A Typicality-Based Algorithm for Partial Trace Estimation

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Application: Quantum Equilibrium Thermodynamics

Total system (with Hamiltonian \mathbf{H}) in thermal equilibrium at inverse temperature β (due to weak coupling with a "superbath") has density matrix:

$$\boldsymbol{\rho}_{\mathrm{t}}(\beta) = \frac{\exp(-\beta \mathbf{H})}{Z_{\mathrm{t}}(\beta)}, \qquad Z_{\mathrm{t}}(\beta) = \operatorname{tr}(\exp(-\beta \mathbf{H}).$$
(1)

The subsystem of interest has density matrix:

$$\boldsymbol{\rho}^*(\beta) = \operatorname{tr}_{b}(\boldsymbol{\rho}_{t}(\beta)). \tag{2}$$

Partial Traces

Let d_s and d_b be the dimension of \mathcal{H}_s and \mathcal{H}_b respectively, so that $d_t = d_s d_b$ is the dimension of $\mathcal{H}_t = \mathcal{H}_s \otimes \mathcal{H}_b$. A generic matrix $\mathbf{A} : \mathcal{H}_t \to \mathcal{H}_t$ can be partitioned as

$$\mathbf{A} = \begin{bmatrix} \mathbf{A}_{1,1} & \mathbf{A}_{1,2} & \cdots & \mathbf{A}_{1,d_s} \\ \mathbf{A}_{2,1} & \mathbf{A}_{2,2} & \cdots & \mathbf{A}_{2,d_s} \\ \vdots & \vdots & \ddots & \vdots \\ \mathbf{A}_{d_s,1} & \mathbf{A}_{d_s,2} & \cdots & \mathbf{A}_{d_s,d_s} \end{bmatrix},$$
(3)

where $\mathbf{A}_{i,j}:\mathcal{H}_{\mathrm{b}}\to\mathcal{H}_{\mathrm{b}}$ for each i,j. The partial trace of \mathbf{A} over \mathcal{H}_{b} is defined as

$$\operatorname{tr}_{b}(\mathbf{A}) := \begin{bmatrix} \operatorname{tr}(\mathbf{A}_{1,1}) & \operatorname{tr}(\mathbf{A}_{1,2}) & \cdots & \operatorname{tr}(\mathbf{A}_{1,d_{s}}) \\ \operatorname{tr}(\mathbf{A}_{2,1}) & \operatorname{tr}(\mathbf{A}_{2,2}) & \cdots & \operatorname{tr}(\mathbf{A}_{2,d_{s}}) \\ \vdots & \vdots & \ddots & \vdots \\ \operatorname{tr}(\mathbf{A}_{d_{s},1}) & \operatorname{tr}(\mathbf{A}_{d_{s},2}) & \cdots & \operatorname{tr}(\mathbf{A}_{d_{s},d_{s}}) \end{bmatrix}. \tag{4}$$

Stochastic (Partial) Trace Estimation

Suppose ${\bf v}$ has independent and identically distributed standard normal entries. Then,

$$\mathbb{E}[\mathbf{v}^\mathsf{T}\mathbf{M}\mathbf{v}] = \mathrm{tr}(\mathbf{M}), \qquad \mathbb{V}[\mathbf{v}^\mathsf{T}\mathbf{M}\mathbf{v}] = 2\|\mathbf{M}\|_\mathsf{F}^2.$$

