilon}(x,t) = u(x,t) + \epsilon |x|^2.$$

Then: $u_{\epsilon} \in C_1^2(\Omega \times (0, T']) \cap C(\bar{\Omega} \times [0, \tilde{T}])$ and $\partial_t u_{\epsilon} - \Delta_x u_{\epsilon} < 0$ in $\Omega \times (0, \tilde{T}]$. By Step 1:

$$\max_{\bar{\Omega}_{\bar{T}}} u_{\epsilon} \leqslant \max_{\bar{\partial}^{\star} \Omega_{\bar{T}}} u_{\epsilon}$$

$$\stackrel{\epsilon \to 0}{\Rightarrow} \qquad \max_{\bar{\Omega}_{\bar{T}}} u \leqslant \max_{\bar{\partial}^{\star} \Omega_{\bar{T}}} u$$

$$\stackrel{\tilde{T} \to T}{\Rightarrow} \qquad \max_{\bar{\Omega}_{T}} u \leqslant \max_{\bar{\partial}^{\star} \Omega_{T}} u$$

Theorem 6.16 (Maximum principle for $\Omega = \mathbb{R}^d$) Let $\Omega_T = \mathbb{R}^d \times (0,T)$, $\bar{\Omega}_T = \mathbb{R}^d \times [0,T]$. Let $u \in C_1^2(\Omega_T) \cap C(\bar{\Omega}_T)$ such that

- $\partial_t u \Delta_x u \leq 0$ in Ω_T
- $u(x,t) \leq Me^{M|x|^2}$ for all $(x,t) \in \bar{\Omega}_T$

Then

$$\sup_{(x,t)\in\bar{\Omega}_T}u(x,t)=\sup_{x\in\mathbb{R}^d}u(x,0).$$

Proof.

Step 1: For all $y \in \mathbb{R}^d$ and $\epsilon > 0$ define

$$v(x,t) = u(x,t) - \frac{\epsilon}{(T+\epsilon-t)^{\frac{d}{2}}} \exp\left(\frac{|x-y|^2}{4(T+\epsilon-t)}\right)$$

This implies

$$\partial_t v - \Delta_x v = \partial_t u - \Delta_x u \leqslant 0$$

in Ω_T . For $U = B(y, r), U_T = U \times (0, T), \bar{U}_T = \bar{U} \times [0, T], \partial^* U_T = (U \times \{0\}) \cup (\partial U \times [0, T])$, by the maximum principle for U bounded we have

$$\max_{\bar{U}_T} v \leq \max_{\partial^* U_T} v$$

Let us bound $\max_{\partial^* U_T} v$.

• On $U \times \{0\}$ we use $v \leq u$ and hence

$$\max_{x \in \bar{U}} v(x,0) \leqslant \max_{x \in \bar{U}} u(x,0) \leqslant \max_{x \in \mathbb{R}^d} u(x,0).$$

• On $\partial U \times [0,T]$ we use $|x-y|=r \Rightarrow |x| \leq |y|+r$.

$$\begin{split} v(x,t) &= u(x,t) - \frac{\epsilon}{(T+\epsilon-t)^{\frac{d}{2}}} \exp\left(\frac{|x-y|^2}{4(T+\epsilon-t)}\right) \\ &\leqslant M e^{M(|y|+r)^2} - \frac{\epsilon}{(T+\epsilon)^{\frac{d}{2}}} \exp\left(\frac{r^2}{4(T+\epsilon)}\right) \xrightarrow{r\to\infty} -\infty \end{split}$$

if $M < \frac{1}{4(T+\epsilon)}$. In particular, we can choose r large s.t.

$$\max_{\substack{x \in \partial U \\ t \in [0,T]}} v(x,t) \leqslant \max_{x \in \mathbb{R}^d} u(x,0).$$

In summary, if $M < \frac{1}{4(T+\epsilon)}$, then:

$$u(y,t) - \frac{\epsilon}{(T+\epsilon-t)^{\frac{d}{2}}} = v(y,t) \leqslant \max_{\bar{U}_T} v \leqslant \max_{x \in \mathbb{R}^d} u(x,0)$$

This holds for all $(y,t) \in \mathbb{R}^d \times [0,T]$. Thus,

$$\max_{\mathbb{R}^d \times [0,T]} u \leqslant \frac{\epsilon}{(T+\epsilon-t)^{\frac{d}{2}}} + \max_{x \in \mathbb{R}^d} u(x,0)$$

Taking $\epsilon \to 0$ we conclude that if $M < \frac{1}{4T}$,

$$\max_{\mathbb{R}^d \times [0,T]} u \le \max_{x \in \mathbb{R}^d} u(x,0)$$

Step 2: For general T, we denote $T_1 = \frac{T}{N}, N \in \mathbb{N}$ s.t. $M < \frac{4}{T_1}$. Then by step 1:

$$\max_{\mathbb{R}^d \times [0,T_1]} u \leqslant \max_{x \in \mathbb{R}^d} u(x,0)$$

$$\max_{\mathbb{R}^d \times [T_1,2T_1]} u \leqslant \max_{x \in \mathbb{R}^d} u(x,T_1) \leqslant \max_{x \in \mathbb{R}^d} u(x,0)$$

$$\vdots$$

$$\max_{\mathbb{R}^d \times [(N-1)T_1,NT_1]} \leqslant \max_{x \in \mathbb{R}^d} u(x,(N-1)T_1) \leqslant \max_{x \in \mathbb{R}^d} u(x,0)$$

$$\Rightarrow \max_{\mathbb{R}^d \times [0,T]} u \leqslant \max_{x \in \mathbb{R}^d} u(x,0)$$

Remark 6.17 The condition $u \leq Me^{M|x|^2}$ is necessary, otherwise there are solutions $u \neq 0$ s.t. u(x,0) = 0

Theorem 6.18 (Uniqueness) If $u \in C_1^2(\mathbb{R}^d \times (0,T)) \cap C(\mathbb{R}^d \times [0,T])$ and

$$u(x,t) \leq Me^{M|x|^2}$$
 in $\mathbb{R}^d \times [0,T]$,
 $\partial_t u - \Delta_x u = 0$ in $\mathbb{R}^d \times (0,T)$,
 $u(x,0) = 0$ in \mathbb{R}^d

Then u = 0 in $\mathbb{R}^d \times [0, T]$.

Proof. Use the maximum principle for u and -u.

Remark 6.19 If $u(\cdot,t) \in L^2(\mathbb{R}^d)$, the proof of uniquness can be done without the maximum principle. Heuristically:

$$\frac{d}{dt} \int_{\mathbb{R}^d} |u(x,t)|^2 dt = 2 \int_{\mathbb{R}^d} (\partial_t u) u \, dx = 2 \int_{\mathbb{R}^d} \Delta_x u u \, dx = -2 \int_{\mathbb{R}^d} |\nabla_x u|^2 \, dx \le 0$$

This implies

$$e(t) := \int_{\mathbb{D}^d} |u(x,t)|^2 dx$$

is descreasing. Hence, if e(0) = 0, then e(t) = 0 for all $t \ge 0$. This argument will be helpful below for the heat backward equation.

Remark 6.20 The heat equation

$$\begin{cases} \partial_t u - \Delta_x u = 0 \\ u(t=0) = g \end{cases}$$

is a well-posed problem:

- Existence
- Uniqueness
- Stability (solution depends continuously on data)

For the latter issue, by the maximum principle we have

$$||u(\cdot,t)||_{L^{\infty}} \le ||u(\cdot,0)||_{L^{\infty}} \quad \forall t$$

or in the L^2 -situation:

$$||u(\cdot,t)||_{L^2} \le ||u(\cdot,0)||_{L^2} \quad \forall t$$

On the other hand, the heat backward equation

$$\begin{cases} \partial_t u - \Delta_x u = 0 \\ u(t = T) = g \end{cases}$$

is *not* well-posed.

- Non-Existence: In general, the existence requires some special property on g, e.g. g is very smooth (only $g \in C(\mathbb{R}^d) \cap L^{\infty}(\mathbb{R}^d)$ or $g \in L^2(\mathbb{R}^d)$ is not enough)
- Uniqueness: On the other hand, the uniqueness still holds.

Lemma 6.21 If $e \in C^2(0,T)$, $e(t) \ge 0$, $e'(t) \le 0$, $e''(t) \ge 0$ and $|e'(t)|^2 \le e(t)e''(t)$ for $t \in [0,T]$ and e(T) = 0, then $e \equiv 0$.

Proof. Since e is monotonly decreasing and e(T) = 0 there is a $t_0 \in [0,T]$ s.t. $e(t_0) = 0$ and e(t) > 0 if $t \le t_0$. We need to prove that $t_0 = 0$. Assume by contradiction $0 < t_0 \le T$, then for $t \in (0,t_0)$ define $f(t) := \log e(t)$. Then

$$f'(t) = \frac{e'(t)}{e(t)}$$

$$\Rightarrow f''(t) = \frac{e''(t)e(t) - |e'(t)|^2}{e(t)^2} \ge 0$$

This means that f is convex, so for all $t_1, t_2 \in (0, t_0)$ and $\tau \in (0, 1)$:

$$f(\tau t_1 + (1 - \tau)t_2) \le \tau f(t_1) + (1 - \tau)f(t_2)$$

$$\Rightarrow e(\tau t_1 + (1 - \tau)t_2) \le e(t_1)^{\tau} e(t_2)^{1 - \tau}$$

Now, $e(\tau t_1 + (1-\tau)t_2) \xrightarrow{t_2 \to t_0} 0$ and $\tau \to 1$ implies $e(t_1) = 0$ for all $t_1 \in (0, t_0)$ which is a contradiction.

Theorem 6.22 If $u \in C_1^2(\mathbb{R}^d \times [0,T]) \cap C^1(H^1(\mathbb{R}^d) \times [0,T])$ and

$$\begin{cases} \partial_t u - \Delta_x u = 0 & \text{in } \mathbb{R}^d \times (0, T) \\ u(x, T) = 0 \end{cases}$$

Then u = 0 in $\mathbb{R}^d \times [0, T]$.

Proof. Recall

$$e(t) = \int_{\mathbb{R}^d} |u(x,t)|^2 dx.$$

Then,

$$\begin{split} e'(t) &= 2 \int_{\mathbb{R}^d} u \partial_t u \, dx = 2 \int_{\mathbb{R}^d} u \Delta_x u \, dx = -2 \int_{\mathbb{R}^d} |\nabla_x u|^2 \, dx \\ e''(t) &= -4 \int_{\mathbb{R}^d} \nabla_x u \nabla_x (\partial_t u) = 4 \int_{\mathbb{R}^d} \Delta_x u \partial_t u \, dx = 4 \int_{\mathbb{R}^d} |\Delta_x u|^2 \, dx \geqslant 0 \end{split}$$

and hence

$$|e'(t)|^2 = 4 \left| \int_{\mathbb{R}^d} u \Delta_x u \, dx \right|^2 \le 4 \left(\int_{\mathbb{R}^d} |u|^2 \, dx \right) \left(\int_{\mathbb{R}^d} |\Delta_x u|^2 \, dx \right) = e(t)e''(t)$$

Then the statement follows with lemma 6.21.