Such estimators can be extended to partial traces [1]. In particular,

$$(\mathbf{I}_{d_{s}} \otimes \mathbf{v})^{\mathsf{T}} \mathbf{A} (\mathbf{I}_{d_{s}} \otimes \mathbf{v}) = \begin{bmatrix} \mathbf{v}^{\mathsf{T}} \mathbf{A}_{1,1} \mathbf{v} & \mathbf{v}^{\mathsf{T}} \mathbf{A}_{1,2} \mathbf{v} & \cdots & \mathbf{v}^{\mathsf{T}} \mathbf{A}_{1,d_{s}} \mathbf{v} \\ \mathbf{v}^{\mathsf{T}} \mathbf{A}_{2,1} \mathbf{v} & \mathbf{v}^{\mathsf{T}} \mathbf{A}_{2,2} \mathbf{v} & \cdots & \mathbf{v}^{\mathsf{T}} \mathbf{A}_{2,d_{s}} \mathbf{v} \\ \vdots & \vdots & \ddots & \vdots \\ \mathbf{v}^{\mathsf{T}} \mathbf{A}_{d_{s},1} \mathbf{v} & \mathbf{v}^{\mathsf{T}} \mathbf{A}_{d_{s},2} \mathbf{v} & \cdots & \mathbf{v}^{\mathsf{T}} \mathbf{A}_{d_{s},d_{s}} \mathbf{v} \end{bmatrix}$$

$$(5)$$

Given independent and identically distributed copies $\mathbf{v}_1, \dots, \mathbf{v}_m$ of \mathbf{v} , we arrive at an estimator

$$\widehat{\mathsf{tr}}_{\mathrm{b}}^{m}(\mathbf{A}) := \frac{1}{m} \sum_{i=1}^{m} (\mathbf{I}_{d_{\mathrm{s}}} \otimes \mathbf{v}_{i})^{\mathsf{T}} \mathbf{A} (\mathbf{I}_{d_{\mathrm{s}}} \otimes \mathbf{v}_{i}). \tag{6}$$

The variance of a random matrix \mathbf{X} can be defined as $\mathbb{V}[\mathbf{X}] := \mathbb{E}[\|\mathbf{X} - \mathbb{E}[\mathbf{X}]\|_{\mathsf{F}}^2]$. For our estimator:

$$\mathbb{V}\left[\widehat{\mathsf{tr}}_{b}^{m}(\mathbf{A})\right] = \frac{1}{m}\mathbb{V}\left[(\mathbf{I}_{d_{s}} \otimes \mathbf{v})^{\mathsf{T}} \mathbf{A} (\mathbf{I}_{d_{s}} \otimes \mathbf{v})\right] = \frac{1}{m} \sum_{i,j=1}^{d_{s}} 2\|\mathbf{A}_{i,j}\|_{\mathsf{F}}^{2} = \frac{2}{m}\|\mathbf{A}\|_{\mathsf{F}}^{2}. \tag{7}$$

A Variance-Reduced Algorithm

The linearity of partial trace implies that, for any matrix $\hat{\mathbf{A}}$,

$$\operatorname{tr}_{b}(\mathbf{A}) = \operatorname{tr}_{b}(\widetilde{\mathbf{A}}) + \operatorname{tr}_{b}(\mathbf{A} - \widetilde{\mathbf{A}}).$$
 (8)

We can then estimate the partial trace of $\mathrm{tr_b}(\mathbf{A})$ by computing the partial trace of the first term exactly, and applying the randomized estimator (5) to the residual term:

$$\operatorname{tr}_{\mathrm{b}}(\mathbf{A}) \approx \operatorname{tr}_{\mathrm{b}}(\widetilde{\mathbf{A}}) + \widehat{\operatorname{tr}}_{\mathrm{b}}^{m}(\mathbf{A} - \widetilde{\mathbf{A}}).$$
 (9)

The variance of such an estimate is entirely due to the variance of $\widehat{\mathbf{tr}}_{b}^{m}(\mathbf{A} - \widetilde{\mathbf{A}})$. If $\|\mathbf{A} - \widetilde{\mathbf{A}}\|_{\mathsf{F}}^{2} < \|\mathbf{A}\|_{\mathsf{F}}^{2}$, then the variance of the estimator on the right of (9) is reduced over that of $\widehat{\mathbf{tr}}_{b}^{m}(\mathbf{A})$.

Splittings similar to (9) have previously been used as a variance reduction technique for regular trace estimation [2, 3, etc.].

Using the fact that $\operatorname{tr}(\mathbf{x}_{(i)}\mathbf{x}_{(i)}^{\mathsf{T}}) = \mathbf{x}_{(i)}^{\mathsf{T}}\mathbf{x}_{(j)}$, we find that

$$\operatorname{tr}_{b}(\mathbf{x}\mathbf{x}^{\mathsf{T}}) = \begin{bmatrix} \mathbf{x}_{(1)}^{\mathsf{T}}\mathbf{x}_{(1)} & \mathbf{x}_{(2)}^{\mathsf{T}}\mathbf{x}_{(1)} & \cdots & \mathbf{x}_{(d_{s})}^{\mathsf{T}}\mathbf{x}_{(1)} \\ \mathbf{x}_{(1)}^{\mathsf{T}}\mathbf{x}_{(2)} & \mathbf{x}_{(2)}^{\mathsf{T}}\mathbf{x}_{(2)} & \cdots & \mathbf{x}_{(d_{s})}^{\mathsf{T}}\mathbf{x}_{(2)} \\ \vdots & \vdots & \ddots & \vdots \\ \mathbf{x}_{(1)}^{\mathsf{T}}\mathbf{x}_{(d_{s})} & \mathbf{x}_{(2)}^{\mathsf{T}}\mathbf{x}_{(d_{s})} & \cdots & \mathbf{x}_{(d_{s})}^{\mathsf{T}}\mathbf{x}_{(d_{s})} \end{bmatrix}.$$
(10)

This allows us to take $\widetilde{\mathbf{A}} = \mathbf{Q}\mathbf{Q}^{\mathsf{T}}\mathbf{A}\mathbf{Q}\mathbf{Q}^{\mathsf{T}} = \sum_{i=1}^{k} \theta_{i}\mathbf{x}_{i}\mathbf{x}_{i}^{\mathsf{T}}$ as a low-rank approximation of \mathbf{A} which gives the approximation,

$$\widehat{\mathsf{tr}}_{\mathrm{b}}^{m}(\mathbf{A}; \mathbf{Q}) := \sum_{i=1}^{k} \theta_{i} \operatorname{tr}_{\mathrm{b}}(\mathbf{x}_{i} \mathbf{x}_{i}^{\mathsf{T}}) + \widehat{\mathsf{tr}}_{\mathrm{b}}^{m}(\mathbf{A} - \mathbf{Q} \mathbf{Q}^{\mathsf{T}} \mathbf{A} \mathbf{Q} \mathbf{Q}^{\mathsf{T}})$$
(11)

We take \mathbf{Q} as the top k eigenvectors of \mathbf{A} , which minimizes the variance of the algorithm and helps with numerical stability in the case that products with \mathbf{A} are inexact, such as if products with $f(\mathbf{H})$ are approximated using the Lanczos algorithm.

In this case, we have:

$$\|\mathbf{A} - \mathbf{Q}\mathbf{Q}^{\mathsf{T}}\mathbf{A}\mathbf{Q}\mathbf{Q}^{\mathsf{T}}\|_{\mathsf{F}}^{2} = \min_{\substack{\text{rank}(\widetilde{\mathbf{A}})=k}} \|\mathbf{A} - \widetilde{\mathbf{A}}\|_{\mathsf{F}}^{2} = \sum_{i=k+1}^{d_{\mathsf{t}}} \sigma_{i}^{2}, \tag{12}$$

where $\{\sigma_i\}$ are the singular values of \mathbf{A} in non-increasing order.