Some remarks about the eat equation in unbounded domains:

$$\begin{cases} \partial_t u - \Delta_x u = 0 & \text{in } \mathbb{R}^d \times (0, \infty) \\ u(x, 0) = 0 & \text{(i.e. } \lim_{t \to 0} u(x, t) = 0 \forall x \in \mathbb{R}^d) \end{cases}$$

There is a classical solution $0 \neq u \in C^1(\mathbb{R}^d \times (0, \infty))$. An example is

$$u(x,t) = \sum_{n=0}^{\infty} \frac{g^{(n)}(t)}{(2n)!} x^{2n}, \quad g(t) = e^{-\frac{1}{t^2}}$$

(s.t. $g \to 0$ as $t \to 0$). Note

$$g(t) = e^{-\frac{1}{t^2}},$$

$$g'(t) = \frac{2}{t^3}g(t)$$

$$g''(t) = \left(\frac{2}{t^3}\right)'g(t) + \frac{2}{t^3}\frac{2}{t^3}g(t)$$

$$g^{(n)}(t) = P_n\left(\frac{1}{t}\right)g(t)$$

where

$$\begin{cases}
P_0 = 1 \\
P_{n+1} \left(\frac{1}{t}\right) = \left(P_n\left(\frac{1}{t}\right)\right)' + \left(\frac{2}{t^3}\right) P_n\left(\frac{1}{t}\right) = A_1 P_n + A_2 P_n, \\
P_{n+1} = (A_1 + A_2) P_n = (A_1 + A_2)(A_1 + A_2) P_{n-1} = A_1 P_n + A_2 P_n
\end{cases}$$

This implies:

$$P_n = (A_1 + A_2)^n P_0 = \sum_{\sigma \in \{1,2\}^n} A_{\sigma(1)} A_{\sigma(2)} \cdots A_{\sigma(n)} P_0$$
 (6.1)

$$A_1\left(\frac{\alpha}{t^s}\right) = \frac{-s\alpha}{t^{s+1}} \to A_1$$

Multiple coefficients by a factors and + power by 1

$$A_2\left(\frac{\alpha}{t^s}\right) = \frac{2\alpha}{t^{s+3}} \to A_2$$

Mul Cof by a factor 2 and + power by 3

$$\left| \underbrace{A_{\sigma(1)} \cdots A_{\sigma(n)}, 1}_{k \text{ times } A_2, \ n-k \text{ times } A_1} \right| \leqslant \frac{2^k}{t^{3k}} \leqslant \frac{2^k}{t^{3k}} \frac{(3n)^{n-k}}{t^{n-k}} = \frac{2^k (3n)^{n-k}}{t^{n+2k}}$$

This implies

$$|P_n\left(\frac{1}{t}\right)| \leqslant \max_{0 \leqslant k \leqslant n} \frac{2^n 2^k (3n)^{n-k}}{t^{n+2k}}$$

Thus:

$$\sum_{n} \left| \frac{g^{(n)}(t)}{(2n)!} x^{2n} \right| \leq \sum_{n} \max_{0 \leq k \leq n} \frac{2^{n} 2^{k} (3n)^{n-k}}{t^{n+2k} (2n)!} \frac{e^{-\frac{1}{t^{2}}}}{1} x^{2n}$$

$$\leq \sum_{n} \max \frac{2^{n} 2^{k} (3n)^{n-k}}{t^{n+2k} (2n)!} (k!) (2t^{2})^{k} e^{-\frac{1}{2t^{2}}} x^{2n}$$

$$= \sum_{n} \frac{2^{n} 2^{k} 2^{k} (3n)^{n-k} (k!)}{(2n)! t^{n}} e^{-\frac{1}{2t^{2}}} x^{2n}$$

$$\leq \sum_{n} \frac{(c_{n})^{n}}{(2n)! t^{n}} e^{-\frac{1}{2t^{2}}} x^{2n}$$

$$\leq \sum_{n} \frac{c^{n}}{n! t^{n}} e^{-\frac{1}{2t^{2}}} x^{2n}$$

$$\leq \sum_{n} \frac{c^{n}}{n! t^{n}} e^{-\frac{1}{2t^{2}}} x^{2n}$$

$$\leq \sum_{n} \frac{c^{n}}{n! t^{n}} e^{-\frac{1}{2t^{2}}} x^{2n}$$

Where we used that

$$e^s = \sum_k \frac{s^k}{k!} \geqslant \frac{s^k}{k!}$$

for all $s \ge 0$ implies

$$e^{-\frac{1}{2t^2}} = \frac{1}{e^{\frac{1}{2t^2}}} \le \frac{1}{\left(\frac{1}{2t^2}\right)\frac{1}{k!}} = k!(2t^2)^k.$$

We conclude:

- u(x,t) is well-defined, $x \in \mathbb{R}^d$, t > 0 real? to heat equation.
- $u(x,t) \to 0$ as $t \to 0$ for all $x \in \mathbb{R}^d$.

Exercise 6.23 (E 12.1) Let $\Omega \subseteq \mathbb{R}^d$ be open and $u \in C^2(\Omega)$. Assume that $x_0 \in \Omega$ is a local maximizer of u, namely there exists some r > 0 such that $u(x_0) \ge u(x)$ for all $x \in B_r(x_0) \subseteq \Omega$.

(a) Prove that the Hessian matrix $H = (D^{\alpha}u(x_0))_{|\alpha|=2}$ is negative semi-definite, namely

$$yHy \leq 0$$

for all $y \in \mathbb{R}^d$.

(b) Prove that $\Delta u(x_0) \leq 0$

Hint: Recall that we used (b) for the maximum principle by Hopf's method.

Solution. (a) In 1D this is obvious. If x_0 is a local minimizer of u, then $u'(x_0) = 0$, $u''(x_0) \leq 0$ (Taylor expansion).