Numerical Experiments

Our experiments focus on Heisenberg spin systems in an isotropic magnetic field oriented with the z-axis

$$\mathbf{H} := \sum_{i,j=1}^{N} \left[J_{i,j}^{\mathsf{x}} \boldsymbol{\sigma}_{i}^{\mathsf{x}} \boldsymbol{\sigma}_{j}^{\mathsf{x}} + J_{i,j}^{\mathsf{y}} \boldsymbol{\sigma}_{i}^{\mathsf{y}} \boldsymbol{\sigma}_{j}^{\mathsf{y}} + J_{i,j}^{\mathsf{z}} \boldsymbol{\sigma}_{i}^{\mathsf{z}} \boldsymbol{\sigma}_{j}^{\mathsf{z}} \right] + \frac{h}{2} \sum_{i=1}^{N} \boldsymbol{\sigma}_{i}^{\mathsf{z}}.$$
(13)

where $\sigma_i^{\text{x/y/z}} = \mathbf{I}_{i-1} \otimes \sigma^{\text{x/y/z}} \otimes \mathbf{I}_{N-i}$ and $\sigma^{\text{x/y/z}}$ are the Pauli matrices

Numerical Experiments cont.

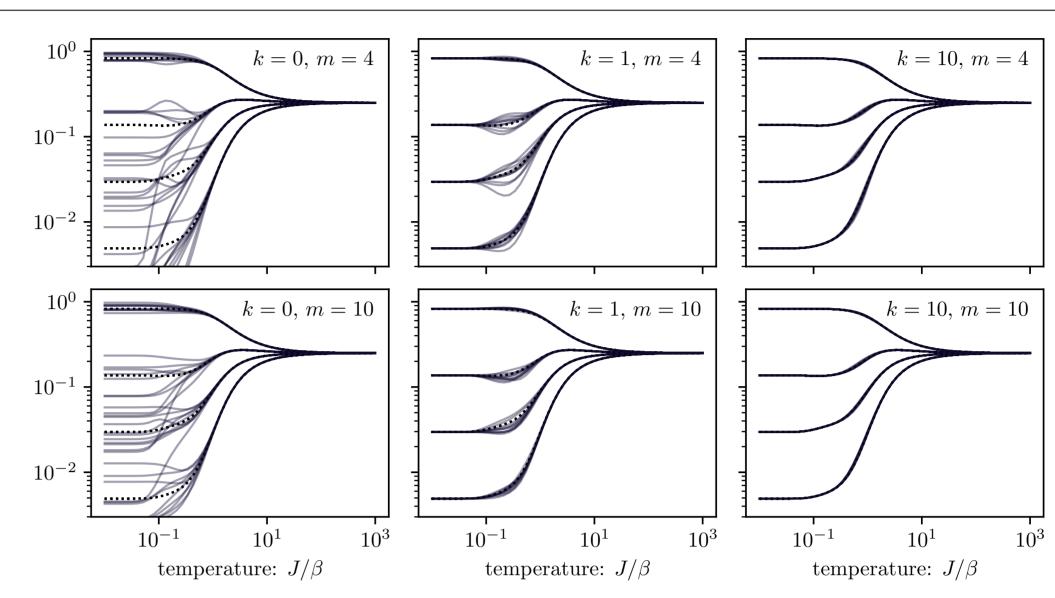


Figure 1. Comparison eigenvalues of $\rho^*(\beta)$ when k=0 (equivalent to [1]), k=1, and k=10. Solid lines correspond to repeated runs of the algorithm and dotted lines the true values.