In d dimensions:

$$\phi(t) = u(x_0 + t\xi) \quad \xi \in \mathbb{R}^d, t \in \mathbb{R}, |t| \text{ small}$$

So 0 is a local maximizer of ϕ . This implies

$$0 = \phi'(0) = \nabla u(x_0)\xi \quad \forall \xi \in \mathbb{R}^d \Rightarrow H \leqslant 0$$

$$\phi''(0) = \lim_{t \to 0} \frac{\phi'(t) - \phi'(0)}{t} = \lim_{t \to 0} \frac{(\nabla u(x_0 + t\xi) - \nabla u(x_0))\xi}{t}$$
$$= \lim_{t \to 0} \sum_{i=1}^d \frac{(\partial_i u(x_0 + t\xi) - \partial_i u(x_0))\xi_i}{t} = \sum_{i=1}^d \sum_{j=1}^d \partial_j \partial_i u(x_0)\xi_j \xi_i = \langle \xi, H\xi \rangle,$$

$$H = (\partial_i \partial_j u(x_0))_{i,j=1}^d.$$

(b) Consequently

$$\Delta u(x_0) = \sum_{i=1}^d \partial_i \partial_i u(x_0) = \text{Tr}(H) \le 0$$

Exercise 6.24 (E 12.2) Let $\Omega \subseteq \mathbb{R}^d$ be open and bouned. We prove the maximum principle for a general elliptic operator

$$Lu(x) = \sum_{i,j=1}^{d} a_{ij}(x)\partial_i\partial_j u(x) + \sum_{i=1}^{d} b_i(x)\partial_j u(x),$$

 $a_{ij}, b_i \in C(\bar{\Omega}), \ A(x) = (a_{ij}(x))_{i,j=1}^d \ge 1$ (as matrices). Prove that if $Lu(x) \ge 0$ for all $x \in \Omega$ and $u \in C^2(\Omega) \cap C(\bar{\Omega})$, then

$$\max_{x \in \bar{\Omega}} u(x) = \max_{x \in \partial \Omega} u(x).$$

Solution.

Step 1: Assume Lu(x) > 0 for all $x \in \Omega$: Since $u \in C(\bar{\Omega})$ there is a $x_0 \in \bar{\Omega}$ s.t.

$$u(x_0) = \max_{x \in \bar{\Omega}} u(x).$$

We prove $x_0 \in \partial \Omega$. Assume by contradiction that $x_0 \notin \partial \Omega$, so $x_0 \in \Omega$ is a local maximizer. We prove $Lu(x_0) \leq 0$. Note:

$$Lu(x_0) = \sum_{i,j=1}^{d} a_{ij}(x_0) \partial_i \partial_j u(x_0) + \sum_{i=1}^{d} b_i(x_0) \partial_i u(x_0)$$

$$= \text{Tr}[A(x_0)H(x_0)] + B(x_0) \underbrace{\nabla u(x_0)}_{=0} \leqslant 0 \quad \text{$\not \bot}$$

$$A(x_0) = (a_{ij}(x_0))_{i,j=1}^d$$
, $B(x_0) = (b_i(x_0))_{i=1}^d$, where $\text{Tr}[AH] = \sum_i (AH)_{ii} = \sum_i \sum_j A_{ij} H_{ij}$

General fact: If $A \ge 0, B \ge 0$ (matrices), then $Tr(AB) \ge 0$.

•
$$A = (\sqrt{A})^2 \Rightarrow \text{Tr}(AB) = \text{Tr}((\sqrt{A})^2 B) = \text{Tr}(\underbrace{\sqrt{A}B\sqrt{A}}_{>0}) \ge 0$$

• Spectral theorem: $A \ge 0$, then there are eigenvectors (α_i) and eigenvalues $\lambda_i \ge 0$ s.t.

$$\operatorname{Tr}(AB) = \sum_{i} \langle \alpha_i, AB\alpha_i \rangle = \sum_{i} \underbrace{\lambda_i}_{\geq 0} \underbrace{\langle \alpha_i, B\alpha_i \rangle}_{\geq 0} \geq 0$$

• General Case: $Lu(x) \ge 0$ for all $x \in \Omega$. Assume that there is a $v \in C^2(\Omega) \cap C(\bar{\Omega})$ s.t. Lv(x) > 0 for all $x \in \Omega$. Define for all $\epsilon > 0$ $u_{\epsilon} = u + \epsilon v$. Then $Lu_{\epsilon}(x) = Lu(x) + \epsilon Lv(x) > 0$ for all $x \in \Omega$. By Step 1,

$$\max_{x \in \bar{\Omega}} u_{\epsilon}(x) \leqslant \max_{x \in \partial \Omega} u_{\epsilon}(x)$$

$$\xrightarrow{\epsilon \to 0} \max_{x \in \bar{\Omega}} u(x) \leqslant \max_{x \in \partial \Omega} u(x)$$

What v? First $v(x) = x^2 = x_1^2 + \cdots + x_d^2$

$$Lv(x) = \sum_{ij} a_{ij}(x) 2\delta_{ij} + \sum_{i} b_{i}(x) 2x_{i}$$

not clear to be ≥ 0 .

$$v(x) = x^{2n} \quad n \text{ large}$$

$$v(x) = x_1^{2n} \longrightarrow Lv(x) = a_{11}(x)2n(2n+1)x_1^{2n-2} + b_1(x)2nx_1^{2n-1}$$

$$\geqslant 2nx_1^{2n-2}[(2n-1) + \underbrace{b_1(x)x_1}_{\text{b.d. in }\bar{\Omega}}] \geqslant 0 \quad \forall x \in \bar{\Omega}$$

if n is large enough.

$$v(x) = (x_1 + R)^{2n}$$

where R > 0 large s.t. $x_1 + R \ge 1$ for all $\forall x \in \bar{\Omega}$. This implies

$$Lv(x) \ge 2n\underbrace{(x_1+R)^{2n-2}}_{>0} [\underbrace{2n-1+b_1(x)(x_1+R)}_{>0}] > 0$$

for all $x \in \bar{\Omega}$ if n is large.

Exercise 6.25 (E 12.3) Consider the inhomogeneous heat equation

$$\begin{cases} \partial_t u - \Delta_x u = f(x, t) & \text{in } \mathbb{R}^d \times (0, T) \\ u(t = 0) = g & \text{in } \mathbb{R}^d \end{cases},$$

 $f \in C^2_1(\mathbb{R}^d \times (0,T))$ and compactly supported and $g \in C(\mathbb{R}^d \times [0,T]) \cap L^\infty(\mathbb{R}^d \times [0,T])$. Assume that there exists a solution $u \in C^2_1(\mathbb{R}^d \times (0,T)) \cap C(\mathbb{R}^d \times [0,T])$ satisfying

$$u(x,t) \leqslant Me^{M|x|^2}, \quad (x,t) \in \mathbb{R}^d \times [0,T].$$

Prove that

$$\max_{(x,t)\in\mathbb{R}^d\times[0,T]}|u(x,t)|\leqslant \|g\|_{L^\infty}+T\|f\|_{L^\infty}.$$

Solution.

Step 1: There is at most one solution u.

Step 2:

$$u(x,t) = \int_{\mathbb{R}^d} \phi(x - y, t) g(y) \, dy + \int_0^t \int_{\mathbb{R}^d} \phi(x - y, t - s) f(y, s) \, dy \, ds$$

This implies:

$$||u||_{L^{\infty}} \leqslant \int_{\mathbb{R}^{d}} \phi(x-y,t) ||g||_{L^{\infty}} dy + \int_{0}^{t} \int_{\mathbb{R}^{d}} \phi(x-y,t-s) ||f||_{L^{\infty}} dy ds$$

$$\Rightarrow ||u||_{L^{\infty}_{x,t}} \leqslant \int_{\mathbb{R}^{d}} \phi(x-y,t) ||g||_{L^{\infty}} dy + \int_{0}^{T} \int_{\mathbb{R}^{d}} \phi(x-y,t-s) ||f||_{L^{\infty}} dy ds$$

$$= ||g||_{L^{\infty}_{x,t}} + T||f||_{L^{\infty}_{x,t}}$$

This is optimal! E.g. g = 0, f = 1, u(x, t) = u(t).

$$\begin{cases} u' = 1 \\ u(0) = 0 \end{cases} \Rightarrow u(t) = t$$

Exercise 6.26 (Bonus 11) Denote for all $u \in C^2(\Omega) \cap C(\bar{\Omega})$:

$$Lu(x) = \sum_{i,j=1}^{d} a_{ij}(x)\partial_i\partial_j u(x)$$

where $a_{ij} \in C(\bar{\Omega})$ s.t. $A(x) = (a_{ij}(x)) \ge 1$. Prove that if $\Omega \subseteq \mathbb{R}^d$ is open and bounded, $u \in C_1^2(\bar{\Omega} \times [0,T])$ and

$$\begin{cases} \partial_t u - Lu \leqslant 0 & \text{in } \Omega \times (0, T) \\ u(t = 0) = 0 \\ u(x \in \partial \Omega) = 0 \end{cases}$$

Prove that $u(x,t) \leq 0$ for all $(x,t) \in \bar{\Omega} \times [0,T]$.

6.3 Backward heat equation

Theorem 6.27 (Instability)

There exist functions $u_{\epsilon} \in C_1^2(\mathbb{R}^d \times (0,T)) \cap C^1(H^1(\mathbb{R}^d) \times [0,T])$ s.t.

$$\partial_t u_{\epsilon} - \Delta_x u_{\epsilon} = 0 \quad \text{in } \mathbb{R}^d \times [0, T]$$

with:

$$\|u_{\epsilon}(\bullet,T)\|_{L^{2}(\mathbb{R}^{d})} \xrightarrow{\epsilon \to 0^{+}} 0, \quad \|u_{\epsilon}(\bullet,0)\|_{L^{2}(\mathbb{R}^{d})} \xrightarrow{\epsilon \to 0^{+}} \infty.$$

Proof. Recall by Fourier Transform

$$\partial_t \hat{u}(k,t) + |2\pi k|^2 \hat{u}(k,t) = 0$$

$$\Leftrightarrow \qquad \partial_t (e^{|2\pi k|^2 t} \hat{u}(k,t)) = 0$$

$$\Rightarrow \qquad e^{|2\pi k|^2 t} \hat{u}(k,t) = u(k,0)$$

$$\Rightarrow \qquad \hat{u}(k,t) = e^{-t|2\pi k|^2} \hat{u}(k,0)$$

$$\Rightarrow \qquad \hat{u}(k,0) = e^{T|2\pi k|^2} (k,T).$$

Now we can take

$$\hat{u}_{\epsilon}(k,t) = \mathbb{1}\left(|k| \leqslant \frac{1}{\epsilon}\right)\epsilon^{d+1} dk$$

Then,

$$\begin{split} \|u(\bullet,T)\|_{L^2(\mathbb{R}^d)}^2 &= \int_{\mathbb{R}^d} \hat{u}_{\epsilon}(k,t) \, dk = \lambda^d (\{|k| \leqslant \epsilon^{-1}\}) \epsilon^{d+1} \sim \epsilon \xrightarrow{\epsilon \to 0} 0 \\ \|u(\bullet,0)\|_{L^2(\mathbb{R}^d)}^2 &= \int_{\mathbb{R}^d} e^{2T|2\pi k|^2} \mathbb{1}(|k| \leqslant \epsilon^{-1}) \epsilon^{d+1} \, dk \\ &\geqslant \int_{\frac{\epsilon}{2} \leqslant |k| \leqslant \frac{\epsilon}{2}} e^{2T|2\pi k|^2} \mathbb{1}(|k| \leqslant \epsilon^{-1}) \epsilon^{d+1} \, dk \tilde{\geqslant} e^{2T\epsilon^{-2}} \epsilon \xrightarrow{\epsilon \to 0} \infty \end{split}$$

Remark 6.28 This means that a small error of the data at t = T may cause a large error of the output t = 0.