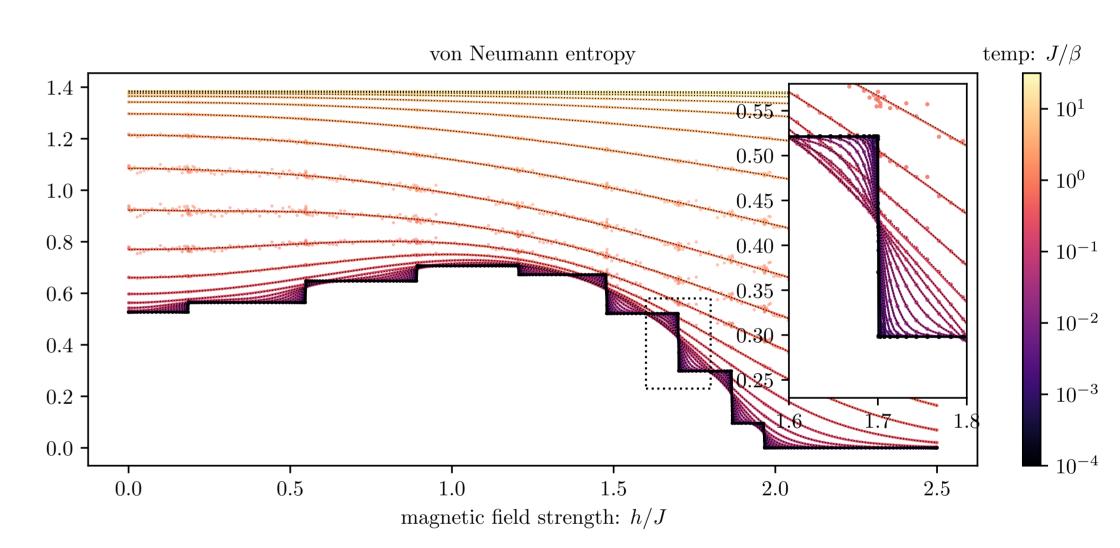


Figure 2. von Neumann Entropy $-\operatorname{tr}(\boldsymbol{\rho}^*(\beta)\ln\boldsymbol{\rho}^*(\beta))$ as a function of magnetic field strength at various temperatures on the solvable model at N = 16.

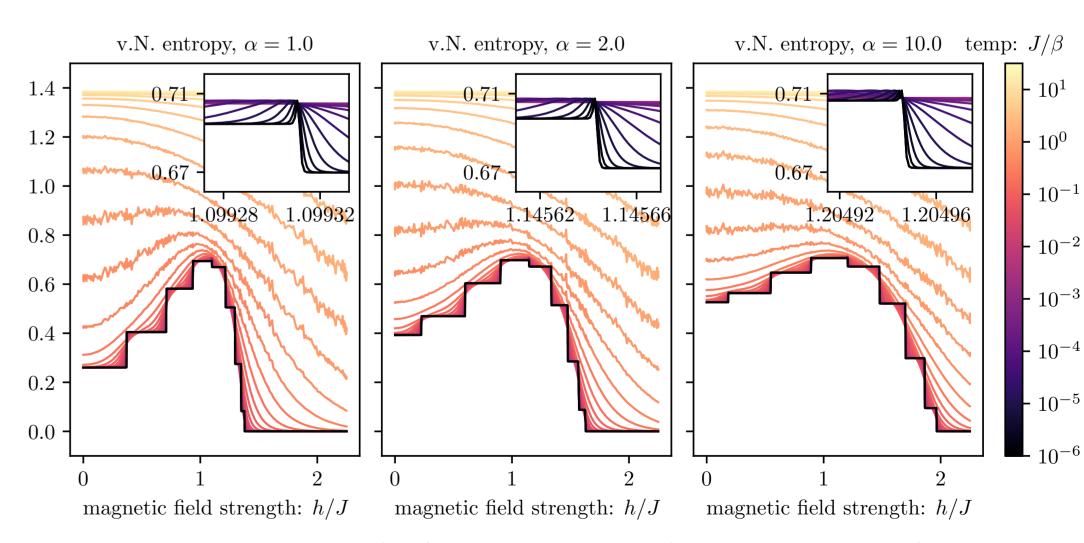


Figure 3. Von Neumann entropy for the long-range spin chain at varying values of α , h, and β . The inset figure shows the transition to the right of the top "plateau".

[1] T. Chen and Y-C. Cheng. Numerical computation of the equilibrium-reduced density matrix for strongly coupled open quantum systems. J. Chem. Phys., 157(6):064106, August 2022. [2] D. Girard. Un algorithme rapide pour le calcul de la trace de l'inverse d'une grande matrice, March 1987. [3] L. Wu, et al. Estimating the trace of the matrix inverse by interpolating from the diagonal of an approximate inverse. J. Comp. Phys., 326:828–844, 2016