Theorem 6.29 (Regularized solution)

Assume that $u \in C_1^2(\mathbb{R}^d \times (0,T)) \cap C^1(H^1(\mathbb{R}^d),[0,T])$

$$\begin{cases} \partial_t u - \Delta_x u = 0 & \text{in } \mathbb{R}^d \times (0, T) \\ u(x, T) = g(x) & \text{in } \mathbb{R}^d \end{cases}$$

Then from given data $g_{\epsilon} \in L^2(\mathbb{R}^d)$ s.t.

$$||g_{\epsilon} - g||_{L^2(\mathbb{R}^d)} \le \epsilon$$

we construct a solution \tilde{u}_{ϵ} s.t.

$$\sup_{t \in [0,T]} \|\tilde{u}_{\epsilon}(\bullet,t) - u(\bullet,t)\|_{L^{2}(\mathbb{R}^{d})} \xrightarrow{\epsilon \to 0} 0$$

Proof. Clearly we should not choose \tilde{u}_{ϵ} to solve

$$\begin{cases} \partial_t u_{\epsilon} - \Delta_x u_{\epsilon} = 0 \\ u_{\epsilon}(t = T) = g_{\epsilon} \end{cases},$$

i.e.

$$\hat{u}_{\epsilon}(k,t) = e^{(T-t)|2\pi k|^2} \hat{g}_{\epsilon}(k).$$

Rather we take

$$\hat{u}_{\epsilon}(k,t) = e^{(T-t)|2\pi k|^2} \hat{g}_{\epsilon}(k) \mathbb{1}(|k| \leq \delta_{\epsilon}^{-1})$$

Where $\delta_{\epsilon} \to 0$ (chosen later). Then we have for all $t \in [0, T]$:

$$\begin{split} \|u_{\epsilon}(\bullet,t) - u(\bullet,t)\|_{L^{2}(\mathbb{R}^{d})}^{2} &= \int_{\mathbb{R}^{d}} e^{2(T-t)|2\pi k|^{2}} |\hat{g}_{\epsilon}(k)\mathbb{1}(|k| \leqslant \delta_{\epsilon}^{-1}) - \hat{g}(k)|^{2} \, dk \\ &\leqslant 2 \int_{\mathbb{R}^{d}} e^{2T|2\pi k|^{2}} |\hat{g}_{\epsilon}(k) - \hat{g}(k)|\mathbb{1}(|k| \leqslant \delta_{\epsilon}^{-1}) \, dk \\ &+ 2 \int_{\mathbb{R}^{d}} \underbrace{e^{2T|2\pi k|^{2}} |\hat{g}(k)|^{2}}_{|\hat{u}(k,0)|^{2}} \mathbb{1}(|k| > \delta_{\epsilon}^{-1}) \, dk = (\mathrm{I}) + (\mathrm{II}) \end{split}$$

We have

$$\begin{split} &(\mathrm{I})\leqslant 2\int_{\mathbb{R}^d}e^{c\delta_{\epsilon}^{-2}}|\hat{g}_{\epsilon}(k)-\hat{g}(k)|^2\,dk=2e^{c\delta_{\epsilon}^{-2}}\epsilon^{-2}\longrightarrow 0 & \text{if }\delta_{\epsilon}\gg \frac{1}{\sqrt{|\log\epsilon|}}\\ &(\mathrm{II})=2\int_{\mathbb{R}^d}|\hat{u}(k,0)|^2\mathbb{1}(|k|\geqslant \delta_{\epsilon}^{-1})\,dk\leqslant 2\int_{\mathbb{R}^d}|k|^2\delta_{\epsilon}^2|\hat{u}(k,0)|^2\,dk \end{split}$$

Thus chosing $\frac{1}{\sqrt{|\log \epsilon|}} \ll \delta_{\epsilon} \ll 1$, e.g. $\delta_{\epsilon} = (|\log \epsilon|)^{-\frac{1}{4}}$.

$$\sup_{t \in [0,T]} \|u_{\epsilon}(\bullet,t) - u(\bullet,t)\|_{L^{2}(\mathbb{R}^{d})} \leq (\mathrm{I}) + (\mathrm{II}) \xrightarrow{\epsilon \to 0} 0$$

Remark 6.30 In application, both u and g are unknown. Only g_{ϵ} is given. So we have to construct \tilde{u}_{ϵ} using only information from g_{ϵ} .

Chapter 7

Wave Equation

7.1 d'Alembert

Wave equation:

$$\begin{cases} \partial_t^2 u - \Delta_x u = 0 & x \in \mathbb{R}^d, t > 0 \\ u = g, \partial_t u = h & x \in \mathbb{R}^d, t = 0 \end{cases}$$

In d = 1:

$$\begin{cases} \partial_t^2 u - \partial_x^2 u = 0, & (x,t) \in \mathbb{R} \times (0,\infty) \\ u = g, \partial_t u = h, & x \in \mathbb{R}, t = 0 \end{cases}$$

Key idea: Factorization:

$$\partial_t^2 - \partial_x^2 = (\partial_t + \partial_x)(\partial_t - \partial_x).$$

Then, if we denote $v = (\partial_t - \partial_x)u$, we get the transport equation

$$(\partial_t + \partial_x)v = 0.$$

This implies

$$v(x,t) = a(x-t), \quad a(x) = v(x,0)$$

From this we get the inhomogeneous transport equation

$$(\partial_t - \partial_x)u = a(x - t).$$

Now we decompose $u = u_1 + u_2$, where

$$\begin{cases} (\partial_t - \partial_x)u_1 = 0\\ (\partial_t - \partial_x)u_2 = a(x - t) \end{cases}.$$

Like above, we get $u_1 = b(x+t)$ and an explicit choice of u_2 is

$$u_2(x,t) = \frac{1}{2} \int_{x-t}^{x+t} a(y) \, dy$$

Thus,

$$u(x,t) = b(x+t) + \frac{1}{2} \int_{x-t}^{x+t} a(y) \, dy$$

Let's compute a and b:

$$b(x) = u(x,0) = g(x)$$

$$a(x) = v(x,0) = (\partial_t u - \partial_x u)_{t=0} = h - g'.$$

Theorem 7.1 (d'Alembert formula) For d=1 let $g\in C^2(\mathbb{R}^d)$, $h\in C^1(\mathbb{R})$ and define u by the d'Alembert formula

$$u(x,t) = \int_{x-t}^{x+t} (h(y) - g'(y)) \, dy + g(x+t)$$
$$= \frac{1}{2} [g(x+t) + g(x-t)] + \frac{1}{2} \int_{x-t}^{x+t} h(y) \, dy.$$

Then:

- $u \in C^2(\mathbb{R} \times (0, \infty))$
- $\bullet \ \partial_t^2 u \partial_r^2 u = 0$
- $u = g, \partial_t u = h \text{ when } t \to 0$

Proof. Exercise.

Remark 7.2 If $g \in C^k$ and $h \in C^{k-1}$, then $u \in C^k$ (but not better).

Now, let's apply the *Reflection Method*. Replace \mathbb{R} by $\mathbb{R}_+ = (0, \infty)$ and assume

$$\begin{cases} \partial_t^2 u - \partial_x^2 u = 0 & \mathbb{R}_+ \times (0, \infty) \\ u = g, \partial_t u = h & \text{on } \mathbb{R}_+ \times \{t = 0\}, g(0) = h(0) = 0 \\ u = 0 & \text{on } \{x = 0\} \times \{t > 0\} \end{cases}$$

Define

$$\tilde{u}(x,t) = \begin{cases} u(x,t), & x \geqslant 0, t \geqslant 0 \\ -u(-x,t), & x \leqslant 0, t \geqslant 0 \end{cases}$$
$$\tilde{g}(x) = \begin{cases} g(x) & x \geqslant 0 \\ -g(-x) & x \leqslant 0 \end{cases}$$
$$\tilde{h}(x) = \begin{cases} h(x) & x \geqslant 0 \\ -h(-x) & h \leqslant 0 \end{cases}$$

Then

$$\begin{cases} \partial_t^2 \tilde{u} - \partial_x^2 \tilde{u} = 0 & \text{in } \mathbb{R} \times (0, \infty) \\ \tilde{u} = \tilde{g}, \partial_t \tilde{u} = \tilde{h} & \text{on } \mathbb{R} \times \{t = 0\} \end{cases}.$$

By d'Alembert formula

$$\tilde{u}(x,t) = \frac{1}{2} \left[\tilde{g}(x+t) + \tilde{g}(x-t) \right] + \frac{1}{2} \int_{x-t}^{x+t} \tilde{h}(y) \, dy$$

This imples

$$u(x,t) = \begin{cases} \frac{1}{2} [g(x+t) + g(x-t)] + \frac{1}{2} \int_{x-t}^{x+t} h(y) \, dy & x \geqslant t \geqslant 0 \\ \frac{1}{2} [g(x+t) - g(t-x)] + \frac{1}{2} \int_{t-x}^{x+t} h(y) \, dy & t \geqslant x \geqslant 0 \end{cases}.$$

This is the solution of the heat equation in $\mathbb{R}_+ \times (0, \infty)$.

7.2 Euler-Poisson-Darboux

$$(\star) \begin{cases} \partial_t^2 u - \Delta_x u = 0 & \text{in } \mathbb{R}^d \times (0, \infty) \\ u = g, \partial_t u = h & \mathbb{R}^d \times \{t = 0\} \end{cases}$$

Idea: Averaging of u over sphere \leadsto 1D problem. Define for $x \in \mathbb{R}^d, t > 0, r > 0$,

$$U_r(x,t) := \int_{\partial B(x,r)} u(y,t) \, dS(y)$$
$$G_r(x) := \int_{\partial B(x,r)} g(y) \, dS(y)$$
$$H_r(x) := \int_{\partial B(x,r)} h(y) \, dS(y)$$

Lemma 7.3 (Euler-Poisson-Darboux equation) If $u \in C^2(\mathbb{R}^d \times [0, \infty))$ solves (\star) , then for all $x \in \mathbb{R}^d$:

$$\bullet \ (r,t) \mapsto U \in C^2([0,\infty) \times [0,\infty)])$$

$$\bullet \begin{cases} \partial_t^2 U - \partial_r^2 U - \frac{d-1}{r} \partial_r U = 0 & \text{in } \mathbb{R}_+ \times \mathbb{R}_+ \\ U = G, \partial_t U = H & \text{on } \mathbb{R}_+ \times \{t = 0\} \end{cases}$$

Note that $\partial_r^2 + \frac{d-1}{r}\partial_r$ is the radial part of Δ .

Proof. We compute for r > 0:

$$\begin{split} \partial_r U_r(x,t) &= \partial_r \int_{\partial B(x,r)} u(y,t) \, dS(y) \\ &= \partial_r \int_{\partial B(0,1)} u(x+rz,t) \, dS(z) \\ &= \int_{\partial B(0,1)} \nabla u(x+rz,t) z \, dS(z) \\ &= \int_{\partial B(x,r)} \nabla u(y,t) \frac{y-x}{r} \, dS(y) \\ &= \int_{\partial B(x,r)} \frac{\partial u(y,t)}{\partial \vec{n}} \, dS(y) \\ &= \int_{\partial B(x,r)} \int_{B(x,r)} \Delta_x u(y,t) \, dy \\ \left(|B(0,r)| = \frac{r}{d} |\partial B(0,r)| \right) &= \frac{r}{d} \int_{B(x,r)} \Delta_x u(y,t) \, dy \end{split}$$

(The computation is similar to the proof of the mean-value theorem for the Poisson equation.) We compute the second derivative

$$\begin{split} \partial_r^2 U_r(x,t) &= \partial_r \left[\frac{r}{d} \oint_{B(x,r)} \Delta_x u(y,t) \, dy \right] \\ \left(|B(0,r)| &= r^d |B(0,1)| \right) &= \partial_r \left[\frac{1}{d|B_1|r^{d-1}} \int_{B(x,r)} \Delta_x u(y,t) \, dy \right] \end{split}$$

Now,
$$\partial_r \frac{1}{d|B_1|r^{d-1}} = \frac{-d+1}{d|B_1|r^d} = -\frac{d-1}{d|B(0,r)|}$$
 and
$$\partial_r \int_{B(x,r)} \Delta_x u(y,t) \, dS(y) = \partial_r \int_{B(0,r)} \Delta_x u(x+ry,t) \, dy$$

$$(\text{Green 2.3}) = \partial_r \int_{\partial B(0,1)} \nabla_x u(x+ry,t) \frac{y}{|y|} \, dS(y)$$

$$(|y|=1) = \int_{\partial B(0,1)} \partial_r \nabla_x u(x+ry,t) y \, dS(y)$$

$$= \int_{\partial B(0,1)} \Delta_x u(x+ry,t) y \cdot y \, dS(y)$$

$$(y \cdot y = |y|^2 = 1) = \int_{\partial B(0,1)} \Delta_x u(x+ry,t) \, dS(y)$$

$$= \int_{\partial B(0,1)} \Delta_x u(y,t) \, dS(y)$$

Now with the product rule we get

$$\partial_r^2 U_r(x,t) = -\left(\frac{d-1}{d}\right) \oint_{B(x,r)} \Delta_x u(y,t) \, dy$$

$$+ \frac{1}{d|B_1|r^{d-1}} \oint_{\partial B(x,r)} \Delta_x u(y,t) \, dS(y)$$

$$= -\left(\frac{d-1}{d}\right) \oint_{B(x,r)} \Delta_x u(y,t) \, dy$$

$$+ \oint_{\partial B(x,r)} \Delta_x u \, dS(y)$$

And, since u is a solution,

$$\partial_t^2 U = \partial_t^2 \oint_{\partial B(x,r)} u \, dS(y) = \oint_{\partial B(x,r)} (\partial_t^2 u) \, dS(y) = \oint_{\partial B(x,r)} (\Delta_x u) \, dS(y).$$

So we can conclude

$$\partial_t^2 U - \partial_r^2 U - \frac{d-1}{d}U = 0$$

the above computation also shows that $U \in C^2(\mathbb{R}_+ \times [0, \infty))$. Moreover

$$\partial_r U_r(x,t) \xrightarrow{r \to 0^+} 0$$

$$\partial_r^2 U_r(x,t) \xrightarrow{r \to 0^+} \left(\frac{1}{d} - 1\right) \Delta_x u + \Delta_x u = \frac{1}{d} \Delta_x u$$

This implies that $U \in C^2([0,\infty) \times [0,\infty))$. Finally, when t=0,

$$\begin{cases} u = g \\ \partial_t = h \end{cases} \Rightarrow \begin{cases} U = G \\ \partial_t U = H \end{cases}.$$

We showed that it is a neccessary condition for u to solve the Euler-Poisson-Darboux equation. Now we try to actually solve the equation. In general, this is easier for odd d than for even d. We will consider the cases d = 2, 3.

7.2.1 Solution in three dimensions

Now, for r > 0 let $\tilde{U} = rU, \tilde{G} = rG, \tilde{H} = rH$. Then

$$\begin{cases} \partial_t^2 \tilde{U} - \partial_r^2 \tilde{U} = 0 & \text{in } \mathbb{R}_+ \times (0, \infty) \\ \tilde{U} = \tilde{G}, \partial_t \tilde{U} = \tilde{H} & \text{when } t = 0 \\ \tilde{U} = 0 & \text{when } r = 0 \end{cases}$$

Then, by d'Alembert's formula, for $0 \le r \le t$ we have

$$\begin{split} \tilde{U}_r(x,t) &= \frac{1}{2} \left[\tilde{G}(r+t) - \tilde{G}(t-r) \right] + \frac{1}{2} \int_{t-r}^{t+r} \tilde{H}(y) \, dy. \\ \Rightarrow \quad U_r(x,t) &= \frac{1}{2} \left[\frac{\tilde{G}(r+t) - \tilde{G}(t-r)}{r} \right] + \frac{1}{2r} \int_{t-r}^{t+r} \tilde{H}(y) \, dy \end{split}$$

Now, taking $r \to 0$ we get

$$\begin{split} u(x,t) &= \tilde{G}'(t) + \tilde{H}(t) \\ &= \partial_t \left(t \oint_{\partial B(x,t)} g(y) \, dS(y) \right) + t \oint_{\partial B(x,t)} h(y) \, dS(y) \end{split}$$

Using

$$\oint_{\partial B(x,t)} g(y) \, dS(y) = \oint_{\partial B(0,1)} g(x+tz) \, dS(z)$$

we get

$$\partial_t \oint_{\partial B(x,t)} g(y) \, dS(y) = \oint_{\partial B(0,1)} \nabla g(x+tz) z \, dz$$

$$= \oint_{\partial B(x,t)} \nabla g(y) \left(\frac{y-x}{t}\right) \, dS(y)$$

$$\Rightarrow \quad \partial_t \left(t \oint_{\partial B(x,t)} g(y) \, dS(y)\right) = \oint_{\partial B(x,t)} \left(g + \nabla g(y-x)\right) \, dS(y)$$

From that we get:

Remark 7.4 (Kirchhoff's formula in 3D) For all $x \in \mathbb{R}^3$, t > 0:

$$u(x,t) = \int_{\partial B(x,t)} (g(y) + \nabla g(y-x) + th(y)) \ dS(y)$$